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# Ni on the $\text{CeO}_2(110)$ and $(100)$ surfaces: adsorption vs. substitution effects on the electronic and geometric structures and oxygen vacancies†

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We report density functional theory (DFT) calculations of the interactions of both Ni adsorbate and substitutional dopant with the ceria (110) and (100) surfaces to explain the origin of the activity of Ni/ceria catalysts. Our results indicate that the Ni adatom on the (110) surface prefers to adsorb on a two-fold bridge site over a hollow site up to 0.25 ML coverage, and the most stable position of a Ni adsorbate on the (100) surface was found to be the bridge site where the Ni atom is coordinated to two surface O atoms. The  $\text{Ni}^+$  oxidation state for the Ni adatom on the (110) surface was found to be more favorable than the  $\text{Ni}^{2+}$  state on the two-fold bridge site while on the (100) surface, a Ni adatom prefers its  $\text{Ni}^{2+}$  oxidation state over the  $\text{Ni}^+$  oxidation state. With increasing coverage, the binding energy of a Ni adatom on the (110) surface was found to decrease from  $-0.45$  eV at 0.083 ML coverage to  $-0.32$  eV at 0.25 ML coverage. Oxidation of the Ni adatom to  $\text{Ni}^+$  reduces one  $\text{Ce}^{4+}$  ion on the ceria surface to  $\text{Ce}^{3+}$  which preferred to be located next to the  $\text{Ni}^+$  ion in the nearest neighbor location. The  $\text{Ce}^{3+}$  ions on the (100) surface also prefer to stay in the vicinity of the adsorbed Ni atom, while they prefer to be located away from the Ni adatom on the (111) surface. No reduction of  $\text{Ce}^{4+}$  ions was observed upon substitution of Ce atoms by Ni atoms. Two Ni substituents preferred to be distributed on adjacent metal ion sites on the (110) surface. Ni adsorbate and substituent on the (110) surface were both found to induce significant structural distortions. In comparison to the pure ceria (110) and (100) surfaces, we show that a Ni adsorbate increases the energy required to create an oxygen vacancy while a Ni dopant reduces it. While multiple dopants on the (110) surface do reduce the vacancy formation energy, the degree of reduction is smaller compared to a single dopant indicating the presence of an optimum level of doping to obtain enhanced catalytic activity.

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## 1. Introduction

Owing to their ability to cycle oxygen through redox processes and high oxygen ion mobility, ceria based compounds have been used in a variety of applications ranging from catalysts for the treatment of automobile exhausts and industrial pollutants<sup>1–3</sup> to the water gas shift (WGS) reaction<sup>3–5</sup> and as solid electrolytes in solid oxide fuel cells.<sup>6</sup> The ability of ceria to cycle oxygen efficiently

originates from the relative ease of the reversible removal of an oxygen ion from the surface. The formation of an oxygen vacancy results in the partial reduction of lattice cerium ions from the  $\text{Ce}^{4+}$  to the  $\text{Ce}^{3+}$  oxidation state. Numerous studies, both theoretical and experimental, have been carried out in the past to characterize the role of ceria in such applications and develop novel catalysts with higher activity. Often ceria has been used in combination with other transition metals (*e.g.* Pd, Pt, Au *etc.*) or transition metal oxides (*e.g.*  $\text{VO}_x$ ) to achieve enhanced catalytic performance. For instance, a higher activity was shown by ceria-supported  $\text{VO}_x$  catalysts for the oxidative dehydrogenation of methanol,<sup>7,8</sup> ceria-supported Au<sup>9</sup> and Pt<sup>10,11</sup> catalysts for CO oxidation, ceria-supported Cu, Au, Pt, Pd, Rh<sup>4,5,12–15</sup> for the WGS reaction *etc.* More recently, it has been shown that the Ni/ $\text{CeO}_2$  system has an exceptional activity and selectivity for both WGS and ethanol steam reforming reactions.<sup>16–20</sup> A detailed review of the theoretical approaches used to study ceria and heteroatom doped ceria surface energies, catalysis by ceria and

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their comparison to experiments was presented in a recent article by Sauer *et al.*<sup>21</sup>

The geometric and electronic structure changes induced by a hetero metal atom on a metal oxide support have been described traditionally *via* 'strong metal support interaction'.<sup>22–24</sup> To explain the enhanced catalytic activity of transition metal/ceria catalysts, firstly it is necessary to gain a fundamental understanding of the interaction of the metal particle with the ceria support. Towards this end, a number of studies in the past have been carried out. Most of these studies have, however, focused on the adsorption of metal ions on the  $\text{CeO}_2(111)$  surface, which is the most stable surface<sup>25</sup> amongst the low-index ceria surfaces (*i.e.*, (100), (110) and (111) surfaces). Metal-atom adsorbates that were studied include Cu,<sup>26–29</sup> Ag,<sup>26,29</sup> Au,<sup>26,29,30</sup> Pt,<sup>31–33</sup> Pd,<sup>33,34</sup> Rh,<sup>33</sup> and Ni.<sup>19,28</sup> All these metal atoms adsorbed strongly on the (111) surface and were found to be oxidized. The oxidation of the metal adsorbates resulted in the partial reduction of one or more lattice  $\text{Ce}^{4+}$  ions to  $\text{Ce}^{3+}$ , with the unpaired electrons localizing on the Ce 4f orbitals. An Au adsorbate was found to have a number of competing states with  $\text{Au}^+$  and  $\text{Au}^0$  oxidation states of the metal lying close in energy.<sup>35</sup> While the (111) surface is the most stable of the low index ceria surfaces, ceria nanorods often expose the more reactive and less stable (110) and (100) surfaces.<sup>36,37</sup> As such, an understanding of the interaction of metal adsorbates with (110) and (100) surfaces is essential to explain the origin of their activity. Very few studies in the past have investigated the interaction of transition-metal adsorbates with the (110) and (100) surfaces. Adsorption of Cu, Ag and Au on ceria (110) was reported by Cui *et al.*<sup>38</sup> while Nolan<sup>39</sup> reported the adsorption of a series of 3rd row transition metals (V, Cr, Fe), d<sup>10</sup> metals (Cu, Ag and Au) and trivalent cations (Al, Ga, In, La). It was found that all the transition metals were oxidized (Cu, Ag and Au to their 1+ oxidation states, V and Cr to their 3+ oxidation states, Fe to its 2+ oxidation state) with the partial reduction of an equivalent number of lattice  $\text{Ce}^{4+}$  ions to  $\text{Ce}^{3+}$ . It was found that amongst the studied transition metals, Cu had the largest adsorption energy. Also, the d<sup>10</sup> metals adsorbed more strongly on the (110) surface compared to other transition metals. Amongst the trivalent cations, Al and La were oxidized to their 3+ oxidation states while Ga and In were only partially oxidized to their 2+ oxidation states. On the other hand, to the best of our knowledge, a detailed investigation of the adsorption of Ni on ceria (110) and (100) is yet to be undertaken. Experimentally it is known that Ni particles anchor more strongly on the (110) and (100) surfaces as compared to the (111) surface.<sup>40</sup> In addition, recent experiments<sup>20</sup> have shown that the oxidation state of Ni particles on ceria (111) could vary from  $\text{Ni}^{2+}$  to  $\text{Ni}^0$  depending on their surface coverage. Thus, any theoretical investigation of a transition-metal adsorbate on ceria surfaces must take into account the possibility of the presence of many competing close-lying oxidation states of the metal adsorbate. Finally, doping of ceria surfaces has also been considered as a way to enhance the oxidative capacity of the ceria surface.<sup>41</sup> Thus, a number of studies in the past have attempted to understand the modifications to the surface properties induced by transition-metal substituents. Doping of ceria surfaces with Mn,<sup>42</sup> Ti, Zr and Hf,<sup>43</sup> La,<sup>44</sup> Ni and

Pd<sup>45</sup> resulted in a strong distortion of the surface induced by the substituents along with a reduction in the energy required for the creation of an oxygen vacancy.

In this article, we have carried out a detailed analysis of the interaction of a Ni adsorbate and a substituent with the ceria (110) and (100) surfaces. Considering the points mentioned above, in this article we have attempted to answer the following questions:

(1) What are the preferred adsorption site and oxidation state for a Ni adsorbate?

(2) What is the adsorption energy of a Ni adsorbate and how does it compare with the adsorption energies of the other transition metal atoms reported in the literature?

(3) What is the nature of the distortions induced in the geometric and electronic structure of the surface by the Ni adsorbate?

(4) How do the above quantities vary as a function of the surface coverage of the Ni adsorbate?

(5) How do multiple Ni substituents prefer to distribute themselves and how do they compare with a mono-Ni substituent?

(6) How is the oxygen vacancy formation energy affected by Ni adsorption and substitution?

Section 2 describes the details of the computational methods used in this work while Section 3 discusses our results. Conclusions from this work are presented in Section 4.

## 2. Computational methods and details

DFT calculations have been performed to understand the interaction of Ni adsorbate/substituents with the ceria surface. All the calculations were performed using the VASP software.<sup>46–49</sup> The valence electronic states were expanded in a basis of plane waves. The core-valence interaction was described through the Projector Augmented Wave (PAW) approach.<sup>50,51</sup> A plane wave kinetic energy cutoff of 400 eV was used. The Perdew–Burke–Ernzerhof (PBE) GGA was used to describe the exchange–correlation interactions. The DFT+ $U$ <sup>52</sup> approach was used to describe the localized 4f and 3d electronic states in Ce and Ni respectively. In line with earlier works,<sup>53,54</sup> we choose a Hubbard- $U$  value of 4.5 eV and 4.6 eV for the Ce 4f and Ni 3d states respectively. The Hubbard- $U$  value results in a lattice constant of 5.44 Å for bulk  $\text{CeO}_2$  (fluorite structure,  $Fm\bar{3}m$  space group), 3.83 Å and 6.01 Å ( $a$  and  $c$  lattice parameter values) for bulk  $\text{Ce}_2\text{O}_3$  ( $P\bar{3}m1$  space group), 4.17 Å for bulk NiO ( $Fm\bar{3}m$  space group) in close agreement with the experimental values of 5.41 Å for  $\text{CeO}_2$ , 3.89 Å and 6.06 Å for  $\text{Ce}_2\text{O}_3$ , and 4.17 Å for NiO. Our computed lattice parameters agree well with earlier reported values for  $\text{CeO}_2$  (5.48 Å),<sup>54</sup>  $\text{Ce}_2\text{O}_3$  (3.87 and 5.93 Å)<sup>55</sup> and NiO (4.19 Å).<sup>56</sup> All our calculations utilized 'accurate' precision settings in VASP. Complete cell optimizations and a re-relaxation for the bulk systems were performed to avoid any artifacts due to Pulay stress. Spin-orbit coupling (SOC) was not considered since our calculations indicated that the inclusion of SOC has only a minor effect on the electronic structure, for instance the band gap of  $\text{CeO}_2$  changes only by 70 meV (see Fig. S1 in ES<sup>†</sup>).



We have also ignored the effect of variations in energy due to different occupations of the  $m$  projections of the 4f orbitals on Ce and 3d orbitals on Ni.<sup>57</sup> Table T1 in the ESI† gives a comparison of the results obtained with DFT/PBE, DFT/PBE+ $U$  and experiment. Clearly, our method of choice (DFT/PBE+ $U$ ) is in much better agreement with experiments as compared to DFT/PBE.

In order to study the effect of Ni adsorbate coverage on Ni/ceria interactions, 3D boundary conditions were used throughout and, hence, the surfaces were modelled using the repeated slab method in which a finite number of crystal layers are used to generate two identical surfaces *via* the introduction of a vacuum gap perpendicular to the surface. In the case of the (110) surface, two different CeO<sub>2</sub>(110) surface models, a (2 × 1) surface and a (3 × 2) surface, were considered. Both the surfaces were generated from a conventional CeO<sub>2</sub> unit cell consisting of 12 atoms (4 Ce atoms and 8 O atoms in a fluorite structure with  $Fm\bar{3}m$  space group). The (2 × 1) surface model consisted of 60 atoms (20 Ce atoms and 40 O atoms) with simulation box sizes of 10.89 Å × 7.7 Å, while the (3 × 2) surface model consisted of 180 atoms (60 Ce atoms and 120 O atoms) with box sizes of 16.33 Å × 15.4 Å. Both the slabs were 5 atomic layers thick and a vacuum gap of 15 Å was used in the direction perpendicular to the surface. Thus, the surface layer of the (2 × 1) slab contained 4 Ce atoms while that of the (3 × 2) slab contained 12 Ce atoms. Adsorption of one Ni atom then corresponds to a coverage of 0.25 ML on the (2 × 1) surface and 0.083 ML on the (3 × 2) surface. Interactions of a Ni substituent with the CeO<sub>2</sub>(110) surface were studied for the larger (3 × 2) surface only. Mono-substitution and di-substitution of surface Ce atoms with Ni atoms were considered to investigate the effect of the Ni substituent concentration. Thus, mono-substitution corresponds to 8.3% Ni concentration, while di-substitution corresponds to 16.6% Ni concentration at the surface. The (100) surface was modelled as an oxygen-terminated surface with half of the oxygen atoms in a chessboard-like arrangement. It is stoichiometric and has been found to be the most stable polar-compensated (100) plane in the past.<sup>58</sup> A (3 × 3) expansion of the (100) surface was used for our calculations, and includes 11 atomic layers and 135 atoms (45 CeO<sub>2</sub> units) with simulation box sizes of 11.54 Å × 11.54 Å. Adsorption of one Ni atom then corresponds to a coverage of 0.111 ML. Mono-substitution corresponds to 11.1% Ni concentration at the surface. These models were further used to compute the formation energies of charge-compensating and active O vacancies.

For the (100) surface, the atomic coordinates of all the atoms were allowed to relax fully, but for the (110) surface, atoms in the bottom two layers were frozen at their bulk positions due to the large size of the system (180 atoms). For both the slabs, the lattice parameters were fixed at the above-mentioned values during optimization. Geometry optimizations were deemed to have converged when the forces dropped below 0.03 eV Å<sup>-1</sup>. For the electronic structure, the SCF energy convergence was set to a threshold of 10<sup>-5</sup> eV. A Pulay mixing scheme<sup>59</sup> as implemented in VASP was used for charge-density mixing in the SCF solution. For the (3 × 3) slab of the (100) surface, during geometry optimization, the Brillouin zone was sampled using a 2 × 2 × 1

Monkhorst-Pack (MP) grid. For the (2 × 1) slab of the (110) surface, the Brillouin zone was sampled using a (1 × 2 × 1) MP  $k$ -point mesh. Later, a finer (4 × 6 × 1) MP  $k$ -point mesh was used to compute the energies at the optimized structure. For the (3 × 2) slab of the (110) surface, the Brillouin zone was sampled using the  $\Gamma$ -point approximation.

CeO<sub>2(s)</sub>, Ni<sub>(s)</sub>, Ce<sub>(s)</sub> and O<sub>2(g)</sub> were used as the reference to compute the formation energies in eqn (1)–(8). Calculations on the face-centered cubic (fcc) Ni<sub>(s)</sub> and  $\gamma$ -Ce<sub>(s)</sub> were performed on a one-atom unit cell using a 16 × 16 × 16 Monkhorst-Pack  $k$ -point mesh, resulting in an optimized lattice parameter of 3.474 Å for Ni<sub>(s)</sub> and 5.208 Å for Ce<sub>(s)</sub>, in good agreement with the experimental value of 3.52 Å and 5.16 Å respectively. The local magnetic moment on the Ni atom in Ni<sub>(s)</sub> was found to be  $\sim 0.78 \mu_B$  while it was  $\sim 1 \mu_B$  on the Ce atom in Ce<sub>(s)</sub>. Calculations on the O<sub>2(g)</sub> molecule were performed in a 10 Å × 10 Å × 10 Å box in the triplet state using the  $\Gamma$ -point approximation to sample the Brillouin zone. The optimized O<sub>2</sub> bond length was found to be 1.234 Å.

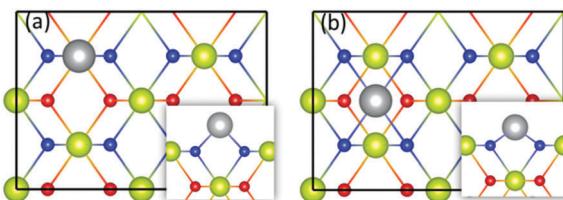
### 3. Results and discussion

In bulk CeO<sub>2</sub> (fluorite structure,  $Fm\bar{3}m$  space group), Ce<sup>4+</sup> and O<sup>2-</sup> are characterized by eight- and four-fold coordination, respectively. The relaxed O–Ce bond length of 2.355 Å is close to the experimental value of 2.343 Å. The CeO<sub>2</sub>(110) surface consists of six-coordinated Ce<sup>4+</sup> ions and three-coordinated O<sup>2-</sup> ions. On the (100) surface, Ce<sup>4+</sup> and O<sup>2-</sup> are six- and two-coordinated, respectively. The relaxed surface was found to have a Ce–O bond length of 2.326 Å on the (110) surface and 2.191 Å on the (100) surface. We have also investigated the CeO<sub>2</sub>(111) surface, in which Ce<sup>4+</sup> and O<sup>2-</sup> are seven- and three-coordinated, respectively. The O–Ce bond length 2.360 Å is close to the bulk O–Ce bond length 2.355 Å. The surfaces were diamagnetic with all the Ce ions in their Ce<sup>4+</sup> oxidation state and O ions in their O<sup>2-</sup> oxidation state. The surface energies of CeO<sub>2</sub>(110), (100) and (111) were found to be 1.06, 1.44 and 0.60 J m<sup>-2</sup>, respectively. They are in very good agreement with earlier reported values (1.06 J m<sup>-2</sup> for (110), 1.44 J m<sup>-2</sup> for (100) and 0.68 J m<sup>-2</sup> for the (111) surface) in the literature.<sup>25,60–62</sup> It is known that the order of stability of CeO<sub>2</sub> low index surfaces is CeO<sub>2</sub>(111) > CeO<sub>2</sub>(110) > CeO<sub>2</sub>(100).<sup>25</sup>

#### 3.1 Ni adsorption on CeO<sub>2</sub> surfaces

**A. Ni adsorption on CeO<sub>2</sub>(110): preferred adsorption site.** A few previous articles<sup>38,39,63</sup> have reported the adsorption of transition metal atoms on the CeO<sub>2</sub>(110) surface. It was shown that the preferred adsorption site was a bridge site, connecting two surface O atoms. We have studied the adsorption of a Ni atom on a bridge site and a hollow site (where the adsorbate is coordinated to four surface O atoms), as shown in Fig. 1. The (2 × 1) CeO<sub>2</sub>(110) slab was used for these calculations, corresponding to a Ni coverage of 0.25 ML. Recent experiments<sup>20</sup> on the Ni/CeO<sub>2</sub>(111) surface indicated that the oxidation state of the adsorbed Ni atom could vary from Ni<sup>2+</sup> to Ni<sup>0</sup> depending on





**Fig. 1** Initial structures of Ni adsorption on the ceria (110) surface. (a) Bridge site (b) Hollow site. Insets show the front view of the structures. Ce atoms are yellow, O atoms are red and blue, Ni atom is grey. In order to distinguish the surface O atoms from other O atoms in the slab, they are shown in blue colour.

the coverage of Ni. At lower coverages, the Ni atom was found to be oxidized to  $\text{Ni}^{2+}$ . As the coverage approached unity, the Ni atom seemed to remain in the  $\text{Ni}^0$  oxidation state.

Thus, in addition to the adsorption site, any detailed analysis of the interaction of a transition-metal adsorbate on  $\text{CeO}_2$  surfaces must consider the oxidation state of the adsorbed transition metal atom. As such, we have studied configurations where the Ni atom is in its  $\text{Ni}^+$  and  $\text{Ni}^{2+}$  oxidation states on both the bridge and hollow sites. In its  $\text{Ni}^+$  oxidation state, the Ni atom donates an electron from its 4s orbital to the 4f orbital of a surface Ce, thereby reducing it from  $\text{Ce}^{4+}$  to  $\text{Ce}^{3+}$ . Similarly, when Ni is in its  $\text{Ni}^{2+}$  oxidation state, two surface Ce are reduced from  $\text{Ce}^{4+}$  to  $\text{Ce}^{3+}$ . It is now well known that the reduced  $\text{Ce}^{3+}$  could be located in one or more of the many surface metal ion sites and various configurations must be investigated to determine the most favorable distribution of  $\text{Ce}^{3+}$  on the surface. In order to do so, we have followed the procedure outlined by Kullgren *et al.*<sup>64</sup> First we perform a geometry optimization using a different PAW potential (where the Ce 4f electron is included in the core) on those Ce atoms that we wish to be reduced to  $\text{Ce}^{3+}$ , while the rest of the Ce atoms use the full valence pseudopotential. In the second step we continue the optimization from the converged structure in the previous step with all the Ce atoms described using the full valence pseudopotential. At the end of each calculation, the magnetizations on the Ce ions were checked to ensure that the reduced  $\text{Ce}^{3+}$  ion was located at the same position as chosen initially. This was done by looking at the local magnetization on the  $\text{Ce}^{3+}$  ion which should be *ca.*  $1 \mu_B$  due to an electron localized in its f orbital while the remaining  $\text{Ce}^{4+}$  ions have a magnetic moment of  $0 \mu_B$ . The reaction energy for the adsorption of a Ni atom on  $\text{CeO}_2(110)$  was computed using eqn (1):

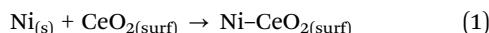
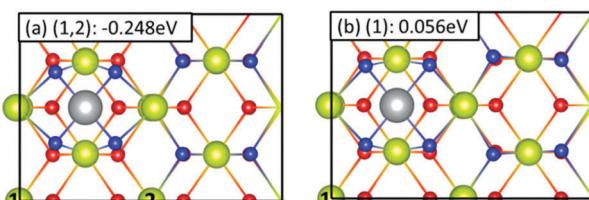
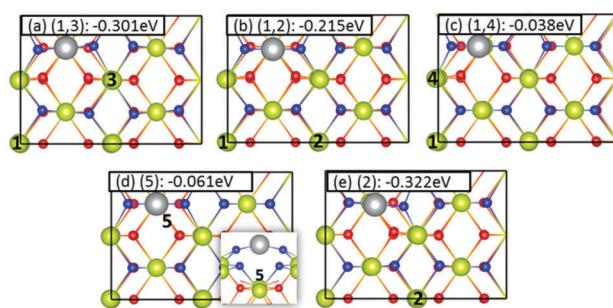


Fig. 2 and 3 show the optimized structures and reaction energies for the adsorption of Ni on hollow and bridge sites respectively. From Fig. 2 we see that on a hollow site, the Ni ion prefers a  $\text{Ni}^{2+}$  oxidation state over a  $\text{Ni}^+$  oxidation state, with the two reduced  $\text{Ce}^{3+}$  ions located at sites (1,2). The reaction energy for the adsorption of a  $\text{Ni}^{2+}$  ion on a hollow site at 0.25 ML coverage was found to be  $-0.248$  eV while  $\text{Ni}^+$  adsorption was disfavored. The local structure around the  $\text{Ni}^{2+}$  ion was nearly planar with the four Ni-O bond lengths equal to  $1.936 \text{ \AA}$ .



**Fig. 2** Adsorption of Ni on a hollow site on the  $\text{CeO}_2(110)$  surface. (a) Ni is in its  $\text{Ni}^{2+}$  oxidation state with the two  $\text{Ce}^{3+}$  ions located at sites 1 and 2. (b) Ni is in  $\text{Ni}^+$  oxidation state with the reduced  $\text{Ce}^{3+}$  ion located at site 1. Ce atoms are yellow, O atoms are red and blue, and Ni atom is grey.

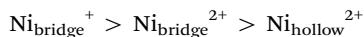


**Fig. 3** Adsorption of Ni on a bridge site on the  $\text{CeO}_2(110)$  surface. (a–c) Ni is in its  $\text{Ni}^{2+}$  oxidation state with the two  $\text{Ce}^{3+}$  ions located at sites (1,3), (1,2) and (1,4) respectively. (d–e) Ni is in its  $\text{Ni}^+$  oxidation state with the reduced  $\text{Ce}^{3+}$  ion located at site 5 and 2 respectively. Inset in (d) shows the front view of the structure. Here, site 5 is the Ce ion beneath the Ni adatom. Ce atoms are yellow, O atoms are red and blue, Ni atom is grey.

Comparing this with the Ni-O bond length of  $2.085 \text{ \AA}$  in solid  $\text{NiO}$ , we see that the Ni-O bond lengths are shortened upon adsorption of a  $\text{Ni}^{2+}$  ion on a hollow site. The four  $\text{O}^{2-}$  ions that coordinate with the adsorbed  $\text{Ni}^{2+}$  are pulled away from their normal positions on the surface, with the  $\text{Ce}^{3+}-\text{O}^{2-}$  distance increased to  $2.835 \text{ \AA}$  from  $2.326 \text{ \AA}$ . On a bridge site, however (see Fig. 3), we see that both  $\text{Ni}^+$  and  $\text{Ni}^{2+}$  oxidation states have stable configurations, with  $\text{Ni}^+$  being the preferred oxidation state. As mentioned earlier, for a given oxidation state of the adsorbate, various distributions of the reduced  $\text{Ce}^{3+}$  ions are possible. Here we have investigated three different  $\text{Ce}^{3+}$  distributions for  $\text{Ni}^{2+}$  adsorbate and two different distributions for  $\text{Ni}^+$  adsorbate, as shown in Fig. 3. For the  $\text{Ni}^{2+}$  adsorbate, two  $\text{Ce}^{3+}$  ions located diagonally across the  $\text{Ni}^{2+}$  ion at the (1,3) site as shown in Fig. 3(a) on the surface was found to be the most favorable distribution, with a reaction energy of  $-0.301$  eV. Upon adsorption, the two O atoms bridged by the  $\text{Ni}^{2+}$  ion are pulled outward from their location on the surface. This results in an elongation of the two  $\text{O}^{2-}-\text{Ce}^{4+}$  bonds to  $2.464 \text{ \AA}$  from  $2.326 \text{ \AA}$ . The two  $\text{O}^{2-}-\text{Ce}^{3+}$  distances were  $2.584 \text{ \AA}$  while the  $\text{Ni}^{2+}-\text{O}^{2-}$  bond distance was  $1.69 \text{ \AA}$ . In comparison to  $\text{Ni}^{2+}$  adsorption on the hollow site, we notice that the distortions in the Ce-O bond lengths are smaller, resulting in a larger binding energy of the  $\text{Ni}^{2+}$  adsorbate on the bridge site. The least favorable distribution of  $\text{Ce}^{3+}$  ions for  $\text{Ni}^{2+}$  adsorbate was found to be at nearest neighbor locations (site (1,4) as shown in Fig. 3c), which was *ca.*  $263$  meV less stable compared to  $\text{Ce}^{3+}$  at site (1,3).



located at the (1,3) site. For the  $\text{Ni}^+$  adsorbate on a bridge site, we have investigated two different locations of the  $\text{Ce}^{3+}$  ion, a nearest-neighbor surface site (site 2, Fig. 3e) and a subsurface site (site 5, Fig. 3d). We find that the subsurface site is unfavorable while the nearest neighbor location for the  $\text{Ce}^{3+}$  ion results in adsorption energy of  $-0.322$  eV. Adsorption of  $\text{Ni}^+$  on a bridge site causes the  $\text{O}^{2-}$  coordinated to the  $\text{Ce}^{3+}$  ion to be pulled away from the surface significantly (by  $1.23$  Å with respect to  $\text{CeO}_2(110)$ ), resulting in a locally asymmetric structure around the  $\text{Ni}^+$  adsorbate. Fig. S2 in the ESI<sup>†</sup> shows the bond lengths surrounding the  $\text{Ni}^+$  adsorbate. Thus, we find that a Ni adsorbate prefers the bridge site over the hollow site. Further, on the bridge site, the Ni adsorbate prefers the  $\text{Ni}^+$  oxidation state, with the reduced  $\text{Ce}^{3+}$  located on the nearest neighbour surface site, over the  $\text{Ni}^{2+}$  oxidation state. The stabilization energy of the  $\text{Ni}^+$  oxidation state over the  $\text{Ni}^{2+}$  oxidation state is *ca.*  $21$  meV. In addition, the  $\text{Ni}^+$  adsorbate on the bridge site is *ca.*  $74$  meV more stable than the  $\text{Ni}^{2+}$  adsorbate on a hollow site. Thus, the preferred adsorption pattern for a Ni adsorbate on the  $\text{CeO}_2(110)$  surface is:



### B. Ni adsorption on $\text{CeO}_2(110)$ : effect of surface coverage.

Results from the previous section indicate that a Ni adsorbate prefers to adsorb on a bridge site in the  $\text{Ni}^+$  oxidation state on the  $\text{CeO}_2(110)$  surface. As mentioned earlier, experiments on Ni clusters adsorbed on the  $\text{CeO}_2(111)$  surface<sup>20</sup> have indicated that the preferred oxidation state of Ni on  $\text{CeO}_2(111)$  varies with the coverage of Ni. In order to study the effect of coverage on the oxidation state of the adsorbed Ni atom, we have investigated the adsorption of a Ni atom on a larger ( $3 \times 2$ )  $\text{CeO}_2(110)$  surface. We have studied adsorption on the bridge site only since results from the previous section indicated that the bridge site is preferred. Again, we have investigated a number of distributions of the  $\text{Ce}^{3+}$  ions on the surface, corresponding to a given oxidation state of the Ni adsorbate (*i.e.*,  $1$   $\text{Ce}^{3+}$  ion for the  $\text{Ni}^+$  oxidation state and two  $\text{Ce}^{3+}$  ions for the  $\text{Ni}^{2+}$  oxidation state). Fig. 4 shows the optimized structures along with the adsorption energies of the Ni adsorbate computed using eqn (1).

From Fig. 4, it is clear that on a larger ( $3 \times 2$ ) surface (corresponding to  $0.083$  ML Ni coverage) as well, a Ni adsorbate prefers the  $\text{Ni}^+$  oxidation state over the  $\text{Ni}^{2+}$  oxidation state. This is in contrast to Ni adsorption on the  $\text{CeO}_2(111)$ <sup>19</sup> and  $\text{CeO}_2(100)$  surfaces, where the Ni adsorbate oxidizes to its  $\text{Ni}^{2+}$  state. The most favourable location of the reduced  $\text{Ce}^{3+}$  ion was found to be next to the adsorbed  $\text{Ni}^+$  (see Fig. 4a) while locations far away from the adsorbed  $\text{Ni}^+$  (see Fig. 4d) or the subsurface location were found to be unfavourable. For the Ni adsorbate in the  $\text{Ni}^{2+}$  oxidation state, similar to the structure found for the ( $2 \times 1$ ) slab above, two  $\text{Ce}^{3+}$  ions located diagonally across the  $\text{Ni}^{2+}$  adsorbate was found to be the most stable configuration. However, amongst the configurations studied here, while the energy difference between the least stable and the most stable  $\text{Ce}^{3+}$  distributions for  $\text{Ni}^+$  adsorbate was *ca.*  $795$  meV, we find this value to be reduced to only  $97$  meV for the  $\text{Ni}^{2+}$

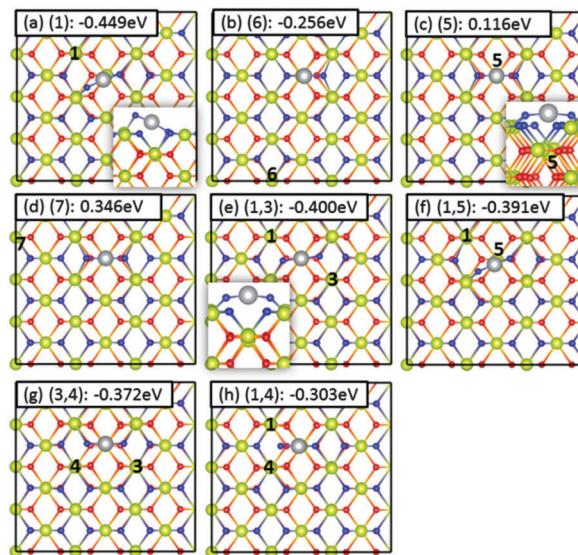


Fig. 4 Adsorption of Ni on a bridge site on the  $\text{CeO}_2(110)$  ( $3 \times 2$ ) surface. (a-d) Ni is in its  $\text{Ni}^+$  oxidation state with the  $\text{Ce}^{3+}$  ion located at sites 1, 6, 5 and 7 respectively. (e-h) Ni is in its  $\text{Ni}^{2+}$  oxidation state with the two reduced  $\text{Ce}^{3+}$  ions located at sites (1,3), (1,5), (3,4) and (1,4) respectively. The insets in panels (a) and (e) show the front view of the structure. Ce ion is yellow, O atoms are red and blue, Ni atom is gray.

adsorbate. Furthermore, comparing the above results with those on the smaller ( $2 \times 1$ ) slab (corresponding to  $0.25$  ML Ni coverage) presented in the previous section, the following trends emerge:

(1) As the coverage of Ni on the  $\text{CeO}_2(110)$  surface increases from  $0.083$  ML to  $0.25$  ML, the adsorption energy for the most favourable configuration, *i.e.*,  $\text{Ni}^+$  on a bridge site, reduces from  $-0.449$  eV to  $-0.322$  eV.

(2) With increase in coverage of Ni from  $0.083$  ML to  $0.25$  ML, the stabilization of the  $\text{Ni}^+$  oxidation state over the  $\text{Ni}^{2+}$  oxidation state on a bridge site decreases from *ca.*  $49$  meV to  $21$  meV for the most preferred  $\text{Ce}^{3+}$  ion distribution in the respective oxidation states of the Ni adsorbate.

(3) Amongst the configurations investigated here, with increase in coverage from  $0.083$  ML to  $0.25$  ML, the energy difference between the least and most stable  $\text{Ce}^{3+}$  distribution for the  $\text{Ni}^+$  adsorbate decreases from *ca.*  $795$  meV to  $383$  meV, while this value increases from *ca.*  $97$  meV to  $263$  meV for the  $\text{Ni}^{2+}$  adsorbate.

In comparison to the adsorption energies of other transition-metal atoms on the  $\text{CeO}_2(110)$  surface,<sup>39</sup> namely Au ( $-2.19$  eV), Ag ( $-1.45$  eV), Cu ( $-3.25$  eV), V ( $-0.97$  eV), Cr ( $-0.64$  eV) and Fe ( $-0.75$  eV), we find that Ni has the lowest adsorption energy amongst the transition metals reported so far. On the other hand, comparing the adsorption energy of a Ni atom on the  $\text{CeO}_2(100)$  ( $-1.490$  eV) and  $\text{CeO}_2(111)$  ( $-0.185$  eV) surfaces (see Table T2 in ESI<sup>†</sup>), we find that the order of stability of a Ni adsorbate on ceria surfaces is  $\text{CeO}_2(100) > \text{CeO}_2(110) > \text{CeO}_2(111)$ . This is in agreement with the experimental finding that the  $\text{CeO}_2(110)$  and  $\text{CeO}_2(100)$  surfaces bind Ni stronger than the  $\text{CeO}_2(111)$  surface.<sup>40</sup>



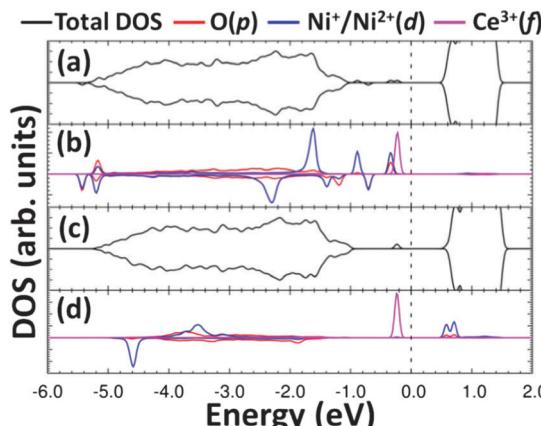


Fig. 5 Total and projected density of states for the adsorption of Ni on  $\text{CeO}_2(110)$  bridge site at 0.083 ML coverage. (a and b) Total and projected density of states for Ni adsorbate in  $\text{Ni}^+$  oxidation state. (c and d) Total and projected density of states for Ni adsorbate in  $\text{Ni}^{2+}$  oxidation state. The Fermi level is set to 0 in all the plots.

Fig. 5 shows the total and projected density of states for a Ni adsorbate on a bridge site at 0.083 ML coverage in both  $\text{Ni}^+$  (Fig. 5a and b) and  $\text{Ni}^{2+}$  (Fig. 5c and d) oxidation states. The Fermi energy is set to 0 eV in all panels of Fig. 5. The common feature in both cases is the appearance of split-off states above the valence band of  $\text{CeO}_2(110)$  and right below the Fermi level. While for the  $\text{Ni}^{2+}$  adsorbate, this state arises almost exclusively from the two electrons localized on the 4f orbitals of the two reduced  $\text{Ce}^{3+}$  ions, for the  $\text{Ni}^+$  adsorbate this state arises from the combination of the 3d orbital of  $\text{Ni}^+$ , 4f orbital on the reduced  $\text{Ce}^{3+}$  ion and the 2p orbitals on the two  $\text{O}^{2-}$  ions coordinating the  $\text{Ni}^+$  adsorbate.

**C. Ni adsorption on  $\text{CeO}_2(100)$ : preferred adsorption site.** Two adsorption sites (a 2-fold site on top of a surface Ce atom and a bridge site on top of two surface O atoms) have been explored for a single Ni adsorbed on the  $\text{CeO}_2(100)$  surface. The most stable position was found to be the bridge site where the Ni atom is coordinated to two surface O atoms (see Fig. 6). In contrast to the (110) surface,  $\text{Ni}^{2+}$  state was found to be more stable than the  $\text{Ni}^+$  state on the (100) surface, resulting in the partial reduction of two surface Ce ions from  $\text{Ce}^{4+}$  to  $\text{Ce}^{3+}$ . Imposing this constraint, we have systematically explored a range of possible configurations for the  $\text{Ce}^{3+}$  pair (see Fig. S9 in ESI†). The most stable distribution of two  $\text{Ce}^{3+}$  ions on the (100) surface is shown in Fig. 6. For this structure, the oxygen atoms surrounding  $\text{Ni}^{2+}$  relax outward, resulting in a  $\text{O}-\text{Ce}^{3+}$  bond length of 2.592 Å. The calculated  $\text{O}-\text{Ce}^{3+}$  bond lengths range from 2.324 to 2.666 Å, which is longer than the  $\text{O}-\text{Ce}^{4+}$  bond length of 2.191 Å on the pure (100) surface. The nearest-neighbour  $\text{Ce}^{4+}$  atoms around  $\text{Ni}^{2+}$  have  $\text{O}-\text{Ce}^{4+}$  bond lengths varying from 2.080 to 2.425 Å, leading to a large distortion of the surface. The two  $\text{O}-\text{Ni}$  bond lengths are 1.727 Å, which is much shorter than the  $\text{O}-\text{Ni}$  bond length (2.085 Å) in  $\text{NiO}$ . Therefore, a strong metal-support interaction is preferred. Fig. 6(b) shows the total and projected density of states for this structure. The main features are similar to the DOS of  $\text{Ni}^{2+}$  on

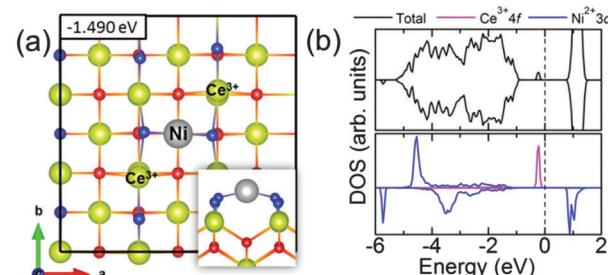


Fig. 6 (a) The optimized atomic structure of Ni adsorbed on the  $\text{CeO}_2(100)$  surface with the two  $\text{Ce}^{3+}$  ions located in the nearest neighbour locations. Inset shows the side view of the outermost four atomic layers. Ce atoms are yellow, O atoms are red and blue, Ni atoms is grey. (b) The total density of states (DOS) in the upper panel and projected DOS of  $\text{Ce}^{3+}$  4f and  $\text{Ni}^{2+}$  3d orbitals in the lower panel. The curves above (below) the horizontal axis correspond to spin-up (down) states. The Fermi level (the vertical dashed line) is set to zero.

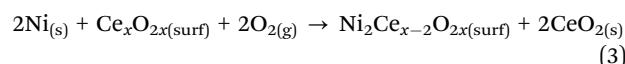
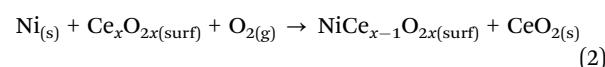
the (110) surface. The two Ni 4s electrons are transferred to the empty Ce 4f orbitals that appear as split-off states from the Ce 4f band right below the Fermi level. The calculated local magnetic moment of  $\text{Ce}^{3+}$  is  $0.97 \mu_B$ , which is equal to the value  $0.96 \mu_B$  of  $\text{Ce}^{3+}$  in  $\text{Ce}_2\text{O}_3$ .

The  $\text{Ce}^{3+}$  ion prefers to be distributed around the  $\text{Ni}^{2+}$  adsorbate on the ceria (110) and (100) surfaces. This trend is in contrast to  $\text{Ni}^{2+}$  adsorbed on the ceria (111) surface,<sup>19</sup> where the  $\text{Ce}^{3+}$  ions strongly prefer to maximize their separation from the adsorbed  $\text{Ni}^{2+}$  ion (see Fig. S14 and S15 in ESI†).

### 3.2 Ni substituent on $\text{CeO}_2$ surfaces

#### A. Ni substituent on $\text{CeO}_2(110)$ : mono- vs. disubstitution and the formation of charge compensating oxygen vacancies.

The previous section investigated the effects of the adsorption of a Ni atom on the  $\text{CeO}_2(110)$  and (100) surfaces. In this section, results for a Ni-substituted  $\text{CeO}_2(110)$  surface are discussed. We have considered the substitution of one surface Ce atom by a Ni atom and also the substitution of two surface Ce atoms by two Ni atoms in order to analyze how multiple substituents prefer to distribute themselves on the surface. The formation energy of mono and di Ni substituted surfaces were calculated using eqn (2) and (3) respectively.



A number of different configurations for Ni substituents on the surface were examined to determine the most favourable location. Fig. 7 shows the most favourable configuration of a mono- and di-Ni substituted surface while Fig. S3 in the ESI† shows the geometries of various other metastable structures along with their energies relative to the most favourable configuration shown in Fig. 7.

Doping the surface with a Ni atom results in a strong distortion of the surface with the Ni substituent moving from



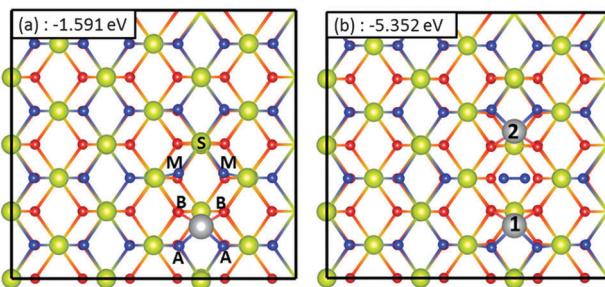
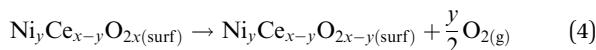


Fig. 7 Ni substituent on  $\text{CeO}_2(110)$  surface. (a) Optimum configuration for a mono substituent with four-coordinated Ni. (b) Optimum configuration for di Ni substituted surface. Ce atoms are yellow, O atoms are red and blue, Ni substituent is gray. 'S', 'M', 'A' and 'B' in (a) correspond to a surface  $\text{Ce}^{4+}$  ion, magnetic oxygen atoms, surface and sub-surface oxygen atoms respectively. 1, 2 in (b) correspond to the two Ni dopants.

the surface metal ion site into the ceria lattice by about  $0.69\text{ \AA}$ , consistent with the observations of Nolan.<sup>45</sup> The local geometry around the Ni substituent was found to be nearly planar with the Ni substituent coordinating with two surface- (designated by 'A' in Fig. 7a) and two subsurface O ions (designated by 'B' in Fig. 7a). The  $\text{Ni}-\text{O}_A$  distance was found to be  $1.806\text{ \AA}$  while the  $\text{Ni}-\text{O}_B$  distance was found to be  $1.870\text{ \AA}$ . For di-Ni substitution, the preferred distribution of the two Ni substituents was found to be at the nearest neighbour positions as shown in Fig. 7b. This configuration resulted in the spontaneous formation of an  $\text{O}_2$  molecule (see Fig. 7b) that remained adsorbed at around  $1.9\text{ \AA}$  above the surface. The  $\text{O}_2$  molecule was found to be in a triplet spin state with a bond length of  $1.238\text{ \AA}$ , very close to the VASP-optimized bond length of  $1.234\text{ \AA}$ . Further we note that once a dopant is introduced, the energy cost associated with the introduction of a second dopant is reduced, compared to the first dopant ( $-1.591\text{ eV}$  vs.  $-5.352\text{ eV}$ ). In  $\text{CeO}_2$ , Ce ions are in an oxidation state of  $\text{Ce}^{4+}$  while O ions are in an oxidation state of  $\text{O}^{2-}$ . Substituting a  $\text{Ce}^{4+}$  ion in  $\text{CeO}_2(110)$  with a divalent cation results in a charge imbalance which is compensated by spin polarization of the oxygen ion (in the case of Ni substituent, mainly on those O ions marked with the label M in Fig. 7a), leading to partially unoccupied 'p' orbitals on these ions. The magnetic moment on the two  $\text{O}_M$  ions was found to be  $0.37\text{ }\mu_B$  while the magnetic moment on the Ni substituent was  $0.92\text{ }\mu_B$ , indicating that the Ni substituent does not get fully oxidized to the  $\text{Ni}^{2+}$  state. The total and projected density of states for the mono-Ni substituted surface is shown in Fig. S4 of the ESI.<sup>†</sup>

Another means to achieve charge neutrality is through the loss of one or more surface O atoms to form 'charge compensating' oxygen vacancies (one and two vacancies for mono- and di-Ni substituted  $\text{CeO}_2(110)$  respectively). We have studied various distributions of charge compensating O vacancies on both the surfaces and computed their formation energy according to eqn (4)



In the eqn (4),  $y = 1$  for the mono-Ni substituted surface and  $y = 2$  for the di-Ni substituted surface. The most favourable

location for the formation of a compensating O vacancy in mono-Ni substituted  $\text{CeO}_2(110)$  was reported<sup>45</sup> to be from one of the  $\text{O}_M$  sites in Fig. 7a. We computed the O vacancy formation energy using eqn (4) to be  $-0.681\text{ eV}$  from the  $\text{O}_M$  site while Nolan<sup>45</sup> reported a value of only  $-0.24\text{ eV}$ . For the di-Ni substituted surface, a number of different locations for the compensating O vacancies were studied for different distributions of Ni dopants on the surface. For the most favourable (1, 2) distribution of the Ni dopants (see Fig. 7b), the compensating O vacancy formation energy was found to be  $-0.533\text{ eV}$ . However, it must be noted that this energy really corresponds to the molecular adsorption energy of  $\text{O}_2$  on the surface rather than O-vacancy formation energy as the (1, 2) substitution of dopants led to a spontaneous ejection of an  $\text{O}_2$  molecule as shown in Fig. 7b. In all the cases we notice that the Ni substituents continue to remain in a four coordinated environment even upon the formation of charge compensating vacancies.

**B. Ni substituent on  $\text{CeO}_2(100)$ : monosubstitution and the formation of charge compensating oxygen vacancies.** We have also considered the structure of Ni substituted at a Ce site on the (100) surface, as shown in Fig. 8. On the (100) surface, the Ni substituent is surrounded by six anions as the surface O is shared by two metal ions while an O in the bulk is shared by four metal ions. It is to be noted that Ni in  $\text{NiO}$  is coordinated to six O atoms, too. Upon relaxation, while Ni retains its six-fold coordination, the  $\text{Ni}-\text{O}$  bond lengths of the six  $\text{Ni}-\text{O}$  bonds are not identical, with the bond lengths ranging from  $1.979\text{ \AA}$  to  $2.410\text{ \AA}$ . Four of the  $\text{Ni}-\text{O}$  bonds (two of type M1 and two of type M2 as shown in the lower-right panel of Fig. 8a) are shortened to  $1.979\text{ \AA}$  while the remaining two  $\text{Ni}-\text{O}$  bonds with the O atoms in the third atomic layer are elongated to  $\sim 2.410\text{ \AA}$ . However, the distortion of the Ni-substituted (100) surface is smaller than that of the Ni-substituted (110) surface. A strong distortion of the (111) surface by a Ni substituent was also indicated by Nolan<sup>45</sup> in which the fluorite structure of  $\text{CeO}_2$

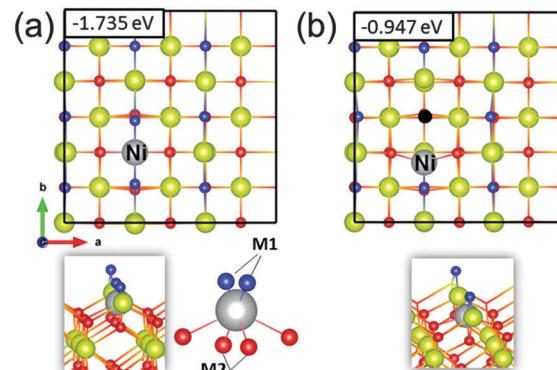
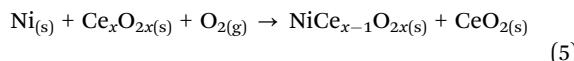


Fig. 8 (a) The optimized atomic structures of the Ni-substituted  $\text{CeO}_2(100)$  surface. The lower-left inset shows the side view of the outermost four atomic layers and the lower-right inset shows the local structure of the  $\text{Ni}^{2+}$  ion. M1 and M2 represent the magnetic O ions surrounding the Ni substituent. Ce atoms are yellow, O atoms are red and blue and Ni atom is grey. (b) The Ni-substituted  $\text{CeO}_2(100)$  surface with a charge compensating oxygen vacancy represented by a black sphere.



does not facilitate a six- or seven- coordination environment for Ni. As a result, the large distortion around the Ni substituent leads to a nearly square planar four-coordinate environment for the Ni substituent (see Fig. S16 in ESI†).

As indicated, on the Ni-substituted (110) surface, substituting a  $\text{Ce}^{4+}$  ion in  $\text{CeO}_2(110)$  with a lower-valent cation results in a charge imbalance which is compensated by spin polarization of the oxygen ion. The Ni substituent on the (100) surface also leads to spin-polarized O atoms, resulting in unoccupied gap states  $\sim 0.5$  eV above the top of the valence band (see Fig. S18 in ESI†). These states have mostly O 2p character and are spatially localized on the four nearest-neighbor O atoms of the Ni substituent (see the atoms marked by M1 and M2 in Fig. 8a). The two O atoms marked by M1 and M2 in Fig. 8a have a local magnetic moment of  $\sim 0.61$  and  $0.23 \mu_B$  per O atom, respectively. Here, the magnetic moment on the Ni substituent was  $1.78 \mu_B$ , indicating that the Ni substituent gets fully oxidized to the  $\text{Ni}^{2+}$  state. For reference, the calculated magnetic moment of  $\text{Ni}^{2+}$  in  $\text{NiO}$  was  $1.65 \mu_B$ . The formation energy of the Ni-substituted (100) surface was found to be  $-1.735$  eV, while this value was found to be  $-0.572$  eV for the (111) surface (see Table T2 in ESI†), indicating that it is easier to dope the (100) surface than the (110) and (111) surfaces, in agreement with their order of stability. The formation energy of a Ni substituent in bulk  $\text{CeO}_2$  with a 96-atom supercell (see Fig. S20 in ESI†) was calculated using eqn (5).

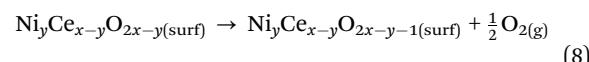
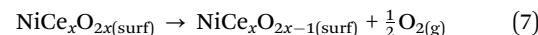
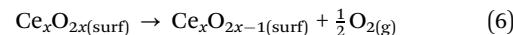


The formation energy was found to be  $0.841$  eV, indicating that the bulk Ni substitutional defect is thermodynamically less stable compared to Ni substitution on ceria surfaces, as shown in Fig. 7 and 8. Recent experiments<sup>65</sup> have indicated that increasing the Ni dopant percentage in a ceria nanorod resulted in the segregation of a  $\text{NiO}$  phase and a reduction in the catalytic activity. However, we have not considered the segregation of the  $\text{NiO}$  phase in our calculations.

As concluded above, forming a 'charge compensating' oxygen vacancy is another means to achieve charge neutrality on a Ni-substituted surface. In order to understand better the relationship between the charge-compensating vacancy and Ni on the  $\text{CeO}_2(100)$  surface, we have considered four possible configurations of the charge-compensating O vacancy on the Ni-substituted (100) surface (see Fig. S11 in ESI†) and the most stable configuration is shown in Fig. 8b. As shown in Fig. 8b, the charge-compensating O vacancy prefers to be incorporated into the host matrix close to the Ni substituent. The Ni atom relaxes outward and coordinates to four O atoms. The local structure around the Ni atom was found to be asymmetric with the Ni–O bond lengths varying between  $1.981$  and  $2.011$  Å, close to the Ni–O bond length of  $2.085$  Å in  $\text{NiO}_{(s)}$ . We computed the O-vacancy formation energy using eqn (4) to be  $-0.947$  eV.

In fact, what is of interest in catalysis is the cost to remove the next so-called "active vacancy". Hence, the formation of a second vacancy on the Ni-substituted (110) and (100) surfaces will be discussed in the next section.

**C. Formation of active oxygen vacancies from pure, Ni adsorbed and Ni doped  $\text{CeO}_2(110)$  and (100) surfaces.** The effectiveness of ceria as a catalyst lies in its ability to participate in redox reactions through the creation and compensation of oxygen vacancies on the surface. Thus, any modification to the surface that results in a reduction in the energy cost associated with the creation of a surface O vacancy is expected to present a larger catalytic activity. We have computed the active O vacancy formation energy from pure, Ni-adsorbed and Ni-substituted  $\text{CeO}_2(110)$  and (100) surfaces according to eqn (6)–(8).



In eqn (8),  $y = 1$  for the mono-Ni substituted surface and  $y = 2$  for the di-Ni substituted surface. The creation of an active O vacancy results in the partial reduction of two of the surface  $\text{Ce}^{4+}$  ions to  $\text{Ce}^{3+}$  ions. As noted earlier, a number of different distributions of  $\text{Ce}^{3+}$  ions on the surface are possible for a given vacancy location. The most stable configurations for the (110) and (100) surfaces are shown in Fig. 9a and 10a, respectively. The formation energy of an oxygen vacancy from the pure  $\text{CeO}_2(110)$  surface computed by us ( $1.414$  eV) compares favourably with the value reported by Kullgren *et al.*<sup>64</sup> ( $1.56$  eV, refer to Table III of their paper).

For the Ni-adsorbed (110) surface, we showed that  $\text{Ni}^+$  and  $\text{Ni}^{2+}$  oxidation states for the adsorbate were close in energy. Thus we have investigated active O-vacancy formation beginning with both  $\text{Ni}^+$  and  $\text{Ni}^{2+}$  configurations for the adsorbate.

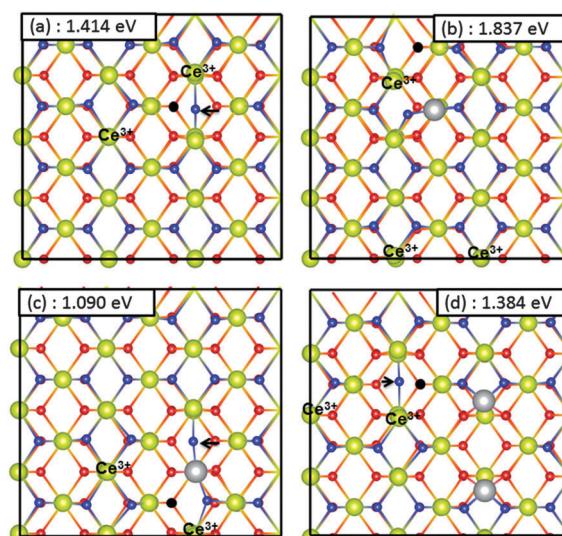
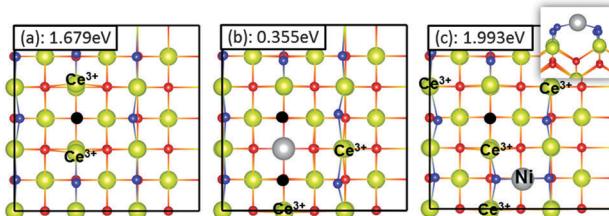


Fig. 9 Active oxygen vacancy on (a) pure, (b) Ni adsorbed, (c) mono-Ni substituted and (d) di-Ni substituted  $\text{CeO}_2(110)$  surface.  $\text{Ce}^{4+}$  ions are yellow,  $\text{Ce}^{3+}$  ions are blue,  $\text{O}^{2-}$  ions are red and  $\text{Ni}^+/\text{Ni}^{2+}$  ions are grey. The position of the oxygen vacancy is shown by black spheres. Arrows in (a), (c) and (d) show the motion of an oxygen atom from its surface site to a bridging position.





**Fig. 10** The optimized atomic structure of (100) surface with active oxygen vacancies. (a) Pure (100) surface, (b) Ni substituted (100) surface, (c) Ni adsorbed (100) surface. Ce atoms are yellow, Ni atoms is grey, O atoms are red and blue. The position of the O vacancy is represented by a black sphere. The inset shows the side view corresponding to the structure in panel (c).

We have investigated two different vacancy locations for the mono-Ni substituted surface and three different locations for the di-Ni substituted surface. For each of these locations various distributions of  $\text{Ce}^{3+}$  ions were investigated. Fig. 9 shows the most favourable configuration of an active O vacancy on all the four (110) surfaces studied here while Fig. S5–S8 in the ESI† show the structures and formation energies for various other vacancy positions.

For the Ni-adsorbed (100) surface, we have considered two possible configurations for the coexistence of a Ni adsorbate and active oxygen vacancy defect. As discussed above, two  $\text{Ce}^{3+}$  ions are created on the  $\text{CeO}_2(100)$  surface by the formation of an oxygen vacancy or the adsorption of a Ni atom due to the transfer of electrons from the adsorbate or the vacancy defect to  $\text{Ce}^{4+}$ . Therefore, the coexistence of a Ni adsorbate and an oxygen vacancy is expected to lead to a partially reduced surface with four  $\text{Ce}^{3+}$  ions. The most stable configuration is shown in Fig. 10c while the other configurations that we have studied are shown in Fig. S12 of the ESI†. The distortions around the oxygen vacancy and Ni on this surface are similar to those found around an oxygen vacancy on the pure (100) surface and Ni-adsorbed (100) surface, as shown in Fig. 6a and 10a.

From Fig. 9 and 10, we see that modifying the  $\text{CeO}_2(110)$  and (100) surfaces by a Ni substitution results in a reduction in the energy cost associated with the creation of an active oxygen vacancy compared to the pristine surface, though a higher doping percentage results in a smaller change in the vacancy formation energy on the (110) surface. On the other hand, the mono-Ni adsorbate on both (110) and (100) surfaces makes it more unfavourable to create an oxygen vacancy from the surface which could lead to suppression of the activity of the catalyst. This indicates that there exists an optimal level of surface doping to achieve enhanced catalytic activity. These results are in agreement with recent experimental data which shows that increasing Ni dopant percentage beyond 7% results in a reduction in the catalytic activity of ceria nanorods.<sup>65</sup>

For mono-Ni substituted (110) and (100) surfaces, we find that the most favourable site for oxygen vacancy formation is an oxygen ion site coordinated to the Ni dopant (see Fig. 9c and 10b). On the mono-Ni substituted (110) surface, this results in significant structural rearrangement with the Ni dopant moving

from a square-planar-like position to underneath a lattice metal ion site. Nolan<sup>45</sup> on the other hand reports the most favourable site for the creation of an active oxygen vacancy on the mono-Ni substituted (110) surface to be one of the next nearest lattice sites to the Ni dopant (see Fig. S7(b) in the ESI†) with a vacancy formation energy of 1.30 eV. For the same configuration, we have computed the vacancy formation energy to be 1.39 eV (see Fig. S7(b) in the ESI†). Furthermore, the formation of an oxygen vacancy results in the migration of a neighbouring  $\text{O}^{2-}$  ion to a bridging position on the (110) surface (Fig. 9a, c and d; arrows indicate the direction of motion of the oxygen atom). This rearrangement of an  $\text{O}^{2-}$  ion was not observed on the (100) surface. Comparing the formation energy of an active oxygen vacancy on the pure (110) with (111) and (100) surfaces (1.983 and 1.679 eV respectively), we find that (110) is the most active amongst the three surfaces. However, upon doping the surface with a Ni substituent, the formation energy of an active O vacancy from the (100) surface reduced dramatically to only 0.355 eV, indicating that surface doping has a larger effect on the  $\text{CeO}_2(100)$  surface as compared to (110).

## 4. Conclusions and concluding remarks

In this work, we have undertaken a systematic study of the interaction of a Ni adsorbate and Ni substituents with the  $\text{CeO}_2(110)$  and (100) surfaces using density functional theory calculations. The following conclusions can be derived from this work:

(1) A Ni atom prefers to adsorb on a bridge site as compared to a hollow site on the (110) surface, and the adsorption of Ni induces strong geometric distortions locally around the adsorbate. On the (100) surface, the most stable position of a Ni adsorbate was found to be the bridge site where the Ni atom is coordinated to two surface O atoms.

(2) On a bridge site, up to 0.25 ML coverage on the (110) surface, a Ni adsorbate prefers its  $\text{Ni}^+$  oxidation state over the  $\text{Ni}^{2+}$  oxidation state with the reduced  $\text{Ce}^{3+}$  ion located next to the  $\text{Ni}^+$  adsorbate. Location of the  $\text{Ce}^{3+}$  ion far away from the  $\text{Ni}^+$  adsorbate or at a subsurface position was found to be unfavourable. For the  $\text{Ni}^{2+}$  oxidation state, the most favourable distribution of  $\text{Ce}^{3+}$  ions was found to be when they were located diagonally across the  $\text{Ni}^{2+}$  adsorbate on the neighbouring sites. On the (100) surface, a Ni adsorbate prefers its  $\text{Ni}^{2+}$  oxidation state over the  $\text{Ni}^+$  oxidation state. The  $\text{Ce}^{3+}$  ions on the (100) surface prefer to stay in the vicinity of the adsorbed Ni or the oxygen vacancy, while they prefer to separate from the Ni adatom or the oxygen vacancy on the (111) surface.

(3) As the coverage of Ni on  $\text{CeO}_2(110)$  surface increases from 0.083 ML to 0.25 ML, the adsorption energy for the most favourable configuration, *i.e.*,  $\text{Ni}^+$  on a bridge site, decreases from  $-0.449$  eV to  $-0.322$  eV.

(4) With increase in coverage of Ni from 0.083 ML to 0.25 ML, the stabilization of the  $\text{Ni}^+$  oxidation state over the  $\text{Ni}^{2+}$  oxidation state on a bridge site decreases from *ca.* 49 meV



to 21 meV for the most preferred Ce<sup>3+</sup> ion distribution in the respective oxidation states of the Ni adsorbate.

(5) Adsorption of Ni introduces split-off-states above the CeO<sub>2</sub>(110) and (100) valence band. In its Ni<sup>+</sup> oxidation state on the (110) surface, these states arise from the f orbital on the reduced Ce<sup>3+</sup> ion, the d orbital of the Ni<sup>+</sup> adsorbate and the p orbital of the oxygen ions that are coordinated to the Ni<sup>+</sup> adsorbate. However, in its Ni<sup>2+</sup> oxidation state on both (110) and (100) surfaces, these states arise from the f orbitals on the two reduced Ce<sup>3+</sup> ions only.

(6) In contrast to Ni adsorption, a Ni substituent does not induce the reduction of a Ce<sup>4+</sup> to Ce<sup>3+</sup>. However, it induces strong structural distortion locally, with the Ni substituent moving into a four-coordinated nearly planar environment from the six-coordinated metal ion site on the (110) surface.

(7) The nearest neighbour locations were found to be the most favourable distribution of the two Ni substituents on the (110) surface. Further, this resulted in spontaneous ejection of an O<sub>2</sub> molecule from the surface.

(8) Substitution of surface Ce atoms by Ni atoms resulted in a reduction in the energy required to create an active oxygen vacancy while a Ni adsorbate caused an increase in the formation energy of an active O vacancy. The degree of reduction in the vacancy formation energy points to the presence of an optimum level of surface doping to achieve enhanced catalytic activity.

The effect of these surface modifications on the catalytic activity of ceria catalysts for WGS and methanation reactions is a subject of an ongoing study in our group.

## Notes

W.Q. Li and S. G. S are joint first authors in this article.

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