



Cite this: *Energy Environ. Sci.*, 2016, 9, 2916

## Temperature dependent behaviour of lead sulfide quantum dot solar cells and films†

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Despite increasing greatly in power conversion efficiency in recent times, lead sulfide quantum dot (PbS QD) solar cells still suffer from a low open circuit voltage ( $V_{OC}$ ) and fill factor (FF). In this work, we explore the temperature dependent behavior of ~9% efficient solar cells. In the temperature range of 290 to 230 K, we find increased  $V_{OC}$  and FF values without significant degradation of the short circuit current, leading to up to 10.3% efficiency at 230 K. The change in  $V_{OC}$  is driven by the decrease of the reverse saturation current which fits the PN-junction model. Using Schottky and single carrier devices, we measure the carrier mobility, diffusion lengths, and doping concentrations of PbS QD films with tetrabutylammonium iodide and ethane dithiol ligands as a function of temperature. Both mobility and diffusion length are found to decrease with decreasing temperature while device performance increases, indicating that the 260 nm thick active layer is fully depleted. Finally, we propose that further optimization of the doping concentrations could help achieve increased device performance at room temperature.

Received 2nd June 2016,  
Accepted 11th August 2016

DOI: 10.1039/c6ee01577h

www.rsc.org/ees

### Broader context

The increasing need for clean, renewable, and cheap energy sources has driven the research of a new generation of low-cost and highly efficient solar cells. Lead sulfide colloidal quantum dots (PbS QDs) are promising candidates for solar cell applications due to their size tunable bandgap, broad absorption profile, and their compatibility with large-scale production techniques. Recent progress has led to efficiencies above 10%, but the performance is still limited by a low open circuit voltage ( $V_{OC}$ ) and fill factor (FF). In this report, we systematically measure the temperature dependent behaviour of the most important physical properties of state-of-the-art solar cells in order to determine the limiting factors of the  $V_{OC}$  and FF. We then provide practical and quantitative guidelines for further efficiency enhancement.

## 1 Introduction

Lead sulfide quantum dots (PbS QDs) have been a topic of intense study for over a decade due to their excellent optoelectronic properties and their large versatility in such applications as infrared sensors,<sup>1–4</sup> infrared photon sources,<sup>5</sup> transistors,<sup>6–8</sup> and solar cells.<sup>9–12</sup> Their versatility stems mainly from the size dependent bandgap and the wide array of possible ligands and surface modifications which allow low cost device processing from solution and facile modification of the electronic properties.

PbS QD solar cells in particular have seen a rapid rise in solar cell performance, from less than 1% in 2005 to 10.6% in 2015.<sup>1,11</sup> This progress has been made possible by several factors. First, improvements in the synthesis of the QDs have increased the monodispersity and scalability of PbS QDs.<sup>13,14</sup> Secondly, a vast body of research has been carried out to improve post-deposition passivation of QD films,<sup>15–18</sup> which has been shown to reduce trap density and allows tuning of the band positions and Fermi levels.<sup>19</sup> Finally, an improved understanding of energy level alignment and band bending has led to increasingly sophisticated device structures which facilitate the extraction of generated charges.<sup>12,20–22</sup> To further improve solar cell efficiencies, the prevailing limitations must be understood. Presently, the open circuit voltage ( $V_{OC}$ ) in particular falls well short of its theoretical maximum and must be improved if PbS QD solar cells are to become a viable technology.<sup>23,24</sup> Studying the temperature dependent properties of PbS films has been very successful at probing both the optical and electronic properties of thiol-capped QDs,<sup>25,26</sup> but as yet no such study

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† Electronic supplementary information (ESI) available. See DOI: 10.1039/c6ee01577h



has been done for state-of-the-art solar cells. Moreover, no such study has been carried out for halide ligands which have become widespread in the QD field and which behave very differently to thiol ligands in many respects.<sup>27,28</sup>

In this work, we report the temperature dependent behavior of highly efficient solar cells comprising a layer of tetrabutylammonium iodide (TBAI) capped PbS and a layer of ethane dithiol (EDT) capped PbS. We observe a large increase in device performance at lower temperatures, mainly due to an increased  $V_{OC}$  and fill factor (FF) with only a slightly decreased short circuit current ( $J_{SC}$ ). We explain the origin of this behavior by measuring important electronic properties of PbS films such as the carrier mobility, permittivity, and doping concentration as a function of temperature. Finally, we provide guidelines to further improve PbS QD solar cell performance in the near future.

## 2 Results and discussion

### 2.1 Temperature dependent $JV$ characteristics

Following the bilayer strategy reported by the Bawendi group,<sup>12</sup> solar cells are fabricated *via* layer-by-layer spin casting of the active layer consisting of 200 nm TBAI capped PbS, which has been shown to be n-type,<sup>19,29,30</sup> and 60 nm EDT capped PbS, which is less n-type or even p-type depending on how much it has been oxidized.<sup>18,19</sup> The full structure including  $TiO_2$  and  $MoO_3/Au$  electrodes is shown in Fig. 1. The  $JV$  curves measured under AM1.5G 1 Sun intensity are shown in Fig. 1a and b and the corresponding parameters are displayed in Fig. 1c; at room temperature the solar cell displays a  $V_{OC}$  of 0.61 V, a  $J_{SC}$  of  $25.8 \text{ mA cm}^{-2}$ , a FF of 0.60, and an overall power conversion efficiency (PCE) of 9.4%, which is higher than previous reports

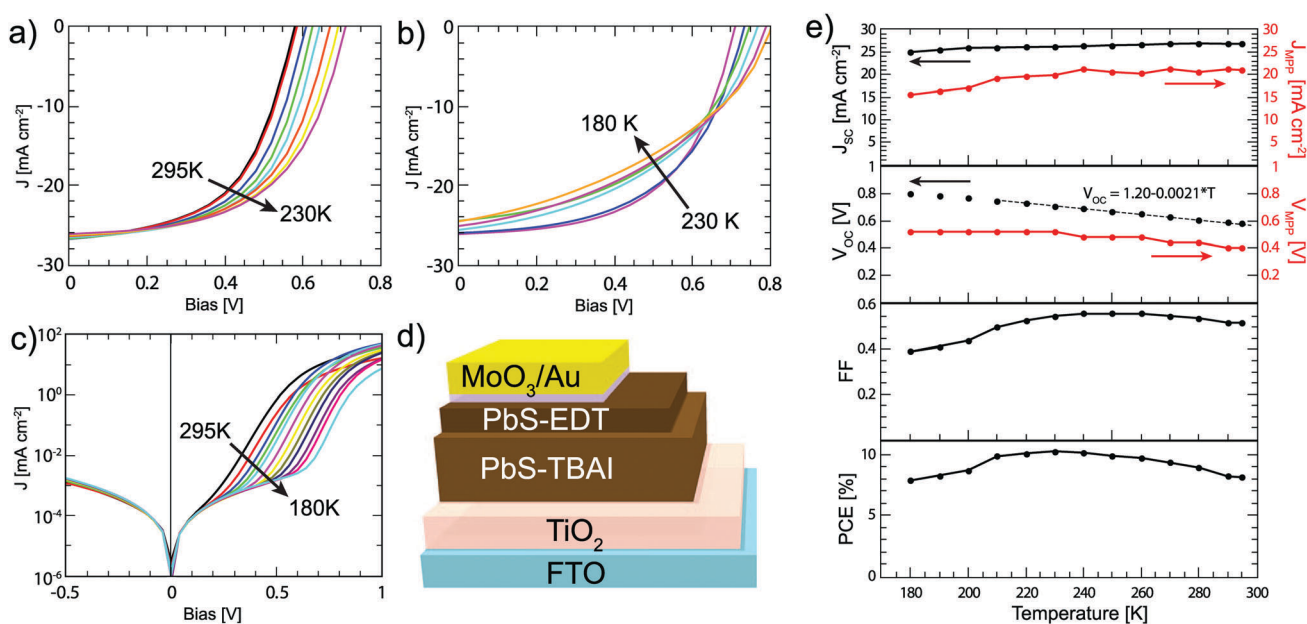
with a similar device structure.<sup>12</sup> The  $J_{SC}$  is in good agreement with the current density of  $24.5 \text{ mA cm}^{-2}$  calculated from the external quantum efficiency (EQE) (Fig. S1, ESI†).

As the temperature is decreased from 295 K to 180 K, the  $J_{SC}$  remains almost constant, decreasing only about 5% over the measured range. The  $V_{OC}$  shows a linear increase until 200 K and then plateaus at 0.80 V, or about 38% above the  $V_{OC}$  measured at room temperature. Extrapolating the  $V_{OC}$  from the linear regime to 0 K, a value of 1.20 eV is obtained. This is close to the bandgap of the active layer, which is  $1.37 \pm 0.05 \text{ eV}$  as determined by the first excitonic peak of the EQE spectrum. This indicates that the  $V_{OC}$  is not limited by Fermi level pinning to the trap states.<sup>31</sup> Below 200 K, the  $V_{OC}$  is likely limited by the energy level alignment with the  $TiO_2$  and Au electrodes.

The FF shows a more complex behaviour; it initially increases with decreasing temperature until a maximum at 250 K and then decreases at lower temperature. Overall, the  $JV$  parameters cause the PCE to initially increase with decreasing temperature to a maximum of 10.3% at 230 K, mainly due to the increasing  $V_{OC}$ . At temperatures lower than 230 K, the FF dominates the behavior of the solar cells, leading to an overall decreasing of the PCE with decreasing temperature.

The increase in  $V_{OC}$  cannot be explained by a changing bandgap. Photoluminescence measurements (Fig. S2, ESI†) show that as the temperature is decreased, the bandgap in fact narrows, and a smaller  $V_{OC}$  would be expected as a result. To explain the behaviour of the solar cell parameters, and in particular the very desirable increase in  $V_{OC}$ , we look at the Shockley diode equation for the current  $J$ ,

$$J = J_0 \exp\left(\frac{e(V + JR_S)}{nkT}\right) - \frac{V + JR_S}{R_{SH}} - J_{PH}. \quad (1)$$



**Fig. 1** (a)  $JV$  curves under 1 Sun illumination in the temperature range of 295–230 K showing increasing PCE, and (b) in the range of 230–180 K, where the PCE decreases. (c)  $JV$  curves measured in the dark. (d) Schematic of the solar cell device structure used. (e) Solar cell parameters extracted from the  $JV$  curves in (a and b) (black dots) as well as the current and voltage at the maximum power points (red dots). The solid lines are added as a guide to the eye; the dashed line is the fit for the linear part of the  $V_{OC}$  temperature dependence.



Here  $J_0$  is the reverse saturation current,  $e$  is the elementary charge,  $R_S$  and  $R_{SH}$  are the series and shunt resistances, respectively,  $n$  is the ideality factor,  $k$  is Boltzmann's constant, and  $V$  is the applied bias. Rearranging this equation to open circuit conditions ( $J = 0$ ), and assuming that the shunt current is much smaller than the photocurrent ( $V/R_{SH} \ll J_{PH}$ ), we find that the open circuit voltage is given by

$$V_{OC} = \frac{nkT}{e} \ln\left(\frac{J_{PH}}{J_0}\right) \quad (2a)$$

$$= \frac{nkT}{e} \alpha \ln(I) + c \quad (2b)$$

For the latter equality, we have made use of the relationship  $J_{PH} \propto I^\alpha$ , where  $I$  is the illumination intensity,  $\alpha$  is an empirical parameter indicating the linearity of the photocurrent with intensity, and  $c$  is a fitting parameter collecting all the terms independent of light intensity. From eqn (2a), we see that the temperature affects the  $V_{OC}$  directly *via* the linear term, and indirectly *via*  $J_0$  and possibly the ideality factor  $n$ . As shown above, the variation of  $J_{PH}$  with temperature is small and thus has a negligible effect on the  $V_{OC}$ . Since the  $V_{OC}$  increases with decreasing temperature, the linear term in eqn (2a) is clearly not the dominating factor. The temperature dependence of the ideality factor was measured under various illuminations under open circuit conditions, as shown in Fig. 2. In our devices  $\alpha$  is equal to  $0.95 \pm 0.3$  at all temperatures measured (Fig. S3, ESI†), which indicates that bimolecular recombination is low in these solar cells. Over a wide range of temperatures, no change is seen in the ideality factor within the measurement error. An ideality factor close to unity indicates that the dominant

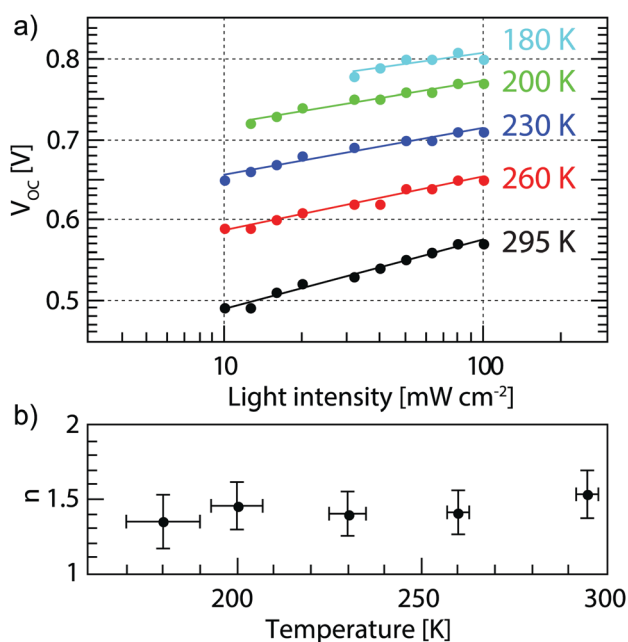


Fig. 2 (a) Light intensity dependent  $V_{OC}$  values (dots) at different temperatures and the fits from eqn (2a) (solid lines) and (b) the ideality factors extracted from (a) showing an approximately constant ideality factor.

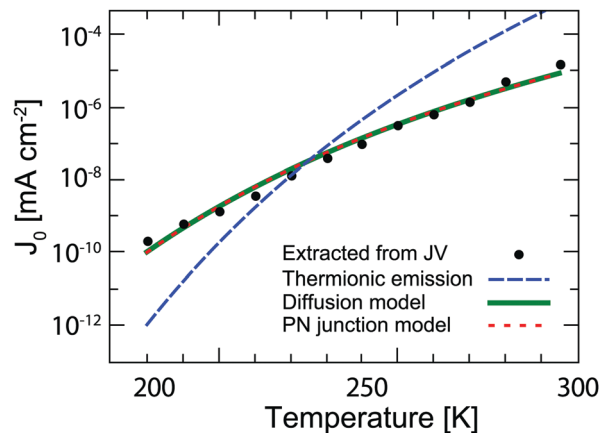


Fig. 3 Temperature dependence of the reverse saturation current extracted from dark JV curves (black dots). The data are fitted with the thermionic emission theory (dashed blue line), diffusion theory (solid green line), and the PN-junction model (dotted red line). Note that the latter two fits virtually overlap.

recombination process is bimolecular, while an ideality factor close to 2 indicates mostly trap-assisted recombination. Thus, from our measurements we can conclude that bimolecular and trap assisted recombination both play a significant role in our devices and that the relative contributions of the two do not change significantly with temperature. Therefore, the only remaining term to explain the increased  $V_{OC}$  as from eqn (2a) is the reverse saturation current, which is known to be temperature dependent. Fig. 3 shows  $J_0$  extracted from the dark JV measurements (Fig. 1d) fitted with the least squares method to the explicit form of the non-ideal Shockley equation using the Lambert-W function.<sup>32</sup> Several models can be used to explain the temperature dependence of  $J_0$ . The most common in the case of QD solar cells are the thermionic emission model<sup>25,33,34</sup> and the diffusion model.<sup>26,35</sup> According to the thermionic emission model,  $J_0$  is given by

$$J_{0,TE} = J_{00} T^2 \exp\left(\frac{-e\phi_B}{kT}\right), \quad (3)$$

where  $J_{00}$  is the effective Richardson's constant, equal to  $120 \text{ A cm}^{-2} \text{ K}^{-2}$ , and  $\phi_B$  is the Schottky barrier height. The diffusion model on the other hand predicts a reverse saturation current given by<sup>36</sup>

$$J_{0,\text{Diff}} = e\mu_{e(h)} N_{C(V)} E_m \exp\left(\frac{-e\phi_B}{kT}\right) \quad (4)$$

where  $\mu_{e(h)}$  is the majority carrier mobility,  $N_{C(V)}$  is the density of states in the conduction (valence) band, and  $E_m$  is the maximum electric field in the device. Finally, the reverse saturation current given by a regular PN junction is given by<sup>36</sup>

$$J_{0,\text{PN}} = eN_C N_V \left( \frac{1}{N_A} \sqrt{\frac{\mu_e}{\tau_e}} + \frac{1}{N_D} \sqrt{\frac{\mu_h}{\tau_h}} \right) \sqrt{\frac{kT}{e}} \exp\left(\frac{-E_g}{kT}\right) \quad (5)$$

where  $N_A$  and  $N_D$  are the acceptor and donor concentrations of the p and n type layers, respectively,  $\tau_e$  and  $\tau_h$  are the minority carrier lifetimes, and  $E_g$  is the bandgap. Without making



assumptions about any of the unknown parameters, the extracted  $J_0$  values are fitted to each of these models in Fig. 3. The diffusion and PN model both provide a good fit within measurement error to the experimental data, while the best fit for the thermionic emission model predicts too high a temperature dependence. The thermionic emission and diffusion models are both based on the presence of a Schottky barrier, which is not present in our devices. Instead, the PN junction model is physically more applicable to the solar cells investigated here. It is important to underline that this model gives us insight into how the reverse saturation current can be decreased in order to increase the  $V_{OC}$  in future solar cells at room temperature. The term in the exponent shows that a trivial method to enhance the  $V_{OC}$  is to increase the bandgap. However, this will simultaneously decrease the amount of photons absorbed and the generated carriers and thus the  $J_{SC}$ . It is also well known that for a high  $V_{OC}$ , recombination should be suppressed. This is reflected in the carrier lifetime terms  $\tau_e$  and  $\tau_h$ , which should be maximized. This can be achieved by reducing the trap density, to which much effort in the colloidal QD field is already being directed. Less attention has so far been directed to increasing the doping concentrations in the p and n layers. This should not only increase the  $V_{OC}$  through the reduction of  $J_0$ , but can help increase the built-in electric field, which facilitates charge extraction. Furthermore, it appears as a linear term in the reverse saturation current as opposed to only a square root dependence of the carrier lifetime, thus we posit that optimizing the doping concentration is a more effective strategy method to further increase the solar cell efficiency of PbS QDs. Finally, from eqn (5) we see that while a high mobility is desirable for charge extraction and thus necessary for high  $J_{SC}$ , it has a weakly detrimental effect on the reverse saturation current, possibly leading to an unavoidable tradeoff between  $V_{OC}$  and  $J_{SC}$ .

## 2.2 Permittivity and mobility

In order to better understand the behaviour of these solar cells at lower temperature, it is important to know whether the permittivity of the material is changing, since the permittivity not only influences recombination rates,<sup>37</sup> but also determines the necessary exciton dissociation energy,<sup>38</sup> plays an important role in the optical field distribution in solar cells,<sup>39</sup> and will subsequently be used in the calculation of the charge carrier mobility. For this, we fabricated Schottky devices with the structure ITO/PEDOT:PSS/PbS/LiF/Al with either EDT or TBAI ligands, and the capacitance of the films is measured at several temperatures between 290 K and 190 K. For the geometric permittivity to be accurately obtained from the capacitance measurements, the device must be fully depleted. For this reason we make these devices thinner ( $\sim 110$ – $150$  nm) than the solar cell active layer, and measure under reverse bias to ensure the device is depleted. To verify this, the capacitance is measured at several voltages and is found not to change significantly between 0 V and  $-0.5$  V, see Fig. S4 (ESI<sup>†</sup>). The permittivity is calculated from the average capacitance in the frequency range of  $10^3$ – $10^4$  Hz, since here the capacitance is well behaved and flat, and using the parallel plate capacitor

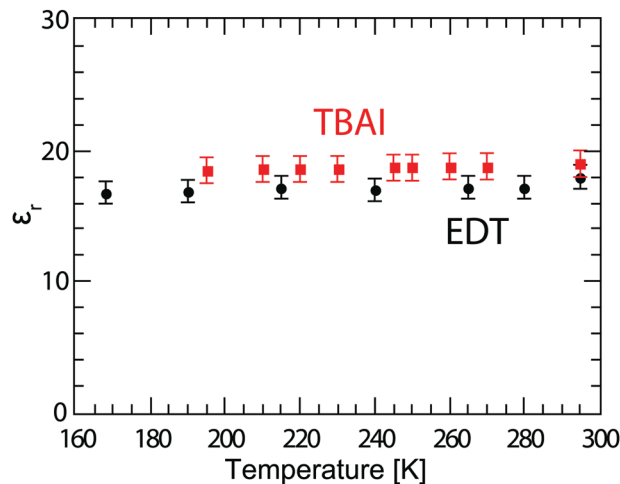


Fig. 4 Permittivity of PbS QD films capped with EDT (black dots) and TBAI (red squares), calculated using the parallel plate model from capacitance measurements averaged over the frequency  $10^3$ – $10^4$  Hz. The error bars represent the measurement uncertainty over a single device, dominated by the uncertainty in film thickness.

model  $C = (A\epsilon_r\epsilon_0)/d$ , where  $A$  is the area defined by the overlap of the electrodes equal to  $0.16$  cm<sup>2</sup> in our devices,  $d$  is the layer thickness, and  $\epsilon_r$  and  $\epsilon_0$  are the relative and vacuum permittivity, respectively. Here we consider the PEDOT:PSS layer to be part of the anode and neglect its contribution to the total capacitance of the device. The permittivity obtained in this way for various temperatures is plotted in Fig. 4. The obtained  $\epsilon_r$  values for EDT decrease slightly from  $18.0 \pm 0.9$  at room temperature to  $16.8 \pm 0.9$  at 168 K. The permittivity of TBAI capped PbS instead remains constant within measurement error at a value of  $18.7 \pm 1.0$ .

Eqn (5) shows that charge carrier mobility, doping concentration, and carrier lifetime are important parameters that determine the overall device performance. It is therefore useful to investigate each of these parameters in turn. To measure the temperature dependent mobility in the active layer, single charge carrier devices are fabricated with EDT and TBAI capped PbS. Unlike field effect transistors (FETs) or Hall effect measurements, which probe the mobility in the plane of the substrate, single charge carrier devices allow the out of plane mobility to be determined, *i.e.* in the same device configuration as the solar cell. Hole only devices were fabricated using device structure ITO/PEDOT:PSS/PbS/MoO<sub>3</sub>/Au, since both PEDOT:PSS and MoO<sub>3</sub>/Au are hole selective contacts aligning favorably to the PbS valence band. Silver or aluminum contacts both allow close alignment with the conduction band of PbS ( $\sim -4.0$  eV), therefore either Ag/PbS/Ag or Ag/PbS/LiF/Al device structures are used for electron only devices.  $JV$  curves of devices exhibiting space charge limited current (SCLC) in the presence of shallow traps show four characteristic regions.<sup>40,41</sup> At low bias, thermally generated charges outnumber the injected charge carriers, and the device follows Ohm's law ( $J \propto V$ ). At a certain voltage  $V_{on}$ , the injected charges exceed the thermally generated charges and they also fill the trap states leading to trap-filling SCLC behaviour:

$$J_{TF,n} = \frac{9}{8} \epsilon_0 \epsilon_r \mu_n \theta \frac{V^2}{d^3} \quad (6)$$



Here,  $\theta$  is the ratio of free charge carriers to the total carrier density,  $V$  is the applied bias, and  $d$  is the device thickness. When all the traps are filled, the device shows an exponential current injection. Since the contacts allow only one charge carrier to be injected into the device, a self-limiting space charge builds up and the device follows Mott–Gurney behaviour

$$J_{MG,n} = \frac{9}{8} \epsilon_0 \epsilon_r \mu_n \frac{V^2}{d^3} \quad (7)$$

The  $JV$  characteristics for hole only devices with EDT capped PbS are shown in Fig. 5a. For these measurements, we perform an analysis similar to the one proposed by Jeong *et al.*<sup>40</sup> Of the four regimes, the Ohmic, injection, and Mott–Gurney regimes are clearly evident; the trap-filling regime is less pronounced, but is taken to be in the region where the log–log slope of the  $JV$  curve equals 2. From the ratio of  $J_{TF,h}$  and  $J_{MG,h}$  we find a  $\theta$  value of 0.70 at room temperature, indicating that approximately 30% of charge carriers are trapped, which increases to 0.98 at 240 K.

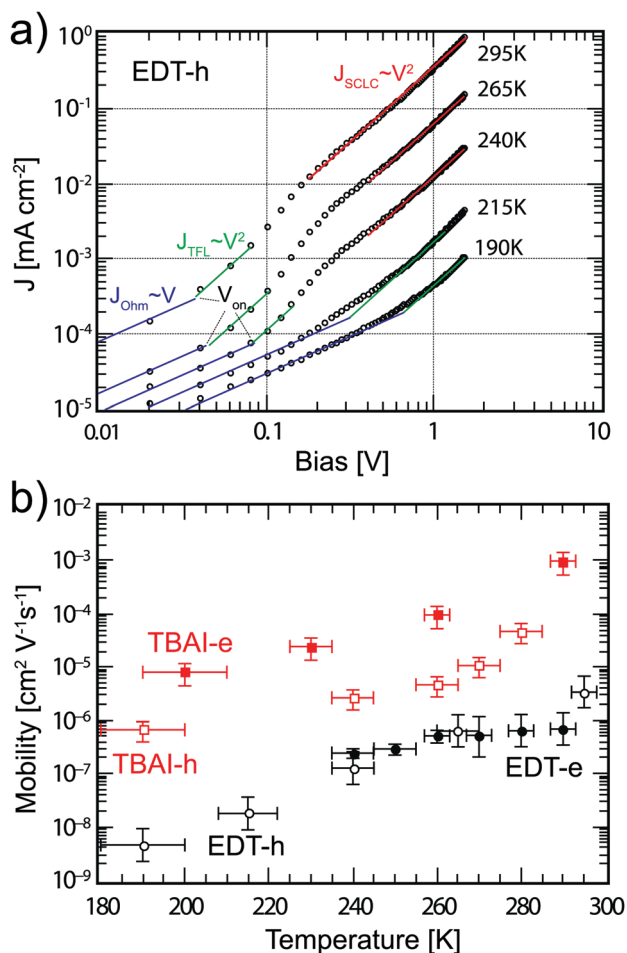


Fig. 5 (a)  $JV$  curves for hole only device with structure ITO/PEDOT:PSS/PbS\_EDT/MoO<sub>3</sub>/Au at various temperatures. The solid lines indicate the Ohmic regime (blue), trap-filling regime (green) and SCLC regime (red). (b) Extracted mobilities from single carrier devices for EDT (black circles) and TBAI (red squares), for both electrons (solid markers) and holes (empty markers). The error bars represent the measurement uncertainty over a single device resulting mainly from the uncertainty in film thickness.

We extract the mobility using the Mott–Gurney region for the temperatures where this is possible. At lower temperatures, only the Ohmic and trap-filling regions can be seen because  $V_{on}$  increases with lower temperature and pushes the Mott–Gurney regime to higher biases outside the range of the  $JV$  curve. For those temperatures, we extract  $\mu_{th}$  from the trap-filling region with the assumption that  $\theta$  remains approximately constant. Since  $\theta$  cannot exceed unity, the error made with this assumption is small. The extracted mobilities are plotted in Fig. 5b. At room temperature, a mobility of  $3 \times 10^{-6} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  is found, lower than the value of  $1 \times 10^4 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  found by Rath *et al.* with the structure ITO/PbS/Au.<sup>42</sup> With decreasing temperature, the mobility decreases almost three orders of magnitude to  $5 \times 10^{-9} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  at 190 K. These results are in agreement with the results found by Jeong *et al.*,<sup>40</sup> though the temperature dependence is slightly more pronounced in our devices. We then perform similar analyses on electron only devices with EDT capped PbS QDs and TBAI capped QDs. The  $JV$  characteristics are shown in Fig. S5 (ESI<sup>†</sup>) while the extracted mobilities are summarized in Fig. 5b. The data show that the EDT capped PbS films are slightly p-type, exhibiting one order of magnitude higher hole mobility than electron mobility at room temperature. In contrast the TBAI capped films display one order of magnitude higher electron mobility indicating slightly n-type behaviour, as has been previously reported in the literature.<sup>19,43</sup> Furthermore, at room temperature the mobility found in TBAI capped films is more than an order of magnitude higher for holes ( $5 \times 10^{-5} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ) and almost four orders of magnitude higher for electrons ( $1 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ) than the one of EDT capped PbS. Both the hole and electron mobility of TBAI capped PbS decreases roughly two orders of magnitude between room temperature and 200 K, while the electron mobility for EDT capped PbS shows the least influence from temperature, decreasing from  $7 \times 10^{-7} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  to  $2 \times 10^{-7} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  between room temperature and 240 K.

Field effect mobilities are also measured for comparison. For this, SiO<sub>2</sub> gated FETs were fabricated by depositing EDT and TBAI capped films with the same recipe as the solar cells and measured in a vacuum. Charge carrier mobility values are extracted from the linear regime according to the gradual channel approximation and using a parallel plate capacitor model for the gate electrode charge accumulation. In contrast to the SCLC devices, the FETs show initially increasing mobility with decreasing temperature, which may indicate band like transport (Fig. S6, ESI<sup>†</sup>). The FET mobility values follow the same trend as previously reported for PbS QDs using the Hall effect and FET measurements,<sup>44</sup> and FET mobilities of CdSe QDs.<sup>45</sup> Hall effect measurements only probe the mobility of the non-trapped charge carriers,<sup>44</sup> and in FETs the Fermi level is raised by the gate bias such that most traps are filled.<sup>6</sup> Therefore, the Hall effect and FET analyses both give mobility values closer to the trap free mobility of PbS QDs, but which are far from the effective mobility in solar cell configuration. The mobility results from our SCLC devices likely indicate a high trap density in PbS QD films that prevents band-like transport under low biasing conditions; given this high trap density, a decrease in



the temperature means less thermal energy is available to re-excite charge carriers from shallow traps back to the transport levels, lowering the overall mobility. Furthermore, we observe that in contrast to the SCLC devices, the FETs with EDT capped PbS display notably higher electron mobility. This can be explained by the vacuum environment used for the temperature dependent FET measurements, since vacuum can reverse the ambient p-doping effect of thiol capped PbS, as observed previously by Balazs *et al.*<sup>18</sup>

### 2.3 Charge carrier lifetimes and diffusion lengths

Charge carrier mobility plays an important role in determining whether or not generated charges can be extracted from solar cells before recombining. It is therefore interesting to note that despite the trend of decreasing mobility, as measured using the single carrier devices, the  $J_{SC}$  value does not decrease very much. The current density in a diode is given by the sum of the drift and diffusion currents,

$$J = qE(n\mu_e + p\mu_h) + q\left(D_e \frac{dn}{dx} + D_h \frac{dp}{dx}\right) \quad (8)$$

where  $n$  and  $p$  are the electron and hole concentration, respectively,  $E$  is the electric field in the device caused by the built-in electric field and the applied bias, and  $D_e$  and  $D_h$  are the electron and hole diffusion coefficients given by  $D_{e/h} = (kT/q)\mu_{e/h}$ . Charges from outside the depletion region must diffuse to the region where there is an electric field to drive their transport. Charges generated in regions further away from the depletion region than their diffusion length recombine before they can be extracted and therefore do not contribute to the photocurrent. To determine the diffusion length in our solar cells, it is necessary to measure not only the mobility, but also the lifetime of photogenerated charge carriers. For this, Schottky devices with structure ITO/PEDOT:PSS/PbS/LiF/Al are fabricated with both EDT and TBAI ligands and the capacitance is measured under 1 Sun illumination under open circuit conditions. Since no current leaves the device at  $V_{OC}$ , all generated charges must recombine, and the characteristic lifetime of the process can be found from the product of the recombination resistance and the chemical capacitance  $\tau_n = R_r C_{\mu}$ .<sup>46,47</sup> The Nyquist plots obtained are shown in Fig. S7 (ESI<sup>†</sup>) and are fitted with the simple three component model shown in the inset. The lifetimes obtained therefrom are displayed in Fig. 6a. It is important to note that these are the lifetimes of the minority carriers in their respective layers, *i.e.*, the electron lifetime for EDT capped PbS and hole lifetime for TBAI capped PbS. EDT capped PbS displays carrier lifetimes of 2  $\mu$ s at room temperature, in close agreement with the values measured by Bawendi *et al.* using transient photovoltage measurements,<sup>24</sup> which increases slightly to 13  $\mu$ s at 215 K. The carrier lifetimes in TBAI capped films are also 2  $\mu$ s at room temperature but do not show significant temperature dependence in the measured range. With the lifetime and mobility known, the diffusion length can easily be calculated from  $L_n = \sqrt{D_n \tau_n}$ , where  $D_n = \frac{kT}{e} \mu_n$  is Einstein's diffusion coefficient, and is shown in Fig. 6b. The electron diffusion length in EDT capped PbS is found to be extremely short,

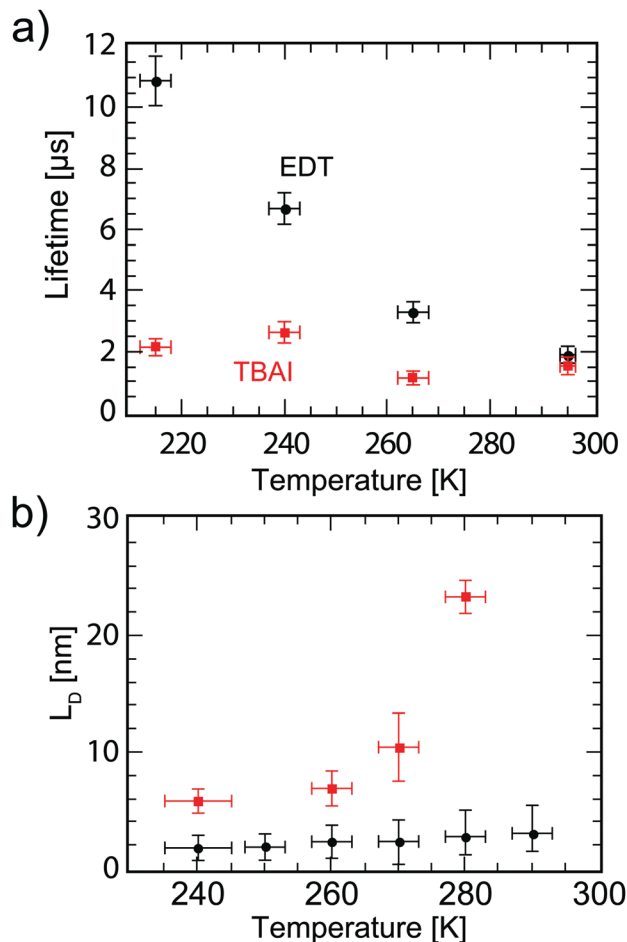


Fig. 6 (a) Carrier lifetime for electrons in EDT capped PbS (black dots), and holes in TBAI capped PbS (red squares) obtained with impedance spectroscopy measurements under 1 Sun illumination under open circuit conditions. (b) Diffusion lengths of electrons in EDT capped PbS (black dots) and holes in TBAI capped PbS (red squares) calculated from  $L_n = \sqrt{D_n \tau_n}$  using the mobilities as reported in Fig. 5b and lifetimes from (a). The error bars represent the measurement uncertainty over a single device following from the uncertainty in mobility and carrier lifetime.

around  $3.0 \pm 1.5$  nm, which is significantly shorter than reported values for PbS coupled with MPA ( $30 \pm 10$  nm) or with MPA +  $\text{CdCl}_2$  ( $80 \pm 10$  nm).<sup>48</sup> Owing to the higher mobility, TBAI capped PbS displays a higher diffusion length at room temperature of  $23 \pm 3$  nm, which decreases to  $5.7 \pm 1$  nm at 240 K. Considering the solar cell thickness of 260 nm and the fact that the  $J_{SC}$  does not drop significantly between room temperature and 240 K (Fig. 1d), the declining diffusion length for TBAI capped PbS and very small diffusion length for EDT capped PbS indicate that the solar cells are practically fully depleted under short circuit conditions.

### 2.4 Doping concentration

The doping concentrations in the n and p layers together determine both the degree of band bending and the distribution of the depletion width  $w$  via  $N_a w_p = N_d w_n$ . In most silicon solar cells, the main absorbing layer has a moderate (usually p-type) doping concentration of  $10^{16}$ – $10^{17}$   $\text{cm}^{-3}$ , followed by a highly



doped (usually n-type) layer. This leads to a large depletion width in the absorbing layer, allowing most photogenerated charges to be swept out of the device without relying on diffusion. To determine the doping concentrations in our solar cells we again investigate Schottky diodes with the structure ITO/PEDOT:PSS/PbS/LiF/Al. From the Mott-Schottky equation,

$$C'^{-2} = \frac{2}{\epsilon_r \epsilon_0 N} \left( V - V_{\text{fb}} - \frac{kT}{e} \right), \quad (9)$$

the doping concentration  $N$  can be obtained from the slope of  $C'^{-2}$  versus the applied bias  $V$ . For this, the capacitance is measured with a 4.6 kHz AC bias with an amplitude of 10 mV superimposed on a DC bias ranging from 0 to 1 V at various temperatures. The Mott-Schottky curves are shown in Fig. S8 (ESI<sup>†</sup>) and the extracted doping concentrations are plotted in Fig. 7a. EDT capped PbS shows a doping concentration of  $(1.0 \pm 0.4) \times 10^{17} \text{ cm}^{-3}$ , which remains essentially constant over the measured temperature range. TBAI capped PbS exhibits a doping concentration at room temperature of  $(4.6 \pm 2.3) \times 10^{16} \text{ cm}^{-3}$ , which decreases to  $(1.5 \pm 0.5) \times 10^{16} \text{ cm}^{-3}$  at 215 K.

With the doping concentration known, the depletion width  $w$  can be calculated by<sup>49</sup>

$$w = \left[ \frac{2\epsilon_r \epsilon_0}{e} \left( \frac{N_A + N_D}{N_A N_D} \right) (V_{\text{bi}} - V) \right]^{\frac{1}{2}}, \quad (10)$$

where  $N_A$  and  $N_D$  are the acceptor and donor concentrations, and  $V_{\text{bi}}$  is the built-in potential. The depletion width is calculated for varying temperature at zero bias and at the maximum operating point, shown in Fig. 7b. At room temperature under short circuit conditions the depletion width is found to be  $234 \pm 15 \text{ nm}$ . Therefore the depletion width plus diffusion length covers almost the whole device thickness of 260 nm. As the temperature is decreased, the depletion region increases leading to full depletion of the device below  $\sim 280 \text{ }^\circ\text{C}$ . This can explain why the short circuit current does not change with decreasing temperature despite the lower mobility. Under maximum power point conditions, the depletion width shows a similar trend but is  $\sim 30\%$  lower, such that  $w$  is only  $168 \pm 17 \text{ nm}$  at room temperature, which limits the collection efficiency of charges outside the depletion region and explains why  $J_{\text{MPP}}$  is significantly lower than the  $J_{\text{SC}}$ , leading to a relatively low FF in these devices.

For depletion widths smaller than the device thickness, it is important to optimize the distribution of the depletion width over the pn-junction. Given that the mobility of TBAI capped PbS is significantly higher than in the EDT capped layer, it is desirable to use the TBAI layer as the main absorber layer, and to limit the thickness of the EDT layer. With these doping concentrations, 64% of the depletion region is located on the TBAI side at room temperature, while at lower temperatures the depletion width shifts increasingly to the TBAI layer, such that 93% of the depletion area is located on the TBAI side at 125 K. This, together with the increasing depletion width could explain the initial increase in the fill factor between room temperature and 260 K, while at lower temperatures the decreasing mobility dominates the fill factor behaviour. Thus, to improve the efficiency

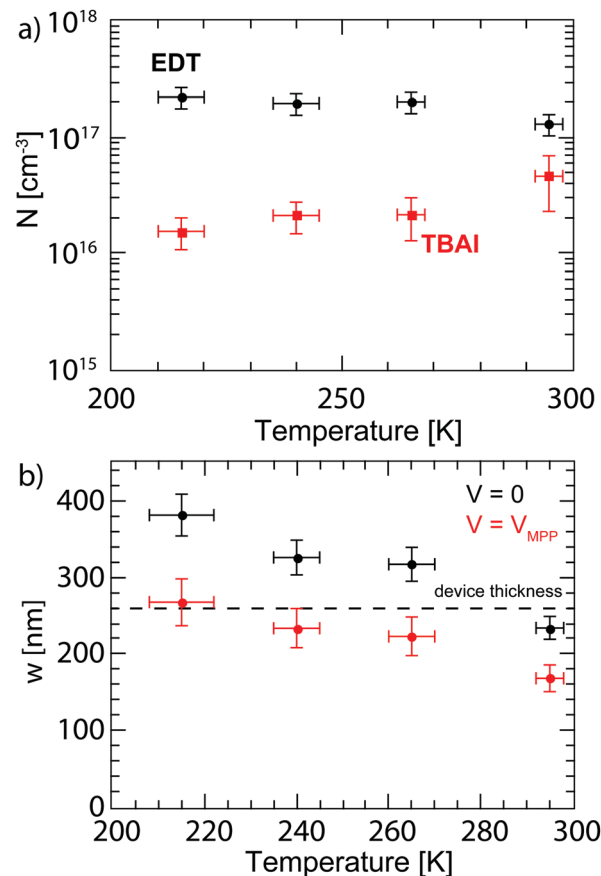


Fig. 7 (a) Doping concentration of EDT (black dots) and TBAI (red squares) capped PbS measured by Mott-Schottky analysis. The error bars are determined by the uncertainty in permittivity, slopes of the Mott-Schottky curves, and temperature. (b) Depletion width calculated at zero bias (black) and at the maximum power point (red). The dashed line shows the device thickness. The error bars originate from the uncertainty in the permittivity and doping concentrations.

at room temperature, an order of magnitude higher doping of the p-type layer is required such that practically all of the depletion width falls in the TBAI layer.

### 3 Conclusion

We have fabricated highly efficient PbS QD solar cells and explored their temperature dependent properties. The  $V_{\text{OC}}$  is found to be governed solely by the reverse saturation current, which can be explained using the PN-junction model. Based on this model, we propose that increasing the doping levels in the PN-junction structure is a promising method for increasing the  $V_{\text{OC}}$  in future QD based solar cells. In addition, the doping concentration of the p-type layer should be at least 1 order of magnitude higher than the n-type layer for a favourable depletion width distribution across the junction. Moreover, we have determined all the most important physical parameters that determine the working mechanism of our solar cells. Electron and hole mobilities were measured for both EDT and TBAI capped QD films, and were found to decrease with decreasing temperature.



TBAI capped QDs showed electron and hole mobility values of  $1 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $5 \times 10^{-5} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ , respectively, which decrease by roughly two orders of magnitude as temperature is decreased to 200 K. EDT capped films show significantly lower electron and hole mobility values of  $7 \times 10^{-7} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  and  $3 \times 10^{-6} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  respectively. Very short diffusion lengths were observed in EDT capped PbS independent of temperature between 290 K and 240 K, while TBAI capped films showed diffusion lengths decreasing from  $26 \pm 3 \text{ nm}$  to  $6 \pm 1 \text{ nm}$ . In spite of these short diffusion lengths the current output of our devices at  $J_{\text{SC}}$  is almost temperature independent. This is explained by the depletion region, which is found to extend almost throughout the active layer of the solar cells at room temperature and to extend completely throughout the device at lower temperatures. Thus, these solar cells are not limited by diffusion, instead they rely on drift-dominated transport for charge carrier extraction.

## 4 Experimental methods

### Device fabrication

TiO<sub>2</sub> is formed by spincoating a 20:2:1 solution of ethanol: titanium(IV) butoxide:HCl onto pre-patterned fluorine doped tin-oxide substrates ( $13 \Omega \text{ sq}^{-1}$ ) and annealed at 450 °C for 30 min. Oleic acid capped PbS QDs are synthesized by a previously described hot injection method.<sup>23,46</sup> QD films are made in a N<sub>2</sub> glovebox by layer-by-layer spincoating a 10 mg ml<sup>-1</sup> solution of PbS in hexane and exposing the film to either EDT (0.01% w in acetonitrile) or TBAI (15 mg ml<sup>-1</sup> in methanol) for 30 s before spin drying. The film is then washed twice in the solvent of the ligand and spin-dried. This process is repeated until the desired thickness is achieved. Under  $< 10^{-6}$  mBar vacuum, the solar cells are finished by thermal evaporation of 5 nm MoO<sub>3</sub> and 80 nm Au at a rate of  $0.2 \text{ \AA s}^{-1}$  and  $2 \text{ \AA s}^{-1}$ , respectively. In total, 30 solar cells with this structure were fabricated on 12 different substrates and the analysis reported here was performed on the champion device. For Schottky devices the active layer is deposited on top of 90 nm poly(3,4-ethylenedioxythiophene)-poly(styrenesulfonate) (PEDOT:PSS; VP Al4083, H. C. Stark) spincoated from water and annealed at 120 °C for 10 min. The Schottky devices are finished by thermal evaporation of 1 nm LiF ( $0.1 \text{ \AA s}^{-1}$ ) and 100 nm Al ( $0.5\text{--}10 \text{ \AA s}^{-1}$ ). The device area is  $0.16 \text{ cm}^2$  as defined by the overlap of electrodes.

### Current-voltage characterization

*JV* sweeps were carried out in an inert environment using a Keithley 2400 source-meter. Simulated AM1.5G illumination was provided by a Steuernagel Solarconstant 1200 metal halide lamp set to  $100 \text{ mW cm}^{-2}$  intensity, measured using a silicon reference cell (SRC-1000-RTD-QZ, VLSI Standards Inc.) and corrected for the spectral mismatch.<sup>50</sup> The illuminated area was limited to  $0.10 \text{ cm}^2$  by a well-defined shadow mask for efficiency calculations. The temperature was controlled by an adjustable N<sub>2</sub> gas flow through a liquid N<sub>2</sub> bath.

### Impedance spectroscopy

Using a SP-200 Bio-Logic potentiostat a forward bias is superimposed with a 30 mV AC perturbation over the frequency range of 1 MHz to 100 Hz. For permittivity measurements, the capacitance is measured in the dark at  $-0.5 \text{ V DC}$  bias. For carrier lifetime measurements, the device is held at open circuit bias under 1 Sun illumination.

### EQE measurements

The EQE is measured under short circuit conditions under monochromatic light. As a white light source a 250 W quartz tungsten halogen lamp (6334NS, Newport) with lamp housing (67009, Newport) is used. Monochromatic light is obtained using narrow band pass filters (Thorlabs) with a full width half maximum (FWHM) of  $10 \pm 2 \text{ nm}$  from 400 nm to 1300 nm and a FWHM of  $12 \pm 2.4 \text{ nm}$  from 1300 nm to 1400 nm. The light intensity is determined by calibrated PD300 and PD300IR photodiodes (Ophir Optics).

### PL measurements

PbS QD films were deposited on quartz substrates using the layer-by-layer technique mentioned above. The samples were excited at 400 nm by the second harmonic of a modelocked Ti:sapphire (Mira 900) laser delivering pulses of 150 fs. An optical pulse selector was used to vary the repetition rate of the exciting pulse. All measurements were performed in an optical cryostat, loaded inside a glovebox to maintain an oxygen-free environment at all times.

## Acknowledgements

M. J. Speirs and M. A. Loi acknowledge the financial support of the Alumnikring Den Haag/Rotterdam through the Ubbo Emmius Fund of the University of Groningen. M. A. L. acknowledges also the support of the ERC Starting Grant "Hybrids Solution Processable Optoelectronic Devices" (Hy-SPOD) (ERC-306983); M. V. Kovalenko acknowledges partial financial support from ERC Starting Grant NANOSOLID (GA No. 306733).

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