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Dipolar molecules inside C₇₀: an electric field-driven room-temperature single-molecule switch†

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We propose a two-state electric field-driven room-temperature single-molecule switch based on a dipolar molecule enclosed inside ellipsoidal fullerene C₇₀. We show that the two low-energy minima of the molecular dipole inside the C₇₀ cage provide distinguishable molecular states of the system that can be switched by application of an external electric field.

Current technologies employ about one million atoms for storing a single bit of data.¹ Recently, numerous attempts have been made to design and manufacture high-density memory devices which would use just a few atoms to store one bit.² However, a shortcoming of such devices is that they remain operative only at very low temperatures, *e.g.*, up to 25 K. In this paper we suggest an alternative principle for data storage on the basis of a neutral dipolar molecule encapsulated inside a fullerene. We demonstrate that the dipolar system inside a fullerene can be manipulated by means of an external electric field (EF) that is within the range of electric fields generated by devices, such as a scanning tunneling microscope (STM), which can be as high as 0.5 au ($2.57110326 \times 10^{12} \text{ V m}^{-1}$).³ The proposed single-molecule switch may store information on the basis of the orientation of the dipole@fullerene. The hypothesized systems are stable and remain operative under room-temperature conditions.

Endohedral fullerene science has gone through rapid development in the past decade.⁴ To the best of our knowledge, the idea of manipulating a neutral dipole enclosed in a fullerene has not been pursued before. Several single-molecule switches based on endohedral fullerenes have been proposed.⁵ Delaney and Greer studied the possibility of manipulating endohedral Li⁺ by EF inside C₆₀.⁶ They noted that fullerene behaves as a partial Faraday cage, about 25% of the EF penetrates inside. Yasutake *et al.* have

shown that the Tb@C₈₂ molecule can be oriented on a gold surface by electric field at 13 K.⁷ Electric current-driven rotation of the endohedral Sc₃N cluster in C₈₀ fullerene has been studied at 4.7 K.⁸ A single-molecule magnetic switch, based on turning on/off the EPR spectra of Sc₃C₂@C₈₀, has been reported.⁹ Levitt *et al.* proposed switches based on *ortho*- and *para*-water molecules encapsulated in C₆₀.^{9,10}

A schematic model of the dipole@fullerene switch proposed here is shown in Fig. 1. The selected fullerene is ellipsoidal, so the internal potential provides two energy minima (states) for the enclosed dipolar molecule, separated by an energy barrier. The barrier should be large enough to prevent the enclosed moiety from spontaneous switching *via* thermal motion but not too large so that it can be overcome by the interaction of the molecular dipole with an external EF.

An ideal dipolar molecule for building the dipole@fullerene should have a large dipole moment and should be stable under high-temperature conditions that are normal for fullerene synthesis.^{4b} Moreover, the harder (in chemical sense) the endohedral dipole, the higher the stability of the system in high-EF with respect to potential electron transfer from the enclosed molecule to the fullerene cage that is inherently a good electron acceptor.^{4b}

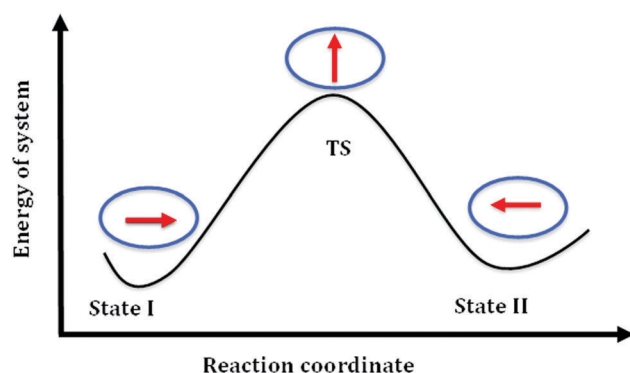


Fig. 1 Schematic outline of the suggested endohedral single-molecule switch.

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A few fullerenes with neutral dipolar molecule(s) inside their cage have been experimentally prepared so far. These include $\text{H}_2\text{O}@C_{60}$,¹¹ $(\text{H}_2\text{O})_2@C_{70}$,¹² $\text{HF}@C_{60}$,¹³ and open-cage fullerenes with CO and H_2O enclosed.¹⁴ A number of systems have been investigated theoretically, such as LiF , BN , HCN , H_2CO , NH_3 , CO_2 , and similar molecules enclosed in C_{60} or C_{70} .¹⁵ In a pre-screening study (not discussed here) we investigated various endohedral systems, such as $\text{LiF}@C_{70}$, $\text{CO}@C_{70}$, $\text{CH}_3\text{F}@C_{70}$, $\text{CH}_3\text{CN}@C_{70}$, and $\text{CH}_2\text{O}@C_{70}$.^{15a,b} The energy barrier for turning these molecules inside the C_{70} cage without an EF was found to be either too low ($< 3 \text{ kcal mol}^{-1}$) or too large ($> 50 \text{ kcal mol}^{-1}$ for $\text{CH}_3\text{CN}@C_{70}$). The barrier was found to be proportional to the size of the molecule. Here we focus on $\text{MX}@C_{70}$ -based systems ($\text{MX} = \text{LiCl}$, NaF , NaCl), which showed the best properties among the tested systems, having the right size and large dipole moments ($> 7.0 \text{ D}$). We note that all alkali metal halides are known to exist in the form of diatomic molecules in the gas-phase at high temperatures necessary for the preparation of fullerenes.¹⁶ Furthermore, encapsulation of MX molecules in C_{70} is an exothermic process, see Table S1 (ESI†). Another potentially possible way to prepare such molecules is *via* chemical opening and closing of a fullerene cage.¹⁴

Building a real device ultimately requires that the fullerene is attached to a surface.¹⁷ Because linkers including N-heterocyclic carbenes¹⁸ and disulfides¹⁹ on a gold surface are prone to dissociation under high-EF conditions (due to their bond polarity), we propose to use a nonpolar C-C linker. A hypothetical target system is a graphene sheet, decorated with $\text{MX}@C_{70}$ systems *via* the Diels-Alder reaction.²⁰

A system in which C_{70} is attached to a hexane ring²¹ serves as the smallest model of the fullerene bonded to a graphene sheet, see Fig. 2. We limited our study to the thermodynamically most-stable Diels-Alder product.²¹ It is worth noting that the EF can catalyze the Diels-Alder reaction^{20,22} and potentially control the regioselectivity of the reaction between $\text{MX}@C_{70}$ and the graphene surface.²³ The reactivity of $\text{MX}@C_{70}$ in the presence of an EF is the subject of our ongoing research to be published subsequently.

Our computations[‡] show that the MX aligns nearly along the main axis of symmetry of C_{70} to minimize the steric clash with the belt region, as expected. Without an EF there are two local minima (LMs), as illustrated in Fig. 1 and 2. With the electric

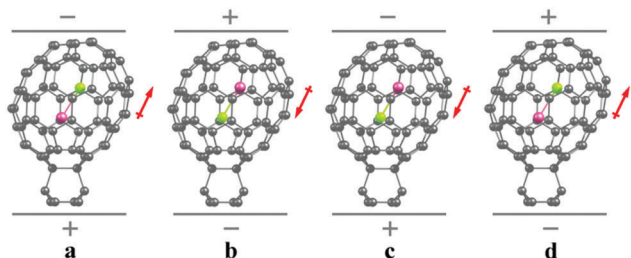


Fig. 2 Schematic representation of the orientation of MX ($\text{M} = \text{Li}$, Na purple and $\text{X} = \text{F}$, Cl green spheres) in C_{70} . The small red vectors represent the direction of the dipole moment of MX . The direction of the applied external electric field is denoted by the + and – signs.

field on, the conformations with the dipole moment of MX along the EF (**c** and **d** in Fig. 2) are stabilized more than the other two (**a** and **b** in Fig. 2). With increasing EF the stabilization should ultimately lead to the removal of the energy barrier and to the reorientation of the dipole molecule, from **a** to **c** (or **b** to **d**) in Fig. 2.

The energy dependence of local minima and their inter-connecting transition state(s) (TS) on the EF is presented in Fig. 3 (numerical values are given in Table S2, ESI†). Without the EF, the two minima differ slightly in their energy which is due to

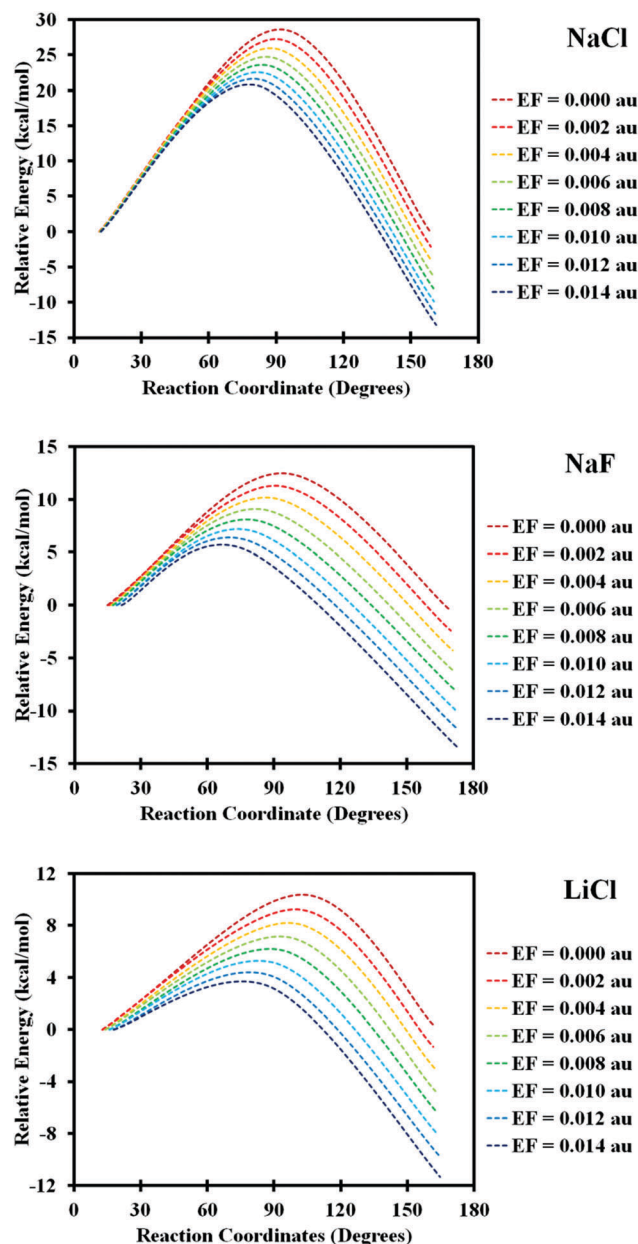


Fig. 3 The effect of EF on the energy barrier for rotation of the $\text{MX}@C_{70}$ memory system. (top) NaCl : TS_{bd} , (middle) NaF : TS_{ac} , and (bottom) LiCl : TS_{ac} . The reaction coordinate is defined by measuring the angle that MX makes with the substituted α -carbon atom (on the pole region of the cage) in the fullerene structure, see Fig. 2.

the presence of the linker. By increasing the EF the energy of the minimum with the dipole aligned parallel to the EF decreases faster as compared to the minimum with the dipole aligned antiparallel to the EF, Fig. 3 and Table S2 (ESI†).

The energy and structure of the high-energy LM become closer to that of the corresponding TS upon increasing the EF, as expected from classical Hammond's postulate.²⁴ This is the case for all studied systems, see Fig. 3 and Table S2 (ESI†). When extrapolating the results towards a larger EF the internal rotation barrier for MX should disappear, see below. As a result, only one LM should remain on the potential energy surface of the system.²⁵ Note that TS_{ac} responds to EF differently from TS_{bd}, Table S2 (ESI†). This is because in the **a** arrangement the dipole moment of the MX is mostly parallel to the dipole moment of the substituted fullerene cage; therefore, conformer **a** is more influenced by the EF and reaches the TS more easily.

The changes in the relative energies of the LMs and the TSs are regular in all cases; see Fig. 3 and Table S2 (ESI†). For NaF the barrier is about 13 kcal mol⁻¹ without an EF and goes down to 5(10) kcal mol⁻¹ at EF = 0.014 au for TS_{ac}(TS_{bd}). Increasing the size of the internal dipole by replacing NaF with NaCl does not affect the relative energies of the conformers but increases the TS energy by about 100% due to the stronger steric clash between the larger NaCl molecule and the belt region of C₇₀. For a smaller LiF molecule only a small barrier of ~2 kcal mol⁻¹ was predicted at zero field. Such a barrier can be easily overcome by an EF; for this reason, optimizations of **a** and **b** conformers ended up in **c** and **d** conformers for EF ≥ 0.004 au and no proper transition states could be located at a higher EF. LiCl “fits” between the largest (NaCl) and the smallest (LiF) systems. Its zero-field barrier is ~10 kcal mol⁻¹, which prevents the thermal motion switching but the system responds well to the applied EF, see Table S2 (ESI†) and Fig. 3.

Because the response to EF could not be computed beyond 0.014 au as the DFT approximation cannot be trusted after that,[‡] we extrapolated the data from Table S2 (ESI†) using the quadratic formula in Fig. 4. While the data for NaF and NaCl diverge, the extrapolated results for LiCl pass the barrier near EF = 0.02 au. The LiCl@C₇₀ system thus appears among studied

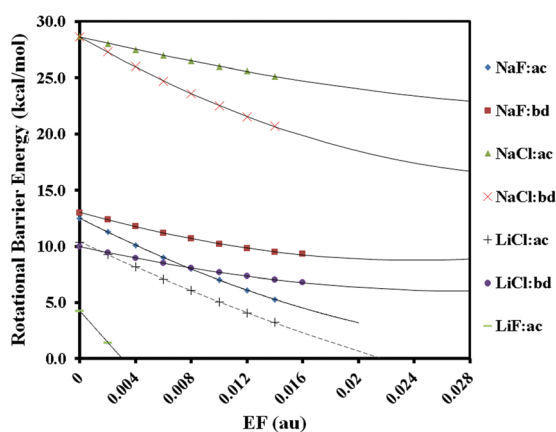


Fig. 4 Rotational barrier energies extrapolated *via* a quadratic equation.

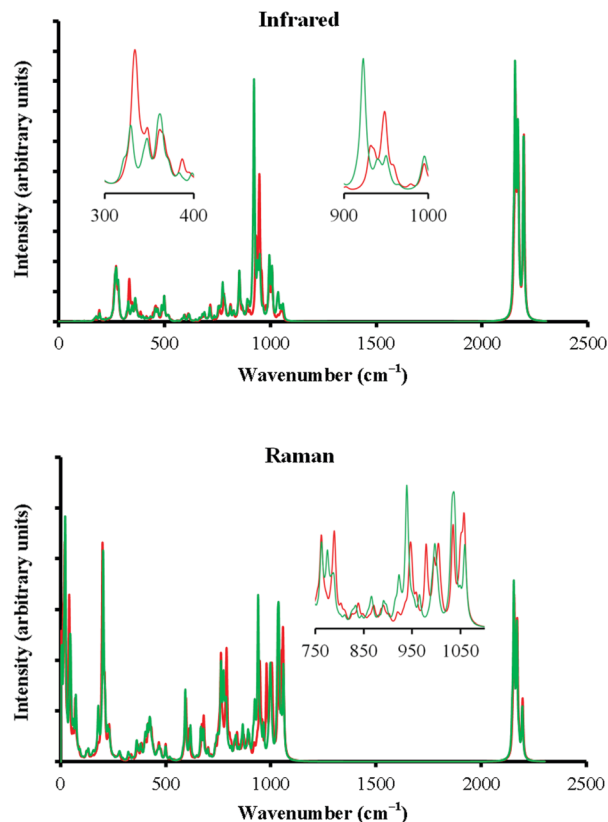


Fig. 5 Comparison of the calculated infrared spectra (top) and Raman spectra (bottom) of **a** (red) and **b** (green) conformers of LiCl@C₇₀ at EF = 0.0 au. The regions with the largest differences in the spectra are magnified for better visibility.

systems to be the best candidate for the suggested single-molecule switch.

The information can thus be encoded in the MX@C₇₀ systems; but are the two states distinct enough to allow for “the reading” of the device? They are. The two minima of MX@C₇₀ systems differ significantly in their electric multipole moments (Table S3, ESI†) and thus should be distinguishable on the basis of their molecular properties either *via* STM or spectroscopic methods. Fig. 5 illustrates the differences between IR and Raman spectra of the two minima for LiCl@C₇₀ without an EF (see Fig. S3 and S4 for NaF@C₇₀ and NaCl@C₇₀ spectra, ESI†). While most of the spectral features of both conformers are similar, the differences near 330 and 950 cm⁻¹ in IR spectra and near 800 and 1050 cm⁻¹ in Raman spectra are large enough to be observable experimentally. The detection of a single molecule is still a complicated task but novel advances in single-molecule spectroscopy offer a potential method for reading the dipole@fullerene-type memory devices in the future.²⁶

In summary, we have presented a single-molecule switch that can store information in terms of different orientations of a dipolar molecule inside an ellipsoidal fullerene using model systems of alkali metal halides enclosed in a C₇₀ cage further substituted with cyclohexane to model a carbon-based linker to a graphene sheet. MX molecules have two local energy minima inside the C₇₀ cage separated by an energy barrier due to the



shape of the fullerene. The external electric field decreases the energy barrier and enables switching between the states (conformers) of the switch. Extrapolation data suggest that the barrier removal can be achieved at a field of about 0.02 au for LiCl. The calculated IR and Raman spectra of the local minima show that they have distinct molecular properties; hence, reading of the switch is in principle possible.

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Notes and references

‡ **Methods:** The potential energy hypersurfaces for rotation of the enclosed molecule were investigated by the Gaussian 09 program.²⁷ The local minima and the corresponding transition states to pass from one minimum to another were calculated at different strengths of the electric field using the dispersion-corrected B97D3 functional²⁸ and 6-31+G(d) basis set as implemented in the Gaussian 09 code.²⁷ The fields up to a value of 0.014 au ($1.028441304 \times 10^{10}$ V m⁻¹) were examined in 0.002 au intervals. In fields stronger than the aforementioned field the HOMO–LUMO gap decreases dramatically as a result of state-mixing and DFT computations with contemporary functionals being no longer valid.²⁹ No charge transfer between the carbonic framework and the encapsulated MX was observed in the studied range of the electric field. An EF influences the energy level of a polar molecule the most when it is applied parallel to the largest component of the dipole moment vector. In our study the direction of the EF was selected to be perpendicular to the hypothetical surface that the fullerene is attached to. The transition states and local minima were confirmed by frequency analysis.

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