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Kinetic isotope effects for fast deuterium and proton exchange rates

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By monitoring the effect of deuterium decoupling on the decay of transverse ¹⁵N magnetization in D–¹⁵N spin pairs during multiple-refocusing echo sequences, we have determined fast D–D exchange rates *k*_D and compared them with fast H–H exchange rates *k*_H in tryptophan to determine the kinetic isotope effect as a function of pH and temperature.

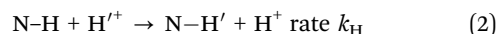
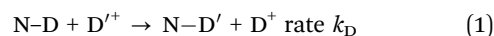
Introduction

In the parlance of magnetic resonance, chemical exchange is a process where a nucleus undergoes a change of its environment.¹ The determination of the exchange rates of labile protons can provide valuable insight into both structural and dynamic aspects of a wide range of molecules,^{2–4} such as the opening of base-pairs in nucleic acids and protection factors in protein–ligand complexes.^{5,6} In this paper, we shall focus on measurements of D–D exchange rates *k*_D and their comparison with H–H exchange rates *k*_H in tryptophan.⁷ The knowledge of kinetic isotope effects, *i.e.*, of the ratio *k*_H/*k*_D that expresses the reduction of D–D exchange rates *k*_D compared to H–H exchange rates *k*_H, may contribute to the characterization of reaction mechanisms.^{8,9} The kinetic isotope effect can give insight into the stability of hydrogen-bonded secondary structures in biomolecules.¹⁰ In this work, we shall consider exchange processes involving labile D^N deuterons and H^N protons that are covalently bound to the nitrogen atom in the indole ring of tryptophan.[†]

Experimental section

We have adapted to the case of deuterium (spin *S* = 1) a scheme that was originally designed to determine fast exchange rates of protons^{7,11,12} (spin *S* = 1/2) by monitoring the effect of

deuterium decoupling on the decay of transverse ¹⁵N magnetization during multiple-refocusing sequences (CPMG).^{‡13,14} The modified pulse sequence is shown in Fig. 1. The scheme requires isotopic enrichment with ¹⁵N and ¹³C, since the ¹⁵N coherence is excited by transfer from neighboring protons through two successive INEPT transfer steps *via* ¹*J*(¹³C, ¹H) and ¹*J*(¹⁵N, ¹³C). The decay of the ¹⁵N coherence is monitored indirectly after transferring the coherence back to the proton of origin. The ¹⁵N, ¹³C-labelled isomers of tryptophan are dissolved in either D₂O or H₂O to determine the kinetic isotope effect *k*_H/*k*_D of the following reactions:



where *k*_D and *k*_H are the pseudo-first order rate constants since the concentration of the solvent D₂O or H₂O, which is the source of the incoming D⁺ or H⁺ ions, is constant and much higher than the concentration of the solute.[§]

The first and last parts of the pulse sequence in Fig. 1 lead to a transfer of the magnetization from the blue non-exchanging 'spy' proton to ¹⁵N and back, *via* the adjacent ¹³C nuclei, by two successive pulse sequences for Insensitive Nuclei Enhanced by Polarization Transfer (INEPT).¹⁵ The first INEPT sequence transforms longitudinal proton magnetization *H*_z into two-spin order 2*H*_z*C*_z. The second INEPT sequence converts 2*H*_z*C*_z into 2*C*_z*N*_z. WALTZ-16 proton decoupling¹⁶ is used to suppress the evolution under ¹*J*(¹H, ¹³C) during the intervals of the INEPT sequences

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[†] Despite IUPAC recommendations, we use the notation ²D rather than ²H.

[‡] We use the symbol ²D when referring to isotopes as in the expressions ¹*J*(¹H, ¹⁵N) or ¹*J*(²D, ¹⁵N).

[§] We shall refer to H or D for atoms that appear in molecular formulae and to H^N or D^N in N–H and N–D groups. For the Cartesian components of angular momentum operators, we have used *H*_x, *H*_y, *H*_z, *D*_x, *D*_y, *D*_z, *N*_x, *N*_y, *N*_z, *C*_x, *C*_y, *C*_z rather than the common notation *I*_x, *R*_x, *S*_x, *etc.*



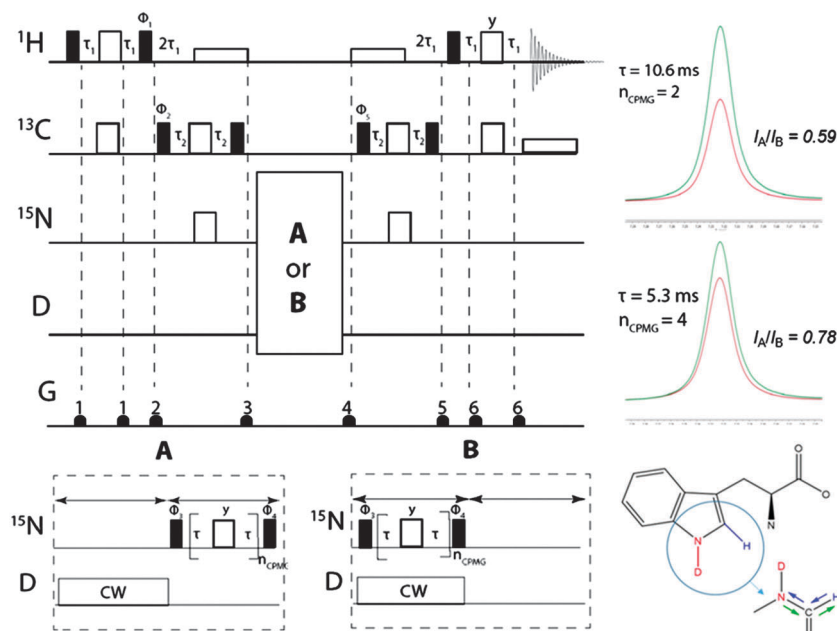


Fig. 1 (left) Pulse sequence for measurements of the indole D–D exchange rate k_D . The $\pi/2$ and π pulses are represented by narrow filled and wide open rectangles respectively while wide open rectangles depict decoupling sequences. All phases are along the x-axis unless indicated otherwise. The phase cycling is: $\Phi_1 = 16(y), 16(-y)$; $\Phi_2 = x, -x$; $\Phi_3 = 2(x), 2(-x)$; $\Phi_4 = 4(x), 4(-x)$; $\Phi_5 = 8(x), 8(-x)$ and the receiver phase is $\Phi_{rec} = x, -x, -x, x, 2(-x, x, x, -x), x, -x, -x, x, -x, x, -x, 2(x, -x, -x, x), -x, x, x, -x$. The delays are: $\tau_1 = 1/(4J_{CH}) = 1.56$ ms and $\tau_2 = 1/(4J_{CN}) = 16$ ms. The gradient pulses G that bracket π -pulses at positions 1 and 6 are of equal strength and polarity to cancel effects of pulse imperfections. The gradients applied at positions 2, 3, 4 and 5 are used to purge any undesired transverse magnetization, as the magnetization of interest is aligned with z during the corresponding intervals. (top right) Proton signals of tryptophan at pD 8.7 and $T = 300$ K. The red lines correspond to experiment A without decoupling while the green lines stem from experiment B with deuterium decoupling. (top) $I_A/I_B = 0.59$ with $\tau = 10.6$ ms and $n_{CPMG} = 2$. (bottom) $I_A/I_B = 0.78$ with $\tau = 5.3$ ms and $n_{CPMG} = 4$. (bottom right) Consecutive coherence transfer steps from the blue 1H to the red ^{15}N via ^{13}C and back in 2D , ^{13}C , and ^{15}N labelled tryptophan.

where the coherence is transferred from ^{13}C to ^{15}N . The anti-phase coherence $2N_yC_z$ excited at the beginning of the multiple-refocusing CPMG interval decays in the course of this pulse train. At this point, two variants (A and B) of the experiments must be performed. In experiment B, continuous wave (CW) deuterium decoupling is applied during the CPMG pulse train, while in experiment A the deuterium irradiation is applied for the same duration but prior to the CPMG pulse train in order to avoid differences in temperature.

The remaining coherence $2N_yC_z$ is transferred back to the 'spy' proton for detection. The intensity of the resulting peak near 7.22 ppm in the proton spectra is proportional to the magnitude of the nitrogen $2N_yC_z$ coherence that remains at the end of the CPMG interval. In order to extract k_D one can determine the ratio I_A/I_B of the peak intensities recorded without decoupling during the CPMG pulse train (experiment A) and with deuterium decoupling (experiment B). The delay τ is defined as one-half of the interval between consecutive nitrogen π -pulses. The τ delays need to be long enough to ensure that the ratio I_A/I_B is significantly different from 1. Typically, values of $\tau = 10.6$ or 21.2 ms have been used. The scalar coupling is $^1J(^{15}N, ^2D) = 15.4$ Hz, smaller than $^1J(^{15}N, ^1H) = 98.6$ Hz by the factor $\gamma(^2D)/\gamma(^1H) \approx 0.15$, but $^1J(^{15}N, ^2D)$ is still large enough to act as an efficient vehicle of scalar relaxation.

We can construct the matrix representations of the $4 \times 9 = 36$ Cartesian operators that span a complete basis set

for a system comprising a ^{15}N nucleus with spin $I = 1/2$ and a 2D nucleus with spin $S = 1$.¹⁷ When a CPMG multiple echo sequence is applied to the ^{15}N spins with an on-resonance rf field at the chemical shift of ^{15}N , while deuterium decoupling is applied with an amplitude ω_1^D at an offset Ω^D with respect to the chemical shift of 2D . The rf pulses applied to the ^{15}N spins are considered to be ideal. Starting from an operator N_y , coherent evolution leads to the following terms:¹⁸

$$\begin{aligned} N_y &\xleftrightarrow{J_{N-D}} N_x D_z \xleftrightarrow{J_{N-D}} N_y (3D_z^2 - 2E) \xleftrightarrow{\omega_1^D} N_y (D_z D_y + D_y D_z) \\ &\xleftrightarrow{\omega_1^D} N_y (D_x^2 - D_y^2) \xleftrightarrow{\Omega^D} N_y (D_x D_y + D_y D_x) N_x D_z \xleftrightarrow{\omega_1^D} N_x D_y \\ &\xleftrightarrow{J_{N-D}} N_y (D_z D_y + D_y D_z) \xleftrightarrow{\Omega^D} N_y (D_z D_x + D_x D_z) N_x D_y \xleftrightarrow{\Omega^D} N_x D_x \end{aligned} \quad (3)$$

Therefore the dimension of the basis set can be reduced from 36, leaving only 9 terms:

$$\begin{aligned} &\left\{ 2\sqrt{\frac{2}{3}} N_y, 2N_x D_z, 2\sqrt{2} N_y (3D_z^2 - 2E), 2N_x D_y, 2\sqrt{2} N_y (D_z D_y + D_y D_z), \right. \\ &2\sqrt{2} N_y (D_x^2 - D_y^2), 2N_x D_x, 2\sqrt{2} N_y (D_z D_x + D_x D_z), \\ &\left. 2\sqrt{2} N_y (D_x D_y + D_y D_x) \right\} \end{aligned} \quad (4)$$

Note that in the experiments of Fig. 1, the single-quantum coherence at the beginning of the CPMG period is an anti-phase operator $2N_yC_z$. Since the presence of the C_z term affects the signal intensities in experiments A and B equally, this C_z term can be omitted without loss of generality. The solution of the Liouville-von Neumann equation¹⁹ up to the n^{th} echo is:

$$\sigma(t = 2n\tau) = [\exp(-L\tau) \cdot R_N \cdot \exp(-L\tau)]^n \sigma(0) \quad (5)$$

The matrix representation of the Liouvillian L in the basis of eqn (4) is:

$$L = \begin{bmatrix} 0 & 2\sqrt{\frac{2}{3}}J\pi & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -2\sqrt{\frac{2}{3}}J\pi & k & -\frac{2}{\sqrt{3}}J\pi & \omega_1^D & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \frac{2}{\sqrt{3}}J\pi & k & 0 & \sqrt{3}\omega_1^D & 0 & 0 & 0 & 0 & 0 \\ 0 & -\omega_1^D & 0 & k & -J\pi & 0 & \Omega_D & 0 & 0 & 0 \\ 0 & 0 & -\sqrt{3}\omega_1^D & J\pi & k & -\omega_1^D & 0 & \Omega_D & 0 & 0 \\ 0 & 0 & 0 & 0 & \omega_1^D & k & 0 & 0 & 0 & 2\Omega_D \\ 0 & 0 & 0 & -\Omega_D & 0 & 0 & k & -J\pi & 0 & 0 \\ 0 & 0 & 0 & 0 & -\Omega_D & 0 & J\pi & k & -\omega_1^D & 0 \\ 0 & 0 & 0 & 0 & 0 & -2\Omega_D & 0 & \omega_1^D & k & 0 \end{bmatrix} \quad (6)$$

The matrix representation of R_N represents a π_y pulse applied to the ^{15}N spins:

$$R_N = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad (7)$$

If the rf field for deuterium decoupling is applied on resonance, the evolution of the density operator can be described in a simplified base comprising only 6 product operators:

$$\left\{ 2\sqrt{\frac{2}{3}}N_y, 2N_xD_z, 2\sqrt{2}N_y(3D_z^2 - 2E), 2N_xD_y, 2\sqrt{2}N_y(D_zD_y + D_yD_z), 2\sqrt{2}N_y(D_x^2 - D_y^2) \right\} \quad (8)$$

In this reduced base, the matrix representation of the Liouvillian is:

$$L = \begin{bmatrix} 0 & 2\sqrt{\frac{2}{3}}J\pi & 0 & 0 & 0 & 0 \\ -2\sqrt{\frac{2}{3}}J\pi & k & -\frac{2J\pi}{\sqrt{3}} & \omega_1^D & 0 & 0 \\ 0 & \frac{2J\pi}{\sqrt{3}} & k & 0 & \sqrt{3}\omega_1^D & 0 \\ 0 & -\omega_1^D & 0 & k & -J\pi & 0 \\ 0 & 0 & -\sqrt{3}\omega_1^D & J\pi & k & -\omega_1^D \\ 0 & 0 & 0 & 0 & \omega_1^D & k \end{bmatrix} \quad (9)$$

For the π_y pulse applied to the ^{15}N spins one obtains in this reduced base:

$$R_N = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix} \quad (10)$$

In the experiment of Fig. 1, the amplitude ν_1^D of the continuous-wave rf field applied to the deuterium spins should be chosen carefully. The higher the rf amplitude ν_1^D , the more efficient the decoupling, although one should avoid excessive heating. On the other hand, if the rf amplitude is too low, the ratio I_A/I_B is affected in a manner that can lead to erroneous measurements of the exchange rates. By way of illustration, at pD 7.7 and $T = 300$ K, where the exchange rate is very low (see Table 1), the ratio I_A/I_B has been determined as a function of the rf amplitude for $\tau = 10.6$ ms and $n_{\text{CPMG}} = 2$. For these experimental conditions, the amplitude can be attenuated as low as $\nu_1^D = \omega_1^D/(2\pi) = 100$ Hz without affecting significantly the ratio I_A/I_B . For lower amplitudes the ratio is very sensitive to the exact amplitude. An rf field with an



Table 1 Pseudo first-order exchange rate constants k_D [s^{-1}] without corrections for contributions due to quadrupolar relaxation as a function of temperature and pD

pD	290 K	pD	300 K	pD	310 K	pD	320 K
1.05	273 ± 21	1.05	491 ± 63	1.05	697 ± 110	1.05	1670 ± 250
1.49	56.9 ± 11	1.49	65.9 ± 13	1.49	81.6 ± 17	1.49	91.6 ± 19
2.18	49.7 ± 4.8	2.18	47.9 ± 4.2	2.18	57.6 ± 11	2.18	61.8 ± 8.7
3.29	39.1 ± 1.9	3.29	30.2 ± 1.8	3.29	26.3 ± 1.5	3.29	23.6 ± 0.7
4.78	36.2 ± 3.0	4.78	26.2 ± 3.3	4.78	20.9 ± 2.9	4.78	16.5 ± 3.2
5.98	37.1 ± 3.1	5.98	27.1 ± 1.9	5.98	22.9 ± 1.9	5.98	18.7 ± 1.4
6.43	38.7 ± 5.7	6.43	28.6 ± 4.2	6.43	23.7 ± 3.1	6.43	19.4 ± 2.6
7.97	41.4 ± 6.4	7.69	31.6 ± 4.3	7.41	27.2 ± 3.4	7.13	23.8 ± 3.1
9.11	82.4 ± 15	8.83	88.1 ± 8.2	8.55	110 ± 11	8.27	182 ± 31
9.68	121 ± 13	9.40	328 ± 17	9.12	507 ± 24	8.84	860 ± 49
10.8	831 ± 42	10.52	1546 ± 90	10.24	2670 ± 330	9.96	4330 ± 380
12.1	3220 ± 840	11.82	8060 ± 2200	11.54	11 800 ± 2600	11.26	18 800 ± 1700
12.97	14 000 ± 3400	12.69	17 600 ± 6300	12.41	40 400 ± 11 000	—	—

amplitude $\nu_1^D = 3$ kHz seems to be a safe value regardless of the exchange rates and can be used for all experiments. The ratio I_A/I_B also depends on the offset Ω^D of the rf carrier with respect to the exchanging 2D spins, since decoupling becomes less efficient when the carrier is off-resonance. The ratio I_A/I_B has the smallest value when the carrier coincides with the chemical shift of the exchanging 2D spins, *i.e.*, when $\Omega^D = 0$ (Fig. 2). The heteronuclear scalar coupling constant $^1J(^{15}N, ^2H) = 15.4$ Hz at pD 7.7 was determined experimentally from the doublet in the 2H spectrum and corresponds to the expected value $^1J(^{15}N, ^2H) = ^1J(^{15}N, ^1H) \gamma(^2H)/\gamma(^1H)$ with $^1J(^{15}N, ^1H) = 98.6$ Hz.

All experiments were performed at 14.1 T (600 MHz for 1H , 151 MHz for ^{13}C , 92 MHz for 2H , and -61 MHz for ^{15}N) using a Bruker Avance III spectrometer equipped with a cryogenically cooled TXI probe. The samples were prepared by dissolving 20 mM tryptophan (fully ^{13}C and ^{15}N enriched) in 100% D_2O buffered with 20 mM citrate, acetate, Tris or phosphate buffer depending on the pH range. We determined k_H in our earlier work⁷ using 97% H_2O and 3% D_2O . The pH was adjusted by DCl or NaOD; the indicated pH values include corrections to take into account that the pH was measured in D_2O with an electrode calibrated for H_2O according to the following equation²⁰

$$pD = pH_{\text{apparent}} + 0.4 \quad (11)$$

Results and discussion

For each pH and temperature, the exchange rates k_D have been determined from three to seven ratios I_A/I_B of the signal intensities corresponding to six to fourteen experiments performed with variable numbers of π -pulses $2 \leq n \leq 8$ in the CPMG trains, and different intervals, $\tau = 2.6, 5.3, 10.6$ and 21.2 ms, but with the same total relaxation time $2\tau n_{\text{CPMG}}$. A minimum of two ratios I_A/I_B at different delays are required for an unambiguous determination of k_D , since two rates can be compatible with a single I_A/I_B ratio. Fig. 3 shows how this ambiguity is lifted by changing the inter-pulse delay 2τ in the CPMG pulse train. The pseudo first-order exchange rate constants were found to lie in a range $0 < k_D < 40\,000$ s^{-1} , depending on pH and temperature (Table 1). At each temperature, the exchange rate k_D was found to be slowest for pD_{min} 4.8.

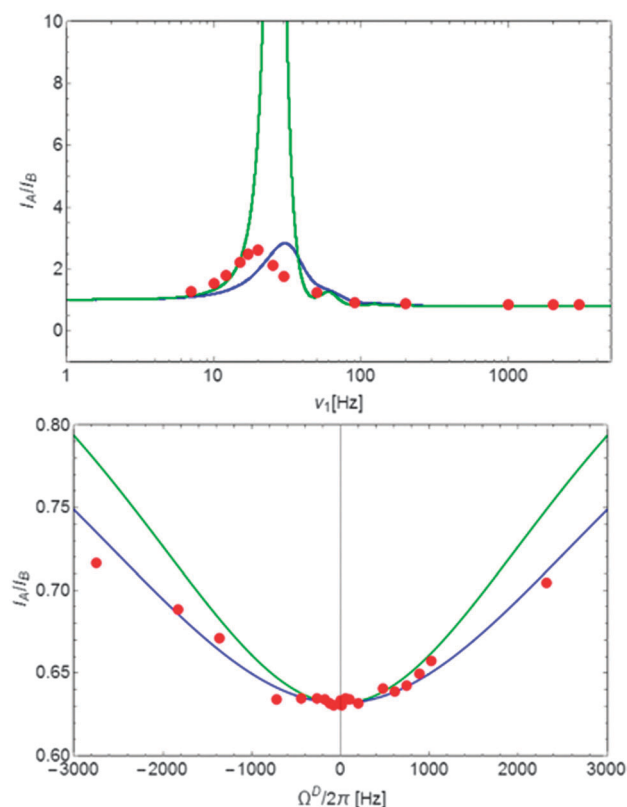


Fig. 2 (top) Experimental ratio I_A/I_B as a function of the amplitude ν_1^D of the rf field applied to the deuterium spins recorded at pD 7.7 and $T = 300$ K with $\tau = 10.6$ ms and $n_{\text{CPMG}} = 2$. Anomalous ratios $I_A/I_B > 1$ only occur when the rf amplitude is too low, in particular in the vicinity of $^1J(^{15}N, ^2H)$. (bottom) Experimental ratio I_A/I_B as a function of the offset Ω^D of the carrier frequency with respect to the deuterium resonance for pD 9.4, $T = 300$ K, $\tau = 10.6$ ms, and $n_{\text{CPMG}} = 2$. The lines correspond to eqn (9) (top) and eqn (6) (bottom). For the blue lines, we have assumed that different operator products involving deuterium terms have distinct quadrupolar relaxation rates that depend on the spectral density. For the green lines, we have assumed that all deuterium terms have the same relaxation rate. For strong on-resonance rf fields, as we have used for the determination of exchange rates, the ratios I_A/I_B do not change significantly if one assumes a single or several distinct relaxation rates.

When the exchange rate k_D is very low, one cannot neglect contributions due to the difference in relaxation rates of the in-phase ^{15}N coherence and other rates in the relaxation matrix



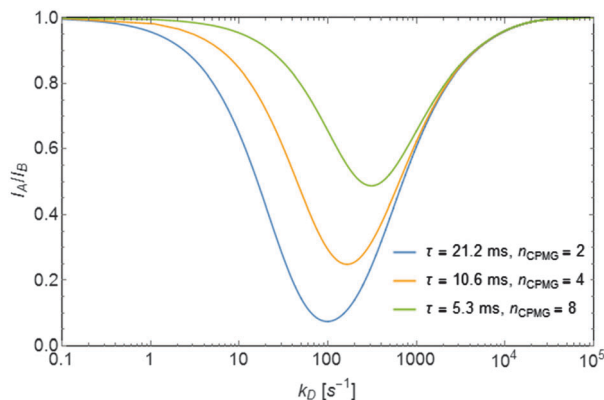


Fig. 3 Simulated ratios I_A/I_B as a function of the exchange rates k_D . The curves correspond to $\tau = 21.2$ ms and $n_{\text{CPMG}} = 2$, $\tau = 10.6$ ms and $n_{\text{CPMG}} = 4$, and finally $\tau = 5.3$ ms and $n_{\text{CPMG}} = 8$, keeping the total time $2\tau n_{\text{CPMG}}$ constant.

of eqn (6). From an earlier study of the exchange of indole protons,⁷ we know that the exchange rate k_H almost vanishes near pH_{min} . On the other hand, as can be seen in Table 1, the exchange rates k_D do not vanish near pD_{min} . Moreover, if one neglects relaxation of deuterium, some apparent exchange rates increase at lower temperatures, which is physically impossible. Hence, we incorporated a temperature-dependent quadrupolar relaxation rate R_Q in eqn (12) and subtracted it from the apparent exchange rates at all pD. The use of a single constant R_Q to describe the effects of deuterium relaxation is rather naive. In particular for weak rf fields or large deuterium offsets, this assumption may lead to errors. We can calculate the relaxation rates of operator products containing terms such as D_z , $(3D_z^2 - 2E)$, D_x , D_y , $(D_x^2 - D_y^2)$, $(D_y D_z + D_z D_y)$ and $(D_x D_z + D_z D_x)$.²¹ However, we have verified that under the conditions for which the rates of Fig. 4 were obtained, i.e., for strong rf fields and vanishing deuterium offsets, the exchange rates are barely affected if we assume that all deuterium terms have a common relaxation rate. The errors in the experimental ratios I_A/I_B were determined from standard deviations. The error propagation was further simulated by the Monte Carlo technique. The errors in the exchange rates k_D were estimated from the curvature around the minima of χ^2 and found to lie in a range between 3 and 28%.

If the exchange rate constants k_D are plotted as a function of pD on a logarithmic scale, one obtains a V-shaped curve that is characteristic of acid catalysis by D^+ ions and basic catalysis by OD^- ions, the latter being more efficient (Fig. 4). In the cationic, zwitterionic and anionic forms of tryptophan, the exchange rates result from sums of acidic and basic contributions. The overall exchange rate constant k_D can be written as:^{2,22}

$$k_D = k_{\text{D}}^c [\text{D}^+]_c + k_{\text{D}}^z [\text{D}^+]_z + k_{\text{OD}}^z [\text{OD}^-]_z + k_{\text{OD}}^a [\text{OD}^-]_a + R_Q \quad (12)$$

where the rate R_Q expresses contributions due to the quadrupolar deuterium relaxation to the decay of antiphase ^{15}N coherences. The indices D and OD represent the contributions of acidic and basic mechanisms (see below) for the cationic, zwitterionic and anionic forms of tryptophan, abbreviated by c , z , and a in Fig. 5.

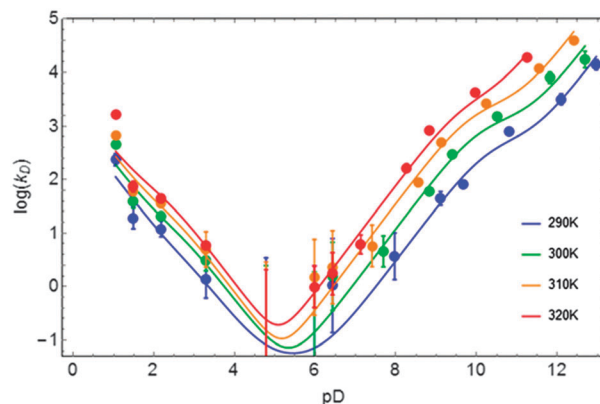


Fig. 4 Exchange rate constants k_D with corrections of Table 2 for the contributions due to quadrupolar relaxation as a function of pD over the temperature range $290 \leq T \leq 320$ K. Solid lines result from fits to eqn (12).

The mole fractions f_c , f_z and f_a of the cationic, zwitterionic and anionic forms of tryptophan are:

$$\begin{aligned} f_c &= (1 + 10^{\text{pD}-\text{pK}_{a1}} + 10^{2\text{pD}-\text{pK}_{a1}-\text{pK}_{a2}})^{-1} \\ f_z &= (1 + 10^{-\text{pD}+\text{pK}_{a1}} + 10^{\text{pD}-\text{pK}_{a2}})^{-1} \\ f_a &= (1 + 10^{-\text{pD}+\text{pK}_{a2}} + 10^{-2\text{pD}+\text{pK}_{a1}+\text{pK}_{a2}})^{-1} \end{aligned} \quad (13)$$

Where $[\text{D}^+] = 10^{-\text{pD}}$, $[\text{OD}^-] = K_W 10^{\text{pD}}$. The auto-ionization constant K_W of D_2O depends on the temperature.²³ In H_2O at 25°C , $\text{pK}_{a1} = 2.46$ for the protonation of the carboxyl group, while $\text{pK}_{a2} = 9.41$ corresponds to the protonation of the amine group. In D_2O at 25°C , we have determined that $\text{pK}_{a1} = 2.60$ and $\text{pK}_{a2} = 10.05$.²⁴ The variation of pK_a with temperature²³ has been taken into account. Fig. 4 and Table 2 show the results of the fitting of the exchange rate constants k_D to eqn (12), which allows one to obtain the catalytic rate constants for the contributions of acid and basic mechanisms for each of the three forms c , z , and a . The basic contribution of the cationic form and the acidic contribution of the anionic form are masked by other terms and can be neglected.

The activation energy E_a of the transition state provides a measure of the strength of N-D or N-H bonds.²⁵ The activation energy E_a is defined by the Arrhenius equation

$$k = A e^{-E_a/RT} \quad (14)$$

where A is an empirical pre-exponential “frequency factor”, R the universal gas constant, T the temperature and k the exchange rate. The dependence of E_a on pH or pD for H-H and D-D exchange processes and the activation energies and pre-exponential frequency factors are shown in Table 3 for protons and in Table 4 for deuterium.

One can speak of a kinetic isotope effect when the exchange rate is affected by isotopic substitution.²⁶ In the present case, we compare the exchange rates of indole protons in tryptophan with H_2O on the one hand, and analogous exchange rates of indole deuterons with D_2O on the other. The kinetic isotope effect is defined as the ratio of the rate constants k_H/k_D . The change in exchange rates results from differences in the vibrational frequencies of the N-H or N-D bonds formed between ^{15}N and ^1H or ^2D .^{27–29} Deuterium will



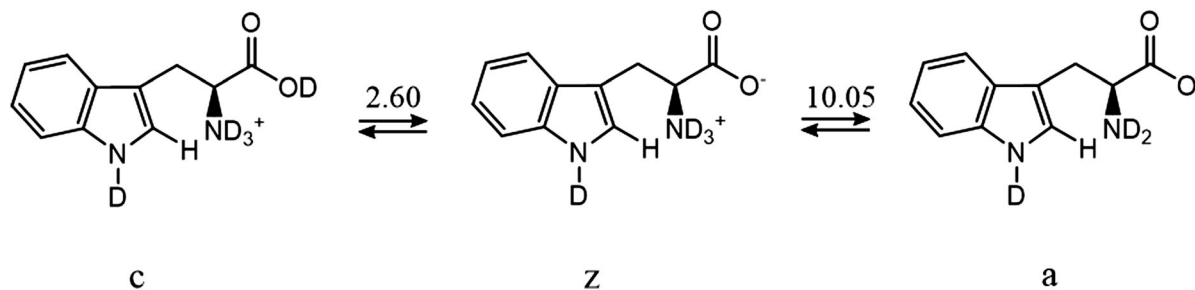


Fig. 5 Tryptophan exists in three forms *c* (cationic), *z* (zwitterionic) and *a* (anionic), with mole fractions f_c , f_z and f_a that depend on pD.

Table 2 Exchange rate constants k_D and k_H [s^{-1}] derived by fitting to eqn (12)

	290 K	300 K	310 K	320 K		300 K	310 K	320 K
R_Q	36.2 ± 5.6	26.51 ± 3.7	19.4 ± 9.7	16.6 ± 3.3	R_Q	0.37 ± 0.03	^a	^a
$\log(k_D^c)$	2.91 ± 0.64	3.31 ± 0.27	3.49 ± 0.32	3.01 ± 1.20	$\log(k_H^c)$	2.91 ± 0.04	^a	^a
$\log(k_D^z)$	3.80 ± 0.95	3.74 ± 0.64	3.70 ± 1.22	4.14 ± 0.24	$\log(k_H^z)$	3.31 ± 0.05	^a	^a
$\log(k_{OD}^z)$	7.89 ± 0.07	8.10 ± 0.06	8.28 ± 0.03	8.41 ± 0.08	$\log(k_{OH}^z)$	8.13 ± 0.02	8.29 ± 0.29	8.47 ± 0.41
$\log(k_{OD}^a)$	6.64 ± 0.20	6.69 ± 0.30	6.87 ± 0.12	7.19 ± 0.19	$\log(k_{OH}^a)$	7.53 ± 0.05	7.72 ± 0.27	7.97 ± 0.37

^a Proton exchange rates were not measured at these temperatures.⁷

Table 3 Activation energies E_a and pre-exponential frequency factors A for the indole proton H^N in tryptophan

pH	E_a [$kJ\ mol^{-1}$]	$\ln(A)$
6.3	88 ± 2	37 ± 1
7.41	84 ± 2	37 ± 1
8.31	83 ± 4	39 ± 1
9.08	82 ± 6	40 ± 2
10.01	86 ± 14	43 ± 5
10.6	94 ± 12	47 ± 4

Table 4 Activation energies E_a and pre-exponential frequency factors A for the indole deuterium D^N in tryptophan. The activation energies and the pre-exponential factors are strongly correlated

pD	E_a [$kJ\ mol^{-1}$]	$\ln(A)$
1.0	38 ± 17	20 ± 7
1.5	40 ± 12	20 ± 5
2.0	35 ± 11	17 ± 4
7.0	87 ± 3	35 ± 1
8.0	86 ± 8	37 ± 3
9.0	82 ± 9	38 ± 3
10.0	73 ± 14	36 ± 6
11.0	72 ± 7	36 ± 3
12.0	87 ± 23	44 ± 9

lead to a lower vibrational frequency because of its heavier mass (lower zero-point energy). If the zero-point energy is lower, more energy is needed to break an N–D bond than to break an N–H bond, so that the rate of the exchange will be slower. Moreover, one expects E_a to be larger for deuterium. The results in Tables 3 and 4 do not support this expectation, but if one assumes the same pre-exponential frequency factor for H and D, E_a is indeed larger for the heavier isotope.

Fig. 6 shows exchange rates k_D and k_H at 300 K. For acid catalyzed exchange, $k_D/k_H > 2.5$ because D_3O^+ is a stronger acid than H_3O^+ . For base catalyzed exchange, $k_D/k_H < 1$. However, to

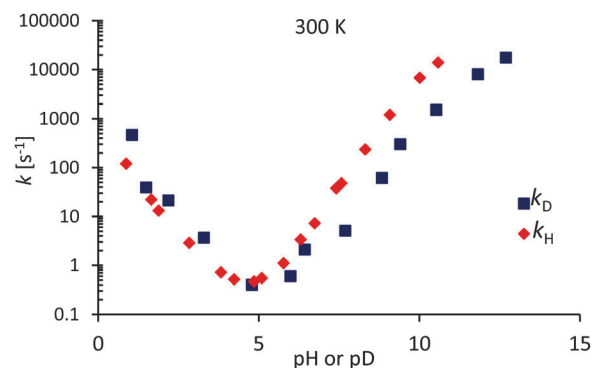


Fig. 6 Deuterium and proton exchange rates k_D (blue) and k_H (red) as a function of pH or pD at 300 K. The pD scale has been corrected according to eqn (11) to take into account the use in D_2O of a glass electrode designed for H_2O .

compare the difference between catalysis by OH^- and OD^- , we need to take into account the difference of the ionization constants: $pK_W(D_2O) = 14.95$ and $pK_W(H_2O) = 13.99$ at 25 °C.

Fig. 7 shows the base-catalyzed exchange rate constants k_D and k_H as a function of pOH or pOD. The exchange rates k_D are slightly lower than k_H , giving the approximate kinetic isotope effects: $k_H/k_D = 2.2 \pm 0.3$, 2.3 ± 0.3 and 2.1 ± 0.3 at 300 K, 310 K and 320 K respectively (Fig. 7). These values result from averages of the exchange rate constants for the zwitterionic and anionic forms (Table 5).

In Table 5 the KIE is defined as k_H^i/k_D^i for acid catalysis or as k_{OH}^i/k_{OD}^i for base catalysis, where $i = c, z$, and a stand for the cationic, zwitterionic, and anionic forms of tryptophan in solution, with the heaviest isotope always in the denominator. If tunneling can be neglected, the KIE depends on the nature of the transition state. The maximum isotope effect for N–H bonds is $k_H/k_D \approx 9$, assuming that the bond is completely broken in the



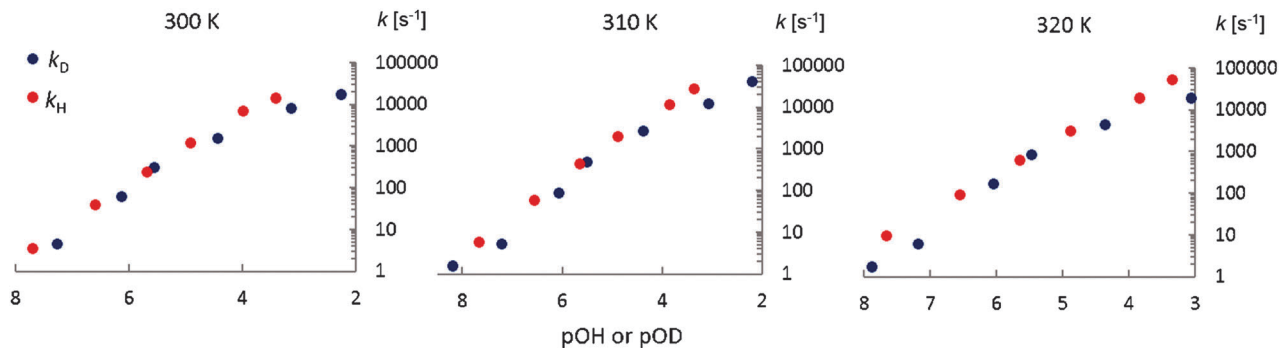


Fig. 7 Base catalyzed exchange rates k_D and k_H as a function of pOH or pOD at different temperatures.

Table 5 Kinetic isotope effects (KIE) k_H/k_D for the exchange rate constants of each of the three forms of tryptophan in solution: c (cationic), z (zwitterionic), and a (anionic)

	300 K	310 K	320 K
k_H^c/k_D^c	0.40 ± 0.04	^a	^a
k_H^z/k_D^z	0.37 ± 0.09	^a	^a
k_{OH}^z/k_{OD}^z	1.1 ± 0.1	1.0 ± 0.3	1.1 ± 0.5
k_{OH}^a/k_{OD}^a	7 ± 2	7 ± 2	6 ± 3

^a Proton exchange rates were not measured at these temperatures.⁷

transition state (TS). The KIE can be reduced if the bonds are not completely broken in the TS. The KIE can be close to 1 if the TS is very similar to the reactant (N–D bond nearly unaffected) or very similar to the product (N–D bond almost completely broken).

The experimental ratio k_{OH}^a/k_{OD}^a is near its maximum when $pH > pK_{a2}$, which suggests that the N–D bond is broken in the rate-limiting step and that the deuteron is half-way between the donor and the acceptor. However the ratio $k_{OH}^z/k_{OD}^z \approx 1$ suggests that the N–D bond is either only slightly or almost completely broken in the TS. The protonation of the amine withdraws electron density and increases the acidity of the H^N group which favors the formation of the anionic form. This explains why $k_{OH}^z/k_{OD}^z > 1$ and $k_{OD}^z/k_{OD}^a > 1$. For the acid-catalyzed exchange constants, we observe an inverse kinetic isotope effect. This can happen when the degree of hybridization of the reactant is lower than that of the reaction center in the TS during the rate-limiting step.

The mechanisms for proton or deuteron exchange have been thoroughly reviewed.^{30–32} Englander³⁰ and his collaborators pointed out that the rate of the exchange of protons attached to nitrogen depends on the ability to form hydrogen-bonded complexes in the transition state involving the donor (tryptophan) and the acceptor (D_2O or OD^-). This occurs in three steps: (i) encounter of the donor and the acceptor, (ii) formation of the transition state involving the donor and acceptor, and (iii) cleavage of the N–D bond. The mechanism of acid-catalyzed exchange consists of the addition onto the nitrogen of a D^+ ion from the solvent, followed by removal of D^+ by D_2O (Fig. 8). The mechanism of the base-catalyzed reaction involves removing the indole deuterium to create the conjugate base, which then abstracts a D^+ from D_2O to regenerate the indole (Fig. 9).

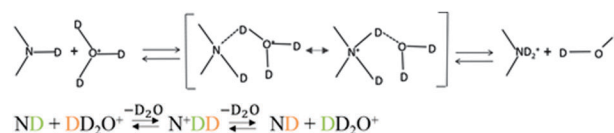


Fig. 8 Acid-catalyzed mechanism of exchange. The transition state is shown in brackets.

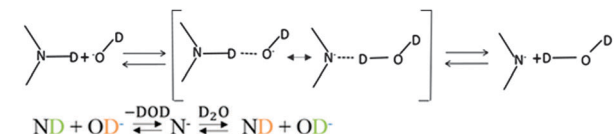


Fig. 9 Base-catalyzed mechanism of exchange. The transition state is shown in brackets.

Altogether we can say that the rate-limiting step in the base-catalyzed mechanism is the removal of the proton or deuteron from the nitrogen. On the other hand, for the acid-catalyzed mechanism, is the donation of a proton or deuteron by H_3O^+ respectively D_3O^+ . Finally, the curves of $\log k_D$ vs. pD and of $\log k_H$ vs. pH show a combination of specific base catalysis at high pH, and a specific acid catalysis at low pH, which becomes more important at higher temperatures.

Conclusions

We have adapted our method that was originally designed for measuring fast H–H exchange rates k_H to the study of D–D exchange rates k . In tryptophan in aqueous solution over a range of pH, respectively pD, the kinetic isotope effect, defined as the ratio k_H/k_D between the H–H and D–D exchange rates, was determined at several temperatures. The dependence of the activation energies on pH provides new insight into the mechanisms of the exchange processes. The results agree with the mechanisms discussed by Englander *et al.*³⁰

Abbreviations

CPMG	Carr Purcell Meiboom Gill
KIE	Kinetic isotope effect
TS	transition state



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Notes and references

- 1 A. D. Bain, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2003, **43**, 63–103.
- 2 C. E. Dempsey, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2001, **39**, 135–170.
- 3 M. Hoshino, H. Katou, Y. Hagihara, K. Hasegawa, H. Naiki and Y. Goto, *Nat. Struct. Biol.*, 2002, **9**, 332–336.
- 4 F. Persson and B. Halle, *Proc. Natl. Acad. Sci. U. S. A.*, 2015, **112**(33), 10383–10388.
- 5 M. Guéron, M. Kochoyan and J. L. Leroy, *Nature*, 1987, **328**, 89–92.
- 6 Y. Paterson, S. W. Englander and H. Roder, *Science*, 1990, **249**, 755–759.
- 7 F. Kateb, P. Pelupessy and G. Bodenhausen, *J. Magn. Reson.*, 2007, **184**, 108–113.
- 8 D. B. Northrop, *Biochemistry*, 1975, **14**(12), 2644–2651.
- 9 R. Sharma, T. J. Strelevitz, H. Gao, A. J. Clark, K. Schilknecht, R. S. Obach, S. L. Ripp, D. K. Spracklin, L. M. Tremaine and A. D. N. Vaz, *Drug Metab. Dispos.*, 2012, **40**, 625–634.
- 10 S. W. Englander, T. R. Sosnick, J. J. Englander and L. Mayne, *Curr. Opin. Struct. Biol.*, 1996, **6**, 18–23.
- 11 T. Segawa, F. Kateb, L. Duma, G. Bodenhausen and P. Pelupessy, *ChemBioChem*, 2008, **9**, 537–542.
- 12 A. A. Sehgal, L. Duma, G. Bodenhausen and P. Pelupessy, *Chem. – Eur. J.*, 2014, **20**, 6332–6338.
- 13 H. Carr and E. Purcell, *Phys. Rev.*, 1954, **94**, 630–638.
- 14 S. Meiboom and D. Gill, *Rev. Sci. Instrum.*, 1958, **29**, 688–691.
- 15 G. A. Morris and R. Freeman, *J. Am. Chem. Soc.*, 1979, **101**(3), 760–762.
- 16 A. J. Shaka and J. Keeler, *Prog. Nucl. Magn. Reson. Spectrosc.*, 1987, **19**, 47–129.
- 17 P. Allard and T. Härd, *J. Magn. Reson.*, 2001, **153**, 15–21.
- 18 R. R. Ernst, G. Bodenhausen and A. Wokaun, *Principles of Nuclear Magnetic Resonance in One and Two Dimensions*, Clarendon Press, New York, 1987.
- 19 A. D. Bain and B. Berno, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2011, **59**, 223–244.
- 20 P. K. Glasoe and F. A. Long, *J. Phys. Chem.*, 1960, **64**(1), 188–190.
- 21 O. Millet, D. R. Muhandiram, N. R. Skrynnikov and L. E. Kay, *J. Am. Chem. Soc.*, 2002, **124**, 6439–6448.
- 22 Y. Bai, J. S. Milne, L. Mayne and S. W. Englander, *Proteins*, 1993, **17**(1), 75–86.
- 23 D. R. Lide, *CRC Handbook of Chemistry and Physics*, CRC Press, Boca Raton, 85th edn, 1999.
- 24 A. Krezel, *J. Inorg. Biochem.*, 2004, **98**, 161–166.
- 25 K. Laidler and C. King, *J. Phys. Chem.*, 1983, **87**(15), 2657.
- 26 S. Scheiner, *Biochim. Biophys. Acta*, 2000, **1458**, 28–42.
- 27 P. F. Cook, *Isot. Environ. Health Stud.*, 1998, **34**, 3–17.
- 28 F. H. Westheimer, *Chem. Rev.*, 1961, **3**, 265–273.
- 29 G. P. Connelly, Y. Bai, M.-F. Jeng and S. W. Englander, *Proteins*, 1993, **17**, 87–92.
- 30 S. W. Englander, N. W. Downer and H. Teitelbaum, *Annu. Rev. Biochem.*, 1972, **41**, 903–924.
- 31 in *Chemical Kinetics: Proton Transfer*, ed. C. H. Bamford, R. G. Compton and C. F. H. Tipper, Elsevier Scientific Publishing Co., Amsterdam, 1977, vol. 8.
- 32 M. Eigen, *Angew. Chem., Int. Ed.*, 1964, **3**(1), 1–72.

