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## Interaction of the ionic liquid [BMP][TFSA] with rutile TiO<sub>2</sub>(110) and coadsorbed lithium<sup>†</sup>

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Aiming at a fundamental understanding of the processes at the electrode|ionic liquid interface in Li ion batteries, we investigated the interaction of the ionic liquid *n*-butyl-*n*-methylpyrrolidinium bis(trifluoromethylsulfonyl)imide [BMP][TFSA] and of Li with a reduced rutile TiO<sub>2</sub>(110) (1 × 1) surface as well as the interaction between [BMP][TFSA] and Li on the TiO<sub>2</sub>(110) surface under ultrahigh vacuum (UHV) conditions by X-ray photoelectron spectroscopy and scanning tunnelling microscopy. Between 80 K and 340 K [BMP][TFSA] adsorbs molecularly on the surface and at higher temperatures decomposition is observed, resulting in products such as S<sub>ad</sub>, F<sub>ad</sub> and TiN<sub>x</sub>. The decomposition pattern is compared to proposals based on theory. Small amounts of Li intercalate even at 80 K into TiO<sub>2</sub>(110), forming Li<sup>+</sup> and Ti<sup>3+</sup> species. The stoichiometry in the near surface region corresponds to Li<sub>7</sub>Ti<sub>5</sub>O<sub>12</sub>. For higher coverages in the range of several monolayers part of the Li remains on the surface, forming a Li<sub>2</sub>O cover layer. At 300 K, Ti<sup>3+</sup> species become sufficiently mobile to diffuse into the bulk. Li post-deposition on a [BMP][TFSA] covered TiO<sub>2</sub>(110) surface at 80 K results in two competing reactions, Li intercalation and reaction with the IL, resulting in the decomposition of the IL. Upon warming up, the Ti<sup>3+</sup> formed at low *T* is consumed by reaction with the IL adlayer and intermediate decomposition products. Post-deposition of [BMP][TFSA] (300 K) on a surface pre-covered with a Li<sub>2</sub>O/Li<sub>7</sub>Ti<sub>5</sub>O<sub>12</sub> layer results in the partial reaction of [BMP][TFSA] with the Li<sup>+</sup> and Ti<sup>3+</sup> species, which gets completed at higher temperatures.

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## 1 Introduction

Ionic liquids, which are molten salts with a melting point below 100 °C,<sup>1</sup> have attracted high interest in the last few decades because of their uncommon physicochemical properties such as their low vapour pressure, their non-flammability or their high electrochemical stability, and the ability to fine tune these properties by combining different anion–cation pairs.<sup>2–4</sup> Because of the latter properties they also emerged as promising new solvents for use as electrolytes in lithium ion and lithium air batteries.<sup>5,6</sup> In lithium ion batteries the interface between the electrolyte and the electrode and the formation of the so-called solid-electrolyte interphase (SEI), which is a protective layer of decomposition products formed in the first cycle of operation and which prevents further decomposition of the electrolyte, are crucial for the performance and stability of the battery. Despite significant efforts, a molecular scale understanding of structure, composition and formation of these protective layers

is still lacking.<sup>7,8</sup> Aiming at a comprehensive understanding of the processes directly at the interface between the solid electrode and the ionic liquid based electrolyte, we have therefore started to investigate the interactions between solids and ionic liquids on a fundamental scale, using well defined model systems and conditions. In a first step we have investigated the interaction of battery relevant ionic liquids such as *n*-butyl-*n*-methylpyrrolidinium bis(trifluoromethylsulfonyl)imide [BMP][TFSA], 1-ethyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide [EMIM][TFSA] and 1-octyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide [OMIM][TFSA] and Ag(111), Au(111) and Cu(111) model substrates under ultrahigh vacuum (UHV) conditions.<sup>9–13</sup> Employing Surface Science tools such as scanning tunnelling microscopy (STM), angle resolved X-ray photoelectron spectroscopy (XPS) and IR spectroscopy, and combining the experimental efforts with theory we could gain detailed insight into the molecule–molecule and molecule–substrate interactions and the structure formation and thermal stability of the IL adlayers on these simplified model systems. We now extended this study to more realistic materials. In the present contribution we report the results of a study on the interaction of [BMP][TFSA] with a single crystalline rutile TiO<sub>2</sub>(110) substrate, and the effect of Li on the IL–TiO<sub>2</sub>(110) adsorption system. TiO<sub>2</sub> with rutile or anatase structure is in addition to the spinel Li<sub>4</sub>Ti<sub>5</sub>O<sub>12</sub> (LTO) a standard electrode material for Li ion batteries.<sup>14–16</sup>

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Despite extensive literature on the physics and chemistry of  $\text{TiO}_2$  surfaces,<sup>17,18</sup> very little is known on the interaction of Li with these surfaces. Therefore, in addition to investigating the interaction between [BMP][TFSA] and  $\text{TiO}_2(110)$  (Section 2.1), we also explored the interaction between Li and that surface, upon vapor deposition of small amounts of Li on that surface and upon subsequent annealing (Section 2.2). Finally, in order to gain information on the effect of Li on the interaction between a [BMP][TFSA] adlayer and  $\text{TiO}_2(110)$ , we investigated the interaction of [BMP][TFSA] with a  $\text{TiO}_2(110)$  surface modified by pre-deposited Li, and the effect of post-deposited Li on a [BMP][TFSA] adlayer covered  $\text{TiO}_2(110)$  substrate, and the thermal stability of these systems (Section 2.3).

Before presenting the present results, we will briefly summarize the main previous findings relevant for this study. The structure of ILs in the bulk and also at interfaces have been reviewed recently by Hayes *et al.*<sup>19</sup> UHV studies on the interaction of Li with  $\text{TiO}_2(110)$  upon Li deposition reported partly contradictory results. In a combined metastable impact electron spectroscopy (MIES) and ultraviolet photoelectron spectroscopy (UPS) study Krischok *et al.* reported that Li adsorbs ionically up to 0.3 ML and 130 K, accompanied by the formation of  $\text{Ti}^{3+}$  species at the surface. Above 160 K the  $\text{Li}2s$  signal in the MIES spectra was found to vanish, pointing to Li insertion into the substrate.<sup>20</sup> In contrast, employing STM, Tatsumi *et al.* reported the observation of small protrusions on a  $\text{TiO}_2(110)$  ( $1 \times 1$ ) surface upon Li adsorption at room temperature, which they attributed to adsorbed Li.<sup>21</sup> In an electrochemical environment, Li intercalation was found to take place in  $\text{TiO}_2$  substrates (rutile as well as anatase) at room temperature.<sup>22–24</sup> For rutile this intercalation was limited by slow diffusion of the Li into

the bulk crystal.<sup>25,26</sup> Based on these findings, lithium intercalation could therefore be a competing process in addition to the expected reaction with the adsorbed IL layer. Former studies of [OMIM][TFSA] in contact with bulk lithium<sup>27</sup> and of co-deposited Li and a [BMP][TFSA] adlayer on  $\text{Cu}(111)$ <sup>28</sup> revealed a decomposition reaction of the respective IL adlayer on the Li substrate or induced by coadsorbed Li, and the formation of various decomposition products such as  $\text{LiF}$ ,  $\text{F}_{\text{ad}}$ ,  $\text{Li}_2\text{S}$ ,  $\text{S}_{\text{ad}}$ ,  $\text{SO}_{x,\text{ad}}$ ,  $\text{CF}_{3,\text{ad}}$ ,  $\text{Li}_2\text{O}$  and  $\text{Li}_x\text{CH}_y$ .

## 2 Results

### 2.1 [BMP][TFSA] on rutile $\text{TiO}_2(110)$

First we focus on the interaction of [BMP][TFSA] with the rutile  $\text{TiO}_2(110)$  surface. In Fig. 1 the spectra represent the XPS signals in the O1s, C1s, N1s, S2p and F1s regions, recorded at room temperature on a  $\text{TiO}_2(110)$  surface covered with  $\sim 1.5$  monolayers (ML) of [BMP][TFSA], where 1 ML corresponds to a closed layer of adsorbed [BMP][TFSA]. Additionally we show a ball and stick model of a [BMP][TFSA] ion pair. The O1s region reveals two peaks at binding energies (BEs) of 530.6 eV and 532.8 eV. The peak at lower BE is related to the oxide ions of the  $\text{TiO}_2$  substrate. Its position agrees with values reported for a clean and fully oxidized rutile  $\text{TiO}_2$  crystal (530.4 eV).<sup>29</sup> The peak at 532.8 eV as well as the peaks in the other regions shown in Fig. 1 appear only after evaporation of [BMP][TFSA] on the substrate. The peaks resemble what is expected for adsorption of the intact [BMP][TFSA] molecules: the peak at 532.8 eV in the O1s region is attributed to the  $\text{SO}_2$  groups of the  $[\text{TFSA}]^-$  anion. In the C1s region three peaks are identified, similar to previous

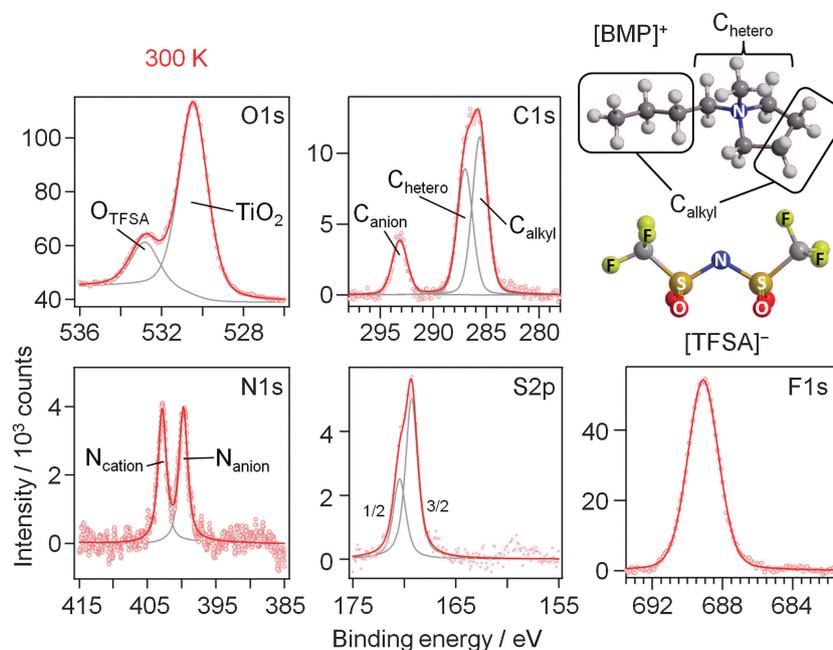


Fig. 1 O1s, C1s, N1s, S2p and F1s XP core level spectra of 1.5 ML [BMP][TFSA] adsorbed on  $\text{TiO}_2(110)$  at 300 K. A structural model of [BMP][TFSA] is displayed on the right hand side.



reports for [BMP][TFSA] multilayer adsorption on Cu(111)<sup>12</sup> and Ag(111):<sup>10</sup> the peaks at  $E_b = 285.6$  eV and 287 eV arise from the [BMP]<sup>+</sup> cations, where the C-atoms directly bonding to the N-atom (“C<sub>hetero</sub>”) have a 1.4 eV higher BE than the alkyl C-atoms (“C<sub>alkyl</sub>”). The peak at 293.2 eV is related to the C-atoms in the [TFSA]<sup>−</sup> anion. As expected for the CF<sub>3</sub> groups with their highly electronegative substituents, this peak is strongly shifted to higher BE compared to the cation related peaks. In the N1s region two peaks with equal intensity appear at 402.8 eV and 399.7 eV. Since the electron density at the positively charged N-atom in the cation is lower than at the N-atom in the anion, the peak with higher BE is attributed to N<sub>cation</sub>, the other one to N<sub>anion</sub>. This assignment resembles earlier ones reported for [BMP][TFSA] on Cu(111)<sup>28</sup> and for [xMIM][TFSA] (xMIM = 1x-3methylimidazolium with x = methyl, ethyl, octyl) on Au(111), Ni(111) and a glass substrate, respectively.<sup>30–32</sup> A doublet peak in the S2p region, at 169.3 eV (S2p<sub>3/2</sub>) and 170.5 eV (S2p<sub>1/2</sub>), and a singlet peak in the F1s region, at 689.1 eV, originate from the [TFSA]<sup>−</sup> anion. After correcting for the different cross sections/sensitivity factors of the different peaks, their intensities largely fit ( $\pm 0.2$ ) to the stoichiometry of the intact [BMP][TFSA] ion pair, with a composition of 4 O<sub>TFSA</sub> : 2 C<sub>anion</sub> : 4 C<sub>hetero</sub> : 5 C<sub>alkyl</sub> : 1 N<sub>cation</sub> : 1 N<sub>anion</sub> : 2 S<sub>TFSA</sub> : 6 F<sub>TFSA</sub>. Only for F<sub>TFSA</sub>, we find a higher intensity (6.6–7) than expected from the stoichiometry, while for O<sub>TFSA</sub> the relative intensity (3.2) is lower than expected. Similar XPS results, including the deviations from a stoichiometric intensity distribution, were reported and discussed by Cremer *et al.* for [OMIM][TFSA] and [MMIM][TFSA] on Au(111)<sup>31</sup> and Ni(111).<sup>32</sup> Following their interpretation these deviations can be used to draw conclusions on the adsorption geometry of the [BMP][TFSA] adsorbates. Elements with signal intensities higher than expected from the stoichiometry are in average positioned slightly higher in the adsorbate layer, those with lower signal intensities closer to the surface with their XPS signal damped by inelastic scattering at the atoms located higher. Accordingly, we assume an adsorption geometry similar to those found for [MMIM][TFSA] on Au(111)<sup>31</sup> and for [BMP][TFSA] on Au(111)<sup>9</sup> and Ag(111)<sup>10</sup> with both the anion and the cation directly adsorbed on the surface. Following our earlier discussion for adsorption on Ag(111),<sup>10</sup> the anion adopts a *cis*-conformation (as shown in Fig. 1), with the O atoms pointing to the surface, and the F atoms to the vacuum. The other atoms are arranged on the main axis of the molecule and are roughly positioned at the same height relative to the substrate surface. In this geometry, the signal of the O atoms is damped relative to the signals of the other atoms, while that of the F atoms is more intense. (In this discussion it should be kept in mind that the XPS data were recorded at a very shallow emission angle of 80° to the surface normal, which makes the measurement very surface sensitive.) Since the intensities of the cation related peaks also fit to the stoichiometry, the cation atoms should be arranged roughly at the same height above the surface as the atoms in the main axis of the anion. For the vertical position of the C<sub>alkyl</sub> atoms, which are situated in the butyl chain and in the ring of the cation, we can only state that in average they are at the same height as the C<sub>hetero</sub> atoms. The individual positions of the different C<sub>alkyl</sub> atoms may be different, if, *e.g.*, a higher

position of the atoms in the alkyl chain is counterbalanced by a lower position of the C<sub>alkyl</sub> atoms in the ring. Note that this means that from the XPS measurements we cannot decide if the alkyl chain is adsorbed flat on the surface or if it is pointing up from the surface.

Additional information on the adsorption behaviour of [BMP][TFSA] on rutile(110) was obtained from STM measurements. STM images of a clean rutile(110) surface, which illustrate also the quality of the surface preparation, are shown in Fig. 2a and b. In Fig. 2a the STM contrast resembles that known from various publications.<sup>17,18,33</sup> There is a consensus that the line structure observed in the STM images is due to a (1 × 1) structure. A detailed model was given in ref. 18. In short, the terminating atoms along the [001] direction are aligned in rows of (a) twofold coordinated “bridge-bonded” O-atoms on top of 6-fold coordinated Ti-atoms, (b) 3-fold coordinated O-atoms and (c) 5-fold coordinated Ti-atoms in the sequence . . . abcbabcb . . . along the [1–10] direction. The bridge-bonded oxygen rows are geometrically highest, but appear as dark rows in STM images, while the (lower lying) 5-fold coordinated Ti-atoms are represented by the rows of bright protrusions. The appearance of the STM images therefore is opposite to the topography of the surface, meaning that it is dominated by electronic effects. The individual bright spots, which are always positioned on top of the dark rows, have been explained as missing bridge-bonded oxygen atoms at the surface (oxygen vacancies).<sup>17</sup>

In addition to the STM images with “normal” contrast (see above), we sometimes also obtained images as shown in Fig. 2b, where the rows of bridge-bonded oxygen atoms appear as bright lines. Images with comparable contrast had also been reported previously;<sup>33,34</sup> by some authors they were termed as STM images with ‘non-reversed contrast’.<sup>34</sup> These images closely resembled images of the TiO<sub>2</sub>(110) surface obtained by non-contact atomic force microscopy (NC-AFM) with atomical resolution.<sup>35</sup> Since the tunnelling parameters are in general not different from those used in the “normal” reversed-contrast images, this different contrast is presumably due to changing tip conditions. Overall, the STM images largely resemble those reported in the literature for (1 × 1) surface structure on a partly reduced sample.<sup>36–40</sup>

Room temperature STM imaging of an [BMP][TFSA] adlayer covered TiO<sub>2</sub>(110) surface, which had been prepared by [BMP][TFSA] evaporation at room temperature, results only in noisy features on top of the (1 × 1) surface structure. We explain this by the formation of a 2D gas/liquid on the surface, similar to previous findings for [BMP][TFSA], [EMIM][TFSA] and [OMIM][TFSA] on Au(111) and Ag(111).<sup>9,10,13</sup> In the latter cases, cooling down to (nearly) liquid nitrogen (LN<sub>2</sub>) temperatures led to the formation of condensed 2D crystalline and 2D glass phases or islands on the surface. For [BMP][TFSA] on TiO<sub>2</sub>(110), even STM imaging at 100 K resulted in very noisy STM images. Since we were reproducibly able to clearly resolve the clean TiO<sub>2</sub>(110) surface, we do not think that this is a technical problem of the STM measurement, but indicates that even at this low temperature the adlayer is rather mobile, forming a mobile 2D gas/liquid phase. In that case, the mobility of [BMP][TFSA] is significantly higher on the TiO<sub>2</sub>(110) surface than on the noble metal surfaces



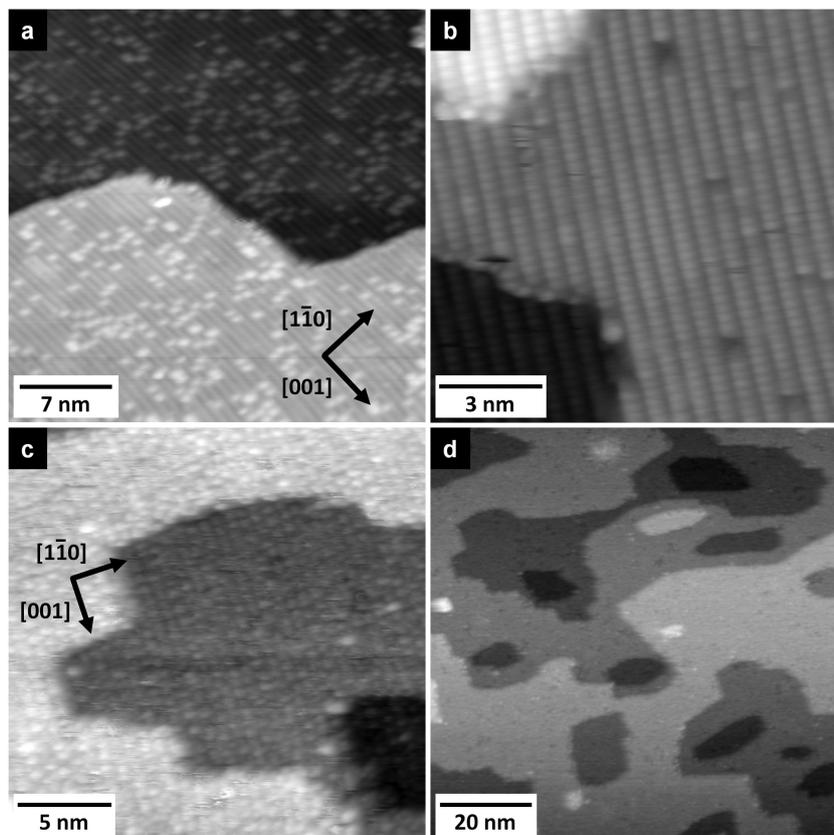


Fig. 2 (a) STM image of a clean  $\text{TiO}_2(110)$  surface. The typical lines of the  $(1 \times 1)$  structure are clearly visible as well as bright spots, indicating oxygen vacancies in the bridge-bonded oxygen rows ( $T = 110$  K,  $U_T = 1.1$  V,  $I_T = 21$  pA). (b) Inverted contrast STM image of  $\text{TiO}_2(110)$  with atomic resolution of the bridge-bonded oxygen atoms and dark spots in these rows resembling missing bridge-bonded oxygen atoms ( $T = 300$  K,  $U_T = 1.7$  V,  $I_T = 650$  pA). Small (c) and large scale (d) STM images of  $\text{TiO}_2(110)$  covered with 1 ML of [BMP][TFSA] (c:  $T = 127$  K,  $U_T = 1.1$  V,  $I_T = 20$  pA; d:  $T = 126$  K,  $U_T = 1.1$  V,  $I_T = 20$  pA).

investigated so far.<sup>9,10</sup> Only at coverages close to 1 ML we could resolve new features, and only at very low tunnelling currents  $< 20$  pA, as shown in Fig. 2c. Most likely, the formation of reasonably stable structures is related to a stabilization of the adsorbates in the closed layers by adsorbate–adsorbate interactions. In the STM image round protrusions can be resolved, which are essentially closely packed along the  $[001]$  oriented rows formed by the  $(1 \times 1)$  surface structure of  $\text{TiO}_2(110)$ . The distance in the  $[1-10]$  direction between the rows was  $0.61 \pm 0.04$  nm. This fits very well to the distance between the bridge-bonded oxygen rows of the  $\text{TiO}_2(110)$  surface of  $0.63 \pm 0.02$  nm (literature:  $0.63 \pm 0.025$  nm<sup>18</sup>). Along the rows the distances seem to vary at around  $0.61 \pm 0.07$  nm. The diameter and density of the round protrusions are measured from several independent STM images to be  $0.56 \pm 0.06$  nm (full width at half maximum) and  $2.6 \pm 0.4$  nm<sup>-2</sup>. The round protrusions observed in these images shall be compared with the structures formed for [BMP][TFSA] on Au(111) and Ag(111),<sup>11</sup> where in addition to round protrusions also pairs of longish protrusions appeared next to each other on the surface at LN<sub>2</sub> temperature. Based also on DFT-D calculations,<sup>10</sup> the round protrusions were identified as [BMP]<sup>+</sup> cations adsorbed on Ag(111), while a pair of longish protrusions reflects a [TFSA]<sup>-</sup> anion. The density of cations was  $0.61$ – $0.79$  nm<sup>-2</sup> (depending on the substrate and

the coverage), and their diameter in the STM images (full width at half maximum) was  $0.70 \pm 0.04$  nm. Our observations shall further be compared with the results of an *in situ* electrochemical study of [BMP][TFSA] adsorption on Au(111), where two different adlayer structures were found: at the negatively charged interface at  $-1.4$  V vs. a Pt quasi reference, only the [BMP]<sup>+</sup> cations adsorb in the first layer in a  $(\sqrt{3} \times \sqrt{13})$  superstructure with a density of  $1.9$  nm<sup>-2</sup> cations, with the alkyl chains lying flat on the surface. For adsorption at  $-1.6$  V they reported a  $(\sqrt{3} \times 2)$  structure with a density of  $3.6$  cations nm<sup>-2</sup>, with the alkyl chains pointing upwards from the surface.<sup>41</sup> Also monolayers of a related IL, [BMP][FAP] ([BMP]<sup>+</sup>-tris(pentafluoroethyl)trifluorophosphate), on Au(111)<sup>42</sup> (same cation, different anion) appeared in STM images as randomly distributed round protrusions with a diameter of  $\sim 0.5$  nm, a mean distance of  $0.45 \pm 0.02$  nm between the protrusions and a density of  $4.5 \pm 0.1$  nm<sup>-2</sup> (imaged under UHV conditions at 210 K). There it was discussed whether the protrusions represent both ion types (1 protrusion = 1 [BMP]<sup>+</sup> or 1 [FAP]<sup>-</sup>) or whether the ions form two layers on top of each other, where STM detects only the upper layer. The possibility that one protrusion represents a complete ion pair was ruled out because the size of the protrusion and their mean distance are too small. The density of protrusions found for [BMP][TFSA] on  $\text{TiO}_2(110)$  is between



that of the adsorbed cations on the charged Au(111) surfaces, where it is higher (138% of that on TiO<sub>2</sub>(110)) at  $-1.6$  V and lower (73%) at  $-1.4$  V. On the other hand, it is much higher (373%) than in the structures of [BMP][TFSA] on Au(111) and Ag(111) at the solid|UHV interface with both the anion and cation directly adsorbed on the surface. It should also be noted that the round protrusions have roughly the same size as the ones representing [BMP]<sup>+</sup> cations adsorbed with upstanding alkyl chains on Au(111) and Ag(111). Considering that based on the XPS results both cations and anions are directly adsorbed on the TiO<sub>2</sub>(110) surface, the present STM results seem to indicate that on TiO<sub>2</sub>(110) one [BMP][TFSA] ion pair is represented by either 3 or 4 round protrusions. In both cases we assume that the anion appears as two protrusions, close to previous findings on Au(111) and Ag(111), where the anion is represented by two protrusions with longish shapes and a lower height compared to the round protrusions representing the cation alkyl chain.<sup>9,10</sup> In the first case the remaining one of the three round protrusions per ion pair would represent an adsorbed [BMP]<sup>+</sup> cation with the alkyl chain pointing upwards. The second case could be rationalized by a model where the cation is imaged as two protrusions, one for the ring and one for the alkyl chain, in this case most likely a flat lying alkyl chain. That would lead to densities of 0.87 or 0.65 ion pairs nm<sup>-2</sup>, respectively, which is in the range of the densities of adsorbed ion pairs with upstanding alkyl chain on the Au(111) and Ag(111) surfaces under UHV conditions. Because of the much higher noise level in these images compared to those obtained for [BMP][TFSA] on Au(111) and Ag(111), a final decision on the adlayer structure on the molecular level is not possible at present.

The 2D growth behaviour of [BMP][TFSA] adsorbates at low temperatures was also examined in large-scale STM images. An example is shown in Fig. 2d. In the first layer [BMP][TFSA] forms a homogeneous layer across the whole surface. Island formation in the submonolayer coverage range, as it was reported on Au(111),<sup>9</sup> Ag(111)<sup>10</sup> and Cu(111),<sup>12</sup> could not be identified in the present work, not even at temperatures of 80 K, which is the lowest temperature accessible for our setup. Hence, attractive molecule-molecule interactions must be weaker on TiO<sub>2</sub>(110) than on the metallic substrates mentioned above, or are even absent.

Finally, we also did not observe any preference for adsorption at the TiO<sub>2</sub>(110) steps (structural effects), whereas this was commonly observed for metallic substrates. In the latter case island growth started preferably from the steps, in the case of Au(111) also at the elbows of the herringbone surface reconstruction.<sup>11</sup>

**Thermal stability of the [BMP][TFSA] adlayer.** The thermal stability of [BMP][TFSA] adlayers was probed by stepwise slow annealing from RT up to 640 K, holding the sample at each temperature for 5 min, respectively, and subsequent XPS characterization, with XPS measurements after each step. One such series of measurements with an adlayer thickness of  $\sim 1.5$  ML is exemplarily shown in Fig. 3 with the topmost spectra (which were already discussed before) at 300 K and the bottom spectra at 640 K. Starting at 300 K, the spectra first slowly lose intensity

in all peaks related to [BMP][TFSA], which is most likely due to multilayer desorption (in the experiment shown in Fig. 3 desorption of molecules in the second adsorption layer). Upon annealing to above 350 K, the decay of the peak intensity is more pronounced. At the same time the peaks related to the [TFSA]<sup>-</sup> anions shift by  $1.0 \pm 0.1$  eV to lower BEs. For the cation related peaks a similar shift appears, but it is only  $0.4 \pm 0.1$  eV. The peak related to the oxygen atoms in the [TFSA]<sup>-</sup> anion, O<sub>TFSA</sub>, shifts from 532.8 eV at 300 K to 531.9 eV at 500 K, which was the highest temperature where a peak could reasonably be fitted into the shoulder of the larger peak generated by bulk TiO<sub>2</sub>. In contrast, the position of the O<sub>TiO<sub>2</sub></sub> peak remains constant, indicating that there are no significant chemical changes of the surface near substrate ions. The intensity of the latter peak is constant until 350 K and rises continuously afterwards up to 145% of its intensity at 640 K. This increase in intensity is due to the loss of adsorbate which reduces the damping of the substrate signals. Nevertheless it only reaches  $41 \pm 1\%$  of the intensity of the adsorbate free surface, which must be due to the presence of decomposition products remaining on the surface (especially F and S containing species, see below). It should be noted that due to the high surface sensitivity of the measurements at 80° detection angle a strong signal damping is obtained already for rather thin layers. In the N1s region the [BMP]<sup>+</sup> related N<sub>cation</sub> peak shifts by 0.3 eV to lower BE before it vanishes at 590 K. The N<sub>anion</sub> related peak shifts by 1.0 eV to lower BE up to a temperature of 440 K, at even higher temperatures the shift amounts to 2.0 eV at 470 K and 2.7 eV at 560 K. This is much more than for the other peaks related to the [TFSA]<sup>-</sup> anion. Therefore we suggest that at 470 K and above this peak arises from an inorganic species which was formed by decomposition of the adsorbed [TFSA]<sup>-</sup> anion. Based on the BE of 397 eV this could be a TiN<sub>x</sub> species, since titanium nitride films show an N1s peak between 397.1 and 397.4 eV.<sup>43</sup> The intensity of this peak has its maximum at 500 K, at higher temperatures it decreases again and vanishes at 590 K (see below). In the S2p region a new peak doublet starts to grow at 500 K. The S2p peaks are at binding energies of 161.3 eV and 162.6 eV. Peaks at exactly the same BEs were also detected upon the decomposition of [BMP][TFSA] on Cu(111) and were attributed to the formation of S<sub>ad</sub> (or Cu<sub>x</sub>S) species.<sup>12</sup> Rodriguez *et al.*<sup>44</sup> reported that adsorption of S<sub>2</sub> on TiO<sub>2</sub>(110) at 300 K result in a peak doublet at 161.6 eV (S2p<sub>3/2</sub>), which they attributed to S adsorbed in (pre-existent) oxygen vacancies in the bridge-bonded oxygen rows. This is only possible until all oxygen vacancies are occupied, and the excess S atoms adsorb on top of the fivefold coordinated surface Ti atoms. These latter S species were proposed to appear as two peak doublets with their S2p<sub>3/2</sub> peaks at 162.8 eV and 163.3 eV. At higher temperatures, the latter doublets vanish and at 600 K only the doublet with the S2p<sub>3/2</sub> peak at 161.6 eV remains.<sup>44</sup> Based on combined XPS and STM experiments between RT and 1073 K, Hebenstreit *et al.* concluded that S<sub>ad</sub> is not only preferentially adsorbed at the oxygen vacancies, but also react with the surface above 390 K by replacing bridge-bonded oxygen with sulphur species. The XPS peak corresponding to these latter species was found



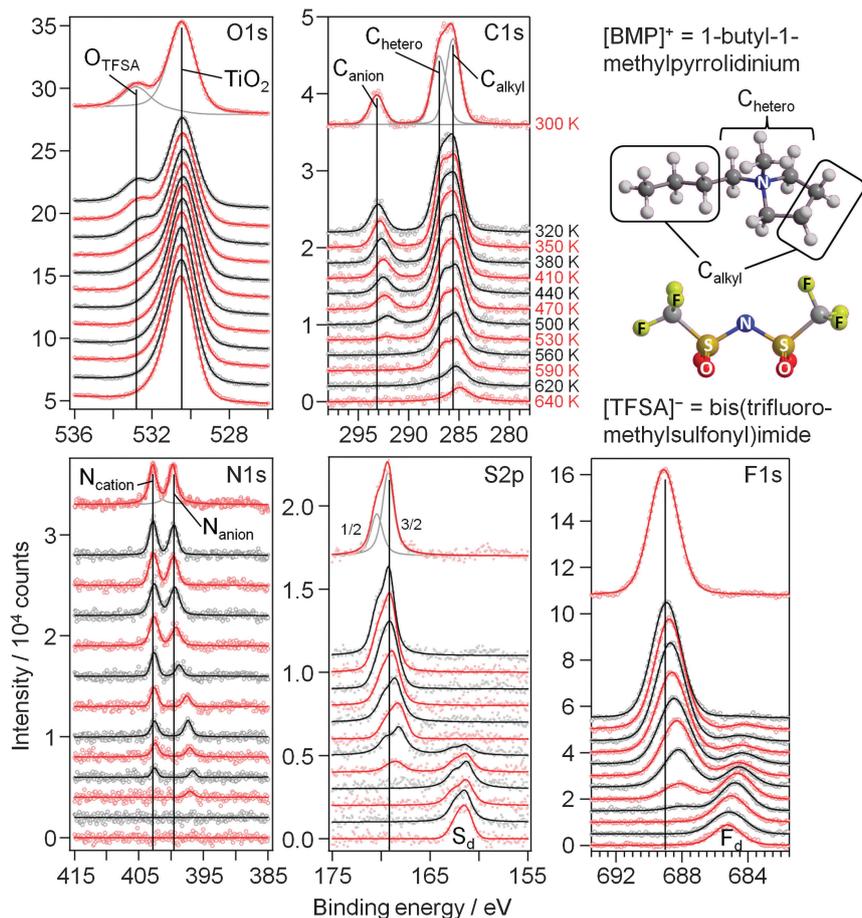


Fig. 3 XPS core level spectra of 1.5 ML [BMP][TFSA] adsorbed on  $\text{TiO}_2(110)$  recorded upon stepwise heating from 300 K to 640 K in steps of 30 K. The structural model of [BMP][TFSA] is displayed on the right hand side.

at 162.4 eV.<sup>45,46</sup> Based on these observations we assume that the  $\text{S}_{\text{ad}}$  atoms generated by [BMP][TFSA] decomposition directly react with the  $\text{TiO}_2(110)$  surface and replace the bridge-bonded oxygen species at temperatures above 500 K, where this peak is formed.

In the F1s region the original  $[\text{TFSA}]^-$  related peak shifts successively by 0.9 eV to lower BE between 300 and 500 K and concomitantly decreases in intensity. At 380 K a new peak starts to grow at 684.6 eV, accompanied by a further intensity decrease of the  $[\text{TFSA}]^-$  related peak. The BE of the new species corresponds to signals measured for adsorbed  $\text{F}^-$  atoms on  $\text{TiO}_2(110)$ .<sup>47–49</sup> We would like to note that  $\text{CF}_3$ -containing molecules are known to decompose *via* generation of  $\text{F}^-$  under reductive conditions.<sup>50</sup> For thermal decomposition of [BMP][TFSA] on Cu(111) no such peak was observed.<sup>28</sup> Between 380 K and 640 K, this peak shifts by 0.6 eV to higher BE. At this point we can only speculate that this shift is either due to a change in the adsorption sites with higher temperature or an incorporation of the  $\text{F}^-$  ions into the subsurface region of the  $\text{TiO}_2(110)$  crystal. The  $[\text{TFSA}]^-$  related peak vanished completely at 560 K.

In the C1s region the  $\text{C}_{\text{anion}}$  related peak follows the trend of the other anion signals, with a shift of 1.0 eV to lower BE and a continuous decay of the peak intensity until it vanished at 560 K.

The signals from the cation carbon shift only by  $\sim 0.3$  eV and persist up to higher temperatures compared to the anion, up to 590 K for the  $\text{C}_{\text{hetero}}$  and up to 640 K for the  $\text{C}_{\text{alkyl}}$  related peak. In addition, the intensities of these signals decrease slower in intensity than all other signals including the  $\text{N}_{\text{cation}}$  signal. Since the BE of  $\sim 285.3$  eV is characteristic also for various kinds of carbonaceous species, we expect that the slower decrease in intensity is related to the formation of carbon containing decomposition products. Since no corresponding signal is found in the N1s region, these compounds do not include any nitrogen any more. A more accurate prediction of the species formed during decomposition is not possible from these XPS measurements (see also the later discussion). Starting at temperatures of  $\sim 410$  K, the intensity of the  $\text{C}_{\text{cation}}$  peaks decays slower than that of the other XP signals during heating, indicating that carbon containing adsorbed decomposition products are formed, while for the other elements the fraction of volatile (desorbing) decomposition products is higher. This discrepancy reaches its maximum at 530/560 K, where the decomposition products account for roughly 80% of this peak-multiplet. At even higher temperature the carbon containing decomposition products start to desorb as well, either directly or *via* further decomposition followed by desorption, and the related signals have essentially vanished at 640 K.



The above data in combination indicate that multilayer desorption occurs between 300 K and 350 K. This is followed by decomposition of both cation and anion species, accompanied by desorption of decomposition products and possibly also of intact cation and anion species. Decomposition starts at  $\sim 380$  K with the formation of adsorbed  $F^-$  and hydrocarbon species. Since the overall peak intensities in the F1s and S2p regions decrease above this temperature, we assume that only part of the  $[TFSA]^-$  anions decomposes, leaving adsorbed decomposition products, while the other part desorbs fully, possibly as intact molecules. At 440 K the amount of decomposition products increases more strongly, and  $S_{ad}$  is formed in addition, which replaces bridge-bonded surface oxygen in the substrate. Finally, nitrides are also formed on the surface, presumably as  $TiN_x$  species. At 530–560 K,  $[BMP][TFSA]$  is completely decomposed. Upon further heating up to 640 K, also the signals of the  $TiN_x$  and of the hydrocarbon decomposition products vanish. In the case of  $TiN_x$  we expect that these species dissolve into the bulk  $TiO_2$  or decompose to  $N_2$  and desorb. The adsorbed carbonaceous species are expected to further decompose and desorb. The  $S_{ad}$  and  $F_{ad}$  species finally remain on the surface.

A similar decomposition pattern of the  $[TFSA]^-$  anion was also found for  $[HMIM][TFSA]$  (1-hexyl-3-methylimidazolium  $[TFSA]^-$ ) on reduced  $CeO_{2-x}$  surfaces.<sup>51</sup>

The decomposition of  $[BMP][TFSA]$  on  $TiO_2(110)$  takes place over a wide temperature range, involving various and changing reaction intermediates/products. This points to a complex multistep mechanism, which is presumably influenced also by the structure of the adsorption site (kink site, different step sites, terrace position). To gain further insight into the decomposition behaviour, further input from experiment and theory

is required, such as high resolution STM data recorded at even lower temperatures or quantum chemical calculations. So far, the decomposition of the  $[BMP]^+$  cation under reductive electrochemical conditions was investigated in DFT-based calculations,<sup>52</sup> and the reductive electrochemical decomposition of the  $[TFSA]^-$  anion in  $[PMP][TFSA]$  ( $PMP = n$ -propyl- $n$ -methyl-pyrrolidinium) was explored by *ab initio* calculations.<sup>53</sup> Furthermore, the thermal decomposition of  $[BMIM][TFSA]$  ( $b =$  butyl) was mapped out *via* DFT based calculations.<sup>54</sup> Finally, the thermal decomposition of bulk  $[BMP][TFSA]$  was characterized by thermal gravimetric analysis, coupled with mass spectrometry detection of the reaction products (TGA-MS).<sup>55</sup> The most likely pathways calculated for the decomposition of  $[TFSA]^-$  are shown in the upper part of Fig. 4a. For the electrochemical decomposition (left side), one of the S–N bonds breaks in the first step, induced by the uptake of one electron at negative electrode potential. The resulting  $SO_2CF_3^-$  and  $NSO_2CF_3^-$  species further react to  $SO_2^{(-)}$ ,  $NSO_2^{(-)}$  and  $CF_3^-$  by breaking the S–C bonds, and they may further decompose to sulfide, oxide, fluoride and carbide anions.<sup>53</sup> For the thermal decomposition (right side), DFT calculations were performed for isolated molecules, where no electron-transfer from the outside is possible. The energetically most favourable decomposition pathway under these conditions proceeds *via* the release of sulfur dioxide.<sup>54</sup> Considering the very different environment in these calculations, one has to be careful when transferring those results to the present experimental findings. Nevertheless, the calculations provide valuable information on the possible decomposition products on the  $TiO_2(110)$  surface. Both decomposition pathways (see Fig. 4a) have in common that they break the same bonds in the molecule, therefore they are most likely also the

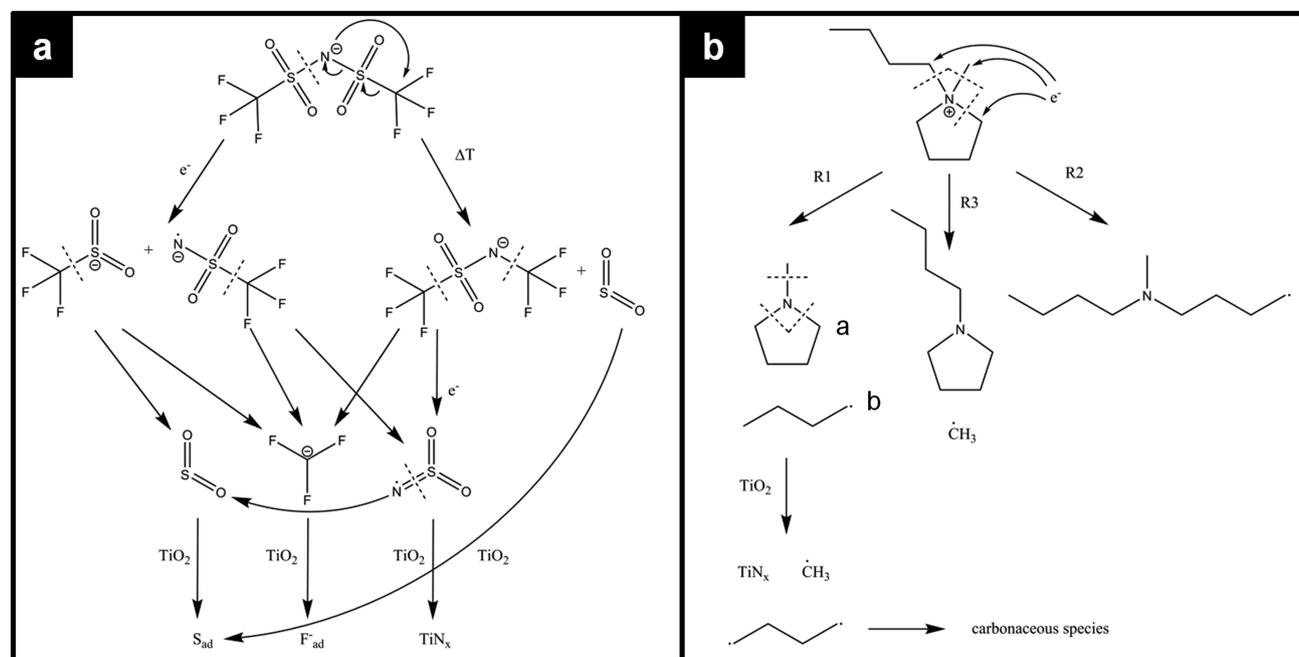


Fig. 4 Possible decomposition mechanisms of (a) the  $[TFSA]^-$  anion and (b) the  $[BMP]^+$  cation according to calculations of ref. 52–54 and thermal decomposition products found on the  $TiO_2(110)$  surface according to our XPS measurements.



ones which will be broken upon decomposition of [BMP][TFSA] on  $\text{TiO}_2(110)$ .

DFT calculations on the decomposition of a  $[\text{BMP}]^+$  cation (under reductive electrochemical conditions) revealed three possible reaction pathways R1–R3, which are shown in Fig. 4b.<sup>52</sup> R1 is the energetically most favoured one, followed by R2 and R3. The TGA-MS measurements (see above) showed that thermal decomposition of [BMP][TFSA] occurs at above 673 K with MS signals at 41, 56, 85 and 69 amu in the gas phase above the solvent, which were attributed to  $\text{C}_3\text{H}_5^+$ ,  $\text{C}_3\text{H}_6^+$ , and  $\text{C}_5\text{H}_{11}^+$  fragments from the cation and  $\text{CF}_3^+$  from the anion. These results agree with the calculations insofar as in both cases  $\text{CF}_3$  species are formed. On the other hand,  $\text{SO}_2$  was proposed as a decomposition product based on the calculations, but was not detected in the TGA-MS measurements. This discrepancy may arise from a too low volatility of these species, such that they are not vaporized and remain in the liquid (possibly as  $\text{SO}_3^-$ ). The  $\text{C}_5\text{H}_{11}^+$  fragment fits to molecule “a” formed by reaction pathway R1 in Fig. 4b. The fragment  $\text{C}_3\text{H}_5^+$  most likely belongs to a butyl-species like molecule “b”, which is also formed in pathway R1. The butyl species is cracked in the ionization zone of the mass spectrometer (among other reactions) into a C3 and a C1 fragment, where the C1 fragment is uncharged and therefore not detected. This decomposition pattern was concluded from comparing the mass spectrum of *n*-butane, where amu 41 and 42 signals are among the most intense signals.<sup>56</sup> Comparing the DFT calculations for electrochemical decomposition with the TGA-MS measurements, it is likely that the decomposition of bulk [BMP][TFSA] proceeds along the reaction pathway R1. Coming back to [BMP][TFSA] on  $\text{TiO}_2(110)$ , thermal decomposition does not stop at the stage described in the three reaction mechanisms. Based on our XPS results, the methyl-pyrrolidinium species is presumably decomposed by breaking the N–C bonds, forming butyl and  $\text{CH}_3$  species. The nitrogen (ad)atoms react with the rutile surface forming  $\text{TiN}_x$ , which is detected in the

XP spectra. If butane is formed, this is likely to desorb at lower temperatures than detected here, since for other surfaces desorption temperatures of  $<200$  K were reported.<sup>57–59</sup> Therefore we assume that the adsorbed reaction products detected by XPS are fragments of the butyl species formed in the decomposition process of  $[\text{BMP}]^+$ . The adsorbed reaction products should solely contain C (and H) atoms based on the XPS results. A more detailed assignment is not possible, since a BE of around 285 eV is typical for a wide variety of  $\text{CH}_x$  containing species, and even graphite shows a peak at around 284.7 eV.<sup>60,61</sup> On the other hand, the  $[\text{TFSA}]^-_{\text{ad}}$  intermediates are further decomposed: for  $\text{SO}_2$  it is known from the literature that it forms sulfite ( $\text{SO}_3^{2-}$ ) and sulfate like species by attachment of  $\text{SO}_2$  to the bridge-bonded oxygen species of the  $\text{TiO}_2(110)$  surface in a temperature range between 135 K and 450 K.<sup>62,63</sup> For  $\text{TiO}_2(110)$  with a high density of surface defects (e.g., after sputtering) or with a high step density (e.g.  $\text{TiO}_2(441)$ ),  $\text{SO}_2$  was reported to react to  $\text{S}_{\text{ad}}$  even at room temperature.<sup>63–65</sup> Our observation of  $\text{S}_{\text{ad}}$  formation at temperatures  $>500$  K, but no formation of  $\text{SO}_x$ , can be explained the following way. The  $\text{SO}_2$  generated by the decomposition of  $[\text{TFSA}]^-$  forms adsorbed sulfite and sulfate species. At the high temperatures these may instantaneously react to  $\text{S}_{\text{ad}}$ , or the decomposition of the  $[\text{TFSA}]^-$  creates a high density of defects on the  $\text{TiO}_2(110)$  surface, which catalyse the reaction  $\text{SO}_4^{2-} \rightarrow \text{S}_{\text{ad}}$ . If adsorbed  $\text{CF}_3$  and  $\text{NSO}_2$  species are formed in the decomposition process of  $[\text{TFSA}]^-$ , these are also decomposed forming  $\text{F}_{\text{ad}}$ ,  $\text{TiN}_x$  and  $\text{SO}_2$ , where the latter again reacts to  $\text{S}_{\text{ad}}$ , which are detected in the XPS measurements. The remaining carbon may finally further react to CO or  $\text{CO}_2$ , which both desorb instantaneously.<sup>18</sup>

## 2.2 Interaction of Li with $\text{TiO}_2(110)$

First we present results of XPS measurements on the interaction of Li with  $\text{TiO}_2(110)$ . Detailed spectra of the Li1s, Ti2p and O1s regions are shown in Fig. 5a. The spectra were recorded after

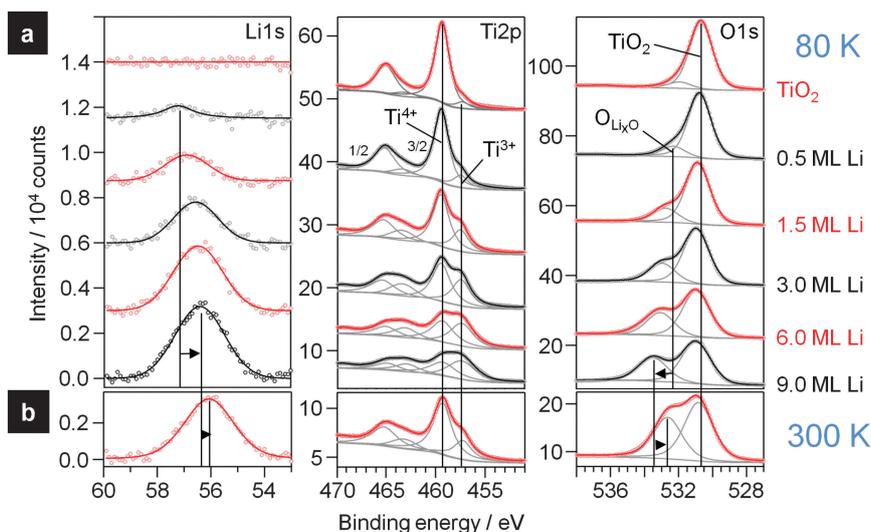


Fig. 5 (a) XP detail spectra of the Li1s, Ti2p and O1s region of clean  $\text{TiO}_2$  (topmost spectra) and upon increasing evaporation of Li onto this surface at 80 K. (b) XP spectra after slowly heating up the sample overnight to 300 K.



stepwise deposition of Li on a clean TiO<sub>2</sub>(110) surface (topmost spectra) at 80 K. The bottom spectra show the highest Li dose (corresponding coverage  $\sim 9.0$  ML Li). In the Li1s region a peak emerges at 57.2 eV after the first Li dose (0.5 ML). Upon increasing the amount of deposited Li, this peak grows in intensity and shifts lower BEs, to 56.4 eV at 9.0 ML. Generally, the intensity of the Li1s peak is very low due to the very small atomic sensitivity factor of this peak. Interestingly, the peak positions are at significantly higher BE than for bulk Li (54.8 eV<sup>66</sup>) or for Li adsorbed on Cu(111) (55.1 eV<sup>28</sup>) and even higher than reported for intercalated Li in anatase<sup>22</sup> or LTO.<sup>67</sup> Therefore we assume that Li has reacted with the surface and exists as a Li<sup>+</sup> species, either as a Li<sub>x</sub>O species at the surface or intercalated into the substrate, rather than being atomically adsorbed on the surface.

This assumption is supported by the corresponding Ti2p and O1s spectra, where new peaks are formed upon deposition of Li. After the first Li dose (0.5 ML), an additional new doublet evolves in the Ti2p region at 2.0 eV lower BE than the doublet of the initial Ti2p state of TiO<sub>2</sub>(110). In the O1s region also a new peak ("O<sub>Li<sub>x</sub>O</sub>") appears, which is shifted by 1.5 eV to higher BE relative to the TiO<sub>2</sub>(110) peak. The new O<sub>Li<sub>x</sub>O</sub> peak shifts to higher BE with increasing Li deposition, by 1.1 eV at 9.0 ML. This is very similar to the results obtained for Li intercalation in anatase at RT, where electrochemical insertion of Li also caused an additional Ti2p peak at lower BE relative to the original TiO<sub>2</sub> peak.<sup>22,24</sup> Note that in our measurement these additional peaks appear already with weak intensities in the spectra of the clean rutile sample. We explain this by small amounts of Ti<sup>3+</sup> species in the crystal due to the reductive pre-treatment (Ti2p state) and to residual contaminations from previous experiments with Li on the same crystal. Nevertheless, the evolution of these peaks clearly demonstrates that the atomic Li reacts with the surface upon interaction with Li even at 80 K. The additional doublet in the Ti2p region is attributed to the formation of Ti<sup>3+</sup> species, which are formed by the transfer of one electron from Li<sup>0</sup> to a Ti<sup>4+</sup> species in the rutile crystal. Since there is also a second peak in the O1s region evolving at lower BE relative to the original TiO<sub>2</sub> peak, we expect that the resulting Li<sup>+</sup> atoms are positioned next to O<sup>2-</sup>-atoms in the crystal which therefore have a different electronic surrounding. Another important question is whether this reaction takes place solely on the surface or if intercalation into the near surface region or bulk rutile crystal occurs as well. For this purpose we concentrate on the changes of the peak intensities. In the spectra recorded on pure TiO<sub>2</sub>(110) the Ti2p:O1s peak intensity ratio (corrected by the sensitivity factors) is  $0.54 \pm 0.02$ , which is slightly higher than the expected ratio of 1:2. We tentatively explain this by the presence of oxygen vacancies and Ti<sup>3+</sup> species. Upon addition of Li this ratio is decreased to 0.32, because the intensity of both O1s peaks together increases to 123% while the total intensities of the Ti2p peaks fades to 69% compared to the pure rutile signals (see below). The intensity ratio between the Ti<sup>4+</sup> and Ti<sup>3+</sup> peaks is very close to 2:3. The high fraction of reduced Ti<sup>3+</sup> species implies that these are not only in the surface layer, but also in deeper layers, meaning that diffusion of Li into the subsurface regions is possible under these conditions and leads to the

formation of Ti<sup>3+</sup> and Li<sup>+</sup> in these regions. The ratio of 2:3 would fit to the stoichiometry of a Li loaded LTO spinel Li<sub>7</sub>Ti<sub>5</sub>O<sub>12</sub>. In LTO the ratio between Ti and O atoms is smaller than in rutile TiO<sub>2</sub>, therefore part of the lower relative intensity of the Ti2p peaks, relative to that of the O1s signals, can be explained by the change in stoichiometry from Ti:O of 1:2 (or 0.5) in rutile to 5:12 (or 0.417) in LTO. The measured Ti:O ratio of 0.32, however, is even lower, in agreement with the existence of a Li oxide (Li<sub>x</sub>O) cover layer on top of the LTO-like layer. In that case the peak at higher BE in the O1s region must originate from two different species: the Li<sub>x</sub>O formed on top of the substrate and oxygen species in the subsurface region which are neighboured to intercalated Li<sup>+</sup>. Based on the higher Li1s BE (bulk Li<sub>2</sub>O: Li1s: 53.8 eV; O1s: 528.5 eV<sup>68,69</sup>), the Li<sub>x</sub>O layer must be electronically different from bulk Li<sub>2</sub>O.

The calculation of the thickness of the Li<sub>x</sub>O overlayer is not simple since the decay of the Ti2p peak intensity originates from both the Li<sub>x</sub>O cover layer and the change in stoichiometry of the TiO<sub>2</sub>(110) substrate in the near surface region. The same combination of effects is responsible also for the higher intensity of the O1s signal (123%) compared to the XPS signal of the Li free sample, and for the evolution of the Li1s peak, which originates both from the Li in the Li<sub>x</sub>O surface layer and from the Li in the LTO like layer. For an estimate of the Li<sub>x</sub>O layer thickness we assume that the Ti2p signals are generated by a LTO layer and calculate the O1s signal which corresponds to this Ti2p signal to achieve the stoichiometry between Ti and O in LTO. Subtraction of this value from the more intense measured O1s signals results in the oxygen fraction which originates from the Li<sub>x</sub>O layer. Details of the calculation can be found in the ESI.† For deposition of 9 ML Li this calculation results in a thickness of the Li<sub>x</sub>O cover layer of *ca.* 0.1–0.2 nm, which is  $\leq 1$  ML of Li<sub>x</sub>O.

The thickness can also be calculated from the damping of the O1s<sub>LTO</sub> or of the Ti2p signals,<sup>70</sup> where the latter is again assumed to fully result from a stoichiometric LTO layer (details on the calculation are found in the ESI†). Both estimates give a result for the thickness of the Li<sub>x</sub>O layer which is in the same range as it was found in the first calculation. The assumption that a fully stoichiometric LTO substrate is used for the calculation appears plausible for the sample with 9 ML Li deposited and an detection angle of 80°, since in this case the layer with LTO like stoichiometry is expected to be clearly thicker than the detection depth (details on the calculation of the LTO layer thickness are given in the ESI†).

The small thickness of the Li<sub>x</sub>O layer is presumably the reason why the peaks are shifted to higher BE compared to bulk Li<sub>2</sub>O because of the interaction to the rutile crystal underneath. Furthermore, it may also be that for these low coverages no homogeneous, closed layer of Li<sub>x</sub>O is formed but instead oxygen and Li<sup>+</sup> are enriched on/in the first surface layer. The mobility of the Ti<sup>m+</sup> ions in rutile seems to be high enough that diffusion is possible from the bulk TiO<sub>2</sub>(110) in the near surface region and to the surface even at 80 K. The Li:O stoichiometry in this Li<sub>x</sub>O film cannot be quantified precisely because of the low intensity of the Li1s peak, leading to problems in the



background subtraction, and since the Li1s peaks of the Li species at the surface and those intercalated in the near surface region appear at the same BE.

To elucidate how thick the LTO layer between the  $\text{Li}_x\text{O}$  layer and the bulk rutile is we also conducted the same XPS measurements as shown in Fig. 5a with a detection angle of  $0^\circ$  with respect to the surface normal (the spectra are shown in the ESI†). Here the mean escape depth is with *ca.* 2.9 nm higher than for  $80^\circ$  with 0.5 nm. For these signals the  $\text{Ti}^{4+}$  peak nearly has double the intensity compared to the  $\text{Ti}^{3+}$  peak at the highest investigated Li loading. Therefore we assume that now a mixture of the  $\text{Li}_x\text{O}$  layer, followed by  $\text{Li}_7\text{Ti}_5\text{O}_{12}$  and presumably a layer with a decreasing Li concentration on top of unchanged bulk rutile is probed. Since we cannot derive information on the type of the concentration gradient of  $\text{Li}^+$  and  $\text{Ti}^{3+}$  from our XPS measurements, because we are always sampling the average of the uppermost atomic layers (and a deep profiling with sputtering is also not possible since this would alter the surface), we cannot determine the thickness of the LTO-like layer except for the fact that for 9 ML Li deposition it is thicker than the sampling depth of the measurement at  $80^\circ$ . Nevertheless, for a rough estimate we calculated the thickness assuming a layer geometry of  $\text{Li}_x\text{O}$  on top of LTO on top of  $\text{TiO}_2(110)$  with no concentration gradients in between. Under these conditions the LTO layer thickness is calculated to be  $2.0 \pm 0.1$  nm. Upon slow warm up of the sample to 300 K overnight the stoichiometry changes. Spectra recorded at the  $80^\circ$  emission angle after annealing are shown in Fig. 5b. The overall damping of the Ti2p peaks stays the same, indicating that the  $\text{Li}_x\text{O}$  layer is still present, but the intensity ratio of the  $\text{Ti}^{4+}:\text{Ti}^{3+}$  peak changed to  $\sim 5:1$  (2.7), indicative of a distinct migration of  $\text{Ti}^{3+}$  species deeper into the bulk. For measurements conducted at the  $0^\circ$  emission angle the relative intensities changed marginally from  $(\text{Ti}^{4+} + \text{Ti}^{3+}):\text{O} = 0.42$ ,  $\text{Ti}^{4+}:\text{O} = 0.33$  and  $\text{Ti}^{3+}:\text{O} = 0.08$  at 80 K to  $(\text{Ti}^{4+} + \text{Ti}^{3+}):\text{O} = 0.43$ ,  $\text{Ti}^{4+}:\text{O} = 0.36$  and  $\text{Ti}^{3+}:\text{O} = 0.06$  at RT. This means that in the uppermost 1–2 nm, which are probed at the  $80^\circ$  emission angle, the composition changes strongly, while the average stoichiometry in the uppermost  $\sim 5$  nm stays about the same. The higher temperature enhances the diffusion into the bulk and the  $\text{Li}^+$  and  $\text{Ti}^{3+}$  species agglomerated in the first 1–2 nm at 80 K move deeper into the material at higher temperature. The  $\text{Li}_7\text{Ti}_5\text{O}_{12}$  layer in the surface region therefore was kinetically stabilized by the low temperatures and upon heating the gradient of  $\text{Li}^+$  and  $\text{Ti}^{3+}$  interstitials is reduced. Nevertheless, since the signal at  $0^\circ$  detection angle changes very little, the Li ions seem to remain in the layer probed here, in the topmost 5 nm.

To summarize our XPS results on Li deposition on  $\text{TiO}_2(110)$ , we found that even at 80 K part of the deposited Li intercalates, while the remaining part forms a thin adlayer of  $\text{Li}_x\text{O}$  like species. Because of diffusion limitations, the Li species stay in the near surface region and form a layer with a LTO-like stoichiometry, as evidenced by the concentration of  $\text{Ti}^{3+}$  species. At higher temperature (300 K), diffusion becomes faster and the  $\text{Li}^+$  and  $\text{Ti}^{3+}$  species are probed not mainly in the first 2 nm, resulting in a lower signal intensity of these species at the  $80^\circ$  detection angle in XPS. Since the signal intensities did not change significantly

for detection at  $0^\circ$ , the  $\text{Li}^+$  and  $\text{Ti}^{3+}$  species are still in the uppermost *ca.* 5 nm which are probed in this geometry.

Comparing the present XPS results on Li interaction with rutile  $\text{TiO}_2(110)$  with previous related findings for Li electro-deposition, but also for Li deposition in UHV on  $\text{TiO}_2$ , both for anatase and for rutile, leads to the following picture: investigating the insertion of Li into thin films of anatase with XPS (synchrotron radiation as the X-ray source) at room temperature, Henningsson *et al.*<sup>22,23</sup> found a new peak in the Ti2p region growing simultaneously with the successive deposition of Li, which appeared at 1.6 eV lower BE compared to the original Ti2p peak related to  $\text{Ti}^{4+}$  species in the substrate. For high amounts of Li ( $> 5$  ML) a third peak emerged at even lower BE. These two peaks were attributed to the formation of  $\text{Ti}^{3+}$  and  $\text{Ti}^{2+}$  species in the substrate upon Li insertion into the material. In the Li1s region they observed peaks at 56.5 eV and 55.5 eV with the peak at higher BE forming at coverages up to 2.8 ML. The peak at lower BE appears at coverages  $> 0.64$  ML and dominates the spectrum above 2.8 ML. The authors explained these results by the formation of two phases in anatase, with a Li rich  $\text{Li}_{0.4}\text{TiO}_2$  phase and a Li poor  $\text{Li}_{<0.07}\text{TiO}_2$  phase. This is in agreement with findings for the electrochemical insertion of Li into anatase.<sup>26</sup> Södergren *et al.*<sup>24</sup> combined electrochemical Li insertion with XPS measurements and found that the formation of the Ti2p peak related to  $\text{Ti}^{3+}$  is reversible when cycling the sample potential, which shows that the original anatase structure is recovered upon deintercalation of Li. Comparison with the studies mentioned before indicates that electrochemical Li insertion and Li insertion upon evaporating Li in UHV on anatase proceed in a comparable process. Li deposition on anatase  $\text{TiO}_2$  allows rapid Li intercalation at room temperature, also under the conditions applied in the present study (Li evaporation in UHV).

On rutile  $\text{TiO}_2$ , intercalation of Li is more difficult. For bulk materials it was thought for a long time that at RT Li intercalation is not possible at all. Zachau-Christiansen *et al.* reported that the electrochemical Li insertion into rutile was inhibited at  $25^\circ\text{C}$ , only for reaction at  $120^\circ\text{C}$  two Li containing phases with stoichiometries of  $\text{Li}_{0.13-0.15}\text{TiO}_2$  and  $\text{Li}_{0.5}\text{TiO}_2$  were observed,<sup>26</sup> comparable to observations for anatase (see above). More recent studies, however, indicated that intercalation is feasible also at RT, *e.g.*, by reducing the size of the rutile particles, *e.g.*, to needles with diameters of 5–20 nm.<sup>25</sup> This was explained by a slow diffusion of Li at  $25^\circ\text{C}$ , especially perpendicular to the *c* channels in the crystal. (The *c* channels run along the [001] direction in the (110) lattice plane, therefore intercalation into a  $\text{TiO}_2(110)$  terminated crystal has to proceed along the slow diffusion direction.)

These results are in good agreement with our findings when keeping in mind that the earlier studies on rutile mentioned above were conducted determining the bulk composition of the  $\text{TiO}_2$  materials rather than surface sensitive techniques. Using surface sensitive measurements we could demonstrate that Li intercalation starts already at much lower temperatures. Under those conditions it is limited, however, to the near surface regions because of the slow diffusion of the Li species.



Our results agree also with the findings by Krischok *et al.*,<sup>20</sup> who used MIES (metastable impact electron spectroscopy) and UPS (HeI) to study exactly the Li interaction with rutile TiO<sub>2</sub>(110) at temperatures between 130 K and 620 K, where in particular MIES is even more surface sensitive than XPS (although it does not allow to quantify how much and how deep the Li is inserted). They reported that at 130 K Li adsorbs ionically up to 0.3 ML and Ti<sup>3+</sup> is formed in parallel. At higher coverages the ionic character of the adsorbed Li atoms is smaller (*e.g.* at 0.7 ML Li<sup>0.4+</sup>). Above 160 K, the Li2s signal in the very surface sensitive MIES spectra disappeared almost completely, which they explained by (almost complete) Li insertion into the bulk. This agrees at least qualitatively with our findings, when using comparable Li coverages of 1.4 ML.

In addition to XPS we performed STM measurements of the surface after evaporating Li at 80 K. Representative STM images are shown in Fig. 6. The first image, recorded after 1.5 ML Li deposition, provides an overview of the surface morphology after deposition, while the smaller scale image in Fig. 6b resolves also details of the surface structure. The steps and terraces of the substrate seem to be unchanged after Li deposition. On the surface we find a number of protrusions which are (nearly) uniformly spread across the surface. At some places also the original TiO<sub>2</sub>(110) surface with its 1 × 1 structure is visible in Fig. 6b (two regions are marked by red circles). The round protrusions may correspond to Li<sub>x</sub>O species, which according to our XPS data are formed upon Li deposition on top of the substrate. Another possibility would be that they result from the change of the substrate itself, reflecting the formation of a LTO like species. The STM images do not indicate any change in appearance between 80 K and RT. Comparing our results with the STM data reported by Tatsumi *et al.* for room temperature adsorption of Li on TiO<sub>2</sub>(110), our STM images correspond to those shown in their work for higher deposition times (600 s in Fig. 2b in ref. 21). Note that the latter authors assigned the new structures observed in their STM images to Li adatoms, which seems to be in contrast to our findings and to those of other authors.<sup>20</sup> We rather expect that these structures reflect the initial stage of Li insertion and formation of Li<sub>x</sub>O. Since the

protrusions detected are mostly situated on the rows of 5-fold coordinated Ti<sup>4+</sup> species,<sup>21</sup> one may speculate that they are generated by Ti<sup>3+</sup> species formed by the reaction with an impinging Li atom or by an adsorption complex of Li<sup>+</sup> on Ti<sup>3+</sup>.

Furthermore we may compare our STM images with AFM images of single crystalline LTO(111) surfaces reported by Kitta *et al.*,<sup>71</sup> who performed AFM measurements on surfaces which were electrochemically charged by Li insertion. They observed a surface roughening, and their AFM images at the initial stage of the Li insertion are very similar to our STM images. This seems to indicate that the initial Li insertion in TiO<sub>2</sub>(110) occurring upon Li deposition under UHV conditions closely resembles the initial stage of electrochemical Li insertion into LTO.

### 2.3 [BMP][TFSA] + Li on rutile TiO<sub>2</sub>(110)

This section focuses on the interaction of coadsorbed [BMP][TFSA] and Li with each other and with the TiO<sub>2</sub>(110) substrate. First we discuss the effect of Li deposition on a TiO<sub>2</sub>(110) surface which was pre-covered with 0.7 ML of [BMP][TFSA]. XP core level spectra recorded after Li deposition at 80 K are shown in Fig. 7. From top to bottom spectra recorded on the clean TiO<sub>2</sub>(110) substrate (a), after deposition (80 K) of 0.7 ML of [BMP][TFSA] (b), and after stepwise deposition of Li (c–h) are presented. In the Ti2p region deposition of [BMP][TFSA] causes a damping of the bulk TiO<sub>2</sub> (Ti<sup>4+</sup> and Ti<sup>3+</sup>) peaks. With increasing Li uptake (c–h) the intensity of the Ti<sup>4+</sup> decreases further, and that of the Ti<sup>3+</sup> signals increases, comparable to the findings for Li deposition on pure TiO<sub>2</sub>(110), without [BMP][TFSA]. Obviously, also under these conditions Li reacts with the TiO<sub>2</sub>(110) substrate and intercalates into the substrate. This is confirmed by an additional peak in the O1s spectral range (O<sub>Li<sub>x</sub>O</sub>), which appears at a BE of 532.3 eV between the two peaks O<sub>TFSA</sub> and O<sub>Ti<sub>2</sub>O</sub> and shifts gradually to 532.7 eV with increasing Li deposition. This peak is at the same BE as the additional peak for Li deposition/intercalation on/in TiO<sub>2</sub>(110) in the absence of a [BMP][TFSA] adlayer, and therefore attributed to the same species (O species with changed electrochemical surrounding due to the neighbouring Li<sup>+</sup> species). The intensity of the O1s rutile peak stays (nearly) constant during Li deposition, while that of the O<sub>TFSA</sub> peak strongly fades to

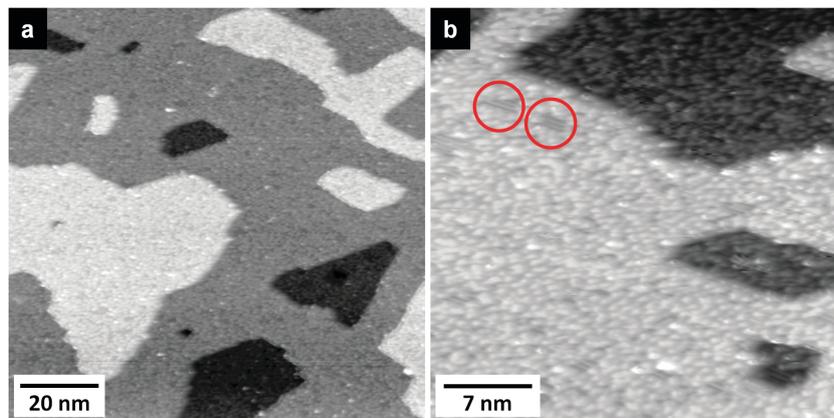


Fig. 6 Large (a) and small scale (b) STM images of a TiO<sub>2</sub>(110) surface after Li deposition at 80 K. At some places (*e.g.*, marked by red circles) parts of the TiO<sub>2</sub>(110) structure are still visible (a:  $T = 108$  K,  $U_T = 1.1$  V,  $I_T = 21$  pA; b:  $T = 102$  K,  $U_T = 1.1$  V,  $I_T = 21$  pA).



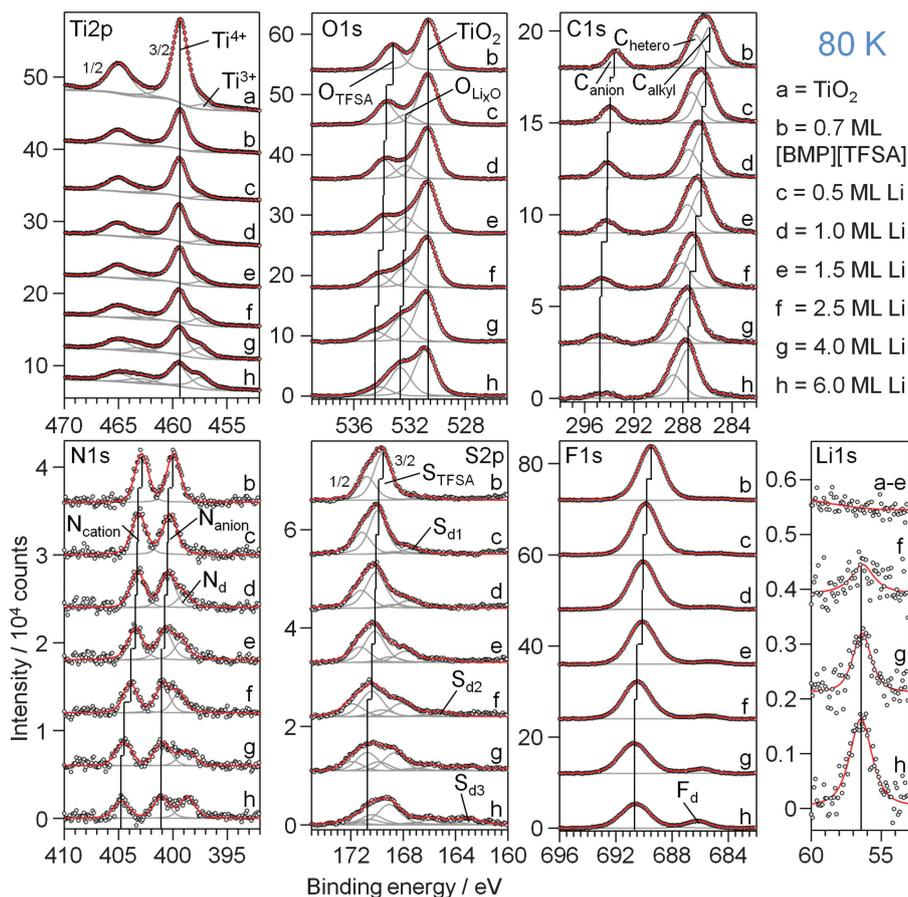


Fig. 7 XPS detail spectra of (a) a clean  $\text{TiO}_2(110)$  surface, (b) covered with 0.7 ML [BMP][TFSA] and (c–h) after increasing amounts of Li post-deposited step by step onto [BMP][TFSA]/rutile(110); the sample was always at 80 K.

37% of its starting value and shifts by 1.1 eV to higher BE. The intensity decay of the  $[\text{TFSA}]^-$  related peaks implies that the deposited Li atoms not only react with and intercalate into the  $\text{TiO}_2$ , but also partly react with the adsorbed [BMP][TFSA], even at low temperature. Therefore the  $\text{O}_{\text{Li}_x\text{O}}$  peak could at least partly result from decomposition products of the IL anion. Possible reaction products would be  $\text{SO}_x$  or  $\text{SO}_x$  containing species, since according to the literature  $\text{SO}_3^{2-}$  adsorbed on  $\text{TiO}_2(110)$  has roughly the same BE, between *ca.* 532 and 533 eV (for related S peaks see below).<sup>64</sup> When comparing, *e.g.*, the peak intensities of  $\text{Ti}2\text{p}_{\text{Ti}^{4+}}$ ,  $\text{Ti}2\text{p}_{\text{Ti}^{3+}}$  and  $\text{O}1\text{s}_{\text{Li}_x\text{O}}$  signals in this measurement, using the highest Li deposition of 6 ML, with those obtained for a similar Li uptake on a pure  $\text{TiO}_2(110)$  substrate, the intensity ratios of  $\text{Ti}^{4+} : \text{O}1\text{s}_{\text{Li}_x\text{O}}$  (0.27) and  $\text{Ti}^{3+} : \text{O}1\text{s}_{\text{Li}_x\text{O}}$  (0.16) are much lower for  $\text{TiO}_2(110) + \text{IL} + \text{Li}$  than for  $\text{TiO}_2(110) + \text{Li}$  (0.53, 0.5). Since the intercalation into the  $\text{TiO}_2(110)$  surface region should result in (roughly) the same stoichiometry for the same Li dose/same temperature, the  $\text{O}1\text{s}_{\text{Li}_x\text{O}}$  peak is too intense relative to the  $\text{Ti}2\text{p}$  peaks to only originate from Li intercalation. Therefore we expect this peak to include also contributions from [BMP][TFSA] decomposition products of the  $\text{SO}_x$  type. For the highest amount of Li deposition (6 ML) the two species should contribute about equally. The formation of a  $\text{SO}_x$  species is supported also by a new peak in the  $\text{S}2\text{p}$  regime. Here an

additional peak doublet (marked as  $\text{S}_{\text{d}1}$  in Fig. 7) increases with Li deposition at BEs of 167.5 ( $\text{S}2\text{p}_{3/2}$ ) and 168.7 eV ( $\text{S}2\text{p}_{1/2}$ ), which shift to 169.1 and 170.4 eV, respectively, with increasing Li deposition. This fits well to  $\text{SO}_3^{2-}$  or  $\text{SO}_4^{2-}$  species adsorbed on  $\text{TiO}_2(110)$ <sup>64,72–74</sup> and  $\text{Li}_2\text{SO}_3$ .<sup>75</sup> Nevertheless, considering also the peaks in the other regions (see also discussion below), this peak also must include contributions from  $\text{F}_3\text{C}-\text{O}_2\text{S}-\text{N}^- \text{Li}^+$  species. In addition to the  $\text{S}_{\text{d}1}$  peak, two new peak doublets develop with higher Li amounts at lower BEs of 166.0 and 162.8 eV (position of the  $\text{S}2\text{p}_{3/2}$  peak of each doublet). The latter one is at 1.5 eV higher BE than the peak obtained upon thermal decomposition of a pure [BMP][TFSA] adlayer on  $\text{TiO}_2(110)$ , which was attributed to the formation of  $\text{S}_{\text{ad}}$ . Accordingly, this may be of similar origin in the present case, but with a slightly higher BE due to a different chemical surrounding. It is also possible, however, that a  $\text{Li}_x\text{S}$  species is formed, whose  $\text{S}2\text{p}_{3/2}$  signal is expected to appear in the same range of BEs (162.0 eV).<sup>28</sup> The origin of the other  $\text{S}2\text{p}_{3/2}$  peak at 166.0 eV is not fully clear. Its BE is too low for a  $\text{SO}_3^{2-}$  or  $\text{SO}_4^{2-}$  species,<sup>64,73,74</sup> but it could be due to a  $\text{F}_3\text{C}-\text{O}_2\text{S}^- \text{Li}^+$  like species following the reaction scheme for [BMP][TFSA] decomposition in Fig. 4a. The original  $\text{S}_{\text{TFSA}}$  peak has lost almost 85% of its initial intensity after the highest Li deposition. This is much more than the loss of the  $\text{O}_{\text{TFSA}}$  peak, which was about 63%. Therefore also the  $\text{O}_{\text{TFSA}}$



peak must contain contributions from decomposition products with similar BE. Possible candidates would be  $F_3C-O_2S-N^-Li^+$  and  $F_3C-O_2S^-Li^+$ . The same is true for the N1s region, where the two peaks originating from intact [BMP][TFSA],  $N_{cation}$  and  $N_{anion}$ , lose about 2/3 of their initial intensity during Li evaporation and shift by 1.9 and 1 eV to higher BE, respectively. Due to the large shift of the  $N_{cation}$  peak this peak is expected to also contain contributions from  $C_xH_yN_z$  species, which appear at similar BE's. Finally, another N1s peak appears at BE 398.7 eV, 2 eV higher than the additional peak observed upon thermal decomposition of [BMP][TFSA] on  $TiO_2(110)$ , which we had tentatively attributed to  $TiN_x$ . Therefore the new signal at 398.7 eV must be due to another reaction product. Signals at this BE were attributed to  $Li_3N$  species in a postmortem XPS analysis of Li-ion batteries,<sup>75</sup> or to reaction products arising from the interaction of [OMIM][TFSA] with Li on a polycrystalline Cu foil<sup>27</sup> and [BMP][TFSA] with Li on Cu(111).<sup>28</sup>

In the F1s region the original  $F_{TFSA}$  peak at 689.5 eV decreases to 50% of its initial intensity and shifts by 1.1 eV to higher BE during Li deposition, while at the same time a new peak grows in at 685 eV, which shifts by 1.2 eV to higher BE with increasing Li coverage. This peak is attributed to adsorbed  $F^-$  species, similar to the assignment made for thermal decomposition of [BMP][TFSA] on  $TiO_2(110)$ , and to the formation of  $LiF$ .<sup>75</sup>

Finally, in the C1s region, no new peaks are formed in addition to the initial [BMP][TFSA] peaks, but the  $C_{anion}$  peak decreases to  $\sim 1/3$  of its initial intensity and shifts by 1.1 eV to higher BE upon Li deposition, whereas the  $C_{hetero}$  peak loses only 20% in intensity together with a shift by 1.8 eV to higher BE. The  $C_{alkyl}$  peak even grows to 146% of its initial intensity, together with a shift by 1.8 eV to higher BE. Based on these data, the  $F_{TFSA}$  and  $C_{anion}$  peaks are expected to contain contributions from  $F_3C-O_2S-N^-Li^+$  like reaction products. For the C1s<sub>cation</sub> peaks, possible reaction products include various  $C_xH_yN_z$  and  $C_xH_y$  species, which cannot be identified in more detail.

Comparing the decomposition products of the  $[TFSA]^-$  anion and their formation found here by XPS with the decomposition mechanism predicted for reductive electrochemical decomposition of  $[TFSA]^-$  by Howlett *et al.*<sup>53</sup> (also summarized on the left hand side in Fig. 4a), it is found that they are in principle identical. This can plausibly be rationalized by assuming that the impinging Li atoms readily donate their valence electron to form  $Li^+$ , which is comparable to the reductive electrochemical conditions employed by Howlett *et al.*<sup>53</sup>

For comparison we also deposited Li on a  $TiO_2(110)$  surface covered by a multilayer film of [BMP][TFSA] (4.7 ML). The spectra are included in Fig. S2, ESI†. In this case, no  $Ti^{3+}$  species are formed, and therefore Li intercalation can be excluded. Instead, Li only reacts with the [BMP][TFSA] film. The decomposition products are identical with those observed for the reaction with [BMP][TFSA] at submonolayer coverage. In particular the peak which was related to  $O_{Li_xO}$  before is formed as well. In this case, however, it does not stem from  $Li_xO$ , but from decomposition products of the  $[TFSA]^-$  anion such as  $F_3C-O_2S-N^-Li^+$  and  $Li_2SO_{3-4}$  (see above). This also supports our above conclusion that this peak originates not only from oxygen

species in the  $TiO_2(110)$  crystal, which are modified by neighbored intercalated Li, but also from these  $[TFSA]^-$  decomposition products. This measurement furthermore demonstrates that the Li induced decomposition of the [BMP][TFSA] adlayer is not catalysed and not even modified by the  $TiO_2(110)$  surface. The presence of the substrate mainly affects the decomposition reaction by removing part of the Li *via* the competing Li intercalation process.

#### 2.4 Thermal stability of a [BMP][TFSA] + Li adlayer on $TiO_2(110)$

The thermal stability of mixed [BMP][TFSA] + Li adlayers on  $TiO_2(110)$  was investigated by RT deposition of Li (2 ML) on a  $TiO_2(110)$  surface covered by a submonolayer film of the IL (deposition also at RT) and subsequent stepwise heating of the mixed adlayer covered sample. XP spectra recorded after the respective preparation and heating steps are shown in Fig. 8. As illustrated by the formation of  $Ti^{3+}$  species, Li reacts with and intercalates into the  $TiO_2(110)$  substrate also upon deposition at 300 K. Furthermore, it also reacts with the IL, and the decomposition products are largely identical with those formed upon interaction at 80 K. Only the  $S_{d2}$  peak in the S2p regime, which was tentatively related to an adsorbed  $F_3C-O_2S^-Li^+$  species, does not appear in the present measurement. Most likely, this species is not stable at this temperature and directly decomposes to  $SO_4^{2-}$ ,  $SO_3^{2-}$  or  $LiSO_3$  ( $S_{d1}$ ).

Heating the sample to higher temperatures the decomposition of the adlayer continues (Fig. 8). In the O1s region the peak related to the intact adsorbed  $[TFSA]^-$  species rapidly loses in intensity between 300 K and 370 K and has completely vanished at 420 K. The same is true for the  $S_{TFSA}$  signal. Hence, upon annealing to 420 K the intact  $[TFSA]^-$  is completely decomposed. Therefore the other (anion related) peaks remaining at this temperature must arise from decomposition products. This includes the  $C_{anion}$  peak (slightly shifted to a lower BE of 293 eV), the  $N_{anion}$  signal (shifted to a BE of 400 eV), as well as the  $S_{d1}$ ,  $O_{Li_xO}$  and  $F_{TFSA}$  signals. As discussed for Li deposition on predeposited [BMP][TFSA] thin layers at 80 K, the latter peaks contain contributions also from a  $[TFSA]^-$  decomposition product, most likely adsorbed  $F_3C-O_2S-N^-Li^+$ . The  $S_{d1}$  signal is additionally attributed to  $Li_2SO_{3-4}$  and  $O_{Li_xO}$  to  $Li_xO$ . The remaining  $N_{cation}$  peak presumably originates from adsorbed  $C_xH_yN_z$  species, following the proposal by Kroon *et al.* (see Fig. 4b).<sup>52</sup> The peak at BE 161.4 eV (marked as  $S_{d3}$  in Fig. 8) arises due to formation of  $Li_2S$ . Interestingly, when looking at the Ti2p region, the  $Ti^{3+}$  signal is fading and reaches a minimum at 420 K. It seems that the decomposition reaction of the IL adlayer is induced by reaction with  $Ti^{3+}$  species that were intercalated in the near surface region (first few nm) of the  $TiO_2(110)$  substrate and start segregating to the surface. Apparently, this goes along with a re-oxidation of the  $Ti^{3+}$  species to  $Ti^{4+}$ . Therefore, at this temperature the  $O_{Li_xO}$  peak is no longer (co-)generated by O atoms in the  $TiO_2(110)$  crystal, but solely due to the decomposition products of the IL. In the Li1s region, there is a small change of the BE from 56.4 eV at 340 and 370 K to 56.1 eV at 420 K. For higher temperatures the Li1s peak shifts



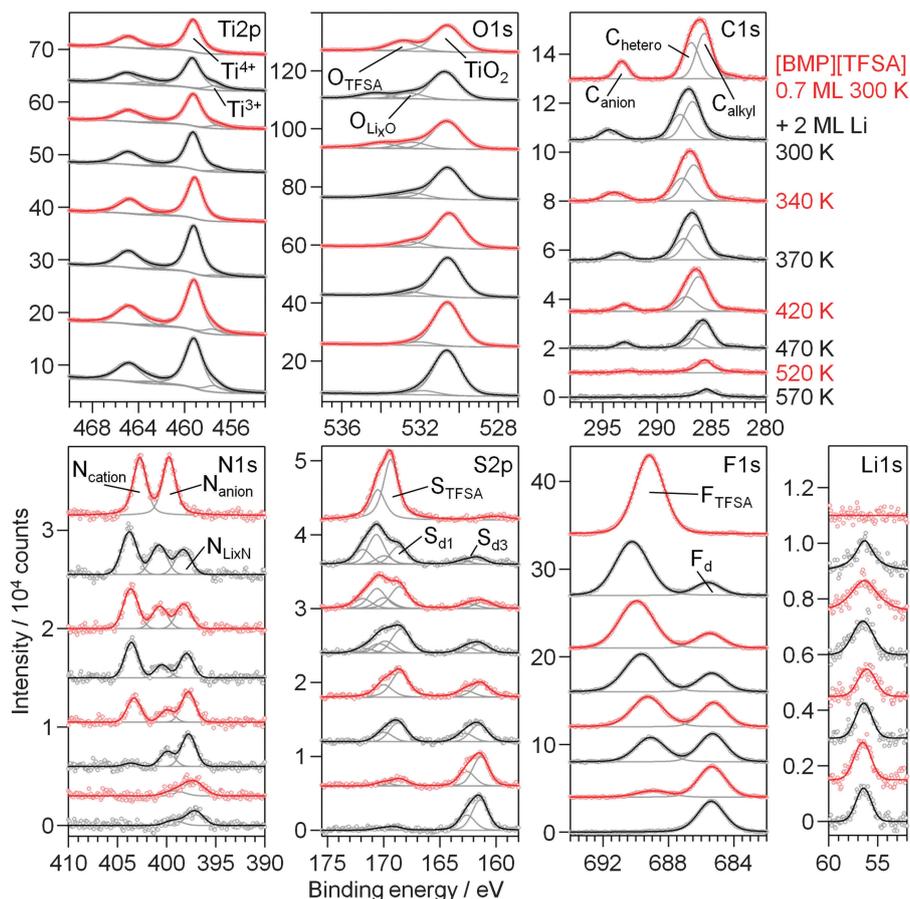


Fig. 8 XP detail spectra of 0.7 ML [BMP][TFSA] on  $\text{TiO}_2(110)$  at 300 K (topmost spectra), after post-deposition of 2 ML Li at 300 K and upon stepwise heating to 570 K.

back to 56.4 eV and the  $\text{Ti}^{3+}$  species forms again. This occurs in parallel to the further decomposition of the adlayer, which is indicated by a strong decrease of the  $\text{N}_{\text{cation}}$  peak between 370 K and 470 K. We explain these changes to a decomposition of the  $\text{C}_x\text{H}_y\text{N}_z$  species related to this peak to H-containing carbonaceous species. At 520 K, also the peaks related to  $\text{F}_3\text{C}-\text{O}_2\text{S}-\text{N}^-\text{Li}^+$  have further decreased in intensity, and they completely vanished after annealing to 570 K. At this point also most of the adsorbed carbonaceous species have decomposed and desorbed. The only remaining species are  $\text{LiF}$  ( $=\text{F}_d$ ),  $\text{Li}_2\text{S}$  ( $=\text{S}_{d3}$ ),  $\text{Li}_x\text{N}$  (=peak at 399 eV) and  $\text{TiN}_x$  (=peak at 397.1 eV). The remaining  $\text{O}_{\text{Li}_x\text{O}}$  peak is due to the interaction of bulk O atoms with intercalated  $\text{Li}^+$  (and presumably  $\text{Li}_x\text{O}$  on the surface).

In a second experiment we tested the stability and decomposition of a mixed [BMP][TFSA] + Li adlayer after the reversed deposition order. In this case Li was deposited on  $\text{TiO}_2(110)$  at 300 K, followed by post-deposition of [BMP][TFSA]. The amount of Li is higher in this case than in the above experiment to achieve saturation of the first substrate layers by intercalated Li. Subsequently, a comparable submonolayer amount of [BMP][TFSA] was evaporated on the Li modified  $\text{TiO}_2(110)$  substrate. XP spectra recorded after IL deposition and subsequent heating steps are shown in Fig. 9. Already after IL deposition at 300 K decomposition products are visible in the spectra (topmost spectra), in contrast to

deposition on pure  $\text{TiO}_2(110)$ . Apparently, the surface species formed during the preceding Li intercalation process,  $\text{Ti}^{3+}$  species and possibly also surface near  $\text{Li}^+$  species, partly react with the [BMP][TFSA]. Here it should be noted that  $\text{Ti}^{3+}$  surface species were reported to be more reactive than  $\text{Ti}^{4+}$  species already mentioned previously.<sup>63</sup> Therefore, in the  $\text{Ti}2p$  regime the  $\text{Ti}^{3+}$  peak intensity decreases relative to the  $\text{Ti}^{4+}$  peak (see also our above discussion). Overall, the new peaks formed upon [BMP][TFSA] decomposition are largely identical to those obtained for the reaction of post-deposited Li with a pre-deposited [BMP][TFSA] adlayer. The total amount of decomposition products, however, is smaller in the present case. In particular the decomposition peaks in the  $\text{S}2p$  region are less intense, and the  $\text{N}_{\text{anion}}$  peak does not show a strong decay. The amount of  $\text{LiF}$  ( $=\text{F}_d$ ) and  $\text{Li}_x\text{N}$  ( $=\text{N}_{\text{Li}_x\text{N}}$ ) species is comparable. These differences may be due to the lower availability of the  $\text{Ti}^{3+}$  species due to diffusion limitations, as compared to direct deposition of Li on predeposited [BMP][TFSA] layers.

Upon sequential heating of the sample (lower spectra in Fig. 9), the  $\text{O}_{\text{TFSA}}$  peak as well as the  $\text{S}_{\text{TFSA}}$  signal vanished at 370–420 K, as it was also the case for  $\text{Li}|\text{IL}|\text{TiO}_2(110)$ . In the same temperature range,  $\text{Ti}^{3+}$  again has a minimum in intensity. The peaks remaining after this heating step, the  $\text{C}_{\text{anion}}$  (BE 293.3 eV),  $\text{N}_{\text{anion}}$  (BE 400.5 eV),  $\text{S}_{d1}$  (BE 168.7 eV),  $\text{O}_{\text{Li}_x\text{O}}$  (BE 532.8 eV) and



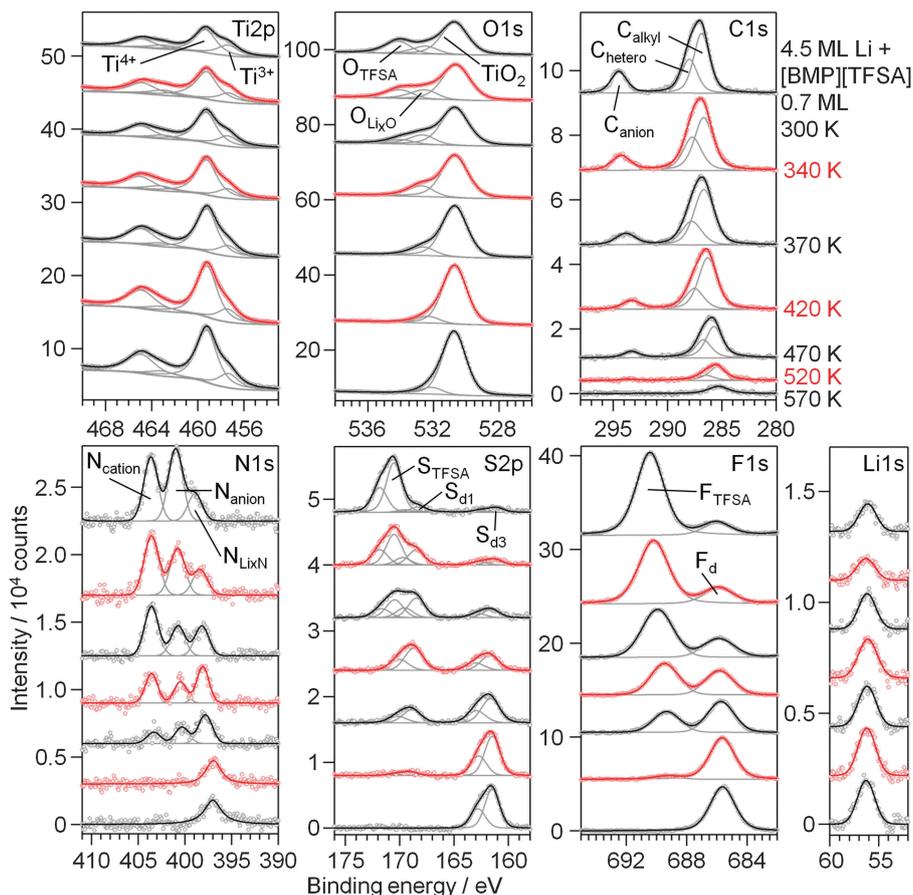


Fig. 9 XPS detail spectra of 0.7 ML [BMP][TFSA] post-deposited on a  $\text{TiO}_2(110)$  surface pre-covered by 4.5 ML Li at 300 K (topmost spectra) and upon stepwise heating to 570 K.

$\text{F}_{\text{TFSA}}$  (BE 689.5 eV) peaks, are again related to the formation of  $\text{F}_3\text{C}-\text{O}_2\text{S}-\text{N}^-\text{Li}^+$ . They also vanished upon heating to 520 K, indicating the further decomposition of this product. The  $[\text{BMP}]^+$  cations are presumably also completely decomposed at 420 K and the remaining  $\text{C}_{\text{alkyl}}$ ,  $\text{C}_{\text{hetero}}$  and  $\text{N}_{\text{cation}}$  peaks are due to remaining  $\text{C}_x\text{H}_y\text{N}_z$  species, which decompose until 520 K. Above this temperature the surface mainly contains adsorbed  $\text{LiF}$  and  $\text{Li}_2\text{S}$  (or/and  $\text{S}_{\text{ad}}$ ) species, and a small amount of remaining carbonaceous species are visible. Finally, a peak remaining in the  $\text{N}1\text{s}$  region with a BE of 397 eV, which at  $T > 520$  K is shifted by 1 eV to lower BE compared to  $\text{N}_{\text{Li}_x\text{N}}$ , is again attributed to  $\text{TiN}_x$  species. For annealing to higher temperatures, e.g. to 630 K, we would expect this species to disappear as well, similar to our findings for annealing of a pure  $[\text{BMP}][\text{TFSA}]$  adlayer on  $\text{TiO}_2(110)$ .

Overall, the thermal decomposition of  $[\text{BMP}][\text{TFSA}]$  in interaction with codeposited Li does not seem to depend on the sequence of evaporation, whether Li is evaporated first followed by  $[\text{BMP}][\text{TFSA}]$  deposition or the other way around.

### 3 Summary and conclusion

Employing XPS and STM measurements for chemical and structural characterization, we have investigated the interaction

of  $[\text{BMP}][\text{TFSA}]$  with a well-defined  $\text{TiO}_2(110)$  surface, and the influence of Li thereon. For comparison, the interaction of Li with the bare  $\text{TiO}_2(110)$  substrate was investigated as well. The results of the present study in combination with previous results lead us to the following conclusions:

(1) In the temperature range between 80 K and 380 K,  $[\text{BMP}][\text{TFSA}]$  adsorbs as intact ion pairs on the substrate, with both anion and cation in direct contact with the surface. At 80 K, the  $[\text{BMP}][\text{TFSA}]$  ion pairs are adsorbed in lines along the  $\text{TiO}_2(110)$   $(1 \times 1)$  [001] direction and uniformly cover the surface. The adsorbed species are highly mobile, molecular resolution STM imaging of the IL was only possible at this temperature and for a coverage of around 1 ML, while for other conditions only frizzy features could be resolved.

(2) Above 380 K, the adsorbed IL starts to decompose in a complex multistep mechanism.  $\text{S}_{\text{ad}}$ ,  $\text{F}_{\text{ad}}$ ,  $\text{TiN}_x$  and also adsorbed  $\text{C}_x\text{H}_y\text{N}_z$  species are detected by XPS, which at  $T > 530$  K further decompose to carbonaceous species. The reaction products fit to earlier proposals for the  $[\text{BMP}][\text{TFSA}]$  decomposition based on quantum chemical calculations.<sup>52–54</sup> According to this mechanism the  $\text{C}_x\text{H}_y\text{N}_z$  species are presumably  $[\text{BMP}]^+$  species with one of the C–N bonds broken, while the carbon containing species remaining at higher temperatures are butyl and methyl fragments, which can further decompose.



The decomposition of the intact [BMP][TFSA] is completed after heating to 560 K.

(3) Li reacts already at 80 K with the TiO<sub>2</sub>(110) substrate, by intercalation into the near surface region of the TiO<sub>2</sub>(110) crystal, which leads to the formation of Li<sup>+</sup> and Ti<sup>3+</sup> species, and, at higher coverages, by the formation of a Li<sub>x</sub>O cover layer. Intercalation results in a stoichiometry which is very close to that of Li loaded LTO (Li<sub>7</sub>Ti<sub>5</sub>O<sub>12</sub>). Upon heating to room temperature, the Ti<sup>3+</sup> and Li<sup>+</sup> species move deeper into the TiO<sub>2</sub>(110) bulk, but for the amounts investigated here (up to 9 ML equivalents of Li) they remain in the near surface region (in the topmost 5 nm), as evidenced by XPS measurements at the normal detection angle of the crystal.

(4) Deposition of Li at 80 K onto a TiO<sub>2</sub>(110) substrate pre-covered by a [BMP][TFSA] submonolayer film results in two competing reactions: reactive interaction of the Li atoms with the [BMP][TFSA] adlayer and intercalation into the TiO<sub>2</sub>(110) substrate forming intercalated Li<sup>+</sup> and Ti<sup>3+</sup> species. Reactive interaction between the IL adsorbate and Li results in decomposition products which are rather similar to those obtained for thermal decomposition of [BMP][TFSA] on TiO<sub>2</sub>(110), such as Li<sub>2</sub>S/S<sub>ad</sub> or LiF/F<sub>ad</sub>, but also additional peaks are found in XPS. Based on the XP binding energies, F<sub>3</sub>C–O<sub>2</sub>S–N<sup>–</sup>Li<sup>+</sup> and F<sub>3</sub>C–O<sub>2</sub>S<sup>–</sup>Li<sup>+</sup> species are formed as well as intermediate products such as SO<sub>x</sub>, in agreement with the theoretical proposals mentioned above. For Li deposition with pre-deposited multilayer [BMP][TFSA] films, the reaction with the IL is more efficient than diffusion of Li to the TiO<sub>2</sub>(110) surface, and reactive interaction with the TiO<sub>2</sub>(110) surface (intercalation) is therefore inhibited.

(5) For temperatures up to ca. 340 K, the adlayer composition changes only little. Mainly the S<sub>d2</sub> peak vanishes, indicating that F<sub>3</sub>C–O<sub>2</sub>S<sup>–</sup>Li<sup>+</sup> is further decomposed. At higher temperatures further decomposition takes place: the remaining intact [BMP][TFSA] as well as the intermediate decomposition products F<sub>3</sub>C–O<sub>2</sub>S–N<sup>–</sup>Li<sup>+</sup>, SO<sub>x</sub> and C<sub>x</sub>H<sub>y</sub>N<sub>z</sub> species react with the Ti<sup>3+</sup> and Li<sup>+</sup> species in the substrate, forming similar reaction products as obtained for the thermal decomposition of [BMP][TFSA] on TiO<sub>2</sub>(110), such as F<sub>ad</sub>/LiF, S<sub>ad</sub>/Li<sub>2</sub>S and TiN<sub>x</sub>/Li<sub>x</sub>N and some residual carbonaceous compounds. The decomposition of [BMP][TFSA] upon reactive interaction with TiO<sub>2</sub>(110) and Li is complete at a temperature of 420 K, the one of F<sub>3</sub>C–O<sub>2</sub>S–N<sup>–</sup>Li<sup>+</sup> at 520–570 K. The Ti<sup>3+</sup> species consumed in the decomposition process are reformed above 420 K. When evaporating [BMP][TFSA] on a TiO<sub>2</sub>(110) surface pre-covered with Li, the IL decomposes partly already at 300 K. Since this was not the case on pure TiO<sub>2</sub>(110), this higher reactivity is mainly attributed to the Ti<sup>3+</sup> species intercalated in the near surface region. Upon heating to higher temperatures the decomposition process of [BMP][TFSA] is similar to that on samples where Li is deposited on TiO<sub>2</sub>(110) pre-covered with [BMP][TFSA].

Overall, this study gained detailed insight on a molecular level on the reactive interaction of [BMP][TFSA] with a TiO<sub>2</sub>(110) substrate and on the effect of Li on this interaction. These insights will be highly valuable and provide a solid basis for subsequent studies on the interaction between IL based electrolytes

and TiO<sub>2</sub> electrodes and on the initial stages of the formation of the solid-electrolyte interphase (SEI) between these compounds in lithium ion battery applications.

## 4 Experimental

The measurements were performed in a commercial UHV system (SPECS GmbH) with a base pressure of  $< 2 \times 10^{-10}$  mbar. It is equipped with a combined Aarhus-type STM/AFM system (SPECS; Aarhus SPM 150 with Colibri Sensor), a non-monochromatic X-ray source (SPECS XR50, Al-K<sub>α</sub> and Mg-K<sub>α</sub>) and a hemispherical analyzer (SPECS, DLSEGD-Phoibos-HAS3500) for XPS measurements, and standard facilities for surface preparation. STM as well as XPS measurements can be carried out with the sample between 90–400 K by cooling with LN<sub>2</sub> and resistive heating; for XPS measurements the sample can also be heated to temperatures up to 1000 K by electron bombardment (from the backside of the sample). All STM images were acquired in constant current mode with tunnelling currents between 20 and 200 pA and tunnelling voltages of 0.1–2 V. For XPS measurements we used an Al K<sub>α</sub> X-ray source (1486.6 eV), operated at a power of 150 W ( $U = 12$  kV,  $I = 12.5$  mA). The spectra were recorded at a pass energy of 100 eV at detection angles of 0° and 80° relative to the surface normal. At 80° the surface sensitivity of the measurement is significantly higher than at normal emission (0° to the surface normal). The spectra shown in the figures are all taken at an electron emission angle of 80° if not stated otherwise. Peak fitting was performed with CasaXPS as well as Igor Pro 6. For background subtraction a Shirley background (if necessary convoluted with a parabolic background) was used, the peaks were fitted using a Pseudo-Voigt type peak shape, which is a convolution of a Gaussian and Lorentzian function. For the background subtraction of the Li peaks it was additionally necessary to subtract the signal of the empty TiO<sub>2</sub> surface before performing the Shirley background subtraction due to the small intensity of the Li1s peak.

Rutile TiO<sub>2</sub>(110) crystals were purchased from MaTeck GmbH with one side polished with an orientation accuracy of  $< 0.1^\circ$  and a surface roughness  $< 0.03$  μm. The initial preparation of the rutile TiO<sub>2</sub> crystals follows procedures reported in the literature:<sup>76</sup> the crystals were cleaned in an ultrasonic bath in 2 × 5 min acetone, 5 min in carotic acid and 2 × 5 min in ultrapure water (in between the steps the crystals were rinsed with ultrapure water). Afterwards the crystals were calcined at 1173 K under air atmosphere and cooled down at 5 K min<sup>–1</sup>. This procedure produces clean and fully oxidized colorless or slightly yellow rutile crystals. These crystals are transferred into the UHV chamber and heated in UHV to 1023 K for 5 h which partly reduces the sample (formation of Ti<sup>3+</sup> species and oxygen vacancies, details can be found elsewhere<sup>18,46</sup>). Upon reduction, the color of the rutile crystal changes to blue. To achieve a properly prepared smooth (110) surface several cycles of Ar<sup>+</sup> ion sputtering (0.5 kV, 45 min) and heating to 973 K for 30 min were carried out. The surface structure and cleanliness were checked by STM and XPS (see *e.g.* Fig. 2a, b and 5) The surfaces show the



typical ( $1 \times 1$ ) reconstruction reported by Diebold *et al.*<sup>17,18</sup> with protruding rows running in the [001] direction. Bright round protrusions situated in between the bright rows are attributed to missing bridge-bonded oxygens on the surface. Terraces range over  $>100$  nm. The XP spectra also show a contamination free substrate.

The ionic liquid [BMP][TFSA] was purchased from Merck in the highest available quality (ultrapure). It was filled in a glass crucible, which prior to use was baked at 870 K in UHV to remove vaporizable impurities from the glass. The crucible with the IL was positioned in a triple Knudsen effusion evaporator (ventiotec, OVD-3) and attached to the UHV chamber. The IL was degassed for several weeks under UHV conditions at room temperature and subsequently for at least 24 hours at 350 K which removes all vaporizable contaminations (such as water) from the IL. Karl-Fischer titration of the IL after room temperature degassing revealed water contents of well below 25 ppm, the water content after 350 K degassing should be considerably lower. Furthermore, the cleanness of the IL vapor was checked using a quadrupole mass spectrometer (QMS; Pfeiffer HiQuadQMG700), which was mounted with its ionization area into the IL beam; also the evaporation rate was calibrated against the evaporator temperature using the QMS. An evaporation temperature of 453 K and 5 min evaporation time results in a coverage of 1 ML. The [BMP][TFSA] adlayers were found to slowly decompose in the X-ray beam during XPS measurements, as it is already known from previous experiments applying non-monochromatized X-ray radiation.<sup>77</sup> We quantified the beam damage occurring on our samples by measuring the decay of all peaks in the XP spectra assigned to [BMP][TFSA] over several hours. A decomposition rate of *ca.*  $0.002 \text{ min}^{-1}$  of the total intensities of the peaks for both the anion as well as the cation was found for a [BMP][TFSA] adlayer with a coverage of 0.7 ML at room temperature and with the X-ray source operated at 150 W. Decomposition products due to beam damage could not be detected in the spectra, which is clearly different to the decomposition due to addition of Li or the thermal decomposition of the adlayers described in Sections 2.1 and 2.3. We assume that the reaction products formed during exposure to X-ray radiation are desorbed right after formation. The decomposition mechanism under X-ray radiation is therefore completely different from the one induced thermally and/or Li incorporation. To minimize the effects of the beam damage in our measurements, the time for irradiation per sample was kept as short as possible, switching off the X-ray source between the XPS measurements. Furthermore, in the annealing experiments the heating was switched off before recording the XP spectra, these were therefore recorded at temperatures considerably lower than the respective annealing temperatures. In the evaluation of the XPS data the loss of intensity due to beam damage is taken into account, using the decomposition rate determined before.

Li was evaporated from a Li metal dispenser (SAES Getters) incorporated into a homemade electrical heating stage. The dispenser was operated at 7.1 A. The sample was placed in line-of-sight in front of the dispenser, the evaporation time was controlled with a shutter in between the sample and the dispenser.

For calibration of the Li evaporation rate a Cu(111) crystal was used, which was cooled to 80 K. Under these conditions Li adsorbs atomically on the surface, it does not react with the substrate and only very slowly with residual contaminations adsorbing from the UHV gas phase.<sup>28</sup> The coverage was evaluated by the damping of the  $\text{Cu}2p_{3/2}$  peak. For our setup a deposition rate of  $0.049 \pm 0.009 \text{ ML min}^{-1}$  was calculated. The cleanness of the deposited Li adlayers both on Cu(111) and on  $\text{TiO}_2(110)$  was checked by XPS and partly by STM.

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## References

- 1 J. S. Wilkes, *Green Chem.*, 2002, **4**(2), 73–80.
- 2 P. Wasserscheid and W. Keim, *Angew. Chem., Int. Ed.*, 2000, **39**(21), 3772–3789.
- 3 G. B. Appetecchi, M. Montanino, D. Zane, M. Carewska, F. Alessandrini and S. Passerini, *Electrochim. Acta*, 2009, **54**(4), 1325–1332.
- 4 J. M. Slattery, C. Dagueuet, P. J. Dyson, T. J. S. Schubert and I. Krossing, *Angew. Chem.*, 2007, **119**(28), 5480–5484.
- 5 M. Armand and J. M. Tarascon, *Nature*, 2008, **451**(7179), 652–657.
- 6 M. Armand, F. Endres, D. R. MacFarlane, H. Ohno and B. Scrosati, *Nat. Mater.*, 2009, **8**(8), 621–629.
- 7 A. A. Kornyshev, *J. Phys. Chem. B*, 2007, **111**(20), 5545–5557.
- 8 F. Silva, C. Gomes, M. Figueiredo, R. Costa, A. Martins and C. M. Pereira, *J. Electroanal. Chem.*, 2008, **622**(2), 153–160.
- 9 B. Uhl, T. Cremer, M. Roos, F. Maier, H. P. Steinrück and R. J. Behm, *Phys. Chem. Chem. Phys.*, 2013, **15**, 17295–17308.
- 10 F. Buchner, K. Forster-Tonigold, B. Uhl, D. Alwast, N. Wagner, H. Farkhondeh, A. Groß and R. J. Behm, *ACS Nano*, 2013, **7**(9), 7773–7784.
- 11 B. Uhl, F. Buchner, D. Alwast, N. Wagner and R. J. Behm, *Beilstein J. Nanotechnol.*, 2013, **4**(1), 903–918.
- 12 B. Uhl, F. Buchner, S. Gabler, M. Bozorgchenani and R. J. Behm, *Chem. Commun.*, 2014, **50**(62), 8601–8604.
- 13 B. Uhl, H. Huang, D. Alwast, F. Buchner and R. J. Behm, *Phys. Chem. Chem. Phys.*, 2015, **17**, 23816–23832.
- 14 Z. Chen, I. Belharouak, Y.-K. Sun and K. Amine, *Adv. Funct. Mater.*, 2013, **23**(8), 959–969.
- 15 S. Goriparti, E. Miele, F. De Angelis, E. Di Fabrizio, R. P. Zaccaria and C. Capiglia, *J. Power Sources*, 2014, **257**, 421–443.



- 16 T. Fröschl, U. Hörmann, P. Kubiak, G. Kucerova, M. Pfanzelt, C. K. Weiss, R. J. Behm, N. Hüsing, U. Kaiser and K. Landfester, *Chem. Soc. Rev.*, 2012, **41**(15), 5313–5360.
- 17 U. Diebold, *Appl. Phys. A: Mater. Sci. Process.*, 2003, **76**(5), 681–687.
- 18 U. Diebold, *Surf. Sci. Rep.*, 2003, **48**(5), 53–229.
- 19 R. Hayes, G. G. Warr and R. Atkin, *Chem. Rev.*, 2015, **115**(13), 6357–6426.
- 20 S. Krischok, J. A. Schaefer, O. Höfft and V. Kempter, *Surf. Interface Anal.*, 2004, **37**(1), 83–89.
- 21 H. Tatsumi, A. Sasahara and M. Tomitori, *J. Phys. Chem. C*, 2012, **116**(25), 13688–13692.
- 22 A. Henningsson, M. P. Andersson, P. Uvdal, H. Siegbahn and A. Sandell, *Chem. Phys. Lett.*, 2002, **360**(1), 85–90.
- 23 A. Henningsson, H. Rensmo, A. Sandell, H. Siegbahn, S. Södergren, H. Lindström and A. Hagfeldt, *J. Chem. Phys.*, 2003, **118**, 5607–5612.
- 24 S. Södergren, H. Siegbahn, H. Rensmo, H. Lindström, A. Hagfeldt and S.-E. Lindquist, *J. Phys. Chem. B*, 1997, **101**(16), 3087–3090.
- 25 M. Pfanzelt, *Development of TiO<sub>2</sub> Rutile as Negative Electrode Material for Lithium Ion Batteries*, PhD thesis, University Ulm, ZSW, 2012.
- 26 B. Zachau-Christiansen, K. West, T. Jacobsen and S. Atlung, *Solid State Ionics*, 1988, **28**, 1176–1182.
- 27 M. Olschewski, R. Gustus, M. Marschewski, O. Höfft and F. Endres, *Phys. Chem. Chem. Phys.*, 2014, **16**(47), 25969–25977.
- 28 F. Buchner, M. Bozorgchenani, B. Uhl, H. Farkhondeh, J. Bansmann and R. J. Behm, *J. Phys. Chem. C*, 2015, **119**(29), 16649–16659.
- 29 U. Diebold and T. E. Madey, *Surf. Sci. Spectra*, 1996, **4**(3), 227–231.
- 30 T. Cremer, M. Killian, J. M. Gottfried, N. Paape, P. Wasserscheid, F. Maier and H. P. Steinrück, *ChemPhysChem*, 2008, **9**(15), 2185–2190.
- 31 T. Cremer, M. Stark, A. Deyko, H. P. Steinrück and F. Maier, *Langmuir*, 2011, **27**, 3662–3671.
- 32 T. Cremer, L. Wibmer, S. K. Calderon, A. Deyko, F. Maier and H. P. Steinrück, *Phys. Chem. Chem. Phys.*, 2012, **14**(15), 5153–5163.
- 33 U. Diebold, J. Lehman, T. Mahmoud, M. Kuhn, G. Leonardelli, W. Hebenstreit, M. Schmid and P. Varga, *Surf. Sci.*, 1998, **411**(1), 137–153.
- 34 S. Fischer, A. W. Munz, K. D. Schierbaum and W. Göpel, *Surf. Sci.*, 1995, **337**(1), 17–30.
- 35 K. i. Fukui, H. Onishi and Y. Iwasawa, *Phys. Rev. Lett.*, 1997, **79**(21), 4202–4205.
- 36 S. Wendt, R. Schaub, J. Matthiesen, E. K. Vestergaard, E. Wahlström, M. D. Rasmussen, P. Thstrup, L. M. Molina, E. Laegsgaard, I. Stensgaard, B. Hammer and F. Besenbacher, *Surf. Sci.*, 2005, **598**, 226–245.
- 37 Z. Zhang, J. Lee, J. T. Yates, R. Bechstein, E. Lira, J. Hansen, S. Wendt and F. Besenbacher, *J. Phys. Chem. C*, 2010, **114**(7), 3059–3062.
- 38 E. Lira, J. Hansen, P. Huo, R. Bechstein, P. Galliker, E. Laegsgaard, B. Hammer, S. Wendt and F. Besenbacher, *Surf. Sci.*, 2010, **604**(21), 1945–1960.
- 39 J. Hansen, P. Huo, U. Martinez, E. Lira, Y. Y. Wei, R. Streber, E. Laegsgaard, B. Hammer, S. Wendt and F. Besenbacher, *Phys. Rev. Lett.*, 2011, **107**(13), 136102.
- 40 U. Martinez, J. Hansen, E. Lira, H. H. Kristoffersen, P. Huo, R. Bechstein, E. Laegsgaard, F. Besenbacher, B. Hammer and S. Wendt, *Phys. Rev. Lett.*, 2012, **109**(15), 155501.
- 41 R. Wen, B. Rahn and O. M. Magnussen, *Angew. Chem.*, 2015, **127**(20), 5889.
- 42 T. Waldmann, H. H. Huang, H. E. Hoster, O. Höfft, F. Endres and R. J. Behm, *ChemPhysChem*, 2011, **12**, 2565–2567.
- 43 B. J. Burrow, A. E. Morgan and R. C. Ellwanger, *J. Vac. Sci. Technol., A*, 1986, **4**(6), 2463–2469.
- 44 J. A. Rodriguez, J. Hrbek, J. Dvorak, T. Jirsak and A. Maiti, *Chem. Phys. Lett.*, 2001, **336**(5), 377–384.
- 45 E. L. D. Hebenstreit, W. Hebenstreit and U. Diebold, *Surf. Sci.*, 2001, **470**(3), 347–360.
- 46 E. L. D. Hebenstreit, W. Hebenstreit and U. Diebold, *Surf. Sci.*, 2000, **461**(1), 87–97.
- 47 A. M. Czoska, S. Livraghi, M. Chiesa, E. Giamello, S. Agnoli, G. Granozzi, E. Finazzi, C. D. Valentin and G. Pacchioni, *J. Phys. Chem. C*, 2008, **112**(24), 8951–8956.
- 48 J. C. Yu, J. Yu, W. Ho, Z. Jiang and L. Zhang, *Chem. Mater.*, 2002, **14**(9), 3808–3816.
- 49 H. Park and W. Choi, *J. Phys. Chem. B*, 2004, **108**(13), 4086–4093.
- 50 R. E. Banks, B. E. Smart and J. C. Tatlow, *Organofluorine Chemistry: Principles and Commercial Applications*, Springer, 1994.
- 51 S. Schernich, M. Laurin, Y. Lykhach, H. P. Steinrück, N. Tsud, T. Skala, K. C. Prince, N. Taccardi, V. Matolin, P. Wasserscheid and J. Libuda, *J. Phys. Chem. Lett.*, 2013, **4**(1), 30–35.
- 52 M. C. Kroon, W. Buijs, C. J. Peters and G. J. Witkamp, *Green Chem.*, 2006, **8**(3), 241–245.
- 53 P. C. Howlett, E. I. Izgorodina, M. Forsyth and D. R. MacFarlane, *Z. Phys. Chem.*, 2006, **220**(10), 1483–1498.
- 54 M. C. Kroon, W. Buijs, C. J. Peters and G. J. Witkamp, *Thermochim. Acta*, 2007, **465**(1), 40–47.
- 55 R. E. Del Sesto, T. M. McCleskey, C. Macomber, K. C. Ott, A. T. Koppisch, G. A. Baker and A. K. Burrell, *Thermochim. Acta*, 2009, **491**(1), 118–120.
- 56 S. E. Stein, “Mass Spectra” by NIST Mass Spec Data Center, NIST Chemistry WebBook, NIST Standard Reference Database Number 69, ed. P. J. Linstrom and W. G. Mallard, National Institute of Standards and Technology, Gaithersburg MD, 20899, <http://webbook.nist.gov>, 2015.
- 57 C. Xu, B. E. Koel and M. T. Paffett, *Langmuir*, 1994, **10**(1), 166–171.
- 58 S. L. Tait, Z. Dohnalek, C. T. Campbell and B. D. Kay, *J. Chem. Phys.*, 2005, **122**(16), 164707.
- 59 R. M. Slayton, C. M. Aubuchon, T. L. Camis, A. R. Noble and N. J. Tro, *J. Phys. Chem.*, 1995, **99**(7), 2151–2154.
- 60 P. E. Pehrsson, P. Armistead and S. Sastri, *Surf. Interface Anal.*, 1993, **20**(7), 573–582.
- 61 K. Hamrin, G. Johansson, U. Gelius, C. Nordling and K. Siegbahn, *Phys. Scr.*, 1970, **1**(5–6), 277–280.



- 62 D. R. Warburton, D. Purdie, C. A. Muryn, K. Prabhakaran, P. L. Wincott and G. Thornton, *Surf. Sci.*, 1992, **269**, 305–309.
- 63 K. E. Smith, J. L. Mackay and V. E. Henrich, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1987, **35**(11), 5822–5829.
- 64 H. Onishi, T. Aruga, C. Egawa and Y. Iwasawa, *Surf. Sci.*, 1988, **193**(1), 33–46.
- 65 K. E. Smith and V. E. Henrich, *J. Vac. Sci. Technol., A*, 1989, **7**(3), 1967–1971.
- 66 S. P. Kowalczyk, L. Ley, F. R. McFeely, R. A. Pollak and D. A. Shirley, *Phys. Rev. B: Solid State*, 1973, **8**(8), 3583–3585.
- 67 L. El Ouatani, R. Dedryvere, C. Siret, P. Biensan and D. Gonbeau, *J. Electrochem. Soc.*, 2009, **156**(6), A468–A477.
- 68 L. Liu, V. E. Henrich, W. P. Ellis and I. Shindo, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1996, **54**(3), 2236–2239.
- 69 A. C. Kozen, A. J. Pearse, C. F. Lin, M. A. Schroeder, M. Noked, S. B. Lee and G. W. Rubloff, *J. Phys. Chem. C*, 2014, **118**(48), 27749–27753.
- 70 M. P. Seah and S. J. Spencer, *Surf. Interface Anal.*, 2002, **33**(8), 640–652.
- 71 M. Kitta, T. Akita, Y. Maeda and M. Kohyama, *Langmuir*, 2012, **28**(33), 12384–12392.
- 72 B. Lindberg, K. Hamrin, G. Johansson, U. Gelius, A. Fahlman, C. Nordling and K. Siegbahn, *Phys. Scr.*, 1970, **1**(5–6), 286–298.
- 73 K. R. Zavadil and N. R. Armstrong, *Surf. Sci.*, 1990, **230**(1), 61–73.
- 74 X. Feng, M. K. Song, W. C. Stolte, D. Gardenghi, D. Zhang, X. Sun, J. Zhu, E. J. Cairns and J. Guo, *Phys. Chem. Chem. Phys.*, 2014, **16**(32), 16931–16940.
- 75 C. Xu, B. Sun, T. Gustafsson, K. Edström, D. Brandell and M. Hahlin, *J. Mater. Chem. A*, 2014, **2**(20), 7256–7264.
- 76 S. Kielbassa, M. Kinne and R. J. Behm, *J. Phys. Chem. B*, 2004, **108**(50), 19184–19190.
- 77 A. Keppler, M. Himmerlich, T. Ikari, M. Marschewski, E. Pachomow, O. Höfft, W. Maus-Friedrichs, F. Endres and S. Krischok, *Phys. Chem. Chem. Phys.*, 2011, **13**, 1174–1181.

