# RSC Advances

# PAPER

Cite this: RSC Adv., 2017, 7, 24188

Received 12th February 2017 Accepted 25th April 2017

DOI: 10.1039/c7ra01748k

rsc.li/rsc-advances

## Introduction

van der Waals epitaxy of large-area continuous ReS<sub>2</sub> films on mica substrate†

Jing-Kai Qin,<sup>ac</sup> Wen-Zhu Shao,<sup>a</sup> Yang Li,<sup>ab</sup> Cheng-Yan Xu, D<sup>\*abc</sup> Dan-Dan Ren,<sup>ac</sup> Xiao-Guo Song<sup>b</sup> and Liang Zhen<sup>\*ac</sup>

Rhenium disulfide (ReS<sub>2</sub>) has attracted scientists' attention for its unique physical properties and potential applications in high-efficiency photodetector devices. Although lots of works have been done to obtain high-quality ReS<sub>2</sub> nanoflakes, in-plane uniform growth is still challenging due to its unique decoupling property between layers. In this work, we successfully realized the epitaxial growth of continuous monolayer ReS<sub>2</sub> films on mica substrate by chemical vapour deposition (CVD). By prolonging the growth time, continuous multilayer ReS<sub>2</sub> films can also be obtained. The growth mechanism of ReS<sub>2</sub> films is proposed based on Stranski-Krastanov theory. Filed effect transistors (FETs) based on multilayer ReS<sub>2</sub> films exhibit typical n-type semiconducting behaviour with a carrier density of 0.27 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> and ON/ OFF ratio of about 4  $\times$  10<sup>3</sup>. The photoresponsivity of the phototransistor could reach up to 0.98 A W<sup>-1</sup> with a light intensity of 0.56 mW cm<sup>-2</sup>, suggesting that ReS<sub>2</sub> is a promising material for electronic and optoelectronic applications. PAPER<br>
(a) Check for updates<br>  $\text{Res}(x_0, y_0)$  2.2438<br>
Since the risk and 2017 (Namely 10 and 2018)<br>
Since the risk and 2012 (Namely 2022)<br>
Since Club Song and Lang 22hen<sup>23</sup><br>
Nao-Guo Song<sup>o</sup> and Lang 22hen<sup>23</sup><br>
Nao-Guo So

Transition metal dichalcogenides (TMDs), two-dimensional materials with tunable bandgaps, have attracted extensive attention in recent years. Their suitable bandgaps ranging from 1–3 eV make them highly promising to fabricate nanoscale electronic and optoelectronic devices.<sup>1</sup>–<sup>7</sup> Most TMDs undergo a crossover of their bandgap from direct to indirect as the number of layers increases, which means that monolayers with a direct bandgap could absorb and emit light more efficiently. This transition can be attributed to a strong interlayer coupling and confinement effect between layers in the stacking direction. $8-10$  Recently, rhenium disulfide (ReS<sub>2</sub>) has attracted scientists' attention due its distinct properties, such as interlayer decoupling effect and anisotropic electronic properties.<sup>11-13</sup> Based on previous reports, the bandgap of  $\text{ReS}_2$  remains direct with layer increase, which means that the interlayer coupling in  $Res<sub>2</sub>$  is much weaker than other two-dimensional materials.<sup>14</sup> This property makes  $\text{ReS}_2$  an ideal candidate to fabricate novel devices irrespective to the thickness. Additionally, the distorted octahedral (1T) phase of  $Res<sub>2</sub>$  could cause a serious symmetry

splitting, leading to an anisotropic optical and electrical properties along in-plane directions.<sup>15,16</sup>

Mechanical exfoliation is a typical method to obtain highly quality nanoflakes.<sup>17-19</sup> However, the obtained mono- or fewlayer  $Res<sub>2</sub>$  are always very small, which make them difficult to fabricate nanoscale devices. Chemical vapor deposition (CVD) is an efficient way to achieve large area monolayer TMDs films.<sup>20–22</sup> K. Keyshar et al.<sup>23</sup> first obtained monolayer ReS<sub>2</sub> flakes using CVD method under relative low temperature (400-500 °C) with NH<sub>4</sub>ReO<sub>4</sub> as Re resource, however, the synthesized monolayers have high level of vacancies due to insufficient reaction under relatively low temperature. He et  $al.^{24}$ achieved highly crystallized pyramid-like few-layer  $\text{ReS}_2$  flakes with Re powder as source, but the growth rate was hindered due to the high melting point of Re metal. The strong interlayer decoupling effect of  $Res_2$  makes its out-of-plane growth predominant, which is very different from other TMDs such as  $MoS<sub>2</sub>$  and WSe<sub>2</sub>. Certain substrates with low surface energy could facilitate the atoms migration along in-plane directions, and would be helpful to obtain uniform films with very flat surface.<sup>25-27</sup> For example, F. Cui et al.<sup>25</sup> realized epitaxial growth of  $\text{ReS}_2$  films on mica substrate. Te powder was used as catalyst, which could form the low melting-temperature Te–Re binary eutectic with Re metal. By this way, the growth efficiency of  $\text{ReS}_2$  can be significantly improved. Besides, the atomically smooth and chemical inert surface of mica could also facilitate the atoms migration on surface, which makes the in-plane growth of  $\text{ReS}_2$  possible. Although some work has been done, the growth of large-area continuous  $\text{ReS}_2$  films has not even been reported.



<sup>&</sup>lt;sup>a</sup>School of Materials Science and Engineering, Harbin Institute of Technology, Harbin 150001, China. E-mail: cy\_xu@hit.edu.cn; lzhen@hit.edu.cn

b Shandong Provincial Key Lab of Special Welding Technology, Harbin Institute of Technology at Weihai, Weihai 264209, China

c MOE Key Laboratory of Micro-Systems and Micro-Structures Manufacturing, Harbin Institute of Technology, Harbin 150080, China

<sup>†</sup> Electronic supplementary information (ESI) available. See DOI: 10.1039/c7ra01748k

In this study, we successfully prepared large-area continuous  $Res<sub>2</sub>$  films by van der Waals epitaxy on mica substrate, and we also explicitly analyzed the growth mechanism based on Stranski-Krastanov theory. Large area monolayer  $\text{ReS}_2$  films with thickness of 0.8 nm can be prepared. Continuous multilayer  $\text{ReS}_2$  films were also obtained by prolonging the growth time. Furthermore, we fabricated field effect transistor (FETs) based on multilayer  $Res_2$  films and evaluated its electrical and optoelectronic properties. Our work provides a new approach to obtain large-area continuous  $\text{ReS}_2$  films and it would be helpful to pave the way for widespread applications of  $\text{ReS}_2$ .

## **Experimental**

### Materials preparation

 $Res<sub>2</sub>$  films were obtained using chemical vapor deposition (CVD). Fresh fluorophlogopite mica  $[KMg_3(AlSi_3O_{10})F_2]$  was used as epitaxy substrate, since it could provide an atomically smooth and chemically inert surface for atoms in-plane diffusion.<sup>26,27</sup> In this work, we use Re metal power and  $Re<sub>2</sub>O<sub>7</sub>$  as precursors.  $Re<sub>2</sub>O<sub>7</sub>$  volatilizes rapidly under temperature above 300 °C,<sup>28</sup> and the partially sulfurized  $Re<sub>2</sub>O<sub>7</sub>$  could absorb on the substrate surface, which provides abundant nucleation sites for following  $Res_2$  growth. As the temperature goes up,  $Res_2$  layers starts to form and expand with incoming Re and S atoms. With this method, continuous  $\text{ReS}_2$  films are obtained under reaction temperature of 600 $\degree$ C through ambient pressure chemical vapor deposition process. Paper<br> **Open Karticle Commonsumers** Article is linear and the main term in the state of the common and the main term in the state of the main term in the state of the main term in the state of the main term in the state of

The synthesis of  $Res_2$  films was conducted in a two-heating zone tube furnace as schematically illustrated in Fig. 1.  $Al_2O_3$ ceramic boat containing  $Re<sub>2</sub>O<sub>7</sub>$  and Re mixture powder is located at the hot centre of zone 1, and fresh fluorphlogopite mica substrate is placed directly on the top of the boat. The



Fig. 1 (a) Schematic of the experiment set-up for CVD growth of monolayer ReS<sub>2</sub> films. (b) Heating curves of the two zones. (c) Schematic for the epitaxial growth process of  $\text{ReS}_2$  on mica. Partially reduced  $Re_2O_{7-x}$  could absorb on the mica surface and act as nucleation sites.

heating zone is heated to 600  $^{\circ}$ C at a ramping rate of 15  $^{\circ}$ C  $min^{-1}$  and kept for 10 min. 0.4 g S powder is placed in another ceramic boat at the upper stream side, where the temperature is maintained at 220  $^{\circ}$ C during the reaction. After growth, the chamber is naturally cooled down to room temperature. Throughout the growth process, argon (at 40 sccm) is used as a carrier gas into the reaction chamber.

#### Materials characterization

The morphology of  $\text{ReS}_2$  films was checked by optical microscope (Zeiss Imager A2m) and atomic force microscope (Bruker Dimension ICON-PT). Raman spectrum was recorded using a LabRAM Raman microscopy system with a 532 nm laser excitation. Si peak at 520  $cm^{-1}$  was calibrated as a reference for wave number calibration. The point size of laser is  $1 \mu m$  with a power of 0.5 W to avoid possible thermal effect introduced damage. The elemental compositions and chemical states of ReS<sub>2</sub> were investigated by XPS (Phi Quantera). PMMA wet transfer method was taken to transfer  $\text{ReS}_2$  thin films onto a copper grid for TEM characterization.

#### Device fabrication and electric measurement

Field effect transistors based on  $\text{ReS}_2$  films were fabricated by evaporating Cr/Au (10/100 nm) electrodes on the top of  $\text{ReS}_2$ layer using shadow mask method. Before test, the devices were annealed at 300 $\degree$ C for 4 h in Ar gas to improve the electrical contact. The electrical and optoelectronic properties were measured using semiconductor analyzer (Keithley 4200) on a Lakeshore probe station. A 500 W xenon lamp was used as the light source, and monochromatic lights of 254–850 nm were obtained using optical filters. The intensities of the incident light source were corrected using a power energy meter (Model 372, Scienteck).

## Result and discussion

Fig. 2a shows the digital photograph of monolayer  $\text{ReS}_2$  films growth on mica.  $\text{ReS}_2$  membrane on mica substrate exhibit grey contrast. The optical microscopy of as-grown  $\text{ReS}_2$  films is given in Fig. 2b, where the monolayer  $\text{Res}_2$  films exhibit a slight colour contrast with mica substrate. After transferred onto Si/  $SiO<sub>2</sub>$  substrate, ReS<sub>2</sub> films could still maintain continuous and intact morphology (Fig. S1, ESI†). Atomic force microscopy (AFM) was employed to determine the surface morphology and thickness. As shown in Fig. 2c, the thickness of  $\text{ReS}_2$  films is about 0.8 nm, corresponding to monolayer structure. With growth time increased to 30 min, the films became thicker with average thickness up to 4.2 nm (Fig. S2, ESI†). It is noteworthy that the surface roughness of monolayer films is very small  $(R<sub>q</sub>$  $= 0.10$  nm), while the roughness of multilayer films is much large  $(R<sub>q</sub> = 0.54$  nm).

In the Raman spectra of monolayer  $\text{ReS}_2$  films, two dominant peaks located at 149.1 and 210.4  $cm^{-1}$  are clearly observed (Fig. 2d), which are assigned to the in-plane  $(E_{\varphi})$  and out-ofplane  $(A_g)$  vibration modes, respectively.<sup>29</sup> Several other Raman modes ranging from 100-400  $\text{cm}^{-1}$  may arise from the



Fig. 2 (a) Digital photograph of as-grown ReS<sub>2</sub> films on mica substrate. Inset is digital picture of mica substrate with and without ReS<sub>2</sub> films. (b) Optical image and (c) AFM tomography image of ReS<sub>2</sub> monolayer films. Inset is the height profile. (d) Raman spectra of as-grown ReS<sub>2</sub> films on mica and SiO<sub>2</sub> substrate. (e) Comparison of Raman spectra between monolayer and multilayer ReS<sub>2</sub> films on mica substrate. (f) Raman mapping of monolayer ReS $_2$  films at E<sub>g</sub> peak 149 cm $^{-1}$ .

symmetry splitting introduced by the distorted 1T phase of ReS<sub>2</sub>.<sup>15</sup> Compared to as-synthesized monolayer films on mica,  $E_g$  and  $A_g$  mode of transferred samples present a slight blueshift by 2.86 and 2.19  $\rm cm^{-1}.$  This could be attributed to the release of tensile strain which arises from the lattice mismatch between  $Res<sub>2</sub>$  and mica. Raman mode shift is considered as an effective approach to evaluate the strain in TMDs.<sup>31,32</sup> For example, the  $E_{2g}$  and  $A_{1g}$  peaks of monolayer MoS<sub>2</sub> grown on mica both are blueshifted about 3–4  $cm^{-1}$  due to compressive strain.<sup>30</sup> In this work, red shift of Raman modes in  $\text{ReS}_2$  films indicates the existence of tensile strain. With the increase of film thickness, tensile strain would accumulate and it has to be released by splitting  $\text{ReS}_2$  films, resulting in lots of cover-like projections as shown in Fig. S3.† Consistent with previous studies,<sup>12,15</sup> Raman peaks of multilayer films do not show any difference compared to monolayer samples due to the weak interlayer coupling effect (Fig. 2e). Raman mapping of monolayer  $Res_2$  films was also performed at  $E<sub>g</sub>$  mode position, and a homogeneous feature in large area up to 100  $\mu$ m<sup>2</sup> scale is presented.

Compared with previous work,<sup>23-25</sup> where only isolated  $\text{ReS}_2$ domains appeared on substrate, the products obtained in our experiment are continuous films.  $Re<sub>2</sub>O<sub>7</sub>$  is used as catalyst, which has the lowest melting point (220  $^{\circ}$ C) among all the rhenium oxide. During the growth process,  $Re<sub>2</sub>O<sub>7</sub>$  would rapidly volatilize into vapour phase, and then be reduced by sulfur vapour into  $\text{Re}_2\text{O}_{7-x}$  species.<sup>23</sup> Subsequently, they are conveyed downstream by Ar gas and absorbed on mica surface, acting as active nucleation sites for further  $Res<sub>2</sub>$  growth.

The morphology of  $\text{ReS}_2$  is very sensitive to the growth time, as shown in Fig. 3. At the beginning, intermediate  $Re<sub>2</sub>O<sub>7-x</sub>$ species absorbed on mica surface could act as nucleation sites. As temperature go up,  $Re<sub>2</sub>O<sub>7-x</sub>$  species are further sulfurized into  $\text{ReS}_2$  with incoming sulfur vapour. Due to low atoms migration energy of mica surface,  $\text{ReS}_2$  could easily grow along the in-plane direction into monolayer domains. These isolated monolayer  $\text{ReS}_2$  domains have a small size ranging from 500 nm to 2  $\mu$ m (Fig. 3a). It is worthy noted that these isolated domains exhibit irregular dendrite shapes, which is very different from traditional TMDs, such as  $MOS_2$  and  $WS_2$ .<sup>22,30</sup> This kind of anisotropic growth is attributed to the anisotropic interfacial energy introduced by distorted 1T crystal structure in  $\text{ReS}_2$ . Volatile  $\text{Re}_2\text{O}_7$  would easily be run out before reaching 600 °C, and it cannot act as a stable reaction source for growing thick ReS<sub>2</sub> films. The main function of  $Re<sub>2</sub>O<sub>7</sub>$  is supplying enough nucleation sites for further  $Res_2$  films growth. When the temperature rises up to 600 $\degree$ C, Re powder start to volatilize and react directly with sulfur into  $\text{ReS}_2$ , which plays the decisive part in multilayer continuous films growth. Isolated  $\text{ReS}_2$  domains start to interconnect with each other into continuous films as reaction undergoing (Fig. 3b). Typically, when two separate  $\text{ReS}_2$ domains gather, the edges tend to interconnect rather than overlap, which is the typical characteristic of 2D materials growth.<sup>30</sup> After the growth of continuous monolayer films is completed, following  $Res_2$  layers start to emerge on the surface of the first layer (Fig. 3c). We believe that these  $Res<sub>2</sub>$  layers should be originated from the zero dimensional or quasi-onedimensional structures as indicated with white circles in Fig. 3c, suggesting that the  $Res_2$  growth obeys Stranski-Krastanov growth mode.<sup>33</sup> In other word, except the first layer, the next layers follow the island-based growth mode, which would result in multilayer films with relatively rough surface. By prolonging the growth time to 30 min, we could obtain multilayer films with thickness about 4.2 nm (Fig. 3d). The surface roughness difference between monolayer and multilayer films further confirms our suggestion  $(R_q \text{ are } 0.10 \text{ and } 0.54 \text{ nm},$ respectively). XRD pattern of multilayer films shows a strong



Fig. 3 (a-d) Schematic (upper) and corresponding AFM topography images (lower row) showing the possible growth mechanism of ReS<sub>2</sub>: (a) 5 min. (b) 10 min. (c) 15 min. (d) 30 min. Scale bar: 1 μm. (e, f) Schematic growth model for different stages of ReS<sub>2</sub> films on mica substrate.

(001) orientation due to the layer-by-layer growth mode on mica substrate (Fig. S4, ESI†).

The substrate plays a crucial role in the growth process. As shown in Fig. 3e, atomically smooth and inert chemically surface of mica could facilitate the atoms migration along the surface due to the low energy barrier,<sup>34,35</sup> which make the first layer growth on mica surface-dominated and result in continuous monolayer structure. Compared with the very flat epitaxial growth of the first layer on mica, following layers always have relatively rough surface. Because the surface of  $\text{ReS}_2$  could not provide such low energy for atoms migration as mica, thus the in-plane growth is hindered (Fig. 3f). Compared with other TMDs, such as  $MOS<sub>2</sub>$  and WS<sub>2</sub>, and weak van der Waals interaction between layers of the  $Res<sub>2</sub>$  could also facilitate the 3D structure growth,<sup>36</sup> and accelerate the atoms migration along vertical direction, leading to the rough surface. As shown in Fig. S5,† the products grown on  $SiO<sub>2</sub>$  substrate exhibit flowerlike structures. This is because the rough surface of  $SiO<sub>2</sub>$  with dangling bonds could increase the energy barrier for atoms migration and prohibit  $\text{ReS}_2$  from growing along in-plane directions.

Re could directly react with S vapour to form  $Res<sub>2</sub>$  under temperature range of 600–900  $^{\circ}$ C.<sup>37</sup> However, the high melting temperature of Re (about 3150  $^{\circ}$ C) leads to a very low nucleation rate during the growth process. In our experiments, no products appear without  $\text{Re}_2\text{O}_7$  under the identical growth conditions. We believe the severe volatilization of  $\text{Re}_2\text{O}_7$ , which provide adequate nucleation sites for the later  $\text{ReS}_2$  growth, plays crucial part in monolayer film forming. The dosage of  $Re<sub>2</sub>O<sub>7</sub>$  is also important for the growth process. In this experiment,  $Re<sub>2</sub>O<sub>7</sub>$  is only one-tenth of the Re power. Only  $Re<sub>2</sub>O<sub>7</sub>$  as precursor would result in numerous nanoparticles without  $Res<sub>2</sub>$  on mica surface. XPS spectrum shows these particles are mainly composed of  $Re<sub>2</sub>O<sub>7-x</sub>$  intermediate species (Fig. S6, ESI†). Meanwhile, more  $Re<sub>2</sub>O<sub>7</sub>$  could lead to thick film consist with lots of tiny grains with large surface roughness (Fig. S7, ESI†).

X-ray photoelectron spectroscopy (XPS) was performed to determine the stoichiometry of  $\text{ReS}_2$  films. Fig. 4a shows the survey spectrum of the monolayer  $\text{Res}_2$  films. The core  $2p_{3/2}$  and  $2p_{1/2}$  level peaks for sulfur are located at 161.37 and 162.63 eV (Fig. 4b). The two dominant level peaks for rhenium are located at 42.6 and 45.1 eV (assigned to Re  $4f_{7/2}$  and Re  $4f_{5/2}$  respectively), consistent with previous study<sup>36</sup> (Fig. 4c). The prominent  $4f_{7/2}$  and  $4f_{5/2}$  peak for  $Re_2O_7$ ,<sup>38</sup> located at 40.1 and 42.3 eV, are not observed, suggesting the  $\text{Re}_2\text{O}_7$  has been completely sulfurized.



Fig. 4  $\times$ PS spectra of monolayer ReS<sub>2</sub> films. (a) XPS survey spectrum of the sample. Three kinds of elements are present: Re and S (from the sample), Si (from the substrate). (b) XPS spectrum of S 2p and (c) XPS spectrum of Re 4f.

Transmission electron microscope (TEM) analysis was also conducted to explore the crystal structure of  $\text{ReS}_2$  films. Using a traditional poly(methyl methacrylate) (PMMA) wet transfer method,  $\text{ReS}_2$  films were transferred onto carbon-supported copper grids. From the low-magnification TEM image, a continuous membrane with obvious folding could be clearly observed (Fig. 5a). Edge folding often appears in TEM analysis of two-dimensional materials, like  $MoS<sub>2</sub>$  and graphene,<sup>30,39</sup> which is probably introduced by the transfer process. HRTEM reveals perfect single-crystal structure of the monolayer films, and clear DS-chains of Re atoms could be seen (Fig. 5b). The angle between  $a[100]$  and  $b[010]$  is about 60°, which is agree well with the date published on exfoliated  $\text{ReS}_2$  samples.<sup>15</sup> Typical (100) and (010) crystal plane spacing are 0.31 and 0.35 nm, respectively, confirming the high crystallinity of monolayer films. As for the multilayer films, it is demonstrated to be polycrystal consist of many tiny grains with random orientations (Fig. 5c). Fig. 5d shows the HRTEM and corresponding SAED pattern. AED pattern shows clear dim diffraction rings, which further prove the films are polycrystal. **BSC Advances**<br>
Transmission electron increase/pe (TEM) amalysis was also and the direction of profit commonly commonly commonly commonly commonly commonly commonly commonly commonly compare and the commonly compare and t

Field effect transistors (FETs) are fabricated based on multilayer  $\text{ReS}_2$  films (Fig. 6a). It should be noted that previous reports only focus on single crystal  $\text{Res}_{2}$ ,<sup>11,15,24</sup> whereas the performance of multilayer polycrystalline films are still unexplored. The gate-dependent transport curves show a typical ntype behaviour with maximum ON/OFF ratio about  $4 \times 10^3$ , and the threshold voltage is about  $-27$  V with  $V_{ds}$  of 3 V. The carrier mobility could be calculated using following equation:<sup>4</sup>

$$
\mu = \frac{L}{W V_{\rm ds} C_{\rm g}} \times \frac{dI_{\rm ds}}{dV_{\rm bg}}
$$

where L, W,  $C_{\rm g}$ ,  $V_{\rm ds}$ ,  $V_{\rm bg}$  and  $I_{\rm ds}$  are the channel length, width, the gate capacitance per unit area, source-drain voltage, gate voltage and source-drain current, respectively. The carrier mobility of the device can be up to 0.27  $\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ . Compared with single-crystal monolayer  $Res_2$  FET with carrier mobility 1.43  $\text{cm}^2\,\text{V}^{-1}\,\text{s}^{-1}$  and ON/OFF ratio 10<sup>6</sup> (Fig. S8, ESI†), FET based on multilayer  $\text{ReS}_2$  films exhibits degraded electrical performance, which is attributed the carrier scattering effect in grain boundary.<sup>41</sup>

The optoelectronic properties of multilayer  $\text{ReS}_2$  films were also evaluated based on the same FET device. Fig. 7a shows the schematic image of the device, a 500 W xenon lamp was used as the light source with illumination power ranging from 0.25– 0.56 mW  $\text{cm}^{-2}$ . Fig. 7b shows the *I-V* curves under different incident light intensities. Obviously, the channel current increases significantly under illumination compared to dark condition. Fig. 7c and d shows the photocurrent and responsitivity as a function of light power. Here, photocurrent  $(\Delta I)$  is the difference between  $I_{on}$  and  $I_{off}$  with a  $V_{ds}$  of 1 V, and the responsitivity is defined as  $R = \Delta I/P \times S$  (P is the light power intensity, and  $S$  is the illuminated area of the devices). Both  $\Delta I$ 



Fig. 5 (a and b) Bright-field TEM and HRTEM images of monolayer ReS<sub>2</sub> films. (c and d) Bright-field TEM and HRTEM (STEM) images of few-layer ReS<sub>2</sub> films. Insets in (b and d) show corresponding SAED patterns.



Fig. 6 Electrical properties of multilayer ReS<sub>2</sub> films. (a) Optical image shows the as-fabricated FET device. (b, c) Output and transfer curve of this device. Back-gate voltage ( $V<sub>a</sub>$ ) sweeping: -30 to 30 V; source-drain bias ( $V<sub>ds</sub>$ ): 1 V to 3 V.



Fig. 7 Optoelectronic performance of the phototransistor based on multilayer  $\text{Re}S_2$  films on SiO<sub>2</sub> substrate. (a) Schematic image of the device (b) I–V curves of device in dark and under 490 nm illumination with different intensities. (c, d) Photocurrent and photoresponsitivity as a function of incident power. (e, f) Time-dependent photoresponse of the device under 490 nm illumination with intensity of 0.56 mW cm<sup>-2</sup> and bias of 1 V

and R raise slightly with the increase of intensity. The responsivity is estimated to be 0.98 A  $W^{-1}$  at an irradiation intensity of 0.56 mW  $\text{cm}^{-2}$ , almost one order smaller than that in singlecrystal monolayer  $\text{ReS}_2$  based devices (Fig. S9, ESI†). Timeresolved photoresponse behaviour was probed by switching the light on and off with a time interval of 100 s under a bias of 1 V. The device exhibits a stable and repeatable response to the illumination after 5 cycles (Fig. 7e). A considerable enhancement of current (19 nA) is measured when the device is illuminated by 490 nm light. The response and recovery time is calculated by considering the time it takes to reach 80% of the final values, which are 17.8 and 14.3 s, respectively. Both of the rising time and recovery time of  $\text{ReS}_2$  films devices are larger than that of exfoliated  $\text{ReS}_2$  nanosheet,<sup>13,40</sup> but comparable to the bilayer polycrystalline film.<sup>41</sup> Generally, the active sites of the material significantly affect the rise and decay times.<sup>27,42</sup> Based on our measurement condition, gas molecules in air such as oxygen can easily absorbed on the surface of films. Under light illumination, photogenerated carriers would first fill the trap states, which

prolongs the response time. We believe the improvement of measurement condition would be effective to improve the responsive speed.<sup>43</sup>

## **Conclusions**

In conclusion, we develop a facile method to obtain large-area continuous  $\text{ReS}_2$  films on mica substrate using CVD method.  $Re<sub>2</sub>O<sub>7</sub>$  acting as catalyst could provide enough seeds for continuous films growth. The multilayer films exhibits a layerby-layer growth mode with strong (001) orientation. FET measurement indicates that multilayer  $\text{ReS}_2$  films exhibit n-type transport behaviour with ON/OFF ratio about  $4 \times 10^3$  and carrier mobility of 0.27 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>. The optoelectronic properties of  $\text{ReS}_2$  films are also investigated, with the highest photoresponsivity of 0.98 A  $W^{-1}$  at an illumination intensity of  $0.56\,\mathrm{mW\,cm}^{-2}$  . Our work suggests that these scalable  $\mathrm{ReS}_2$  films may be an ideal candidates for future high performance electronic and optoelectronic devices.

## Acknowledgements

This work was financially supported by National Natural Science Foundation of China (No. 51572057) and Shandong Provincial Key Lab of Special Welding Technology, Harbin Institute of Technology at Weihai. The authors thanks Prof. Ping-An Hu for assistance on electrical measurements. **SCARTIGGEREE ARTIST**<br> **Access Artisle. Common Access Artisle. Published on 03 May 2017. This article. Access Artisle. Published and the common and** 

## Notes and references

- 1 N. R. Pradhan, D. Rhodes, S. Feng and Y. Xin, ACS Nano, 2014, 8, 5911–5920.
- 2 H. Fang, S. Chuang, T. C. Chang, K. Takei, T. Takahashi and A. Javey, Nano Lett., 2012, 12, 3788–3792.
- 3 Z. Yin, H. Li, H. Li, L. Jiang, Y. Shi and Y. Sun, ACS Nano, 2012, 6, 74–80.
- 4 B. Radisavljevic, A. Radenovic, J. Brivio, V. Giacometti and A. Kis, Nat. Nanotechnol., 2011, 6, 147–150.
- 5 Y. Huang, et al., ACS Nano, 2014, 8, 10743–10755.
- 6 Y. H. Huang, C. C. Peng, R. S. Chen, Y. S. Huang and C. H. Ho, Appl. Phys. Lett., 2014, 105, 093106.
- 7 O. Lopez-Sanchez, D. Lembke, M. Kayci, A. Radenovic and A. Kis, Nat. Nanotechnol., 2013, 8, 497–501.
- 8 I. G. Lezama, A. Arora, A. Ubaldini, C. Barreteau, E. Giannini, M. Potemski and A. F. Morpurgo, Nano Lett., 2015, 15, 2336–2342.
- 9 Y. Li, et al., Adv. Funct. Mater., 2016, 26, 293–302.
- 10 Y. Li, C.-Y. Xu and L. Zhen, Appl. Phys. Lett., 2013, 102, 143110.
- 11 Y.-C. Lin, H.-P. Komsa, C.-H. Yeh, T. Björkman, Z.-Y. Liang, C.-H. Ho, Y.-S. Huang, P.-W. Chiu, A. V. Krasheninnikov and K. Suenaga, ACS Nano, 2015, 9, 11249–11257.
- 12 S. Tongay, et al., Nat. Commun., 2014, 5, 3252.
- 13 F. Liu, et al., Adv. Funct. Mater., 2016, 26, 1169–1177.
- 14 Q. Zhang, et al., Adv. Mater., 2016, 28, 2616–2623.
- 15 D. A. Chenet, O. B. Aslan, P. Y. Huang, C. Fan, A. M. van der Zande, T. F. Heinz and J. C. Hone, Nano Lett., 2015, 15, 5667-5672.
- 16 O. B. Aslan and D. A. Chenet, ACS Photonics, 2016, 3, 96–101.
- 17 N. R. Pradhan, et al., Nano Lett., 2015, 15, 8377–8384.
- 18 Y. Li, C.-Y. Xu, P. Hu and L. Zhen, ACS Nano, 2013, 7, 7795–7804.
- 19 Y. Li, C.-Y. Xu, B.-Y. Zhang and L. Zhen, Appl. Phys. Lett., 2013, 103, 033122.
- 20 Y. Zhang, et al., ACS Nano, 2013, 7, 8963–8971.
- 21 A. L. Elías, et al., ACS Nano, 2013, 7, 5235-5242.
- 22 W. Zhang, M.-H. Chiu, C.-H. Chen and W. Chen, ACS Nano, 2014, 8, 8653–8661.
- 23 K. Keyshar, et al., Adv. Mater., 2015, 27, 4640–4648.
- 24 X. He, F. Liu, P. Hu, W. Fu, X. Wang, Q. Zeng, W. Zhao and Z. Liu, Small, 2015, 11, 5423–5429.
- 25 F. Cui, et al., Adv. Mater., 2016, 28, 5019–5024.
- 26 Y. Zhou, Y. Nie, Y. Liu, K. Yan, J. Hong, C. Jin, Y. Zhou, J. Yin, Z. Liu and H. Peng, ACS Nano, 2014, 8, 1485–1490.
- 27 Q. Wang, M. Safdar, K. Xu, M. Mirza, Z. Wang and J. He, ACS Nano, 2014, 8, 7497–7505.
- 28 I. R. Beattie, T. R. Gilson and P. J. Jones, Inorg. Chem., 1996, 35, 1301–1304.
- 29 X.-F. Qiao, J.-B. Wu, L. Zhou, J. Qiao and W. Shi, Nanoscale, 2016, 8, 8324–8332.
- 30 Q. Ji, et al., Nano Lett., 2013, 13, 3870–3877.
- 31 Z. H. Ni, T. Yu, Y. H. Lu, Y. Y. Wang, Y. P. Feng and Z. X. Shen, ACS Nano, 2008, 2, 2301–2305.
- 32 M. Huang, H. Yan, T. F. Heinz and J. Hone, Nano Lett., 2010, 10, 4074–4079.
- 33 E. Bauer and J. H. van der Merwe, Phys. Rev. B: Condens. Matter Mater. Phys., 1986, 33, 3657.
- 34 Q. Wang, K. Xu, Z. Wang, F. Wang, Y. Huang, M. Safdar, X. Zhan, F. Wang, Z. Cheng and J. He, Nano Lett., 2015, 15, 1183–1189.
- 35 X. Li, F. Cui, Q. Feng, G. Wang, X. Xu, J. Wu, N. Mao, X. Liang, Z. Zhang and J. Zhang, Nanoscale, 2016, 8, 18956–18962.
- 36 J. Gao, L. Li, J. Tan, H. Sun, B. Li, J. C. Idrobo, C. V. Singh, T.-M. Lu and N. Koratkar, Nano Lett., 2016, 16, 3780–3787.
- 37 C.-H. Ho, Opt. Express, 2005, 13, 8–19.
- 38 Y. Jia, L. Duan, D. Zhang, J. Qiao and G. Dong, J. Phys. Chem. C, 2013, 117, 13763–13769.
- 39 J. C. Meyer, A. K. Geim, M. I. Katsnelson, K. S. Novoselov, T. J. Booth and S. Roth, Nature, 2007, 446, 60–63.
- 40 E. Liu, et al., Adv. Funct. Mater., 2016, 26, 1938–1944.
- 41 M. Hafeez, L. Gan, H. Li, Y. Ma and T. Zhai, Adv. Funct. Mater., 2016, 26, 4551–4560.
- 42 Y. Jiang, W. J. Zhang, J. S. Jie, X. M. Meng, X. Fan and S. T. Lee, Adv. Funct. Mater., 2007, 17, 1795–1800.
- 43 J. Zhou, Y. Gu, Y. Hu, W. Mai, P.-H. Yeh, G. Bao, A. K. Sood, D. L. Polla and Z. L. Wang, Appl. Phys. Lett., 2009, 94, 191103.