



Cite this: *Dalton Trans.*, 2016, **45**, 7158

Received 12th February 2016,  
Accepted 8th March 2016

DOI: 10.1039/c6dt00588h

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# Isomerization of the osmium–tellurium cluster $\text{Os}_3(\mu\text{-TeR})_2(\text{CO})_{10}$ : a kinetic and computational study†

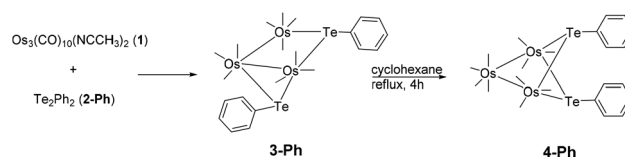
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The kinetics for the isomerization of the 50e cluster  $\text{Os}_3(\mu\text{-TeTol-}p)_2(\text{CO})_{10}$  (**3**), where the tellurides bridge two different Os–Os edges, to one in which the tellurides bridge the same open Os...Os edge (**4**) have been measured experimentally by  $^1\text{H}$  NMR spectroscopy. The determined activation parameters are  $\Delta H^\ddagger = 77 \pm 9 \text{ kJ mol}^{-1}$  and  $\Delta S^\ddagger = -12 \pm 28 \text{ J mol}^{-1} \text{ K}$ . The conversion of **3** to **4** has been computationally investigated by electronic structure calculations using the model compound  $\text{Os}_3(\mu\text{-TeMe})_2(\text{CO})_{10}$ . The computed isomerization pathway is consistent with the kinetic data and the thermodynamic preference for the product stereoisomer that possesses a slipped, eclipsed conformation for the two *p*-tolyl groups.

## Introduction

An interesting aspect of organometallic cluster chemistry is the possibility of fluxionality and isomerization. These could involve the ligand sphere or the metal core. Such processes involving heavier main group elements are expected to be less facile, often because of unfavourable factors such as their higher atomic mass and high inversion barrier. Few studies on the kinetics of such processes exist, partly because of the lack of suitable systems.<sup>1</sup> A system suitable for such studies is the isomeric pair of di-telluride substituted clusters **3-Ph** and **4-Ph**, with the molecular formula  $\text{Os}_3(\mu\text{-TePh})_2(\text{CO})_{10}$ , which were obtained from the reaction of  $\text{Os}_3(\text{CO})_{10}(\text{NCCH}_3)_2$  (**1**) with diphenyl ditelluride (**2-Ph**).<sup>2</sup> To our knowledge, that remains the only report on osmium–tellurium carbonyl clusters, in contrast to the plethora of reports on the chemistry of the congeners sulfur and selenium, and of their neighbour in the periodic table, antimony.<sup>3</sup> In that study, it was shown that **3-Ph** isomerized to **4-Ph** in refluxing cyclohexane (Scheme 1).

Cluster **4-Ph** adopts an unusual conformation in which the two phenyl rings are parallel to each other and stacked in a slipped, eclipsed orientation (**A1**). In principle, three different stereoisomers are possible, but the other two stereoisomers are not observed (Fig. 1). The important points that we wish to address in this report are: (i) the mechanism associated with



Scheme 1 Synthesis and isomerization of  $\text{Os}_3(\mu\text{-TePh})_2(\text{CO})_{10}$ .

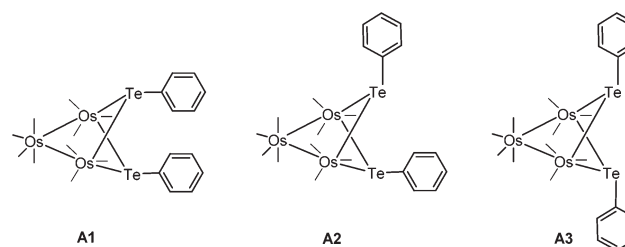


Fig. 1 Different stereoisomers of **4-Ph**.

the transformation of **3** to **4**, and (ii) the relative stability of the three stereoisomers of **4** and their role, if any, in the isomerization reaction.

## Results and discussion

In order to study the isomerization, we prepared the *p*-tolyl analogs  $\text{Os}_3(\mu\text{-Te-}p\text{-tolyl})_2(\text{CO})_{10}$ , **3-Tol-}p** and **4-Tol-}p**, using the same synthetic route previously reported;<sup>2</sup> the singlet  $^1\text{H}$  resonances for the methyl groups would enable the straightforward monitoring of the isomerization by NMR spectroscopy. The structure of **4-Tol-}p** was also confirmed *via* a

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†Electronic supplementary information (ESI) available. CCDC 1421965 for **4-Tol-}p**. For ESI and crystallographic data in CIF or other electronic format see DOI: 10.1039/c6dt00588h



single-crystal X-ray diffraction study to have the tolyl groups arranged as in the isomer **A1**; the ORTEP plot depicting its molecular structure, together with selected bond parameters, is given in Fig. 2. A variable-temperature  $^1\text{H}$  NMR experiment conducted up to 90  $^\circ\text{C}$  showed no evidence of isomerization or the presence of any of the other isomers, thus suggesting that the stereoisomer **A1** is thermodynamically the most stable.

This is also supported by a computational study, using density functional theory (DFT), on the three possible isomers. To facilitate the computations and the subsequent study on the mechanism of the isomerization of **3**  $\rightarrow$  **4**, we replaced the *p*-tolyl groups in **4-Tol-p** with methyl groups; the optimized structures are depicted in Fig. 3. The lowest energy isomer is **A1** where the two methyl groups adopt an equatorial-like disposition with respect to the  $\text{Os}_3$  core and oriented similar to the aryl groups in the solid-state structures of clusters **4-Ph** and **4-Tol-p**. The bond distances and angles in the optimized structure of **A1** are in good agreement with the solid-state structures of **4-Ph** and **4-Tol-p**. Isomer **A2** lies 7.0  $\text{kJ mol}^{-1}$  above **A1**, and the destabilization can be attributed to the proximity of the axial methyl group to the  $\text{Os}(\text{CO})_4$  moiety. The energetic penalty for two axial methyl groups is severe, with **A3** lying 40.5  $\text{kJ mol}^{-1}$  above **A1**.

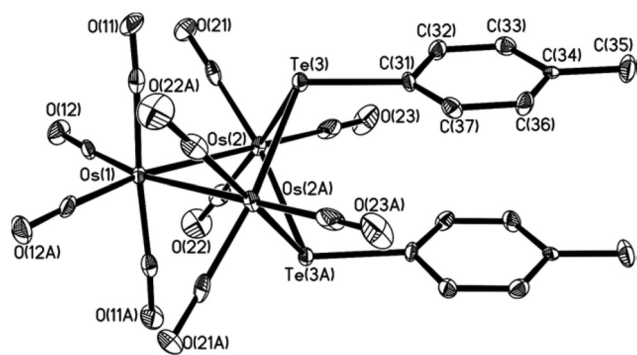


Fig. 2 ORTEP diagram of cluster **4-Tol-p**. Thermal ellipsoids are drawn at the 50% probability level. All H atoms have been omitted for clarity. Selected bond lengths (Å) and angles ( $^\circ$ ):  $\text{Os1-Os2} = 2.9343(4)$ ;  $\text{Os2-Te3} = 2.7330(6)$ ;  $\text{Os2-Te3A} = 2.7375(6)$ ;  $\text{Os2}\cdots\text{Os2A} = 3.607$ ;  $\text{Os2-Te3-Os2A} = 82.509(15)$ ;  $\text{Os2-Os1-Os2A} = 75.859(14)$ .

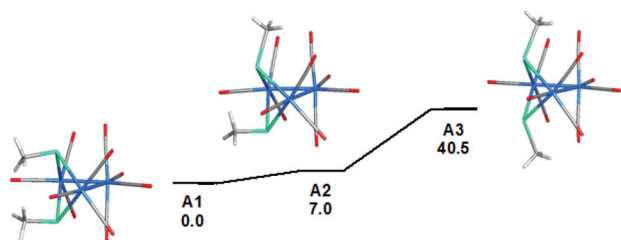


Fig. 3 Geometry-optimized M06 structures and ground-state energy ordering for the stereoisomeric  $\text{Os}_3(\mu\text{-TeMe})_2(\text{CO})_{10}$  clusters **A1**–**A3**. The reported free energies are in  $\text{kJ mol}^{-1}$  relative to **A1**.

The charge distribution and bonding in species **A1**–**A3** were also examined, and the data are consistent with our published results on related 50e clusters (Table S2†).<sup>4</sup> In brief, the osmium atoms and the  $\mu\text{-TeMe}$  ligands have opposing natural charges ( $\sim -1.5$  and 0.8, respectively). The formal Os–Os and Os–Te single bonds are characterized by Wiberg bond indices (WBI) of  $\sim 0.41$  and 0.76, respectively, while the values for the corresponding nonbonding vectors are  $\sim 0.11$  and 0.06, respectively. Elongation of the  $\text{Os}_1\text{--Os}_3$  vector in **A1**–**A3** follows from the Polyhedral Skeletal Electron Pair (PSEP) theory, the bonding in the cluster core being derived from seven skeletal electron pairs.<sup>5</sup>

The kinetics for the isomerization of **3-Tol-p** to **4-Tol-p** were followed by monitoring the methyl resonances in the  $^1\text{H}$  NMR spectrum; the former shows two distinct resonances at 2.379 and 2.362 ppm while the latter shows a single resonance at 2.289 ppm. The rate constants at five different temperatures were determined (Fig. 4), and an Eyring plot gave the kinetic parameters as  $\Delta H^\ddagger = 77 \pm 9 \text{ kJ mol}^{-1}$  and  $\Delta S^\ddagger = -12 \pm 28 \text{ J mol}^{-1} \text{ K}$ , which corresponds to a  $\Delta G^\ddagger$  of  $80 \pm 17 \text{ kJ mol}^{-1}$  at 298 K.

This isomerization reaction was investigated computationally with the methyl analog of cluster **3**, viz. **B**. The geometry-optimized structure of **B** (Fig. S7†) is in good agreement with the experimentally determined structure of **3-Ph**.<sup>2</sup> The proposed pathway from **B** to the final product **A1** is depicted in Scheme 2, together with the free energy change (in  $\text{kJ mol}^{-1}$ ) for each step; the potential energy surface for the isomerization is depicted in Fig. 5.

The first step involves the conversion of **B** to an isomer **D** via the intermediate species **C** which contains an  $\eta^1\text{-TeMe}$  group. While the bridging  $\text{Te}_2$  moiety that spans the non-bonded  $\text{Os}_1\cdots\text{Os}_3$  vector in **B** contains a single stereochemically-active lone pair, the tellurium centre in the  $\eta^1\text{-TeMe}$  group in **C** possesses two stereochemically-active lone pairs. Species **C** lies 90.6  $\text{kJ mol}^{-1}$  above **B**, and points to the thermodynamic preference of a bridging  $\text{TeMe}$  moiety over the  $\eta^1$  mode.<sup>6</sup> The transition state **TSBC** for the conversion from **B** to **C** has been located and lies 103.0  $\text{kJ mol}^{-1}$  above **B**; the

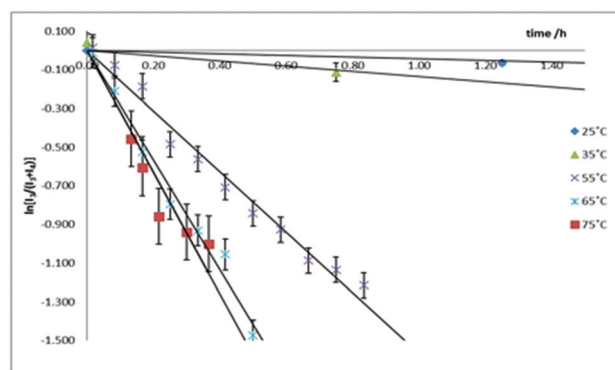
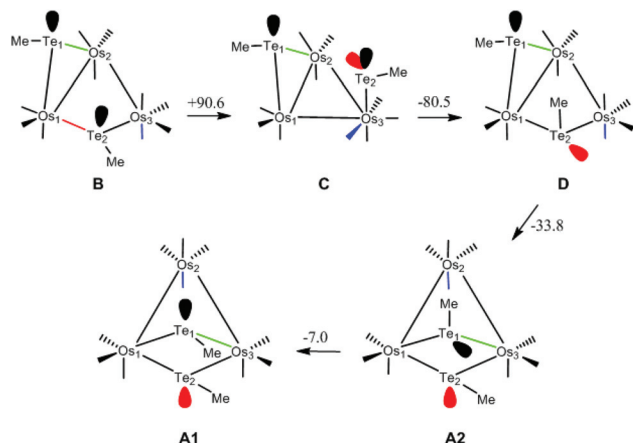
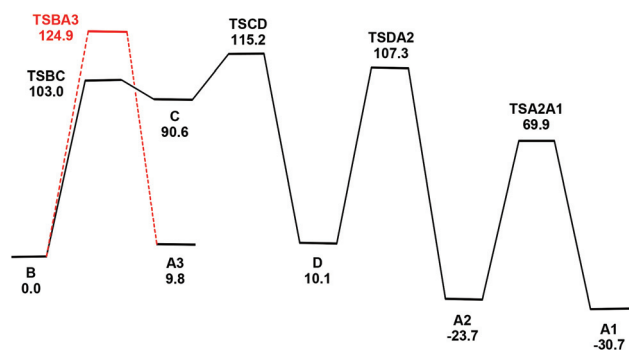


Fig. 4 Plots of  $\ln[I_3/(I_3 + I_4)]$  against time at various temperatures. The two lowest temperature plots have been truncated, and all the plots shifted to a common origin.





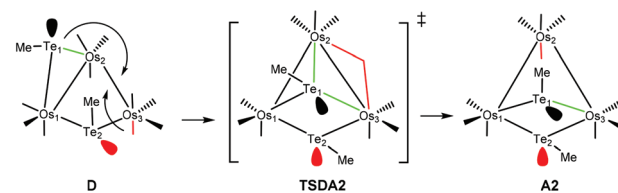
**Scheme 2** Isomerization of **B** to **A1** with atom labels and the stereochemically active lone pair(s) at each tellurium centre highlighted. Energy values given for each step are  $\Delta G$  (kJ mol<sup>-1</sup>).



**Fig. 5** Potential energy surface for the isomerization of **B** to **A1** and the conversion of **B** to **A3**. Energy values are  $\Delta G$  in kJ mol<sup>-1</sup> with respect to **B**.

internuclear distance between the non-bonded Os<sub>1</sub> and Os<sub>3</sub> atoms in **B** contracts by 1.054 Å as a bond develops between them.

Rotation of the  $\eta^1$ -TeMe ligand in **C**, coupled with expansion of the adjacent Os<sub>1</sub>–Os<sub>2</sub> bond and reformation of the  $\mu$ -TeMe ligand using the original lone pair on Te<sub>2</sub> in **B**, furnishes **D**. This transformation corresponds formally to an exchange of the electron pair in the Os<sub>1</sub>–Te<sub>2</sub> bond with the lone pair on the Te<sub>2</sub> centre.<sup>7</sup> The transition state **TSCD** for this step lies 115.2 kJ mol<sup>-1</sup> above **B** and represents the rate-limiting step for the overall process. The stereoisomers **B** and **D** differ in the disposition of their methyl groups. The methyl group on the Te<sub>2</sub> centre in **D** is in close contact with the Te<sub>1</sub> centre and the adjacent axial CO ligand on the Os(CO)<sub>4</sub> moiety. Accordingly, **D** lies 10.1 kJ mol<sup>-1</sup> above **B**. We have also checked the possible inversion at the Te<sub>2</sub> centre that would give **D** directly from **B**; this leads to an increase in energy above 165 kJ mol<sup>-1</sup> and without a feasible transition state. This rise in energy can be attributed to the eclipsing interactions between the methyl group and the two adjacent equatorial CO groups as the former migrates to its final orientation in **D**.



**Scheme 3** Conrotation of the CO and Te<sub>1</sub> groups about the Os<sub>2</sub>–Os<sub>3</sub> vector in **D** through the transition state **TSDA2** to furnish **A2**.

The conversion of **D** to **A2** proceeds *via* a merry-go-round of five CO groups and the  $\mu$ -Te<sub>1</sub> ligand, with the latter migrating across the non-tellurium-ligated Os<sub>2</sub>–Os<sub>3</sub> vector (Scheme 3).<sup>8</sup> The optimized structure of the transition state **TSDA2**, which lies 107.3 kJ mol<sup>-1</sup> above **B**, clearly shows the presence of the capping  $\mu_3$ -Te<sub>1</sub> and bridging  $\mu_2$ -CO groups; the Wiberg bond indices for the Os<sub>1</sub>–Te<sub>1</sub> (0.79), Os<sub>2</sub>–Te<sub>1</sub> (0.44), and Os<sub>3</sub>–Te<sub>1</sub> (0.50) vectors are consistent with the face-capping TeMe moiety.

The final step in the reaction involves a formal inversion of the Te<sub>1</sub> centre in **A2**, through the transition state **TSA2A1**. The energy barrier for this inversion is relatively low, lying only 46.2 kJ mol<sup>-1</sup> above **A2**. The sum of the angles around the Te<sub>1</sub> (*ca.* 360°) and Te<sub>2</sub> (*ca.* 293°) centres in **TSA2A1** are consistent with the inversion at Te<sub>1</sub>. The overall reaction of **B** → **A1** is exergonic by 30.7 kJ mol<sup>-1</sup>.

Migration of the Te<sub>1</sub> moiety across the Os<sub>2</sub>–Os<sub>3</sub> bond in **B** to give **A3** was also examined. This process, which involves a merry-go-round of CO and the Te<sub>1</sub> ligands, is analogous to that for the **D** → **A2** transformation. The transition state **TSBA3** for this presents a barrier of 124.9 kJ mol<sup>-1</sup>, which is 9.7 kJ mol<sup>-1</sup> higher than the rate-limiting step (**TSCD**) in the four-step pathway outlined above. While not evaluated computationally due to the unfavorable thermodynamics associated with **TSBA3**, a stepwise inversion sequence at each tellurium centre in **A3** is expected to yield **A2** and subsequently **A1**.

## Conclusions

In this work, we have investigated the polyhedral rearrangement attendant in the conversion of cluster **3** to **4** by NMR spectroscopy and determined the activation parameters for the process. The reaction pathway has been examined computationally with the model cluster Os<sub>3</sub>( $\mu$ -TeMe)<sub>2</sub>(CO)<sub>10</sub>, and the rate-limiting step is shown to involve a merry-go-round in which a  $\mu$ -TeMe ligand migrates across an Os–Os bond to an open Os...Os edge in the cluster, coupled with the permutation of five CO ligands around the same Os–Os bond.

## Experimental

All reactions and manipulations were carried out under an argon atmosphere using standard Schlenk techniques. All solvents employed in this study were distilled over an appropriate



drying agent under argon before use. Infrared (IR) spectra were recorded in a solution IR cell with NaCl windows and a path length of 0.1 mm, at a resolution of 2 cm<sup>-1</sup>. <sup>1</sup>H NMR spectra were recorded on a Bruker AV300 at 300 MHz while variable temperature <sup>1</sup>H NMR spectra were recorded on a JEOL ECA 400 at 400 MHz, referenced to the residual proton resonance of the solvent. Compounds **1** and **2** were synthesized according to literature methods.<sup>2,9</sup>

### Kinetics measurements

Cluster **3-Tol-p** (8 mg, 0.0062 mmol) was dissolved in C<sub>6</sub>D<sub>6</sub> and 1,3,5-trimethoxybenzene was added as an internal standard. The solutions were placed in an oil bath set at various temperatures (25 °C, 35 °C, 55 °C or 75 °C) and the conversion of **3-Tol-p** was monitored *via* <sup>1</sup>H NMR spectroscopy. The mass balance was checked by plotting (*I*<sub>3</sub> + *I*<sub>4</sub>) against the reaction time (*t*), where *I*<sub>3</sub> and *I*<sub>4</sub> are the scaled, integrated intensities of **3-Tol-p** and **4-Tol-p**, respectively. The rate constant (*k*) at each temperature was obtained from a plot of  $\ln\left(\frac{I_3}{I_3 + I_4}\right)$  against *t*. An Eyring plot was then obtained to obtain Δ*H*<sup>‡</sup> and Δ*S*<sup>‡</sup>.

### X-ray crystal structure determination

A crystal of **4-Tol-p** was mounted on a quartz fiber. X-ray data were collected on a Bruker AXS APEX system, using Mo Kα radiation, with the SMART suite of programs.<sup>10</sup> The data were processed and corrected for Lorentz and polarization effects with SAINT,<sup>11</sup> and for absorption effects with SADABS.<sup>12</sup> Structural solution and refinement were carried out with the SHELXTL suite of programs.<sup>13</sup> The structure was solved by direct methods, followed by difference maps to complete the structure for all the non-hydrogen atoms. The structure was found to be a racemic twin. Organic hydrogen atoms were placed in the calculated positions and refined with a riding model. All non-hydrogen atoms were generally given anisotropic displacement parameters in the final model.

### Computational studies

All DFT calculations were carried out with the Gaussian 09 package of programs,<sup>14</sup> using M06 as the DFT functional.<sup>15</sup> The osmium and tellurium atoms were described with the Stuttgart–Dresden effective core potential and SDD basis set,<sup>16</sup> and the 6-31G(d') basis set was employed for all the remaining atoms.<sup>17</sup>

All the reported geometries were fully optimized, and analytical second derivatives were evaluated at each stationary point to verify whether the geometry was an energy minimum (positive eigenvalues) or a transition structure (one negative eigenvalue). Unscaled vibrational frequencies were used to make zero-point and thermal corrections to the electronic energies. The resulting free energies are reported in kJ mol<sup>-1</sup> relative to the specified standard. All transition states on the potential energy surface were evaluated by intrinsic reaction coordinate (IRC) calculations, in order to establish the reactant and product species associated with each transition-state struc-

ture. The natural charges and Wiberg bond indices were computed using Weinhold's natural bond orbital (NBO) program, as executed by Gaussian 09.<sup>18</sup> Geometry-optimized structures have been drawn with the JIMP2 molecular visualization and manipulation program.<sup>19</sup>

## Acknowledgements

This work was supported by Nanyang Technological University, the Ministry of Education (Research Grant No. M4011158), and the Robert A. Welch Foundation (Grant B-1093-MGR). Computational resources through the High-Performance Computing Services and CASCAM at the University of North Texas are acknowledged. We thank Dr Rakesh Ganguly for assistance with the X-ray crystallographic data collection and processing and Prof. Michael B. Hall (TAMU) for providing us a copy of his JIMP2 program, which was used to prepare the geometry-optimized structures reported here.

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