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# A probe of steric ligand substituent effects on the spin crossover of Fe(II) complexes†

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Identifying and quantifying the individual factors affecting the temperature and properties of the spin crossover in transition metal complexes is a challenging task, because many variables are involved. While the most decisive factor is the crystal field imparted by ligands around the active metal center, some less common actors are intramolecular steric repulsions or non-covalent interactions. A series of three Fe(II) complexes of 1,3bpp derivatives of 2-(pyrazol-1-yl)-6-(1H-pyrazol-3-yl)pyridine have been prepared and characterized crystallographically to probe these effects: [Fe(1,3bpp)<sub>2</sub>](ClO<sub>4</sub>)<sub>2</sub> (**1**), [Fe(met1,3bpp)<sub>2</sub>](ClO<sub>4</sub>)<sub>2</sub> (**2**) and [Fe(dimet1,3bpp)<sub>2</sub>](ClO<sub>4</sub>)<sub>2</sub> (**3**). The ligands exhibit none, one or two methyl substituents on the pyrazol-1-yl heterocycle. These groups exert a dramatic effect on the SCO temperature in the solid state, and, most significantly, in solution (with  $T_{SCO}$  (**3**) >  $T_{SCO}$  (**1**) >  $T_{SCO}$  (**2**)). Extensive DFT calculations have unveiled the origin of these effects which lie in the intramolecular non-covalent or steric interactions rather than resulting from crystal field effects.

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## Introduction

Octahedral transition metal complexes in the configurations d<sup>4</sup> to d<sup>7</sup> may be in two different spin states, depending on the energy of the t<sub>2g</sub> vs. e<sub>g</sub> orbital splitting in relation to the energy necessary for pairing two electrons within one d orbital. If both energies are comparable, the system is likely to exhibit spin crossover (SCO) phenomena following small external perturbations.<sup>1–3</sup> The ensuing spin transitions cause important changes not only to the magnetic properties, but also to the structure and to a number of physical properties.<sup>4,5</sup> For this reason, this phenomenon is considered as a promising entry into molecule-based switching materials for potential applications in nanotechnology.<sup>6,7</sup> Among the suitable metals, Fe(II) is especially interesting because the transition toggles the complex between a diamagnetic (*S* = 0) and a paramagnetic (*S* = 2) state, causes dramatic colour changes and leads to Fe-to-ligand bond distance variations of 10% or larger.<sup>8,9</sup> For a given

metal ion, the temperature and dynamics of the SCO are affected by many factors, most often superimposed, dependent on the ligands,<sup>10</sup> the crystal system,<sup>11</sup> intermolecular interactions<sup>5,12,13</sup> or secondary bonding interactions.<sup>14–16</sup> Many synthetic chemists are dedicating efforts to designing and creating complexes with the challenging goal of unveiling the specific influence of each factor, if possible with independence of any other effect. The many ways in which the specific nature of the ligands affects the thermodynamics of the SCO have been recently reviewed.<sup>10</sup> This is best investigated, whenever possible, in solution rather than in the solid state, since the latter situation involves often the existence of solvatomorphs<sup>17–26</sup> and sometimes polymorphs,<sup>11</sup> usually affecting dramatically the SCO properties. Ligand field effects, as conveyed through the incorporation of ligand substituents, have been analysed extensively, leading to results that sometimes may appear conflicting.<sup>27–30</sup> Some of the earlier reports already point out to the difference between σ-donating and π-accepting properties of the ligands to explain the complexity of the substituent effects on the SCO properties.<sup>30,31</sup> Very recently, the study in solution of an extensive family of 1bpp/Fe(II) complexes (1bpp = 2,6-bis-(pyrazol-1-yl)-pyridine) featuring a variety of substituents on the central pyridine or on the pyrazole rings demonstrated the coexistence and the opposing effect of both, σ and π bonding properties, as well as their differing relative importance depending on the position of the substituent.<sup>32</sup> Another way for substituents to influence the SCO temperature is through variations in intramolecular attractive or repulsive interactions resulting from the structural changes accompanying the spin transition. In the

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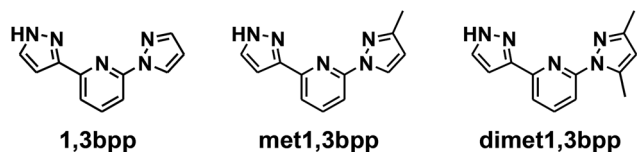
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**Scheme 1** Molecular structure of ligands 1,3bpp, met1,3bpp and dimet1,3bpp.

mentioned study, this is not the case since the substituents investigated are located at distal positions, thus not taking part in such interactions. Nonetheless, several ligand families have been shown experimentally to stabilize or even trap the high spin (HS) state of Fe(II) complexes by introducing a steric constraint to the shrinking bond to occur upon SCO to the low spin (LS) state.<sup>33–37</sup> This effect has also been rationalized through DFT calculations.<sup>37,38</sup> In much rarer cases, intramolecular interactions block or impede the transition to the HS state. This occurs in an Fe(II) scorpionate complex exhibiting a bulky group that exerts steric intra-ligand repulsion with the lengthening of the Fe–N bond distances accompanying the SCO.<sup>39</sup> It has also been suggested to happen with the cation  $[\text{Fe}(\text{Me}_4\text{-1bpp})_2]^{2+}$  ( $\text{Me}_4\text{-1bpp}$  = 2,6-bis-(3,5-dimethylpyrazol-1-yl)-pyridine), as a result of the interaction of the methyl groups in position 5 of pyrazolyl with the central pyridine in the 1bpp core.<sup>10</sup> The same effect was invoked to explain the decrease in the SCO temperature in an Fe(II) complex of an indazolylpyridine derivative.<sup>40</sup> We have now designed a ligand system to probe these two not ligand field related opposite effects within an analogous series of Fe(II) complexes. Thus, three ligands have been prepared (Scheme 1); 2-(pyrazol-1-yl)-6-(1H-pyrazol-3-yl)pyridine (1,3bpp),<sup>41</sup> 2-(3-methylpyrazol-1-yl)-6-(1H-pyrazol-3-yl)pyridine (met1,3bpp) and 2-(3,5-dimethylpyrazol-1-yl)-6-(1H-pyrazol-3-yl)pyridine (dimet1,3bpp), showing zero, one and two methyl substituents on the pyrazol-1-yl rings. The corresponding homoleptic Fe(II) complexes  $[\text{Fe}(\text{1,3bpp})_2](\text{ClO}_4)_2$  (previously published,<sup>41</sup> **1**),  $[\text{Fe}(\text{met1,3bpp})_2](\text{ClO}_4)_2$  (**2**) and  $[\text{Fe}(\text{dimet1,3bpp})_2](\text{ClO}_4)_2$  (**3**) were prepared and their magnetic properties determined in the solid state and in solution. The results show that the methyl in position 3 favours the HS state compared to the unsubstituted system, and that the same substituent on position 5 favours the LS state, with a stronger incidence. These results have been fully rationalized and quantified with the help of DFT calculations in terms of inter- and intra-ligand interaction effects. The computational method has been used to investigate the potential analogue with only one methyl on position 5 (“4”), not accessible experimentally with our synthetic procedure. This has allowed confirming the observed trends and their interpretation.

## Results and discussion

### Synthesis

The ligands met1,3bpp and dimet1,3bpp were prepared in three steps using a procedure analogous to the synthesis of 1,3bpp.<sup>41</sup> Thus, the appropriate substituted pyrazole ring was

first coupled through one N atom to 2-acetyl-6-bromopyridine. The pyrazole with only one substituent couples through the less crowded nitrogen atom, leading ultimately to the ligand met1,3bpp with the substituent on position 3. This in fact prevents using this procedure to prepare met1,3bpp with the methyl on position 5. The product is then functionalized at the carbonyl end with *N,N*-dimethylformamide-dimethyl acetal into the corresponding 3-(dimethylamino)prop-2-en-1-one moiety which is readily converted by ring closure with hydrazine into the pyrazole-3-yl substituent of the central pyridine group, common to all the 1,3bpp ligands. The complex  $[\text{Fe}(\text{1,3bpp})_2](\text{ClO}_4)_2$  (**1**) was prepared as published by our group.<sup>41</sup> It could be obtained following two distinct procedures as two different polymorphs, **1a** and **1b** (see below). Complexes  $[\text{Fe}(\text{met1,3bpp})_2](\text{ClO}_4)_2$  (**2**) and  $[\text{Fe}(\text{dimet1,3bpp})_2](\text{ClO}_4)_2$  (**3**) were obtained by direct reaction of the hydrated  $\text{Fe}(\text{ClO}_4)_2$  salt with the met1,3bpp and dimet1,3bpp ligands, respectively, in the presence of a catalytic amount of ascorbic acid to prevent the oxidation of Fe(II) to Fe(III). Compound **2** was obtained from a reaction in absolute ethanol that produced a yellow solution, using hexane as the crystallization medium. The pair of solvents used to obtain **3** as red crystals are acetone and diethyl ether. The reaction with met1,3bpp using acetone/toluene produces a solvatomorph,  $[\text{Fe}(\text{met1,3bpp})_2](\text{ClO}_4)_2 \cdot \text{H}_2\text{O}$  (**2b**), incorporating one molecule of water per complex unit (ESI†).

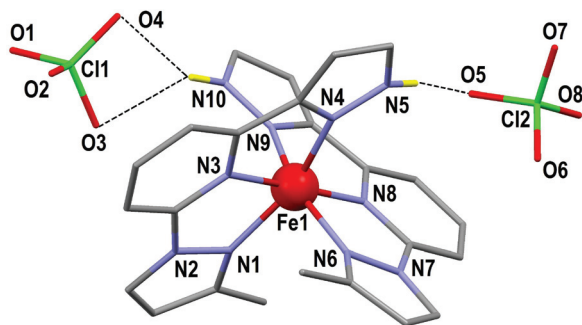
### Description of the structures

The structure of complex **1** has been already described in a previous publication.<sup>41</sup> This compound can be obtained following two different procedures as two polymorphs, **1a** and **1b**, showing two different organizations of the Fe(II) complex cations closely related to these observed for compounds **2** and **3**, respectively (see below). The structure of the solvatomorph **2b** is briefly described in the ESI (Table S1 and Fig. S1 and S2†).

$[\text{Fe}(\text{met1,3bpp})_2](\text{ClO}_4)_2$  (**2**). The structure of **2** was determined at 100 K, on crystals that had turned red from their original yellow color at room temperature. Their solvent free lattice is found in the monoclinic space group *C2/c*. The asymmetric unit consists of one formula unit, with eight such moieties present in the unit cell. The complex cation features a distorted octahedral Fe(II) center coordinated to two met1,3bpp tris-imine ligands lying approximately perpendicular to each other (Fig. 1). As a result of the asymmetric character of the ligands, this complex is chiral, both enantiomers being present in the lattice, which is racemic. The average of the Fe–N bond distances is 1.96(4) Å, corroborating the LS state of the compound at this temperature. The spin state is also evident from the distortion parameters  $\Sigma$  and  $\Theta$ ,<sup>42–44</sup> which here amount to 93.2° and 367.8°, respectively, within the region expected for LS compounds.<sup>9</sup>

As expected in solvent free structures of Fe/bpp complexes with at least one pyrazol-3-yl ring per ligand, the  $\text{ClO}_4^-$  anions establish hydrogen bonding interactions with their N–H groups (Fig. 1). The complexes organize in the lattice as one of

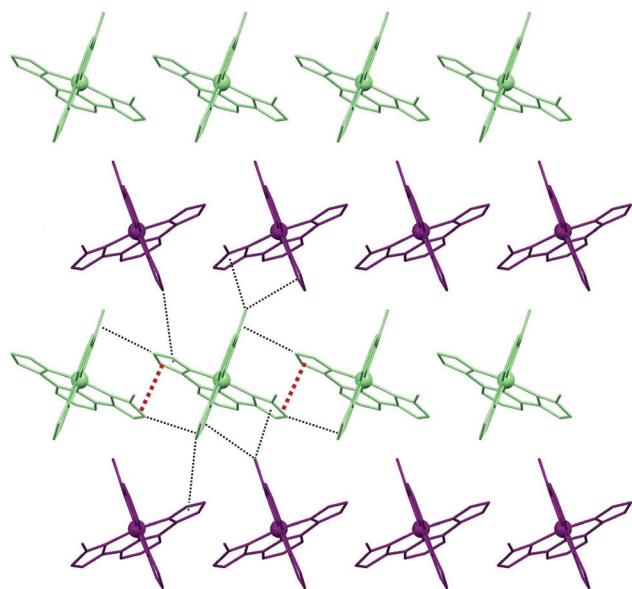




**Fig. 1** Molecular representation of  $[\text{Fe}(\text{met1,3bpp})_2](\text{ClO}_4)_2$  (**2**) at 100 K with heteroatoms labelled. Only H atoms of N–H groups shown (in yellow). Dashed lines are H-bonds.

the polymorphs previously reported of compound **1** (**1a**).<sup>41</sup> Thus, they are disposed as sheets containing arrays of  $[\text{Fe}(\text{met1,3bpp})_2]^{2+}$  cations. Within the sheets, each complex interacts with two neighbours *via* two  $\pi\cdots\pi$  and six C–H $\cdots\pi$  interactions. In between the sheets, each cation establishes a total of six weaker C–H $\cdots\pi$  contacts with two nearby congeners (Fig. 2). The parallel arrays within these layers alternate complexes of opposed coordination chirality and also two different orientations. The angle between complexes in these two orientations (measured using idealized planes of two equivalent ligands) is 41.40°. There are two types of (very similar) inter-layer separations (Fig. S3†), 9.640 Å and 9.887 Å.

The structure of **2** was also determined at 300 K, on crystals that had turned pale yellow, as a result of a LS to HS conversion. Following the SCO and the thermal expansion, the unit cell experiences an isotropic growth (Table S2†), with a volume



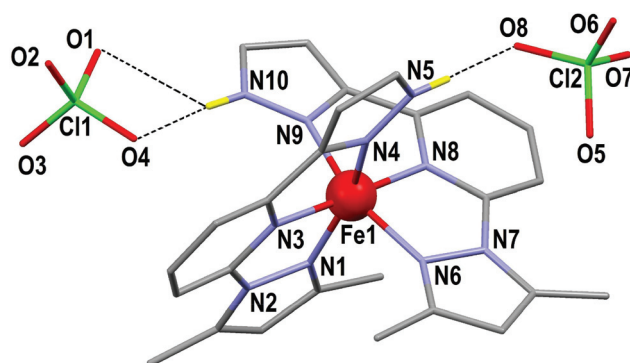
**Fig. 2** Sheet organization of the cations in **2**, emphasizing their two different orientations and the  $\pi\cdots\pi$  and C–H $\cdots\pi$  interactions formed by each complex with its immediate neighbours within the sheet.

expansion of 5%, translating into an increase of the separation between layers of complex cations to 9.767 and 9.894 Å, respectively. Furthermore, the average Fe–N at this temperature is 2.16(2) Å, while the distortion parameters are  $\Sigma = 147.5^\circ$  and  $\theta = 378.2^\circ$ , respectively, confirming that at 300 K, the compound is in its HS configuration.

$[\text{Fe}(\text{dimet1,3bpp})_2](\text{ClO}_4)_2$  (**3**). The molecular structure of **3** was determined at 100 and 298 K. At both temperatures, the lattice exhibits the monoclinic  $P2_1/n$  space group, the asymmetric unit coinciding with the empirical formula and the unit cell enclosing four such moieties. The complex cation  $[\text{Fe}(\text{dimet1,3bpp})_2]^{2+}$  is analogous to that of **2**, now with the ligand dimet1,3bpp (Fig. 3). The average of the Fe–N distances (1.95(4) Å at both temperatures) and the  $\Sigma/\theta$  distortion parameters ( $88.7/367.2^\circ$  and  $90.0/367.1^\circ$  respectively) show that the Fe(II) centres are in the LS state at 100 and 296 K. Attempts to obtain the structure of **3** in the HS (>400 K, see below) were unsuccessful because of crystal damage. In the LS, the lattice is related to the layered organization of **2**. One main difference is that within layers of complex cations, the chirality and orientations of the complexes are identical, whereas the handedness of the cations alternate in moving through adjacent layers. Within each sheet, each complex has four first neighbours, establishing four  $\pi\cdots\pi$  and eight C–H $\cdots\pi$  interactions (Fig. 4). Interestingly, this particular arrangement is the same as that of the other polymorph of compound **1** (**1b**) also previously characterized.<sup>41</sup>

The layers feature separations (Fig. S4†) of 9.420/9.499 and 9.636/9.598 Å at both temperatures, respectively. In going from 100 K to 296 K, the cell dimensions experience an isotropic expansion purely of thermal origin, with a volume increase of 4%.

The homogeneity of the phases described above was established by means of powder X-ray diffraction (PXRD) methods. The purity and homogeneity of compound **1**, in its two polymorphic forms, respectively, had been established previously using this method.<sup>41</sup> For the cases of **2** and **3**, PXRD experiments were also conducted. The results prove that in both



**Fig. 3** Molecular representation of  $[\text{Fe}(\text{dimet1,3bpp})_2](\text{ClO}_4)_2$  (**3**) at 100 K with heteroatoms labelled. Only H atoms of N–H groups shown (in yellow). Dashed lines are H-bonds.





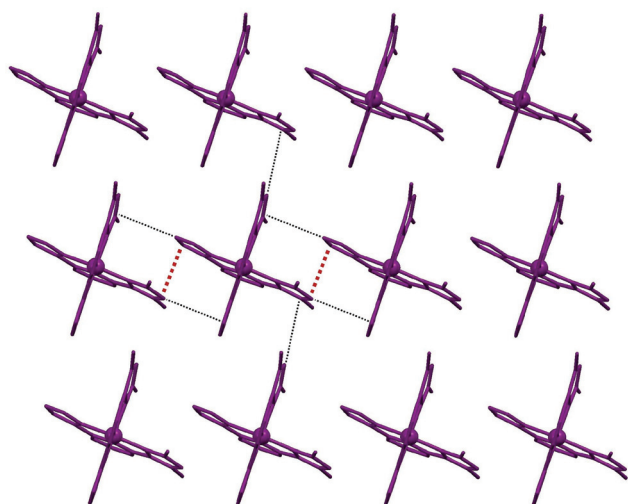


Fig. 4 Sheet organization of the cations in **3**, showing their sole orientation and the  $\pi\cdots\pi$  and C–H $\cdots\pi$  interactions formed by each complex with its immediate neighbours within the sheet.

cases the bulk material corresponds to the compound unveiled by single crystal X-ray diffraction (Fig. S5†).

### Solid state magnetic and thermal properties

The influence of the methyl substituents on the SCO of the 1,3bpp/Fe(II) complexes was first assessed through bulk magnetic susceptibility measurements. Data from polycrystalline samples of **2** and **3** were collected between 5 and 400 K in the warming and cooling modes under a constant magnetic field (see the ESI†) and were compared to these from polymorphs **1a** and **1b**. All the results are displayed in Fig. 5 in the form of  $\chi_M T$  vs.  $T$  plots ( $\chi_M$  is the molar paramagnetic susceptibility). At low temperature, all complexes are essentially diamagnetic, with  $\chi_M T$  values at 100 K ranging 0.06 to 0.17 cm<sup>3</sup> K mol<sup>−1</sup>. In all cases, an abrupt increase of  $\chi_M T$  occurs upon warming up to nearly constant values of 3.01 (**2**), 3.26 (**1a**) and 3.4 cm<sup>3</sup> K mol<sup>−1</sup> (**1b**), while for **3** the product (2.64 cm<sup>3</sup> K mol<sup>−1</sup>) was still increasing at 400 K, the maximum temperature reached by the magnetometer. The high temperature values show that the Fe(II) centers reach the HS ( $S = 2$ ), consistent with the data from SCXRD, and with the occurrence of SCO. The SCO profiles in the cooling mode are quasi-superimposable to these in the warming mode, indicative of the absence of hysteresis, even though the transitions of **2** and **1b** are clearly more abrupt than those of **1a** and **3** (see below). The various systems, in addition, exhibit dramatically different transition temperatures, with  $T_{1/2}$  values of 183 (**2**), 278 (**1a**), 314 (**1b**) and 378 K (**3**). These observations are fully consistent with the temperature-dependence of the molar heat capacity. Indeed, anomalies are observed at temperatures coinciding with those of the SCO processes, which are very sharp, sharp and relatively broad, respectively, for **2**, **1b** and **1a/3** (Fig. 5, Fig. S6† and Table 1). The excess enthalpy and entropy associated with SCO,  $\Delta H_{SCO}$  and  $\Delta S_{SCO}$  (as derived from the excess heat

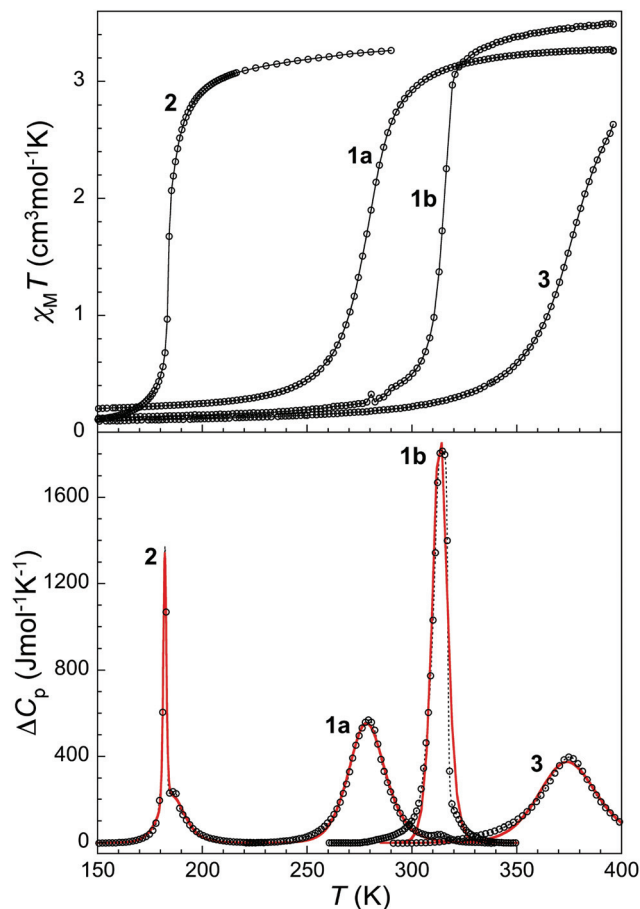


Fig. 5 Thermal SCO properties of compounds **1a**, **1b**, **2** and **3**:  $\chi_M T$  vs.  $T$  plots (top) and molar heat capacity at constant pressure as derived from DSC measurements (bottom). All data shown correspond to the warming mode.

Table 1 Thermodynamic parameters of compounds **1a**, **1b**, **2** and **3**

	<b>1a</b>	<b>1b</b>	<b>2</b>	<b>3</b>
$\Delta H_{SCO}$ (kJ mol <sup>−1</sup> )	13.57	17.74	5.87	14.11
$\Delta S_{SCO}$ (J mol <sup>−1</sup> K <sup>−1</sup> )	48.7	56.8	31.7	38.0
$n^a$	8.5(1)	20.0(2)	128.7(3)	9.0(1)
$T_{SCO}$ (K) <sup>a</sup>	278.8(3)	313.4(3)	182.1(1)	375.0(3)
$T_{1/2}$ (K) <sup>b</sup>	278	314	184	378
$T_{1/2}$ (K) <sup>c</sup>	262	262	232	281

<sup>a</sup> Solid-state, from fit of  $\Delta C_p$  vs.  $T$  to the domain model (see the ESI).

<sup>b</sup> Solid-state, from  $\chi_M T$  vs.  $T$ . <sup>c</sup> Solution, from NMR (in solution, **1a** and **1b** become, to a very good approximation, the same system).

capacity  $\Delta C_p$ ; Table 1 and the ESI†) give a qualitative measure of the cooperativeness of a SCO process. Here, the excess entropies turn out to be much larger than the electronic component  $R \ln 5$ , which is indicative of significant coupling of the SCO with lattice phonons. These thermodynamic parameters are however affected by the temperature of the SCO processes, which varies dramatically in the present compounds. Therefore, a more quantitative measure of the cooperativity



has been obtained by fitting the experimental  $\Delta C_p$  vs.  $T$  data to the so-called domain model (see the ESI† for details).<sup>45,46</sup> The derived number of interacting molecules per domain,  $n$ , is similar for **1a** and **3**, and is 8.5 and 9.0, respectively, while for **1b** it is about double, characteristic of the medium to high cooperative character of the SCO (values of  $n$  close to unity are expected for gradual SCO while values above 20 are found for strongly cooperative systems).<sup>46–48</sup> On the contrary, the very large  $n$  obtained for **2** ( $n = 128.7$ ) ranges among the largest reported,<sup>46</sup> thus depicting a highly cooperative system. This is likely due to a strong coupling between the SCO and the induced structural modifications, in agreement with a sharp variation of cell parameters at the SCO. In fact, the anomaly in the  $C_p$  vs.  $T$  curve exhibits clearly two components (Fig. S7†), an extremely sharp peak on top of a broader feature, most likely reflecting both processes.

While the substituents on the 1,3bpp ligand-core certainly have an impact on  $T_{1/2}$ , the marked disparity between polymorphs **1a** and **1b** (of about 40 K) demonstrates that the crystal packing alone is very influential. Thus, while solid-state measurements are essential to investigate the latter effects, especially on the cooperativity, this technique is not appropriate to quantify with independence the influence of the nature and location of the methyl substituents on the SCO temperature.

### <sup>1</sup>H-NMR spectroscopy

To identify the influence of the ligand on the temperature of the spin transition, excluding solid state effects, the best choice is the use of a solution methodology, such as NMR. The variable temperature paramagnetic susceptibility of a soluble substance may be calculated by this technique, using the Evans method.<sup>49–51</sup> The latter is based on the relationship between  $\chi_M$  of a paramagnetic substance and the effect that it causes to the chemical shift of any species in solution. This effect is called the paramagnetic shift and it may be measured directly using a diamagnetic reference (such as TMS) by collecting the NMR spectrum with a coaxial tube made of two separate compartments that contain both, the solvent and the reference. One of the compartments must also contain the magnetic species in solution. The paramagnetic shift can thus be obtained directly from the composite spectrum, extracting the difference between the signals of the reference in both compartments. The room temperature <sup>1</sup>H NMR spectra of complexes **1**, **2** and **3** (Fig. S8†) in CD<sub>3</sub>OD show sets of paramagnetically shifted and broadened peaks lacking hyperfine splitting, consistent with the symmetry, number of protons and integration expected in each case (considering that the signals from the N–H groups are broadened beyond detection because of their proximity to Fe(II) and their ability to exchange). In these spectra, the remainder of the peaks is due to TMS, residual MeOH, H<sub>2</sub>O and other solvents, in addition to the free ligand. In view of their stability in solution, the spectra of the three compounds at various temperatures were recorded between 193 and 300 K. The variations of the TMS paramagnetic shift (Table S6†) provide the temperature depen-

dence of  $\chi_M$  in solution for **1**, **2** and **3**. The corresponding plot of  $\chi_M T$  vs.  $T$  (Fig. 6) reveals that the three complexes exhibit gradual SCO processes, with approximately  $T_{1/2}$  values of (in K) of 232 (**2**), 262 (**1**) and 281 (**3**). These values are correlated with the temperatures obtained from bulk measurements, while the differences are ascribed to the contribution from solid-state effects. These can be very important, as illustrated by the almost 100 K difference in the SCO temperature shown by **3**.

From the solution experiments, it is deduced that a methyl group at position 3 of the pyrazolyl ring stabilizes the HS state with respect to the LS state (thus reducing the SCO temperature). The presence of methyl groups on positions 3 and 5 leads to the opposite result, an increase of the SCO temperature. Therefore, a substituent on position 5 not only opposes the influence of the 3-methyl group but also has a dominant impact. This could be corroborated if the compound with only one methyl group on position 5 was accessible experimentally; however, it is not the case (see above). In any case, the above data demonstrate that the choice of the specific 1,3bpp derivative opens a means of tuning the SCO temperature. In order to rationalize the causes of the observed effects, in addition to investigating what would be the net impact of a hypothetical 5-methyl derivative, DFT calculations were performed.

### DFT calculations

The relative stability of the HS and LS state formed in the gas phase was calculated by means of DFT+U+D2 for complexes **2** (3-methyl derivative), **3** (3,5-dimethyl derivative) and for a hypothetical 5-methyl derivative (“**4**”). The results were compared with those previously published for the bare complex cation of **1**.<sup>37</sup> The energies of the optimized structures furnished the electronic contributions to the enthalpy difference existing between both states ( $\Delta H_{\text{elec}}$ ). The computed values

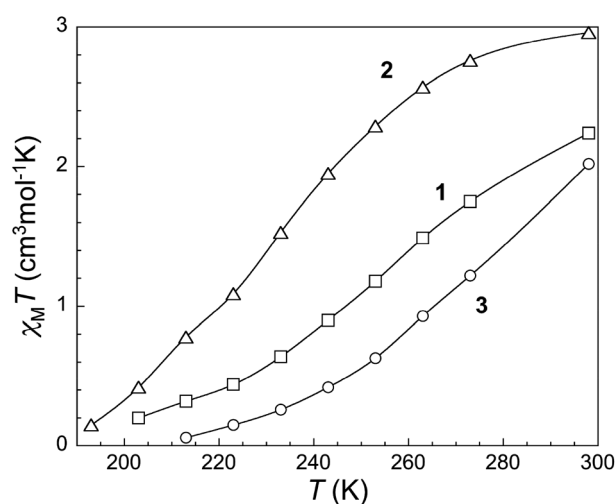


Fig. 6  $\chi_M T$  vs.  $T$  plots as derived from solution <sup>1</sup>H NMR spectra in d<sub>4</sub>-MeOH of **1**, **2** and **3** by employing the Evans method (see the text for details). Calculations have been performed from the temperature where the paramagnetic shift was large enough to be measured.



**Table 2** Average energy of the two sets of 3d-orbitals for compounds **1–3** and “**4**”, energy difference between them ( $\Delta E$ , in eV), and electronic enthalpy ( $\Delta H$ , in kJ mol<sup>-1</sup>)

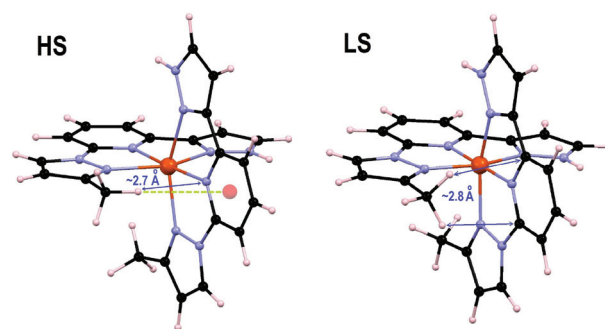
	1	2	3	“4”
$t_{2g}$	-1.48	-1.45	-1.41	-1.48
$e_g$	2.26	2.28	2.31	2.28
$\Delta E$	-3.74	-3.73	-3.72	-3.76
$\Delta H$	11.9	9.9	16.4	19.6

(Table 2) are fully consistent with the experimental results. Thus, compared with  $\Delta H_{\text{elec}}$  for complex **1**, with a bare 1,3bpp ligand (of 11.9 kJ mol<sup>-1</sup>), the effect of adding methyl groups depends drastically on the position of this substituent. When it is located at position 3 (complex **2**),  $\Delta H_{\text{elec}}$  diminishes by *ca.* 2.0 kJ mol<sup>-1</sup> ( $\Delta H_{\text{elec}} = 9.9$  kJ mol<sup>-1</sup>) whereas adding it at position 5 (“**4**”) increases its value notably by 7.7 kJ mol<sup>-1</sup>, up to  $\Delta H_{\text{elec}} = 19.6$  kJ mol<sup>-1</sup>. The addition of two methyl substituents, one at each position, results in a value of  $\Delta H_{\text{elec}} = 16.4$  kJ mol<sup>-1</sup>. This value is 4.5 kJ mol<sup>-1</sup> larger than that of **1**, which is close to the combination of both individual effects, calculated separately for **2** and “**4**” (7.7–2.0 = 5.7 kJ mol<sup>-1</sup>). It is clear that the opposite effects exerted by methyl groups at positions 3 and 5, respectively, partially cancel each other when both are present.

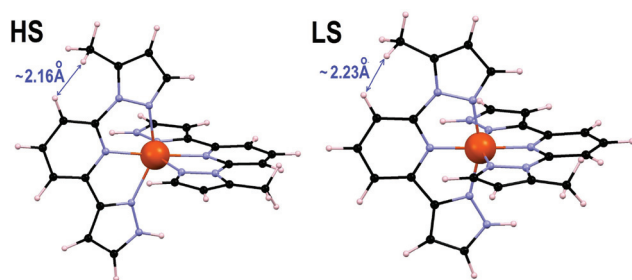
The origin of these  $\Delta H_{\text{elec}}$  values was investigated by analyzing first the effect of the substituents on the  $t_{2g}$  and  $e_g$  orbitals of the Fe ion (Table 2) as was done recently on a family of 1bpp/Fe(II) complexes.<sup>32</sup> Here, the addition of one (**2** and “**4**”) or two (**3**) methyl substituents seems to destabilize the  $e_g$  orbitals (Table 2). However, a clear pattern is not observed for the  $t_{2g}$  set. In any case, the orbital energy splitting does not show any correlation with  $\Delta H_{\text{elec}}$ ; therefore, the effect of the methyl groups ascribed to the ligand field is at best, very minor. Indeed, the previously reported differences in  $\Delta H_{\text{elec}}$  when changing two H atoms by two 4-methyl groups on 1bpp-pyrazoles (*i.e.*, only causing ligand field effects) are less than one order of magnitude smaller than the values in Table 2.<sup>32</sup> This indicates that the differences seen here must be associated, to a large extent, to inter- and/or intra-ligand interactions *within the complex* involving the methyl groups, linked to the changes in Fe–N distances occurring upon SCO. These effects contribute to  $\Delta H_{\text{elec}}$  in two ways: (i) by causing a strain to the overall structure of the [Fe(1,3-bpp)<sub>2</sub>]<sup>2+</sup> core, and (ii) through direct inter- and intra-ligand interactions. In order to analyze the influence of the 3-methyl group, single-point calculations were performed on the optimized structures of **2**<sup>LS</sup> and **2**<sup>HS</sup>, after substituting the methyl group by an H atom, keeping the rest of the geometry untouched. These species, termed here **2**<sup>core</sup>, are equivalent to **1**, but have different HS and LS nuclear configurations. The computed  $\Delta H_{\text{elec}}$  for **2**<sup>core</sup> (12.9 kJ mol<sup>-1</sup>) is *ca.* 1 kJ mol<sup>-1</sup> larger than for **1**; therefore, the strain of the backbone caused by the substituent destabilizes the HS more than the LS state (thus, it opposes the observed overall effect for this substituent). Indeed, comparing the optimized geometries

of **2** and **1** (Fig. S9 and S10†) reveals that the HS structures are much more distant from each other than the LS geometries. Specifically, in the HS, the plane of the 1,3bpp core exhibits a rotation of  $\sim 10^\circ$  around its N<sub>pz</sub>–Fe–N<sub>pz</sub> axis when moving from **1** to **2**. This difference may be due to the steric effect of the 3-methyl, pushing the other ligand back, perhaps also favoring an attractive C–H... $\pi$  interaction between the methyl group and the central pyridine of that ligand (see also Fig. 7). In fact, the optimized structure of **2** also shows a shorter Fe–N bond for the 1-pyrazole than for the 3-pyrazole ring, which could be due to such favorable contact. The difference between **2** and **2**<sup>core</sup> in terms of  $\Delta H_{\text{elec}}$  must be then traced back to the direct interaction between the methyl group and the other ligand. Inspection of the optimized structures **2**<sup>HS</sup> and **2**<sup>LS</sup> shows that the former exhibits the closest contact between the 3-methyl group and the other 1,3bpp ligand of *ca.* 2.7 Å, with this substituent well positioned for the mentioned C–H... $\pi$  contact. The contraction of the Fe(II) coordination sphere caused by the SCO to the LS state forces the methyl group to rotate and exhibit two closest C–H... $\pi$  contacts with the other ligand (now of *ca.* 2.8 Å each), instead of one. The attractive interaction may not be now so favorable or have turned repulsive (the Fe–N bond distance is now shorter for the 3-pyrazole than for the 1-pyrazole ring). This would explain an overall stabilization of the HS state. The *ca.* 3 kJ mol<sup>-1</sup> difference in  $\Delta H_{\text{elec}}$  between **2**<sup>core</sup> and **2** (12.9 *vs.* 9.9 kJ mol<sup>-1</sup>) is in any case the consequence of going from an attractive inter-ligand interaction to a less favorable one.

The individual effect of the methyl group at position 5 is studied by analyzing the hypothetical compound “**4**”. In analogy with the above procedure, we have used the methyl-free “**4**<sup>core</sup>” complex to quantify (i) the strain of the [Fe(1,3-bpp)<sub>2</sub>]<sup>2+</sup> core and (ii) the direct intramolecular interactions. First, the comparison between “**4**<sup>core</sup>” and **1** shows that the HS state is 4.7 kJ mol<sup>-1</sup> less stable in the former case (with  $\Delta H_{\text{elec}}$  of 16.6 kJ mol<sup>-1</sup> in “**4**<sup>core</sup>” compared to 11.9 kJ mol<sup>-1</sup> in **1**). This must be traced back to the presence of the 5-methyl group causing the strain of the [Fe(1,3-bpp)<sub>2</sub>]<sup>2+</sup> core to accom-

**Fig. 7** DFT+U+D2 optimized geometries of complex **2** in both spin states. The HS state structure (left) emphasizes the closest contact between the 3-methyl and the other ligand, as well as the possible C–H... $\pi$  with that ligand. The LS state structure highlights the two closest inter-ligand interactions.





**Fig. 8** DFT+U+D2 optimized geometries of complex “4” in both spin states showing the closest contact between the 5-methyl and the central pyridine of the same ligand.

moderate part of the steric congestion between the 5-methyl and the central pyridine. Second, the difference between “4” (19.6 kJ mol<sup>-1</sup>) and “4<sup>core</sup>” quantifies the direct impact of the intraligand interactions associated with the 5-methyl (Fig. 8), which account for an effect of 3 kJ mol<sup>-1</sup> on  $\Delta H_{\text{elec}}$ , thus completing the difference in  $\Delta H_{\text{elec}}$  of 7.7 kJ mol<sup>-1</sup> between “4” and 1. Therefore, the stabilization of the LS state as a result of an intra-ligand repulsion is here shown and proved theoretically for the first time. This effect was invoked to explain the increase in the SCO temperature of the Fe(II) complex of an indazolepyridine derivative,<sup>40</sup> but a subsequent computational analysis suggested that the HS vs. LS state stability was instead influenced by inter-ligand interactions altering the FeN<sub>6</sub> coordination sphere, and not by such steric hindrance.<sup>37</sup>

## Experimental

### Synthesis

The ligand 2-(pyrazol-1-yl)-6-(1*H*-pyrazol-3-yl)pyridine (1,3bpp) was synthesized as published,<sup>41</sup> using a slight modification of a previously reported procedure.<sup>52</sup> The corresponding complex [Fe(1,3bpp)<sub>2</sub>](ClO<sub>4</sub>)<sub>2</sub> (1) was prepared as previously published.<sup>41</sup> **Caution:** Perchlorate salts of metal complexes are potentially explosive. Only small quantities of material should be prepared and the samples should be handled with care.

**1-(6-(3-Methylpyrazol-1-yl)pyridin-2-yl)ethanone.** To a solution of 1-(6-bromopyridin-2-yl)ethanone (2.5 g, 12.5 mmol) in toluene (15 mL) were added, under a N<sub>2</sub> atmosphere, 3-methylpyrazole (1.53 g, 18.75 mmol), 1,10-phenanthroline monohydrate (0.5 g, 2.5 mmol), CuI (0.24 g, 1.25 mmol) and K<sub>2</sub>CO<sub>3</sub> (1.9 g, 12.5 mmol). The resulting black mixture was heated to reflux and vigorously stirred overnight. After cooling to room temperature, ethyl acetate (20 mL) and water (20 mL) were added and the organic layer was isolated. The aqueous solution was extracted two additional times with ethyl acetate and the organic phases were recombined, washed with brine, dried with MgSO<sub>4</sub> and evaporated under vacuum to afford the product as a brown liquid (2.4 g, 96%). <sup>1</sup>H NMR (400 MHz, CDCl<sub>3</sub>, ppm):  $\delta$  2.37 (s, 3H), 2.72 (s, 3H), 6.22 (d, *J* = 2.5 Hz, 1H), 7.91–7.77 (m, 2H), 8.05–8.02 (m, 1H), 8.45–8.42 (m, 1H).

**1-(6-(3-Methylpyrazol-1-yl)pyridin-2-yl)-3-(dimethylamino)prop-2-en-1-one.** *N,N*-Dimethylformamide-dimethyl acetal (2.5 mL, 24 mmol) was added to 1-(6-(3-methylpyrazol-1-yl)pyridin-2-yl)ethanone (2.4 g, 11.9 mmol) and the mixture was heated to reflux (120 °C) and stirred overnight. After cooling to room temperature, the resulting dark yellow solution was concentrated under vacuum to obtain the product as a brown powder (2.77 g, 91%). <sup>1</sup>H NMR (400 MHz, CDCl<sub>3</sub>, ppm):  $\delta$  2.32 (s, 3H), 2.93 (s, 3H), 3.13 (s, 3H), 6.24–6.11 (m, 1H), 6.46–6.34 (m, 1H), 7.97–7.77 (m, 4H), 6.42 (m, 1H), 8.46 (d, *J* = 2.0 Hz, 1H).

**2-(3-Methylpyrazol-1-yl)-6-(1*H*-pyrazol-3-yl)pyridine (met1,3bpp).** A large excess of hydrazine monohydrate (2.5 mL, 50 mmol) was added to a methanolic solution (25 mL) of 1-(6-(3-methylpyrazol-1-yl)pyridin-2-yl)-3-(dimethylamino)prop-2-en-1-one (2.77 g, 10.8 mmol) and the mixture stirred and refluxed overnight. Subsequently, the resulting solution was cooled to room temperature, producing a pale yellow precipitate that was filtered, washed with water and diethyl ether and dried in air to afford the product as a white powder (1.5 g, 62%). <sup>1</sup>H NMR (400 MHz, CDCl<sub>3</sub>, ppm):  $\delta$  2.33 (s, 1H), 6.22 (d, *J* = 2.5 Hz, 1H), 6.78 (s, 1H), 7.51 (s, 1H), 7.60 (d, *J* = 1.9 Hz, 1H), 7.83–7.72 (m, 1H), 8.45 (d, *J* = 2.4 Hz, 1H), 11.11–10.20 (m, 1H).

**1-(6-(3,5-Dimethylpyrazol-1-yl)pyridin-2-yl)ethanone.** To a solution of 1-(6-bromopyridin-2-yl)ethanone (2.25 g, 11.3 mmol) in toluene (5 mL) were added, under a N<sub>2</sub> atmosphere, 3,5-dimethylpyrazole (0.9 g, 9.36 mmol), 1,10-phenanthroline monohydrate (0.37 g, 1.9 mmol), CuI (0.1 g, 0.5 mmol) and K<sub>2</sub>CO<sub>3</sub> (3.2 g, 23.15 mmol). The resulting dark brown mixture was heated to reflux and vigorously stirred overnight. After cooling to room temperature, ethyl acetate (20 mL) and water (20 mL) were added and the organic layer isolated. The aqueous solution was extracted two additional times with ethyl acetate and the organic phases were recombined, washed with brine, dried with MgSO<sub>4</sub> and evaporated under vacuum. Column chromatography (7 : 3 hexanes/ethyl acetate) provided 0.66 g (32% yield) of the title compound as a white solid. <sup>1</sup>H NMR (400 MHz, CDCl<sub>3</sub>, ppm):  $\delta$  2.31 (s, 3H), 2.70 (s, 3H), 2.76 (s, 3H), 6.04 (s, 1H), 7.97–7.8 (m, 2H), 8.00–7.87 (m, 2H), 8.13 (dd, *J* = 7.8, 1.4 Hz, 1H).

**1-(6-(3,5-Dimethylpyrazol-1-yl)pyridin-2-yl)-3-(dimethylamino)prop-2-en-1-one.** *N,N*-Dimethylformamide-dimethyl acetal (0.3 mL, 2.23 mmol) was added to 1-(6-(3,5-dimethylpyrazol-1-yl)-(1*H*-pyrazol-1-yl)pyridin-2-yl)ethanone (0.66 g, 3.1 mmol) and the mixture was heated to reflux (110 °C) and stirred overnight. After cooling to room temperature, the resulting dark yellow solution was concentrated under vacuum to yield the product as a pale brown solid (0.83 g, 87%). <sup>1</sup>H NMR (400 MHz, CDCl<sub>3</sub>, ppm):  $\delta$  2.31 (s, 3H), 2.78 (s, 3H), 2.96 (s, 3H), 3.19 (s, 3H), 6.02 (s, 1H), 6.45 (d, *J* = 12.7 Hz, 1H), 8.02–7.86 (m, 4H).

**2-(3,5-Dimethylpyrazol-1-yl)-6-(1*H*-pyrazol-3-yl)pyridine (dimet1,3bpp).** An excess of hydrazine monohydrate (0.1 mL, 1.35 mmol) was added to an ethanolic solution (5 mL) of 1-(6-(3,5-dimethylpyrazol-1-yl)pyridin-2-yl)-3-(dimethylamino)prop-2-en-1-one (0.72 g, 0.27 mmol) and the mixture was stirred and



refluxed overnight. The resulting solution was cooled to room temperature and water (5 mL) was added. The organic phase was separated by decantation. The aqueous phase was extracted three times with  $\text{CH}_2\text{Cl}_2$  (10 mL), and the organic layers were then combined, washed with brine, dried with  $\text{MgSO}_4$  and evaporated under vacuum to afford the product as a pale brown solid (0.52 g, 80%).  $^1\text{H}$  NMR (400 MHz,  $\text{CDCl}_3$ , ppm):  $\delta$  2.32 (s, 3H), 2.74 (s, 3H), 6.04 (s, 1H), 6.82 (d,  $J = 2.0$ , 1H), 7.6 (s, 1H), 7.66 (d,  $J = 2.1$ , 1H), 7.83–7.82 (m, 2H), 10.67 (s, 1H).

**[Fe(met1,3bpp) $_2$ ](ClO $_4$ ) $_2$  (2).** To a solution of  $\text{Fe}(\text{ClO}_4)_2 \cdot 6\text{H}_2\text{O}$  (0.023 g, 0.065 mmol) and ascorbic acid (~2 mg) in absolute ethanol (10 mL) was added dropwise a solution of met1,3bpp (0.027 g, 0.12 mmol) in absolute ethanol (10 mL). The resulting dark yellow solution was stirred for 40 minutes at room temperature. The solution was then filtered and layered with hexane (1 : 1 vol.). Yellow crystals of the product suitable for single crystal X-ray diffraction were obtained after 4 days. Yield: 43.2%. EA, calcd (%) for  $\text{C}_{24}\text{H}_{22}\text{Cl}_2\text{FeN}_{10}\text{O}_8$  (found): C, 40.87 (41.04); H, 3.14 (3.01); N, 19.86 (19.06).

**[Fe(dimet1,3bpp) $_2$ ](ClO $_4$ ) $_2$  (3).** To a solution of  $\text{Fe}(\text{ClO}_4)_2 \cdot 6\text{H}_2\text{O}$  (0.023 g, 0.065 mmol) and ascorbic acid (~2 mg) in dry acetone (10 mL) was added dropwise a solution of 2met1,3bpp (0.029 g, 0.12 mmol) in dry acetone (10 mL). In this case, the resulting red solution was stirred for 40 minutes, filtered, and layered with diethyl ether. 3–4 days later, red crystals of the product of good quality for single crystal X-ray diffraction were obtained. Yield: 60.1%. EA, calcd (%) for  $\text{C}_{26}\text{H}_{26}\text{Cl}_2\text{FeN}_{10}\text{O}_8$  (found): C, 42.59 (43.03); H, 3.57 (3.22); N, 19.10 (18.96).

### Single-crystal X-ray diffraction

Data for **2** and **3** were collected on a Bruker APEXII QUAZAR diffractometer equipped with a microfocus multilayer monochromator with  $\text{MoK}\alpha$  radiation ( $\lambda = 0.71073 \text{ \AA}$ ), at 100 and 298/300 K for both compounds. Data for **2b** were collected at 100 K on Beamline 11.3.1 at the Advanced Light Source, on a Bruker D8 diffractometer equipped with a PHOTON 100 CCD detector and using silicon (111) monochromated synchrotron radiation ( $\lambda = 0.7749 \text{ \AA}$ ). Data reduction and absorption corrections were performed with SAINT and SADABS, respectively.<sup>53</sup> All structures were solved by intrinsic phasing with SHELXT<sup>54</sup> and refined by full-matrix least-squares on  $F^2$  with SHELXL-2014.<sup>55</sup> The structure of **2b** was refined as a 2-component twin, using a twin law found through PLATON.<sup>56</sup> The structure was first solved in  $P_1$  on HKLF4 data and then transformed to the monoclinic  $Cc$  with PLATON/ADDSYM. All details can be found in CCDC 1534003–1534004 (**2**), 1534005 (**2b**) and 1534006–1534007 (**3**)<sup>†</sup> that contain the supplementary crystallographic data for this paper. Crystallographic and refinement parameters are summarized in Table S1<sup>†</sup> together with average Fe–N bond lengths and distortion parameters. Selected bond lengths and angles and intermolecular distances are given in Tables S2–S4.<sup>†</sup>

### Physical measurements

**Magnetic measurements.** These measurements were performed with either an MPMS5 or an MPMS-XL SQUID magnetometer through the “Unitat de mesures Magnètiques” of the Universitat de Barcelona or the Servicio General de Apoyo a la Investigación-SAI, Universidad de Zaragoza. Diamagnetic corrections for the sample holder were applied as well as a correction for the diamagnetic contribution of the sample, as derived from Pascal’s constants.

**Differential scanning calorimetry (DSC).** These experiments were done with a Q1000 calorimeter from TA Instruments equipped with the LNCS accessory. Calibration of the temperature and enthalpy scales was achieved with a standard sample of In, using its melting transition (156.6 °C, 3296 J mol $^{-1}$ ). Mechanically crimped Al pans with an empty pan as a reference were used. All reported data were obtained at a scanning rate of 10 K min $^{-1}$ . Measurements on **2** were also done at scanning rates down to 0.5 K min $^{-1}$  to confirm that no hysteresis was present. For heat capacity, a synthetic sapphire was measured in the same temperature range. By comparison, an overall accuracy of 0.2 K for the temperature and up to 10% for the heat capacity was estimated over the whole studied temperature range.

**Powder X-ray diffraction (PXRD).** Patterns were recorded through the X-ray diffraction and fluorescence unit of the Servicio General de Apoyo a la Investigación-SAI, Universidad de Zaragoza, using a D-Max Rigaku diffractometer equipped with a Cu rotating anode and a graphite monochromator to select the Cu  $\text{K}\alpha_{1,2}$  wavelength.

### Computational details

All energy evaluations were performed on molecular geometries optimized in the HS and LS states using the Quantum Espresso package (QE),<sup>57</sup> the PBE +  $U$  functional with a Hubbard-like  $U$  parameter of 2.65 eV on the “d” orbitals of iron, the D2 correction of Grimme,<sup>58</sup> and Vanderbilt pseudopotentials.<sup>59</sup> The molecules were introduced in a cubic cell of 60 Bohr $^3$  to isolate them from the virtual counterparts, which means that all calculations simulate gas-phase conditions. This has been done with the help of the Makov–Payne approximation to treat the charged unit cells.<sup>60</sup> The Hubbard term has been used to cure the incomplete cancellation of the electronic self-interaction in the PBE functional, which results in an unrealistic delocalization of orbitals.<sup>61,62</sup> The value  $U = 2.65$  eV has been found to be adequate to describe  $\Delta H_{\text{elec}}$  in  $\text{FeN}_6$ -based compounds.<sup>63</sup> The  $t_{2g}$  (and  $e_g$ ) orbitals were identified by projecting the density of states on the Fe atom, and the values given in Table 1 correspond to the average value for the three (two) non-degenerate orbitals of the LS species of molecules **1–4**.

### Conclusions

By preparing a family of analogous  $[\text{Fe}(1,3'\text{bpp}')_2](\text{ClO}_4)_2$  complexes, with ‘bpp’ being non-substituted, 3-methyl or 3,5-





dimethyl substituted 1,3-bis-pyrazolylpyridine ligands, it is shown that these remote substituents have a dramatic effect on the SCO temperature of the Fe(II) spin carrier. This influence is manifested on the solid-state thermal behaviour of the concerned systems and most significantly on their SCO in solution, where packing effects are absent. DFT calculations show that these dramatic effects are due to intramolecular steric or non-covalent interactions, which favour either the LS or the HS state, depending on the position of the substituent.

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