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# Behavior of interactions between hydrogen chalcogenides and an anthracene $\pi$-system elucidated by QTAIM dual functional analysis with QC calculations $\dagger$ † 

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#### Abstract

The nature of $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ interactions ( $\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te ) of an anthracene system was elucidated by applying QTAIM dual functional analysis (QTAIM-DFA) after clarification of the structural features with quantum chemical (QC) calculations. $\pi-\mathrm{HB}$ (hydrogen bond) interactions were detected for $\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te , whereas $\pi$-EB (chalcogen bond) interactions were observed for $\mathrm{E}=\mathrm{O}$ in ( $\mathrm{EH}_{2}$ )-*- $\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$, where the bond paths connected H in $\mathrm{EH}_{2}$ to $\mathrm{C}_{14} \mathrm{H}_{10}$ in $\pi-\mathrm{HB}$, and they connected E in $\mathrm{EH}_{2}$ to $\mathrm{C}_{10} \mathrm{H}_{8}$ in $\pi$-EB. The QTAIM-DFA parameters of $(R, \theta)$ and ( $\theta_{p}, \kappa_{p}$ ) were evaluated for the interactions via analysing the plots of $H_{b}\left(r_{c}\right)$ versus $H_{b}\left(r_{c}\right)-V_{b}\left(r_{c}\right) / 2$ for the interactions at the bond critical points. Data obtained from the perturbed structures around the fully optimized structures were employed for the plots, in addition to the fully optimized structures. Data obtained from the fully optimized structures were analysed using $(R, \theta)$, which corresponded to the static nature, and those obtained from the perturbed structures were analysed using ( $\theta_{\mathrm{p}}, \kappa_{\mathrm{p}}$ ), which represented the dynamic nature of the interactions, where $\theta_{\mathrm{p}}$ corresponds to the tangent line of the plot and $\kappa_{\mathrm{p}}$ is the curvature. The $\theta$ and $\theta_{\mathrm{p}}$ values are less than $90^{\circ}$ for all the interactions examined, except for the ${ }^{i} \mathrm{H}-*-{ }^{11} \mathrm{C}(\pi)$ interaction in $\mathrm{TeH}_{2}-*-\mathrm{C}_{14} \mathrm{H}_{10}\left(\mathrm{C}_{1}\right.$ : $\left.\| \mathrm{B}_{\text {Atc }}\right)$, where ' H is located closer to the centre of $\mathrm{C}_{14} \mathrm{H}_{10}$. Therefore, the interactions examined were predicted to have vdW nature, appeared in the pure-CS (closed shell) interaction region, although ${ }^{\mathrm{i}} \mathrm{H}-*-{ }^{11} \mathrm{C}(\pi)$ was predicted to have the pure-CS/typical- HB nature without covalency. Additionally, the $\pi-\mathrm{HB}$ interaction seems to be slightly stronger than $\pi$-EB in $\left(\mathrm{OH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$.


## Introduction

Significant attention has been paid to the interactions between hydrogen chalcogenides $\left(\mathrm{EH}_{2}\right.$ : $\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$ and aromatic $\pi$-systems, and some structures have been reported for benzene adducts formed through $\pi$-interactions. ${ }^{1-8}$ The interactions in benzene $\pi$-systems are mainly characterized by the hydrogen bonds formed between $\mathrm{EH}_{2}$ and the benzene $\pi$-system, which act as a proton donor and proton acceptor, respectively, and have been called hydrogen bonds of the $\pi$-type ( $\pi$-HBs) herein. ${ }^{1-8}$ The $\pi$-HB (hydrogen bond) interactions were also investigated for 1:1 and 2:1 adducts between $\mathrm{EH}_{2}(\mathrm{E}=\mathrm{O}$ and

[^0]S) and the naphthalene $\pi$-system. ${ }^{9-20}$ Indeed, there have been many investigations on $\pi$-adducts, which are mainly based on a theoretical background; however, these investigations seem to be rather limited to structural features. The nature of the $\pi-\mathrm{HB}$ interactions has seldom been reported, especially for the anthracene $\pi$-system, $\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$, to the best of our knowledge.

Very recently, we reported the nature of the $\mathrm{EH}_{2} \cdots \pi$ interactions for benzene and naphthalene $\pi$-systems, together with their structural features. ${ }^{1,2,10,11}$ The H atom(s) in $\mathrm{EH}_{2}$ is (are) connected to the benzene $\pi$-system, $\mathrm{EH}_{2} \cdots \pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$, via bond paths (BPs). Through careful examination of the BPs in the adducts, another type of interaction was also detected in $\mathrm{H}_{2} \mathrm{E} \cdots$ $\pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$, where E in $\mathrm{EH}_{2}$ is joined to the benzene $\pi$-system via a BP. Such interaction was called a chalcogen $\pi$-type ( $\pi$-EB) bond. $\mathrm{EH}_{2}$ is connected to the C atoms or BCPs (bond critical points: $\left.r_{\mathrm{c}}: *\right)^{21}$ on the $\mathrm{C}=\mathrm{C}$ bonds of $\mathrm{C}_{6} \mathrm{H}_{6}$ by BPs. In the case of the benzene $\pi$-system, $\pi$-EBs seem more important relative to $\pi$-HBs, although the predicted importance may change depending on the calculation system. ${ }^{2}$

What happens if $\mathrm{EH}_{2}(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$ interacts with the anthracene $\pi$-system? The anthracene $\pi$-system contains two types of (three) benzene $\pi$-systems and one type of (two)
naphthalene $\pi$-system, in addition to the original anthracene $\pi$ system. What are the differences and similarities in the $\mathrm{EH}_{2} \cdots \pi$ interactions between the anthracene $\pi$-system and the $\pi$ systems of naphthalene and benzene? The nature of the interactions in $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$ is elucidated together with its structural feature. Scheme 1 illustrates some of the structures expected for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ of the $1: 1$ adducts, which are inferred from the structures of $\left(\mathrm{EH}_{2}\right)$ -$*-\pi\left(\mathrm{C}_{10} \mathrm{H}_{8}\right)$ and $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$. Type $\mathrm{IIA}_{\text {Atc }}$ and type $\mathrm{IIC}_{\text {Atc }}$ in $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ are defined for the adduct, where $\mathrm{EH}_{2}$ locates on the central and outer benzene rings of anthracene, respectively, whereas it is above the ${ }^{11} \mathrm{C}-{ }^{12} \mathrm{C}\left(\right.$ or ${ }^{13} \mathrm{C}^{14} \mathrm{C}$ ) bond in type $\mathrm{IIB}_{\text {Atc }}$. The structures for the 1:1 adduct of $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$ and $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{10} \mathrm{H}_{8}\right)$ are also shown in Scheme 1 for convenience of discussion. The optimized structures are called type $\mathrm{Ia}_{\mathrm{Bzn}}, \mathrm{Ib}_{\mathrm{Bzn}}$ and type $\mathrm{II}_{\mathrm{Bzn}}$ for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$ and type $\mathrm{I}_{\text {Nap }}$ and type $\mathrm{II}_{\text {Nap }}$ for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{10} \mathrm{H}_{8}\right)$.

The QTAIM (quantum theory of atoms-in-molecules) approach, introduced by Bader ${ }^{22,23}$ enables the analysis of the nature of chemical bonds and interactions. ${ }^{24-30}$ Interactions seem to be defined by the corresponding BPs, but we must be careful to use the correct terminology with the concept. ${ }^{31}$ BCP is an important concept in QTAIM, which is a point along the BP at the interatomic surface, where charge density, $\rho(\boldsymbol{r})$, reaches a minimum. ${ }^{32}$ It is denoted by $\rho_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$, in addition to other QTAIM functions at BCPs, such as Laplacians of $\rho(\boldsymbol{r})\left(\nabla^{2} \rho_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)\right)$, total electron energy densities $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$, potential energy densities $V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$ and kinetic energy densities $G_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$, together with $k_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)\left(=V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) /\right.$ $G_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$ ).

Recently, we proposed QTAIM dual functional analysis (QTAIM-DFA), ${ }^{33-36}$ according to QTAIM. ${ }^{22-30,37}$ QTAIM-DFA enables experimental chemists to analyse their own results concerning chemical bonds and interactions with their own image. ${ }^{38}$ QTAIM-DFA provides an excellent possibility for evaluating, classifying and understanding weak to strong interactions in a unified form. ${ }^{33-36}$ To elucidate the nature of the interactions in question with QTAIM-DFA, $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$ is plotted






Scheme 1 Structures expected for $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te), based on that optimized for $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$ and $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{10} \mathrm{H}_{8}\right)$.
versus $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)-V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) / 2\left[=\left(\hbar^{2} / 8 m\right) \nabla^{2} \rho_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)\right]$, where both the $x$ and $y$ axes are given units of energy. In our treatment, data for perturbed structures around fully optimized structures are employed for the plots, in addition to the fully optimized structures. ${ }^{33-36}$ We propose the concept of "the dynamic nature of interactions" which originated from the data containing the perturbed structures. ${ }^{33 a, 34-36,39}$

Data from the fully optimized structures correspond to the static nature of interactions. QTAIM-DFA is applied to typical chemical bonds and interactions and rough criteria are established. The rough criteria can distinguish the chemical bonds and interactions in question from others. QTAIM-DFA and the criteria are explained in the ESI with Schemes S1 and S2, Fig. S1 and eqn (S1)-(S7). $\ddagger$ The basic concept of the QTAIM approach is also surveyed.

QTAIM-DFA is applied to elucidate the dynamic and static nature of the interactions in $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and $\mathrm{Te})$. The discussion is limited to the $1: 1$ adducts of $\left(\mathrm{EH}_{2}\right)$ -$*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ for simplicity. Herein, we present the results of the investigations on the nature of the interactions in question. The interactions are classified and characterized as a reference by employing the criteria.

## Methodological details in calculations

Structures were optimized using the Gaussian 09 programs. ${ }^{40}$ The $6-311+G(3 d f)$ basis set ${ }^{41}$ was employed for O, S and Se and the basis set of the ( $7433111 / 743111 / 7411 / 2+1 s 1 p 1 d 1 f)$ type from Sapporo Basis Set Factory ${ }^{42}$ was used for Te with the 6$311++G(d, p)$ basis set ${ }^{41}$ for C and H. The basis set system (BSS) is called BSS-F according to examinations of the BSSs in a previous study. The Møller-Plesset second order energy correlation (MP2) level ${ }^{43}$ was applied to the calculations. Optimized structures were confirmed by the frequency analysis.

QTAIM functions were calculated by employing the wfn files using the Gaussian 09 program package ${ }^{40}$ with the same method for optimizations, and the data were analysed with the AIM2000 program. ${ }^{44}$ The normal coordinates of internal vibrations (NIV) obtained by the frequency analysis were employed to generate the perturbed structures, ${ }^{35,36}$ which is explained in eqn (1). The $k$-th perturbed structure in question $\left(\boldsymbol{S}_{\boldsymbol{k w}}\right)$ was generated by the addition of the normal coordinates of the $k$-th internal vibration $\left(\boldsymbol{N}_{\boldsymbol{k}}\right)$ to the standard orientation of a fully optimized structure ( $\boldsymbol{S}_{\mathbf{0}}$ ) in the matrix representation. ${ }^{35}$ The coefficient $f_{k w}$ in eqn (1) controls the difference in the structures between $\boldsymbol{S}_{\boldsymbol{k} \boldsymbol{w}}$ and $\boldsymbol{S}_{\mathbf{0}}: f_{k w}$ is determined to satisfy eqn (2) for the interaction in question, where, $r$ and $r_{0}$ stand for the distances in question in the perturbed and fully optimized structures, respectively, with $a_{0}$ of Bohr radius ( $0.52918 \AA$ ). The perturbed structures with NIV correspond to that with $r$ being elongated or shortened by $0.05 a_{0}$ or $0.1 a_{0}$, relative to $r_{0}$, as shown in eqn (2). $\boldsymbol{N}_{\boldsymbol{k}}$ of five digits are used to predict $\boldsymbol{S}_{\boldsymbol{k w}}$. The selected vibration must contain the motion of the interaction in question most effectively among all the zero-point internal vibrations.

$$
\begin{equation*}
\boldsymbol{S}_{\boldsymbol{k} w}=\boldsymbol{S}_{\mathbf{0}}+f_{k w} \boldsymbol{N}_{\boldsymbol{k}} \tag{1}
\end{equation*}
$$

$$
\begin{equation*}
r=r_{0}+w a_{0}\left(w=(0), \pm 0.05 \text { and } \pm 0.1 ; a_{0}=0.52918 \AA\right) \tag{2}
\end{equation*}
$$

In the QTAIM-DFA treatment, $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$ is plotted versus $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)-$ $V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) / 2$ for data of five points of $w=0, \pm 0.05$ and $\pm 0.1$, as shown in eqn (2). Each plot is analysed using a regression curve of the cubic function, as shown in eqn (3), where $(x, y)=\left(H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)\right.$ $\left.-V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) / 2, H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)\right)\left(R_{\mathrm{c}}{ }^{2}>0.99999\right.$ usually $) .{ }^{45}$

## Results and discussion

Before the final structural optimizations for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ ( $\mathrm{E}=\mathrm{O}, \mathrm{S}$, Se and Te ), the minima were searched for systematically with MP2/BSS-F. The search was started assuming typically type $\mathrm{IIA}_{\text {Atc }}$, type $\mathrm{IIB}_{\text {Atc }}$ and type $\mathrm{IIC}_{\text {Atc }}$ structures of the $C_{1}$ symmetry. The type $\mathrm{IIA}_{\text {Atc }}$ and type $\mathrm{IIB}_{\text {Atc }}$ structures of $\left(\mathrm{EH}_{2}\right)$ symmetry. The type $\mathrm{IIA}_{\text {Atc }}$ and type $\mathrm{IB}_{\text {Atc }}$ Structures of $\left(\mathrm{EH}_{2}\right)^{-}$
$*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ were optimized for $\mathrm{E}=\mathrm{O}$ and S , whereas only type $\mathrm{IIB}_{\text {Atc }}$ was optimized for $\mathrm{E}=\mathrm{Se}$ and Te . The processes of the convergence are summarized in Table S1 of the ESI. $\ddagger$

The type IIA $_{\text {Atc }}$ structure optimized for $\left(\mathrm{OH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ with MP2/BSS-F was apparently different from type IIB $_{\text {Atc }}$, as with MP2/BSS-F was apparently different from type IIB $_{\text {Atc }}$, as
confirmed by the $r_{1}$ values. In the case of $\left(\mathrm{SH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$, one structure was optimized, which was close to the type IIA $_{\text {Atc }}$ structure if initially assuming a typical type $\mathrm{IIA}_{\text {Atc }}$ structure. The type $\mathrm{IIA}_{\text {Atc }}$ structure optimized for $\left(\mathrm{SH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left[\mathrm{SH}_{2} \cdots\right.$ $\left.\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIA}_{\text {Atc }}\right)\right]$ seems somewhat distorted from the $C_{\mathrm{s}}$ symmetry. ${ }^{46}$ The search converged to another one, if started assuming typical type $\mathrm{IIB}_{\text {Atc }}$ and type $\mathrm{IIC}_{\text {Atc }}$ structures, which are close to type $\mathrm{IIA}_{\text {Atc }}$. Indeed, their $r_{1}$ values are very close to each other, but the $\phi_{2}$ and $\phi_{3}$ values seem meaningfully different for the two types (Fig. 1c and d, respectively). Therefore, the second structure is (tentatively) called type $\mathrm{IIB}_{\text {Atc }}$, herein. No type $\mathrm{I}_{\text {Atc }}$
$\left(R_{\mathrm{c}}{ }^{2}\right.$ : square of correlation coefficient).

$$
\begin{equation*}
y=c_{\mathrm{o}}+c_{1} x+c_{2} x^{2}+c_{3} x^{3} \tag{3}
\end{equation*}
$$

## Structural feature of $\left(\mathbf{E H}_{\mathbf{2}}\right) \cdots \boldsymbol{\pi}\left(\mathbf{C}_{\mathbf{1 4}} \mathbf{H}_{\mathbf{1 0}}\right)(\mathrm{E}=\mathbf{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$

structures were detected after similar treatment, even if optimizations started from those very close to type $\mathrm{IA}_{\text {Atc }}$.

Table 1 collects the structural parameters selected for the optimized structures of $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$, $r_{1}, r_{2}, r_{3}, \theta_{1}, \theta_{2}, \theta_{3}, \phi_{1}, \phi_{2}$ and $\phi_{3}$, which are defined in Scheme 2. The optimized structures are not shown in the figures, but they can be found in molecular graphs, which are drawn on the optimized structures (see Fig. 1). What factors appear to control the optimized structures? We compared the $\mathrm{H} \cdots \mathrm{H}$ distance in $\mathrm{EH}_{2}, r\left(\mathrm{H}, \mathrm{H}: \mathrm{EH}_{2}\right)$, with the ${ }^{11} \mathrm{C} \cdots{ }^{13} \mathrm{C}$ distance in $\mathrm{C}_{14} \mathrm{H}_{10}, r\left({ }^{11} \mathrm{C}\right.$, $\left.{ }^{13} \mathrm{C}: \mathrm{C}_{14} \mathrm{H}_{10}\right)$. Indeed, $r\left(\mathrm{H}, \mathrm{H}: \mathrm{EH}_{2}\right)$ for $\mathrm{E}=\mathrm{O}(1.522 \AA), \mathrm{S}(1.931$ $\AA)$, $\mathrm{Se}(2.088 \AA)$ and $\mathrm{Te}(2.361 \AA)$ is shorter than $r\left({ }^{11} \mathrm{C},{ }^{13} \mathrm{C}\right.$ : $\left.\mathrm{C}_{14} \mathrm{H}_{10}\right)(2.452 \AA$ for the central benzene ring), but the differences in $r\left(\Delta r=r\left(\mathrm{C}, \mathrm{C}: \mathrm{C}_{14} \mathrm{H}_{10}\right)-r\left(\mathrm{H}, \mathrm{H}: \mathrm{EH}_{2}\right)\right)$ are larger than $0.5 \AA$ for $\mathrm{O}(\Delta r=0.9303 \AA)$ and $\mathrm{S}(0.5215 \AA)$, whereas they are smaller than $0.4 \AA$ for $\mathrm{Se}(0.3644 \AA)$ and $\mathrm{Te}(0.0913 \AA)$ (see Table S2 of the ESI $\ddagger)$. The type IIB $_{\text {Atc }}$ structures of $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}$ $=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te ) would form with no limitations of $\Delta r$, whereas it may be necessary for $\Delta r$ larger than around $0.5 \AA$ to give the type $\mathrm{IIA}_{\text {Atc }}$ structures $(\mathrm{E}=\mathrm{O}$ and S$)$. In the case of type $\mathrm{IIB}_{\mathrm{Atc}}$, one H in $\mathrm{EH}_{2}$ seems to be slightly above the ${ }^{11} \mathrm{C} \cdots{ }^{13} \mathrm{C}$ bond for $\mathrm{E}=\mathrm{S}$, Se and Te , but $\mathrm{OH}_{2}$ seems to exist almost right above the ${ }^{11} \mathrm{C} \cdots{ }^{13} \mathrm{C}$ bond, which may also be controlled by $\Delta r$. Indeed, the difference in $\Delta r$ between $\mathrm{E}=\mathrm{S}$ and Se seems small, at a first glance, but the small difference would play an important role in the appearance of the type IIA $_{\text {Atc }}$ structure.

What are the stabilization energies in the formation of the adducts? Table 1 contains the $\Delta E\left(\Delta E_{\mathrm{ES}}\right.$ and $\left.\Delta E_{\mathrm{Ent}}\right)$ values for $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$, where, $\Delta E_{\text {ES }}$ and $\Delta E_{\text {Ent }}$ are the $\Delta E$ values on the energy surfaces and those with the thermal corrections to enthalpies, respectively $\left[\Delta E=\left(E\left(\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\right)-\left(E\left(\mathrm{EH}_{2}\right)+\right.\right.\right.$ $\left.\left.E\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\right)\right]$. An excellent correlation was obtained in the plot of $\Delta E_{\text {Ent }}$ versus $\Delta E_{\mathrm{ES}}\left(y=5.262+1.047 x ;{R_{\mathrm{c}}{ }^{2}=0.998(n=6 \text { : number }}^{2}\right.$ of data points)), although is not shown in the figure. Therefore,


Fig. 1 Molecular graphs for $\mathrm{OH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(\mathrm{C}_{1}\right.$ : type $\left.I I \mathrm{~A}_{\text {Atc }}\right)(\mathrm{a}), \mathrm{H}_{2} \mathrm{O}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(\mathrm{C}_{1}\right.$ : type $\left.I I \mathrm{~B}_{\text {Atc }}\right)(\mathrm{b}), \mathrm{SH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(\mathrm{C}_{1}\right.$ : type IIA $\left.\mathrm{A}_{\text {Atc }}\right)(\mathrm{c})$, $\mathrm{SH} \mathrm{H}_{2}-$ $*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(\mathrm{C}_{1}\right.$ : type IIB Atc ) (d), $\mathrm{SeH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(\mathrm{C}_{1}\right.$ : type IIB Atc ) (e) and $\mathrm{TeH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}\right.$ : type IIB Atc$)$ (f), calculated with MP2/BSS-F. The bond critical points (BCPs) are denoted by red dots, ring critical points (RCPs) by yellow dots and cage critical points (CCPs) by green dots, together with bond paths by pink lines. Carbon atoms are in black and hydrogen atoms are in grey, with oxygen, sulphur, selenium and tellurium atoms in red, yellow, pink and purple, respectively.

Table 1 Structural parameters for $\left(E H_{2}\right) \cdots \pi\left(C_{14} \mathrm{H}_{10}\right)(E=O, S$, Se and $T e)$ optimized at the MP2 level with BSS- $F^{a b c}$

| Species (X-Y) (symmetry: type) | $r_{1}(\mathrm{~A})$ | $r_{2}(\mathrm{~A})$ | $r_{3}(\mathrm{~A})$ | $\theta_{1}\left({ }^{\circ}\right)$ | $\theta_{2}\left({ }^{\circ}\right)$ | $\theta_{3}\left({ }^{\circ}\right)$ | $\phi_{1}\left({ }^{\circ}\right)$ | $\phi_{2}\left({ }^{\circ}\right)$ | $\phi_{3}\left({ }^{\circ}\right)$ | $\Delta E_{\mathrm{ES}}{ }^{d}\left(\mathrm{~kJ} \mathrm{~mol}^{-1}\right)$ | $\Delta E_{\mathrm{Ent}^{e}}\left(\mathrm{~kJ} \mathrm{~mol}{ }^{-1}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{OH}_{2} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}:^{f} \mathrm{IIA}_{\text {Atc }}\right)$ | 3.2865 | 0.9629 | 0.9629 | 89.9 | 51.8 | 103.5 | -92.5 | -1.7 | 2.8 | -22.7 | -18.1 |
| $\mathrm{H}_{2} \mathrm{O} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)$ | 3.6682 | 0.9631 | 0.9624 | 63.4 | 69.6 | 103.6 | -93.8 | -14.3 | 3.5 | -21.9 | -17.6 |
| $\mathrm{SH}_{2} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIA}_{\text {Atc }}\right)^{g}$ | 3.6162 | 1.3389 | 1.3400 | 80.5 | 60.6 | 92.2 | -92.1 | -6.4 | -3.3 | -28.7 | -25.0 |
| $\mathrm{SH}_{2} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)^{g}$ | 3.6165 | 1.3389 | 1.3400 | 80.5 | 60.6 | 92.2 | -91.8 | -2.1 | -1.9 | -28.7 | -25.0 |
| $\mathrm{SeH}_{2} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)$ | 3.6959 | 1.4597 | 1.4603 | 78.9 | 61.5 | 91.3 | -91.5 | 3.0 | -0.2 | -31.7 | -28.3 |
| $\mathrm{TeH}_{2} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)$ | 3.7703 | 1.6576 | 1.6578 | 78.1 | 64.2 | 91.5 | -89.9 | -11.4 | -17.5 | -40.5 | -36.8 |

${ }^{a}$ BSS-F: the $6-311+G(3 d f)$ basis set was employed for $O$, $S$ and Se and the basis set of the $(7433111 / 743111 / 7411 / 2+1$ s1p1d1f) type for Te with the 6 $311++\mathrm{G}(\mathrm{d}, \mathrm{p})$ basis set for C and $\mathrm{H} .{ }^{b}$ See Scheme 2 for the structural parameters. ${ }^{c}$ Optimized structures are not given in the figures but can be found in the molecular graphs drawn on the optimized structures (see Fig. 1). ${ }^{d} \Delta E_{\mathrm{ES}}=E_{\mathrm{ES}}\left(\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\right)-\left(E_{\mathrm{ES}}\left(\mathrm{EH}_{2}\right)+E_{\mathrm{ES}}\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\right)$ on the energy surface. ${ }^{e} \Delta E_{\text {Ent }}=E_{\text {Ent }}\left(\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\right)-\left(E_{\mathrm{Ent}}\left(\mathrm{EH}_{2}\right)+E_{\mathrm{Ent}}\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\right)$ with the thermal corrections to enthalpies. ${ }^{f}$ Very close to $C_{\mathrm{s}} .{ }^{g}$ The structures are very close to each other; however, they are analysed as two different structures here since the differences in $\phi_{2}$ and $\phi_{3}$ seem meaningful. They are called IIA $_{\text {Atc }}$ and IIB $_{\text {Atc }}$.


Scheme 2 Structural Parameters Illustrated for Type $\|_{\text {Atc }}$ of $\left(E H_{2}\right)$ -*- $\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$.
the $\Delta E_{\text {ES }}$ values will be employed for the discussion of $\Delta E$. A good correlation was also obtained in the plot of $\Delta E_{\mathrm{ES}}$ versus $r_{1}$, if the data for $\mathrm{H}_{2} \mathrm{O} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)$ are neglected $(y=$ $-1410.42+820.87 x-121.30 x^{2} ;{R_{c}}^{2}=0.975(n=5)$ ), where the structure of $\mathrm{H}_{2} \mathrm{O} \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}:\right.$ IIB $\left._{\text {Atc }}\right)$ seems different from the others.

What are the relations between $\Delta E_{\mathrm{Atc}}, \Delta E_{\mathrm{Nap}}$ and $\Delta E_{\mathrm{Bzn}}$ ? The $\Delta E_{\text {Nap }}$ values are linearly proportional to $\Delta E_{\text {Bzn }}$ in type II. The ratio of $\Delta E_{\text {Nap }} / \Delta E_{\mathrm{Bzn}}$ becomes larger in the order of $\mathrm{E}=\mathrm{O}$ $\left(\Delta E_{\text {Nap }} / \Delta E_{\text {Bzn }}=1.34\right)<\mathrm{S}(1.60)<\operatorname{Se}(1.68)<$ Te (1.75), which shows a substantial chalcogen dependence. On the other hand, the $\Delta E_{\text {Atc }} / \Delta E_{\text {Nap }}$ ratio seems almost constant for $\mathrm{E}=\mathrm{O}\left(\Delta E_{\mathrm{Atc}} /\right.$ $\left.\Delta E_{\text {Nap }}=1.01\right)<\mathrm{S}(1.07)<\mathrm{Se}(1.09)<\mathrm{Te}(1.11)$. The different basis sets for Te may somewhat affect the evaluated values for E $=\mathrm{Te}$, which are from the Sapporo Basis Set Factory, whereas the others are from the Gaussian09 program. What mechanisms operate to control the ratios? It is difficult to clarify this, however, we examined the ratios of $r_{1: \mathrm{Bzn}} / r_{1: \mathrm{Nap}}$ and $r_{1: \mathrm{Atc}} /$ $r_{1: \text { Nap. }}$. The $r_{1: \mathrm{Bzn}} / r_{1: \text { Nap }}$ ratios become larger in the order of $1.01(\mathrm{E}=\mathrm{O})<1.13(\mathrm{~S})<1.15(\mathrm{Se})<1.19(\mathrm{Te})$ and the $r_{1}:$ Atc $/ r_{1}:$ Nap ratios increase similarly in the order of $1.00(\mathrm{E}=\mathrm{O})<1.09(\mathrm{~S})<$ $1.11(\mathrm{Se})<1.13(\mathrm{Te})$. However, the chalcogen dependence of the ratios seems somewhat smaller for $r_{1}$ : Atc $/ r_{1}$ : Nap, relative to the case for $r_{1: \mathrm{Bzn}} / r_{1}$ : Nap. These results would be responsible for the observed $\Delta E / \Delta E$ ratios, although the vdW and/or covalent radii of chalcogens must also be carefully examined for the discussion.

Before the application of QTAIM-DFA to $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$, the molecular graphs and contour plots are examined next.

## Molecular graphs and contour plots of $\rho(r)$ for $\left(\mathbf{E H}_{2}\right) \cdots$ $\pi\left(\mathrm{C}_{\mathbf{1 4}} \mathrm{H}_{\mathbf{1 0}}\right)$

Fig. 1 illustrates the molecular graphs for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}$ $=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te ), which were calculated with MP2/BSS-F. The BCPs are clearly detected, which contain that expected between $\mathrm{EH}_{2}$ and $\mathrm{C}_{14} \mathrm{H}_{10}$. Each BP with a BCP connects each H in $\mathrm{EH}_{2}$ and the carbon atom of ${ }^{11} \mathrm{C}$ or ${ }^{13} \mathrm{C}$ in $\mathrm{C}_{14} \mathrm{H}_{10}$ for all optimized structures in Table 2, except for the type IIB $_{\text {Atc }}$ of $\mathrm{H}_{2} \mathrm{O}$ -$*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$. The BP with BCP joins O in $\mathrm{OH}_{2}$ and a carbon atom of ${ }^{11} \mathrm{C}$ or ${ }^{13} \mathrm{C}$ in $\mathrm{C}_{14} \mathrm{H}_{10}$ of $\mathrm{OH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$. Therefore, the interactions between $\mathrm{EH}_{2}$ and $\mathrm{C}_{14} \mathrm{H}_{10}$ are classified by $\pi$ - HBs for all the adducts in Table 2, except for the type $\mathrm{IIB}_{\text {Atc }}$ of $\mathrm{H}_{2} \mathrm{O}-$ $*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$, the interaction of which should be classified by the $\pi$-EB type. Ring critical points (RCPs) and cage critical points (CCPs) are also detected in Fig. 1. Fig. 2 shows the contour plots of $\rho(r)$ for $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$ calculated with MP2/BSS-F. The contour plots of $\rho(r)$ are drawn on a plane containing E, BCP $(*)$ of the $\left(\mathrm{EH}_{2}\right)-*-\pi$ type and an atom or BCP suitable for the contour plots of $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$. The BCPs are well located at the (three dimensional) saddle points of $\rho(r)$ in the species. Negative Laplacians and trajectory plots are drawn in Fig. S2 and S3 of the ESI. $\ddagger$ It is well visualized how the BCPs are classified through $\nabla^{2} \rho(r)$ and the space around the species is well divided by the atoms in it.

## Survey of the interactions in $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(E=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te)

As shown in Fig. 1 and 2, some BPs are apparently curved. Therefore, the lengths of the BPs ( $r_{\mathrm{BP}}$ ) will be substantially different from the straight-line distances $\left(R_{\mathrm{SL}}\right)$ in some cases. The $r_{\mathrm{BP}}$ and $R_{\mathrm{SL}}$ values are collected in Table S 2 of the ESI $\ddagger$ for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$, which were evaluated with MP2/BSS-F, together with the differences between $r_{\mathrm{BP}}$ and $R_{\mathrm{SL}}\left(\Delta r_{\mathrm{BP}}=r_{\mathrm{BP}}-R_{\mathrm{SL}}\right)$. There are two types of BPs in each adduct since each adduct has $C_{1}$ symmetry. Therefore, it is necessary to define the BPs to distinguish the two. Letting ${ }^{\mathrm{i}} \mathrm{H}$ and ${ }^{\circ} \mathrm{H}$ in $\mathrm{EH}_{2}$ connect to $\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ through BPs, which are placed more inside and more outside, respectively, in relation to the centre of $\mathrm{C}_{14} \mathrm{H}_{10}$, then ${ }^{\mathrm{i}} \mathrm{BP}$ and ${ }^{o} \mathrm{BP}$ can be defined as ${ }^{\mathrm{i}} \mathrm{BP}\left({ }^{\mathrm{i}} \mathrm{H}-*-\pi\right)$ and ${ }^{o} \mathrm{BP}\left({ }^{\circ} \mathbf{H}-*-\pi\right)$, respectively.

Table 2 QTAIM functions and QTAIM-DFA parameters for the interactions in $\left(E H_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$ evaluated with MP2/BSS-F ${ }^{a}$

| Interaction $(\mathrm{X}-*-\mathrm{Y})^{e}$ | Type ${ }^{\text {b }}$ | $\begin{aligned} & \rho_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) \\ & \left(e a_{\mathrm{o}}^{-3}\right) \end{aligned}$ | $\begin{aligned} & c \nabla^{2} \rho_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)^{c} \\ & \text { (au) } \end{aligned}$ | $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$ <br> (au) | $k_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)^{d}$ | $R$ (au) | $\theta\left({ }^{\circ}\right)$ | Freq $\left(\mathrm{cm}^{-1}\right)$ | $k_{\text {f }}$ (unit ${ }^{\text {f }}$ | $\theta_{\mathrm{p}}\left({ }^{\circ}\right)$ | $\kappa_{\mathrm{p}}\left(\mathrm{au}^{-1}\right)$ | Classification characterization |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| OHH $*$ - ${ }^{\mathbf{1 1}(13)} \mathbf{C}(\pi)$ | IIA Atc ${ }^{\text {g }}$ | 0.0060 | 0.0024 | 0.0007 | -0.814 | 0.0025 | 72.6 | 108.7 | 0.0411 | 74.5 | 8.0 | p-CS/vdW |
| $\mathrm{H}_{2} \mathrm{O}-*-{ }^{13} \mathrm{C}(\tau)$ | $\mathrm{IIB}_{\text {Atc }}$ | 0.0069 | 0.0029 | 0.0009 | -0.831 | 0.0031 | 73.9 | 104.5 | 0.0346 | 79.3 | 18.1 | p-CS/vdW |
| $\mathrm{S}^{0} \mathrm{H}^{\mathrm{i}} \mathbf{H}-*-{ }^{\mathbf{1 1}} \mathrm{C}(\pi)^{h}$ | $\mathrm{IIA}_{\text {Atc }}$ | 0.0072 | 0.0027 | 0.0009 | -0.807 | 0.0029 | 72.1 | 65.9 | 0.0151 | 72.2 | 1.1 | p-CS/vdW |
| $\mathrm{S}^{\mathrm{i}} \mathrm{H}^{\mathbf{o}} \mathbf{H}-*-{ }^{\mathbf{1 3}} \mathbf{C}(\pi)^{h}$ | $\mathrm{IIA}_{\text {Atc }}$ | 0.0065 | 0.0024 | 0.0007 | -0.835 | 0.0025 | 74.2 | 65.9 | 0.0151 | 74.3 | 145.7 | p-CS/vdW |
| $\mathrm{S}^{o} \mathrm{H}^{\mathrm{i}} \mathbf{H}-*-{ }^{\mathbf{1 1}} \mathbf{C}(\pi)^{h}$ | $\mathrm{IIB}_{\text {Atc }}$ | 0.0070 | 0.0027 | 0.0009 | -0.799 | 0.0029 | 71.5 | 65.8 | 0.0161 | 73.6 | 37.3 | p-CS/vdW |
| $\mathrm{S}^{\mathrm{i}} \mathrm{H}^{\mathbf{o}} \mathbf{H}-*-{ }^{\mathbf{1 3}} \mathrm{C}(\pi)^{h}$ | $\mathrm{IIB}_{\text {Atc }}$ | 0.0065 | 0.0024 | 0.0007 | -0.836 | 0.0025 | 74.2 | 65.8 | 0.0161 | 73.4 | 61.8 | p-CS/vdW |
| $\mathrm{Se}^{o} \mathrm{H}^{\mathbf{i}} \mathbf{H}-*-{ }^{\mathbf{1 1}} \mathbf{C}(\pi)^{h}$ | $\mathrm{IIB}_{\text {Atc }}$ | 0.0072 | 0.0028 | 0.0010 | -0.789 | 0.0029 | 70.8 | 54.6 | 0.0158 | 72.7 | 30.4 | p-CS/vdW |
| $\mathrm{Se}^{\mathrm{i}} \mathrm{H}^{0} \mathbf{H}-*-{ }^{13} \mathrm{C}(\pi)^{h}$ | $\mathrm{IIB}_{\text {Atc }}$ | 0.0070 | 0.0025 | 0.0007 | -0.846 | 0.0026 | 75.0 | 54.6 | 0.0158 | 74.2 | 37.1 | p-CS/vdW |
| $\mathrm{Te}^{o} \mathbf{H}^{\mathrm{i}} \mathbf{H}-*-{ }^{\mathbf{1 1}} \mathbf{C}(\tau){ }^{h}$ | $\mathrm{IIB}_{\text {Atc }}$ | 0.0077 | 0.0028 | 0.0008 | -0.829 | 0.0029 | 73.7 | 50.5 | 0.0086 | 93.8 | i | $\mathrm{p}-\mathrm{CS} / t-\mathrm{HB}_{\mathrm{nc}}$ |
| $\mathrm{Te}^{\mathrm{i}} \mathrm{H}^{\mathbf{o}} \mathbf{H}-*-{ }^{\mathbf{1 3}} \mathbf{C}(\tau){ }^{h}$ | $\mathrm{IIB}_{\text {Atc }}$ | 0.0088 | 0.0031 | 0.0008 | -0.849 | 0.0032 | 75.3 | 23.9 | 0.0028 | 82.9 | 299.8 | p-CS/vdW |

${ }^{a}$ BSS-F: the $6-311+G(3 \mathrm{df})$ basis set employed for $\mathrm{O}, \mathrm{S}$ and Se and the basis set of the $(7433111 / 743111 / 7411 / 2+1$ s1p1d1f) type for Te with the 6 $311++\mathrm{G}(\mathrm{d}, \mathrm{p})$ basis set for C and H. ${ }^{b}$ See Scheme 1 and Fig. $1 .{ }^{c} c \nabla^{2} \rho_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)=H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)-V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) / 2$, where $c=\hbar^{2} / 8 m$. ${ }^{d} k_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)=V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) / G_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) .{ }^{e}$ The optimized structure has $C_{1}$ symmetry and atoms taking part in the interaction are shown in bold with BCP denoted by $*^{f}{ }^{f} \mathrm{mDyne}^{\circ} \AA^{-1}$. ${ }^{g}$ Very close to $C_{\mathrm{s}}$ symmetry. ${ }^{h \mathrm{i}} \mathrm{H}$ and ${ }^{o} \mathrm{H}$ in $\mathrm{EH}_{2}$ stand for the atoms taking part in the interactions, which are placed more inside and more outside, respectively, in relation to the centre of $\mathrm{C}_{14} \mathrm{H}_{10} .{ }^{i}$ Vary large value of $1259 \mathrm{au}^{-1}$.


Fig. 2 Contour plots of $\rho(r)$ for $\mathrm{OH}_{2}-*-\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}\right.$ : type $\left.\| \mathrm{A}_{\text {Atc }}\right)(\mathrm{a}), \mathrm{H}_{2} \mathrm{O}-*-\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(\mathrm{C}_{1}\right.$ : type IIB Atcc$)(\mathrm{b}), \mathrm{SH}_{2}-*-\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(\mathrm{C}_{1}\right.$ : type $\left.I I \mathrm{~A}_{\text {Atc }}\right)(\mathrm{c}), \mathrm{SH}_{2}-$
 on the plane are denoted by red dots, those outside of the plane in dark pink dots, RCPs by blue squares, CCPs by green dots and bond paths on the plane by black line and those outside of the plane are by grey lines. Carbon atoms are in black and hydrogen atoms are in grey, with the other atoms in black. The contours $\left(e a_{0}{ }^{-3}\right)$ are at $2^{l}(l= \pm 8, \pm 7, \ldots, 0)$ with 0.0047 (heavy line).

The $\Delta r_{\mathrm{BP}}$ values are small (ca. $0.03 \AA$ Aㅇ for ${ }^{\mathrm{i}} \mathrm{BP}\left(\mathrm{S}^{o} \mathrm{H}^{\mathrm{i}} \mathbf{H}-{ }^{-{ }^{11}} \mathrm{C}(\pi)\right)$ and ${ }^{\mathrm{i}} \mathrm{BP}\left(\mathrm{Se}^{o} \mathrm{H}^{\mathrm{i}} \mathbf{H}-*-{ }^{11} \mathrm{C}(\pi)\right)$ in $\left(\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\right)\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)$. Thus, the ${ }^{\mathrm{i}} \mathrm{BPs}$ can be approximated as straight lines. However, the $\Delta r_{\mathrm{BP}}$ value is large ( $0.41 \AA$ ) for ${ }^{o} \mathrm{BP}\left(\mathrm{Se}^{\mathrm{i}} \mathrm{H}^{o} \mathbf{H}-*-{ }^{11} \mathrm{C}(\pi)\right.$ ), which seems difficult to be approximated by a straight line. The values seem moderate ( $0.09-0.17 \AA$ ) for most of the $\pi$-HB and $\pi$-EB interactions in $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ for $\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te , other than three cases. The BPs could be approximated as almost straight lines to gentle curves. The $r_{\mathrm{BP}}$ values are plotted versus $R_{\mathrm{SL}}$ in Fig. 3. A very good correlation was obtained for the case of $0.09 \leq \Delta r_{\mathrm{BP}} \leq 0.17 \AA$, which is shown in the figure. The data for
three BPs deviate from the correlation, although two of them can be approximated as straight lines, as mentioned above.

QTAIM-DFA was applied to the interactions between $\mathrm{EH}_{2}$ and $\mathrm{C}_{14} \mathrm{H}_{10}$ in $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te$)$ and the QTAIM functions were calculated for the interactions at BCP. The results are given in Table 2. Fig. 4 shows the plot of $H_{b}\left(\boldsymbol{r}_{\mathrm{c}}\right)$ versus $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)-V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) / 2$ for $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}$, Se and Te$)$, which contains the data for the perturbed structures evaluated with MP2/BSS-F. All the data in Fig. 4 appear in the area of $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)$ $-V_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right) / 2>0$ and $H_{\mathrm{b}}\left(\boldsymbol{r}_{\mathrm{c}}\right)>0$, which belongs to the pure-CS (closed shell: p-CS) region. The plots were analysed according


Fig. 3 Plots of $r_{\mathrm{BP}}$ versus $R_{\mathrm{SL}}$ for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}$, Se and Te), evaluated with BSS-F at the MP2 level.


Fig. 4 Plots of $H_{b}\left(r_{\mathrm{c}}\right)$ versus $H_{\mathrm{b}}\left(r_{\mathrm{c}}\right)-V_{\mathrm{b}}\left(r_{\mathrm{c}}\right) / 2$ for $\left(\mathrm{EH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$, evaluated with MP2/BSS-F. Compounds with marks are shown in the figure.
to eqn (S3)-(S6) of the ESI $\ddagger$ by applying QTAIM-DFA. These results are discussed next.

## Nature of $\pi$-HBs and $\pi$-EBs in $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}$, Se and Te)

Table 2 collects the QTAIM-DFA parameters of $(R, \theta)$ and $\left(\theta_{\mathrm{p}}, \kappa_{\mathrm{p}}\right)$ for the $\pi$-HB and $\pi$-EB interactions, together with the frequencies and the force constants ( $k_{\mathrm{f}}$ ) correlated to NIV, which are employed to generate the perturbed structures. The nature of $\pi$-HBs and $\pi$-EBs in $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ is discussed based on the ( $R, \theta, \theta_{\mathrm{p}}$ ) values employing the standard values as a reference (see Scheme S2 of the ESI $\ddagger$ ). It should be instructive to survey the criteria related to that herein before a detailed discussion.

The criteria tell us that $45^{\circ}<\theta<90^{\circ}$ for pure-CS interactions. The $\theta_{\mathrm{p}}$ value predicts the character of the interactions. In the pCS region, the character of the interactions will be the vdW type for $45^{\circ}<\theta_{\mathrm{p}}<90^{\circ}$, whereas it will be the $t$-HB (typical hydrogen bonds) type with no covalency ( $t-\mathrm{HB}_{\mathrm{nc}}$ ) for $90^{\circ}<\theta_{\mathrm{p}} \leq 125^{\circ}$, where $\theta_{\mathrm{p}}=125^{\circ}$ tentatively corresponds to $\theta=90^{\circ}$.

The $\theta$ and $\theta_{\mathrm{p}}$ values are less than $90^{\circ}$ for all the $\pi$-HBs and $\pi$ EB in $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ examined herein, except for the $\mathrm{Te}^{\mathrm{i}} \mathrm{H}^{o} \mathrm{H}-$ ${ }^{-11} \mathrm{C}(\pi)$ interaction in $\mathrm{TeH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$. Therefore, the $\pi$-HBs and $\pi$-EB interactions in $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ in Table 2 are all classified by the pure-CS interactions and characterized as the vdW type ( $\mathrm{p}-\mathrm{CS} / \mathrm{vdW}$ ), except for $\mathrm{Te}^{\mathrm{i}} \mathrm{H}^{o} \mathrm{H}_{-*-}{ }^{11} \mathrm{C}(\pi)$ in $\mathrm{TeH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$. Although ${ }^{o} \mathrm{BP}\left(\mathrm{Te}^{\mathrm{i}} \mathrm{H}^{o} \mathrm{H}-{ }^{-}{ }^{13} \mathrm{C}(\pi)\right)$ is predicted to have the nature of ( $\mathrm{p}-\mathrm{CS} / \mathrm{vdW}$ ), ${ }^{\mathrm{i}} \mathrm{BP}\left(\mathrm{Te}^{\mathrm{i}} \mathrm{H}^{o} \mathrm{H}-*-{ }^{11} \mathrm{C}(\pi)\right)$ is predicted to have the nature of $\left(\mathrm{p}-\mathrm{CS} / t-\mathrm{HB}_{\mathrm{nc}}\right)$. The $\pi$-EB interaction in the $\mathrm{H}_{2} \mathrm{O}-{ }^{-11} \mathrm{C}(\pi)$ type $\left[\left(\theta, \theta_{\mathrm{p}}\right)=\left(73.9^{\circ}, 79.3^{\circ}\right)\right]$ is predicted to be somewhat stronger than the $\pi-\mathrm{HB}$ interaction of the $\mathrm{OHH}-*-{ }^{11} \mathrm{C}(\pi)$ type $\left[\left(\theta, \theta_{\mathrm{p}}\right)=\left(72.6^{\circ}, 74.5^{\circ}\right)\right]$ for $\mathrm{OH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$. The $\pi$-HB interactions in the anthracene system are predicted to be very similar to that in the naphthalene system, which seem slightly stronger than that in the benzene system. The results are in accordance with that derived from the energies for the formation of the adducts.

The delocalization indexes and ellipticity are the important parameters to clarify and understand the nature of the $\left(\mathrm{EH}_{2}\right) \cdots \pi$ interactions. ${ }^{47}$ The nature of the $\left(\mathrm{EH}_{2}\right) \cdots \pi$ interactions will be discussed elsewhere based on these parameters for the series of $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right),\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{10} \mathrm{H}_{8}\right)$, and $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=$ $\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te ).

## Conclusions

The behaviour of the interactions for $\mathrm{EH}_{2}$ adducts with the anthracene $\pi$-system ( $\mathrm{E}=\mathrm{O}, \mathrm{S}, \mathrm{Se}$ and Te ) in a $1: 1$ ratio is elucidated by applying QTAIM-DFA. The structures were optimized with MP2/BSS-F. Two types of structures were optimized for $\mathrm{E}=\mathrm{O}$ and S , whereas one was optimized for $\mathrm{E}=\mathrm{Se}$ and Te . The BCPs are clearly detected in the molecular graphs. The interactions are the $\pi$-HB type for $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIA}_{\text {Atc }}\right)(\mathrm{E}$ $=\mathrm{O}$ and S$)$ and $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIB}_{\mathrm{Atc}}\right)(\mathrm{E}=\mathrm{S}$, Se and Te $)$, whereas they are the $\pi$-EB type for $\mathrm{H}_{2} \mathrm{O}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)$. The QTAIM-DFA parameters of $(R, \theta)$ and $\left(\theta_{\mathrm{p}}, \kappa_{\mathrm{p}}\right)$ are calculated for the species, according to eqn (S3)-(S6) $\ddagger$ The $\theta$ values are less than $90^{\circ}$ for all interactions examined in this work. The $\theta_{\mathrm{p}}$ values are also less than $90^{\circ}$ for all the interactions in question, except for ${ }^{\mathrm{i}} \mathrm{H}-*{ }^{11} \mathrm{C}(\pi)$ in $\mathrm{TeH}_{2}-*-\mathrm{C}_{14} \mathrm{H}_{10}\left(C_{1}: \mathrm{IIB}_{\text {Atc }}\right)$. Therefore, all the $\pi$-HB and $\pi$-EB interactions examined in this work are classified by pure-CS interactions and characterized as vdW in nature ( $\mathrm{p}-\mathrm{CS} / \mathrm{vdW}$ ), although ${ }^{\mathrm{i}} \mathrm{H}-*{ }_{-}^{11} \mathrm{C}(\pi)$ in $\mathrm{TeH}_{2}-*-\mathrm{C}_{14} \mathrm{H}_{10}\left(C_{1}\right.$ : $\mathrm{IIB}_{\mathrm{Atc}}$ ) is predicted to have the $\mathrm{p}-\mathrm{CS} / t-\mathrm{HB}$ nature without covalency. The $\pi$-EB interaction is predicted to be somewhat stronger than $\pi$ - HB in $\left(\mathrm{OH}_{2}\right)-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$. It is demonstrated that the predicted nature of $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ is closer to that of $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{10} \mathrm{H}_{8}\right)$ rather than that of $\mathrm{EH}_{2}-*-\pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$. The $\pi$-HB and $\pi$-EB interactions in the anthracene $\pi$-system are well elucidated by applying QTAIM-DFA.

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31 Interactions would be easily imaged by means of QTAIM if they can be defined as the corresponding bond paths (BPs), especially for experimental chemists. However, it is demonstrated that the detection of the BPs between two atoms in a molecule emerging from natural alignment of the gradient vector field of the one-electron density of a molecule is neither necessary nor a sufficient condition for the presence of a chemical bond between those atoms. ${ }^{48}$ In this connection, it is pointed out that the terms line paths (LPs) and line critical points (LCPs) should be used in place of BPs and BCPs, respectively. ${ }^{48 b}$ Consequently, the dynamic and static nature in this work should be regarded as the investigation performed at LCPs on LPs corresponding to the $\mathrm{EH}_{2}-\pi$ interactions in the anthracene- $\pi$ system. Nevertheless, BPs and BCPs are used in this paper for the discussion of the $\mathrm{EH}_{2}-\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ interactions, similarly to the case of the previous discussion for the $\mathrm{EH}_{2}-\pi\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$ and $\mathrm{EH}_{2}-\pi\left(\mathrm{C}_{10} \mathrm{H}_{8}\right)$ interactions. The interactions expected between $\mathrm{EH}_{2}$ and
$\pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)$ are clearly detected by BPs with BCPs which is another reason to use BPs and BCPs in this work.
32 Critical points (CPs) are characterized by the rank ( $\omega$ ) and the signature $(\sigma)$. The CPs of the species in the threedimensional space are classified by $\omega=3$, which generally corresponds to all species. On the other hand, $\sigma$ is defined by the simple algebraic sum of the signs of $\partial^{2} \rho_{\mathrm{b}}\left(r_{\mathrm{c}}\right) / \partial r_{i}^{2}\left(r_{i}\right.$ $=x, y$ and $z$ for $i=1,2$ and 3 , respectively), where the + and - signs of $\partial^{2} \rho_{\mathrm{b}}\left(r_{\mathrm{c}}\right) / \partial r_{i}^{2}$ are counted as +1 and -1 , respectively. Therefore, $\sigma=-3,-1,1$ and 3 correspond to attractors (nuclei), bond critical points (BCPs), ring critical points (RCPs) and cage critical points (CCPs), respectively. Namely, BCP is characterized by $(\omega, \sigma)=(3,-1) .{ }^{17,18}$
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39 The concept of a dynamic molecular graph was recently proposed by Corte's-Guzma'n and co-workers through the investigation of the Born-Oppenheimer molecular dynamics (BOMD), which was exemplified by $[\mathrm{Fe}\{\mathrm{C}-$ $\left.\left.\left(\mathrm{CH}_{2}\right)_{3}\right\}(\mathrm{CO})_{3}\right]$. The investigation illustrates the change in the behaviour of the molecular graph. ${ }^{49}$ The concept of the dynamic molecular graph would be closely related to that of the dynamic nature of interactions predicted by employing the perturbed structures generated with the normal coordinate of internal vibrations (NIV). In this treatment, the selected vibration for NIV must contain the motion of the interaction in question most effectively among all the zero-point internal vibrations. A structural catastrophe is confirmed not to occur for the perturbed and fully optimized structures for the elucidation of the dynamic nature of the interaction in question with NIV. ${ }^{35}$.
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    $\dagger$ Dedicated to Professor Marian Mikołajczyk (Professor at the Centre of Molecular and Macromolecular Studies, Polish Academy of Sciences in Lodz, Poland) on the occasion of his 80th birthday.
    $\ddagger$ Electronic supplementary information (ESI) available: QTAIM-DFA approach, Cartesian coordinates for optimized structures of $\left(\mathrm{EH}_{2}\right) \cdots \pi\left(\mathrm{C}_{14} \mathrm{H}_{10}\right)(\mathrm{E}=\mathrm{O}, \mathrm{S}$, Se and Te). See DOI: 10.1039/c7ra04224h

