



Cite this: *RSC Adv.*, 2017, 7, 24188

van der Waals epitaxy of large-area continuous ReS₂ films on mica substrate†

Jing-Kai Qin,^{ac} Wen-Zhu Shao,^a Yang Li,^{ab} Cheng-Yan Xu,^{ID *abc} Dan-Dan Ren,^{ac} Xiao-Guo Song^b and Liang Zhen^{*ac}

Rhenium disulfide (ReS₂) has attracted scientists' attention for its unique physical properties and potential applications in high-efficiency photodetector devices. Although lots of works have been done to obtain high-quality ReS₂ nanoflakes, in-plane uniform growth is still challenging due to its unique decoupling property between layers. In this work, we successfully realized the epitaxial growth of continuous monolayer ReS₂ films on mica substrate by chemical vapour deposition (CVD). By prolonging the growth time, continuous multilayer ReS₂ films can also be obtained. The growth mechanism of ReS₂ films is proposed based on Stranski–Krastanov theory. Field effect transistors (FETs) based on multilayer ReS₂ films exhibit typical n-type semiconducting behaviour with a carrier density of 0.27 cm² V⁻¹ s⁻¹ and ON/OFF ratio of about 4 × 10³. The photoresponsivity of the phototransistor could reach up to 0.98 A W⁻¹ with a light intensity of 0.56 mW cm⁻², suggesting that ReS₂ is a promising material for electronic and optoelectronic applications.

Received 12th February 2017

Accepted 25th April 2017

DOI: 10.1039/c7ra01748k

rsc.li/rsc-advances

Introduction

Transition metal dichalcogenides (TMDs), two-dimensional materials with tunable bandgaps, have attracted extensive attention in recent years. Their suitable bandgaps ranging from 1–3 eV make them highly promising to fabricate nanoscale electronic and optoelectronic devices.^{1–7} Most TMDs undergo a crossover of their bandgap from direct to indirect as the number of layers increases, which means that monolayers with a direct bandgap could absorb and emit light more efficiently. This transition can be attributed to a strong interlayer coupling and confinement effect between layers in the stacking direction.^{8–10} Recently, rhenium disulfide (ReS₂) has attracted scientists' attention due its distinct properties, such as interlayer decoupling effect and anisotropic electronic properties.^{11–13} Based on previous reports, the bandgap of ReS₂ remains direct with layer increase, which means that the interlayer coupling in ReS₂ is much weaker than other two-dimensional materials.¹⁴ This property makes ReS₂ an ideal candidate to fabricate novel devices irrespective to the thickness. Additionally, the distorted octahedral (1T) phase of ReS₂ could cause a serious symmetry

splitting, leading to an anisotropic optical and electrical properties along in-plane directions.^{15,16}

Mechanical exfoliation is a typical method to obtain highly quality nanoflakes.^{17–19} However, the obtained mono- or few-layer ReS₂ are always very small, which make them difficult to fabricate nanoscale devices. Chemical vapor deposition (CVD) is an efficient way to achieve large area monolayer TMDs films.^{20–22} K. Keyshar *et al.*²³ first obtained monolayer ReS₂ flakes using CVD method under relative low temperature (400–500 °C) with NH₄ReO₄ as Re resource, however, the synthesized monolayers have high level of vacancies due to insufficient reaction under relatively low temperature. He *et al.*²⁴ achieved highly crystallized pyramid-like few-layer ReS₂ flakes with Re powder as source, but the growth rate was hindered due to the high melting point of Re metal. The strong interlayer decoupling effect of ReS₂ makes its out-of-plane growth predominant, which is very different from other TMDs such as MoS₂ and WSe₂. Certain substrates with low surface energy could facilitate the atoms migration along in-plane directions, and would be helpful to obtain uniform films with very flat surface.^{25–27} For example, F. Cui *et al.*²⁵ realized epitaxial growth of ReS₂ films on mica substrate. Te powder was used as catalyst, which could form the low melting-temperature Te–Re binary eutectic with Re metal. By this way, the growth efficiency of ReS₂ can be significantly improved. Besides, the atomically smooth and chemical inert surface of mica could also facilitate the atoms migration on surface, which makes the in-plane growth of ReS₂ possible. Although some work has been done, the growth of large-area continuous ReS₂ films has not even been reported.

^aSchool of Materials Science and Engineering, Harbin Institute of Technology, Harbin 150001, China. E-mail: cy_xu@hit.edu.cn; lzhen@hit.edu.cn

^bShandong Provincial Key Lab of Special Welding Technology, Harbin Institute of Technology at Weihai, Weihai 264209, China

^cMOE Key Laboratory of Micro-Systems and Micro-Structures Manufacturing, Harbin Institute of Technology, Harbin 150080, China

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c7ra01748k



In this study, we successfully prepared large-area continuous ReS₂ films by van der Waals epitaxy on mica substrate, and we also explicitly analyzed the growth mechanism based on Stranski–Krastanov theory. Large area monolayer ReS₂ films with thickness of 0.8 nm can be prepared. Continuous multilayer ReS₂ films were also obtained by prolonging the growth time. Furthermore, we fabricated field effect transistor (FETs) based on multilayer ReS₂ films and evaluated its electrical and optoelectronic properties. Our work provides a new approach to obtain large-area continuous ReS₂ films and it would be helpful to pave the way for widespread applications of ReS₂.

Experimental

Materials preparation

ReS₂ films were obtained using chemical vapor deposition (CVD). Fresh fluorophlogopite mica [KMg₃(AlSi₃O₁₀)F₂] was used as epitaxy substrate, since it could provide an atomically smooth and chemically inert surface for atoms in-plane diffusion.^{26,27} In this work, we use Re metal powder and Re₂O₇ as precursors. Re₂O₇ volatilizes rapidly under temperature above 300 °C,²⁸ and the partially sulfurized Re₂O₇ could absorb on the substrate surface, which provides abundant nucleation sites for following ReS₂ growth. As the temperature goes up, ReS₂ layers starts to form and expand with incoming Re and S atoms. With this method, continuous ReS₂ films are obtained under reaction temperature of 600 °C through ambient pressure chemical vapor deposition process.

The synthesis of ReS₂ films was conducted in a two-heating zone tube furnace as schematically illustrated in Fig. 1. Al₂O₃ ceramic boat containing Re₂O₇ and Re mixture powder is located at the hot centre of zone 1, and fresh fluorophlogopite mica substrate is placed directly on the top of the boat. The

heating zone is heated to 600 °C at a ramping rate of 15 °C min⁻¹ and kept for 10 min. 0.4 g S powder is placed in another ceramic boat at the upper stream side, where the temperature is maintained at 220 °C during the reaction. After growth, the chamber is naturally cooled down to room temperature. Throughout the growth process, argon (at 40 sccm) is used as a carrier gas into the reaction chamber.

Materials characterization

The morphology of ReS₂ films was checked by optical microscope (Zeiss Imager A2m) and atomic force microscope (Bruker Dimension ICON-PT). Raman spectrum was recorded using a LabRAM Raman microscopy system with a 532 nm laser excitation. Si peak at 520 cm⁻¹ was calibrated as a reference for wave number calibration. The point size of laser is 1 μm with a power of 0.5 W to avoid possible thermal effect introduced damage. The elemental compositions and chemical states of ReS₂ were investigated by XPS (Phi Quantera). PMMA wet transfer method was taken to transfer ReS₂ thin films onto a copper grid for TEM characterization.

Device fabrication and electric measurement

Field effect transistors based on ReS₂ films were fabricated by evaporating Cr/Au (10/100 nm) electrodes on the top of ReS₂ layer using shadow mask method. Before test, the devices were annealed at 300 °C for 4 h in Ar gas to improve the electrical contact. The electrical and optoelectronic properties were measured using semiconductor analyzer (Keithley 4200) on a Lakeshore probe station. A 500 W xenon lamp was used as the light source, and monochromatic lights of 254–850 nm were obtained using optical filters. The intensities of the incident light source were corrected using a power energy meter (Model 372, Scientec).

Result and discussion

Fig. 2a shows the digital photograph of monolayer ReS₂ films growth on mica. ReS₂ membrane on mica substrate exhibit grey contrast. The optical microscopy of as-grown ReS₂ films is given in Fig. 2b, where the monolayer ReS₂ films exhibit a slight colour contrast with mica substrate. After transferred onto Si/SiO₂ substrate, ReS₂ films could still maintain continuous and intact morphology (Fig. S1, ESI†). Atomic force microscopy (AFM) was employed to determine the surface morphology and thickness. As shown in Fig. 2c, the thickness of ReS₂ films is about 0.8 nm, corresponding to monolayer structure. With growth time increased to 30 min, the films became thicker with average thickness up to 4.2 nm (Fig. S2, ESI†). It is noteworthy that the surface roughness of monolayer films is very small ($R_q = 0.10$ nm), while the roughness of multilayer films is much large ($R_q = 0.54$ nm).

In the Raman spectra of monolayer ReS₂ films, two dominant peaks located at 149.1 and 210.4 cm⁻¹ are clearly observed (Fig. 2d), which are assigned to the in-plane (E_g) and out-of-plane (A_g) vibration modes, respectively.²⁹ Several other Raman modes ranging from 100–400 cm⁻¹ may arise from the

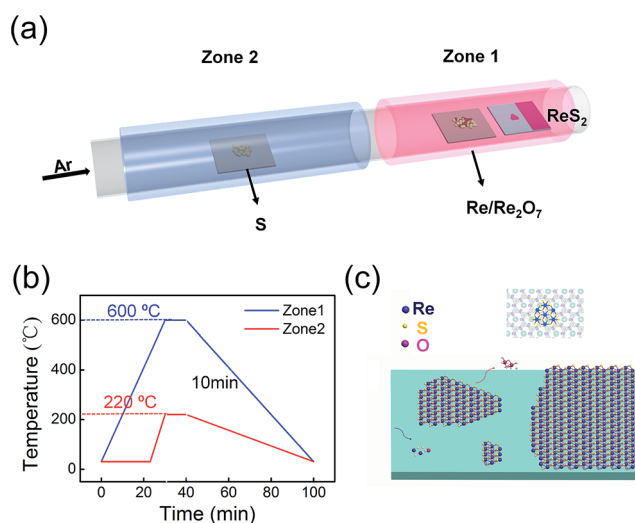


Fig. 1 (a) Schematic of the experiment set-up for CVD growth of monolayer ReS₂ films. (b) Heating curves of the two zones. (c) Schematic for the epitaxial growth process of ReS₂ on mica. Partially reduced Re₂O_{7-x} could absorb on the mica surface and act as nucleation sites.



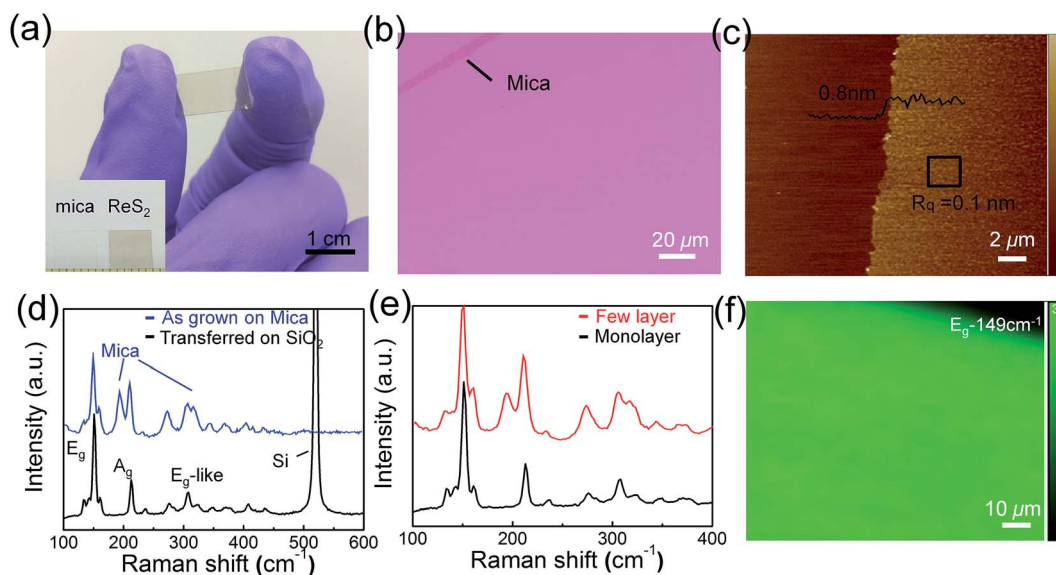


Fig. 2 (a) Digital photograph of as-grown ReS₂ films on mica substrate. Inset is digital picture of mica substrate with and without ReS₂ films. (b) Optical image and (c) AFM tomography image of ReS₂ monolayer films. Inset is the height profile. (d) Raman spectra of as-grown ReS₂ films on mica and SiO₂ substrate. (e) Comparison of Raman spectra between monolayer and multilayer ReS₂ films on mica substrate. (f) Raman mapping of monolayer ReS₂ films at E_g peak 149 cm⁻¹.

symmetry splitting introduced by the distorted 1T phase of ReS₂.¹⁵ Compared to as-synthesized monolayer films on mica, E_g and A_g mode of transferred samples present a slight blueshift by 2.86 and 2.19 cm⁻¹. This could be attributed to the release of tensile strain which arises from the lattice mismatch between ReS₂ and mica. Raman mode shift is considered as an effective approach to evaluate the strain in TMDs.^{31,32} For example, the E_{2g} and A_{1g} peaks of monolayer MoS₂ grown on mica both are blueshifted about 3–4 cm⁻¹ due to compressive strain.³⁰ In this work, red shift of Raman modes in ReS₂ films indicates the existence of tensile strain. With the increase of film thickness, tensile strain would accumulate and it has to be released by splitting ReS₂ films, resulting in lots of cover-like projections as shown in Fig. S3.† Consistent with previous studies,^{12,15} Raman peaks of multilayer films do not show any difference compared to monolayer samples due to the weak interlayer coupling effect (Fig. 2e). Raman mapping of monolayer ReS₂ films was also performed at E_g mode position, and a homogeneous feature in large area up to 100 μm² scale is presented.

Compared with previous work,^{23–25} where only isolated ReS₂ domains appeared on substrate, the products obtained in our experiment are continuous films. Re₂O₇ is used as catalyst, which has the lowest melting point (220 °C) among all the rhenium oxide. During the growth process, Re₂O₇ would rapidly volatilize into vapour phase, and then be reduced by sulfur vapour into Re₂O_{7-x} species.²³ Subsequently, they are conveyed downstream by Ar gas and absorbed on mica surface, acting as active nucleation sites for further ReS₂ growth.

The morphology of ReS₂ is very sensitive to the growth time, as shown in Fig. 3. At the beginning, intermediate Re₂O_{7-x} species absorbed on mica surface could act as nucleation sites. As temperature goes up, Re₂O_{7-x} species are further sulfurized into ReS₂ with incoming sulfur vapour. Due to low atoms

migration energy of mica surface, ReS₂ could easily grow along the in-plane direction into monolayer domains. These isolated monolayer ReS₂ domains have a small size ranging from 500 nm to 2 μm (Fig. 3a). It is worthy noted that these isolated domains exhibit irregular dendrite shapes, which is very different from traditional TMDs, such as MoS₂ and WS₂.^{22,30} This kind of anisotropic growth is attributed to the anisotropic interfacial energy introduced by distorted 1T crystal structure in ReS₂. Volatile Re₂O₇ would easily be run out before reaching 600 °C, and it cannot act as a stable reaction source for growing thick ReS₂ films. The main function of Re₂O₇ is supplying enough nucleation sites for further ReS₂ films growth. When the temperature rises up to 600 °C, Re powder starts to volatilize and react directly with sulfur into ReS₂, which plays the decisive part in multilayer continuous films growth. Isolated ReS₂ domains start to interconnect with each other into continuous films as reaction undergoing (Fig. 3b). Typically, when two separate ReS₂ domains gather, the edges tend to interconnect rather than overlap, which is the typical characteristic of 2D materials growth.³⁰ After the growth of continuous monolayer films is completed, following ReS₂ layers start to emerge on the surface of the first layer (Fig. 3c). We believe that these ReS₂ layers should be originated from the zero dimensional or quasi-one-dimensional structures as indicated with white circles in Fig. 3c, suggesting that the ReS₂ growth obeys Stranski-Krastanov growth mode.³³ In other words, except the first layer, the next layers follow the island-based growth mode, which would result in multilayer films with relatively rough surface. By prolonging the growth time to 30 min, we could obtain multilayer films with thickness about 4.2 nm (Fig. 3d). The surface roughness difference between monolayer and multilayer films further confirms our suggestion (R_q are 0.10 and 0.54 nm, respectively). XRD pattern of multilayer films shows a strong



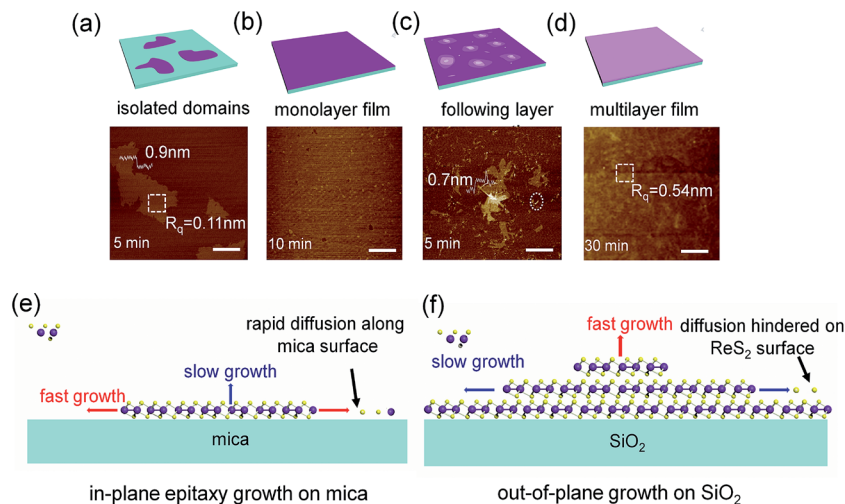


Fig. 3 (a–d) Schematic (upper) and corresponding AFM topography images (lower row) showing the possible growth mechanism of ReS_2 : (a) 5 min. (b) 10 min. (c) 15 min. (d) 30 min. Scale bar: $1 \mu\text{m}$. (e, f) Schematic growth model for different stages of ReS_2 films on mica substrate.

(001) orientation due to the layer-by-layer growth mode on mica substrate (Fig. S4, ESI†).

The substrate plays a crucial role in the growth process. As shown in Fig. 3e, atomically smooth and inert chemically surface of mica could facilitate the atoms migration along the surface due to the low energy barrier,^{34,35} which make the first layer growth on mica surface-dominated and result in continuous monolayer structure. Compared with the very flat epitaxial growth of the first layer on mica, following layers always have relatively rough surface. Because the surface of ReS_2 could not provide such low energy for atoms migration as mica, thus the in-plane growth is hindered (Fig. 3f). Compared with other TMDs, such as MoS_2 and WS_2 , and weak van der Waals interaction between layers of the ReS_2 could also facilitate the 3D structure growth,³⁶ and accelerate the atoms migration along vertical direction, leading to the rough surface. As shown in Fig. S5,† the products grown on SiO_2 substrate exhibit flower-like structures. This is because the rough surface of SiO_2 with dangling bonds could increase the energy barrier for atoms migration and prohibit ReS_2 from growing along in-plane directions.

Re could directly react with S vapour to form ReS_2 under temperature range of $600\text{--}900^\circ\text{C}$.³⁷ However, the high melting temperature of Re (about 3150°C) leads to a very low nucleation rate during the growth process. In our experiments, no products appear without Re_2O_7 under the identical growth conditions. We believe the severe volatilization of Re_2O_7 , which provide adequate nucleation sites for the later ReS_2 growth, plays crucial part in monolayer film forming. The dosage of Re_2O_7 is also important for the growth process. In this experiment, Re_2O_7 is only one-tenth of the Re power. Only Re_2O_7 as precursor would result in numerous nanoparticles without ReS_2 on mica surface. XPS spectrum shows these particles are mainly composed of $\text{Re}_2\text{O}_{7-x}$ intermediate species (Fig. S6, ESI†). Meanwhile, more Re_2O_7 could lead to thick film consist with lots of tiny grains with large surface roughness (Fig. S7, ESI†).

X-ray photoelectron spectroscopy (XPS) was performed to determine the stoichiometry of ReS_2 films. Fig. 4a shows the survey spectrum of the monolayer ReS_2 films. The core $2p_{3/2}$ and $2p_{1/2}$ level peaks for sulfur are located at 161.37 and 162.63 eV (Fig. 4b). The two dominant level peaks for rhenium are located at 42.6 and 45.1 eV (assigned to $\text{Re } 4f_{7/2}$ and $\text{Re } 4f_{5/2}$ respectively), consistent with previous study³⁶ (Fig. 4c). The prominent $4f_{7/2}$ and $4f_{5/2}$ peak for Re_2O_7 ,³⁸ located at 40.1 and 42.3 eV, are not observed, suggesting the Re_2O_7 has been completely sulfurized.

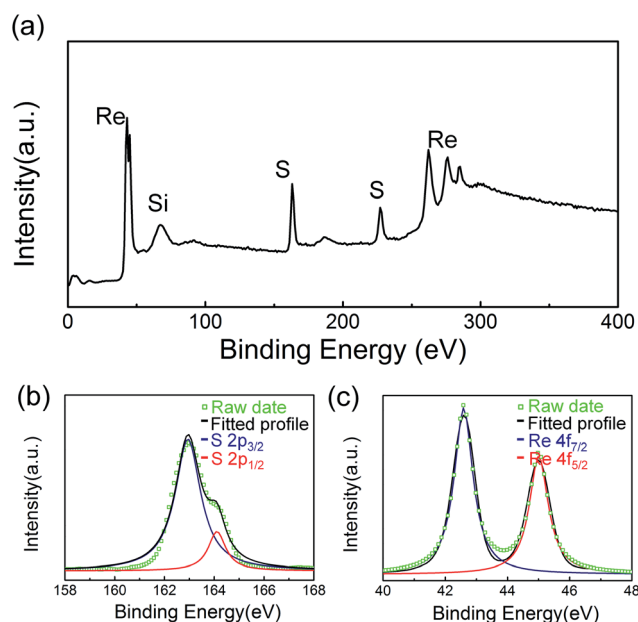


Fig. 4 XPS spectra of monolayer ReS_2 films. (a) XPS survey spectrum of the sample. Three kinds of elements are present: Re and S (from the sample), Si (from the substrate). (b) XPS spectrum of S $2p$ and (c) XPS spectrum of Re $4f$.



Transmission electron microscope (TEM) analysis was also conducted to explore the crystal structure of ReS₂ films. Using a traditional poly(methyl methacrylate) (PMMA) wet transfer method, ReS₂ films were transferred onto carbon-supported copper grids. From the low-magnification TEM image, a continuous membrane with obvious folding could be clearly observed (Fig. 5a). Edge folding often appears in TEM analysis of two-dimensional materials, like MoS₂ and graphene,^{30,39} which is probably introduced by the transfer process. HRTEM reveals perfect single-crystal structure of the monolayer films, and clear DS-chains of Re atoms could be seen (Fig. 5b). The angle between *a*[100] and *b*[010] is about 60°, which is agree well with the date published on exfoliated ReS₂ samples.¹⁵ Typical (100) and (010) crystal plane spacing are 0.31 and 0.35 nm, respectively, confirming the high crystallinity of monolayer films. As for the multilayer films, it is demonstrated to be polycrystal consist of many tiny grains with random orientations (Fig. 5c). Fig. 5d shows the HRTEM and corresponding SAED pattern. AED pattern shows clear dim diffraction rings, which further prove the films are polycrystal.

Field effect transistors (FETs) are fabricated based on multilayer ReS₂ films (Fig. 6a). It should be noted that previous reports only focus on single crystal ReS₂,^{11,15,24} whereas the performance of multilayer polycrystalline films are still unexplored. The gate-dependent transport curves show a typical n-type behaviour with maximum ON/OFF ratio about 4×10^3 ,

and the threshold voltage is about -27 V with V_{ds} of 3 V. The carrier mobility could be calculated using following equation:⁴

$$\mu = \frac{L}{WV_{ds}C_g} \times \frac{dI_{ds}}{dV_{bg}}$$

where L , W , C_g , V_{ds} , V_{bg} and I_{ds} are the channel length, width, the gate capacitance per unit area, source-drain voltage, gate voltage and source-drain current, respectively. The carrier mobility of the device can be up to $0.27 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$. Compared with single-crystal monolayer ReS₂ FET with carrier mobility $1.43 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ and ON/OFF ratio 10^6 (Fig. S8, ESI†), FET based on multilayer ReS₂ films exhibits degraded electrical performance, which is attributed the carrier scattering effect in grain boundary.⁴¹

The optoelectronic properties of multilayer ReS₂ films were also evaluated based on the same FET device. Fig. 7a shows the schematic image of the device, a 500 W xenon lamp was used as the light source with illumination power ranging from 0.25–0.56 mW cm^{-2} . Fig. 7b shows the I - V curves under different incident light intensities. Obviously, the channel current increases significantly under illumination compared to dark condition. Fig. 7c and d shows the photocurrent and responsivity as a function of light power. Here, photocurrent (ΔI) is the difference between I_{on} and I_{off} with a V_{ds} of 1 V, and the responsivity is defined as $R = \Delta I/P \times S$ (P is the light power intensity, and S is the illuminated area of the devices). Both ΔI

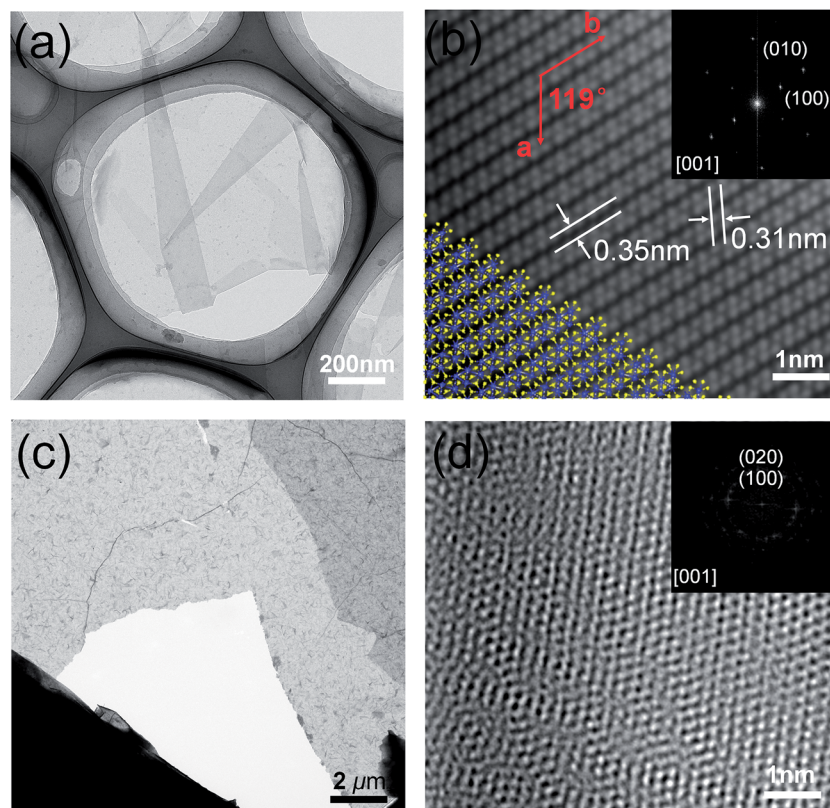


Fig. 5 (a and b) Bright-field TEM and HRTEM images of monolayer ReS₂ films. (c and d) Bright-field TEM and HRTEM (STEM) images of few-layer ReS₂ films. Insets in (b and d) show corresponding SAED patterns.



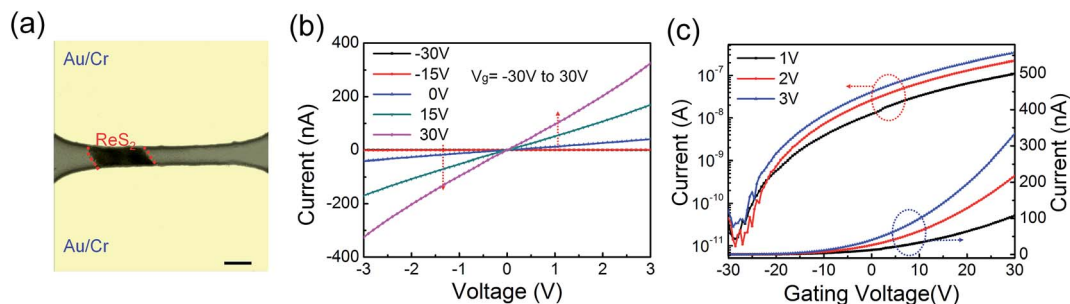


Fig. 6 Electrical properties of multilayer ReS₂ films. (a) Optical image shows the as-fabricated FET device. (b, c) Output and transfer curve of this device. Back-gate voltage (V_g) sweeping: -30 to 30 V; source-drain bias (V_{ds}): 1 V to 3 V.

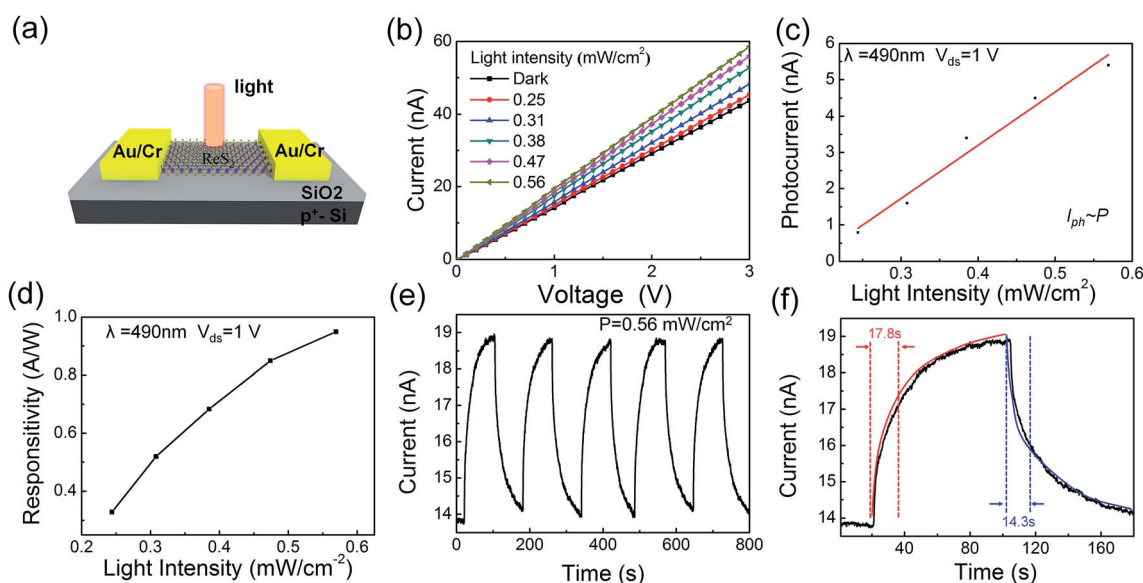


Fig. 7 Optoelectronic performance of the phototransistor based on multilayer ReS₂ films on SiO₂ substrate. (a) Schematic image of the device. (b) I - V curves of device in dark and under 490 nm illumination with different intensities. (c, d) Photocurrent and photoresponsivity as a function of incident power. (e, f) Time-dependent photoresponse of the device under 490 nm illumination with intensity of 0.56 mW cm^{-2} and bias of 1 V.

and R raise slightly with the increase of intensity. The responsivity is estimated to be 0.98 A W^{-1} at an irradiation intensity of 0.56 mW cm^{-2} , almost one order smaller than that in single-crystal monolayer ReS₂ based devices (Fig. S9, ESI†). Time-resolved photoresponse behaviour was probed by switching the light on and off with a time interval of 100 s under a bias of 1 V. The device exhibits a stable and repeatable response to the illumination after 5 cycles (Fig. 7e). A considerable enhancement of current (19 nA) is measured when the device is illuminated by 490 nm light. The response and recovery time is calculated by considering the time it takes to reach 80% of the final values, which are 17.8 and 14.3 s, respectively. Both of the rising time and recovery time of ReS₂ films devices are larger than that of exfoliated ReS₂ nanosheet,^{13,40} but comparable to the bilayer polycrystalline film.⁴¹ Generally, the active sites of the material significantly affect the rise and decay times.^{27,42} Based on our measurement condition, gas molecules in air such as oxygen can easily absorbed on the surface of films. Under light illumination, photo-generated carriers would first fill the trap states, which

prolongs the response time. We believe the improvement of measurement condition would be effective to improve the responsive speed.⁴³

Conclusions

In conclusion, we develop a facile method to obtain large-area continuous ReS₂ films on mica substrate using CVD method. Re₂O₇ acting as catalyst could provide enough seeds for continuous films growth. The multilayer films exhibits a layer-by-layer growth mode with strong (001) orientation. FET measurement indicates that multilayer ReS₂ films exhibit n-type transport behaviour with ON/OFF ratio about 4×10^3 and carrier mobility of $0.27 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$. The optoelectronic properties of ReS₂ films are also investigated, with the highest photoresponsivity of 0.98 A W^{-1} at an illumination intensity of 0.56 mW cm^{-2} . Our work suggests that these scalable ReS₂ films may be an ideal candidates for future high performance electronic and optoelectronic devices.



Acknowledgements

This work was financially supported by National Natural Science Foundation of China (No. 51572057) and Shandong Provincial Key Lab of Special Welding Technology, Harbin Institute of Technology at Weihai. The authors thanks Prof. Ping-An Hu for assistance on electrical measurements.

Notes and references

- N. R. Pradhan, D. Rhodes, S. Feng and Y. Xin, *ACS Nano*, 2014, **8**, 5911–5920.
- H. Fang, S. Chuang, T. C. Chang, K. Takei, T. Takahashi and A. Javey, *Nano Lett.*, 2012, **12**, 3788–3792.
- Z. Yin, H. Li, H. Li, L. Jiang, Y. Shi and Y. Sun, *ACS Nano*, 2012, **6**, 74–80.
- B. Radisavljevic, A. Radenovic, J. Brivio, V. Giacometti and A. Kis, *Nat. Nanotechnol.*, 2011, **6**, 147–150.
- Y. Huang, *et al.*, *ACS Nano*, 2014, **8**, 10743–10755.
- Y. H. Huang, C. C. Peng, R. S. Chen, Y. S. Huang and C. H. Ho, *Appl. Phys. Lett.*, 2014, **105**, 093106.
- O. Lopez-Sanchez, D. Lembke, M. Kayci, A. Radenovic and A. Kis, *Nat. Nanotechnol.*, 2013, **8**, 497–501.
- I. G. Lezama, A. Arora, A. Ubaldini, C. Barreateau, E. Giannini, M. Potemski and A. F. Morpurgo, *Nano Lett.*, 2015, **15**, 2336–2342.
- Y. Li, *et al.*, *Adv. Funct. Mater.*, 2016, **26**, 293–302.
- Y. Li, C.-Y. Xu and L. Zhen, *Appl. Phys. Lett.*, 2013, **102**, 143110.
- Y.-C. Lin, H.-P. Komsa, C.-H. Yeh, T. Björkman, Z.-Y. Liang, C.-H. Ho, Y.-S. Huang, P.-W. Chiu, A. V. Krasheninnikov and K. Suenaga, *ACS Nano*, 2015, **9**, 11249–11257.
- S. Tongay, *et al.*, *Nat. Commun.*, 2014, **5**, 3252.
- F. Liu, *et al.*, *Adv. Funct. Mater.*, 2016, **26**, 1169–1177.
- Q. Zhang, *et al.*, *Adv. Mater.*, 2016, **28**, 2616–2623.
- D. A. Chenet, O. B. Aslan, P. Y. Huang, C. Fan, A. M. van der Zande, T. F. Heinz and J. C. Hone, *Nano Lett.*, 2015, **15**, 5667–5672.
- O. B. Aslan and D. A. Chenet, *ACS Photonics*, 2016, **3**, 96–101.
- N. R. Pradhan, *et al.*, *Nano Lett.*, 2015, **15**, 8377–8384.
- Y. Li, C.-Y. Xu, P. Hu and L. Zhen, *ACS Nano*, 2013, **7**, 7795–7804.
- Y. Li, C.-Y. Xu, B.-Y. Zhang and L. Zhen, *Appl. Phys. Lett.*, 2013, **103**, 033122.
- Y. Zhang, *et al.*, *ACS Nano*, 2013, **7**, 8963–8971.
- A. L. Elías, *et al.*, *ACS Nano*, 2013, **7**, 5235–5242.
- W. Zhang, M.-H. Chiu, C.-H. Chen and W. Chen, *ACS Nano*, 2014, **8**, 8653–8661.
- K. Keyshar, *et al.*, *Adv. Mater.*, 2015, **27**, 4640–4648.
- X. He, F. Liu, P. Hu, W. Fu, X. Wang, Q. Zeng, W. Zhao and Z. Liu, *Small*, 2015, **11**, 5423–5429.
- F. Cui, *et al.*, *Adv. Mater.*, 2016, **28**, 5019–5024.
- Y. Zhou, Y. Nie, Y. Liu, K. Yan, J. Hong, C. Jin, Y. Zhou, J. Yin, Z. Liu and H. Peng, *ACS Nano*, 2014, **8**, 1485–1490.
- Q. Wang, M. Safdar, K. Xu, M. Mirza, Z. Wang and J. He, *ACS Nano*, 2014, **8**, 7497–7505.
- I. R. Beattie, T. R. Gilson and P. J. Jones, *Inorg. Chem.*, 1996, **35**, 1301–1304.
- X.-F. Qiao, J.-B. Wu, L. Zhou, J. Qiao and W. Shi, *Nanoscale*, 2016, **8**, 8324–8332.
- Q. Ji, *et al.*, *Nano Lett.*, 2013, **13**, 3870–3877.
- Z. H. Ni, T. Yu, Y. H. Lu, Y. Y. Wang, Y. P. Feng and Z. X. Shen, *ACS Nano*, 2008, **2**, 2301–2305.
- M. Huang, H. Yan, T. F. Heinz and J. Hone, *Nano Lett.*, 2010, **10**, 4074–4079.
- E. Bauer and J. H. van der Merwe, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1986, **33**, 3657.
- Q. Wang, K. Xu, Z. Wang, F. Wang, Y. Huang, M. Safdar, X. Zhan, F. Wang, Z. Cheng and J. He, *Nano Lett.*, 2015, **15**, 1183–1189.
- X. Li, F. Cui, Q. Feng, G. Wang, X. Xu, J. Wu, N. Mao, X. Liang, Z. Zhang and J. Zhang, *Nanoscale*, 2016, **8**, 18956–18962.
- J. Gao, L. Li, J. Tan, H. Sun, B. Li, J. C. Idrobo, C. V. Singh, T.-M. Lu and N. Koratkar, *Nano Lett.*, 2016, **16**, 3780–3787.
- C.-H. Ho, *Opt. Express*, 2005, **13**, 8–19.
- Y. Jia, L. Duan, D. Zhang, J. Qiao and G. Dong, *J. Phys. Chem. C*, 2013, **117**, 13763–13769.
- J. C. Meyer, A. K. Geim, M. I. Katsnelson, K. S. Novoselov, T. J. Booth and S. Roth, *Nature*, 2007, **446**, 60–63.
- E. Liu, *et al.*, *Adv. Funct. Mater.*, 2016, **26**, 1938–1944.
- M. Hafeez, L. Gan, H. Li, Y. Ma and T. Zhai, *Adv. Funct. Mater.*, 2016, **26**, 4551–4560.
- Y. Jiang, W. J. Zhang, J. S. Jie, X. M. Meng, X. Fan and S. T. Lee, *Adv. Funct. Mater.*, 2007, **17**, 1795–1800.
- J. Zhou, Y. Gu, Y. Hu, W. Mai, P.-H. Yeh, G. Bao, A. K. Sood, D. L. Polla and Z. L. Wang, *Appl. Phys. Lett.*, 2009, **94**, 191103.

