Augmented pH-sensitivity absorbance of a ruthenium(II) bis(bipyridine) complex with elongation of the conjugated ligands: an experimental and theoretical investigation

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An absorbance-based sensor employing ruthenium bipyridyl with a phenanthroline-fused benzoylthiourea moiety formulated as [Ru(II)(bpy)₂(phen-nBT)](PF₆)₂ (bpy = 2,2'-bipyridine, phen = 1,10-phenanthroline, nBT = n-benzoylthiourea) has been synthesized and characterized by elemental analyses, mass spectrometry, and infrared, ultraviolet-visible, luminescence and nuclear magnetic resonance spectroscopy. The changes in the intensity of absorption and emission of the complex induced by functionalization of the benzoylthiourea ligands with amino and carbonyl in their protonated and deprotonated forms were studied experimentally. The absorption and emission properties of the complex exhibit a strong dependence on the pH (1–11) of the aqueous medium. This work highlights the pH-sensitivity augmentation of the absorption band by elongating the conjugation length in the structure of the ruthenium bipyridine complex. The principle of this work was to design the title compound to be capable of enhancing the differences in the absorption sensitivity responses towards pH between the protonated and deprotonated complexes in the absorption measurement. Along with significant and noticeable changes in the absorption spectra, subsequent theoretical investigations specifically on the electronic and absorbance properties of the title compound were carried out in this study. Protonation of the molecule significantly stabilized the lowest-unoccupied molecular orbital (LUMO), whereas the highest-occupied molecular orbital (HOMO) is greatly destabilized upon deprotonation. A time-dependent density functional theory (TDDFT) calculation in the linear-response (-LR) regime was performed to clarify the origin of the experimentally observed linear dependence of absorption intensity upon pH (1–11). The MLCT band exhibits hyperchromic shift at low pH as indicated by the large transition dipole moment and a wider distribution of the response charge of the molecule, which is induced by the stabilization of the electrostatic potential at the carbonyl moiety by protonation. This study provides the possibility of employing theoretical information to gain insight into the origin of the optical absorption obtained experimentally. The ruthenium complex was designed with an elongated ligand conjugation length and exhibited a tremendously large change in the absorption intensity of the protonated and deprotonated forms, which therefore demonstrates its feasibility as an indicator molecule especially for absorbance measurements.

1. Introduction

Ruthenium polypyridine complexes have received extensive interest in the past 45 years due to their strong and broad visible light absorption, high chemical stability, tunable electronic properties, high molar absorptivity, efficient emissions and long emission lifetime.¹ ² Their unique characteristics have led to them being used in various potential advantages, making them good candidates for sensitive and efficient receptors for optical sensing applications such as optical pH sensors,³ carbon dioxide sensors, oxygen sensors, temperature sensors⁴ ⁵ and chemical sensors.⁶ An optical pH sensor uses the interaction of light with ruthenium(II)
polypyridyl complexes to determine the pH and offer better thermal performance and faster response time. The design of the sensor is based on the variations in the physical behaviour of the complexes such as the absorbance and fluorescence properties that can be monitored by an optical detector.3

Protonation or deprotonation of basic or acidic sites in metal complexes, respectively, exhibits a dependence on the pH of the medium. The presence of amino,7 carboxylic,8 hydroxyl9 or imidazole-functionalized ligands10 acting as protonatable and deprotonatable moieties in the Ru[n] complexes could regulate the photophysical properties effectively for a pH-responsive on–off switch, and therefore would be useful in the design of a pH sensor. At the equilibrium point, the protonated and deprotonated forms of the Ru(n) complex exhibit variations in their emission and/or absorption spectra, and the pH of the medium may affect the ground state or the excited state species.11 The transfer of protons is essentially a simple chemical signal for monitoring the pH changes.

Even though extensive studies have been performed and numerous pH sensors based on ruthenium compounds have been reported, the changes in the absorption spectra are rather insensitive to the variations of pH. The ruthenium(n) complexes generally have greater sensitivity in the emission than absorption in response to the pH changes.12 Herein, we define the pH sensitivity as an on–off ratio at a fixed wavelength. The on–off ratio is interpreted as the ratio between the absorption intensity of the protonated form to the deprotonated form (or vice versa) at a fixed wavelength. This paper mainly focuses on the measurements of absorption rather than those of emission, owing to their low cost of ownership and ease of handling and maintenance. In addition, the absorption spectrum can be used to detect additional peaks that reveal the presence of impurities in the molecules.

Herein, we give two examples from previous studies related to pH sensitivity of the absorption spectra of ruthenium complexes. For the sake of comparison, we calculated the on–off ratio of the reported complexes in a water solution only. The first study13 of complex \([\text{dip}]_2\text{Ru-mopip}^{2+}\) (dip = 4,7-diphenyl-1,10-phenanthroline and mopip = 2-phenyl-OCH3-imidazo[4,5-f][1,10]phenanthroline) [Fig. 1(a)] revealed an on–off factor of 2.33 at 500 nm in the absorption spectra. More recently, Baba's group14 presented results of pH-based sensing of bis(1,10-phenanthroline)ruthenium(n) complexes bearing hydroxy-substituted ligands \([\text{phen}]_2\text{Ru-(phen-OH)}]^{2+}\) [Fig. 1(b)] that have an on–off factor of 1.60 at 480 nm.

As a result of these observations on pH sensitivity, we conclude that structurally, the protonatable or deprotonatable moieties of a ruthenium complex that are positioned closer to the ruthenium metal have a remarkably large effect on the pH sensitivity in absorption spectra. Therefore, in our endeavour to achieve greater pH sensitivity in the optical intensities, we have attempted to synthesize and characterize a novel ruthenium complex of the type \([\text{Ru}(n)(\text{bpy})_2(\text{phen}-n\text{BT})]^{2+}\), where bpy = 2,2′-bipyridine, phen = phenanthroline, and nBT = n-benzoylthiouracil. The complex was developed as an on–off switch via the protonation or deprotonation of the BT molecules, and amino and carbonyl moieties are contained in the complex. It was designed with an elongated conjugated length in the structure, in which the protonatable and deprotonatable moieties were at a distance from the ruthenium metal. Such a design became important in finding the pH sensitivity of absorption.

We report the synthesis, characterization and effect of pH on the optical intensities of the complex \([\text{Ru}(n)(\text{bpy})_2(\text{phen}-n\text{BT})][\text{PF}_6]_2\). The experimental results were analyzed and compared with quantum chemical calculations. The complex's response to protonation and deprotonation depends significantly on the site that was perturbed by the proton transfer. Thus, the assignment of the deprotonation site may be ambiguous in the complex. In addition, the effect also depends on the configuration of surrounding solvent molecules. The local atomic structures should be accurately taken into account in the theoretical models to calculate the absorption spectra of the complex in different environments. Therefore, we performed first-principles molecular dynamics (FPMD) simulations of the protonated, neutral and deprotonated complexes in aqueous solutions in which the solvent molecules were explicitly taken into account. Subsequently, by using the stable geometry taken from the FPMD simulations, we calculated the optical absorption spectra of the complex using time-dependent density functional theory (TDDFT) and compared them with the experimental results. We also calculated the first-order density response by TDDFT calculation in the linear response (-LR) regime to determine the electronic factors that affect the optical absorption intensities by protonation and deprotonation of the complex.

2. Experimental

2.1 Materials and chemicals

All the chemicals used in this study were of analytical-reagent grade and were used as received without further purification. RuCl3·3H2O and lithium chloride (LiCl) were purchased from Merck, 1,10-phenanthroline-5-amine and benzoyl chloride were obtained from Fluka (Malaysia) and 2,2′-bipyridine (bpy) was obtained from Aldrich Chemical Co.

2.2 Spectroscopic analyses

The infrared spectra were recorded on an Agilent Cary 630 spectrophotometer in the range of 4000–600 cm⁻¹ in the ATR Diamond mode. The micro elemental analyses (C, H, N and S)
were obtained with a Leco 932 elemental analyzer. The 1H and 13C NMR spectra were collected in deuterated dimethyl sulfoxide using a Bruker Ascend 400 spectrophotometer. The chemical shift values were scaled to parts per million (ppm) with reference to tetramethylsilane (TMS). The mass spectrum was recorded in a DMS 2010 Shimadzu mass spectrometer. The magnetic susceptibility measurements were determined on a Sherwood Scientific MSB-AUTO, at room temperature (25.5 °C) using HgCo(SCN)4 as a calibrant.

2.2.1 Absorption and luminescence spectra. The electronic absorption spectra were measured on a Perkin-Elmer Lambda 35 spectrophotometer in the range of 700–200 nm. The steady-state emission spectra were recorded on an Edinburgh FL920 time-resolved fluorescence spectrometer at 298 K in H2O, using a 1 cm path length quartz cell. Time-correlated single-photon-counting (TCSPC) measurements were used for the luminescence decay of the ruthenium complex at different pHs (1–11). The luminescence quantum yields (\(\eta\)) were calculated by employing eqn (1), using [Ru(bpy)3]2+ as the standard.

\[
\eta = \frac{\eta_{\text{std}}}{\eta_{\text{std}}} \frac{A_{\text{std}}}{A_{\text{r}}} \frac{I_{\text{r}}}{I_{\text{std}}} \frac{\tau_{\text{r}}}{\tau_{\text{std}}} \frac{\eta_{\text{r}}}{\eta_{\text{std}}}
\]

where \(\eta\) and \(\eta_{\text{std}}\) are the quantum yield of the ruthenium complex and the standard quantum yields, respectively. The \(\eta_{\text{std}}\) of the [Ru(bpy)3]2+ is 0.042. \(A_{\text{r}}\) and \(A_{\text{std}}\) (<0.1) are the solution absorbance at the excitation wavelength (\(\lambda_{\text{ex}} = 455\) nm), respectively. \(I_{\text{r}}\) and \(I_{\text{std}}\) are the integrated emission intensities, and \(\eta_{\text{r}}\) and \(\eta_{\text{std}}\) are the refractive indices of the solvent. In addition, the radiative lifetime \(\tau_{\text{r}}\) was also calculated by dividing the measured lifetime \(\tau_m\) with \(\eta_{\text{r}}\). The radiative \(k_r\) and non-radiative \(k_{nr}\) rate constants were obtained using eqn (2) and (3).

\[
k_r = \frac{\eta_r}{\tau_m} \quad (2)
\]

\[
k_{nr} = \frac{1}{\tau_m} - \frac{1}{\tau_r} \quad (3)
\]

To investigate the effect of pH on the absorbance and emission intensities, 0.6 mg of [Ru(n)][bpy]3[phen-NBT][PF6]2 was dissolved in 100 mL of deionized water. The resulting solution was recorded at pH 7 (neutral). Five solutions with various pH values (pH = 1, 3, 5, 9, 11) were then prepared by adding drops of 0.1 M NaOH or 0.1 M HCl to the stock solution. The UV-vis spectra of the samples were recorded in the range of 600–300 nm. The emission spectra were then recorded in the range of 800–500 nm, with the excitation measurement being performed at the absorption maximum for the metal-to-ligand charge transfer (MLCT) transition band.

2.3 Computation details

In FPMD simulations, due to the high computational cost to calculate large numbers of complex clusters in an aqueous solution, the protonated, neutral and deprotonated ligands from their respective Ru[n] complexes were treated to determine their optimal structures. The atomic configurations of the ligands in the aqueous solutions were calculated by using the program package STATE (simulation tool for atom technology), which supports the FPMD within the density functional theory. \(^6\)\(^7\)\(^8\) A generalized gradient approximation (GGA) for the exchange-correlation potential as proposed by Perdew, Burke and Ernzerhof (PBE)\(^9\)\(^10\) was used. We also used the ultrasoft pseudopotentials to describe the electron–ion interactions in the system. Wave functions and the augmentation charge were expanded by a plane wave basis set with cutoff energies of 25 and 225 Ry, respectively. The MD time step was 2.4 femtoseconds (fs), and the total simulation time was 3.82 picoseconds (ps). Increasing the hydrogen mass by a factor of twelve in the calculation steps enables the time step to be increased from 0.5 fs to 2.4 fs.

After obtaining the equilibrium geometries of the ligand, we take the ligand and a number of water solvent molecules near the ligand to form a computational model of the Ru complex with the respective ligand for quantum chemical calculations. Their full geometry optimization and electronic structures were calculated by using the Gaussian 09 software package employing the DFT method with Becke’s three-parameter hybrid functional and restricted Lee–Yang–Parr’s gradient corrected functional (RB3LYP).\(^11\)\(^12\) The LANL2DZ basis set was used for Ru atoms and the correlation-consistent basis with polarized valence double zeta (pVDZ) for the remaining atoms. The vibrational frequency calculations were performed to assure that the optimized geometries represented the local minima with only positive eigenmodes. All the orbitals computed were displayed with an isovalue of 0.036 e\(^{-1}\) Bohr\(^{-3}\). In addition, the optical absorption spectra with time-dependent DFT (TDDFT) were calculated following the reported method.\(^23\)\(^24\) 100 singlet excitation energies were calculated with the TDDFT method considering the solvent effect (H2O) with the polarizable continuum model (PCM).\(^25\) The vertical excitation energies at the ground state in TDDFT appeared as excitation peaks in the calculated absorption spectra. The absorption spectra of the complex were simulated by using Gaussian broadening of the calculated oscillator strength at the respective excitation energy. Gaussian function with a bandwidth of 45 nm was used to match the experimentally observed absorption band shapes. Furthermore, a TDDFT calculation in the linear-response (LR) regime was performed to clarify the origin for the experimentally observed linear dependence of absorption intensity upon pH. The excitation energy \(\omega\) and oscillator strength are determined from the square root of an eigenvalue of a matrix eigenvalue problem, as shown in eqn (4).

\[
\Omega F_1 = \omega_r^2 F_1
\]

where \(F_1\) is an eigenvector to illustrate the contribution of occupied-unoccupied orbital pairs on the excitation and to describe the character of the excited state wave function.\(^23\)

2.4 Synthesis

2.4.1 Synthesis of the cis-ruthenium(II) bis(bipyridyl) dichloride, cis-Ru(bpy)2Cl2·H2O. A cis-Ru(bpy)2Cl2·H2O was synthesized according to a previously reported preparation.\(^26\) The apparatus was thoroughly degassed and performed under a gentle stream of argon by using standard Schlenk techniques.
RuCl₃·3H₂O (3.82 mmol) was mixed with bipyridine (7.65 mmol) and lithium chloride (23.59 mmol) in 15 mL of dimethylformamide and the mixture was brought to a refluxing temperature under stirring in a dark environment for 8 hours. The reaction mixture was cooled to room temperature prior to the addition of acetone and the violet reaction mixture was left at 0°C overnight. A crude complex was recovered by filtration to give a black crystalline solid precipitate that was washed twice with cold distilled water (2× 15 mL), followed by cold diethyl ether (2× 15 mL), to yield a pure solid product. Yield: 68%. FTIR (cm⁻¹): ν(CH=CH) aromatic 3031–3099, ν(C=O) aromatic 1670, 1560, 1458, 1417, 1311, 1265, ν(C–H) aromatic 765, 726. λ_max nm [CH₃CN]: 240 (19175), 297 (45383), 376 (86868), 550 (8989). ¹H NMR (400 MHz, DMSO-d₆): δ: 9.95 (d, 2H), 8.64 (d, 2H), 8.50 (d, 2H), 8.07 (t, 2H), 7.77 (t, 2H), 7.67 (t, 2H), 7.49 (d, 2H), 7.12 (t, 2H). ¹³C NMR (400 MHz, DMSO-d₆): δ: 160.73, 158.58, 153.52, 152.80, 125.80, 125.37, 123.77, 124.29, 124.01. Mass spectrum (m/z): ν(NH) 3408, ν(C–Netherocyclic) 1446, ν(C–C) 1587–1484, ν(C–O) 1670, ν(C=O) 1243. ν(uncordinated PF₆) 829. λ_max nm [H₂O], ε (L mol⁻¹ cm⁻¹): 246 (15872), 284 (21870), 451 (4888). Luminescence, wavelength, nm [H₂O]: 617 nm. ¹H NMR (400 MHz, DMSO-d₆), δ: 12.88 (s, 1H), 12.05 (s, 1H), 8.92–8.81 (m, 4H), 8.76 (d, 1H), 8.70 (s, 1H), 8.23 (t, 2H), 8.20–8.10 (m, 4H), 8.11–8.05 (m, 2H), 7.97–7.88 (m, 2H), 7.85 (d, 2H), 7.75–7.65 (m, 2H), 7.64–7.58 (m, 4H), 7.57–7.47 (m, 2H), 7.45–7.37 (m, 2H). ¹³C NMR (400 MHz, DMSO-d₆), δ: 182.40, 168.99, 157.29, 157.09, 153.16, 147.59, 146.35, 138.48, 135.48, 133.88, 133.28, 132.49, 132.13, 129.35, 129.02, 128.52, 127.17, 125.73, 125.00. Mass spectrum (m/z): 917.05 (917.10 theory), [Ru(n)bpy]⁺[phen-NBT][PF₆]³⁻, 771.09 (771.10 theory), [Ru(n)bpy]⁺[phen-NBT]⁺[Na]⁺. 3. Result and discussion
3.1 General
An outline of the synthetic route to complex Ru(bpy)₂Cl₂, ligand phen-NBT and [Ru(n)bpy]⁺[phen-NBT][PF₆]³⁻ is presented in Fig. 2 and 3 shows the protonated, neutral and deprotonated complex structures. The complex [Ru(bpy)₂(phen-NBT)][PF₆]³⁻ was obtained by an in situ reaction of cis-[Ru(bpy)₂Cl₂·2H₂O] with ligand benzoyl-N'-1,10-phenanthroline-5-ylthiourea. The structures of the ligand and complex were elucidated based on elemental analysis and IR, UV-Vis, mass spectrometry, and ¹H and ¹³C NMR spectroscopic measurements. The ruthenium complexes exhibited the expected magnetic susceptibility value of 0 suggesting that the Ru(n) metal centre adopts a diamagnetic d⁶ configuration. This observation was in agreement with the previously reported magnetic properties of similar [Ru(bpy)₂(thio-based imine)][ClO₄]₂ complexes.³⁰
3.2 Experiment on absorption and luminescence spectroscopy
The Ru(n) complex displayed strong absorption bands in the ultra-violent region at 246 nm (ε = 15 872 L mol⁻¹ cm⁻¹) and 284 nm (ε = 21 870 L mol⁻¹ cm⁻¹) that were unambiguously attributed to the π → π* ligand-centred (LC) transitions. This transition basically denotes the excitation of electrons from filled π to vacant π* orbitals of the bpy, phen and NBT moieties.³¹ The complex also exhibited an absorption band in the visible region at 451 nm (ε = 4488 L mol⁻¹ cm⁻¹) that was attributed to the spin-allowed Ru(n)(d₆) → bpy(π*), Ru(n)(d₆) → phen(π*) and Ru(II)(d₆) → BT(π*) MLCT transitions.³²,³³ Excitation of the complex at its MLCT wavelength typically leads to visible light emission, which shows an asymmetric band at 617 nm. This emission band was assigned to radioactive transitions from a Ru(n)(d₆) → π* (ligand) transition.³² Typically, the emission energy of the Ru(n) complex will be smaller than the energy of RuCl₃·3H₂O (3.82 mmol) dissolved in cold distilled water was added into the filtrate, yielding a bright brown precipitate. The precipitate was then filtered through a Buchner funnel, washed with cold H₂O (3× 10 mL), followed by cold diethyl ether (3× 10 mL) and dried in vacuo. Yield: 63%. Anal. calc: C, 45.25; H, 2.85; N, 10.55, S, 3.02. Found: C, 44.42, H, 3.35, N, 9.97, S, 2.58. FTIR (cm⁻¹): ν(NH) 3408, ν(NH) 3281, ν(C–Netherocyclic) 1470, ν(C–C) 1587–1484, ν(C–O) 1670, ν(C=O) 1243. ν(uncordinated PF₆) 829. λ_max nm [H₂O], ε (L mol⁻¹ cm⁻¹): 246 (15872), 284 (21870), 451 (4888). Luminescence, wavelength, nm [H₂O]: 617 nm. ¹H NMR (400 MHz, DMSO-d₆), δ: 12.88 (s, 1H), 12.05 (s, 1H), 8.92–8.81 (m, 4H), 8.76 (d, 1H), 8.70 (s, 1H), 8.23 (t, 2H), 8.20–8.10 (m, 4H), 8.11–8.05 (m, 2H), 7.97–7.88 (m, 2H), 7.85 (d, 2H), 7.75–7.65 (m, 2H), 7.64–7.58 (m, 4H), 7.57–7.47 (m, 2H), 7.45–7.37 (m, 2H). ¹³C NMR (400 MHz, DMSO-d₆), δ: 182.40, 168.99, 157.29, 157.09, 153.16, 147.59, 146.35, 138.48, 135.48, 133.88, 133.28, 132.49, 132.13, 129.35, 129.02, 128.52, 127.17, 125.73, 125.00. Mass spectrum (m/z): 917.05 (917.10 theory), [Ru(n)bpy]⁺[phen-NBT][PF₆]³⁻, 771.09 (771.10 theory), [Ru(n)bpy]⁺[phen-NBT]⁺[Na]⁺. 3.2 Experiment on absorption and luminescence spectroscopy
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the first absorption peak, as the emission of the ruthenium complex is expected to stem from its low-energy triplet state which is phosphorescence, rather than the excited singlet state.34

3.2.1 The effect of pH on the spectroscopic properties. In acid–base equilibrium, the intensity of emission and/or absorption is influenced by pH variation, irrespective of whether the pH affects the ground state or the excited state. If the excited state of the complex is affected by the pH changes, the emission spectra will notably change with the pH and vice versa.11 Fig. 4 depicts the absorption spectra of the ruthenium complex in different pH values and the absorption data of the complex over the pH range of 1 to 11 have been compiled and shown in Table 1. The absorption spectrum of the Ru(n) complex in water exhibits a strong MLCT absorption band at 455 nm that decreases with increasing pH. Four distinct isobestic points were observed which can be divided into two sets – 394 and 505 nm; and 406 and 484 nm. The two sets of isobestic points may suggest that there coexist 2 configurations each for acidic (3–5) and basic (9–11) pH. The result shows that the absorption intensity was high in an acidic medium (low pH) and the intensity was at
least twice that of the basic condition (high pH). This sensitivity stemmed from the long-distance transmission of electronic interactions in the complex and these observations will be explained later from the TDDFT-LR calculations.

We also conducted a luminescence experiment to determine the behaviour of the emission of the complex in water at various pHs (1–11). Fig. 5(a) shows the quenching of the emission spectra of the complex upon pH changes. The maximum peak of the emission intensity is centred at 620 nm and the intensities decreased significantly as the pH increased from 1 to 11, in consonance with the absorption spectral behaviour. The complex shows a stronger emission at low pH and this was attributed to the reduction in the intramolecular reductive quenching effect from the lone-pair electrons of the oxygen atom, as the availability of the lone-pair of electrons for the electron-transfer process is reduced by protonation.\(^{35}\) In addition, there is a significant decrease in the emission lifetime and quantum yields as the pH increases. A longer emission lifetime was observed in a low pH solution due to a deceleration in the thermally activated decay of the low-lying MLCT state to the short-lived metal centre state as a consequence of the stabilization of the MLCT and the destabilization of the metal-centred (MC) state.\(^{36}\)

### Table 1 Absorption and luminescence data for Ru(II) complexes in aqueous solution at different pH values

<table>
<thead>
<tr>
<th>Complex</th>
<th>Absorption (298 K), (\lambda_{abs} = 455 \text{ nm})</th>
<th>Luminescence (298 K), (\lambda_{em} = 620 \text{ nm})</th>
</tr>
</thead>
<tbody>
<tr>
<td>pH 1</td>
<td>7738</td>
<td>0.044 0.501 11.39 0.088 1.908</td>
</tr>
<tr>
<td>pH 3</td>
<td>5292</td>
<td>0.041 0.446 10.88 0.092 2.150</td>
</tr>
<tr>
<td>pH 5</td>
<td>4508</td>
<td>0.039 0.419 10.74 0.093 2.294</td>
</tr>
<tr>
<td>Neutral</td>
<td>4488</td>
<td>0.038 0.391 10.29 0.097 2.460</td>
</tr>
<tr>
<td>pH 9</td>
<td>3392</td>
<td>0.032 0.339 10.59 0.094 2.855</td>
</tr>
<tr>
<td>pH 11</td>
<td>2887</td>
<td>0.027 0.262 9.70 0.103 3.714</td>
</tr>
</tbody>
</table>

\(^a\) Excitation at 455 nm. \(^b\) Quantum yield of complex. \(^c\) Radiative lifetime. \(^d\) Measured lifetime. \(^e\) Radiative rate constants. \(^f\) Non-radiative rate constants.

The ruthenium complex is excited to singlet \(^1\)MLCT states upon optical excitation and will efficiently cross-over to the light emitting triplet \(^3\)MLCT (phosphorescence) state, which then returns to the ground state either by radiative \((k_r)\) or non-radiative \((k_{nr})\) decay processes. The values of \(\varphi, \tau_r, \tau_{nr}, k_r\) and \(k_{nr}\) at different pH (1–11) are catalogued and shown in Table 1. The emission quantum yield decreases as the pH increases owing to the increase in \(k_{nr}\) and decrease in \(k_r\). The deprotonated complex decays more efficiently and rapidly to the ground state by a deactivation pathway compared to the protonated complex, resulting in a shorter emission lifetime and shorter quantum yields. The emission decay profile of complexes at different pH (1–11) is shown in Fig. 5(b).

### 3.3 Theoretical calculations

#### 3.3.1 Structural optimization

To illustrate the electronic structure of the complex under protonated, neutral and deprotonated conditions, DFT and TDDFT calculations were performed. First, the STATE code was used to determine the local geometries of the protonated, neutral and deprotonated ligands in aqueous solution, in which the solvation effect (\(H_2O\)) was explicitly taken into consideration by including 96 water molecules in a periodic box of \(37.8 \times 22.7 \times 22.7 \text{ Å}\). The protonation and neutral ligands, the proton dissociation did not occur and the acidic proton remained at the original site during the molecular dynamics simulations. The protonation site in the optimized structure was identified at the carbonyl oxygen atom. However, there are two possibilities for the deprotonation of the ligands in which the deprotonated site might occur at one of the two N atoms in the structure. Hence, we investigated the preferred deprotonation site by considering the two deprotonated ligands. Two deprotonated ligand structures (phen-nBt_N1 and phen-nBt_N2) were introduced in an aqueous solution, respectively [Fig. 6(a) and (b)]. To observe the relative stability of both the ligands, we calculated the total and average energy of the two geometries during the FPMD simulations and illustrated...
them as a function of time in Fig. 7. A small difference of only 0.04 eV in the average energy was observed and this indicates that both of the deprotonated structures may coexist at the equilibrium.

The optimized structures of the protonated, neutral and both deprotonated (phen-nBT_N1 and phen-nBT_N2) ligands with eight neighbouring H2O molecules at the respective minimums of the potential energy surface were used to be bound with ruthenium bipyridine to form the corresponding complexes. The molecular structures of the complexes \([\text{Ru(II)}(\text{bpy})_2(\text{phen-nBT})/\text{C}_1\text{H}_2\text{O}]^+\) (protonated), \([\text{Ru(II)}(\text{bpy})_2(\text{phen-nBT})/\text{H}_2\text{O}]^0\) (neutral) and \([\text{Ru(II)}(\text{bpy})_2(\text{phen-nBT})/\text{C}_0\text{H}_2\text{O}]^{1-}\) (deprotonated_N1 and deprotonated_N2) were then optimized by using DFT calculation with the B3LYP exchange–correlation functional. Consequently, the protonation site in the optimized structure was identified at the \(-\text{C}=\text{O}–\) atom, whereas deprotonation occurred at the \(-\text{N}–\) atom in the complex. The optimized molecular structures of the protonated, neutral, deprotonated_N1 and deprotonated_N2 complexes are shown in Fig. 8.

3.3.2 Theoretical absorption spectra. The calculated lowest-energy absorption bands are presented in the visible region (380–440 nm) in Fig. 9. As shown in Fig. 9, the protonated complex has a stronger absorption intensity than the deprotonated one, agreeing well with the experimental observation. Table 2 shows the frontier molecular orbital isosurface of dominant transitions that comprise the lowest-energy absorption bands.

The three highest occupied MOs (HOMO, HOMO−1 and HOMO−2) for the protonated and neutral complexes essentially have ruthenium d(t2g) characteristics. Nonetheless, HOMO and HOMO−1 are mainly composed of the \(\pi\)-orbital of the [−C(S)−N−C(phen)=] and [−C(S)−N−C(O)=] moieties for the deprotonated_N1

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**Fig. 6** Molecular dynamics (MD) simulations of a phen-nBT_N1 (a) and phen-nBT_N2 (b) ligand in aqueous solution. The red circle shows the deprotonated site of the nitrogen atom.

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**Fig. 7** Energy level diagrams of the ligands. The horizontal line indicates average energy obtained.

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**Fig. 8** Theoretical molecular structures of the complexes.
and deprotonated_N2 complexes, respectively. This has been attributed to the increase in the electrostatic potential of the carbon atom (phen), carbonyl and thiocarbonyl moieties upon deprotonation of amino group in the complexes. As a result, the π-orbitals of the \([-\text{C(S)–N–C(phen)}–]\) and \([-\text{C(S)–N–C(O)–}]\) moieties are destabilized. However, different deprotonated sites of the amino group in the complexes made little difference to the energy level of the π-orbital, where the deprotonated_N2 complex shows slightly higher energy (\(-5.635\) eV) compared to the deprotonated_N1 complex (\(-5.905\) eV). As for the protonated complex, the \(d_{\text{aa}}\) orbitals (\(d_{\text{xy}}, d_{\text{xz}}\) and \(d_{\text{yz}}\)) are predominantly located in the frontier HOMO–2, HOMO–3 and HOMO–4 orbitals. However, the LUMO is altered from \(\pi^*\) (bpy) to \(\pi^*\) of the benzoylthiourea moiety by protonation. This is due to the significant stabilization of the \(\pi^*\) orbitals by protonation of the carbonyl moiety. In contrast, the LUMO and three subsequent virtual orbitals for the deprotonated and neutral complexes are essentially \(\pi^*\) of the bpy and phen moieties, hence, the energy levels of the LUMO orbitals are almost similar. It is interesting to note that protonation strongly affects the LUMO orbitals, whereas deprotonation affects the HOMO orbitals. Thus, the frontier molecular orbitals of the lower-energy occupied orbitals in the protonated complex and the frontier molecular orbitals of the lower-energy unoccupied orbitals in deprotonated complexes are similar to those of the neutral complex.

Fig. 9 shows the calculated absorption spectra and the density response mapped over the electronic isodensity surface for the protonated, neutral and deprotonated complexes. As mentioned above, the trend in the calculated spectra agrees very well with the experimental observation. The protonated complex displays a higher intensity in the absorption, whereas both the deprotonated complexes show a relatively lower absorption band of similar intensities. The experimentally observed hyperchromic effect in the absorption spectra at low pH was demonstrated by the theoretical calculation to be related to the large transition dipole moment of the molecules. This is ascribed to the high perturbation electronic density of the molecule oscillators between the electronic states. In accordance with published work by Yanagisawa et al., the first order density response equation [eqn (5)] was revealed to be the most appropriate to illustrate the transition dipole moments.

\[
\rho_{\text{II}}(r) = \sum_{i,a} \phi_i(r)\phi_a(r)F_{ia}
\]  

Eqn (5) describes the density response where \(\rho_{\text{II}}\) are the final and initial states of the density response, \(i\) and \(a\) indicate occupied (HOMO) and unoccupied (LUMO) Kohn–Sham orbitals, respectively, and \(F_{ia}\) depicts the eigenvector component of the transition between \(\phi_i\) and \(\phi_a\). The density responses of the complexes are calculated on the basis of the highest oscillator strength of the calculated optical spectra. There are smaller albeit noticeable differences in the isosurface structures of the protonated, neutral and deprotonated complexes. We found that the electronic charge was distributed over almost the entire ligand in...
the protonated complex [see Fig. 9(a)]. In particular, the stabilization of electrostatic potential at the carbonyl moiety upon protonation leads to a higher charge distribution on the protonated molecule. Therefore, a greater distribution of the density response in the protonated complex leads to a larger transition dipole moment and hence a higher absorption intensity. Likewise, the localization of the density response in the neutral and both the deprotonated complexes is in accordance with the destabilization of the electrostatic potential at the carbonyl moiety. The electronic charge is distributed more equally in the neutral complex and leads to a slightly greater absorption intensity compared with the deprotonated complexes. We believed that the position of the protonatable or deprotonatable moiety in the ruthenium complex inevitably exerts an influence on the pH sensitivity of absorption. The stabilization of the electrostatic potential by protonation accelerates a wider distribution of response charge of the molecule and thus generates a hyperchromic effect on the MLCT absorption band at a low pH.

4. Conclusion

We reported experimental and theoretical investigations on the photophysical properties of novel ruthenium bipyridine complexes bearing a phenanthroline moiety with a benzoylthiourea unit. A simple synthetic route and purification protocol for the ruthenium complexes was demonstrated. Owing to the presence of the functionalised benzoylthiourea with the carbonyl and amino moieties, the complex can exist in three different forms and can readily be protonated and deprotonated under acidic and basic conditions, respectively. The experimental result shows a remarkable improvement in the absorbance response of the complex in the pH range of 1–11. Therefore, a theoretical study using DFT and TDDFT calculations was undertaken on the absorbance behaviour of the complex in the protonated and deprotonated forms to account for the observations. The FPMD simulations revealed that the protonation site in the optimized structure was at the carbonyl oxygen atom, whereas the coexistence of the two deprotonation structures was observed at the equilibrium. The observed trends in the absorption spectra can effectively be reproduced by TDDFT in the linear-response regime using complex geometries optimized using FPMD simulations. Theoretical analyses determined the relationship between the distribution of the electrostatic potential and the absorption intensities under acidic and basic conditions. It was observed that a delocalized distribution of the density response leads to a higher absorption intensity. It is worth underlining that the elongation of the
Table 2  Isosurfaces of frontier molecular orbitals (from HOMO–4 to LUMO+3), which dominate transitions corresponding to the lowest-energy absorption, and the orbital energy values for protonated, neutral and deprotonated (N1 and N2) complexes. Water molecules in complexes are omitted for brevity.

<table>
<thead>
<tr>
<th>Orbitals</th>
<th>Protonated</th>
<th>Neutral</th>
<th>Deprotonated_N1</th>
<th>Deprotonated_N2</th>
</tr>
</thead>
<tbody>
<tr>
<td>LUMO+3</td>
<td>(-2.524 eV)</td>
<td>(-2.435 eV)</td>
<td>(-2.106 eV)</td>
<td>(-2.337 eV)</td>
</tr>
<tr>
<td>LUMO+2</td>
<td>(-2.541 eV)</td>
<td>(-2.474 eV)</td>
<td>(-2.365 eV)</td>
<td>(-2.449 eV)</td>
</tr>
<tr>
<td>LUMO+1</td>
<td>(-2.637 eV)</td>
<td>(-2.503 eV)</td>
<td>(-2.496 eV)</td>
<td>(2.495 eV)</td>
</tr>
<tr>
<td>LUMO</td>
<td>(-2.905 eV)</td>
<td>(-2.595 eV)</td>
<td>(-2.574 eV)</td>
<td>(-2.583 eV)</td>
</tr>
<tr>
<td>HOMO</td>
<td>(-6.180 eV)</td>
<td>(-6.130 eV)</td>
<td>(-5.905 eV)</td>
<td>(-5.635 eV)</td>
</tr>
<tr>
<td>HOMO–1</td>
<td>(-6.325 eV)</td>
<td>(-6.265 eV)</td>
<td>(-5.995 eV)</td>
<td>(-5.932 eV)</td>
</tr>
<tr>
<td>HOMO–2</td>
<td>(-6.355 eV)</td>
<td>(-6.312 eV)</td>
<td>(-6.063 eV)</td>
<td>(-6.119 eV)</td>
</tr>
<tr>
<td>HOMO–3</td>
<td>(-6.852 eV)</td>
<td>(-6.410 eV)</td>
<td>(-6.102 eV)</td>
<td>(-6.261 eV)</td>
</tr>
<tr>
<td>HOMO–4</td>
<td>(-7.183 eV)</td>
<td>(-6.729 eV)</td>
<td>(-6.267 eV)</td>
<td>(-6.304 eV)</td>
</tr>
</tbody>
</table>
conjugated path length in the structure of the phenanthroline-fused benzothiourea moiety in the ruthenium bipyridine complex eventually results in greater pH sensitivity in the absorption spectra. The present communication aims to provide new insights and useful information for the design of new materials for an absorption-based pH sensor.

Conflicts of interest

There are no conflicts to declare.

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References

24 M. E. Casida, C. Jamorski, K. C. Casida and D. R. Salahub, Molecular excitation energies to high-lying bound states from time-dependent density-functional response theory:


