Field-induced slow relaxation in magnetization of dinuclear and trinuclear Co\textsuperscript{III}⋯\textsuperscript{Mn} complexes\textsuperscript{†}

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Two new dinuclear, [Co(3EtO-L5)(μ-CN)Mn(L4a)Cl] (1) and [Co(3EtO-L5)(μ-CN)Mn(L4a)Br] (2), and two trinuclear [(Co(3EtO-L5)(μ-CN))\textsubscript{2}Mn(L4a)] (3) and [(Co(3EtO-L5)(μ-CN))\textsubscript{2}Mn(L4a)][NO\textsubscript{3}] (4) complexes were prepared and thoroughly characterized. (H\textsubscript{2}3EtO-L5 = N,N'-bis(3-ethoxy-2-hydroxybenzylidene)-1,6-diamino-3-azahexane, L4a\textsuperscript{2−} = N,N'-ethane-bis(salicylideneiminate)dianion). The crystal structures were determined for all four compounds, while the static and dynamic magnetic properties were studied only for compounds 1–3. It has been revealed by simultaneous fitting of temperature and field dependent magnetic data and by using the spin Hamiltonian formalism involving the axial anisotropy term that the manganese(III) atoms possess relatively large and negative axial magnetic anisotropy in 1–3, with $D = -3.9(2)$ cm\textsuperscript{−1} in 1, $-4.9(2)$ cm\textsuperscript{−1} in 2, and $-4.1(1)$ cm\textsuperscript{−1} in 3. These results were supported by ab initio CASSCF calculations which were in good agreement with the experimental ones, however, a small rhombicity was calculated contrary to the experimental evaluations: $D_{\text{calc}}/D_{\text{calc}} = -3.2$ cm\textsuperscript{−1} (0.04) in 1, $-3.0$ cm\textsuperscript{−1} (0.03) in 2, and $-3.6$ cm\textsuperscript{−1} (0.04) in 3. The measurements of dynamic magnetic data confirmed that compounds 1–3 represent a new type of Mn(III) field-induced single-ion magnet. The peak maxima of the frequency dependent out-phase susceptibility were below the lowest accessible temperature in all the three cases (i.e. below 1.9 K) and this prevented us from the construction of the Argand diagram. Nevertheless, the approximate procedure for extracting the spin-reversal barrier ($U_{\text{eff}}$) and relaxation time ($\tau_0$) was used, and the values of $U_{\text{eff}}$ ranging from 11 to 20 K and $\tau_0$ from 0.1 to 19 $\times$ 10\textsuperscript{−7} s were obtained.

Introduction

Compounds exhibiting slow relaxation of magnetization (SRM), so called single-molecule magnets (SMMs), attract great attention in the field of molecule-based magnetic materials research. Their potential utilization in molecular spintronics and information storage devices has been intensively studied.\textsuperscript{3} SMMs can behave as molecular nanomagnets possessing distinct magnetic states separated by an energy barrier for spin reversal, while the height of this barrier is controlled by magnetic anisotropy of paramagnetic centre/centres of a molecule. Therefore, the rational design of such materials consists mostly of tuning of uniaxial magnetic anisotropy to large negative values (described by the zero-field splitting parameter $D$); however, there are several examples of mononuclear compounds exhibiting SRM and having positive $D$ values.\textsuperscript{3}

The Mn(III) atom in a hexacoordinate environment is due to its electronic structure in the high-spin state $[\text{HS}, S = 2, \text{i}_{2g} \text{e}_{g}]$. For $E$ ground spectroscopic term) subject of Jahn–Teller distortion.\textsuperscript{4} As a result, a significant prolongation of two opposite chromophore bond lengths and thus formation of an axially elongated tetragonal bipyramidal coordination polyhedron is observed. This geometry is favourable for the occurrence of relatively large and negative uniaxial magnetic anisotropy, as it was observed on numerous examples previously.\textsuperscript{4} Therefore, paramagnetic compounds involving the Mn(III) atoms are the promising objects for the preparation and study of compounds exhibiting SRM. Especially, the Mn(III) compounds involving tetradentate salen-like Schiff base ligands $[\text{L}_4\text{ff}]$ stands for a general abbreviation of a tetradeinate salen-like ligand, salen$^{2−} = N,N'$-ethane-bis(salicylideneiminate) dianion, further abbreviated as L4a$^{2−}$) are the auspicious candidates for the preparation of magnetically interesting compounds. This is due to the ability of L4$^{2−}$ compounds to act as ligands forming the equatorial plane in complexes with axially elongated coordination geometries with expectable negative magnetic anisotropy. Previous reports on the observation of SRM in [Mn$^{III}$L4]\textsuperscript{2+}-type compounds dealt with three basic structural types, i.e. poly-nuclear compounds (a) involving the $[\text{Mn}^{III}\text{L}_4]$\textsuperscript{2+} subunit coordinated by bridging ligands or complexes,\textsuperscript{3} (b) involving the phenoxo bridged dimeric subunits $[(\text{Mn}^{III}\text{L}_4)\text{L}]$\textsuperscript{2+} which are further coordinated in apical positions by ligands/complexes,\textsuperscript{4}
and (c) involving the [Mn^{III}[L4](Solv)]^+ subunits (Solv = water or alcohol molecule) coordinated by cyanidometallates. In the latter group, the very first example of the Mn^{III} SMM with only one paramagnetic centre (so-called Single-Ion Magnet, SIM) has been reported recently. The heterobimetallic Co^{III}-CN-Mn^{III} compound shows relatively small negative $D = -3.3$ cm$^{-1}$ and SRM occurs below 3 K.

Previously, we reported on heterobimetallic [Co^{III}(3EtO-L5)-($\mu$-CN)Fe^{III}(L4a)Cl] compound (further abbreviated as CoFeCl, where H$_3$3EtO-L5 = N,N-bis(3-ethoxy-2-hydrobenzylidene)-1,6-diamino-3-azahexane), having relatively large magnetic anisotropy ($D = +1.47$ cm$^{-1}$) on a usually very isotropic Fe^{III} atom. This was most probably caused by a rather long N$_{CN}$-Fe-Cl axial bond formed by bond lengths Fe-N$_{CN}$ = 2.236(2) and Fe-Cl = 2.4075(7) Å (N$_{CN}$ = the cyanido nitrogen atom). Inspired by these findings we attempted to prepare the Mn^{III} analogues of the aforementioned compounds, which should possess large uniaxial magnetic anisotropy accentuated by the expected axially elongated tetragonal bipyramidal coordination polyhedron of the Mn^{III} atom. Furthermore, in order to analyse the influence of the halido ligand exchange on magnetic properties, we tried to prepare a series of CoMnX compounds (with X = F, Cl, Br, I). However, we were able to prepare only two isostructural compounds with the formulas [Co(3EtO-L5)[($\mu$-CN)Mn(L4a)Cl]] (1) and [Co(3EtO-L5)[($\mu$-CN)Mn(L4a)Br]] (2). The other syntheses involving iodide and fluoride anions resulted in [[Co(3EtO-L5)[($\mu$-CN)]$_2$Mn(L4a)I] (3) and [[Co(3EtO-L5)[($\mu$-CN)]$_2$Mn(L4a)]($\mu$-NO$_3$)] (4). Herein, we report on the syntheses, crystal structures (1–4) and thorough investigations of magnetic properties of the compounds 1–3. While our main intention was to study influence of the different axial ligand fields (and different axial lengths) on magnetic anisotropy of the Mn(n) atom, the magnetic properties of 4 were not studied due to the structural similarity with compound 3.

Results

Synthesis

Compounds 1, 2 and 3 were prepared in a similar way. The [Co(3EtO-L5)[CN]] precursor complex was mixed with the corresponding [Mn(L4a)X] or [Mn(L4a)[H$_2$O]] compounds (X = Cl in 1, Br in 2) in methanol and they were stirred under heating for 15 minutes. The solution was filtered to a beaker and left to cool down. Then, the beaker was sealed with aluminium foil and placed in a big flask with diethyl ether. Single crystals of high quality were obtained by slow diffusion of diethyl ether to the mother liquor. The products of the first two reactions were identified as [Co(3EtO-L5)[($\mu$-CN)Mn(L4a)Cl]] (1) and [Co(3EtO-L5)[($\mu$-CN)Mn(L4a)Br]] (2), while the described reaction resulted in an unexpected trinuclear product [[Co(3EtO-L5)[($\mu$-CN)]$_2$Mn(L4a)] (3, Scheme 1) in the case of the iodido precursor complex. The attempts to synthesise the fluoride analogous of 1 and 2 were performed in two distinct ways. Firstly, the previously reported polymeric [Mn(L4a)[($\mu$-F)$_n$] compound was reacted with [Co(3EtO-L5)[CN]] in methanol. The length of heating and stirring varied from 15 minutes to 24 hours. However, the slow diffusion of diethyl ether led only to crystallization of the precursor complexes in every case. Secondly, in order to avoid the formation of the [Mn(L4a)[($\mu$-F)$_n$]$_n$ polymer we used the [Mn(L4a)Cl] precursor, which was subsequently reacted with AgNO$_3$ (AgCl was filtered off) in methanol. Then, the solution was mixed with the solid of [Co(3EtO-L5)[CN]] and NaF was added after 10 minutes of stirring. After 30 minutes of heating and stirring, the reaction mixture was filtered, cooled down and left to undergo slow-diffusion of diethyl ether as in the previous cases. The resulting crystalline product did not contain fluoride anions and again as in the case of 3, it was a trinuclear complex with the formula of [[Co(3EtO-L5)[($\mu$-CN)]$_2$Mn(L4a)]($\mu$-NO$_3$)] (4).

Crystal structure of 1 and 2

The crystal structures of 1 and 2 are isostructural, crystallizing in monoclinic space group $P2_1$ (Table 1) and they are also isostructural with the previously reported compound [Co(3EtO-L5)[($\mu$-CN)Fe(L4a)Cl]].

The crystal structures of 1 and 2 consist of dinuclear [Co(3EtO-L5)[($\mu$-CN)Mn(L4a)X]] (X = Cl, and Br, respectively) molecules in which the metal atoms are bridged by the cyanido ligand (Co-CN-Mn linkage, Fig. 1). The arrangement of the metal atoms and the bridging ligand is slightly bent with the Co–C–Mn angles equal to 163.6° in 1 and 163.7° in 2. The Co^{III} centre is coordinated by three nitrogen and two oxygen atoms.

Scheme 1 Schematic representations of complexes 1–3 and their preparations. Note: the preparation of complex 4 is described in the text below.
Table 1  Crystal data and structure refinements for 1–4

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$^a R_1 = \sum(|F_o| - |F_c|)/\sum|F_o|$. $^b wR_2 = \left\{ \sum[w(F_o^2 - F_c^2)]/\sum[w(F_o^2)] \right\}^{1/2}$.

Fig. 1  Molecular structures of (a) [Co(3EtO-L5)(μ-CN)Mn(L4a)Cl] (1), (b) [Co(3EtO-L5)(μ-CN)Mn(L4a)Br] (2), (c) [Co(3EtO-L5)(μ-CN)2Mn(L4a)] [3], (d) [(Co(3EtO-L5)(μ-CN)L4a)2Mn(L4a)][NO3] (4). Selected bond lengths (in Å): Mn(1)–N(5) = 1.991(3), Mn(1)–N(6) = 1.999(3), Mn(1)–N(4) = 2.380(3), Mn(1)–O(5) = 1.887(3), Mn(1)–O(6) = 1.887(3), Mn(1)–C(1) = 2.550(1), Co(1)–N(1) = 1.926(3), Co(1)–N(2) = 1.893(3), Co(1)–N(3) = 1.962(3), Co(1)–O(1) = 1.992(3), Co(1)–O(2) = 1.908(2) for 1, Mn(1)–N(5) = 1.992(4), Mn(1)–N(6) = 1.989(4), Mn(1)–N(4) = 2.325(4), Mn(1)–O(5) = 1.877(4), Mn(1)–O(6) = 1.884(4), Mn(1)–Br(1) = 2.739(4), Co(1)–N(1) = 1.925(4), Co(1)–N(2) = 1.899(4), Co(1)–N(3) = 1.990(4), Co(1)–O(1) = 1.908(4) for 2, Mn(1)–N(5) = 1.994(4), Mn(1)–N(6) = 1.971(5), Mn(1)–N(4) = 2.283(5), Mn(1)–N(7) = 2.347(3), Mn(1)–O(5) = 1.873(3), Mn(1)–O(6) = 1.888(3), Co(1)–N(1) = 1.890(4), Co(1)–N(2) = 1.920(5), Co(1)–N(3) = 1.971(3), Co(1)–O(1) = 1.900(3), Co(1)–O(2) = 1.915(2), for 3 and Mn(1)–N(5) = 1.978(2), Mn(1)–N(6) = 1.978(2), Mn(1)–N(4) = 2.316(2), Mn(1)–N(7) = 2.316(2), Mn(1)–O(5) = 1.890(4), Mn(1)–O(6) = 1.890(4), Co(1)–N(1) = 1.926(2), Co(1)–N(2) = 1.894(2), Co(1)–N(3) = 1.971(2), Co(1)–O(1) = 1.899(2), Co(1)–O(2) = 1.920(2) for 4.
from the 3EtO-L5$^{2-}$ ligand. The sixth coordination site is occupied by the carbon atom from the bridging cyanido ligand. The chromophore bond lengths are typical for the low-spin (LS) state with the lengths ranging from 1.89 to 1.91 Å for Co–C, from 1.91 to 1.96 Å for Co–N, and from 1.90 to 1.91 Å for Co–O bonds (Table 2). The coordination polyhedrons of the Co$^{III}$ atoms can be described as slightly distorted octahedrons in both compounds. The Mn$^{III}$ central atoms are also hexacoordinate with four equatoril donor atoms (N$_2$O$_2$) coming from the L4$^{3-}$ ligand. The axial ligands are represented by one cyndio nitrogen (N$_{CN}$) and one halogen atom, thus forming the {MnN$_2$O$_2$Cl} (1) and {MnN$_2$O$_2$Br} (2) chromophores. The bond lengths within the equatorial plane are significantly shorter than those in axial due to the Jahn–Teller effect on the Mn$^{III}$ atoms. This can be documented by the comparison of the Mn–N bond lengths: the equatorial manganese–imino nitrogen (Mn–N$_{im}$) bond lengths are much shorter (1.91–1.95 Å) than the axial Mn–N$_{CN}$ bond lengths (2.380(3) in 1 and 2.325(4) Å in 2). The Jahn–Teller elongation of the bond lengths along the axis of the d$_z^2$ orbitals is further accentuated by long Mn–X lengths (2.5498(10) in 1 and 2.7394(8) Å in 2). When comparing these values with the bond lengths in the previously reported CoFeCl$_4$ value calculated from four bond lengths.

The crystal structures of 3 and 4 with growing atomic radius and polarizability. The crystal structures of 3 and 4 can be described as slightly distorted octahedrons with the {CoN$_2$O$_2$C} chromophores. The chromophore bond lengths correspond to the LS nature of the Co$^{III}$ centre with the lengths ranging from 1.89 to 1.90 Å for Co–C, from 1.91 to 1.97 Å for Co–N, and from 1.90 to 1.91 Å for Co–O bonds (Table 2). The Mn$^{III}$ atom is hexacoordinate with the {MnN$_2$O$_2$} chromophores in both compounds. The bond lengths within the equatorial planes are very similar to those found in 1 and 2 (d(Mn–N$_{im}$) = 1.97–1.98 Å, d(Mn–O) = 1.88–1.89 Å, Table 2). The axial is formed by two N$_{im}$ atoms from two coordinated [Co3EtO-L5(CN)]$^{2-}$ chromophores. Again, due to the Jahn–Teller effect long axial bonds are observed as follows: d(Mn–N$_{CN}$) = 2.283(3), 2.347(3) Å in 3 and 2.316(2) in 4 (two Mn–N$_{CN}$ bonds are equivalent due to symmetry). It can be concluded that compounds 3 and 4 have a slightly shorter axial than compounds 1 and 2 (Table 2).

In the crystal structure of 3 the trinuclear complex molecules are assembled in 1D supramolecular chains through N–H$\cdots$O hydrogen bonding between the nitrogen atoms from the amine groups and oxygen atoms from the ethoxy groups; both located on the [Co3EtO-L5(CN)]$^{3+}$ moieties of neighbouring trinuclear complex molecules. Each complex molecule is held by its in-chain neighbour by two N–H$\cdots$O hydrogen bonds thus forming a synthon with the R$_2$h$_{10}$ ring topology (Fig. 2).

The distances of these non-covalent contacts are rather long: d(N–O) = 3.307(4) and 3.168(4) Å.

The crystal structure of 4 consists of [{Co3EtO-L5[µ-CN]}$_2$Mn(L4a)]$^{2-}$ molecules held together by N–H$\cdots$O hydrogen bonding between the nitrogen atom from the amine group and ethoxy group with the distance d(N–O) = 3.173(4) Å. Packing of the [{Co3EtO-L5[µ-CN]}$_2$Mn(L4a)]$^{2-}$ molecules in the crystal structure of 4 is rather different from packing observed for 3; the complex molecules form a 2D layer structure in which each trinuclear complex interacts with four other complex molecules via four N–H$\cdots$O hydrogen bonds (Fig. 3). Due to the two fold rotational axis on the manganese$^{II}$ atom together with inversion centres located in between the molecules only one of these four hydrogen bonds is symmetrically independent. Its length is rather long again: d(N–O) = 3.173(4) Å.

### Static magnetic properties

The temperature and field dependent magnetic properties of compounds 1–3 were measured on a SQUID magnetometer and are depicted in Fig. 3. The room temperature value of the effective magnetic moment $\mu_{eff}$ is found in the range of 4.8–4.9 $\mu_B$, which is close to the spin-only value for S = 2 equal to 4.90 (g = 2.0). The $\mu_{eff}$ is almost constant down to 30 K, and then decreases to the value $\approx 4$ $\mu_B$ which may indicate zero-field splitting on the Mn$^{III}$ atoms and also the possibly of weak
intermolecular interactions. Moreover, the isothermal magnetizations do not saturate to $M_{\text{mol}}/N_A\mu_B = 4.0$ at maximum applied field for any of 1–3, which can be also explained by significant zero-field splitting. For that reason, the subsequent spin Hamiltonian, involving the single ion axial ($D$) and rhombic ($E$) ZFS parameters, Zeeman term and molecular field correction $z\eta$, were postulated\(^\text{12}\)

$$\hat{H} = D(S_z^2 - 3/3) + E(S_x^2 - S_y^2) + \mu_B g_\theta S_\theta - z\eta(S_\theta^2)$$

where $\eta$ defines the orientation of the magnetic field vector in polar coordinates as $B_\theta = B(\sin \theta \cos \varphi, \sin \theta \sin \varphi, \cos \theta)$. The molar magnetization in the $a$-direction of the magnetic field, $M_a$, was numerically calculated as

$$M_a = -\frac{N_A \sum \left( \sum \sum C_k C_\theta (Z_\theta) \exp(-\varepsilon_{a\theta}/kT) \right)}{\sum \exp(-\varepsilon_{a\theta}/kT)}$$

Fig. 2 A perspective view on supramolecular chains held by N–H⋯Cl hydrogen bonds (dotted lines) in 1 (a) and N–H⋯Br in 2 (b). A perspective view on supramolecular chains held by N–H⋯O hydrogen bonds in 3 (c). The donor⋯acceptor distances for selected non-covalent contacts (in Å): $d$(N⋯Cl) = 3.189(3) in 1, $d$(N⋯Br) = 3.306(4) in 2, $d$(N⋯O) = 3.168(4) and 3.307(4) in 3 and 3.173(4) in 4 (d).
where $Z_a$ is the matrix element of the Zeeman term for the $a$-direction of the magnetic field and $C$ are the eigenvectors resulting from the diagonalization of the complete spin Hamiltonian matrix. The presence of $z_j$ means that the iterative procedure must be applied. The final calculated molar magnetization was calculated as an integral average in order to properly simulate the powder sample signal.

$$M_{\text{mol}} = \frac{1}{4\pi} \int_0^{2\pi} \int_0^{\pi} M_s \sin \theta d\theta d\phi$$

Both, temperature and field dependent magnetic data were fitted simultaneously in order to obtain a consistent set of parameters; however, we found that the fitting procedure was not sensitive enough to the rhombicity of the systems ($E/D$ ratio), and so we fixed $E/D = 0$. This is also supported by multi-configurational complete active space self-consistent field (CASSCF) calculations, where $E/D$ was found to be less than 0.04 (Table 3). Consequently, the single-ion axial ZFS parameter $D$ was found in the range from $-3.9$ to $-4.4$ cm$^{-1}$ (Fig. 3). These values are close to other Mn$^{III}$ complexes listed in Table 3.

**Dynamic magnetic properties**

The AC susceptibility data measured in zero static magnetic field showed no out-of-phase susceptibility signal (Fig. S1, ESI†), but the field dependent measurement at the lowest temperature (1.9 K) confirmed field-induced slow relaxation of magnetization in 1–3 (Fig. S2†). Therefore, the AC susceptibility data were measured at non-zero static magnetic field, $B_{DC} = 0.2$ and 0.4 T, for various frequencies of small AC field as depicted in Fig. 3, and S3 in ESI† respectively. Clearly, the out-of-phase signal is frequency dependent, but lower temperature would be needed to see maxima on out-of-phase susceptibilities.
This prevented us from constructing the Argand (Cole–Cole) diagram and analysing the data to obtain information about the spin reversal barrier by a one-component Debye model. Nevertheless, at least the approximate relationship\textsuperscript{14}

\[
\ln(\chi' / \chi) = \ln(2\pi f_0) + \frac{U}{kT}
\]  

for the extraction of the relaxation time ($\tau_0$) and spin reversal barrier ($U$) was applied to low temperature data and higher applied frequencies as depicted in Fig. S4 in ESL.\textsuperscript{†} In such a way we obtained sets of parameters, which are listed in detail in Table S1\textsuperscript{†} for both static and dynamic magnetic properties of Mn\textsuperscript{III} single-ion magnets.

### Theoretical calculations of spin Hamiltonian parameters

With the aim to support the outcomes of the analysis of magnetic data, multi-reference \textit{ab initio} calculations based on state-averaged complete-active-space self-consistent field (SACASSCF) wave functions complemented by N-electron valence second-order perturbation theory (NEVPT2) were done using ORCA software. The active space was set to the metal-based d-orbitals, CAS(4,5). Firstly, the calculations were performed on the molecular geometries of the whole complex units for [Co(3EtO-L5)(μ-CN)Mn(L4a)Cl] of 1 and [Co(3EtO-L5)(μ-CN)Mn(L4a)Br] of 2. Due to the fact that such calculations are time-consuming, especially for compound 3, which has a larger number of atoms than compounds 1 and 2, we decided to test the possibility to truncate the molecular geometry of complex compounds 1–3 into geometrical fragments without the diamagnetic Co\textsuperscript{III} subunits as [(CN)Mn(L4a)Cl]− (1′), [(CN)Mn(L4a)Br]− (2′) and [(CN)\textsubscript{2}Mn(L4a)]− (3′). The achieved parameters of g-tensors and D-tensors are summarized in Table 4. It is evident that the geometry truncation led only to minor changes in D-values, \textit{i.e.} less than 0.15 cm\textsuperscript{−1} for 1 and 2,
and this is usually within the experimental error. The calculated values are in good agreement with the experimentally derived ones, we observed slightly lower calculated $D$-values, which may be attributed to the neglecting of spin–spin contribution to the $D$-tensor. Next, the rhombicity of the studied compounds was found very low, $E/D < 0.05$ justifying the neglecting of the $E$-term in fitting the magnetic data. Furthermore, the main axes of the $D$-tensor are visualized in Fig. 4 for the molecular fragment 1’ showing that the easy axis of magnetization can be identified with the Jahn–Teller axis, hence with the CN–Mn–Cl bonds.

Conclusions

In this article we reported on the synthesis and properties of new heterobimetallic Co$^{II}$–Mn$^{III}$ compounds. Two novel dinuclear, $[\text{Co}(3\text{EtO-L5})\mu\text{CN}]\text{Mn}(\text{L4a})\text{Cl}]$ (1) and $[\text{Co}(3\text{EtO-L5})\mu\text{CN}]\text{Mn}(\text{L4a})\text{Br}]$ (2), and two new trinuclear $[\text{Co}(3\text{EtO-L5})\mu\text{CN}]_3\text{Mn}(\text{L4a})\text{I}]$ (3) and $[\text{Co}(3\text{EtO-L5})\mu\text{CN}]_2\text{Mn}(\text{L4a})\text{I}]\text{(NO}_3\text{)}$ (4) compounds were prepared. The crystal structures were determined for all four herein presented compounds but magnetic properties were studied only for compounds 1–3 in greater details due to the structural similarity of compounds 3 and 4. The magnetic data analysis, involving the spin Hamiltonian with the axial zero-field splitting term, revealed a relatively large and negative value of the axial magnetic anisotropy parameter $D$ (in cm$^{-1}$): $-3.9(2)$ in 1, $-4.9(2)$ in 2, and $-4.1(1)$ in 3. The experimental study was supported by ab initio calculations at the CASSCF level of theory and the values of $D_{\text{calc}}$, which are in good agreement with the experimentally derived parameters, were obtained. Furthermore, small rhombicity ($E_{\text{calc}}$) for compounds 1–3 was proposed: $D_{\text{calc}} (E_{\text{calc}}/D_{\text{calc}}) = -3.2$ cm$^{-1}$ (0.04) in 1, $-3.0$ cm$^{-1}$ (0.03) in 2 and $-3.6$ cm$^{-1}$ (0.04) in 3. In summary, it has been found that the value of the $D$ parameter does not depend neither on the axial length or lowering of the ligand-field strength by the ligand exchange significantly. This can be documented by two contrasting examples. The longest axial length is observed for 2 (N–Mn–Br type, $d(\text{N}--\text{Br}) = 5.058(4)$ Å, angle (N–Mn–Br) = 174.5(1)$^\circ$) and this compound possesses also the largest negative $D$ value (−4.9 cm$^{-1}$). The shortest axial length is found in 3 (N–Mn–N type, $d(\text{N}--\text{N}) = 4.617(5)$ Å, angle (N–Mn–N) = 171.3(1)$^\circ$) but the value of the $D$ parameter (−4.1 cm$^{-1}$) does not adopt a value much different from that observed for 2.

The measurements of dynamic magnetic data confirmed that compounds 1–3 behave as field-induced single-ion magnets. The peak maxima of frequency dependent out-of-phase susceptibility (in static magnetic field of 0.4 T) were in all three cases below the lowest accessible temperature (1.9 K) and thus, we were not able to construct the Argand (Cole–Cole) diagram. Therefore, we used an approximate relationship for the extraction of the relaxation time ($\tau_0$) and spin reversal barrier ($U$) and this was applied to low temperature data and higher applied frequencies and as a result the values of $U_{\text{eff}}$ ranging from 11 to 20 K and $\tau_0$ from 0.1 to 19 $\times 10^{-7}$ s were obtained.

Experimental

Materials

All the starting chemicals were of analytical reagent grade and were used as received. CoCl$_2$·6H$_2$O, MnCl$_2$·4H$_2$O, MnBr$_2$, MnI$_2$, NaCN, AgNO$_3$, solvents and the organic compounds 2-hydroxybenzaldehyde, 3-ethoxy-2-hydroxybenzaldehyde, ethane-1,2-diamine, N-(2-aminoethyl)-1,3-propanediamine, and triethylamine (Et$_3$N) were obtained from commercial sources (Sigma-Aldrich Co., Acros Organics Co., Lachema Co. and Fluka Co.).

Synthesis of precursor complexes

The manganese(III) complexes were prepared according to the literature procedures. From the conductivity measurements it is apparent that iodide anion in [Mn(L4a)] is not coordinated to Mn$^{III}$ atom.

$[\text{Mn(L4a)Cl}]$. Yield: 93% (663 mg), Anal. calced for C$_{48}$H$_{48}$N$_{8}$O$_{8}$Cl$\cdot$Mn: C, 53.87; H, 3.95; N, 7.85. Found: C, 53.96; H, 3.74; N, 7.65%. $\Lambda_m$ (in dimethyl formamide, further abbr. as DMF, S cm$^{-2}$ mol$^{-1}$): 11.7.

$[\text{Mn(L4a)Br}]$. Yield: 82% (658 mg), Anal. calced for C$_{48}$H$_{48}$N$_{8}$O$_{8}$Br$\cdot$Mn: C, 47.90; H, 3.51; N, 6.98. Found: C, 47.88; H, 3.69; N, 6.57%. $\Lambda_m$ (DMF, S cm$^{-2}$ mol$^{-1}$): 15.8.

$[\text{Mn(L4a)J}]$. Yield: 74% (663 mg), Anal. calced for C$_{48}$H$_{48}$N$_{8}$O$_{8}$Cl$\cdot$Mn: C, 42.88; H, 3.14; N, 6.25. Found: C, 42.58; H, 3.41; N, 6.44%. $\Lambda_m$ (DMF, S cm$^{-2}$ mol$^{-1}$): 76.8.

Synthesis of complexes 1, 2, 3 and 4

Dark crystals of complexes 1–3 were obtained from a methanol solution (10 mL) of the corresponding precursor complexes (0.1 mmol) combined with a methanol/acetonitrile (10/10 mL) solution of $[\text{Co}(3\text{EtO-L5})(\text{CN})]$ (1) (0.1 mmol). The solutions were mixed and the resulting mixture was stirred at room temperature for 1 h. Slow diffusion of diethyl ether into the resulting solution in a closed bottle afforded black single crystals. Black crystals were filtered off, washed twice with cold methanol, diethyl ether and dried in vacuum.

The single crystals of 4 were prepared by the method involving NaF and AgNO$_3$ as described in the Synthesis chapter. In order to prepare a pure phase system the crystals of complex 4 were obtained by a different procedure. A methanol solution (10 mL) of complex 1 was mixed with a water solution (5 mL) of AgNO$_3$ (stoichiometric amount). A white powder (AgCl) was filtered off. Black crystals of 4 were obtained by slow diffusion of diethyl ether into the mother liquor in a closed bottle after a week.
[Co(3ETOL)(μ-CN)Mn(L4a)Br] (2). Yield: 77% (69 mg), Anal. calcd for C64H72N10O10I1Co2Mn1: C, 53.34; H, 5.03; N, 11.33%. Found: C, 53.09; H, 4.68; N, 9.53%. As (DMF, S cm⁻³ mol⁻¹): 10.5. FT-IR (cm⁻¹): 3096w r(C-H)ar; 2960w r(C-H)alip; 2928w r(C-H)alip; 2863w; 2139m r(C≡N); 1647w r(C≡N); 1621s r(C≡N); 1574m r(C≡N); 1496m r(C≡N); 1432m r(C≡N); 1415m r(C≡N); 1396s r(C≡N); 1385m r(C≡N); 1321w; 1296wm; 1213mw; 1150w; 1130w; 1108w; 1079w; 970w; 899w; 877w; 853w; 797w; 751m; 744m.

[(Co(3ETOL)(μ-CN)2Mn(L4a))I] (3). Yield: 66% (99 mg), Anal. calcd for C64H72N10O10I1Co2Mn1: C, 53.34; H, 5.03; N, 9.71. Found: C, 53.48; H, 5.14; N, 9.48%. As (DMF, S cm⁻³ mol⁻¹): 89.2. FT-IR (cm⁻¹): 3052w r(C-H)ar; 3022w r(C-H)ar; 2971w r(C-H)alip; 2936w r(C-H)alip; 2815w; 2713w; 2137m r(C≡N); 2131s r(C≡N); 1622s r(C≡N); 1574m r(C≡N); 1552m r(C≡N); 1466m r(C≡N); 1432w; 1415w; 1396w; 1354m; 1332w; 1251m; 1200w; 1174wm; 1163w; 1128w; 1085w; 1008w; 970m; 885w; 742w; 738m; 724w; 663w.

[(Co(3ETOL)(μ-CN)L4a)](NO3) (4). Yield: 71% (91 mg), Anal. calcd for C64H72N10O10I1Co2Mn1: C, 55.85; H, 5.27; N, 11.19. Found: C, 55.63; H, 5.41; N, 11.33%. As (DMF, S cm⁻³ mol⁻¹): 71.5. FT-IR (cm⁻¹): 3462w r(C≡N); 3215w r(C≡N); 3165w r(C≡N); 2935w r(C≡N); 2921w r(C≡N); 2815w; 2713w; 2137m r(C≡N); 2131s r(C≡N); 1622s r(C≡N); 1574m r(C≡N); 1552m r(C≡N); 1466m r(C≡N); 1432w; 1415w; 1396w; 1354m; 1332w; 1251m; 1200w; 1174wm; 1163w; 1128w; 1085w; 1008w; 970m; 885w; 742w; 738m; 724w; 663w.

**General methods**

Elemental analysis (CHN) was performed on a FLASH 2000 CHN Analyser (ThermoFisher Scientific). Infrared spectra of the complexes were recorded on a ThermoNicolet NEXUS 670 FT-IR spectrometer using the ATR technique on a diamond plate in the range 600–4000 cm⁻¹. The reported FT-IR intensities were defined as w = weak, m = medium, s = strong, and vs = very strong. The magnetic data were measured on powdered samples pressed into pellets using a MPMS XL-7 Quantum Design SQUID magnetometer. The experimental data were corrected for the diamagnetism of the constituent atoms by using Pascal’s constants.

**Single-crystal X-ray analysis details**

X-ray measurements on the selected crystals of 1, 2, 3 and 4 were performed on an Oxford Diffraction Xcalibur² Gemini equipped with a Sapphire2 CCD detector using Mo-Kα radiation at 100 K. The Crysalis program package (version 1.171.33.52, Oxford Diffraction) was used for data collection and reduction. The molecular structures were solved by direct methods SHEXL-2014 and all non-hydrogen atoms were refined anisotropically on F² using full-matrix least-squares procedure SHELS-2014. All the hydrogen atoms were found in differential Fourier maps and their parameters were refined using a riding model with Uiso(H) = 1.2 (CH, CH₂OH) or 1.5 (CH₃) Ueq. Non-routine aspects of the structure refinement are as follows: in 3, two ethoxy groups are disordered over two positions; in 4, the electron density from heavily disordered nitrate anion was removed by SQUEEZE procedure.

**Theoretical methods**

All theoretical calculations were performed with an ORCA 3.0 computational package. All calculations employed the chain-of-spheres (RIJ/COSX) approximation to exact exchange. The polarized triple-ζ-quality basis set def2-TZVP-f was used for all atoms except for carbon and hydrogen, for which the def2-SVP basis set was chosen. The ZFS and g tensors were calculated by using self-consistent field (SA-CASSCF) wave functions complemented by N-electron valence second order perturbation theory (NEVPT2). The active space of the CASSCF calculation was set to five d-orbitals of Mn (CAS(4,5)). The state-averaged approach was used, in which all 5 quintet, 45 triplet and 50 singlet states were equally weighted. The ZFS parameters, based on dominant spin–orbit coupling contributions from excited states, were calculated through quasi-degenerate perturbation theory (QDPT), in which approximations to the Breit–Pauli form of the spin–orbit coupling operator (SOMF approximation) and the effective Hamiltonian theory were utilized. The energies of calculated excited states and their contributions to the D-tensors are listed in Tables S1 and S2 (ESI†).

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**Notes and references**


13 The standard deviations were calculated as $\sigma_i = (P_i^{-1} \cdot S((N - k)^{1/2})$ where $P_i = \sum(\delta \mu_i \delta a_i \cdot \delta \mu_i \delta a_i)$ and $S = \sum(\mu_n - \mu_{n_{\exp}})^2$ with $n = 1$ to $N$; $a_i$ and $q_i$ are fitted parameters, $N$ is number of experimental points (sum of temperature and field dependent data), $\mu_n$ and $\mu_{n_{\exp}}$ are the calculated and experimental effective magnetic moments for given temperature and magnetic field. The $\sigma_i$ was then multiplied by Student’s $t_{0.05\%}$ to provide confidence limits with 95% probabilities listed in text.


