RSC Advances



View Article Online

View Journal | View Issue

PAPER

Check for updates

Cite this: RSC Adv., 2017, 7, 29042

Received 11th April 2017 Accepted 26th May 2017 DOI: 10.1039/c7ra04120a

rsc.li/rsc-advances

1. Introduction

For decades, the corrosion of steel has seriously hindered the development of industrial economy, and even resulted in major accident and bodily injury. Many methods have been used to decrease the corrosion of steel, such as surface coating technology, minimizing the concentration of oxygen, and adding alloyed elements. Additionally, to prevent disasters, Li and colleagues call for open data infrastructures to collate information on material failures.1 The use of organic molecules as corrosion inhibitors has become an increasingly popular strategy for avoiding or reducing the metal corrosion, especially for steel in acidic medium. It is generally accepted that these inhibitors prevent the corrosion by adsorbing onto the metal surface and forming a protective barrier.²⁻⁴ Undoubtedly, a knowledge of the surface chemistry of the adsorbed organics can have profound implications for designing new inhibitors. As reviewed by Dwivedi and his coauthors,⁵ the surface

Toward understanding the adsorption mechanism of large size organic corrosion inhibitors on an Fe(110) surface using the DFTB method⁺

Lei Guo, ⁽¹⁾*^a Chengwei Qi,^b Xingwen Zheng,^c Renhui Zhang,^a Xun Shen^a and Savaş Kaya^d

One of the effective methods developed to inhibit the corrosion of steel is the use of organic molecules as corrosion inhibitors. In particular, the design and synthesis of large size organic corrosion inhibitors draws more and more attention. Unfortunately, an atomic-level insight into the inhibition mechanism is still lacking, and regular density functional theory method is found to be inefficient in dealing with large inhibitor–metal adsorption systems. Given this background, in this work, density functional based tight binding (DFTB) approach was used to investigate the adsorption properties of three large size inhibitors (*i.e.*, chalcone derivatives) on an iron surface. The molecular activity of free chalcone derivatives was studied by means of Frontier molecular orbital theory. The growth characteristics of α -Fe habits were observed using the "Morphology" software. Some adsorption parameters such as charge density difference, density of states, and changes of molecular orbital were described in detail. The present study is helpful to understand the anticorrosive mechanism of similar organic inhibitors and provides a feasible way to develop novel corrosion inhibitors.

analytical techniques, which were used for characterization of steels in corrosive media especially when the presence of corrosion inhibitors, receive more and more attentions. Actually, it is far from perfect to only rely on experiments, and there are still some difficulties in imaging the distribution of molecular orientation with conventional experimental techniques on the nanoscale.⁶ Fortunately, molecular simulation techniques can make a beneficial supplement to the inhibitor-metal interface study. An increasing number of scientists have turned to molecular simulation technique for understanding the anticorrosive mechanism of organic inhibitors.

It is apparent from the studies in the literature, many researchers7-9 investigated the adsorption of organic inhibitors on metal surface employing molecular dynamics simulation method by Forcite or Lammps software, which is based on force field, such as COMPASS,10 cvff,11 Dreiding.12 However, one can only perform qualitative analysis by this approach, their reliability in the modeling of the adsorption of organic molecules is limited. Force fields are not flexible to various chemical situations or to predictions of spectroscopic data which rely quantitatively on the involved knowledge of electronic structure. We hold that more attention should be paid to how these organic inhibitors bind to metal surfaces, and why sometimes minor modifications in the geometry structures can have a great influence on their efficacy. On the other hand, density functional theory (DFT) enables very excellent approximations for the complicated components of electronic motion called exchange and correlation.13 DFT that considers the electron

^aCollege of Material and Chemical Engineering, Tongren University, Tongren 554300, China. E-mail: cqglei@163.com; Tel: +86-0856-5592925

^bSchool of Petroleum Engineering, Chongqing University of Science and Technology, Chongqing 401331, China

^cMaterial Corrosion and Protection Key Laboratory of Sichuan Province, Zigong 643000, China

^dCumhuriyet University, Faculty of Science, Department of Chemistry, Sivas 58140, Turkey

[†] Electronic supplementary information (ESI) available. See DOI: 10.1039/c7ra04120a



Fig. 1 A comparative schematic for small and large size organic inhibitors adsorption on metal substrate.

density to explain the reactivity or stability is helpful to develop novel corrosion inhibitor molecules because it has already proven to be very helpful in quantitatively elucidating the structure–activity relationships.¹⁴ Unfortunately, DFT can only handle small to medium sized organics–metal adsorption systems since it usually accompanies the complex and timeconsuming calculations. For example, Kokalj^{15–17} has systematically studied the adsorption of azoles on bare or oxidized copper surfaces using Quantum Espresso package. In our previous work, we made some investigations on small molecule adsorbate–metal substrate system by Dmol³ code.^{18,19} Despite the success of DFT in terms of the explanation of reactivity or stability of chemical compounds, accurate DFT calculations on molecules including hundreds of atoms are very expensive and time-consuming.

Many studies showed that large size organic corrosion inhibitors tend to present better inhibiting property than small ones, which can be attributed to that the former usually have more active site relative to the latter.^{20,21} They are able to produce a multi-anchoring adsorption and exhibit a high geometric covering effect, as displayed in Fig. 1. Therefore it is very necessary to introduce a new and unusual approach to solve the adsorption problem of large size inhibitors.

In this research, we focused on the adsorption of some large size organic corrosion inhibitors on iron surface using a relatively simple method, *i.e.*, DFTB (see below for description). The main objective of this study is to obtain a qualitative description of these adsorption systems and a better understanding of the anticorrosive mechanism. Then, the most stable adsorption configurations, density of states, Mulliken charge as well as charge density difference will be theoretically analyzed.

2. Methods and computational details

The density functional based tight binding (DFTB) method is based on a second-order expansion of the Kohn–Sham total energy in DFT with respect to charge density fluctuations. The zeroth order approach is equivalent to a common standard nonself-consistent scheme, while at second order a transparent, parameter-free, and readily calculable expression for generalized Hamiltonian matrix elements can be derived. The Coulombic interaction between partial atomic charges was determined by the self-consistent charge (SCC) formalism.²² It considerably reduces the computational cost, and offers a high degree of transferability as well as universality for both groundstate and excited-state properties. DFTB operates with the comparative efficiency and accuracy for organic matters, insulators, solids, clusters, semi-conductors, and metals, even biological systems are involved.²³

As seen in Fig. 2, the selected large size organic corrosion inhibitors are three chalcones derivatives, i.e., 3-(4-(dimethylamino)phenyl)-1-(4-hydroxyphenyl)prop-2-en-1-one (Inh1), 3-(3-hydroxyphenyl)-1-(4-hydroxyphenyl)prop-2-en-1-one 1-(4-hydroxyphenyl)-3-phenylprop-2-en-1-one (Inh2), and (Inh3), which have been identified as good inhibitors for carbon steel corrosion in hydrochloric acid solution.24 According to the electrochemical impedance spectroscopy experiment, for the same inhibitor concentration, the optimal inhibitive efficiency of Inh1, Inh2, and Inh3 were 91.3%, 87.0%, and 85.9%, respectively.24 The iron surface model consisted of a slab of four iron atomic layers. The unit cell was extended to a 7×7 supercell in x- and y-directions. Finally, a vacuum slab with 30 Å thickness was built on the iron surface. It is particularly worth mentioning here that we used the "Morphology" module²⁵ in Materials Studio software (Version 7.0) to investigate the iron substrate in order to choose an appropriate adsorption surface.

All other computations were performed with the DFTB+ program package of Materials Studio. We used the trans3d Slater-Koster library set for all possible pair interactions between C, N, O, H, and Fe. Convergence tolerance for geometry optimization was 0.01 kcal mol⁻¹ for energy, 0.1 kcal mol⁻¹ Å⁻ for force, and 0.001 Å for displacement. The relative SCC tolerance was set to 10^{-8} au elementary charge. To reduce the number of SCC iterations, Broyden mixing²⁶ is used. Besides, fractional orbital occupation with a small thermal smearing parameter (0.001 Ha) was used to speed up SCC convergence. The k-points of $4 \times 4 \times 1$ was selected. During the calculations, the two bottom layers of the slab were kept frozen to the ideal crystalline position, while the upper half of the slab and the inhibitor molecules were allowed to fully relax. Our DFTB calculation yields a lattice parameter of 2.899 Å for bulk bcc Fe, which agrees well with the experimental data (2.866 Å).²⁷

To calculate the adsorption energies E_{ads} , for all systems, the following equation was used as defined in our previous work:^{18,19,28}

$$E_{\rm ads} = E_{\rm complex} - (E_{\rm Fe} + E_{\rm inh}) \tag{1}$$

where E_{complex} is the total energy of the iron slab and adsorbed inhibitors, E_{Fe} the energy of a clean iron slab, and E_{inh} the



Fig. 2 Molecular structures for the selected large size organic corrosion inhibitors.

energy of an isolated inhibitor molecule. With this definition, a negative value of E_{ads} corresponds a stable adsorbate-substrate system.

Obviously, there usually exists more complicated circumstances in the real adsorption situation, for example, the iron surface can be positively charged in acid medium (charge effects),^{29,30} the substrate may be defective, stepped or oxidized (surface geometry effects),^{31,32} as well as that there exists solvent effect. It should be noted that this work is just a primary inquiry on adsorption issues of large size inhibitors. We just performed the simulations at vacuum/metal interface. Unquestionably, the calculations in the vacuum condition cannot perfectly reflect the real environment, but it is suitable for qualitatively illustrating the trend of the chemical interactions of interest. Some other factors will be gradually taken into account in our future work.

3. Results and discussion

3.1. Free chalcones derivatives

To identify the active characteristics of individual inhibitor molecule, the DFTB calculations for the chalcones derivatives in free forms were firstly performed. It is generally acknowledged that the bonding interaction between inhibitor molecules and metal surface depends on the Frontier orbital energetic position and the Fermi energy of metal.³³ Fig. 3 shows the Frontier orbital energies for Inh1, Inh2, and Inh3 molecules.

By careful examination of Fig. 3, it could be noticed that all the inhibitors exhibit very similar values of $E_{\rm HOMO}$ (the energy of the highest occupied molecular orbital) and $E_{\rm LUMO}$ (the energy of the lowest unoccupied molecular orbital). The $E_{\rm HOMO}$ values are much higher than the Fermi level of iron (-13.16 eV, obtained by DFTB approach), and then the electrons can easily flow from the HOMO orbitals of chalcones derivatives to the empty iron orbitals until an equilibrium state was reached. As a global electronic index, the energy gap $\Delta E (E_{\rm LUMO} - E_{\rm HOMO})$ is often adopted to express the chemical reactivity of the inhibitor molecules toward the substrate surface. However, there was some controversy about this descriptor. Some researchers think that a low ΔE value usually corresponds to a high inhibition efficiency in the inhibitor molecule.^{14,34} Conversely, some



Fig. 3 Molecular Frontier orbital energetic positions for three chalcones derivatives.

authors point out that the analysis of molecular electronic structure of organic inhibitors cannot straightforwardly predict the trend of their inhibition performance.^{35,36} Our calculated ΔE values are 2.56 eV, 2.57 eV, and 2.46 eV for Inh1, Inh2, and Inh3 respectively, in which the order is not in good agreement with experimental inhibitive efficiency mentioned above.²⁴ Obviously, these results are consistent with the latter opinion. In fact, the factors that affect inhibition efficiency are complex, for instance, the molecule absorption orientation is sensitive to the dipole–dipole interaction.³⁷ An approach solely focused on the inhibitors while ignoring the metal substrate is not perfect. This emphasizes the necessity of a rigorous modeling of the interactions between the corrosion inhibitors and the metal surface.

3.2. An investigation of iron surfaces

Crystal morphology study for α -Fe was performed in Morphology module using the equilibrium morphology model. It is known that the equilibrium morphology of a crystal is determined by the minimum of the surface free energy for a given volume and temperature.³⁸ If the surface free energies are known for all relevant crystal faces, the morphology of a crystal in equilibrium with its surroundings can be visualized using the Wulff construction.³⁹ In Morphology, the surface energy (E_{surf}) is calculated from the energy of a slab of finite thickness:

$$E_{\text{surf}} = \frac{1}{2} \lim_{M \to \infty} \frac{[E_{\text{latt}}(M) - E_{\text{slice}}(M)]}{A_{hkl}}$$
(2)

where $E_{\text{latt}}(M)$ is the energy of a slab M layers thick inside the infinite crystal, $E_{\text{slice}}(M)$ is the energy of a slab M layers thick in a vacuum, and A_{hkl} is the surface area of a plane with Miller indices (hkl). The factor 1/2 in accounts for the fact that the slab has two surfaces. Another basic recipe is that:

$$D_{hkl} = kE_{\text{surf}(hkl)} \tag{3}$$

where D_{hkl} is the distance from the origin of the coordinate system to the face (*hkl*). It is proportional to the surface energy (*k* is a constant) of the face.

Calculated morphology parameters for α -Fe crystal using the equilibrium morphology method are given in Table 1. The symmetry multiplicities for the (110), (100) and (111) growth forms are 12, 6, and 8, respectively. The interplanar distance decreases gradually from (110) to (111) crystal face. The calculated surface energies (E_{surf}) , which mostly originate from van der Waals force, decrease in the following sequence: (111) > (100) > (110), and this is in good agreement with reported trend by Arya and Carter.⁴⁰ We hold that magnitude difference is acceptable, resulting from the different calculation methods. We also notice that the total facet area (TFA) of Fe(110) accounts for more than 63% of the crystal surface. The clean (110) and (100) surfaces possess only one type of surface atom with coordination numbers (N for short) 6 and 4, respectively, while the (111) surface has two kinds of surface atoms (where N = 4 or 6).

For demonstration purposes, the habit information and three typical surface structures of α -Fe were given in Fig. 4.

hkl	Multiplicity	$d_{\mathrm{hkl}}{}^{a}\left(\mathrm{\AA} ight)$	$A_{hkl}({ m \AA}^2)$	$E_{\rm surf}$ (J m ⁻²)	$D_{hkl}({ m \AA})$	$\mathrm{TFA}^{c}\left(imes 10^{8}\mathrm{\AA}^{2} ight)$	%TFA ^d	Ν
(110)	12	2.02	5.81	$6.36, 2.28^b$	9159.66	7.98	63.31	6
(100)	6	1.43	8.21	$6.50, 2.30^{b}$	9373.48	3.07	24.38	4
(111)	8	0.82	14.23	$6.83, 2.66^b$	9852.12	1.55	12.29	4,6

^{*a*} Interplanar distance. ^{*b*} Ref. 40, obtained from DFT-GGA approach. ^{*c*} The total area of all symmetry images of the facet. ^{*d*} The percentage of the total habit surface area occupied by all symmetry images of the facet.



Fig. 4 (a-c) Top and side views of three iron surface models, (d) Wulff plot showing the equilibrium crystal morphology based on the relaxed surface formation energies. Balls in pink color denote the atoms at the first surface layer.

Obviously, the Fe(110) surface shows a density packed structure, it has a lower surface energy but a greater surface atom coordination number. Therefore, the Fe(110) surface was chosen to address the adsorption of chalcones derivatives, which is primarily because that it has more contact sites to interact with corrosion inhibitors. Actually, Fe(110) surface is also chosen as model for exploring the anticorrosive mechanism of corrosion inhibitors in conventional molecular dynamics simulations.^{41,42} But this morphology method can have an important significance for the determination of oxidized iron surfaces such as α -Fe₂O₃, γ -Fe₃O₄, and FeCO₃.

3.3. Adsorption of chalcones derivatives on the Fe(110) surface

Firstly, one can argue that the nitrogen or oxygen atoms in the chalcones derivatives of interest can be protonated in acid media. Thus, as an example, we did the adsorption study of Inh3H (the protonated form of Inh3 molecule) on the Fe(110) surface. The adsorption configurations of Inh3H before and after the geometry optimization are given in Fig. S1 (ESI data†). We found that there existed a proton transfer process form the nitrogen atom to nearby carbon atom, and compared to the neutral adsorption, a slightly lower adsorption energy (-1.84 eV) was obtained. So the protonated state is metastable, it could be existed unless there are negatively-charged particles such as chloride ion. From this prospective, in this section, we just elaborate a discussion for the adsorption behaviour of inhibitors in neutral forms.

3.3.1. Adsorption geometries and energies. The optimized adsorption configurations of the inhibitors on the Fe(110) surface are shown in Fig. 5. We can see that that all three inhibitors are adsorbed on the Fe(110) surface with a nearby flat orientation. This parallel adsorption can be considered as the result from donor-acceptor type interaction between the inhibitor molecules and metal surfaces. All the three inhibitors contains N, O donor heteroatoms as well as π -electrons in their structures to provide large numbers of electrons to the vacant dorbitals of iron. On the other hand, π -antibonding orbitals of the phenyl rings have sufficient accommodation to accept electron from the 4s or 3d orbitals of iron to form feedback bonds. This parallel configuration is very common in many large size organic corrosion inhibitors, such as reported in ref. 43–45, in which the π –d hybridization plays a dominant role. However, some small inhibitor molecules (benzimidazole for instance) can also chemisorb perpendicularly to the iron surface with unsaturated heteroatoms through σ-molecular orbitals.46

Table 2 lists some geometric parameters and adsorption energy values. The results show that the E_{ads} values of three adsorption systems are negative, which reveals the adsorption is energetically favorable. Moreover, the stability of adsorption system can be measured by the magnitude of E_{ads} , that is, the bigger the absolute value of E_{ads} was, the stronger the inhibitor– Fe(110) interactions and the higher inhibition efficiency were. It is clear from Table 2 that the Inh3@Fe(110) adsorption system is more stable than the other two system and then exhibited the



Fig. 5 Equilibrium adsorption configurations performed for three chalcones derivatives on Fe(110) surface. Top: top view, down: side view.

Table 2 Calculated geometric parameters, adsorption energies (E_{ads}), as well as the experimentally measured corrosion rate (CR) and inhibition efficiency (IE)

Inhibitor@Fe(110)	Bond lengths ^{<i>a</i>} (Å)		$\Delta {h_{{ m Fe}}}^b \left({ m \AA} ight)$	$d_{\mathrm{inh-Fe}}^{c}(\mathrm{\AA})$	$E_{\rm ads}$ (eV)	$CR^{d} (mmy^{-1})$	IE^d %
Inh1	R _{O-Fe}	1.925	0.195	2.620	-1.35	38.7	82.5
Inh2	R _{O-Fe}	1.927	0.238	2.573	-1.68	29.2	86.9
Inh3	$R_{ m O-Fe}$ $R_{ m N-Fe}$	1.943 2.269	0.295	2.514	-1.94	20.9	90.6

 a R_{O-Fe} and R_{N-Fe} represent the minimum N/O–Fe distance. b Δh_{Fe} , maximum vertical displacement of the surface Fe in contact with the inhibitor. c d_{inh-Fe} , the vertical distance from the centroid of inhibitor molecule to the first layer iron atoms. d Expt. ref. 24.

highest inhibition efficiency, which efficiently reduced the corrosion rate of carbon steel.²⁴ The minimum N/O–Fe bond lengths are in the range between 1.9–2.2 Å, which is comparable to the sum of N/O and Fe atomic covalent radii ($r_{cov}^{N} + r_{cov}^{Fe} = 0.71 + 1.32 = 2.03$ Å, $r_{cov}^{O} + r_{cov}^{Fe} = 0.66 + 1.32 = 1.98$ Å),⁴⁷ so this indicates that there the exists chemisorption between the studied inhibitors and Fe(110) surface.

Since the modification of inhibitors to iron substrate, the iron atoms in the first layer relax outward with respect to the computed bulk-terminated geometry. The maximum vertical displacement, $\Delta h_{\rm Fe}$, increased gradually from Inh1 to Inh3. While the vertical distance (d_{inh-Fe}) decreased in the sequence Inh1 > Inh2 > Inh3. These characteristics can be attributed to the number of N/O heteroatoms in the molecular structure. Owing to the oxygen atom has a larger electronegativity (χ) value than nitrogen atom (i.e., $\chi_{\rm O}=$ 3.44, $\chi_{\rm N}=$ 3.04),⁴⁸ then as depicted in Fig. 6, its outer electrons are somewhat firmly held around the oxygen atom. Looked at another way, according to the Hard and Soft Acid and Bases (HSAB) theory, soft interacts better with soft. As bulk metals are chemically the softest materials, this would imply that the smaller is the electronegativity of coordination atoms the stronger is their interaction with metal surfaces. Consequently, oxygen has a slightly weaker electron-donating ability than nitrogen atom. This can be seen the main reason why Inh3 has a better inhibiting property than Inh2.

3.3.2. Charge density difference and population analysis. In order to gain an insight into the electronic properties of the most stable relaxed inhibitor@Fe(110) interfaces, we have examined the charge density difference, $\Delta \rho$, which is defined as:⁴⁹

$$\Delta \rho = \rho_{\text{total}} - (\rho_{\text{inh}} + \rho_{\text{Fe}}) \tag{4}$$



Fig. 6 Schematic diagram of the electron-donating ability comparison between O and N elements.

Paper



Fig. 7 Charge density difference for the adsorption systems of chalcones derivatives on Fe(110). The blue region denotes electronic accumulation, while the red region denotes electronic loss.

where ρ_{total} , ρ_{Fe} , and ρ_{inh} are the total density difference of the system, the clean surface of iron, and the isolated chalcones derivative molecule, respectively. The charge redistribution for the most stable configuration of the parallel adsorption is shown in Fig. 7, which depicts the density difference of several adsorbate/Fe(110) systems with a cutoff value of ± 0.002 e Å⁻³. Notably, there exist clear charge increases in the first layer iron atoms. The aforementioned π -electron–iron interactions are also visible. The charge redistributions between the inhibitor molecules and iron atoms seem to be caused by the electron-donating effects of the electron rich groups.

We also made quantitative analysis about the changes of Mulliken charge for some typical atoms between the free and adsorbed chalcones derivatives, and the results are collected in Table 3. Based on the tabulated atomic charges, a striking feature is that the charges for O1 atoms (in carbonyl group) increase while decrease for other atoms (such as O2, O3), especially N atom in the dimethylamino group. This confirms previous conclusion that N atom possesses greater electrondonating ability.

3.3.3. Projected density of states. To illustrate the effects of adsorption of inhibitor molecules on the electronic properties of Fe(110), the projected density of states (PDOS) for clean Fe(110) and various adsorption systems are analyzed. Gaussian broadening scheme of width was set to be 0.1 eV, and the Fermi level is set to zero for all of the DOS plots.

In Fig. 8, we show the total and partial density of states for clean Fe(110) surface with 7 \times 7 geometry. We observe that the total density of states (TDOS) are mainly dominated by the

 Table 3
 Mulliken charge population analysis for the free and adsorbed inhibitor molecules

	Inh1		Inh2		Inh3		
Atoms	Free	Adsorbed	Free	Adsorbed	Free	Adsorbed	
01	-0.419	-0.479	-0.419	-0.489	-0.419	-0.494	
O2	-0.452	-0.382	-0.452	-0.387	-0.452	-0.381	
O3			-0.452	-0.382			
N					-0.384	-0.065	

cooperative contributions of Fe s, p, and d orbitals, among which Fe 3d is dominant. The Fe 3d state is distributed mainly in energy ranges from -5.0 to 4.8 eV. By comparison with previous Özcan and co-worker's reports, we find that the whole TDOS profile is in fairly good agreement with their calculated curve,⁵⁰ and the locations of several main peaks are also basically consistent. Apparently, Fe has a large density of d-states at the $E_{\rm F}$, then it can be regarded as a more reactive transition metal with d-states partially occupied. This makes it much easier for the chalcones derivatives bonding with π -system to iron surface.

Fig. 9 shows the PDOS of the inhibitor@Fe(110) systems before and after molecule-surface interaction. The top panel shows the PDOS for the molecule located 7 Å above the surface, *i.e.*, in an unbound configuration, while the lower one shows the adsorbed state. Obviously, the plots are similar for all examined derivatives. The top figure can be used to present the alignment of the high-lying molecular valence states with the Fe 3d-band before the interaction occurs. It is evident that several molecular states (sharp peaks in the -4.0 to 0 eV regions) lie at the position of the Fe 3d band, which can be seen to describe the molecule-surface bonding upon adsorption. When the chalcones derivatives adsorb on the Fe(110) surface, the molecular PDOS is relatively unstructured in the vicinity of the Fermi



Fig. 8 Calculated total and partial density of states for clean Fe(110) surface. Fermi energy (E_F) is adopted as the zero energy level (green vertical line).



Fig. 9 Density of states for the inhibitor@Fe(110) systems projected on the inhibitor molecules and Fe 3d orbitals. The upper panel shows the inhibitor located 7 Å above the surface, whereas it is adsorbed in the lower panel.

energy for three flat configurations, and the intensity significantly decreases. A strong orbital hybridization occurred between the reactive sites in inhibitor molecules and Fe atoms. While the PDOS change for Fe 3d is not evident since we have performed the calculations at low coverages. After interacting with the surface, the molecular PDOS of Inh3 is considerably less spiky compared with Inh1 and Inh2, implying a stronger molecule-surface bond.

3.3.4. Molecular orbital change. The adsorption induced Frontier molecular orbital changes for three chalcones derivatives are depicted in Fig. 10, S2 and S3 (ESI data[†]). An inspection of Fig. 10 shows that the LUMO of Inh3 show little change and hence it has small contribution to the earlier mentioned molecule–surface bond. The HOMO and HOMO–1 are π -type orbitals and they take up nearly the whole molecular skeleton. Their distributions have noticeable changes after adsorbing on the Fe(110) surface, this is consistent with the former DOS analysis.

Further quantitative analysis reveals that the adsorption of inhibitor molecule on Fe(110) surface reduces the energy gap of the chalcones derivatives involved. As shown in Fig. 11, the ΔE of isolated Inh3 molecule is 2.46 eV, whereas in the adsorbed system the ΔE is 1.42 eV. This is a pretty common phenomenon in the organic–metal interfaces. For example, Hahn and Kang⁵¹ reported that the energy gap of pyridine molecule decreases about 0.9 eV while interacting with Ag(110) surface. One way of explaining this is that the charging process of the topmost metal surface appeared to affect the electronic structure of the adsorbate in such a way that its molecular orbitals were fully stabilized by adsorption.

3.4. Mechanism of adsorption and inhibition

Finally, after investigating the adsorption mechanism at molecular/atomic level, a clarification of the anticorrosive mechanism in real environment medium is important. In the medium of hydrochloric acid, due to the specific adsorption



Fig. 10 Molecular orbital plots for the Inh3 molecule from the HOMO-1 to LUMO under the free (upper row) and adsorbed state (lower row). (isovalue: $\pm 0.005 \text{ e} \text{ Å}^{-3}$).



Fig. 11 Energy gap values for three chalcones derivatives in free and adsorbed sates.

effect of chloride ions, the pitting corrosion of iron bulk material can proceed by the following steps:^{29,52}

$$Fe + Cl^- \leftrightarrow (FeCl)_{ads} + e^-$$
 (5)

 $(\text{FeCl})_{\text{ads}} \leftrightarrow (\text{FeCl}^+) + e^-$ (6)

$$\operatorname{FeCl}^+ \leftrightarrow \operatorname{Fe}^{2+} + \operatorname{Cl}^-$$
 (7)

As shown in Fig. 12, for iron specimens in HCl solutions inhibited by the chalcones derivatives, these organic surface modifiers can react rapidly with the surface Fe atoms. Then a strong protective layer (FeInh) was formed in the noncorroded (flat) area, which is very thin and could be a single monolayer.

While in the corroded (low-lying) area, chloride ions may be preferred to adsorb onto the positively charged iron surface. On the other hand, some heteroatoms in the chalcones derivatives, such as nitrogen and oxygen, can be partially protonated (abbreviated as InhH) in a low pH value media. Thus, Cl^- plays the role of a bridge joining the InhH and iron surface, then a thick and protective [FeClInhH] complex was formed. Ultimately, the diffusion of corrosive ions such as H_3O^+ can be blocked by the protective film and then the steel was effectively protected from corrosion.



Fig. 12 Proposed schematic for the adsorption mechanism of inhibitor on iron surface under hydrochloric acid medium.

4. Conclusions

The chalcones derivatives have proven experimentally to be efficient against the corrosion of carbon steel. In the present work, the interactions of three chalcones derivatives and iron surface have been simulated and analyzed by DFTB method. The Inh3 molecule present strong interaction toward Fe(110) surface with larger adsorption energy and smaller interaction distance, compared with Inh1 and Inh2. The hetero atoms in the chalcones derivatives are found to have high electron-donating ability and their inhibition efficiency increases in the order O < N. For all the molecules studied, they can chemisorb onto Fe(110) with their molecular plain being nearly parallel to the surface. The π -molecular orbitals play a decisive role for the adsorption performance.

Acknowledgements

This research was sponsored by the Science and Technology Program of Guizhou Province (QKHJC 2016-1149), the Guizhou Provincial Department of Education Foundation (QJHKYZ2016-105, 2016-009), the Research Fund for the Doctoral Program of Tongren University (trxyDH1510), the National Natural Science Foundation of China (51462030), the Student's Platform for Innovation and Entrepreneurship Training Program (2016106665), and the Opening Project of Sichuan University of Science and Engineering (2016CL06).

References

- 1 X. G. Li, D. W. Zhang, Z. Y. Liu, Z. Li, C. W. Du and C. F. Dong, *Nature*, 2015, **527**, 441–442.
- 2 F. M. Donahue and K. Nobe, J. Electrochem. Soc., 1965, 112, 886–891.
- 3 J. O. Mendes, E. C. da Silva and A. B. Rocha, *Corros. Sci.*, 2012, 57, 254–259.
- 4 H. Lu, S. Zhang, W. Li, Y. Cui and T. Yang, *ACS Appl. Mater. Interfaces*, 2017, 9, 4034–4043.
- 5 D. Dwivedi, K. Lepková and T. Becker, *RSC Adv.*, 2017, 7, 4580–4610.
- 6 E. A. Muller, B. Pollard, H. A. Bechtel, P. van Blerkom and M. B. Raschke, *Sci. Adv.*, 2016, 2, e1601006.
- 7 M. Shahraki, M. Dehdab and S. Elmi, *J. Taiwan Inst. Chem. Eng.*, 2016, **62**, 313–321.
- 8 Z. Wang, Q. Lv, S. Chen, C. Li, S. Sun and S. Hu, ACS Appl. Mater. Interfaces, 2016, 8, 7499–7508.
- 9 M. A. Chidiebere, E. E. Oguzie, L. Liu, Y. Li and F. Wang, *Ind. Eng. Chem. Res.*, 2014, 53, 7670–7679.
- 10 H. Sun, P. Ren and J. R. Fried, *Comput. Theor. Polym. Sci.*, 1998, **8**, 229–246.
- 11 P. Hobza, M. Kabelac, J. Sponer, P. Mejzlik and J. Vondrasek, J. Comput. Chem., 1997, 18, 1136–1150.
- 12 S. L. Mayo, B. D. Olafson and W. A. Goddard, *J. Phys. Chem.*, 1990, **94**, 8897–8909.
- 13 C. K. Skylaris, Science, 2016, 351, 1394-1395.
- 14 G. Gece, Corros. Sci., 2008, 50, 2981-2992.

- 15 S. Peljhan, J. Koller and A. Kokalj, *J. Phys. Chem. C*, 2014, **118**, 933–943.
- 16 S. Peljhan and A. Kokalj, *Phys. Chem. Chem. Phys.*, 2011, **13**, 20408–20417.
- 17 N. Kovačević, I. Milošev and A. Kokalj, *Corros. Sci.*, 2015, **98**, 457–470.
- 18 L. Guo, W. P. Dong and S. T. Zhang, RSC Adv., 2014, 4, 41956-41967.
- 19 L. Guo, I. B. Obot, X. Zheng, X. Shen, Y. Qiang, S. Kaya and C. Kaya, *Appl. Surf. Sci.*, 2017, **406**, 301–306.
- 20 P. Singh, E. E. Ebenso, L. O. Olasunkanmi, I. B. Obot and M. A. Quraishi, *J. Phys. Chem. C*, 2016, **120**, 3408–3419.
- 21 M. Finšgar and J. Jackson, Corros. Sci., 2014, 86, 17-41.
- M. Elstner, D. Porezag, G. Jungnickel, J. Elsner, M. Haugk,
 T. Frauenheim, S. Suhai and G. Seifert, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1998, 58, 7260–7268.
- 23 T. Frauenheim, G. Seifert, M. Elstner, T. Niehaus, C. Kohler, M. Amkreutz, M. Sternberg, Z. Hajnal, A. Di Carlo and S. Suhai, *J. Phys.: Condens. Matter*, 2002, 14, 3015–3047.
- 24 A. S. Fouda, M. A. Elmorsi and A. Elmekkawy, *Afr. J. Pure Appl. Chem.*, 2013, 7, 337–349.
- 25 R. F. P. Grimbergen, H. Meekes, P. Bennema, C. S. Strom and L. J. P. Vogels, *Acta Crystallogr., Sect. A: Found. Crystallogr.*, 1998, 54, 491–500.
- 26 G. Kresse and J. Furthmuller, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**, 11169–11186.
- 27 E. P. Yelsukov, E. V. Voronina and V. A. Barinov, J. Magn. Magn. Mater., 1992, 115, 271–280.
- 28 X. W. Zheng, S. T. Zhang, W. P. Li, L. L. Yin, J. H. He and J. F. Wu, *Corros. Sci.*, 2014, **80**, 383–392.
- 29 S. A. Umoren and M. M. Solomon, J. Ind. Eng. Chem., 2015, 21, 81–100.
- 30 M. Mobin and M. Rizvi, *Carbohydr. Polym.*, 2017, **156**, 202–214.
- 31 J. Radilla, M. Boronat, A. Corma and F. Illas, *Phys. Chem. Chem. Phys.*, 2010, **12**, 6492–6498.

- 32 A. Kokalj, Faraday Discuss., 2015, 180, 415-438.
- 33 B. Hammer and J. K. Norskov, Adv. Catal., 2000, 45, 71-129.
- 34 Y. Gong, Z. Wang, F. Gao, S. Zhang and H. Li, *Ind. Eng. Chem. Res.*, 2015, 54, 12242–12253.
- 35 A. Kokalj, *Electrochim. Acta*, 2010, **56**, 745–755.
- 36 L. Guo, S. Zhu, S. Zhang, Q. He and W. Li, *Corros. Sci.*, 2014, 87, 366–375.
- 37 N. Kovacevic and A. Kokalj, J. Phys. Chem. C, 2011, 115, 24189–24197.
- 38 J. W. Gibbs, Collected works, Longman, New York, 1928.
- 39 D. Winn and M. F. Doherty, AIChE J., 2000, 46, 1348-1367.
- 40 A. Arya and E. A. Carter, J. Chem. Phys., 2003, 118, 8982-8996.
- 41 M. A. Chidiebere, E. E. Oguzie, L. Liu, Y. Li and F. H. Wang, *Ind. Eng. Chem. Res.*, 2014, **53**, 7670–7679.
- 42 L. L. Liao, S. Mo, H. Q. Luo and N. B. Li, *J. Colloid Interface Sci.*, 2017, **499**, 110–119.
- 43 S. K. Saha, A. Dutta, P. Ghosh, D. Sukul and P. Banerjee, *Phys. Chem. Chem. Phys.*, 2016, **18**, 17898–17911.
- 44 N. Ding, X. Chen and C.-M. L. Wu, *Environ. Sci.: Nano*, 2014, 1, 55–63.
- 45 A. Dutta, S. K. Saha, P. Banerjee and D. Sukul, *Corros. Sci.*, 2015, **98**, 541–550.
- 46 N. Kovacevic and A. Kokalj, *Mater. Chem. Phys.*, 2012, 137, 331-339.
- 47 P. Pyykko and M. Atsumi, Chem.-Eur. J., 2009, 15, 186-197.
- 48 R. G. Parr and R. G. Pearson, J. Am. Chem. Soc., 1983, 105, 7512–7516.
- 49 T. A. Baker, C. M. Friend and E. Kaxiras, *J. Am. Chem. Soc.*, 2008, **130**, 3720–3721.
- 50 M. Ozcan, D. Toffoli, H. Ustunel and I. Dehri, *Corros. Sci.*, 2014, **80**, 482–486.
- 51 J. R. Hahn and H. S. Kang, Surf. Sci., 2010, 604, 258-264.
- 52 R. Solmaz, G. Kardas, M. Culha, B. Yazici and M. Erbil, *Electrochim. Acta*, 2008, **53**, 5941–5952.