RSC Advances



View Article Online

View Journal | View Issue

PAPER



Cite this: RSC Adv., 2017, 7, 22094

Received 9th January 2017 Accepted 13th April 2017

DOI: 10.1039/c7ra00342k

rsc.li/rsc-advances

1. Introduction

3-d transition metal nitrides, such as CrN and TiN, combine attractive physico-chemical, electronic, optical, magnetic, mechanical and thermal properties.¹ These composite materials have shown a wide range of technical applications in various key industrial areas. For example, due to their extreme hardness, high melting point and superior chemical stabilities, nitride-based ceramic coatings have been used for high-speed machining,² cold forming, moulding dies,³ sliding parts and rocket nozzles.⁴ It is noted that CrN has significant superiority over TiN in wear, corrosion resistance,⁵ and anti-oxidative properties at temperatures up to 700 °C.⁶

Experimental and predicted mechanical properties of $Cr_{1-x}AI_xN$ thin films, at high temperatures, incorporating *in situ* synchrotron radiation X-ray diffraction and computational modelling⁺

Ehsan Mohammadpour,^a Zhong-Tao Jiang, ^[10]*^a Mohmmednoor Altarawneh, ^[10]^b Nicholas Mondinos,^a M. Mahbubur Rahman,^a H. N. Lim,^c N. M. Huang, ^[10]^d Zonghan Xie,^e Zhi-feng Zhou^f and Bogdan Z. Dlugogorski^b

Cr_{1-x}Al_xN coatings, synthesised by an unbalanced magnetic sputtering system, showed improved microstructure and mechanical properties for ~14–21% Al content. *In situ* SR-XRD analysis indicated various crystalline phases in the coatings that included: CrN, AlN, α -Cr with small amounts of AlO₂ and Al₂O₃ over the 25–700 °C range. Al doping improves resistance to crystal growth, stress release and oxidation resistance of the coatings. Al doping also enhances the coating hardness (*H*) from 29 to 42 GPa, elastic modulus (*E*) from 378 to 438 GPa and increased the resistance to deformation. First-principles and quasi-harmonic approximation (QHA) studies on bulk CrN and AlN were incorporated to predict the thermo-elastic properties of Cr_{1-x}Al_xN thin film coatings in the temperature range of 0–1500 °C. The simulated results at *T* = 1500 °C give a predicted hardness of *H* = ~41.5 GPa for a ~21% Al doped Cr_{1-x}Al_xN coating.

Vacuum sputtering technique is an outstanding method for synthesising high quality CrN ceramic thin film coatings with superior structural, mechanical and tribological properties.^{3,4} These studies have been extended to ternary combinations of CrN with doping of other elements, e.g., Al,⁷ Si,^{8,9} Ti,¹⁰ V,¹¹ Nb¹² or Cu.13 The doping elements mainly distribute in the amorphous layers which surrounds the columnar structured nanocrystalline (nc-) CrN grains. These amorphous layers can effectively enhance many mechanical properties of such thin film coatings. In particular, Al has shown great potential in improving the hardness¹⁴ as well as corrosion and oxidation resistances.15,16 This Cr-Al-N framework is of particular interest as they can form aluminium oxides on the top surface layers which suppress the diffusion of oxygen into the coating film. Analysis indicated the thin film coatings have enhanced chemical stability, good mechanical properties, good tribological properties, high temperature stability (up to 1000 °C) and exhibit excellent oxidation resistance.3,11,17

Characterisation at room temperature indicated that Al significantly affected the preferred grain orientation of the thin films.¹⁸ TEM and XPS studies indicated that amorphous AlN existed at the columnar CrN grain boundaries. The incorporation of Al introduced compressive residual stress to the coating material lattice. Moreover, the mechanical properties such as hardness, elastic strain and plastic deformation resistance of the coatings showed significant improvement over non doped CrN coating. These clearly indicated the crucial role of Al in

[&]quot;Surface Analysis and Materials Engineering Research Group, School of Engineering & Information Technology, Murdoch University, Murdoch, WA 6150, Australia. E-mail: Z. Jiang@murdoch.edu.au; Tel: +61 8 9360 2867

^bSchool of Engineering & Information Technology, Murdoch University, Murdoch, WA 6150, Australia

^cDepartment of Chemistry, Faculty of Science, University Putra Malaysia, 43400 UPM Serdang, Selangor, Malaysia

^dFaculty of Engineering, Xiamen University of Malaysia, Jalan Sunsuria, Bandar Sunsuria, 43900 Sepang, Selangor Darul Ehsan, Malaysia

^eSchool of Mechanical Engineering, University of Adelaide, SA 5005, Australia

¹Department of Mechanical and Biomedical Engineering, City University of Hong Kong, Kowloon, Hong Kong, China

[†] Electronic supplementary information (ESI) available. See DOI: 10.1039/c7ra00342k

producing a combination of superior hardness and excellent damage resistance in nc-CrN frameworks.¹⁸

In the present work, *in situ* synchrotron radiation powder diffraction (SR-XRD) beamline was employed to quantitatively investigate the changes in phase composition and microstructure of $Cr_{1-x}Al_xN$ thin film coatings as a function of Al content at different measurement temperatures in an atmospheric environment. Rietveld analysis of high resolution SR-XRD data and nanoindentation measurements will provide quantitative results, hence making a connection between microstate phases and mechanical properties of the coating. Density Functional Theory (DFT) calculations and Quasi-Harmonic Approximation (QHA) are employed to estimate the cell-volume, thermal expansion and hardness of $Cr_{1-x}Al_xN$ system at temperatures up to 1500 °C.

2. Experimental

2.1 Thin film coatings preparation

 $Cr_{1-x}Al_xN$ coatings with a total film thickness of 2 µm were deposited on M2 high speed steel substrates at 550 °C using a closed field unbalanced magnetron sputtering system (Teer Coatings Ltd, UK) with a four-target configuration. The coating process were described in detail by Li *et al.*¹⁸ and Wo *et al.*¹⁹

2.2 In situ SR-XRD characterizations

In situ SR-XRD experiments were performed on the Powder Diffraction beamline at the Australian Synchrotron with monochromatic X-rays, $\lambda = 0.827$ Å, verified by standard reference material (LaB6 660b) as supplied by the USA National Institute of Standards and Technology (NIST). The incidence angles of the X-ray beam (ω) were in range of 4.1 to 5.9°. The diffraction data were collected by a Mythen microstrip detector over the range of $10^\circ \le 2\theta \le 89^\circ$ in flat-plate asymmetric reflection geometry. The coated substrates were mounted on a Pt heating stage in an Anton Paar HTK-2000 high temperature furnace and aligned to the centre of rotation on the diffractometer. The heating rate was 10 °C min⁻¹ from ambient temperature up to 700 °C in air atmosphere.20 The temperature difference between the sample surface and Anton Parr furnace was calibrated manually before the in situ experiments carried out. Diffraction patterns were obtained at three temperatures (25, 200 and 500 $^{\circ}$ C) below and two temperatures (600 and 700 °C) above the deposition temperature of 550 °C due to time constrains at the synchrotron facilities. Each data acquisition at designated temperature was 120 s.

The diffraction patterns were analysed with the TOPAS v5 academic software²¹ using the Rietveld method.²² The atomic positions for the structural models, based on the main identified phases of the coating, were extracted from Crystallography Open Database.²³ Modelled peak shapes were used to determine the volume-weighted mean crystallite sizes (L_{vol}) and microstrain (e_0). An accurate description of the peak shapes, in all of the diffraction patterns, was obtained by applying correctional functions for a flat plate in fixed incident beam geometry.^{24–26} The initial parameters for the refinement, as well as phase determination, were obtained from the JCPDS cards.

2.3 Nanoindentation measurements

Young's modulus and hardness of the coatings were measured, at room temperature (\sim 25 °C), by a calibrated Ultra-Micro Indentation System, UMIS-2000 equipped with a Berkovich indenter. A maximum load of 20 mN was applied in 10 increments. The loading rate used was 2.5 mN s⁻¹ and represented the static response of the materials. The area function of the indenter tip was calibrated using a standard fused silica specimen. Load control method with a maximum loading of 5 mN was employed for calibration. The peak loading is based on the considerations that the maximum displacement during indentation should be no more than 10% of the coating thickness. For better resolutions, the number of test points were 20 for loading and 20 for unloading.

2.4 DFT computational simulations

All structural optimizations and calculations were performed with the generalized gradient approximations (GGA) using the PW91 functional²⁷ as implemented in the VASP package.²⁸ Energy cut-off for plane waves included in the expansion of wave functions was 500 eV. Pseudo-atomic calculations were performed for Cr: $3d^54s^1$, Al $3s^23p^1$ and N $2s^22p^3$. *k*-Point sampling for reciprocal space integration was optimised and adopted a 21 \times 21 \times 21 Gamma scheme for cubic structures. CrN in contrast to other early transition metal (Ti, Sc, V, ...) nitrides, needed additional effort due to the local magnetic moments of the Cr atoms in CrN structure. The calculations for bulk CrN was carried out for both paramagnetic (CrN-m) and non-magnetic (CrN-nm) structures.

The thermodynamic properties of cubic phases were explored using the quasi-harmonic Debye–Slater model as implemented in Gibbs2 code.²⁹ In all temperature dependent property calculations, the Vinet's equation of state was adopted to fit the energy–volume curves.³⁰ To calculate the elastic constants and determine the second derivatives, we used a finite differences method as implemented in VASP package (IBRION = 6).

3. Results and discussion

3.1 Structural analysis *via in situ* synchrotron radiation XRD technique

The surface elemental composition and chemical bonding state, obtained from previous XPS analysis and full discussion of the coatings by Li *et al.*¹⁸ indicated the existence of CrN, Cr_2O_3 , AlN and Al_2O_3 phases in the coating surface region, are summarised in Table 1.

The JCPDS cards used for phase identification in the SR-XRD spectra: 00-076-2494 (CrN), 00-025-1495 (AlN), 00-901-4565 (AlO₂), 00-900-9676 (Al₂O₃), 00-035-0803 (Cr₂N), 00-901-3484 (α -Fe) and 00-041-1351 (Fe₃W₃C). The atomic positions for Fe₃W₃C were taken from.³¹ The *in situ* full patterns of all samples, for all temperatures, are included in the ESI section (Fig. S1–S4†). For clarity the spectra for each temperature are combined in one plot for each sample.

 $\mbox{Table 1}$ Elemental composition of the $\mbox{Cr}_{1-x}\mbox{Al}_xN$ coatings determined by XPS

Coating notation	Element (at%)				
	Cr	Al	Ν	0	
C1	33.4	0	34.9	31.7	
C2-14% Al	27.5	14.2	31.8	26.6	
C3-17% Al	16.6	17.2	29.4	36.8	
C4-21% Al	19.1	21.4	30.4	29.0	

The Rietveld refinement of the diffraction spectra resulted in *R*-weighted pattern (R_{wp}) values in the range of 15–20. The largest errors in the refinement models are mainly due to the incomplete crystallographic data for the substrate. An example of a typical fitted spectrum is indicated by Fig. 1, which shows the Rietveld refinement for the C1 sample at 700 °C with $R_{wp} = 15$. The allowed *hkl* values for each phase are indicated by the small vertical bars. The difference between experimental and modelled results is indicated by the difference plot above the *hkl* bars.

Basically all 20 spectra shown in Fig. S1–S4 of the ESI† section were analysed by Rietveld analysis with similar plots as Fig. 1 producing data used for further analysis as indicated in Fig. 2–6. The phases of the C1 sample at different temperatures was similar to the sample (denoted as CrN-M2ST) recorded in Mohammadpour *et al.*²⁰

SR-XRD diffraction patterns at ~25 °C of as-deposited $Cr_{1-x}Al_xN$ coatings are depicted in Fig. 2. For clarity the 2θ range in Fig. 2 is from 18° to 35° 2θ consisting of the main peaks for each individual phase.

The compositions of the main phase in as-deposited CrAlN thin film coatings were determined to be CrN and AlN phases.

Thin film coating C1 showed strong intensity in (200), (220) directions at $\sim 22.8^{\circ}$ and $\sim 32.4^{\circ}$ 2θ , respectively for the CrN phase. In the Al doped samples the dominant orientations of the CrN phase are in the (200) and (111) planes. These diffraction peaks became broader and of lower intensity as the Al content increased (C3-17% Al and C4-21% Al) indicating a combination of small grain size and high concentration of possible amorphous nature. The peaks at 20.2°, 23.4° and 33.3° 2θ in the C4-21% Al sample are identified as diffractions from c-AlN (111), (200) and (220) planes, with some overlap with substrate (α -Fe) diffraction peaks. The confirmation of the AlN phase in SR-XRD pattern agrees with the NEXAFS results by Mahbubur et al.32 and XPS analysis by Li et al.18 from similar coatings. All samples in asdeposited condition (~25°) had Al₂O₃ and Cr₂O₃ phases detected by XPS but SR-XRD did not show any noticeable peaks related to those oxides. This could be due to (i) most of oxygen accumulate in the top surface layers of thin films, and (ii) possible amorphous state of the oxide phases. However, at temperatures >500 °C minor amounts of Al₂O₃ and metastable AlO₂ were observed from the diffraction patterns (Fig. 3).

Fig. 3 indicates the main phases of these coatings at 700 °C. Diffraction spectra of CrAlN coatings at 700 °C show α -Cr, AlO₂ and Al₂O₃ phases. Peaks at 22.60°, 26.16° and 37.33° 2 θ can be determined as α -Cr (111), (200) and (220), respectively. Agouram *et al.*³³ concluded that the phase content of CrN_x-based coatings is strongly influenced by deposition processes. The presence of α -Cr in CrN_x coatings can reduce the hardness; however, it may improve the coating toughness or as a metallic source which diffuse outward and contribute in the formation of surface oxidation layers. Although hexagonal-structured Cr₂N (h-Cr₂N) exhibits higher hardness than cubic CrN, it is more susceptible to oxidation which can deteriorate the thermal stability of the



Fig. 1 Experimental SR-XRD and Rietveld refinement profile of C1 coating at 700 °C. The allowed *hkl* values for each phase are indicated by the small vertical bars.



Fig. 2 SR-XRD diffraction patterns at \sim 25 °C of as-deposited Cr_{1-x}Al_xN coatings with various Al contents. The dashed vertical lines indicate substrate peaks.



Fig. 3 SR-XRD diffraction patterns at 700 °C of Cr_{1-x}Al_xN coatings with various Al contents. The dashed vertical lines indicate substrate peaks.

C1 thin film with considerable h-Cr₂N content especially at high temperatures.²⁰ It seems Al doping successfully prevents the formation of h-Cr₂N in the CrAlN coatings even at lowest Al percentage (C2-14% Al) in the temperature range of study. Our observations did not confirm any possible existence of Cr_2O_3 crystalline phases in the column structure of coatings; that may indicate (i) small amorphous structures distributing in grain boundary regions, and (ii) either amorphous or crystalline

phase structure in the very top surface layers of thin films where normal XRD cannot fully determine the phases (Grazing incidence SR-XRD would be able to provide more precise surface crystalline information). The shifts of diffraction peaks from substrate, especially for Fe (111) at 23.8° is probably due to the thermal expansion of the unit cells at elevated temperatures. Peaks related to CrN and AlN phases exhibited some sharpening with increasing temperature. From the Rietveld refined SR-XRD data, the lattice constants of the CrN and c-AlN phases in the $Cr_{1-x}Al_xN$ coatings were calculated and plotted as a function of temperature in Fig. 4. The lattice constant of CrN decreases continuously from 4.21 Å to 4.17 Å for C1 and to 4.19 Å for the Al doped samples. The overall changes of the lattice constant for the CrN phase are slightly smaller in the samples with Al doping than without doping. Overall, the lattice parameter values agree with the experimental values from other sources, such as Barshilia *et al.*³⁴ (4.09 to 4.16 Å) and Elangovan *et al.*³⁵ (4.15 to 4.17 Å). Lattice constants of c-AlN (in sample C4-21% Al) increased steadily from approximately 4.09 Å at room temperature to 4.1 Å at 700 °C indicating unit cell expansion by 0.5%. The lattice constant of α -Cr phase, in all samples at 700 °C, were 3.65 Å \pm 10%.

The crystallite size and microstrain as a function of temperature, from Rietveld analysis, for the CrN phase are shown in Fig. 5. The Effect of adding Al onto the microstructure of the as-deposited coatings is more pronounced in C3-17% Al and C4-21% Al coatings exhibiting about 30% smaller crystallite sizes compared to Al free and C2-14% Al samples. The crystallite sizes in all samples gradually increased in the temperature range 25 to 600 °C and then abruptly increased from 600 to 700 °C. C1 and C2-14% Al show the highest crystallite size growth by $\sim 100\%$ and $\sim 67\%$, respectively. The higher Al content in C3-17% Al and C4-21% Al effectively prevented the crystallite growth of CrN during heating below 700 °C. Above, 600 °C the CrN crystallite size in C4-21% Al increases by \sim 34%. The most stable crystallite size for CrN is exhibited by C3-17% Al with an 8% crystallite growth. The crystallite sizes of the α-Cr phase are 1832 (C2-14% Al), 1780 (C3-17% Al) and 828 (C4-21% Al) nm in the 700 °C $Cr_{1-x}Al_xN$ samples.

Our observations showed that microstrain in C1 sample decreased by ~67% during heating which is almost ~1.7× of the microstrain release in samples with Al contents. The residual stress in these types of coatings could be due to two mechanisms (1) thermal stress and/or (2) accumulation of defects (dislocations, vacancies, *etc.*) in the coating layer. Thermal stress is a resultant of the difference in the coefficient



Fig. 4 CrN and c-AlN lattice constants, from Rietveld analysis, in $Cr_{1-x}Al_xN$ coatings as a function of the *in situ* experiment temperature.



Fig. 5 Effect of temperature on the CrN phase in the $Cr_{1-x}Al_xN$ coatings, from Rietveld analysis: (a) crystallite size (b) microstrain.

of thermal expansion (CTE) value for coating and substrate. The decline of internal stress in the coatings could be a result of absorption and annihilation of defects during annealing. It allows the defects to diffuse through the lattice and to be absorbed at grain boundaries and dislocation cores. This seems to be more pronounced in the binary CrN coating case. The presence of amorphous content in $Cr_{1-x}Al_xN$ coatings at grain boundaries release in the coatings which resist against coating softening.

The stress in the coating structure, from the changes in *d* space for CrN (220), is shown in Fig. 6. Similar plots for the (111), (200) and (311) planes of CrN phase are included in Fig. S5–S7 of the ESI section.† The stress (σ) in the coating were calculated based on eqn (1) using the elastic properties of CrN (elastic modulus, E = 300 GPa and Poisson ratio, $\nu = 0.28$).³⁶

$$\sigma = \frac{E}{(1-\nu)} \times \frac{d-d_0}{d_0} \tag{1}$$

where *d* and d_0 are from the experimental diffraction patterns and JCPDS card of CrN, respectively. Refinement data show that there is considerable amount of residual stress in the asdeposited coating structure. Stress values in Al free coating decreases progressively as temperature increases from 25 to



Fig. 6 Effect of temperature on the residual stress, from eqn (1), for CrN (200) orientation in $Cr_{1-x}Al_xN$ coating.

700 °C by ~56%. The stress release in samples containing Al is slower as Al content increases and reduced to only ~11% in C4-21% Al coating. Based on observations in Fig. 5(b) and 6, there is a direct correlation between residual strain, stress and Al at% in the coating composition.

3.2 Experimental mechanical properties of CrAIN coatings

Nano-indentation tests, for $Cr_{1-x}Al_xN$ coatings at as-deposited conditions and after the completion of the *in situ* SR-XRD experiment (heated samples), were undertaken at room temperature. The hardness (*H*) and elastic modulus (*E*) for these coatings are set out in Table 2 and Fig. S8 and S9 of the ESI section.† Analysis of Table 2 and the plots clearly indicate that there is no difference in *H* and *E*, within the experimental errors, between the as-deposited and heated sample. The resistance to deformation is indicated by H^3/E^2 and plotted in Fig. S10 of the ESI section.† The coatings with smaller crystallite size/higher microstrain present greater hardness (C3-17% Al and C4-21% Al samples) in as deposited condition. Improvement in coating hardness confirms the microstructure refining effect of Al dopants.

These findings are consistent with the quantitative SR-XRD measurements of crystallite size and micro strains in the Al doped coatings that exhibited very little hardness drop after heating. The calculated ratio of H^3/E^2 improved significantly in samples containing Al compared to the reference sample (C1).

Table 2 Nanoindentation, at 25 $^\circ\text{C},$ of $\text{Cr}_{1-x}\text{Al}_xN$ coatings: as-deposited and heated samples

	As-deposite	ed	Heated		
Sample	H (GPa)	E (GPa)	H (GPa)	E (GPa)	
C1	29 ± 2	378 ± 21	28 ± 2	376 ± 22	
C2-14% Al	35 ± 3	401 ± 27	35 ± 3	396 ± 25	
C3-17% Al	39 ± 3	421 ± 38	39 ± 3	421 ± 30	
C4-21% Al	42 ± 3	438 ± 32	42 ± 3	436 ± 32	

However, there is no change in resistance to deformation between as deposited and heated conditions of the samples. The combination of microstructure and nanoindentation studies indicate that the hardness of the as-deposited nanocrystalline coatings can be controlled by several parameters such as phase composition, crystallographic directions, crystallite size and residual stress/microstrain.³⁷

3.3 Density functional theory modelling

3.3.1 Structural optimization. The energy and lattice constant of bulk CrN-nm, CrN-m and c-AlN were optimized to find the most stable ground state. Fig. 7 shows changes in energy of the unit cell ($\Delta E = E - E_0$) *versus* primitive unit cell volumes for each structure used for predictions of thermodynamic properties.

Table 3 lists experimental lattice parameter values from SR-XRD data and DFT calculations from present work and available literature. The calculated lattice parameter for CrN-m (spin-polarised) is closer to SR-XRD results. The magnetic moment for each Cr atoms in CrM-m structure converged to 2.044 $\mu_{\rm B}$ which is within range of 1.9 to 2.28 $\mu_{\rm B}$ from calculations of Filippetti *et al.*³⁸

The QHA method (Debye model) was used to predict the thermal expansion of the studied bulk structures as a function of temperatures. Fig. 8(a) and (b) are plots of the experimental and theoretical lattice parameters, with respect to temperature, for bulk CrN and c-AlN phases in the temperature range 0-1500 °C, respectively. Fig. 8(a) also includes lattice constant calculated from the experimental thermal expansion coefficient, α (T), of Zhou *et al.*⁴⁶ with $a_0 = 4.14$ Å (standard CrN phase) at room temperature. Fig. 8(a) clearly shows the CrN-m model is highly compatible to the experimentally obtained data. The CrN lattice constant at room temperature for all samples differ from the simulation values within \sim 1.6%. Interestingly these converge to the simulated results with increasing temperature resulting in a \sim 0.4% difference at 700 °C. The most significant drop in lattice constant values of CrN occurs to C1 sample which is very close to QHA model at 700 °C.



Fig. 7 The change in energy of primitive unit cells for bulk CrN-nm, CrN-m and c-AlN as a function of cell volume.

	Lattice constant, a (Å)				
Phase	Experimental	Calculated			
CrN-nm	4.17-4.20 4.12 (ref. 39) 4.18 (ref. 40) 4.145 (ref. 41)	4.05 4.051 (ref. 42)			
CrN-m		4.12			
c-AlN	4.07	4.173 (ref. 43 and 44) 4.068 4.06 (ref. 45)			



Fig. 8 Theoretical and experimental lattice constants: (a) CrN and (b) c-AlN phases.

Quasi-binary systems like $Cr_{1-x}Al_xN$ coatings with a mixture of NaCl-structured CrN/AlN phases can form (Cr,Al)N solid solutions up to x = 0.75.⁴⁷ Phase segregation during deposition results in a mixture of fcc-structured Cr-rich (Cr,Al)N and AlN^{28,48} which will increase the coating hardness due to activation of solid-solution hardening mechanism. Theoretical⁴⁹ and experimental^{47,50} studies show that substitution of Cr atoms (r =1.4 Å) with Al atoms (r = 1.25 Å) can progressively decrease lattice constant of CrN phase as Al content increases. However, smaller Al atoms can incorporate into interstitial positions in CrN structure which can expand the lattice constant as observed for CrSiN coating. $^{\rm 19,51}$

The c-AlN lattice parameter in C4-21% Al are in good agreement with experimental data from Tkadletz *et al.*⁵² as indicated in Fig. 8(b). The experimental results are within $\sim \pm 0.3\%$, of the simulated QHA results up to 700 °C.

Overall the differences between the simulated and experimental results is due to the use of pure homogeneous bulk material in the Debye (QHA) model calculations, and the solid solution of different nitride phases of the thin film coatings.

3.3.2 Prediction of elastic constants. Three independent elastic constants are needed to describe the behaviour of cubic-structured materials to stress and can be used to describe the nature of the bonds between atoms in the crystals. These constants (C_{11} , C_{12} and C_{44}) were calculated for CrN-nm, CrN-m and c-AlN at zero K and zero pressure and listed in Table 4. The mechanical stability of these phases can be ascertained by using the following criterion for cubic systems:⁵³

$$C_{11} + 2C_{12} > 0, C_{11} - C_{12} > 0, C_{44} > 0, C_{11} > 0, C_{11} > B > C_{12}$$
(2)

All of these conditions were satisfied by the CrN and c-AlN phases indicating the mechanical stability of these structures. The overview of available literature as shown in Table 4 indicates that the reported calculated values for elastic constants and modulus are widely scattered and depend on the choice of *k*-points, cut-off energies, functional (*s*) and simulation packages. As an example, Yan *et al.*⁴² and Brik *et al.*⁵⁴ used the CASTEP module of Materials Studio package with different settings of *k* points grids and cut off energies for CrN-nm analysis. With 8 × 8 × 8 (400 eV) and 10 × 10 × 10 (320 eV) *k* points grids the authors calculated *C*₁₁ values of 526 and 546 GPa, respectively.

ESI section Tables S2 and S3[†] show the influence of *k*-points grids selection on the elastic constants of CrN-m and CrN-nm for the present project. To calculate the elastic constants for CrN-m and CrN-nm, *k*-point grids of $21 \times 21 \times 21$ and $31 \times 31 \times 31$ were chosen for the simulations. The results of these analyses are tabulated in Table 4. The calculated elastic constants for CrN-nm in present study are comparable (~2%) with calculations of Zhou⁴⁶ who used similar approach *via* VASP package and projector-augmented wave (PAW) method. The values of C_{11} and C_{12} for c-AlN underestimate by ~9% the calculations of Wang *et al.*⁴⁵ carried out by VASP package and PAW method.

The positive difference between C_{11} and C_{44} ($C_{11} - C_{44} = 582$ for CrN-nm), known as Cauchy pressure, indicates the metallic bonding nature of Cr-N structure⁵⁵ which is comparable to much smaller difference for AlN ($C_{11} - C_{44} = 83$). Cauchy pressure is also a sign of the ductility of materials, since it indicates the angular character of atomic bonding in metals. A positive Cauchy pressure reveals ductility of crystal, while a brittle material exhibits a negative value.⁵⁶ Bulk modulus (*B*), shear modulus (*G*) and Young's modulus (*E*) for each phase were estimated by applying Voigt–Reuss equations⁵⁷ and compared with available literature.

Paper

Table 4 Calculated elastic constants (C_{ij}), bulk modulus (B), Young's modulus (E), Poisson ratio (ν) of CrN (non-magnetic and spin-polarized) and c-AlN phases. Calculated and experimental data from literature are denoted by c and e, respectively

Phase	C_{11}	C_{12}	C_{44}	B (GPa)	G (GPa)	E (GPa)	ν
CrN-nm	$588.0 \\ 526^{c,42} \\ 502^{c,54} \\ 580^{c,46} \\ 546^{c,58}$	$215.1 \\ 226^{c,42} \\ 214^{c,54} \\ 210^{c,46} \\ 184^{c,58}$	$5.9 \\ 5^{c,42} \\ 4.1^{c,54} \\ 8^{c,46} \\ 20^{c,58}$	$340 \\ 326^{c,42} \\ 310^{c,54} \\ 327^{c,58} \\ 340^{c,59}$	44 35 ^{c,42} 58 ^{c,58}	126 101 ^{c,42} 163 ^{c,58} 100–250 ^{e,36}	$\begin{array}{c} 0.43 \\ 0.43^{\mathrm{c},42} \\ 0.41^{\mathrm{c},58} \end{array}$
CrN-m	$352 \\ 369^{c,42} \\ 348^{c,46}$	$88.7 \\ 125^{c,42} \\ 117^{c,46}$	$61.0 \\ 114^{c,42} \\ 74^{c,46}$	176.4 192 ^{c,42}	93.9 111 ^{c,25}	239.0 279 ^{c,42}	0.27 0.25 ^{c,42}
c-AlN	387 423 ^{c,45}	149 167 ^{c,45}	305 306 ^{c,45}	$254.3 \\ 251^{c,45} \\ 221^{e,60}$	209	303	0.2788

The calculated results, at 0 K, from Table 4 are spread out but have similar magnitudes including the 2 experimental values at \sim 25 °C. The *k*-point grids that optimised the calculation of the elastic constants and unit cell volumes will be used in the DFT and QHA simulations.

QHA method for CrN and c-AlN are used to predict the hardness of two binary compounds (CrN and c-AlN) in a wide temperature range (25–1500 °C). The method for calculating the hardness of materials is based on atomic potentials and bond strength developed by Šimůnek and Vackář⁶¹ for single crystals and then extended to binary compounds. The hardness of a crystal can be calculated by introducing its bond strength (S_{ij}):

$$S_{ij} = \frac{\sqrt{e_i e_j}}{d_{ij} n_{ij}} \tag{3}$$

where e_i and e_j are atomic potentials of *i* and *j* atoms, respectively, and are defined as the valence electron number divided by the radius of the atom; d_{ij} is the distance of the nearest atom *j* to atom *i* and n_{ij} is the number of bonds between atom *i* and its neighbouring atoms *j*.

For binary compounds, the difference between atomic potentials is:

$$f_e = \left[\frac{e_i - e_j}{e_i + e_j}\right]^2 \tag{4}$$

The hardness of such a system with volume (*V*) is defined as:

$$H = \left(\frac{C}{V}\right) \frac{\sqrt{e_i e_j}}{d_{ij} n_{ij}} e^{-\sigma f_e} \tag{5}$$

where *C* and σ are constants (*C* = 1450 and σ = 2.8).⁶¹ The DFT results (at *T* = 25 °C) for c-CrN and c-AlN are listed in Table 5. The difference between the calculated value and the experimental value for the CrN phase in C1 sample is ~15%. This difference can be explained by considering the columnar orientation of the coating which is completely different from an ideal single crystal.

Temperature dependent variations in unit cell volume for each structure were obtained from QHA modelling up to 1500 °C (Fig. 8). Subsequently the bond length, d_{ij} , as a function of temperature was deduced from the cell volumes and by use of eqn (5) the hardness values for CrN (CrN-nm and CrN-m) and c-AlN phases. The results of hardness over a wide range of temperatures are shown in Fig. 9. Also plotted in Fig. 9 are the high temperature (25–500 °C) *in situ* nanoindentation data for commercially available $Al_{0.7}Cr_{0.3}N$ coatings studied by Fox-Rabinovich *et al.*⁶³ and Beake *et al.*⁶⁴ Room temperature nanoindentation results for the current projects' $Cr_{1-x}Al_xN$ coatings are also incorporated in Fig. 9.

Based on predicted material behaviours, hardness values for both phases reduced constantly at elevated temperatures. However, there are considerable differences in the reduction slopes depending on the phase type. Hardness values for CrNnm structure showed the highest resistance to temperature while AlN was more susceptible to temperature. Experimental hardness in Table 2 of $Cr_{1-x}Al_xN$ samples showed that nanocomposites of CrN in amorphous AlN matrix can effectively improve the resistance to grain boundary sliding65 which makes these samples harder than C1 sample. In addition, the deposition process and doping significantly distorts the lattice of CrN/ AlN mixture, as is evident in microstrain and residual stress plots, which can also contribute to improve the deformation resistance of doped samples. According to Šimůnek and Vackář61 the hardness of multicomponent compound systems is expected to be a geometrical sum of the values of hardness of all binary systems in the solid (given by their eqn (6)). Interestingly for the current multicomponent system (at 25 °C), consisting of CrN and AlN phases, $H = H_1 + H_2 = 22.5 + 21 = 43.5$ GPa (H_1 for

Table 5 DFT calculated hardness, at T = 25 °C, of bulk CrN, AlN based on d_{ij} . Current project and literature experimental data included for comparison

Phase	d_{ij}	e _i	e_j	f_e	$H_{\rm DFT}$	$H_{\rm EXPT}$
CrN-nm	2.028	3.125	4.904	0.0490	24.1	28 19–28.6 (ref. 36) 24–25 (ref. 62)
CrN-m c-AlN	2.071 2.034	2.793	4.545	0.056	$\begin{array}{c} 22.4\\ 21.1 \end{array}$	



Fig. 9 Simulated hardness values as a function of temperature for bulk CrN-m, CrN-nm and c-AlN. Nanoindentation hardness data, at high temperature for $Al_{0.7}Cr_{0.3}N$ thin film coatings and room temperature for current $Cr_{1-x}Al_xN$ coatings.

CrN-m and H_2 for AlN) compares remarkably well with the experimental value of 42 GPa for the C4 sample. Here the values correspond to a numerical sum of eqn (5) as opposed to a geometric sum. To show this point more clearly, the simulated data for CrN-m and AlN are added and the plot included in Fig. 9.

The experimental results of Fox-Rabinovich *et al.*,⁶³ are in some agreement with the simulated CrN-m, CrN-nm and AlN results for temperatures up to ~150 °C while they deviate significantly for higher temperatures up to ~500 °C, whilst Beake *et al.*⁶⁴ results are in agreement over the same temperature range. It should be noted that the materials in Beake's investigation has a high percentage Al content (~70%) whilst our materials have a maximum of ~21% Al content (C4). The hardness of experimental result from C4 and computational modelling from the summation of the CrN-m and AlN components are very similar at room temperature, see Fig. 9. If this is a real relationship, then the magnetic properties of the Cr are important for the formation of stable and hard $Cr_{1-x}Al_xN$ coatings.

Based on the results from both experiments and computational calculation, the temperature dependent hardness of the synthesized CrN/AlN thin film nanocomposites can be predicted in the temperature range from 25 °C to 1500 °C, refer to Fig. 9. Furthermore a numeric sum for the hardness of the 2 components, at T = 1500 °C give a predicted hardness $H \sim 41.5$ GPa. Only real experimental measurements at the 200–1500 °C range will show if the hardness of the CrN/AlN multicomponent thin film has the predicted direct sum values.

4. Conclusions

Synthesised $Cr_{1-x}Al_xN$ coatings, with a maximum Al content of ~21%, by unbalanced magnetic sputtering system showed improved microstructure and mechanical properties compared to undoped CrN coating. Analysis of *in situ* synchrotron X-ray diffraction at temperatures range of 200–700 °C established: (i) remarkable resistance to crystal growth and stress release up

to 700 °C for coatings with Al doping greater than \sim 14%. (ii) Aldopants improve oxidation resistance of the coatings by preventing the formation of Cr_2N at high temperatures.

Nano indentation results at ~ 25 °C show: (i) Al doping improved the coating hardness (*H*) and elastic modulus (*E*) from 29 to 42 GPa and 378 to 438 GPa respectively, and (ii) resistance to deformation increased as the Al content increased.

First-principles and quasi-harmonic approximations (QHA) studies for bulk magnetic and non-magnetic CrN (CrN-m and CrN-nm) and AlN over the temperature range of 0–1500 °C were undertaken to calculate elastic constants (C_{ij}) and hardness (H). Comparison of DFT-QHA models and the nanoindentation results indicate that hardness of $Cr_{1-x}Al_xN$ coating with ~21% Al content may be the sum of values from CrN-m and AlN phase. Our computational modelling predicts that the hardness (H) of $Cr_{1-x}Al_xN$ coating with the Al content of ~21% can reach to about 41.5 GPa at T = 1500 °C.

Acknowledgements

Ehsan Mohammadpour is highly grateful to the Murdoch University for the awarded scholarship (MIPS). The authors wish to acknowledge the support provided by the Australian Synchrotron with the beamline proposal AS141/PD/7582. The study has also been supported by grants of computing time from the National Computational Infrastructure (NCI), Canberra, Australia and the Pawsey Supercomputing Centre, Perth, Australia. We acknowledge the support from Dr M. Rowles (Curtain University) for invaluable discussions.

References

- 1 S. Zhang, Y. Li, T. Zhao and Q. Wang, *Sci. Rep.*, 2014, 4(5241), 1–7.
- 2 J. Y. Cheong, X. Z. Ding, B. K. Tay and X. T. Zeng, *Key Eng. Mater.*, 2010, 447–448, 725–729.
- 3 S. R. Pulugurtha, D. G. Bhat, M. H. Gordon, J. Shultz, M. Staia, S. V. Joshi and S. Govindarajan, *Surf. Coatings Technol.*, 2007, 202, 1160–1166.
- 4 J. Musil, Surf. Coatings Technol., 2012, 207, 50-65.
- 5 X. Guan, Y. Wang, G. Zhang, J. Xin, L. Wang and Q. Xue, *RSC Adv.*, 2016, **6**, 87003–87012.
- 6 P. Z. Shi, J. Wang, C. X. Tian, Z. G. Li, G. D. Zhang, D. J. Fu and B. Yang, *Surf. Coatings Technol.*, 2013, 228, S534–S537.
- 7 N. Bagcivan, K. Bobzin and S. Theiß, *Thin Solid Films*, 2013, **528**, 180–186.
- 8 H. N. Shah, R. Jayaganthan and A. C. Pandey, *Mater. Des.*, 2011, **32**, 2628–2634.
- 9 Z. Wu, F. Zhou, Q. Ma, Q. Wang, Z. Zhou and L. Kwok-Yan Li, *RSC Adv.*, 2016, **6**, 76724–76735.
- 10 L. Lu, Q. M. Wang, B. Z. Chen, Y. C. Ao, D. H. Yu, C. Y. Wang, S. H. Wu and K. H. Kim, *Trans. Nonferrous Met. Soc. China*, 2014, 24, 1800–1806.
- P. Panjan, A. Drnovšek, J. Kovač, P. Gselman, T. Bončina, S. Paskvale, M. Čekada, D. Kek Merl and M. Panjan, *Thin Solid Films*, 2015, **591**, 323–329.

- 12 E. Bemporad, C. Pecchio, S. De Rossi and F. Carassiti, *Surf. Coatings Technol.*, 2004, **188–189**, 319–330.
- 13 S. Tan, X. Zhang, X. Wu, F. Fang and J. Jiang, *Appl. Surf. Sci.*, 2011, **257**, 5595–5600.
- 14 Q. Luo, S. C. Wang, Z. Zhou and L. Chen, *J. Mater. Chem.*, 2011, **21**, 9746–9756.
- 15 J. Xu, L. Liu, P. Munroe, Z. Xie and Z. Jiang, *J. Mater. Chem. A*, 2013, **1**, 10281–10291.
- 16 J. Xu, D. Lai, Z. Xie, P. Munroe and Z.-T. Jiang, J. Mater. Chem., 2012, 22, 2596–2606.
- 17 Y. Long, J. Zeng, D. Yu and S. Wu, Ceram. Int., 2014, 40, 9889–9894.
- 18 Z. Li, P. Munroe, Z. Jiang, X. Zhao, J. Xu, Z. Zhou, J. Jiang,
 F. Fang and Z. Xie, *Acta Mater.*, 2012, 60, 5735–5744.
- 19 P. C. Wo, P. R. Munroe, Z. Li, Z.-T. Jiang, Z. H. Xie, Z. F. Zhou and K. Y. Li, *Mater. Sci. Eng. A*, 2012, **534**, 297–308.
- 20 E. Mohammadpour, Z.-T. Jiang, M. Altarawneh, Z. Xie, Z. Zhou, N. Mondinos, J. Kimpton and B. Z. Dlugogorski, *Thin Solid Films*, 2015, **599**, 98–103.
- 21 A. A. Coelho, J. Appl. Crystallogr., 2003, 36, 86-95.
- 22 R. W. Cheary, A. a. Coelho and J. P. Cline, *J. Res. Natl. Inst. Stand. Technol.*, 2004, **109**, 1–25.
- 23 S. Gražulis, A. Daškevič, A. Merkys, D. Chateigner, L. Lutterotti, M. Quirós, N. R. Serebryanaya, P. Moeck, R. T. Downs and A. Le Bail, *Nucleic Acids Res.*, 2012, 40, 420–427.
- 24 H. Toraya, T. C. Huang and Y. Wu, *J. Appl. Crystallogr.*, 1993, 26, 774–777.
- 25 O. Masson, R. Guinebretière and A. Dauger, J. Appl. Crystallogr., 1996, 29, 540–546.
- 26 M. R. Rowles and I. C. Madsen, *J. Appl. Crystallogr.*, 2010, 43, 632–634.
- 27 J. P. Perdew and Y. Wang, *Phys. Rev. B*, 1992, 45, 13244-13249.
- 28 G. Kresse and J. Hafner, Phys. Rev. B, 1993, 47, 558-561.
- 29 A. Otero-De-La-Roza and V. Luaña, *Comput. Phys. Commun.*, 2011, **182**, 1708–1720.
- 30 P. Vinett, J. H. Rose, J. Ferrante and J. R. Smith, *Phys. Rev. B*, 1996, **53**, 5252–5258.
- 31 Pierre Villars, PAULING FILE, in Inorganic Solid Phases, SpringerMaterials (online database), Springer, Heidelberg Springer Materials Fe3W3C (W3Fe3C) Crystal Structure http://materials.springer.com/isp/crystallographic/docs/sd_ 1707470.
- 32 M. Mahbubur Rahman, A. Duan, Z.-T. Jiang, Z. Xie, A. Wu, A. Amri, B. Cowie and C.-Y. Yin, *J. Alloys Compd.*, 2013, **578**, 362–368.
- 33 S. Agouram, F. Bodart and G. Terwagne, Surf. Coatings Technol., 2004, 180–181, 164–168.
- 34 H. C. Barshilia, N. Selvakumar, B. Deepthi and K. S. Rajam, Surf. Coat. Technol., 2006, 201, 2193–2201.
- 35 T. Elangovan, P. Kuppusami, R. Thirumurugesan,
 V. Ganesan, E. Mohandas and D. Mangalaraj, *Mater. Sci. Eng. B*, 2010, 167, 17–25.
- 36 L. Cunha, M. Andritschky, K. Pischow and Z. Wang, *Thin Solid Films*, 1999, 355–356, 465–471.

- 37 D. Rafaja, C. Wüstefeld, C. Baehtz, V. Klemm, M. Dopita, M. Motylenko, C. Michotte and M. Kathrein, *Metall. Mater. Trans. A*, 2011, 42, 559–569.
- 38 A. Filippetti, W. E. Pickett and B. M. Klein, *Phys. Rev. B*, 1999, 59, 7043–7050.
- 39 Z. G. Zhang, O. Rapaud, N. Bonasso, D. Mercs, C. Dong and C. Coddet, *Vacuum*, 2008, 82, 501–509.
- 40 S. Yang, E. Wiemann and D. G. Teer, *Surf. Coatings Technol.*, 2004, **188–189**, 662–668.
- 41 F. Sanchette, C. Ducros, T. Schmitt, P. Steyer and a. Billard, *Surf. Coatings Technol.*, 2011, **205**, 5444–5453.
- 42 M. F. Yan and H. T. Chen, *Comput. Mater. Sci.*, 2014, 88, 81–85.
- 43 I. Scanavino and M. Prencipe, *RSC Adv.*, 2013, 3, 17813–17821.
- 44 B. Alling, T. Marten and I. A. Abrikosov, *Phys. Rev. B*, 2010, 82, 1–9.
- 45 A. J. Wang, S. L. Shang, Y. Du, Y. Kong, L. J. Zhang, L. Chen, D. D. Zhao and Z. K. Liu, *Comput. Mater. Sci.*, 2010, 48, 705– 709.
- 46 L. Zhou, F. Körmann, D. Holec, M. Bartosik, B. Grabowski, J. Neugebauer and P. H. Mayrhofer, *Phys. Rev. B*, 2014, 90, 184102.
- 47 Y. Makino and K. Nogi, *Surf. Coatings Technol.*, 1998, 98, 1008–1012.
- 48 S.-K. Tien and J.-G. Duh, *Thin Solid Films*, 2006, **494**, 173–178.
- 49 B. Alling, T. Marten, I. a. Abrikosov and A. Karimi, *J. Appl. Phys.*, 2007, **102**, 1–8.
- 50 A. Sugishima, H. Kajioka and Y. Makino, *Surf. Coat. Technol.*, 1997, **97**, 590–594.
- 51 F. L. E. Martinez, R. Sanjines and O. Banakh, *Thin Solid Films*, 2004, **446**, 18–22.
- 52 M. Tkadletz, N. Schalk, R. Daniel, J. Keckes, C. Czettl and C. Mitterer, *Surf. Coatings Technol.*, 2016, **285**, 31–46.
- 53 B. Kocak, Y. O. Ciftci, K. Colakoglu and E. Deligoz, *Phys. B Condens. Matter*, 2012, **407**, 316–323.
- 54 M. G. Brik and C.-G. Ma, *Comput. Mater. Sci.*, 2012, 51, 380-388.
- 55 M. E. Eberhart and T. E. Jones, *Phys. Rev. B Condens. Matter Mater. Phys.*, 2012, **86**, 1–7.
- 56 D. Wu, H. Wang, L. Wei, R. Pan and B. Tang, *J. Magnes. Alloy.*, 2014, **2**, 165–174.
- 57 A. Lesar, M. Hodos and I. Introduction, *Chem. Phys.*, 1998, 109, 9410–9416.
- 58 Y. Liang, X. Yuan and W. Zhang, Solid State Commun., 2010, 150, 2045–2048.
- 59 F. Rivadulla, M. Bañobre-López, C. X. Quintela, A. Piñeiro, V. Pardo, D. Baldomir, M. A. López-Quintela, J. Rivas, C. a. Ramos, H. Salva, J.-S. Zhou and J. B. Goodenough, *Nat. Mater.*, 2009, 8, 947–951.
- 60 Q. Xia, H. Xia and A. L. Ruoff, *J. Appl. Phys.*, 1993, 73, 8198-8200.
- 61 A. Šimůnek and J. Vackář, Phys. Rev. Lett., 2006, 96, 1-4.
- 62 H. Ichimura and I. Ando, *Surf. Coatings Technol.*, 2001, 145, 88–93.

Paper

- 63 G. S. Fox-Rabinovich, B. D. Beake, J. L. Endrino, S. C. Veldhuis, R. Parkinson, L. S. Shuster and M. S. Migranov, *Surf. Coatings Technol.*, 2006, 200, 5738–5742.
- 64 B. D. Beake, G. S. Fox-Rabinovich, S. C. Veldhuis and S. R. Goodes, *Surf. Coatings Technol.*, 2009, **203**, 1919–1925.
- 65 S. Veprek and A. S. Argon, *Surf. Coatings Technol.*, 2001, **146**–**147**, 175–182.