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Colloidal synthesis and magnetic properties of anisotropic-shaped spinel CuCr₂Se₄ nanocrystals†

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Anisotropic-shaped $CuCr_2Se_4$ nanocrystals have been synthesized by thermal decomposition and reaction of novel mixed metal-oleate complexes with selenium in a high-boiling point organic solvent, trioctylamine (TOA). The synthesized $CuCr_2Se_4$ nanocrystals exhibit close to triangular and hexagonal morphology, with an average size of 20 nm. X-ray diffraction patterns and XPS spectral analysis confirm the formation of the pure spinel phase without any impurities. A possible reaction mechanism is suggested and formation pathways for the triangular and hexagonal shaped $CuCr_2Se_4$ nanocrystals are proposed. Magnetic studies indicate that the anisotropic-shaped $CuCr_2Se_4$ nanocrystals are superparamagnetic near room temperature but exhibit ferromagnetic behavior at lower temperatures, with magnetization values of 31 and 43 emu g⁻¹ at 300 and 5 K, respectively.

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Introduction

Controlling the shape of magnetic nanomaterials has been a subject of intensive research in recent years. This is primarily motivated by the fact that the physicochemical and magnetic properties of the nanomaterials are highly dependent on the particle morphology and have important applications in catalysis, recording media, biomedical imaging, magneto-optics, etc. 1-4 Solution-based methods are primarily utilized to synthesize inorganic nanocrystals with controlled shape, size, composition, and structure. A number of approaches have been reported for preparing magnetic nanomaterials (transition metals, metal alloys, and metal oxides) with well-controlled shapes, which generally involve the thermal decomposition of metallic precursors, electrochemical reduction, and chemical reduction.5-8 Nevertheless, the synthesis and properties of magnetic chromium-based chalcogenide nanoparticles with anisotropic shapes remain largely unexplored. As a remarkable class of magnetic materials with diverse electronic transport properties (metallic, semiconducting, and insulating), chromium-based chalcogenide spinels (chalcospinels) ACr₂X₄ (A = Cu, Cd, Hg, Fe, and Co; X = S, Se, and Te) have gained significant attention in recent years.9-12 However, due to complications related to control of the oxidation state of chromium and the limited known nonvolatile chalcogen precursors for selective synthesis of the spinel phase, the exploration of novel routes for preparing Cr-based chalcospinel nanomaterials with desired morphology remains a challenge.¹³

An important class of magnetic chalcospinels, CuCr₂S₄, CuCr₂Se₄, and CuCr₂Te₄, are ferromagnetic metals with Curie temperatures (T_c) of 377, 430, and 360 K, respectively, and are also metallic. Because of their high Curie temperature, these materials are potentially attractive candidates for spin-based electronics.14 They have also been reported to exhibit a pronounced room temperature magneto-optical Kerr effect. 15 Moreover, based on band structure calculations, we predicted the exciting possibility of inducing half-metallicity in some of the cation-substituted Cu, Cd chalcospinels (CuxCd1-xCr2-S(Se)₄), which makes them promising candidates for spintronic applications.16,17 Because of these unique characteristics, there have been several efforts in developing synthetic routes for obtaining CuCr₂S(Se,Te)₄ nanocrystals. ^{18,19} In previous studies, we reported on the colloidal syntheses of various ternary (CuCr₂S(Se,Te)₄, CdCr₂S₄, CuSbS₂, and CuInS₂) and quaternary $(Culn_xGa_{1-x}S(Se)_2 \text{ and } Cu_2ZnSnS_4) \text{ chalcogenide nanocrystals}$ with tunable size and shape control. 18-23 However, for the magnetic Cr-based chalcospinels, the morphologies of the kinetically favorable nanocrystals are mostly spherical or cubic, with the exposed surfaces not being specific crystal planes that can significantly influence their physicochemical and magnetic properties. The preparation of magnetic Cr-based chalcospinel nanomaterials with anisotropic shapes thus remains rather limited.

One of the more efficient routes for the synthesis of anisotropic-shaped magnetic nanocrystals involves thermal decomposition of organometallic precursors, such as metal

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acetylacetonates, metal carbonyls, *etc.*, in high-boiling point solvents by hot-injection in the presence of suitable surfactants. This procedure has been utilized for the synthesis of a variety of magnetic metals, magnetic oxides, and some non-chromium-based chalcogenide nanoparticles with good crystallinity and anisotropic shapes. 24,25 Particularly noteworthy is the recent development of inexpensive, environmental-friendly metal–oleate complexes, which have been shown to be promising precursors for ultra-large-scale thermolysis synthesis of magnetic nanocrystals. 26 Herein, we report on a simple, highly efficient synthesis route for preparing single-phase anisotropic-shaped ${\rm CuCr_2Se_4}$ nanocrystals. The process involves the thermal decomposition and reaction of novel mixed metal–oleate complexes with selenium in a high-boiling point organic solvent trioctylamine (TOA). We have further investigated in

detail the structural and magnetic properties of the as-

synthesized triangular and hexagonal CuCr₂Se₄ nanocrystals.

Experimental

Materials

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Copper(II) chloride dihydrate (CuCl $_2 \cdot 2H_2O$, 99.99%), sodium oleate (C $_{18}H_{33}O_2Na$, 99.99%), elemental selenium (Se, 99.99%), and trioctylamine (C $_{24}H_{51}N$, 90%) were purchased from Aladdin Chemical and used without further purification. Chromium(III) chloride hexahydrate (CrCl $_3 \cdot 6H_2O$, 99.99%) was obtained from Xi Long Chemical. Absolute ethanol, deionized water, and hexane were supplied by Aladdin Chemical as commercially available reagents and used without further purification.

Preparation of metal-oleate complexes $Cu(C_{18}H_{33}O_2)_2$ and $Cr(C_{18}H_{33}O_2)_3$

Synthesis of the metal–oleate $Cu(C_{18}H_{33}O_2)_2$ and $Cr(C_{18}H_{33}O_2)_3$ complexes followed the procedure reported by Park et~al., 26 by reaction of sodium oleate and Cu^{2+} or Cr^{3+} chlorides. In a typical synthesis, 5 mmol of $CuCl_2 \cdot 2H_2O$ (5 mmol of $CrCl_3 \cdot 6H_2O$), 10 mmol (15 mmol) of sodium oleate $(C_{18}H_{33}O_2Na)$, 50 mL of H_2O , 50 mL of ethanol, and 80 mL of hexane were mixed and refluxed at 60 °C for 3 h. The metal–oleate $Cu(C_{18}H_{33}O_2)_2$ and $Cr(C_{18}H_{33}O_2)_3$ complexes were obtained by separation of the water phase and subsequent evaporation of the residual ethanol and hexane at 70 °C and water at 110 °C.

Synthesis of anisotropic-shaped CuCr₂Se₄ nanocrystals

All the experiments were carried out in a fume hood using standard Schlenk techniques, and any excess elemental Se was disposed of using safe laboratory techniques. In a typical synthesis of triangular and hexagonal shaped $CuCr_2Se_4$ nanocrystals, 2 mmol Se powder was dissolved in 15 mL of TOA at 320 °C and maintained at this temperature for 1 h under N_2 before cooling down to room temperature. The mixture was subsequently heated to 160–165 °C under vacuum and then to 300 °C under N_2 . Separately, 0.25 mmol of $Cu(C_{18}H_{33}O_2)_2$ and 0.5 mmol of $Cr(C_{18}H_{33}O_2)_3$ complexes, along with 2.5 mL of TOA, were heated to 175 °C and magnetically stirred under flowing N_2 for 1 h. At 300 °C, 2.5 mL of hot metal–oleate

complexes and TOA mixture were rapidly injected into the Se solution under nitrogen atmosphere and the reaction mixture was maintained at this temperature for 30 min while being continually stirred, during which a brown-black solution was formed. After the mixture cooled down to room temperature, a mixture of hexane and ethanol (1:3) was added to the solution. The resulting black precipitate was collected *via* centrifugation (6500 rpm, 5 min) and washed three times before drying.

Characterization

XRD patterns were recorded on a Rigaku-Smart Lab Advance system equipped with Cu Ka radiation source operated as a rotating anode at 40 kV and 100 mA. Transmission electron microscopy (TEM) analysis was performed using a JEOL JEM-2010 UHR, 200 kV transmission electron microscope equipped with a CCD camera. TEM image non-linear processing was carried out using Gatan digital micrograph version 3.9. Scanning electron microscopy (SEM) analysis was performed using a Hitachi S-4800 FESEM equipped with EDX detector. X-ray photoelectron spectroscopy (XPS) analysis was carried out on a Thermo Scientific K-Alpha spectrometer. Magnetic measurements were performed using a Quantum Design Dynacool Physical Properties Measurement System (PPMS) equipped with Vibrating Sample Magnetometer (VSM). Continuous wave (CW) EPR spectra were recorded in the temperature range of 260-370 K on a Bruker ELEXYS E540 X-band spectrometer equipped with an ER 4102 ST resonator (Bruker-Biospin, Billerica, MA) and an ER 4111VT variable temperature unit (Bruker). CW simulations data processing was performed using MATLAB EasySpin program.

Results and discussion

For the synthesis of anisotropic-shaped CuCr₂Se₄ nanostructures, we used TOA as a reducing agent for the reduction of selenium. Based on nuclear magnetic resonance spectra and mass spectra (NMR and mass data, ESI), we propose a mechanism for the reduction process. The reaction primarily leads to the formation of TOASe and some tri-octylamine-*N*-selenide. Besides, we determined that reaction temperatures below 300 °C predominantly favors the formation of the nonmagnetic phase CuCrSe₂, as confirmed by X-ray diffraction (Fig. S1, ESI†). Based on X-ray diffraction studies, it is apparent that the reaction initially leads to the formation of CuCrSe₂ and subsequently to CuCr₂Se₄ (Fig. S2, ESI†). Referring to the reaction mechanism suggested by Steckel *et al.*,²⁷ we propose a possible mechanism for the formation of CuCr₂Se₄ in Scheme 1 shown below.

$$\begin{split} (R_aCOO)_2\text{-}Cu + 2(R_aCOO)_3\text{-}Cr + 4R_bSeH \rightarrow \\ R_bSe\text{-}Cu\text{-}Cr\text{-}(OOCR_a)(SeR_b) + Cr\text{-}(OOCR_a)(SeR_b)_2 + 6R_aCOOH \rightarrow \\ CuCr_2Se_4 + 4R_aCOOH + 4R_bOOCR_a \end{split}$$

 $R_a = CH_3(CH_2)_7CH = CH(CH_2)_7CH_3; R_bH = TOA$

Scheme 1 Proposed reaction mechanism for the formation of $CuCr_2Se_4$ nanocrystals.

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2 0.26 nm < 400>
5 nm

Cu(C₁₈H₃₃O₂)₂ & Cr(C₁₈H₃₃O₂)₃ in TOA

Thermal Decomposition

TOASe

Aging

Assembly

Fig. 1 (a) TEM and (b and c) HRTEM images of the $CuCr_2Se_4$ nanocrystals. (d) Schematic illustration of the possible formation pathway of the $CuCr_2Se_4$ nanocrystals.

The morphology and structure of the CuCr₂Se₄ nanocrystals have been investigated by transmission electron microscopy (TEM). As shown in Fig. 1, the synthesized CuCr₂Se₄ nanocrystals exhibit a mixture of close to triangular and hexagonal morphology, with an average size of 20 nm. Fig. 1b and c show high-resolution TEM (HRTEM) images of the two distinct shape nanocrystals that exhibit lattice fringes. Lattice spacing's of d =0.27 nm for triangular morphology and 0.26 nm for hexagonal morphology are observed, corresponding to the (400) lattice planes of the face-central cubic chalcospinel phase. Based on thorough analysis of our results and suggested formation pathways reported in the literature,28 we propose a likely formation pathway for the triangular and hexagonal CuCr2Se4 nanocrystals, as illustrated in Fig. 1d. In the reaction, metaloleate complexes have lower decomposition temperatures (~300 °C) than previously utilized metal-acetylacetonate precursors (\sim 350 °C). Our earlier results have shown that the metal-acetylacetonate precursors favor the formation of cubic morphology of CuCr₂Se₄ nanocrystals.¹⁸ Due to the slower decomposition rate at the lower temperature than using metalacetylacetonate precursors reacted in a solution of octadecene and oleylamine, the nucleation of CuCr₂Se₄ nanocrystals using the oleate precursors will be slower resulting in a lower concentration of monomers. Besides, the role of solvent as stabilizer and surfactant in the colloidal synthesis of

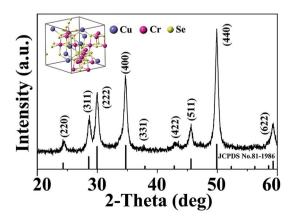


Fig. 2 XRD pattern of synthesized $CuCr_2Se_4$ nanocrystals. Inset shows the crystal structure of $CuCr_2Se_4$ and the standard XRD stick pattern of bulk $CuCr_2Se_4$ (JCPDS no. 81-1986) is also shown.

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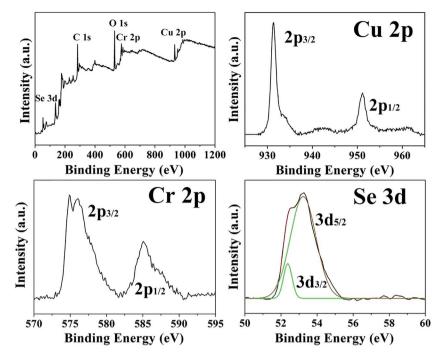


Fig. 3 XPS survey spectrum of the CuCr₂Se₄ nanocrystals, and high resolution XPS spectra of Cu 2p, Cr 2p, and Se 3d of the CuCr₂Se₄ nanocrystals.

nanocrystals will also have a significant influence. These factors kinetically favor formation of the thermodynamically unstable triangular and hexagonal morphology. However, limited by the relatively higher aging temperature and the imperfection of the assembly, some nanocrystals deviating from triangular and hexagonal shapes are also formed.

X-ray diffraction (XRD) measurements confirm the phase purity of the as-synthesized $CuCr_2Se_4$ nanocrystals, as shown in Fig. 2, and the diffraction pattern matches very well with the face-centered cubic spinel structure with $Fd\bar{3}m$ space group (shown in Fig. 2 inset). The major diffraction peaks are indexed as (220), (311), (222), (400), (331), (422), (511), (440), and (622) planes of cubic $CuCr_2Se_4$ phase (JCPDS no. 81-1986). Based on the peak positions, the lattice parameter is calculated to be 10.17 Å for the $CuCr_2Se_4$ nanocrystals, which is close to the value reported for bulk $CuCr_2Se_4$ (10.337 Å). Additionally, we have used energy dispersive X-ray analysis (EDX) to determine the elemental ratio of the $CuCr_2Se_4$ nanocrystals, which shows that the Cu: Cr: Se concentration ratio is close to the expected value of 1: 2: 4 (Fig. S3, ESI†).

X-ray photoelectron spectroscopy (XPS) studies have been carried out to investigate the chemical states of the ions in the CuCr_2Se_4 nanocrystals. A survey spectrum of the synthesized nanocrystals helps identify the presence of Cu, Cr, Se, O and C, and high-resolution spectra of Cu 2p, Cr 2p, and Se 3d are measured to determine the oxidation states of the constituent elements (Fig. 3). In the Cu 2p spectrum, peaks at 931.4 and 951.2 eV are ascribed to Cu 2p_{3/2} and Cu 2p_{1/2}, which can be assigned to Cu(II) with peak splitting of 19.8 eV.²⁹ The Cr 2p spectrum exhibits a spin–orbit separation of 9.7 eV between the Cr 2p_{3/2} and Cr 2p_{1/2} states, indicative of Cr(III).³⁰ The Cr 2p_{3/2} states are split into two peaks at 574.9 and 576.0 eV. The peak

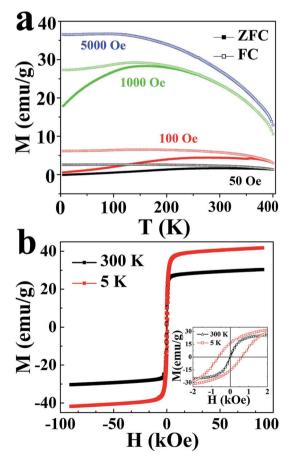


Fig. 4 (a) Zero-field-cooled (ZFC) and field-cooled (FC) magnetization (M) versus temperature (T) of the $CuCr_2Se_4$ nanocrystals measured at different magnetic fields. (b) Magnetization as a function of field measured at 5 and 300 K. The inset in (b) displays an enlarged portion of the hysteresis loops measured at 5 and 300 K.

separation with the binding energy difference ΔE of about 1 eV is typical of the 3d elements. Such a value of ΔE has been reported for a number of chalcogenide spinels.³¹ The peaks identified at 52.4 and 53.2 eV in the high-resolution spectrum of Se are generally attributed to Se $3d_{3/2}$ and Se $3d_{5/2}$, respectively.³² These results further confirm that the as-prepared nanocrystals are primarily CuCr₂Se₄.

The magnetic properties of the phase-pure mixture of triangular and hexagonal-shaped CuCr2Se4 nanocrystals have been measured using a Quantum Design Dynacool Physical Properties Measurement System (PPMS) equipped with a Vibrating Sample Magnetometer (VSM). A series of zero-fieldcooled (ZFC) and field-cooled (FC) scans are shown in Fig. 4a, in which the magnetization is plotted as a function of temperature for several different applied magnetic fields. The ZFC and FC curves measured at magnetic fields of 50, 100 and 1000 Oe are distinct at lower temperatures but merge together at higher temperatures, which is different from the curves measured at a higher magnetic field (5000 Oe). This is a typical characteristic of superparamagnetism resulting from the very small crystalline size of the nanocrystals.33 The Curie temperature is much higher than room temperature, with the blocking temperature obtained from the low-field (50 Oe) ZFC-FC measurement being \sim 390 K. This is higher than the blocking temperature reported for 25 nm cubic $CuCr_2Se_4$ nanocrystals ($T_b = 250$ K, 50 Oe). It is known that the blocking temperature $T_{\rm b}$ depends on interparticle interactions,34,35 which will vary with morphology. The $T_{\rm b}$ behavior can be modeled for single-domain particles by the Stoner-Wohlfarth equation $T_b = KV \sin^2 \theta / k_B \ln(\tau_m / \tau_0)$, where K is the nanoparticle's magnetic anisotropy energy density; θ is the angle between the net nanoparticle moment and the

magnetic easy axis; V is the nanoparticle volume; $k_{\rm B}$ is the Boltzmann constant; $T_{\rm b}$ is the blocking temperature; and $\ln(\tau_{\rm m}/\tau_{\rm 0})$ is invariable for dc magnetic measurements, when the measurement time is kept constant.³⁶ Thus, the larger total anisotropy energy will lead to a higher blocking temperature.

Fig. 4b shows plots of the magnetization as a function of applied magnetic field at 5 and 300 K, respectively. As indicated, the triangular and hexagonal CuCr2Se4 nanocrystals exhibit superparamagnetic behavior (no hysteresis and zero remanence) at room temperature but exhibit ferromagnetic behavior with a coercivity value of \sim 1 kOe at 5 K. The low temperature saturation magnetization value is \sim 43 emu g⁻¹ at a field of 90 kOe, which is similar to the value of 25 nm cubic CuCr₂Se₄ nanocrystals reported previously.18 The superparamagnetic nanocrystals at 300 K exhibit a saturation magnetization value of \sim 31 emu g⁻¹. Nanocrystals with narrow size dispersion show the peak and bifurcation point at the same temperature in FC-ZFC measurements, corresponding to the superparamagnetic blocking temperature. However, for our nanocrystals the peak and bifurcation point occur at different temperatures, suggesting a broad size distribution as confirmed from TEM images in Fig. 1a. In this case, we can refer the peak to an ensemble average blocking temperature and the bifurcation point as the largest blocking temperature. Thus, the synthesized triangular and hexagonal CuCr2Se4 nanocrystals exhibit superparamagnetic behavior at 300 K, which is below the largest blocking temperature of 390 K.

To further investigate the magnetic properties of anisotropic-shaped CuCr₂Se₄ nanocrystals, electron paramagnetic resonance (EPR) spectroscopy has been used to detect the coupling interactions. The EPR signal is produced by localized 3d electrons of

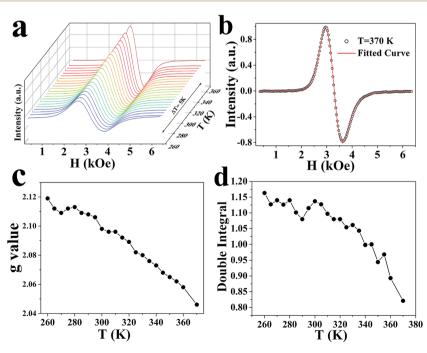


Fig. 5 (a) The differential EPR spectra from 260 to 370 K with steps of 5 K. (b) EPR spectrum at 370 K fitted by a convolution of Gaussian and Lorentzian line shapes. (c) The g values as a function of temperature. (d) Double integral of CW EPR spectra showing the spin susceptibility for the anisotropic-shaped $CuCr_2Se_4$ nanocrystals.

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Cr³⁺ ions. Fig. 5a shows the EPR spectra from 260 to 370 K with steps of 5 K, which is below the magnetic transition temperature (T_c) . The EPR first-derivative spectra become stronger and sharper at temperatures below T_c and can be fitted well by a convolution of Lorentzian and Gaussian line shapes (detailed fitting parameters shown in Fig. S4, ESI.†). Fig. 5b shows the spectrum data and fitted curves at 370 K. The extracted g values are close to 2, and increase as the temperature decreases, as shown in Fig. 5c. Fig. 5d shows the double integral of the CW EPR spectra, which gives the spin susceptibility for the anisotropic CuCr₂Se₄ nanocrystals. These behaviors suggest that the ferromagnetism of the sample adds an internal magnetic field to the applied field. Thus, a weaker applied field is required to achieve resonance at the lower temperatures and the apparent g values are larger at lower temperatures, similar to results reported for bulk CuCr₂Se₄.³⁷ This may be due to the random orientation distribution of the nanocrystals present during measurement. The magnetic anisotropy in this system is thus dominated by the intrinsic lattice magnetocrystalline anisotropy, with only a minor shape anisotropy contribution. A detailed investigation of the influence of surface spins may be necessary and is underway.

Conclusions

In summary, we have developed a simple and highly efficient synthesis route for preparation of single-phase anisotropicshaped CuCr₂Se₄ nanocrystals. The process involves the thermal decomposition and reaction of novel mixed metaloleate complexes with selenium in a high-boiling point organic solvent trioctylamine. We have investigated and proposed a possible reaction mechanism and formation pathway for the triangular and hexagonal CuCr₂Se₄ nanocrystals. Furthermore, we studied the magnetic properties of anisotropic-shaped CuCr₂Se₄ nanocrystals. We believe the process can be extended to the synthesis and studies of other anisotropicshaped magnetic chromium-based chalcogenide nanocrystals.

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Notes and references

- 1 R. F. Ziolo, E. P. Giannelis, B. A. Weinstein, M. P. O'Horo, B. N. Ganguly, V. Mehrotra, M. W. Russel and D. R. Huffman, Science, 1992, 257, 219.
- 2 S. H. Sun, C. B. Murray, D. Weller, L. Folks and A. Moser, Science, 2000, 287, 1989.

- 3 A. M. Whyte, Y. Shuku, G. S. Nichol, M. M. Matsushita, K. Awaga and N. Robertson, J. Mater. Chem., 2012, 22, 17967.
- 4 W. Y. Zhao, Z. Y. Liu, P. Wei, Q. J. Zhang, W. T. Zhu, X. L. Su, X. F. Tang, J. H. Yang, Y. Liu, J. Shi, Y. M. Chao, S. Q. Lin and Y. Z. Pei, Nat. Nanotechnol., 2017, 12, 55.
- 5 L. He, M. Janner, Q. P. Lu, M. S. Wang, H. Ma and Y. D. Yin, Adv. Mater., 2015, 27, 86.
- 6 M. Planells, L. X. Reynolds, U. Bansode, S. Chhatre, S. Ogale, N. Robertson and S. A. Haque, Phys. Chem. Chem. Phys., 2013,
- 7 F. Zhao, M. Rutherford, S. Y. Grisham and X. G. Peng, J. Am. Chem. Soc., 2009, 131, 5350.
- 8 F. Liu, J. H. Zhu, W. L. Yang, Y. H. Dong, Y. L. Hou, C. Z. Zhang, H. Yin and S. H. Sun, Angew. Chem., Int. Ed.,
- 9 J. Hemberger, P. Lunkenheimer, R. Fichtl, H.-A. Krug von Nidda, V. Tsurkan and A. Loidl, Nature, 2005, 434, 364.
- 10 M. Muroi, R. Street and P. G. McCormick, Phys. Rev. B: Condens. Matter Mater. Phys., 2001, 63, 052412.
- 11 M. Esters, A. Liebig, J. J. Ditto, M. Falmbigl, M. Albrecht and D. C. Johnson, J. Alloys Compd., 2016, 671, 220.
- 12 I. Efthimiopoulos, Z. T. Y. Liu, S. V. Khare, P. Sarin, V. Tsurkan, A. Loidl, D. Popov and Y. Wang, Phys. Rev. B, 2016, 93, 174103.
- 13 V. A. Fedorov, Y. A. Kesler and E. G. Zhukov, Inorg. Mater., 2003, 39, S68.
- 14 A. P. Ramirez, R. J. Cava and J. Krajewski, Nature, 1997, 386,
- 15 V. N. Antonov, V. P. Antropov, B. N. Harmon, A. N. Yaresko and A. Y. Perlov, Phys. Rev. B: Condens. Matter Mater. Phys., 1999, 59, 14552.
- 16 Y.-H. A. Wang, A. Gupta, M. Chshiev and W. H. Butler, Appl. Phys. Lett., 2008, 92, 062507.
- 17 Y.-H. A. Wang, A. Gupta, M. Chshiev and W. H. Butler, Appl. Phys. Lett., 2009, 94, 062515.
- 18 Y.-H. A. Wang, N. Z. Bao, L. M. Shen, P. Padhan and A. Gupta, J. Am. Chem. Soc., 2007, 129, 12408.
- 19 K. Ramasamy, D. Mazumdar, Z. Y. Zhou, Y.-H. A. Wang and A. Gupta, J. Am. Chem. Soc., 2011, 133, 20716.
- 20 Y.-H. A. Wang, X. Y. Zhang, N. Z. Bao, B. P. Lin and A. Gupta, J. Am. Chem. Soc., 2011, 133, 11072.
- 21 K. Ramasamy, H. Sims, W. H. Butler and A. Gupta, J. Am. Chem. Soc., 2014, 136, 1587.
- 22 N. Z. Bao, X. M. Qiu, Y.-H. A. Wang, Z. Y. Zhou, X. H. Lu, C. A. Grimes and A. Gupta, Chem. Commun., 2011, 47, 9441.
- 23 P. Chao, L. Gao, A. Chaturvedi, N. Z. Bao, K. Yanagisawa, L. M. Shen and A. Gupta, J. Mater. Chem. C, 2015, 3, 12077.
- 24 Y. J. Kang, J. B. Pyo, X. C. Ye, R. E. Diaz, T. R. Gordon, E. A. Stach and C. B. Murray, ACS Nano, 2012, 7, 645.
- 25 N. Z. Bao, L. M. Shen, W. An, P. Padhan, C. H. Turner and A. Gupta, Chem. Mater., 2009, 21, 3458.
- 26 J. Park, K. An, Y. Hwang, J.-G. Park, H.-J. Noh, J.-Y. Kim, J.-H. Park, N.-M. Hwang and T. Hyeon, Nat. Mater., 2004, 3, 891.
- 27 J. S. Steckel, B. K. H. Yen, D. C. Oertel and M. G. Bawendi, J. Am. Chem. Soc., 2006, 128, 13032.

Paper

- 28 Y. J. Xiong, J. M. McLellan, J. Y. Chen, Y. D. Yin, Z.-Y. Li and Y. N. Xia, *J. Am. Chem. Soc.*, 2005, **127**, 17118.
- 29 T. Ghodselahi, M. A. Vesaghi, A. Shafiekhani, A. Baghizadeh and M. Lameii, *Appl. Surf. Sci.*, 2008, **255**, 2730.
- 30 E. Malicka, A. Waskowska, J. Heimann, R. Sitko and D. Kaczorowski, *J. Alloys Compd.*, 2012, **513**, 353.
- 31 V. Tsurkan, S. Plogmann, M. Demeter, D. Hartmann and M. Neumann, *Eur. Phys. J. B*, 2000, **15**, 401.
- 32 B. Canava, J. Vigneron, A. Etcheberry, J. F. Guillemoles and D. Lincot, *Appl. Surf. Sci.*, 2002, **202**, 8.
- 33 C. B. Murray, S. H. Sun, H. Doyle and T. Betley, *MRS Bull.*, 2001, 26, 985.
- 34 H. L. Sun, H. T. Shi, F. Zhao, L. M. Qi and S. Gao, *Chem. Commun.*, 2005, **34**, 4339.
- 35 D. Farrell, Y. H. Cheng, R. W. McCallum, M. Sachan and S. A. Majetich, *J. Phys. Chem. B*, 2005, **109**, 13409.
- 36 E. C. Stoner and E. P. Wohlfarth, *Philos. Trans. R. Soc., A*, 1948, **240**, 599.
- 37 L. Zhang, W. Tong, J. Y. Fan, C. J. Zhang, R. W. Li and Y. H. Zhang, *Eur. Phys. J. B*, 2011, **83**, 325.