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Cite this: *RSC Adv.*, 2016, 6, 62218

## Highly stable supercapacitive performance of one-dimensional (1D) brookite TiO<sub>2</sub> nanoneedles†

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We report the highly stable supercapacitive performance of one-dimensional (1D) nanoneedles of brookite (β) TiO<sub>2</sub> synthesized on a conducting glass substrate. The 1D β-TiO<sub>2</sub> nanoneedles synthesized over a large area array utilizing hot-filament metal vapor deposition (HFMD) were ~24–26 nm wide, ~650 nm long and tapered in a downward direction. X-ray photoemission spectroscopy (XPS) revealed their chemical properties and stoichiometric Ti and O composition. The 1D β-TiO<sub>2</sub> nanoneedles execute as parallel units for charge storage, yielding a specific capacitance of 34.1 mF g<sup>-1</sup>. Electrochemical impedance spectroscopy revealed that the large surface area and brookite crystalline nature of the 1D nanoneedles provided easy access to Na<sup>+</sup> ions, and resulted in low diffusion resistance, playing a key role in their stable charging–discharging electrochemical mechanism. Moreover, the non-faradic mechanism of these nanoneedles delivered better durability and high stability up to 10 000 cycles, and a columbic efficiency of 98%. Therefore, 1D β-TiO<sub>2</sub> nanoneedles hold potential as an electrode material for highly stable supercapacitive performance with long cycle lifetime.

Received 2nd May 2016  
 Accepted 20th June 2016

DOI: 10.1039/c6ra11348f

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## Introduction

Supercapacitors, known as electric double-layer capacitors, store charges at the interface of an electrolyte and an electrode of choice through reversible electrochemical diffusion of ions. With the increase in demand for high-performance energy storage technologies, supercapacitor electrode materials with fast recharging ability, high cycling stability, reversibility and long cycle lifetime are of importance. At present, supercapacitors provide already higher power density than batteries, and energy density than conventional capacitors. However, their stability and cycle lifetime are inferior to those of batteries and conventional capacitors. Carbon-based materials and conducting polymers known for high specific capacitance values and high power densities suffer from high energy densities even after achieving larger specific surface areas and rational porous morphology.<sup>1</sup> Moreover, their swelling and shrinking during charging and discharging induces stress and results in short lifetime and faster degradation. Therefore, when it comes to the stability and long cycle lifetime, transition metal oxides seem to be a better alternative. Although the life cycle of transition metal

oxides is limited from a few hundreds to a few thousand,<sup>2–4</sup> efforts have been made for the improvement by reducing the structural dimensions of synthesized electrode materials. The nanoscale dimensions with well-defined geometry and colossal surface area are expected to improve the supercapacitive performance of the metal-oxides. Among all available morphologies of transition metal oxides, one-dimensional (1D) nanostructures plays an important role in the enhancement of charge storing ability, stability and cycle lifetime by providing large surface area and therefore more capillary pathways for ion diffusion/transport.<sup>5</sup>

Over the past few decades, titanium dioxide (TiO<sub>2</sub>) is one of the most fascinating functional materials in the fabrication of supercapacitor because of the low processing cost, high chemical stability compared to most alternative materials, abundance and non-toxic nature. In the form of ultra-thin films or nanoparticles of size below 10 nm, mesoporous TiO<sub>2</sub> films dramatically increase the ion transport and provide the high specific capacitance of 90–120 μF cm<sup>-2</sup> at the normalized surface area.<sup>6</sup> Highly ordered arrays of TiO<sub>2</sub> nanotubes provide large surface area were found interesting alternative morphology for charge storage.<sup>3,7,8</sup> However, the discontinuous and polycrystalline morphology of TiO<sub>2</sub> nanotubes offers relatively poor conductivity and low electrochemical activity, limit its use in high-performance supercapacitors. Therefore, significant efforts have been made to synthesize TiO<sub>2</sub> nanorods, nanowires, nanotubes and nanofibers, *etc.* and improve their supercapacitive performance in combination with various electrolytes such as H<sub>2</sub>SO<sub>4</sub>, NaOH, Na<sub>2</sub>SO<sub>4</sub>, LiClO<sub>4</sub>, and LiPF<sub>6</sub>, *etc.*<sup>3,5,7–11</sup> Mainly three different polymorphs, rutile, anatase and

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† Electronic supplementary information (ESI) available. See DOI: 10.1039/c6ra11348f



brookite ( $\beta$ ) are described regarding crystal structural arrangements of distorted  $\text{TiO}_6$  octahedra. Chains of edge-shared distorted octahedra running parallel to the  $c$ -axis yielding in  $\beta$ -phase of 1D nanoneedles. The shared edges of distorted  $\text{TiO}_6$  octahedra provide abundant enough vacant sites to accommodate  $\text{Na}^+$  ions. Therefore, especially  $\beta$ - $\text{TiO}_2$  promises a electrode material for electrochemical processes. In nanocrystalline form, the  $\beta$ -phase is thermodynamically most stable with dimensions between 11 and 35 nm,<sup>12</sup> and crystallographic data in brookite are intermediate between rutile and anatase phases.<sup>13</sup> Moreover,  $\beta$ - $\text{TiO}_2$  is promising dielectrics with much larger static dielectric constant<sup>14</sup> and also known to exhibit markedly higher catalytic activities<sup>12</sup> than anatase and rutile  $\text{TiO}_2$ . Therefore, the  $\beta$ -phase is expected to offer high stability and long cycle lifetime.

So far, supercapacitor properties of anatase and rutile structures are studied extensively<sup>15–17</sup> and seem to have reached their limit. Because of the difficulties encountered in the synthesis, most of the times  $\beta$ -phases was accompanied with anatase or/and rutile phases.<sup>12,18,19</sup> Therefore, high temperature calcinations<sup>19</sup> (700 °C) and annealing<sup>18</sup> (800 °C) treatment was employed to obtain  $\beta$ - $\text{TiO}_2$  nanorods. Nevertheless, the properties of the nanostructure of pure  $\beta$ - $\text{TiO}_2$  phase seem to be interesting. To the best of our knowledge, supercapacitor properties of pure 1D  $\beta$ - $\text{TiO}_2$  nanoneedles have not been reported yet in the literature. Therefore, in this paper, we report stable supercapacitive performance of 1D  $\beta$ - $\text{TiO}_2$  nanoneedles synthesized in a large area *via* hot-filament metal vapor deposition (HFMDV) technique without additional heat treatment. The HFMDV technique is unique and useful for the synthesis of the variety of metal oxide nanostructures with diverse morphologies and crystalline phases.<sup>20–24</sup> The structural morphology and chemical composition of the 1D  $\beta$ - $\text{TiO}_2$  nanoneedles was examined utilizing field-emission scanning electron microscopy (FESEM) and X-ray photoemission spectroscopy (XPS), respectively. The electrochemical characterizations of the large-area array of the 1D  $\beta$ - $\text{TiO}_2$  nanoneedles were carried out by cyclic voltammetry (CV), galvanostatic charge/discharge method, and AC impedance analysis in the electrolyte of  $\text{Na}_2\text{SO}_4$ . The results strongly indicated that 1D  $\beta$ - $\text{TiO}_2$  nanoneedles provide columbic efficiency of 98%, high stability, greater retentivity ( $\sim 88.2\%$ ) and long cycle lifetimes up to 10 000 cycles in  $\text{Na}_2\text{SO}_4$  electrolyte. Therefore, the large area array of 1D  $\beta$ - $\text{TiO}_2$  nanoneedles is an attractive material for fabrication of highly stable supercapacitors of long cycle lifetime.

## Experimental

Large area arrays of  $\text{TiO}_2$  nanoneedles were synthesized utilizing hot-filament metal vapor deposition (HFMDV) technique. A clean titanium (Ti) wire (99.9% pure) with a diameter of 1 mm was passed through a pure graphite disc fixed on two supporting Cu electrodes mounted in a vacuum chamber. Once the pressure of the vacuum chamber was pumped down to  $1 \times 10^{-2}$  Torr, the Ti wire was heated to  $\sim 1150$  °C for 15 min to generate hot titanium vapor. The hot titanium vapor reacted

with the residual oxygen (or leaking air) to form the metal oxide vapor of  $\text{TiO}_x$  ( $x \leq 2$ ). When the  $\text{TiO}_x$  vapor imparted the rather cold glass substrate coated with a conducting ITO thin film, it condensed into one-dimensional (1D)  $\text{TiO}_2$  nanoneedles. The surface morphology of the as-synthesized large-area array of  $\text{TiO}_2$  nanoneedles was characterized using field-emission scanning electron microscope (FESEM, JEOL JSM-6500F). The electronic structure and chemical state of these 1D  $\text{TiO}_2$  nanoneedles were analyzed using an X-ray photoelectron spectrometer (XPS, Thermo Scientific Inc. K-alpha) with a microfocus monochromated Al  $K\alpha$  X-ray. The 1D  $\beta$ - $\text{TiO}_2$  nanoneedles based supercapacitor device constituted of the transparent 1D  $\beta$ - $\text{TiO}_2$  nanoneedle array on a conducting ITO coated glass substrate as working electrode, a platinum foil as the counter electrode and saturated calomel electrode (SCE) stored in saturated KCl as the reference electrode was used to corroborate electrochemical properties. All the electrochemical measurements were performed with a potentiostat (Bio-Logic, SP-150) in 1 M  $\text{Na}_2\text{SO}_4$  aqueous electrolyte at room temperature to investigate the potential of large area arrays of 1D  $\text{TiO}_2$  nanoneedles as a highly stable electrochemical supercapacitor. The cycling stability and performance of the device was confirmed with two terminal measurement method.

## Results and discussion

The FESEM images in Fig. 1 show the surface morphology of the large-area array of 1D  $\text{TiO}_2$  nanoneedles. In Fig. 1(a) and (b) a section of a large area array of 1D  $\text{TiO}_2$  nanoneedles is shown in the top and side view, respectively. The self-assembled 1D  $\text{TiO}_2$  nanoneedles are conical in shape and displays various diameters at the tip, body, and top. The array contains roughly 425 nanoneedles per square micrometer. The textural boundaries of these randomly dispersed  $\text{TiO}_2$  nanoneedles are clearly visible throughout the needle body. The high-magnification FESEM image in the inset of Fig. 1(a) shows single  $\text{TiO}_2$  nanoneedles. The conical shaped 1D  $\text{TiO}_2$  nanoneedles are on average  $\sim 650$  nm long and tapered in the downward direction, forming a film with an average thickness of  $\sim 470$  nm (Fig. 1(b)). The diameter of the nanoneedles is thinnest at the bottom (smaller than 12 nm) and  $\sim 50$  nm at the top for the thickest nanoneedles. However, a statistical histogram of the diameter at the top of the 1D  $\text{TiO}_2$  nanoneedles (Fig. 1(c)) clearly shows that the diameter at the top of most of the  $\text{TiO}_2$  nanoneedles falls in the range of 24 to 26 nm. The variation of the diameters of the  $\text{TiO}_2$  nanoneedles can be fitted with a log-normal distribution function as follows,

$$f_N(w) = \frac{A}{d\sigma\sqrt{2\pi}} \exp\left[-\frac{\ln(d/\bar{d})^2}{2\sigma^2}\right] \quad (1)$$

where,  $d$  is the diameter at the top of the nanoneedle,  $\bar{d}$  [ $\equiv 26.82 \pm 0.52$  nm] is the mean (average) diameter of the various nanoneedles,  $A$  [ $\equiv 111.41 \pm 10.53$ ] is the initial constant, and  $\sigma$  [ $\equiv 0.19 \pm 0.03$ ] is the standard deviation of the diameters of the nanoneedles. The log-normal distribution of the nanoneedles diameters is asymmetrical. The small standard deviation of the



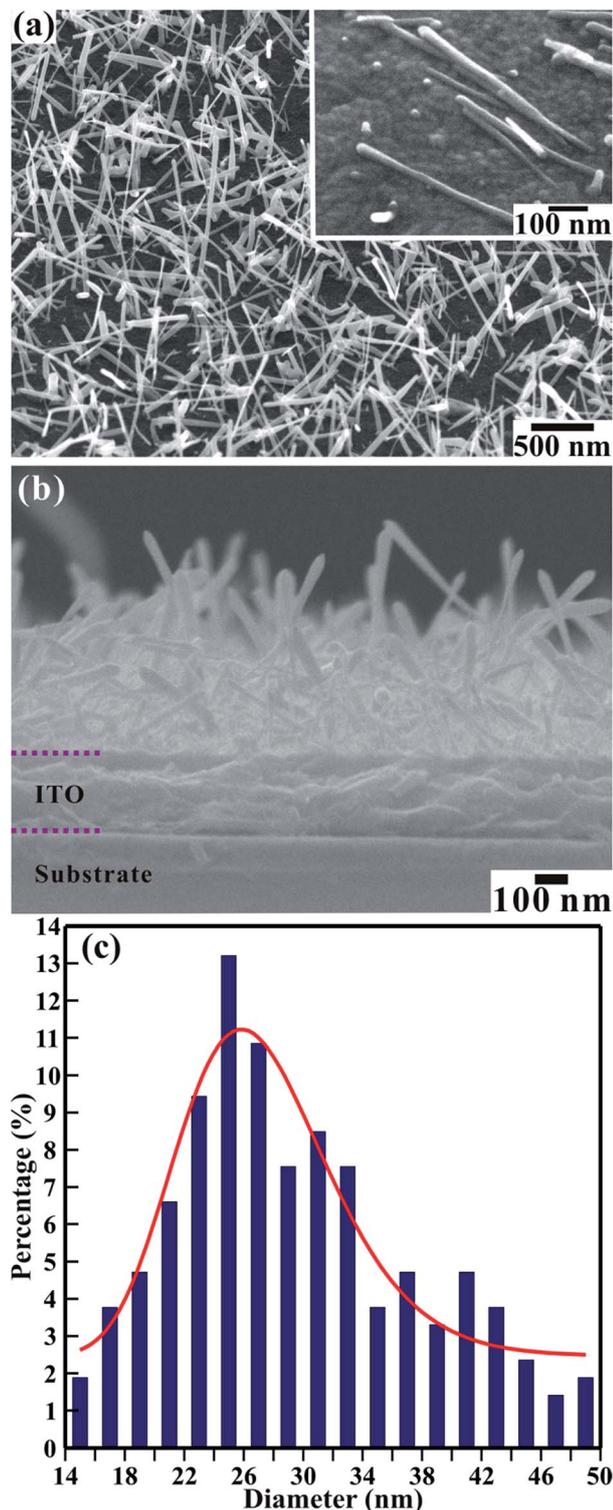


Fig. 1 FESEM image showing the (a) top and (b) side views of the large-area array of 1D TiO<sub>2</sub> nanoneedles on conducting glass substrate. Inset of (a) shows high-magnification FESEM image of the single 1D TiO<sub>2</sub> nanoneedles separated from the array. (c) Statistical histogram of the distribution of diameter at the top of the TiO<sub>2</sub> nanoneedles, fitted by a log-normal distribution function.

diameter distribution ( $\sigma \leq 0.25$ ) indicates that the nanoneedles are well confined to a limited diameter range.

Further, the crystalline structure of these well-defined TiO<sub>2</sub> nanoneedles was confirmed by XRD and discussed elsewhere.<sup>25</sup> The corresponding XRD spectra evidences that the 1D TiO<sub>2</sub> nanoneedles are exclusively of orthorhombic crystal structure in brookite phase, which is assigned to the space group *Pbca* (JCPDS-761936) with lattice constants  $a = 0.919$  nm,  $b = 0.546$  nm,  $c = 0.516$  nm and  $\alpha = \beta = \gamma = 90^\circ$ . Altogether, it demonstrates that our HFMVD technique can be successfully employed for the synthesis of 1D TiO<sub>2</sub> nanostructures with optimal dimensions and required thickness. According to our literature survey, rutile and anatase phases are commonly synthesized and widely studied TiO<sub>2</sub> phases. Only brookite is investigated rarely, even though it has exciting properties. Therefore, the large area array of 1D TiO<sub>2</sub> nanoneedles of brookite phase reported here is of special importance in scientific and technological application point of view.

XPS was used to perform quantitative analysis of the electronic structure and chemical properties of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles. Fig. 2 illustrate two Ti(2p) and O(1s) XPS spectra obtained for the large area array of 1D TiO<sub>2</sub> nanoneedles. Fig. 2(a) shows the double peak features in the Ti(2p) XPS spectrum. To precisely determine the features of the double peaks of Ti(2p<sub>3/2</sub>) and Ti(2p<sub>1/2</sub>), the Ti(2p) XPS spectra was decomposed *via* Voigt curve fitting with the Shirley background. The deconvolution shows a perfect fit for two peaks located at binding energies of 458.56 and 464.24 eV, respectively, corresponding to the Ti(2p<sub>3/2</sub>) and Ti(2p<sub>1/2</sub>) core levels of Ti<sup>4+</sup> cations and not of Ti<sup>3+</sup>.<sup>17,26</sup> The energy separation between Ti(2p<sub>3/2</sub>) and Ti(2p<sub>1/2</sub>) peaks is 5.68 eV, and their area ratio is 2.48, which reflects a strong bonding between the Ti and O atoms. The full width at half-maximum (FWHM) of the Ti(2p<sub>3/2</sub>) and Ti(2p<sub>1/2</sub>) peak are 1.29, 2.27, respectively, indicating the high resolution of the Ti(2p) XPS spectrum in comparison with previous studies.<sup>26,27</sup> Similarly, the O(1s) XPS spectrum of the 1D nanoneedles shown in Fig. 2(b) was decomposed *via* Voigt curve fitting with the Shirley background. The results demonstrate a perfect fit to two peaks located at 529.89 and 531.31 eV, with FWHM's of 1.38 and 2.19 eV, respectively. The lower binding energy peak at 530.14 eV corresponds to the O(1s) core level of the O<sup>2-</sup> anions in the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles. However, the higher binding peak at 531.61 eV is ascribed to surface contamination, such as carbon oxides or hydroxides<sup>20,21</sup> of the 1D nanoneedles. The O(1s) peak observed at a binding energy of 529.89 eV is associated with the Ti-O chemical bonding.<sup>17</sup> The atomic ratio of titanium and oxygen (*i.e.* Ti/O ratio) estimated by integrating the area beneath the decomposed peaks of O(1s) (530.14 eV) and Ti(2p<sub>3/2</sub>) (458.74 eV) is  $\sim 0.51$  (*i.e.* Ti : O = 1.02 : 2), which is very close to the stoichiometric ratio (*i.e.* 1 : 2) of pure TiO<sub>2</sub>. This analysis confirms that all 1D nanoneedles in the large area array were fully oxidized, and were composed of pure stoichiometric TiO<sub>2</sub> without any traces of sub-oxides (TiO<sub>x</sub>). Moreover, the binding energy difference ( $\Delta E$ ) of 71.33 eV between the O(1s) and Ti(2p<sub>3/2</sub>) peaks is very close to that of 71.5 eV for TiO<sub>2</sub>, and significantly smaller than 73.4 eV for Ti<sub>2</sub>O<sub>3</sub> and 75.0 eV for TiO.<sup>28</sup> This confirms again that the 1D



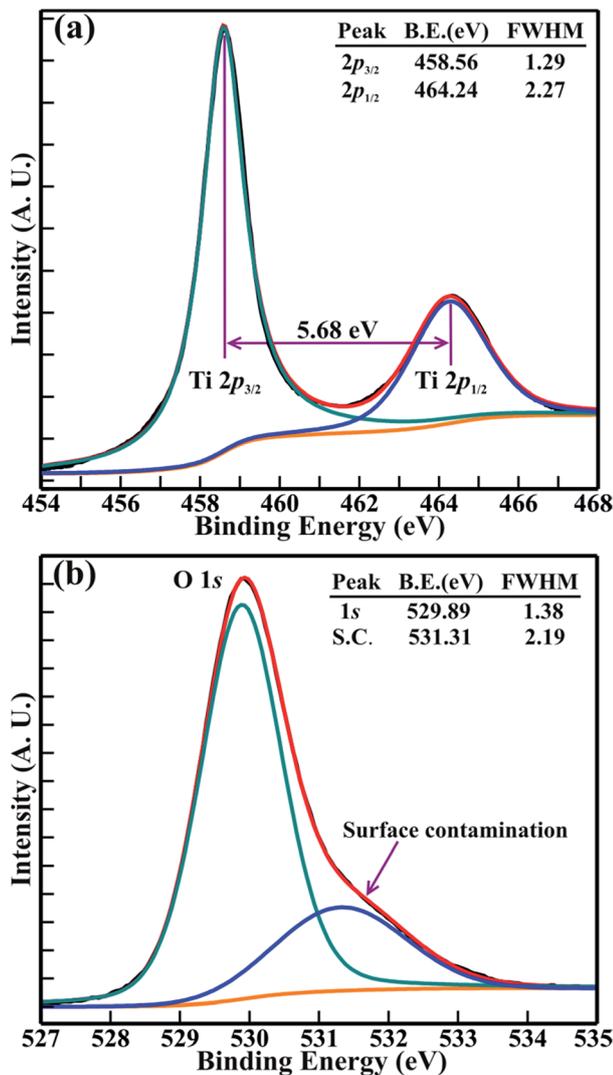


Fig. 2 Typical high-resolution XPS spectra of the (a) Ti(2p) and (b) O(1s) core levels of the large area array of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles. The XPS spectra were decomposed via Voigt curve function fittings.

nanoneedle array over a large area is formed exclusively of stoichiometric TiO<sub>2</sub>.

The supercapacitor behavior of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles is confirmed with CV, galvanostatic charge/discharge and impedance measurements. The CV measurements useful to investigate supercapacitor behavior are recorded for 1D  $\beta$ -TiO<sub>2</sub> nanoneedles within the operating potential range of 0 to 0.8 V and at various scan rates ranging from 15 to 150 mV s<sup>-1</sup>. Fig. 3(a) shows the CV graphs of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles. Irrespective of variation in the scan rates, the CV graph shows well rectangular shape in 1 M Na<sub>2</sub>SO<sub>4</sub> electrolyte, indicating good capacitive behavior and high rate capability of the  $\beta$ -TiO<sub>2</sub> nanoneedles. No evidence of any faradic reaction on well rectangular CV curve is found to ensure the reduction of Ti<sup>4+</sup> to Ti<sup>3+</sup>. This confirms the pure double layer capacitor behavior of the  $\beta$ -TiO<sub>2</sub> nanoneedles. Moreover, the well rectangular shape of CV graph confirms that the large surface area of the nanoneedle

morphology of  $\beta$ -TiO<sub>2</sub> with clearly visible textural boundaries offers abundant diffusion of Na<sup>+</sup> ions and charge transport during non-faradic reaction along the nanoneedle sides. The negative and positive current density occurred due to the insertion and extraction of Na<sup>+</sup> ions on the surface of  $\beta$ -TiO<sub>2</sub> nanoneedles. A little bit more negative current density is observed than that of positive current density at applied potential. It is accepted that the large surface area of nanostructure morphology acts as traps to capture free electrons or ions when they pass through. Likewise, the large surface area of 1D TiO<sub>2</sub> nanoneedles possibly provides a number of trap levels. Therefore, trapping–detrapping of free electrons at trap levels enhances the ionic conductivity<sup>17</sup> and results in a more negative current density at applied potential. The current density of CV increased with increase in scanning rate from 15 to 150 mV s<sup>-1</sup>. Likewise, the area under CV graph increased proportionally with the scan rate. The variation in current density can be well understood via the Randles–Sevcik equation–

$$i_p = 2.69 \times 10^5 n^{3/2} C_0 \nu^{1/2} D^{1/2} \quad (2)$$

where,  $i_p$  is the current density (mA g<sup>-2</sup>),  $n$  is the number of electrons,  $C_0$  is the concentration of Na<sup>+</sup> ions in the electrolyte (mol cm<sup>-2</sup>),  $\nu$  is the scanning rate (mV s<sup>-1</sup>), and  $D$  is the diffusion coefficient for Na<sup>+</sup> ion in the electrolyte. This equation depicts that the current density ( $i_p$ ) of the electrochemical reaction mechanism is proportional to the concentration of Na<sup>+</sup> ions ( $C_0$ ) and the square root of the scan rate ( $\nu$ ). The relation of the peak current density ( $i_p$ ) of the anodic reduction with the square root of the scan rate ( $\nu$ ) is presented graphically in the Fig. 3(b). The linear behavior observed between maximum current density ( $i_p$ ) and scan rate ( $\nu$ ) shown by blue line passes through the origin (extrapolated dotted blue line) and confirms that the process associated with the Na<sup>+</sup> ion insertion is pure diffusion controlled.<sup>29</sup>

The specific capacitance is shown in the Fig. 3(c) is calculated from CV graph at various scan rates (Fig. 3(a)) from the equation,

$$C_s = \frac{\int idV}{2ms\Delta V} \quad (3)$$

where,  $C_s$  is the specific capacitance of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles (F g<sup>-1</sup>),  $m$  is the mass of active material (g),  $s$  is the scan rate (V s<sup>-1</sup>),  $\Delta V$  is the potential range (V), and  $\int idV$  is an integrated area of the CV curve. The calculated specific capacitance ( $C_s$ ) varied linearly with the scan rate. The  $C_s$  value of 34.1 mF g<sup>-1</sup> gained at the scan rate of 15 mV s<sup>-1</sup> was retained to 31.6 mF g<sup>-1</sup> after the scan rate of 150 mV s<sup>-1</sup>. The  $C_s$  retention of 92.7% was observed even after large variation in the scan rate. The capacitance of 34.1 mF g<sup>-1</sup> might have resulted from maximum diffusion of Na<sup>+</sup> ions over the large surface area provided by 1D nanoneedle array. The  $C_s$  value of 34.1 mF g<sup>-1</sup> gained from 1D  $\beta$ -TiO<sub>2</sub> is substantially larger than reported for TiO<sub>2</sub> nanostructures and thin films such as amorphous and air annealed TiO<sub>2</sub> nanotube arrays,<sup>7</sup> polycrystalline anatase TiO<sub>2</sub>



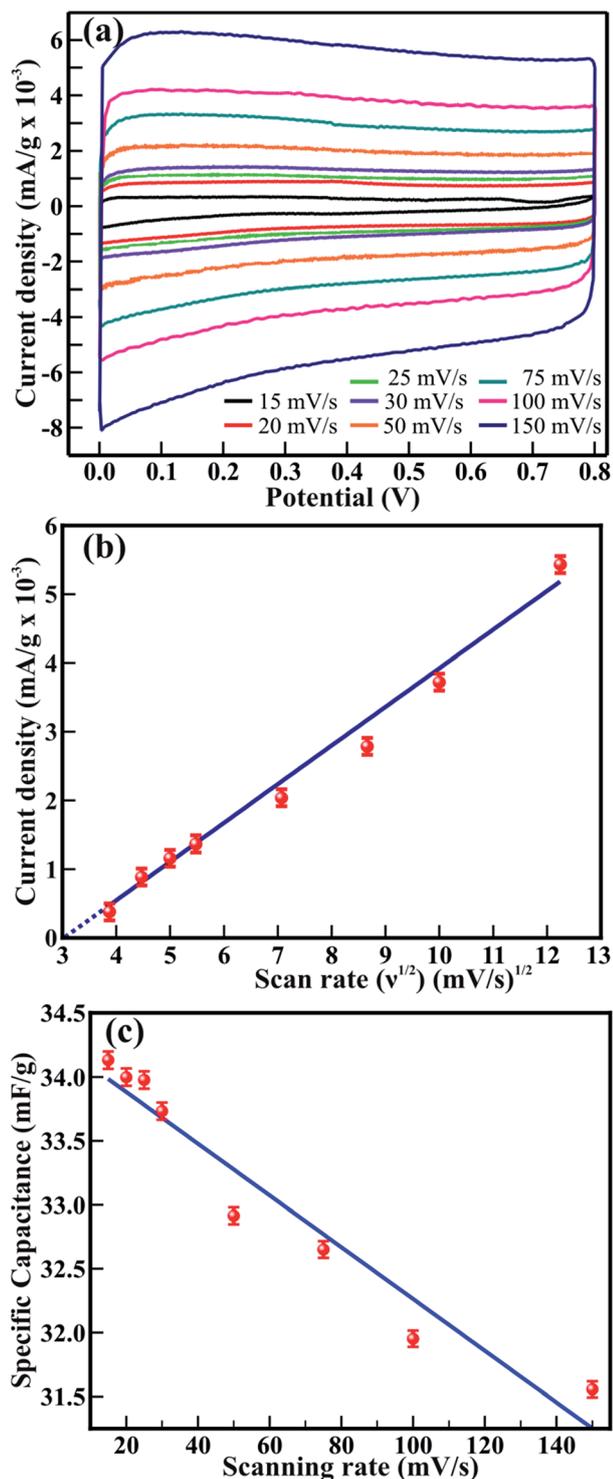


Fig. 3 (a) Steady-state cyclic voltammograms of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles obtained in the voltage range of 0 to 0.8 V at various scan rates from 15 to 150 mV s<sup>-1</sup> in 1 M Na<sub>2</sub>SO<sub>4</sub> electrolyte. (b) The graphical representation of the peak current density ( $i_p$ ) of the anodic reduction with the square root of the scan rate ( $\nu$ ). (c) The specific capacitance calculated from those cyclic voltammograms for various scan rates.

nanowires,<sup>30</sup> anatase/rutile mixed phase TiO<sub>2</sub> nanocrystals,<sup>31</sup> and rutile TiO<sub>2</sub> nanorods.<sup>9</sup> Although Salari *et al.*<sup>32</sup> have reported enhancement of the  $C_s$  in anatase to rutile transformed TiO<sub>2</sub>

nanotubes *via* controlling oxygen vacancies, the very slow scan rate of 1 mV s<sup>-1</sup> was necessary to achieve these values. Furthermore, 7.3% capacitance reduction of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles observed with an increased scan rate from 15 to 150 mV s<sup>-1</sup>, is a much smaller loss than the significant reduction reported in anodically oxidized self-organized TiO<sub>2</sub> nanotubes (80.8% at 100 mV s<sup>-1</sup>),<sup>27</sup> TiO<sub>2</sub> nanowires and nanocrystalline powder (40.9 and 81.8% at 100 mV s<sup>-1</sup>, respectively),<sup>31</sup> and seed-assisted hydrothermally grown rutile TiO<sub>2</sub> nanowires (65% at 100 mV s<sup>-1</sup>).<sup>33</sup> Moreover, TiO<sub>2</sub> thin film (~23.7% at 100 mV s<sup>-1</sup>) showed a reduction of ~35.8% after preparing a multilayer film with graphene.<sup>34</sup> The reason for this smaller reduction (*i.e.* 7.3%) in  $C_s$  of 1D  $\beta$ -TiO<sub>2</sub> is instant diffusion/transport of ion along the textural boundaries of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles irrespective of scanning rate. Therefore, it is assumed that nanosize dimensions and increased textural boundaries of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles play a significant role in the diffusion of ion leading to the observed improvement.

The galvanostatic charging–discharging (Fig. S1, ESI<sup>†</sup>) of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles was studied at various current densities of 166.7, 250, 333.3 and 416.7  $\mu$ A g<sup>-1</sup>. Obviously, the charging curves were relatively symmetric to their discharge counterpart implying that a highly reversible ion transportation is efficiently taking place along the textural boundaries of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles. Furthermore, the overall performance of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles was illustrated with a Ragone plot (ESI, Fig. S2<sup>†</sup>). A Ragone plot manifests a high energy density and power density of 3.04 W h kg<sup>-1</sup>, and 1683 W kg<sup>-1</sup>, respectively.

Not only the amount of active material and electrolyte plays a critical role in the capacitive behavior, but a dominant role is attributed to the electrical properties, textural boundaries, and the core of the nanostructures. Therefore, electrochemical impedance spectroscopy (EIS) was employed to investigate the mechanistic aspects such as electrical resistance involved in the ion diffusion of 1D  $\beta$ -TiO<sub>2</sub> nanoneedle arrays. Typical Nyquist plots of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedle in 1 M Na<sub>2</sub>SO<sub>4</sub> solutions at various electrode potentials are presented in Fig. 4. All the impedance spectra consist of partial semicircles (arc) and straight lines having slopes at the higher and lower frequencies, respectively. In high-frequency region, the distorted semicircle corresponds to the charge-transfer resistance of the interface between the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles and electrolyte. The intercept of this semicircle yields the electrolyte resistance ( $R_e$ ), and the diameter provides the charge transfer resistance ( $R_{ct}$ ). As noticed from the diameter of the semi-circle, the value of the  $R_{ct}$  increased with the applied potential which reduces  $C_s$ . However, the low-frequency regime in the form of the straight line corresponds to diffusion resistance from the textural boundaries. Extrapolation of this gives larger resistance compare to that of the core of the nanoneedle since the 1D nanoneedle morphology of  $\beta$ -TiO<sub>2</sub> provides large surface area. These observations are within the expectation and are supporting the calculated values of  $C_s$ . Moreover, the straight line observed in low-frequency regime represents Warburg behavior. The angle made by the low-frequency data on the real axis decreases gradually from 67.4° to 50.03° on increasing the applied potential from 0.2 to 0.8 V. This indicates gradual



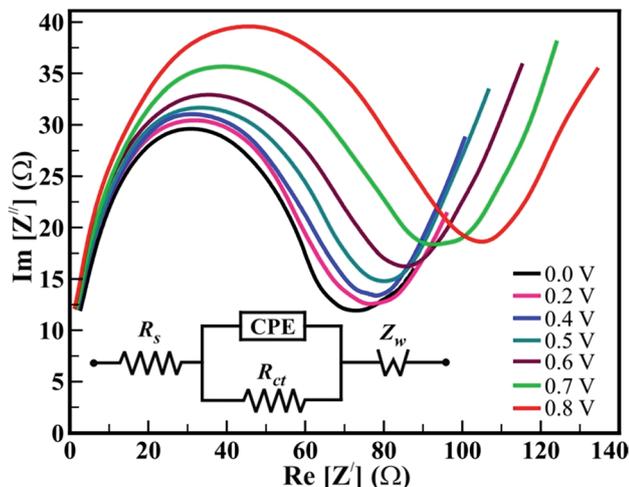


Fig. 4 Electrochemical impedance spectra (Nyquist plot) of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles acquired at various electrode potentials from 0 to 0.8 V. Inset shows best fitted equivalent circuit model for the impedance spectra of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles.

transition from capacitance to Warburg behavior between 0.2 and 0.8 V. The Warburg behavior (Warburg resistance  $Z_w$ ) results from the diffusion-controlled insertion and extractions of anions/cations in the electrode. This trend is in good agreement with the reports based on rutile TiO<sub>2</sub> nanowires and TiO<sub>2</sub>@C core-shell nanowires,<sup>35</sup> and anatase TiO<sub>2</sub> nanotubes.<sup>8</sup> The equivalent circuit for these electrochemical impedance measurements is composed of electrolyte resistance ( $R_e$ ), a charge transfer resistance ( $R_{ct}$ ), finite length Warburg diffusion element ( $Z_w$ ), and a constant phase element (CPE). It is represented in the inset of Fig. 4. Thus, nanoneedle morphology with large surface area and clearly visible textural boundaries improve the accessibility of the Na<sup>+</sup> ions from the electrolyte. The time constant ( $\tau$ ) is calculated from the equation-

$$\tau = \frac{1}{2\pi f^*} \quad (4)$$

where,  $f^*$  is the frequency corresponding to the maximum of the imaginary component ( $\text{Im}[Z'']$ ) of the semicircle. The calculated values of time constant ( $\tau$ ) are in the range of 3.26 to 4.85  $\mu\text{s}$ , which is fast, and clearly show that the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles provide easy access for transport of Na<sup>+</sup> ions or charges. This is an essential requirement for electrochemical supercapacitors to ensure fast charging-discharging mechanism.

Good cycling stability, durability and lifetime are important characteristics for the highly stable performance of supercapacitors. The 1D  $\beta$ -TiO<sub>2</sub> nanoneedles were tested at a scan rate of 100  $\text{mV s}^{-1}$  and the current density of 250  $\mu\text{A g}^{-1}$  for 10 000 cycles and 5000 cycles, respectively (Fig. 5). The  $C_s$  evaluated from the CV measurements at a scan rate of 100  $\text{mV s}^{-1}$  drops from 32.1 to 28.3  $\text{mF g}^{-1}$  after 10 000 cycles. Significantly, the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles exhibit long-term stability with only 11.8% reduction of  $C_s$  after 10 000 continuous cycles (Fig. 5(a)). The CV graphs of selected cycles (Fig. S3†) used to evaluate the  $C_s$  retention are almost identical, indicating

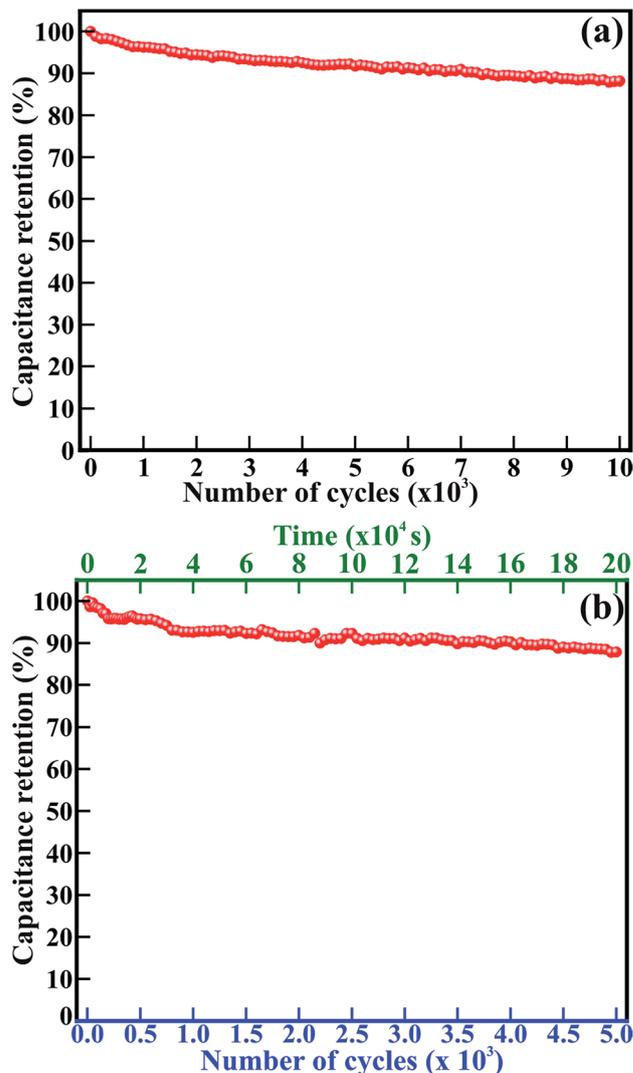


Fig. 5 Cycling stability and capacitance retention of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles observed by (a) cyclic voltammograms at a scan rate of 100  $\text{mV s}^{-1}$  for 10 000 cycles, and (b) galvanostatic charging-discharging at the current density of 250  $\mu\text{A g}^{-1}$  for 5000 cycles of 55.6 h. The percentage numbers denote the specific capacitance retention with the increase in a number of cycles.

extremely stable performances of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles. Furthermore,  $C_s$  measurements were performed for 5000 continuous cycles of galvanostatic charging-discharging at current density of 250  $\mu\text{A cm}^{-2}$  in 200 000 s (55.6 h) (Fig. 5(b)). Galvanostatic charging-discharging cycles were relatively symmetric to their discharge counterpart. First, 50 cycles are shown in Fig. S4.† With increased number of charging-discharging cycles and time, the calculated  $C_s$  drops from 109.9 to 96.7  $\text{mF g}^{-1}$  after 5000 cycles. This small relative decrease indicates that 1D  $\beta$ -TiO<sub>2</sub> nanoneedles exhibit long-term stability with only 12.04% reduction of  $C_s$  after continuous galvanostatic charging-discharging for 55.6 h (*i.e.* 5000 cycles). This  $\sim 12\%$  reduction of  $C_s$  of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles observed is much better than the value for amorphous TiO<sub>2</sub> nanotubes (34.8% reduction after 10 000 cycles) which were further annealed at 400  $^{\circ}\text{C}$  to form anatase TiO<sub>2</sub>



nanotubes (44.3% reduction after 10 000 cycles),<sup>7</sup> hydrothermally synthesized rutile TiO<sub>2</sub> nanorods (20% reduction even after only 1000 cycles),<sup>9</sup> sonochemically carbon activated TiO<sub>2</sub> ( $\beta$ ) nanowires mixed with polyaniline (57.5% reduction after 2000 cycles),<sup>4</sup> polythiophene infiltrated TiO<sub>2</sub> nanotubes (11% reduction after 1100 cycles),<sup>3</sup> and arrays of anodic TiO<sub>2</sub> nanotubes layered with Co(OH)<sub>2</sub> by cathodic deposition (18% reduction after 1000 cycles).<sup>8</sup> Moreover, these nanoneedles showed comparatively better stability than other metal oxides listed in ESI; Table S1.† Even after such a large number of cycles continued continuously for couples of days, 1D  $\beta$ -TiO<sub>2</sub> nanoneedles were still intact and did not appear degraded or damaged by electrochemical interactions. This remarkable cycling performance of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles confirms their superior durability, the longer lifetime, high stability, and excellent electrochemical reversibility in Na<sub>2</sub>SO<sub>4</sub> electrolyte. This is ascribed to their nanostructure, large surface area, clearly visible textural boundaries, enhanced electrical conductivity, and well stable electrochemical reactions. Therefore, the core-shell formation of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles with carbon based materials or conducting polymers may be helpful to overcome their swelling and shrinking and to further improve their stability at large extent.

## Conclusions

Transparent large area array of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles were synthesized on conducting glass substrates *via* hot-filament metal vapor deposition. The 1D  $\beta$ -TiO<sub>2</sub> nanoneedles of the orthorhombic crystalline structure of brookite were conical in shape and tapered in the downward direction with a diameter of about 24–26 nm and a length of ~650 nm. The XPS double-peak features of the Ti(2p<sub>3/2</sub>) and Ti(2p<sub>1/2</sub>) core levels revealed that the large area array consisted exclusively of pure stoichiometric TiO<sub>2</sub> nanoneedles (*i.e.* Ti : O = 1.02 : 2). The large area array of 1D  $\beta$ -TiO<sub>2</sub> nanoneedles with large surface area and clearly visible textural boundaries provides distinct advantages of highly active surface sites and significantly facilitate for Na<sup>+</sup> ion diffusion and charge transfer. The excellent electrochemical performance of large area arrays of 1D  $\beta$ -TiO<sub>2</sub> nanoneedle in the electrolyte of Na<sup>+</sup> ions delivered specific capacitance of 192.2 mF g<sup>-1</sup> and 34.1 mF g<sup>-1</sup> from CC and CV, respectively, high columbic efficiency (98%), energy density of 3.04 W h kg<sup>-1</sup>, and power density of 1683 W kg<sup>-1</sup>. The EIS tests confirmed the important role of textural boundaries of the 1D  $\beta$ -TiO<sub>2</sub> nanoneedles in the enhancement of Na<sup>+</sup> ion diffusion for highly stable non-faradic capacitance behavior. Additionally, the excellent cycle stability up to 10 000 cycles makes the 1D  $\beta$ -TiO<sub>2</sub> nanoneedle arrays an excellent material for the fabrication of supercapacitors with highly stable performance. Overall, the dimension and size-dependent electrochemical properties described in this work have opened up a new era in nanotechnology for the fabrication of supercapacitors with long cycle lifetime.

## Acknowledgements

The authors would like to thank the Department of Science and Technology (DST), Ministry of Science and Technology of India,

for their financial support of this research under INSPIRE Faculty award No. DST/INSPIRE Faculty Award/2013/IFA13-PH-63.

## References

- Z. S. Wu, K. Parvez, X. L. Feng and K. Mullen, *Nat. Commun.*, 2013, **4**, 2487.
- L. H. Bao, J. F. Zang and X. D. Li, *Nano Lett.*, 2011, **11**, 1215–1220.
- R. B. Ambade, S. B. Ambade, N. K. Shrestha, Y. C. Nah, S. H. Han, W. Lee and S. H. Lee, *Chem. Commun.*, 2013, **49**, 2308–2310.
- Q. Q. Tan, Y. X. Xu, J. Yang, L. L. Qiu, Y. Chen and X. X. Chen, *Electrochim. Acta*, 2013, **88**, 526–529.
- R. S. Devan, R. A. Patil, J. H. Lin and Y. R. Ma, *Adv. Funct. Mater.*, 2012, **22**, 3326–3370.
- T. Brezesinski, J. Wang, J. Polleux, B. Dunn and S. H. Tolbert, *J. Am. Chem. Soc.*, 2009, **131**, 1802–1809.
- X. H. Lu, G. M. Wang, T. Zhai, M. H. Yu, J. Y. Gan, Y. X. Tong and Y. Li, *Nano Lett.*, 2012, **12**, 1690–1696.
- G. G. Zhang, C. J. Huang, L. M. Zhou, L. Ye, W. F. Li and H. T. Huang, *Nanoscale*, 2011, **3**, 4174–4181.
- A. Ramadoss and S. J. Kim, *J. Alloys Compd.*, 2013, **561**, 262–267.
- Q. Wang, Z. H. Wen and J. H. Li, *Adv. Funct. Mater.*, 2006, **16**, 2141–2146.
- H. Zhou and Y. R. Zhang, *J. Power Sources*, 2013, **239**, 128–131.
- T. A. Kandiel, A. Feldhoff, L. Robben, R. Dillert and D. W. Bahnemann, *Chem. Mater.*, 2010, **22**, 2050–2060.
- M. Posternak, A. Baldereschi, E. J. Walter and H. Krakauer, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2006, **74**, 125113.
- S. D. Mo and W. Y. Ching, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1995, **51**, 13023–13032.
- J. Wang, J. Polleux, J. Lim and B. Dunn, *J. Phys. Chem. C*, 2007, **111**, 14925–14931.
- J. Z. Chen, W. Y. Ko, Y. C. Yen, P. H. Chen and K. J. Lin, *ACS Nano*, 2012, **6**, 6633–6639.
- H. Zhou and Y. Zhang, *J. Phys. Chem. C*, 2014, **118**, 5626–5636.
- M. H. Yang, P. C. Chen, M. C. Tsai, T. T. Chen, I. C. Chang, H. T. Chiu and C. Y. Lee, *CrystEngComm*, 2014, **16**, 441–447.
- H. Kominami, M. Kohno and Y. Kera, *J. Mater. Chem.*, 2000, **10**, 1151–1156.
- R. A. Patil, C. P. Chang, R. S. Devan, Y. Liou and Y. R. Ma, *ACS Appl. Mater. Interfaces*, 2016, **8**, 9872–9880.
- R. A. Patil, R. S. Devan, J. H. Lin, Y. Liou and Y. R. Ma, *Sci. Rep.*, 2013, **3**, 3070.
- R. S. Devan, J. H. Lin, Y. J. Huang, C. C. Yang, S. Y. Wu, Y. Liou and Y. R. Ma, *Nanoscale*, 2011, **3**, 4339–4345.
- R. S. Devan, C. L. Lin, S. Y. Gao, C. L. Cheng, Y. Liou and Y. R. Ma, *Phys. Chem. Chem. Phys.*, 2011, **13**, 13441–13446.
- R. S. Devan, J. H. Lin, W. D. Ho, S. Y. Wu, Y. Liou and Y. R. Ma, *J. Appl. Crystallogr.*, 2010, **43**, 1062–1067.
- R. A. Patil, R. S. Devan, Y. Liou and Y. R. Ma, *Sol. Energy Mater. Sol. Cells*, 2016, **147**, 240–245.



- 26 Y. Wang, H. J. Sun, S. J. Tan, H. Feng, Z. W. Cheng, J. Zhao, A. D. Zhao, B. Wang, Y. Luo, J. L. Yang and J. G. Hou, *Nat. Commun.*, 2013, **4**, 2214.
- 27 M. Salari, S. H. Aboutalebi, A. T. Chidembo, I. P. Nevirkovets, K. Konstantinov and H. K. Liu, *Phys. Chem. Chem. Phys.*, 2012, **14**, 4770–4779.
- 28 W. B. Hu, L. P. Li, G. S. Li, C. L. Tang and L. Sun, *Cryst. Growth Des.*, 2009, **9**, 3676–3682.
- 29 A. M. Zaky, S. S. A. El-Rehim and B. M. Mohamed, *Int. J. Electrochem. Sci.*, 2006, **1**, 17–31.
- 30 J. Lee, J. Choi, J. Lee, S. K. Choi and H. D. Chun, *Nanotechnology*, 2005, **16**, 1449–1453.
- 31 M. Salari, S. H. Aboutalebi, K. Konstantinov and H. K. Liu, *Phys. Chem. Chem. Phys.*, 2011, **13**, 5038–5041.
- 32 M. Salari, K. Konstantinov and H. K. Liu, *J. Mater. Chem.*, 2011, **21**, 5128–5133.
- 33 M. H. Yu, Y. X. Zeng, C. Zhang, X. H. Lu, C. H. Zeng, C. Z. Yao, Y. Y. Yang and Y. X. Tong, *Nanoscale*, 2013, **5**, 10806–10810.
- 34 W. W. Liu, X. B. Yan and Q. J. Xue, *J. Mater. Chem. C*, 2013, **1**, 1413–1422.
- 35 H. M. Zheng, T. Zhai, M. H. Yu, S. L. Xie, C. L. Liang, W. X. Zhao, S. C. I. Wang, Z. S. Zhang and X. H. Lu, *J. Mater. Chem. C*, 2013, **1**, 225–229.

