

Cite this: *CrystEngComm*, 2017, 19, 1809

Methyl substituent effect on structure, luminescence and semiconducting properties of furan/phenylene co-oligomer single crystals†

Maxim S. Kazantsev,^{*ab} Alina A. Beloborodova,^{ab} Ekaterina S. Frantseva,^{ab} Tatyana V. Rybalova,^{ab} Vladislav G. Konstantinov,^c Inna K. Shundrina,^{ab} Dmitry Yu. Paraschuk^c and Evgeny A. Mostovich^{ab}

Single crystals of furan/phenylene co-oligomers are among the most promising highly-emissive materials for applications in various optoelectronic devices. In this work, we synthesized and studied furan/phenylene co-oligomers with the same conjugated core 1,4-bis(5-phenylfuran-2-yl)benzene and methyl substituents at *p*- and *m*-positions of the terminal phenyls. The effect of substituents on the crystal packing, charge transport and luminescence of the single crystals was studied. Compared to the unsubstituted compound, the methyl-substituted co-oligomers demonstrated improved thermostability and enhanced photo-luminescence, which we assign to J-aggregation resulting from the strong inclination of the molecules against the main crystal facet. The charge mobility in single crystal organic field-effect transistors decreased upon the inclination of the molecules. We conclude that the molecular tilt angle, intermolecular distances and interactions in crystals of heteroaryl-containing linear conjugated oligomers can be controlled by the introduction of end methyl groups in the appropriate positions.

Received 14th December 2016,
Accepted 28th February 2017

DOI: 10.1039/c6ce02565j

rsc.li/crystengcomm

1. Introduction

Organic highly emissive semiconducting single crystals attract much attention because of their high performance in organic light-emitting devices, such as light-emitting transistors and lasers.^{1–4} Crystals for light-emitting transistors should have, on the one hand, high charge mobility, which is highest for the perpendicular orientation of the molecules against the crystal basal plane along which the electrical current flows and, on the other hand, efficient light outcoupling. The latter is usually impeded by the strong waveguide effect, which is first of all dictated by the orientation of the molecular transition dipole moments defining the direction of primary emission.^{1,5} Evidently, for linear conjugated molecules the light outcoupling efficiency will be maximal for the horizontal molecular orientation with respect to the basal plane. Therefore, one may expect that there is an inclination of the molecules

corresponding to the optimum balance between the charge mobility and luminescence. Accordingly, control of molecular inclination against the main crystal facet is a promising tool for the development of high-performance single crystal light-emitting devices.

Single crystals of small conjugated molecules, due to their highly ordered and low-defect structure, also provide a unique opportunity for the investigation of the fundamental structure–property relationship *via* systematic design and the study of libraries of related compounds with peculiar molecular packing. For example, by varying a number of aromatic rings in thiophene/phenylene co-oligomers, single crystals with high mobility⁶ and high luminescence efficiency were found.^{7,8} The introduction of terminal substituents in thiophene/phenylenes was demonstrated to be a powerful tool for solubility and morphology control.^{9,10} Distyrylbenzenes (DSB) were also shown to give highly emissive single crystals and besides terminal substitution could be modified using vinylene units resulting in a tremendous library of DSB derivatives having different intra- or intermolecular interactions and demonstrating specific optical and charge transport properties.¹ Aggregation motifs in DSB single crystals were intensively explored by substitution of the backbone by methyl,^{11,12} *tert*butyl,¹³ methoxy,¹⁴ cyano,^{14,15} fluorine,^{13,14,16} trifluoromethyl^{15,17} and other¹ groups. Among the variety of substituents, the methyl group is the most compact and minimally disturbs the molecular electronic structure; therefore, end-group methylation might be a tool for fine

^a N.N. Vorozhtsov Novosibirsk Institute of Organic Chemistry, Lavrentieva 9, Novosibirsk, 630090, Russian Federation. E-mail: maximkazantsev1988@gmail.com

^b Novosibirsk State University, Pirogova 2, Novosibirsk, 630090, Russian Federation

^c Faculty of Physics and International Laser Center, Lomonosov Moscow State University, Leninskie Gory 1/62, Moscow, 119991, Russian Federation

† Electronic supplementary information (ESI) available: Synthetic procedures, spectra of synthesized compounds, cyclic voltammetry and optical spectroscopy data, microscopy images, additional X-ray data and OFET data. CCDC 1520469–1520471. For ESI and crystallographic data in CIF or other electronic format see DOI: 10.1039/c6ce02565j



tuning the crystal packing, and hence the intermolecular interactions. For example, 4,4'-methyl-substituted DSB single crystals demonstrated high photoluminescence quantum yields (PL QYs), amplified spontaneous emission and efficient charge transport.¹¹ Varghese *et al.* reported unusual co-facial arrangement, high PL QYs and stimulated emission in 3,5,3',5'-tetramethyl DSB single crystals.¹² However, methyl-substituted heteroaryl-containing co-oligomers have not yet been reported.

Recently, we have demonstrated a furan/phenylene co-oligomer, 1,4-bis(5-phenylfuran-2-yl)benzene (BPFB, Fig. 1a), outperforming its thiophene/phenylene analogue in PL efficiency.¹⁸ Here, we report on the dramatic effect of methyl substituents on the crystal packing and optoelectronic properties of the furan/phenylene co-oligomer BPFB. We compare vapor-grown single crystals of co-oligomers substituted at *p*- and *m*-positions of the terminal phenyls and demonstrate that the tilt angle and intermolecular distances could be effectively tuned without distortion of the molecular electronic structure.

2. Experimental

2.1 Synthesis

All reagents and solvents were purchased from commercial sources (Sigma-Aldrich, Acros) and were used without additional purification. Reaction mixtures were monitored by TLC using Sorbfil HPTLC-AF-V-UV plates and hexane as the eluent. For column chromatography, Masherey-Nagel Kieselgel 60 and pure hexane as the eluent were used. The detailed synthetic procedures are presented in the ESI.†

2.2 Characterization

Combustion analysis was performed using a CHN-analyzer (EURO EA). ¹H NMR spectra were recorded on a Bruker AV

400 (400.13 MHz) spectrometer in CDCl₃. Mass spectra were obtained with a Thermo Electron Corporation DFS mass spectrometer (70 eV) using direct injection; the temperature of the ionization chamber was 220–270 °C. IR spectra were recorded in transmission mode on a Bruker Tensor 27 FT-IR spectrometer in potassium bromide pellets. Thermogravimetric analysis and differential scanning calorimetry measurements were performed using a NETZSCH STA 409 instrument at a heating rate of 10 °C min⁻¹ in an inert (He) atmosphere. Cyclic voltammetry measurements were performed in CH₂Cl₂ solution with a computer-controlled P-8nano potentiostat/galvanostat (Elins, Russia) in combination with a three-electrode cell (Gamry), and 0.1 M tetrabutylammonium hexafluorophosphate was used as a supporting electrolyte. Pt, a Pt wire and Ag/AgCl were used as the working, counter and reference electrodes, respectively. The reference electrode was calibrated by measuring the redox potential of ferrocene. UV/VIS spectra were recorded in diluted (10⁻⁵ M) THF solutions and solid samples in potassium bromide pellets on a Varian Cary 5000 UV-VIS-NIR spectrophotometer. The PL QY in THF solution was measured with a Varian Cary Eclipse fluorescence spectrophotometer according to a standard procedure, and a diluted solution of 1,4-bis-(5-phenyloxazole-2-yl)benzene in ethanol (PL QY = 0.885) was used as a reference standard.¹⁹ The PL QY of solid samples was measured using a calibrated 3.3 inch diameter integrating sphere (Newport 819C-SL-3.3) coupled to a spectrometer (InVia, Renishaw). Solid samples on glass substrates were excited at 405 nm using a semiconductor laser. The detailed experimental procedure was described in our previous work.¹⁸

2.3 Crystal growth and analysis

For the crystal growth, physical vapour transport in a stream of high-purity helium (99.995%) was used. Powders were sublimed in a glass tube (13 mm i.d.) placed into a copper tube with a heater to create a temperature gradient.²⁰ The temperature was set near the *T*_{10%} of weight loss (Table 1). The crystal growth time typically was overnight (~14 h). The crystals were examined using an optical microscope (MP-7, "Lomo", Russia) using transmitted light through crossed polarizers and under blue laser irradiation (405 nm). The X-ray diffraction experiments were performed using a Bruker KAPPA APEX II diffractometer with graphite monochromated MoK α radiation. Integration and scaling of the intensity data were accomplished by means of SAINT.²¹ Absorption corrections were applied using SADABS.²² The structures were solved by direct methods with SHELXS-97.²³ Refinement was carried out by the full-matrix least-squares technique with SHELXL.²³ All non-hydrogen atoms were refined anisotropically. The hydrogen atom positions were calculated geometrically and refined isotropically according to the riding model. The crystal structures were analysed for intermolecular interactions using PLATON.^{24,25} The asymmetric unit in all the three cases contains half of a molecule because of a molecule lying on

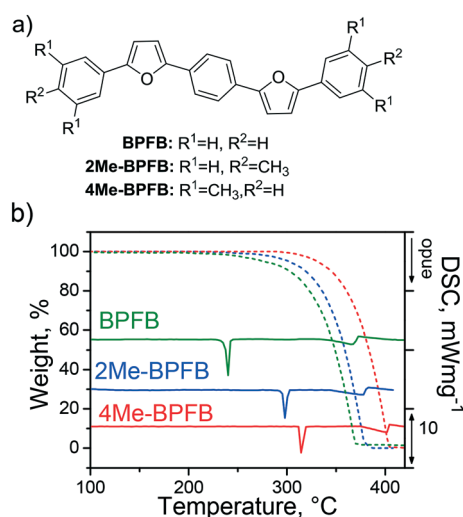


Fig. 1 a) Chemical formulae of the studied furan/phenylene co-oligomers; b) thermogravimetric (dashed lines) and differential scanning calorimetry (solid lines) analyses of BPFB (olive), 2Me-BPFB (blue) and 4Me-BPFB (red) in an inert (He) atmosphere.



Table 1 Summarized electronic and thermal properties. E_{homo} – energy of highest occupied molecular orbital (estimated from CVA) and E_{lumo} – energy of lowest unoccupied molecular orbital (calculated via $E_{\text{lumo}} = E_{\text{homo}} + E_{\text{g}}$, where E_{g} is the optical gap corresponding to the absorption spectrum edge); T_{m} is the melting temperature; $T_{10\%}$ is the temperature of 10% weight loss under sublimation in a He atmosphere; ΔH is the melting enthalpy

	E_{homo} , eV	E_{lumo} , eV	E_{g} , eV	T_{m} , °C	$T_{10\%}$, °C	ΔH , kJ mol ⁻¹
BPFB	-5.26 ^a	-2.26	3.00	238 ^a	302 ^a	46
2Me-BPFB	-5.19	-2.20	2.99	299	318	57
4Me-BPFB	-5.19	-2.20	2.99	316	344	66

^a Ref. 18.

the centre of symmetry. Crystallographic data for the structures have been deposited at the Cambridge Crystallographic Data Centre as supplementary publication CCDC no. 1520469–1520471.

2.4 Single crystal field-effect transistors

Organic field-effect transistors (OFETs) were fabricated on glass substrates using a top-contact/top-gate configuration reported elsewhere.²⁶ The crystals were attached to the substrates using cyanoacrylate glue. Source and drain contacts were painted on the surface of the crystals using a colloidal graphite suspension (PELCO, TED PELLA, INC). The OFET channel width and length were measured using an optical microscope. As a top dielectric layer, 1.5 μm thick parylene N ($C = 1.6 \text{ nF cm}^{-2}$) was deposited from [2,2]paracyclophane (Sigma Aldrich), as described elsewhere.²⁰ Finally, the gate electrode was painted using graphite paste. The OFET current-voltage characteristics were recorded using a source-meter instrument (Agilent B2902A) under dark conditions, ambient atmosphere and room temperature. The voltage sweep rate was 0.3 V s^{-1} . For the determination of field-effect mobility, the saturation regime was used, and the source-drain current (I_{sd}) was fitted according to the equation:²⁷

$$I_{\text{sd}}^{1/2} = \left(\frac{W}{2L} \mu C \right)^{1/2} (V_{\text{g}} - V_{\text{th}})$$

where C is the capacitance per unit area of the insulating layer; W and L are the channel width and length; μ is the carrier mobility; V_{g} and V_{th} are the gate and threshold voltages, respectively.

3. Results and discussion

3.1 Synthesis and characterization

Fig. 1a demonstrates the chemical formulae of the studied 1,4-bis(5-phenylfuran-2-yl)benzene (BPFB) and its derivatives with multiple methyl substituents in the 4,4'-(2Me-BPFB) and 3,5,3',5'-(4Me-BPFB) positions of the terminal phenyls. The target compounds were synthesized according to a previously described procedure¹⁸ (see the ESI†). The methyl substituents only slightly affect the electronic structure (Table 1), as monitored from the cyclic voltammetry measurements (Fig. S8†) and optical spectroscopy data (Fig. S9a†). Therefore, mainly intermolecular interactions affect the properties

and optoelectronic performance of substituted BPFB single crystals.

Thermogravimetric analysis and differential scanning calorimetry (Fig. 1b) reveal a significant increase in the melting and sublimation temperatures for 2Me- and 4Me-BPFB compared to the parent compound (Table 1). The melting enthalpy also increases upon the addition of methyl groups, suggesting the amplification of intermolecular interactions. Indeed, 2Me-BPFB and 4Me-BPFB are unexpectedly very poorly soluble in toluene (lower than 0.01 g l^{-1}) in contrast to the previously reported 1 g l^{-1} for unsubstituted BPFB.¹⁸ As a result, we used physical vapor transport for the crystal growth.

3.2 Single crystal analysis

Fig. 2 demonstrates the optical images of 2Me- (a and b) and 4Me-BPFB (c and d) single crystals which have an elongated plate shape in contrast to solution-¹⁸ and vapor-grown BPFB single crystals grown for comparison (Fig. 2e and f). In-plane rotation of the crystals in crossed polarizers results in completely dark images demonstrating linear birefringence (Fig. S10†). The lengths of 2Me- and 4Me-BPFB single crystals

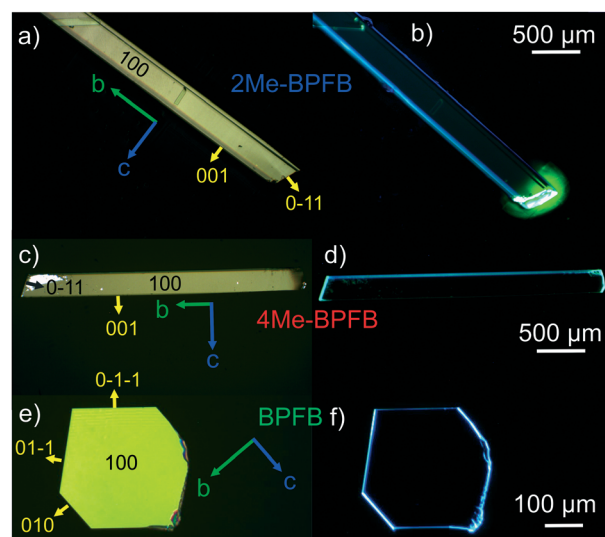


Fig. 2 Microscopy images of single crystals: 2Me-BPFB (attached to a substrate by cyanoacrylate glue) a) in crossed polarizers and b) under blue irradiation; 4Me-BPFB c) in crossed polarizers and d) under blue irradiation; BPFB e) in crossed polarizers and f) under blue irradiation; the arrows indicate the orientation of crystallographic axes and Miller indices of the crystal facets.



(up to 1 cm) are typically much longer than their width (~ 200 μm) and thickness (~ 50 – 100 μm). PL in crystals was observed mainly at the edges and defects (Fig. 2b and d) due to a strong waveguide effect previously reported for thiophene/ and furan/phenylene co-oligomer single crystals.^{8,18,28,29}

X-ray diffraction experiments on the single crystals revealed the same crystal group $P2_1/c$ for all three compounds. The main crystal facet corresponds to the (100) plane for all the crystals, and the 2Me- and 4Me-BPFB ones are extended along the b -axis (Fig. 2a and c). The π -system of 2Me-BPFB is the most planar, as monitored from the values of torsion angles ϕ_1 and ϕ_2 and the maximal deviation d_{max} from the average plane of the π -system demonstrated by the C10 atom (Table S1†). The molecules have a typical herringbone packing motif controlled mainly by C–H $\cdots\pi$ interactions. While BPFB has an almost right tilt angle between the long molecular axes (connecting two carbons of terminal phenyls) and the main crystal facet, the methyl-substituted compounds demonstrate remarkable inclination of the molecules (Fig. 3). This inclination results in the decrease in crystal density and increase in intermolecular distances (Table S1†). Therefore, one may expect a decrease in charge mobility for methyl-substituted BPFBs (*vide infra*).

In contrast to BPFB crystals, where the C–H $\cdots\pi$ interactions are within the layer, in 2Me-BPFB crystals similar interactions between the neighboring layers are observed (Fig. 4). These interlayer interactions in 2Me-BPFB crystals stabilize the structure and increase the melting temperature (Table 1). In 4Me-BPFB, due to steric hindrance, similar interlayer interactions are restricted. Nevertheless, the additional interactions (Fig. 4d) between the methyl groups and the π -systems of neighboring molecules within the layer contribute to the increase in thermal stability (Table 1).

Note that 2Me- and 4Me-BPFB have tilt angles as low as 28 and 36°, respectively, which are among the smallest for the thiophene/phenylene^{6,10,30–32} and furan/phenylene²⁹ co-oligo-

mers. Note that a similar inclination is not so pronounced in methyl-substituted DSBs.^{11,12} Such inclination of molecules is usually achieved by the introduction of bulky substituents^{31,33} or highly electron-withdrawing groups^{1,14,16,32,34} leading to strong intermolecular interactions. It is very unusual that in the case of furan/phenylene co-oligomers, terminal methyl substituents lead to high inclination which could be exploited for the design of materials with desired orientation of luminophores without distortion of their molecular electronic structure. Remarkably, by introducing the methyl groups in appropriate positions, the inclination angle could be varied. And this effect is not directly interrelated with the bulkiness of the substituent because the relatively small terminal methyl substituents in 2Me-BPFB induce extremely high inclination (tilt angle 28°) which is even higher than those in more sterically hindered 4Me-BPFB (tilt angle 36°). Note that assuming the direction of transition dipole moments is parallel to the long molecular axes, the tilt angles of 2Me- and 4Me-BPFB correspond to J-aggregation,³⁵ which is beneficial for high luminescence.¹

3.3 Luminescence and charge transport properties

The optical properties of BPFB, 2Me- and 4Me-BPFB were studied in diluted solutions and the solid state. The solution absorption and emission spectra of all the co-oligomers are almost identical (Fig. S9a†) except for a small bathochromic shift for the methyl-substituted compounds caused presumably by a weak electron-donor effect of methyl groups. All the compounds are highly emissive in diluted solutions with a

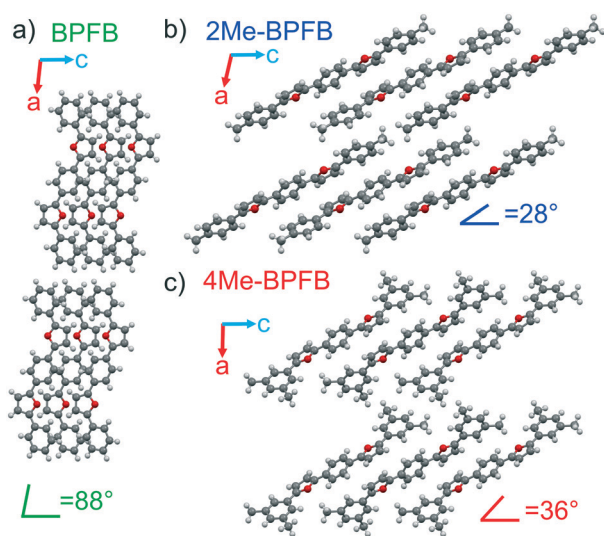


Fig. 3 Crystal structures of the studied co-oligomers: a) BPFB, b) 2Me-BPFB and c) 4Me-BPFB; view along the b -axis.

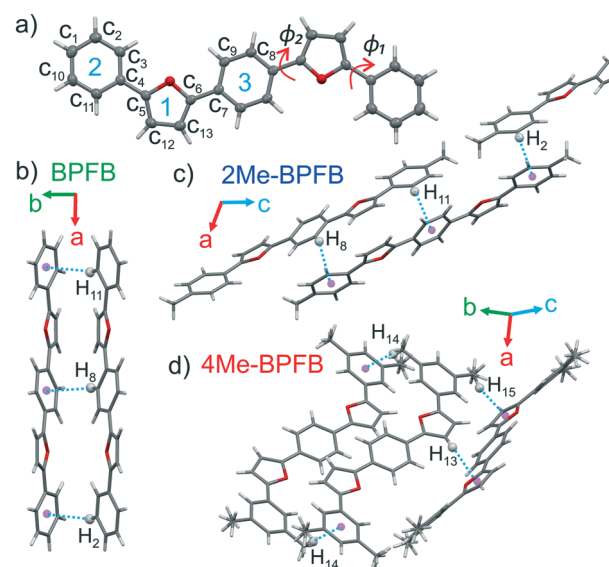


Fig. 4 a) Molecular structure, atom and cycle numbering of BPFB; intermolecular C–H $\cdots\pi$ interactions in the crystals of: b) BPFB, c) 2Me-BPFB, and d) 4Me-BPFB (the corresponding parameters are presented in Table S2†); the numbering for 2Me-BPFB and 4Me-BPFB is the same as for BPFB, except for methyl groups marked as C₁₄ for 2Me-BPFB and C₁₄ and C₁₅ for 4Me-BPFB; ϕ_1 and ϕ_2 correspond to the twisting angles between phenyl and furan rings.



PL QY of about 90% (Table 2). The crystals of methyl substituted BPFBs were also strongly emissive with high quantum yields of 76 and 80% for 2Me- and 4Me-BPFB, respectively (Table 2 and Fig. S9b†), in contrast to BPFB single crystals showing a PL QY of only 45% (Table 2). Therefore, the methyl substituents enhance the PL QY of BPFB-based single crystals. The high PL QY and red shift of the PL spectra of 2Me- and 4Me-BPFB compared to BPFB single crystals (Table 2) are consistent with the J-aggregation model. On the other hand, this red shift is partially caused by a strong reabsorption effect³⁶ so that the longer the crystal, the more red-shifted the emission observed. For example, green emission was observed at the ends of the long crystals (Fig. 2b). To highlight the effect of aggregation on optical absorption, we recorded the optical density spectra of the ground crystals pressed in potassium bromide pellets (Fig. S9c†). All the spectra of the solid samples show red shifts of the 0–0 transition compared to the solution spectra. These shifts can be partially assigned to the higher polarizability of the molecular environment in crystals. Meanwhile, 2Me- and 4Me-BPFB demonstrate considerably higher red shifts in agreement with their J-aggregation, and the values of the red shifts are correlated with the tilt angles of the molecules.

To study the charge transport properties of the crystals, we fabricated top-contact/top-gate OFETs (due to the high thickness of the crystals, they could not be laminated on hard SiO₂ substrates to study devices in common top-contact/bottom-gate architectures). An optical image of a typical single crystal with painted source-drain contacts is shown in Fig. S11†. We did not select any specific charge transport direction for the BPFB single crystals, whereas for the 2Me- and 4Me-BPFB crystals the charge transport was studied along the *b*-axis. Note that this charge transport direction corresponds to the shortest intermolecular distances (Table S1†). Fig. 5 demonstrates the output (a and c) and transfer (b and d) current-voltage characteristics corresponding to p-type conductivity in 4Me-BPFB and BPFB single crystals, respectively.

The mobility of a 4Me-BPFB single crystal OFET is 0.039 cm² V⁻¹ s⁻¹, which is almost 3 times lower than that of an unsubstituted BPFB single crystal (Fig. 5d). Therefore, despite the inclination of the molecules, 4Me-BPFB still demonstrates reasonable charge transport. We fabricated a series of 3 devices for BPFB and 4Me-BPFB (Table S3†) demonstrating reproducible OFET characteristics. However, for all 6 fabricated devices on 2Me-BPFB single crystals no field-effect was observed at all. This lack of OFET performance could be due to high inclination of the molecules

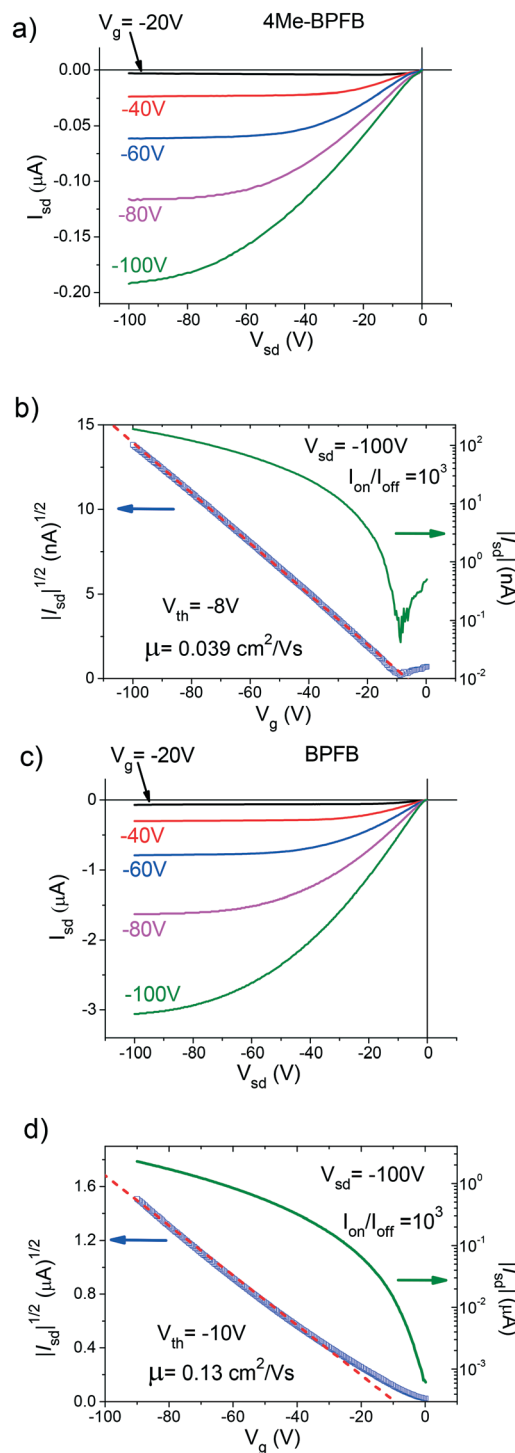


Fig. 5 a) Output and (b) transfer characteristics of a 4Me-BPFB single crystal OFET; c) output and d) transfer characteristics of a BPFB single crystal OFET; V_{sd} , V_g and V_{th} are the source-drain, gate and threshold voltages, respectively; μ is the field-effect mobility extracted in the saturation regime.

Table 2 Optical properties of BPFB, 2Me-BPFB, and 4Me-BPFB; λ_{abs} and λ_{em} are the maxima of absorption and PL spectra, respectively

	Solution				Single crystal
	λ_{abs} , nm	λ_{em} , nm	PL QY, %	λ_{em} , nm	PL QY, %
BPFB	376	410, 434	91	498	45
2Me-BPFB	379	413, 438	86	502	76
4Me-BPFB	379	413, 438	89	495, 521	80

(Fig. 3). This inclination in methyl-substituted BPFBs results in an increase in intermolecular distances as well as a decrease in the molecular stacking angle along the charge transport direction (Table S1†). Both lower the overlapping of



electronic wavefunctions and hence should reduce the charge mobility.

Note that in contrast to previous studies of bulky substituent effect on charge transport,^{33,37} the charge mobility of substituted BPFB single crystals does not correlate with the steric effect of the substituents. Indeed, the relatively small methyl group at the *p*-position of the terminal phenyl in 2Me-BPFB results in complete loss of mobility, whereas the sterically hindered 4Me-BPFB still demonstrates the field-effect in devices.

4. Conclusions

In summary, we have demonstrated a series of highly-emissive single crystals based on methyl-substituted furan/phenylene co-oligomers. The introduction of terminal methyl groups lowers the solubility, enhances the thermal stability and results in high inclination of the molecules against the basal crystal plane. Accordingly, the crystal packing of methyl-substituted derivatives corresponds to J-aggregation, which is beneficial for high luminescence. The tilt angles as low as 28 and 36° for 2-Me and 4Me-BPFB, respectively, are among the lowest for thiophene/ and furan/phenylene co-oligomers and very unusual for linear conjugated molecules. We have shown that the charge mobility in the crystals decreases upon the inclination of the molecules. Therefore, the crystal packing of heteroaryl-containing linear conjugated oligomers can be effectively tuned by methyl substituents, and this is a useful approach for obtaining desired intermolecular distances and interactions without distortion of the molecular electronic structure.

Acknowledgements

This work was supported by the Russian Foundation for Basic Research project 16-33-60011 mol_a_dk. The work on PL in crystals was supported by Russian Science Foundation (project 15-12-30031). This work was done using equipment purchased under the Lomonosov Moscow State University Program of Development. The authors acknowledge Novosibirsk State University program “5-100” and Multi-Access Chemical Service Center SB RAS for the spectral and analytical measurements.

Notes and references

- J. Gierschner and S. Y. Park, *J. Mater. Chem. C*, 2013, **1**, 5818–5832.
- H.-H. Fang, J. Yang, J. Feng, T. Yamao, S. Hotta and H.-B. Sun, *Laser Photonics Rev.*, 2014, **8**, 687–715.
- S. Hotta, T. Yamao, S. Z. Bisri, T. Takenobu and Y. Iwasa, *J. Mater. Chem. C*, 2014, **2**, 965–980.
- J. Gierschner, S. Varghese and S. Y. Park, *Adv. Opt. Mater.*, 2016, **2**, 348–364.
- S. Varghese and S. Das, *J. Phys. Chem. Lett.*, 2011, **2**, 863–873.
- S. Hotta and T. Yamao, *J. Mater. Chem.*, 2011, **21**, 1295–1304.
- T. Komori, H. Nakanotani, T. Yasuda and C. Adachi, *J. Mater. Chem. C*, 2014, **2**, 4918–4921.
- L. G. Kudryashova, M. S. Kazantsev, V. A. Postnikov, V. V. Bruevich, Y. N. Luponosov, N. M. Surin, O. V. Borshchev, S. A. Ponomarenko, M. S. Pshenichnikov and D. Y. Paraschuk, *ACS Appl. Mater. Interfaces*, 2016, **8**, 10088–10092.
- R. Hirase, M. Ishihara, T. Katagiri, Y. Tanaka, H. Yanagi and S. Hotta, *Org. Electron.*, 2014, **15**, 1481–1492.
- V. A. Postnikov, Y. I. Odarchenko, A. V. Iovlev, V. V. Bruevich, A. Y. Pereverzev, L. G. Kudryashova, V. V. Sobornov, L. Vidal, D. Chernyshov, Y. N. Luponosov, O. V. Borshchev, N. M. Surin, S. A. Ponomarenko, D. A. Ivanov and D. Y. Paraschuk, *Cryst. Growth Des.*, 2014, **14**, 1726–1737.
- R. Kabe, H. Nakanotani, T. Sakanoue, M. Yahiro and C. Adachi, *Adv. Mater.*, 2009, **21**, 4034–4038.
- S. Varghese, S. K. Park, S. Casado, R. C. Fischer, R. Resel, B. Milian-Medina, R. Wannemacher, S. Y. Park and J. Gierschner, *J. Phys. Chem. Lett.*, 2013, **4**, 1597–1602.
- J. Gierschner, M. Ehni, H. J. Egelhaaf, B. M. Medina, D. Beljonne, H. Benmansour and G. C. Bazan, *J. Chem. Phys.*, 2005, **123**, 144914.
- G. P. Bartholomew, G. C. Bazan, X. Bu and R. J. Lachicotte, *Chem. Mater.*, 2000, **12**, 1422–1430.
- S. K. Park, J. H. Kim, S.-J. Yoon, O. K. Kwon, B.-K. An and S. Y. Park, *Chem. Mater.*, 2012, **24**, 3263–3268.
- M. L. Renak, G. P. Bartholomew, S. Wang, P. J. Ricatto, R. J. Lachicotte and G. C. Bazan, *J. Am. Chem. Soc.*, 1999, **121**, 7787–7799.
- S. Varghese, S. K. Park, S. Casado, R. Resel, R. Wannemacher, L. Lüer, S. Y. Park and J. Gierschner, *Adv. Funct. Mater.*, 2016, **26**, 2349–2356.
- M. S. Kazantsev, E. S. Frantseva, L. G. Kudryashova, V. G. Konstantinov, A. A. Mannanov, T. V. Rybalova, E. V. Karpova, I. K. Shundrina, G. N. Kamaev, M. S. Pshenichnikov, E. A. Mostovich and D. Y. Paraschuk, *RSC Adv.*, 2016, **6**, 92325–92329.
- A. M. Brouwer, *Pure Appl. Chem.*, 2011, **83**, 2213–2228.
- R. W. I. De Boer, M. E. Gershenson, A. F. Morpurgo and V. Podzorov, *Phys. Status Solidi*, 2004, **201**, 1302–1331.
- SAINT, *Data Reduction and Frame Integration Program for the CCD Area-Detector System*, Bruker Analytical X-ray Systems, Madison, Wisconsin, USA, 1997–2006.
- G. M. Sheldrick, *SADABS, Program for area detector adsorption correction*, Institute for Inorganic Chemistry, University of Goettingen, Goettingen, Germany, 1996.
- G. M. Sheldrick, *Acta Crystallogr., Sect. A: Found. Crystallogr.*, 2008, **64**, 112–122.
- A. L. Spek, *PLATON, A Multipurpose Crystallographic Tool, version 10M*, Utrecht University, Netherlands, 2003.
- A. L. Spek, *J. Appl. Crystallogr.*, 2003, **36**, 7–13.
- V. Podzorov, V. M. Pudalov and M. E. Gershenson, *Appl. Phys. Lett.*, 2003, **82**, 1739–1741.
- V. Podzorov, *MRS Bull.*, 2013, **38**, 15–24.
- Y. Yomogida, T. Takenobu, H. Shimotani, K. Sawabe, S. Z. Bisri, T. Yamao, S. Hotta and Y. Iwasa, *Appl. Phys. Lett.*, 2010, **97**, 173301.



- 29 K. Oniwa, T. Kanagasekaran, T. Jin, M. Akhtaruzzaman, Y. Yamamoto, H. Tamura, I. Hamada, H. Shimotani, N. Asao, S. Ikeda and K. Tanigaki, *J. Mater. Chem. C*, 2013, **1**, 4163–4170.
- 30 S. Hotta, Y. Shimizu, T. Yamao, M. Goto and R. Azumi, *Chem. Lett.*, 2009, **38**, 294–295.
- 31 T. Nicolini, A. Famulari, T. Gatti, J. Martí-Rujas, F. Villafiorita-Monteleone, E. V. Canesi, F. Meinardi, C. Botta, E. Parisini, S. Valdo Meille and C. Bertarelli, *J. Phys. Chem. Lett.*, 2014, **5**, 2171–2176.
- 32 H. Mizuno, T. Maeda, H. Yanagi, H. Katsuki, M. Aresti, F. Quochi, M. Saba, A. Mura, G. Bongiovanni, F. Sasaki and S. Hotta, *Adv. Opt. Mater.*, 2014, **2**, 529–534.
- 33 G. Schweicher, V. Lemaire, C. Niebel, C. Ruzié, Y. Diao, O. Goto, W. Lee, Y. Kim, J.-B. Arlin, J. Karpinska, A. R. Kennedy, S. R. Parkin, Y. Olivier, S. C. B. Mannsfeld, J. Cornil, Y. H. Geerts and Z. Bao, *Adv. Mater.*, 2015, **27**, 3066–3072.
- 34 S.-J. Yoon, J. W. Chung, J. Gierschner, K. S. Kim, M.-G. Choi, D. Kim and S. Y. Park, *J. Am. Chem. Soc.*, 2010, **132**, 13675–13683.
- 35 E. G. Mcrae and M. Kasha, *J. Chem. Phys.*, 1958, **28**, 721–722.
- 36 R. Katoh, K. Suzuki, A. Furube, M. Kotani and K. Tokumaru, *J. Phys. Chem. C*, 2009, **113**, 2961–2965.
- 37 C. Reese, M. E. Roberts, S. R. Parkin and Z. Bao, *Adv. Mater.*, 2009, **21**, 3678–3681.

