


 Cite this: *Phys. Chem. Chem. Phys.*,
2016, 18, 8364

Coupled optical absorption, charge carrier separation, and surface electrochemistry in surface disordered/hydrogenated TiO₂ for enhanced PEC water splitting reaction†

 Dilip Kumar Behara,^a Ashok Kumar Ummireddi,^a Vidyasagar Aragonda,^a
Prashant Kumar Gupta,^a Raj Ganesh S. Pala^{*ab} and Sri Sivakumar^{*abc}

The central governing factors that influence the efficiency of photoelectrochemical (PEC) water splitting reaction are photon absorption, effective charge-carrier separation, and surface electrochemistry. Attempts to improve one of the three factors may debilitate other factors and we explore such issues in hydrogenated TiO₂, wherein a significant increase in optical absorption has not resulted in a significant increase in PEC performance, which we attribute to the enhanced recombination rate due to the formation of amorphization/disorderness in the bulk during the hydrogenation process. To this end, we report a methodology to increase the charge-carrier separation with enhanced optical absorption of hydrogenated TiO₂. Current methodology involves hydrogenation of non-metal (N and S) doped TiO₂ which comprises (1) lowering of the band gap through shifting of the valence band via less electronegative non-metal N, S-doping, (2) lowering of the conduction band level and the band gap via formation of the Ti³⁺ state and oxygen vacancies by hydrogenation, and (3) material processing to obtain a disordered surface structure which favors higher electrocatalytic (EC) activity. This design strategy yields enhanced PEC activity (%ABPE = 0.38) for the N–S co-doped TiO₂ sample hydrogenated at 800 °C for 24 h over possible combinations of N–S co-doped TiO₂ samples hydrogenated at 500 °C/24 h, 650 °C/24 h and 800 °C/72 h. This suggests that hydrogenation at lower temperatures does not result in much increase in optical absorption and prolonged hydrogenation results in an increase in optical absorption but a decrease in charge carrier separation by forming disorderness/oxygen vacancies in the bulk. Furthermore, the difference in double layer capacitance (*C_{dl}*) calculated from electrochemical impedance spectroscopy (EIS) measurements of these samples reflects the change in the electrochemical surface area (ECSA) and facilitates assessing the key role of surface electrochemistry in PEC water splitting reaction. Additionally, we observed a blue-shift of the absorption spectrum and a decrease in both electrochemical (EC) and photoelectrochemical (PEC) activities after the removal of surface layers through focused ion beam (FIB) sputtering suggesting the importance of surface defects and photon absorption.

 Received 18th July 2015,
Accepted 4th February 2016

DOI: 10.1039/c5cp04212g

www.rsc.org/pccp

1. Introduction

Recently, black hydrogenated titania was reported to enhance photocatalytic performance by creating surface disorderness that helps in improving visible light absorption.¹ Furthermore,

the surface defects (oxygen vacancies/Ti³⁺ states) created upon amorphization will enhance the electrical conductivity which is essential for wide band gap materials like TiO₂.^{1–8} Therefore, stabilization of surface defects is essential for better utilization of TiO₂ material for photocatalytic applications in the long run. However, the stabilization of these defects is challenging and control of these oxygen vacancies is highly difficult as they depend on parameters such as temperature, surface states and synthesis conditions. Recently, Lin *et al.* made an attempt to make better utilization of surface defects by addition of dopants to the surface disordered TiO₂ material.⁴ They synthesized oxygen deficient TiO_{2–x} by hydrogenation followed by post-doping of non-metals (H, N, and I) and observed enhanced activity in photocatalytic H₂ generation which is attributed to reduction in

^a Department of Chemical Engineering, Indian Institute of Technology Kanpur, Kanpur, UP-208016, India. E-mail: rpala@iitk.ac.in, srisiva@iitk.ac.in

^b Material Science Programme, Indian Institute of Technology Kanpur, UP-208016, India

^c Centre for Environmental Science & Engineering, Thematic Unit of Excellence on Soft Nanofabrication, Indian Institute of Technology Kanpur, Kanpur, UP-208016, India

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c5cp04212g



surface defects due to doping. The formation of surface defects can be further controlled by predoping of the non-metal during the synthesis of TiO₂ followed by controlled hydrogenation. In the proposed methodology, we have done non-metal (N and S) doping followed by hydrogenation. Since the dopants (N and S) chosen in this methodology are less electro-negative in comparison to oxygen, lower is the energy required for hydrogen to remove these dopants (than oxygen) and create/stabilize the anion vacancies. We optimized the hydrogenation temperature and time (800 °C and 24 h) such that hydrogen diffuses on the surface but not to the bulk (lower hydrogenation temperature and time result in the hydrogenation of the surface and higher hydrogenation temperature and time result in the hydrogenation of the bulk). In this methodology, anion vacancies are predominantly present on the surface rather than in the bulk which reduce the charge carrier recombination and the element doped in the bulk increases the absorption cross-section.

To this end, we report the synthesis of N-S co-doped TiO₂ followed by hydrogenation for enhanced PEC water splitting. In these doped hydrogenated TiO₂ nanostructures, interaction of trapped electrons in oxygen vacancies (V_O) present on the surface with adjacent Ti⁴⁺ centers leads to the formation of Ti³⁺ states. These electronic levels form below the conduction band and contribute to the enhanced optical absorption cross-section.⁹ Further, the electrochemical and photoelectrochemical properties of these doped hydrogenated TiO₂ nanostructures will be improved due to the following reasons: (1) the presence of under co-ordinated or strained bonds on the surface helps in increasing electrocatalytic properties^{10–13} and (2) the synergistic effect between the added dopant and hydrogen passed to the system will stabilize the oxygen vacancies and improve the performance.^{14–18}

In this work, we report (1) multi-element (N-S) doped hydrogenated TiO₂ nanostructures fabricated on a stainless steel (SS) substrate *via* “click” chemistry for enhanced PEC water splitting reaction and (2) the role of surface defects (oxygen vacancies/Ti³⁺ states) in the amorphous layer of hydrogenated TiO₂ material in charge-carrier transport. Nitrogen and sulphur are chosen as dopants as they are less electronegative than oxygen and expected to reduce the band-gap by introducing electronic states above valence band (VB) levels of TiO₂ upon substitutional doping. Furthermore, hydrogenation will lead to stabilization of Ti³⁺ states which is expected to lower the band gap through lowering of the conduction band and also help in achieving amorphization of the surface layer which will lead to increased electrocatalytic properties. From EC/PEC measurements, it is observed that doped hydrogenated TiO₂ nanostructures perform better than unhydrogenated undoped/doped and hydrogenated undoped TiO₂ nanostructures, indicating the importance of the combined effect of doping and hydrogenation. The calculated efficiency metrics of the hydrogenated N-S co-doped TiO₂@800 °C/24 h sample yields an applied bias photon to current conversion efficiency (ABPE) of 0.38%@0.63 V applied bias, an intrinsic solar to chemical conversion (ISTC) efficiency of 0.031 (or 3.1%)@1.27 V vs. RHE (reversible hydrogen electrode) and an electrical and solar power-to-hydrogen conversion (ESPH) efficiency of 6.6%@1.56 V applied bias, respectively, which is higher than undoped/doped

TiO₂ nanostructures hydrogenated at different temperatures and time periods. The enhancement in PEC performance in doped hydrogenated TiO₂ is attributed to the combined effect of doping, oxygen vacancy states, and Ti³⁺ states. We note that the 24 h hydrogenated sample possesses a higher bandgap (~2.95 eV), but shows enhanced efficiency in comparison to the 72 h hydrogenated sample (~2.94 eV). Though, there is no significant difference in the optical bandgap measured from Tauc plots of 24 h and 72 h hydrogenated samples, a noticeable difference in double layer capacitance (Table S6b, ESI†) is observed. This suggests the changes in the electrochemical surface area (ECSA) available for electrochemical reaction, which signifies the prominent role of surface electrochemistry in PEC water splitting reaction. Furthermore, the interfacial kinetics of surface disordered TiO₂ nanostructures calculated from electrochemical impedance spectroscopy (EIS) studies under light illumination also show that optical absorption and charge carrier separation are not mutually exclusive pathways in enhancing the PEC performance of surface disordered/hydrogenated TiO₂ and it is not enough to just increase optical absorption while trying to increase the PEC efficiency, but increase of PEC efficiency requires a finer balance of optimizing optical absorption without severely debilitating electron-hole separation. The role of the disordered surface generated due to hydrogenation in enhancing PEC activity is demonstrated by the removal of a few surface atomic layers from doped hydrogenated TiO₂ nanostructures by focused ion beam (FIB) sputtering upon which a decrease in both EC/PEC activity and photon absorption cross-section in comparison to the un-sputtered sample is observed. This result gives credence to the need for disordered surfaces in TiO₂ for obtaining enhanced PEC performance.

2. Experimental section

2.1 Synthesis of undoped and doped TiO₂ nanoparticles

Titanium(IV) iso-propoxide (TTIP, 9 ml) was added dropwise into a homogeneous mixture of ethanol (150 ml) and water (9.75 ml). The reaction mixture was kept at 85 °C for 4 h under agitation followed by evaporation of ethanol. The sample was dried at 60 °C for 30 min to get TiO₂ nanopowder which was finely ground and calcined at 400 °C for 3 h to yield the anatase TiO₂ phase.²¹ For the synthesis of doped TiO₂ nanoparticles, different precursor salts/solutions were used for different dopants. Nitrogen doping was achieved by addition of 5 ml of aq. NH₃ solution. Similarly, a solution of 9.3 g of thiourea in 3.75 ml of DI water was used as a precursor for sulphur doping. For N-S co-doped TiO₂ nanoparticle synthesis, a mixture of 2.5 ml of aq. NH₃ and 5.4 g of thiourea was used as precursor solution. The above-mentioned experimental protocol for undoped TiO₂ was followed except for the addition of the respective precursor salt/solution before the addition of TTIP to the ethanol-water mixture.

2.2 Hydrogenation of undoped/doped TiO₂ nanoparticles

Calcined samples (undoped and doped TiO₂ nanoparticles) were pelletized and hydrogenated at different temperatures in



a range of 500 °C to 800 °C for 24 h and 72 h at 20 bar pressure in the hydrogenation chamber. The digital photograph of the reactor used for the hydrogenation process is shown in Fig. S1a (ESI†). It is also important to eliminate the possibility of “blackness” in TiO₂ due to carbon intercalation, which is not easy to remove because of reducing hydrogenation conditions. In the present study, the possibility of carbon induced blackness was eliminated by a careful XPS study. We note that reaction conditions have to be maintained precisely to reproduce the results.

2.3 Fabrication of nanostructured TiO₂ electrodes

Scheme 1 shows a schematic representation of the synthesis of hydrogenated N-S co-doped TiO₂ nanostructures followed by clicking onto an alkyne functionalized stainless steel (SS) substrate using click chemistry. It involves four steps: (1) synthesis of doped TiO₂ nanostructures; (2) hydrogenation of doped TiO₂ nanostructures at various temperatures (500 °C, 650 °C, and 800 °C) and time periods (24 h and 72 h); (3) surface functionalization of TiO₂ nanostructures with the azide group (Section 1.2, ESI†); (4) surface functionalization of the SS substrate with the alkyne group (Section 1.3, ESI†) respectively; and (5) coating of TiO₂ nanostructures onto SS substrates *via* click chemistry (Section 1.4, ESI†).

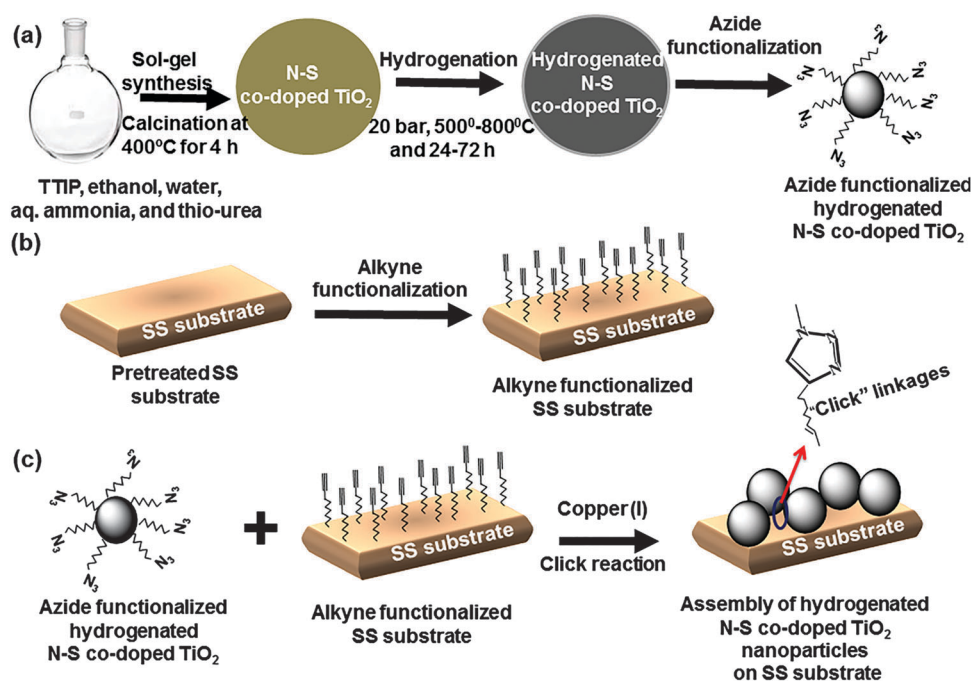
3. Results and discussion

3.1 Material characterization

3.1.1 XRD analysis. XRD data of hydrogenated doped TiO₂ nanostructures prepared at various temperatures and for different time durations are shown in Fig. 1a and Fig. S2 (ESI†).

XRD data of samples hydrogenated at 500 °C for 24 h (*i.e.* hydrogenated anatase TiO₂@500 °C/24 h, hydrogenated N-doped TiO₂@500 °C/24 h, hydrogenated S-doped TiO₂@500 °C/24 h, and hydrogenated N-S co-doped TiO₂@500 °C/24 h) confirm the presence of the anatase phase with JCPDS card no: 00-004-0477 (Fig. 1 and Fig. S2b, ESI†). Up to the hydrogenation temperature of 650 °C, there is no appearance of the rutile phase as evident from the hydrogenated N-S co-doped TiO₂@650 °C/24 h sample which shows a pure anatase phase (Fig. 1a). Further increase of hydrogenation temperature leads to the formation of the rutile phase as the major phase along with a small amount of anatase phase in hydrogenated N-doped TiO₂@700 °C/24 h (Fig. S2b and c, ESI†), hydrogenated N-S co-doped TiO₂@800 °C/24 h, and hydrogenated N-S co-doped TiO₂@800 °C/72 h samples, respectively, (Fig. 1). The formation of a mixed phase is further supported by HRTEM/SAED analysis (Fig. 3 and 4). The N-doped, S-doped and N-S co-doped TiO₂ nanostructures calcined at 400 °C possess crystallite sizes of ~11.5 nm, 7.8 nm, and 4.2 nm, respectively, which are less than undoped TiO₂ (~17.7 nm) (Fig. S2a, ESI†). The decrease in the crystallite size of single element (N or S) and multi-element (N-S co-doped) TiO₂ samples compared to un-doped TiO₂ suggests dopant incorporation in the crystalline matrix.^{19,20} Furthermore, broadening of diffraction peaks of doped (N or S or N-S) TiO₂ nanostructures in comparison with un-doped TiO₂ nanostructures supports the modification of crystalline lattice parameters upon doping.

3.1.2 Electron microscopy studies. To study the morphology and subsequent changes in undoped/doped TiO₂ nanostructures during the hydrogenation step, scanning electron microscopy (SEM) images were taken (Fig. 2, Fig. S3 and S4, ESI†). The N-S co-doped TiO₂ sample hydrogenated at 500 °C temperature for



Scheme 1 Schematic representation of the synthesis and fabrication of hydrogenated N-S co-doped TiO₂ nanostructures on an alkyne functionalized stainless steel (SS) substrate *via* click chemistry.



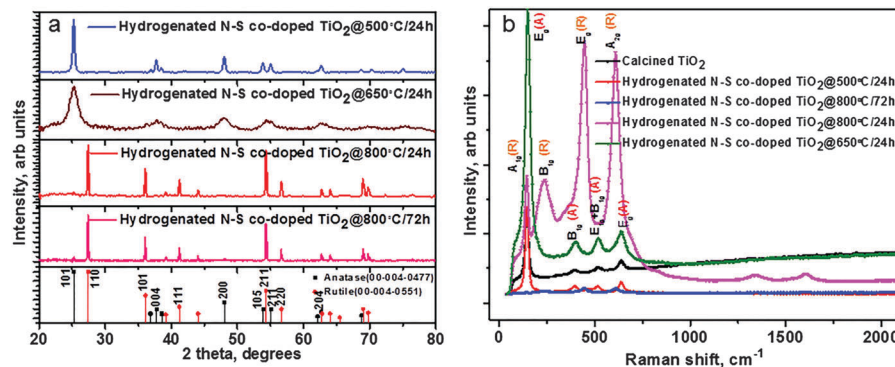


Fig. 1 (a) XRD patterns and (b) Raman spectra of undoped TiO_2 and hydrogenated N-S co-doped TiO_2 samples.

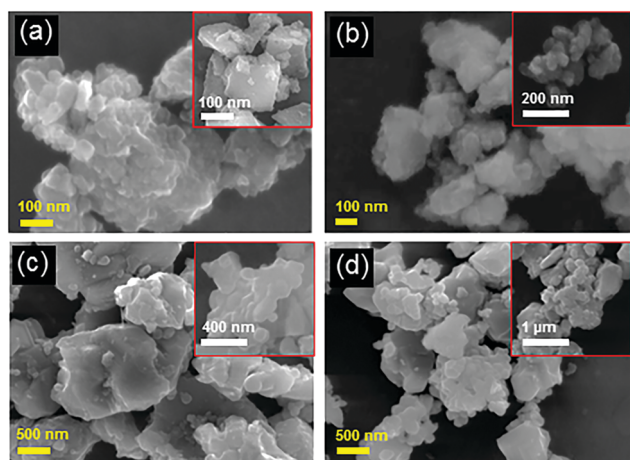


Fig. 2 SEM images of (a) hydrogenated N-S co-doped TiO_2 @500 °C/24 h; (b) hydrogenated N-S co-doped TiO_2 @650 °C/24 h; (c) hydrogenated N-S co-doped TiO_2 @800 °C/24 h; and (d) hydrogenated N-S co-doped TiO_2 @800 °C/72 h.

24 h shows irregular morphology with increased aggregation as compared to the non-hydrogenated sample (Fig. 2a). Furthermore, the same sample hydrogenated at elevated temperatures (650 °C and 800 °C) for 24 h/72 h shows larger clusters with irregular morphology (Fig. 2b–d). Therefore, it can be concluded that the surface morphology of multi-element doped TiO_2 nanostructures enormously changes with increase in hydrogenation temperature. The amorphous TiO_2 (uncalcined) and crystalline TiO_2 (calcined) nanoparticles show spherical morphology with the size ranging from 110–160 nm (Fig. S3a and b, ESI†). N-doped and S-doped TiO_2 nanoparticles also possess a spherical shape with the particle size of $\sim 100 \pm 10$ nm and 175 ± 10 nm, respectively, (Fig. S3c and d, ESI†). However, co-doping TiO_2 with nitrogen and sulphur leads to anisotropic morphology (Fig. S3e, ESI†). Moreover, undoped/single element-doped (N or S) TiO_2 samples hydrogenated at 500 °C for 24 h lead to an increase in the particle size compared to non-hydrogenated samples, while retaining spherical morphology (Fig. S4, ESI†). Additionally, the SEM image of N-doped TiO_2 clicked over the SS substrate (Fig. S4d, ESI†) shows good surface coverage. Furthermore, EDAX analysis (Fig. S3f and Table S1, ESI†)

confirms the presence of Ti and O elements along with dopants nitrogen (N) and sulphur (S). We believe that the photoefficiency can be further enhanced by the precise optimization of the dopant and hydrogenation conditions.

The morphology, crystallinity and formation of an amorphous layer upon hydrogenation have been further confirmed by transmission electron microscopy (TEM), high resolution transmission electron microscopy (HRTEM), and selective area electron diffraction (SAED) studies (Fig. 3 and 4, and Fig. S5, ESI†). TEM images of hydrogenated N-S co-doped TiO_2 @500 °C/24 h (Fig. 3a) and N-S co-doped TiO_2 @650 °C/24 h (Fig. 3c) show the particles with irregular morphology and the indexed SAED pattern (Fig. 3b and d) confirms the presence of the anatase phase which matches with XRD data (Fig. S2d, ESI†). Furthermore, N-S co-doped TiO_2 samples hydrogenated at 800 °C for 24 h and 72 h also show irregular morphology (Fig. 3e and g) but possess both anatase and rutile crystalline phases. This is further supported by the indexed SAED pattern (Fig. 3f and h) and XRD data (Fig. 1a). Fig. 4 shows the HRTEM information of N-doped TiO_2 hydrogenated at 700 °C for 24 h. The contrast in the TEM image clearly demarcates the rutile and anatase phases with lattice fringes (Fig. 4a). Furthermore, the HRTEM image (Fig. 4b) confirms the formation of a disordered/amorphous layer over the crystalline core with a thickness of ~ 0.5 nm, which indicates that hydrogenation has created the surface defects (Fig. 4b) in the TiO_2 material. The calculated d -spacing values of 0.1462, 0.2313 and 0.2145 nm correspond to the anatase phase and (130), (200), (111) planes correspond to rutile phases respectively (Fig. 4b). Furthermore, Fig. S5 (ESI†) shows the TEM images and the SAED pattern of crystalline TiO_2 and N-doped TiO_2 hydrogenated at 500 °C for 24 h. TEM images (Fig. S5a and c, ESI†) of both samples show particles with an irregular shape and SAED indexed patterns (Fig. S5b and d, ESI†) show the pure anatase phase.

3.1.3 Raman spectral analysis. Raman spectra of undoped, doped, and hydrogenated doped samples are shown in Fig. 1b and Fig. S6 (ESI†). Raman spectra of unhydrogenated undoped/doped samples (Fig. S6a, ESI†) depict four major intense peaks observed at 143 cm^{-1} , 395 cm^{-1} , 513 cm^{-1} and 639 cm^{-1} , attributed to the E_g , B_{1g} , $A_{1g} + B_{1g}$, and E_g bands of the anatase phase of TiO_2 respectively. Furthermore, hydrogenated TiO_2 @500 °C/24 h,



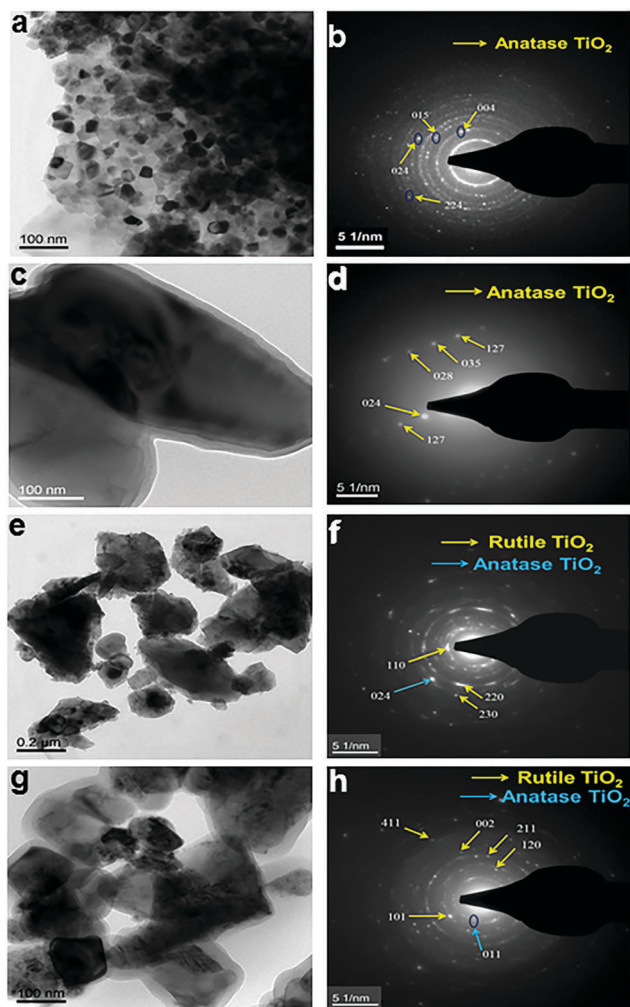


Fig. 3 TEM images of N-S co-doped TiO_2 hydrogenated at temperatures of (a) 500 °C; (c) 650 °C; and (e) 800 °C for 24 h; and (g) at 800 °C for 72 h; (b), (d), (f) and (h) show the corresponding SAED patterns.

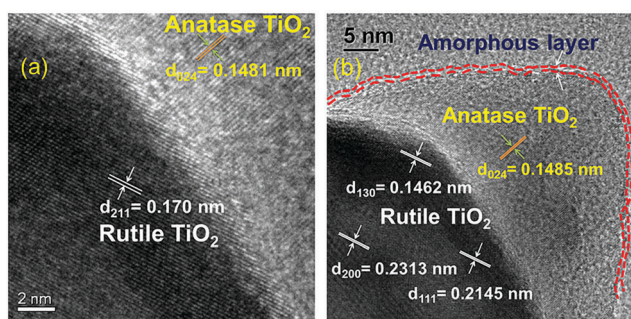


Fig. 4 HRTEM image of the hydrogenated N-doped TiO_2 at 700 °C/24 h sample showing (a) anatase–rutile phase interfaces and (b) the amorphous layer over the crystalline core.

hydrogenated N-doped TiO_2 @500 °C/24 h, hydrogenated S-doped TiO_2 @500 °C/24 h, hydrogenated N-S co-doped TiO_2 @500 °C/24 h, and hydrogenated N-S co-doped TiO_2 @650 °C/24 h possess anatase bands (Fig. S6b, ESI[†] and Fig. 1b) which confirm that the anatase crystalline phase is consistent with XRD data

(Fig. S2b, ESI[†]). However, N-S co-doped TiO_2 @800 °C/24 h and N-S co-doped TiO_2 @800 °C/72 h samples show three peaks located at $\sim 234 \text{ cm}^{-1}$, 443 cm^{-1} , and 610 cm^{-1} , attributed to B_{1g} , E_g , and A_{1g} bands of the rutile TiO_2 phase respectively, which is consistent with its XRD data (Fig. 1a). Furthermore, N-doped TiO_2 @700 °C/24 h possess both anatase and rutile bands indicating the formation of both anatase and rutile crystalline phases, which is consistent with its XRD data (Fig. S2b, ESI[†]). Additionally, a small shift along with broadening of the peak is observed at 141 cm^{-1} (E_g) (attributed to both anatase and rutile phases) with N-doped hydrogenated TiO_2 @500 °C/24 h and N-doped TiO_2 @700 °C/24 h nanostructures in comparison to N doped TiO_2 nanostructure (Fig. S6c, ESI[†]). Furthermore, the same observation is found for the hydrogenated N-S co-doped TiO_2 @650 °C/24 h sample (B_{1g}) as shown in Fig. S6d (ESI[†]). The shift and broadening of the peak indicate the non-stoichiometry or reduction in the original symmetry of the TiO_2 lattice due to surface disorder.²¹ Additionally, Raman spectra of the hydrogenated TiO_2 @500 °C/24 h sample show two peaks at 1339 cm^{-1} (D band) and 1606 cm^{-1} (G band), which are attributed to Ti–H vibrations or surface disorder defects created due to the presence of carbon, wherein D and G bands (which are due to the presence of carbon in the SS substrate) will also emerge at the same wavenumbers (Fig. S18a, ESI[†]).

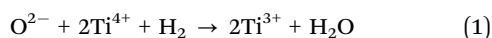
3.1.4 FTIR analysis. FTIR spectra of undoped/doped/hydrogenated TiO_2 nanostructures are shown in Fig. S7 and S8, and Table S2 (ESI[†]). FTIR spectra (Fig. S8b, ESI[†]) of hydrogenated TiO_2 nanostructures show reduction in –OH peak intensity in comparison to unhydrogenated nanostructures suggesting the reduction of –OH groups present in the TiO_2 sample. The presence of dopants N and S in the TiO_2 matrix is confirmed by Ti–S and Ti–N characteristic peaks as shown in Fig. S7c and d and Table S2 (ESI[†]). Furthermore, the peaks at ~ 3645 , 3670 , and 3685 cm^{-1} of the hydrogenated S-doped TiO_2 @500 °C/24 h sample are attributed to the tetrahedral co-ordinated vacancies (Fig. S8a, ESI[†]).²² FTIR spectra of azide functionalized undoped/doped and doped-hydrogenated TiO_2 particles show a peak at $\sim 2110 \text{ cm}^{-1}$ (Fig. S7b, ESI[†]) which confirms the presence of the azide group on particles. Additionally, FTIR spectra of the alkyne functionalized SS substrate show a peak at $\sim 2156 \text{ cm}^{-1}$ confirming the presence of an alkyne group (Fig. S8b, ESI[†]). All functionalized TiO_2 nanostructures show the presence of strong –CH₂ stretching vibrations at $\sim 2950 \text{ cm}^{-1}$ confirming the surface modification. Furthermore, an increase in the contact angle ($\sim 97.50^\circ$) in the case of the alkyne functionalized electrode (hydrophobic) compared to the bare SS substrate/electrode ($\sim 60^\circ$) (Fig. S8c, ESI[†]), suggests the presence of a terminal alkyne moiety on the surface which matches with our earlier report.²³

3.1.5 UV-Vis spectral analysis. UV-Vis spectra of undoped/doped and hydrogenated doped TiO_2 samples are shown in Fig. S9 (ESI[†]). The approximate band gaps of all TiO_2 nanostructures calculated from Tauc plots (undoped/doped/hydrogenated doped TiO_2) are summarized in Table S3 (ESI[†]). The decrement in the optical bandgap of doped and hydrogenated samples is clearly visualized with the color of samples (Fig. S1b–n, ESI[†]).



Both amorphous (uncalcined) and crystalline (calcined) TiO₂ samples show absorption bands at around 387 nm and 392 nm (Fig. S9a, ESI†) which correspond to the bandgaps of 3.66 eV and 3.46 eV, respectively, measured from the corresponding Tauc plots (Fig. S9b, ESI†). A red-shift in the absorption band is observed when the TiO₂ matrix is doped. The measured band gaps for N-doped TiO₂, S-doped TiO₂ and N-S co-doped TiO₂ are found to be ~3.38 eV, 3.37 eV, and 3.09 eV respectively. The S dopant due to its lower electronegativity can shift the VB to more upward direction in comparison to the N dopant, although both seem to possess same absorption cross-sections. Furthermore, N-doping leads to the stabilization of Ti³⁺ states.^{14,24} For this reason, co-doping leads to lowering of the band-gap more than individual doping (S or N). Furthermore, absorption spectra of hydrogenated TiO₂ nanostructures (Fig. S9c, ESI†) show a red-shift with less optical band gap compared to their unhydrogenated counterparts. This is attributed to the formation of mid-gap states between the CB and the VB due to hydrogenation and the formation of oxygen vacancy and Ti³⁺ states.

3.1.6 Electron paramagnetic resonance (EPR) analysis. Ti³⁺ surface defects are quite different from oxygen vacancies (V_O).¹⁹ It is to be noted that Ti³⁺ states will be generally formed by the reduction of Ti⁴⁺ in either of the two ways: (1) the Ti⁴⁺ ion receives photoelectrons usually by UV light absorption on the TiO₂ surface. The electrons can be trapped and tend to reduce Ti⁴⁺ cations to Ti³⁺ states.^{19,20} (2) Ti⁴⁺ reduction to Ti³⁺ is usually accompanied by loss of oxygen from TiO₂ surface. Similarly, oxygen vacancies (V_O) are responsible for the formation of color (F) centers, as the electron pair remains trapped in the cavity (V_O), upon loss of an oxygen atom from the surface of metal oxide.²⁵ Electron paramagnetic resonance (EPR) spectra are used to confirm the presence of surface defects (oxygen vacancies/Ti³⁺ states) in the surface disordered/hydrogenated TiO₂ sample. It is known that for paramagnetic materials containing “F-centers” or oxygen vacancies, the resonance at $g = 2.004$ is the most prominent and characteristic electron signal.²⁶ Undoped TiO₂ nanostructures show the g -values of ~2.003–2.004 (Fig. S10a, ESI†), which is attributed to the existence of oxygen vacancies. Furthermore, N-S co-doped TiO₂ samples hydrogenated at 500 °C, 650 °C, and 800 °C for 24 h result in additional peaks at $g \sim 1.996$ to 1.999 along with a small shift in the original oxygen vacancy state ($g \sim 2.003$ –2.007) (Fig. S10b, e, and f, ESI†). The peak that emerged at $g = 1.996$ to 1.999 in these samples is attributed to Ti³⁺ states due to the interaction of H atoms and the Ti⁴⁺ state of TiO₂ leading to the formation of Ti³⁺ states (eqn (1)).



This observation clearly confirms the presence of oxygen vacancies and Ti³⁺ states upon hydrogenation which is consistent with earlier reports.^{21,27} Furthermore, the doped unhydrogenated TiO₂ sample (N-S co-doped TiO₂) also shows a characteristic signal closer to the g value of 1.997 indicating the presence of Ti³⁺ states (Fig. S10c and d, ESI†). Therefore, Ti³⁺ states are present along with oxygen vacancies before hydrogenation (*i.e.* in doping stage) which is consistent with previous reports.²⁸

Although, it is highly difficult to demarcate the presence and the individual role of oxygen vacancies relative to Ti³⁺ states in the enhanced PEC performance of doped and hydrogenated TiO₂ samples,²⁹ we attempted to find out the relative change in the concentration of oxygen vacancies from Mott-Schottky plots (Fig. S11, ESI†).^{28,30} Since oxygen vacancies act as electron donors, an increase in the donor density can be attributed to the increase in the concentration of oxygen vacancies. The donor density can be calculated from the relation

$$N_d = \left(\frac{2}{e_0 \epsilon \epsilon_0} \right) \left(\frac{d \frac{1}{C^2}}{dV} \right)^{-1} \quad (2)$$

where N_d is the donor density, e_0 is the electron charge, C is the capacitance, ϵ is the dielectric constant of TiO₂ (~31 in the case of anatase TiO₂),³¹ ϵ_0 is the permittivity of vacuum, and V is the potential applied at the electrode. The calculated slopes from Mott-Schottky plots, for anatase TiO₂, hydrogenated N-S co-doped TiO₂@800 °C/24 h, and hydrogenated N-S co-doped TiO₂@800 °C/72 h are found to be $6.87 \times 10^{10} \text{ cm}^4 \text{ F}^{-2} \text{ V}^{-1}$, $4.44 \times 10^{10} \text{ cm}^4 \text{ F}^{-2} \text{ V}^{-1}$, and $4.34 \times 10^{10} \text{ cm}^4 \text{ F}^{-2} \text{ V}^{-1}$ respectively. Furthermore, the donor densities of the anatase TiO₂, hydrogenated N-S co-doped TiO₂@800 °C/24 h, and hydrogenated N-S co-doped TiO₂@800 °C/72 h are calculated to be 0.66×10^{20} , 1.02×10^{20} , and $1.05 \times 10^{20} \text{ cm}^{-3}$ respectively. An increase in the donor density of hydrogenated N-S co-doped TiO₂@800 °C/72 h and hydrogenated N-S co-doped TiO₂@800 °C/24 h samples with respect to anatase TiO₂ suggests that there is an increment in the concentration of oxygen vacancies in doped hydrogenated samples (hydrogenated N-S co-doped TiO₂@800 °C/72 h, and hydrogenated N-S co-doped TiO₂@800 °C/24 h) over the undoped unhydrogenated sample (anatase TiO₂).

3.1.7 X-ray photoelectron spectroscopy (XPS) analysis. X-ray photoelectron spectroscopy (XPS) reveals the presence of elements and the associated chemical bonds on the surface within few atomic layers of the material. Two peaks of Ti 2p_{3/2} (456.8 eV) and Ti 2p_{1/2} (462.7 eV) are observed for hydrogenated N-S co-doped TiO₂@650 °C/24 h nanostructures (Fig. S12, ESI†), which confirms the presence of Ti³⁺ states (Ti⁴⁺ + e⁻ → Ti³⁺). The observed binding energies of 284.05 eV, 162.5 eV, 232.6 eV, and 397.7 eV represent C 1s, S 2p, S 2s, and N 1s signals, respectively (Fig. S12a, ESI†). Furthermore, XPS measurements were taken at slow scan rates to magnify the Ti and O peaks. Two nearby peaks are observed at binding energies of ~55.1 eV and 61.6 eV which correspond to Ti 3s states. Additionally, the O 1s peak can be resolved into two peaks at about 530.39 eV and 531.2 eV, wherein the peak at 530.39 eV is attributed to Ti–O–Ti groups on the surface and the broader peak at 531.2 eV can be attributed to Ti–OH (Fig. S12b and c, ESI†). These observations confirm the presence of dopants as well as the defect states in the TiO₂ matrix. Furthermore, deconvoluted XPS spectra of hydrogenated N-S co-doped TiO₂@800 °C/24 h nanostructures are shown in Fig. 5. Four peaks at 557.7 eV, 558.6 eV, 559.3 eV, and 564.6 eV were observed from deconvoluted XPS Ti 2p spectra. The peaks



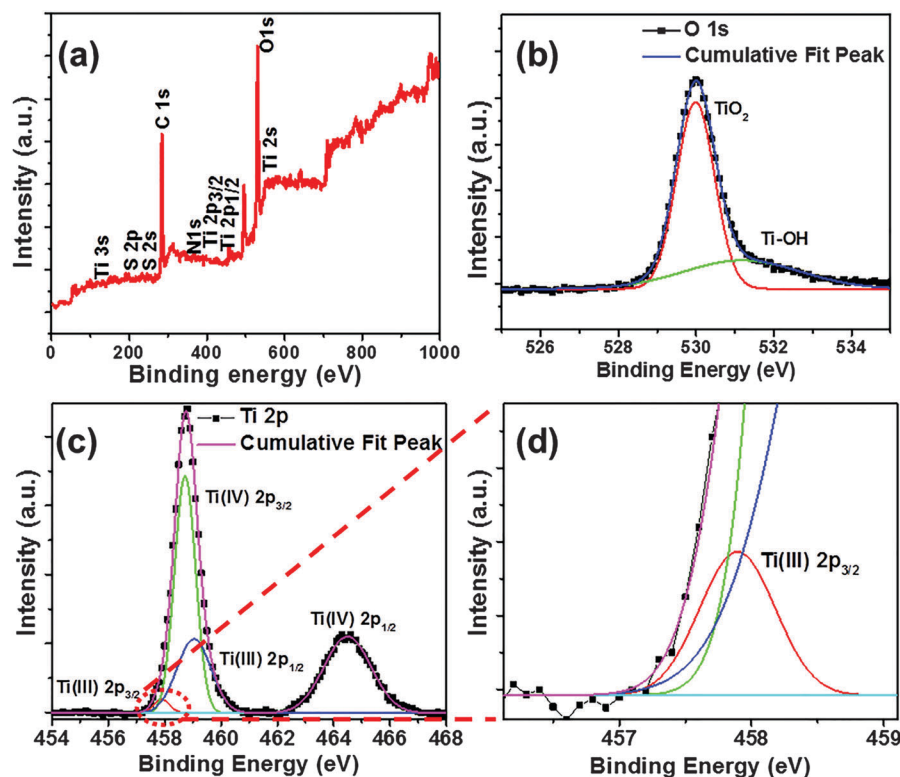


Fig. 5 XPS spectra of hydrogenated N-S co-doped $\text{TiO}_2@800\text{ }^\circ\text{C}/24\text{ h}$: (a) whole scan, (b) O 1s, (c) Ti 2p, and (d) the enlarged/resolved portion (red dotted lines of c) of the Ti 2p XPS peak.

at 558.6 eV and 564.6 eV can be attributed to $2p_{3/2}$ and $2p_{1/2}$ of Ti^{4+} ion states (Fig. 5c), and peaks at 557.7 eV (Fig. 5d) and 559.3 eV (Fig. 5c) can be attributed to $2p_{3/2}$ and $2p_{1/2}$ of Ti^{3+} states, respectively, which confirm the presence of Ti^{3+} states and defect states in the TiO_2 matrix (Fig. 5d) consistent with earlier reports.^{32–34}

We have performed XPS measurements to determine the valence band edge position of anatase TiO_2 and hydrogenated N-S co-doped $\text{TiO}_2@800\text{ }^\circ\text{C}/24\text{ h}$ samples (Fig. S13, ESI†). The calculated valence band edge position of anatase TiO_2 is about 1.29 eV, whereas it is 0.86 eV in the case of the hydrogenated N-S co-doped $\text{TiO}_2@800\text{ }^\circ\text{C}/24\text{ h}$ sample. This confirms the shift of the valence band edge position of the hydrogenated N-S co-doped $\text{TiO}_2@800\text{ }^\circ\text{C}/24\text{ h}$ sample by 0.43 eV towards vacuum. Since the difference in the optical bandgap between anatase TiO_2 and N-S co-doped $\text{TiO}_2@800\text{ }^\circ\text{C}/24\text{ h}$ is about 0.51 eV, the rest of the fraction can be considered as a shift in the conduction band edge towards the valence band.

3.2 Electrochemical (EC)/photoelectrochemical (PEC) water splitting

The performance of fabricated nanostructures towards photoelectrochemical (PEC) water splitting is evaluated using both two-probe and three-probe methods.³⁵ Current vs. applied bias/applied potential plots were recorded for all fabricated undoped/doped/hydrogenated doped TiO_2 nanostructured electrodes in 2-probe (Fig. 6, Fig. S14, S15 and Table S4, ESI†) and 3-probe modes (Table S5, ESI†) in 0.1 M NaOH solution

(pH ~ 13.6). The observed onset working electrode potential of fabricated nanostructured electrodes is 0.52–0.6 V vs. Ag/AgCl (1.52–1.6 V vs. RHE) which corresponds to the oxygen evolution reaction (OER) potential. The best performance among all the prepared hydrogenated doped TiO_2 samples is found to be for the hydrogenated N-S co-doped $\text{TiO}_2@800\text{ }^\circ\text{C}/24\text{ h}$ sample with a current density of 6.3 mA cm^{-2} (under light) and 4.7 mA cm^{-2} (under dark) at 1.6 V applied bias, respectively, as measured in 2-probe mode (Fig. 6d). Furthermore, the same electrode shows current densities of 31.4 mA cm^{-2} (under light) and 25.6 mA cm^{-2} (under dark) in three-probe mode at a working electrode potential of 1.2 V vs. Ag/AgCl (2.2 V vs. RHE) applied potential (Table S5, ESI†). Two-probe measurements of undoped/doped/hydrogenated doped TiO_2 nanostructured electrodes are shown in Fig. S14 and S15 (ESI†).

From Table S5 (ESI†), it is observed that hydrogenated TiO_2 nanostructures show enhanced PEC activity in comparison to unhydrogenated TiO_2 nanostructures. From the table, N doping in TiO_2 improves EC activity but S doping does not. This can be rationalized from the fact that aliovalent N-doping leads to the formation of the Ti^{3+} state which helps in increasing EC activity. However, isovalent S-doping does not lead to the formation of any intermediate Ti^{3+} state, which contributes towards improvement in EC activity. Furthermore, due to less electronegativity of S, there would be less charge polarization over the Ti-atom which slightly lowers EC activity towards water splitting as evident from the lower EC activity of S-doped TiO_2 in comparison to TiO_2 . However, co-doped TiO_2 shows higher EC activity than both



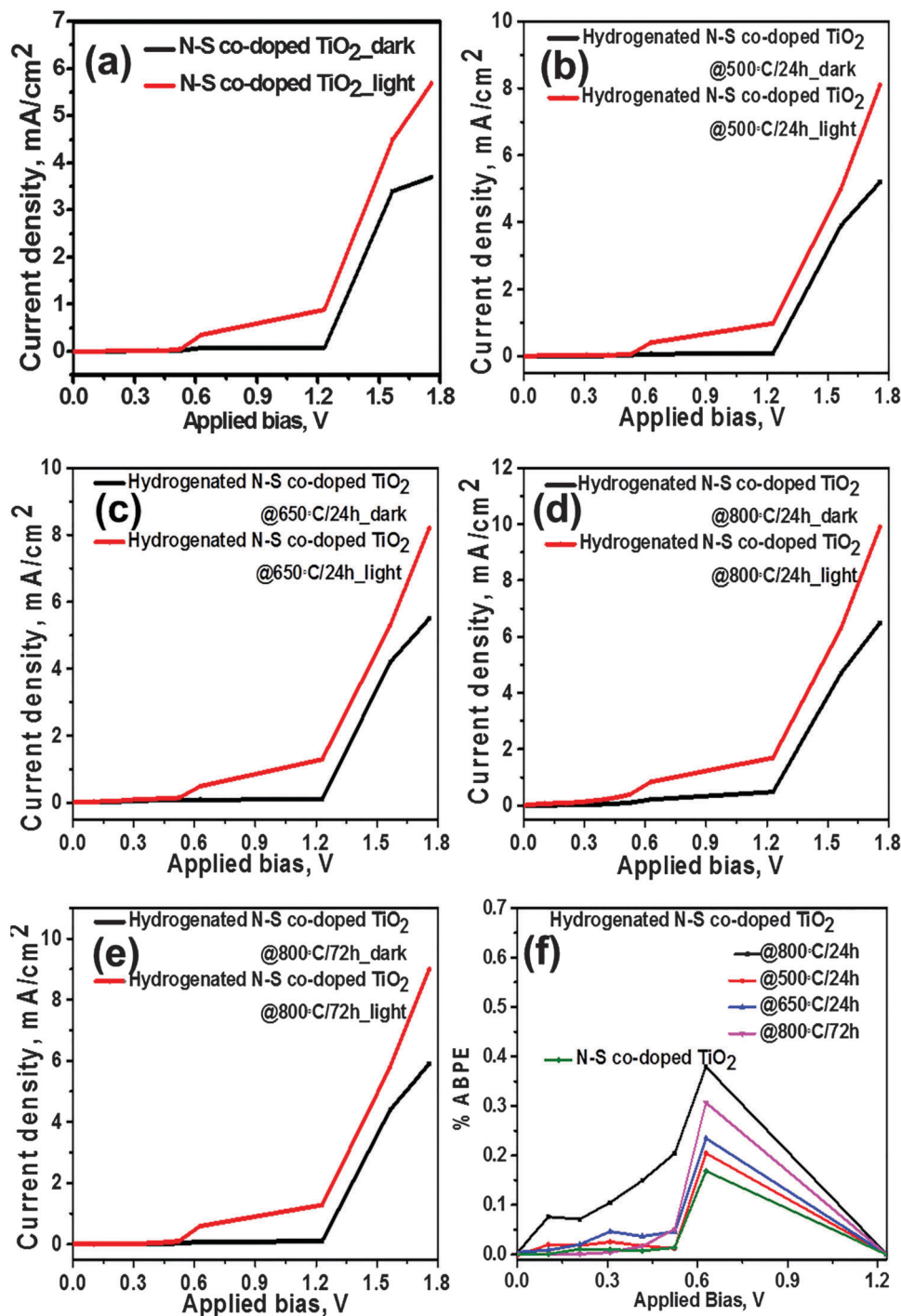


Fig. 6 Current vs. applied bias plots of (a) N-S co-doped TiO₂; (b) hydrogenated N-S co-doped TiO₂@500 °C/24 h; (c) hydrogenated N-S co-doped TiO₂@650 °C/24 h; (d) hydrogenated N-S co-doped TiO₂@800 °C/24 h; (e) hydrogenated N-S co-doped TiO₂@800 °C/72 h, and (f) the applied bias to photo-conversion efficiency (ABPE) of all N-S co-doped hydrogenated TiO₂ samples performed in 0.1 M NaOH solution (two-probe measurements).

monodoped N- and S-doped TiO₂ due to either a synergistic interplay of both dopants in enhancing the activity and/or improved electronic conductivity.^{36–38} It is also observed that all the hydrogenated samples show better EC activity in comparison to their unhydrogenated counterparts. This result implies that amorphization due to the hydrogenation process improves EC activity towards water splitting. However, EC activity of the

hydrogenated samples significantly depends on the hydrogenation temperature which determines the crystalline phase. Among the co-doped hydrogenated TiO₂ samples, the EC activity is found to be in the order of hydrogenated N-S co-doped TiO₂@800 °C/24 h > hydrogenated N-S co-doped TiO₂@800 °C/72 h > hydrogenated N-S co-doped TiO₂@650 °C/24 h > hydrogenated N-S co-doped TiO₂@500 °C/24 h. The increased EC



activity of hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ may be attributed to the synergistic effect of rutile and a small amount of anatase present in the sample. In this regard, it is pointed out that a mixed phase shows better EC activity in comparison to any single phase.^{39,40} However, an increase of the rutile phase beyond a percentage leads to a reduction in EC activity.⁴¹ Furthermore, the decrease in the dark current of the hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/72\text{ h}$ sample in comparison to the hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ sample (Table S4, ESI†) may be attributed to the decrease in the electrochemical surface area (ECSA) upon hydrogenation for longer time (72 h), which is supported by a decrease in the double layer capacitance (C_{dl}) calculated from EIS circuits (Table S6b, Section 3.3, ESI†).

In order to compare the photo-performance of the electrodes, the applied bias photon to current conversion efficiency (ABPE) has been calculated using the following relation,

$$\text{ABPE} = \left[\frac{|j_{ph}(\text{mA cm}^{-2})| \times ((1.23 - |V_b|)(\text{V}))}{P_{\text{total}}(\text{mW cm}^{-2})} \right]_{\text{AM 1.5G}} \quad (3)$$

where j_{ph} is the photocurrent, V_b is the applied bias between working and counter electrodes, and P_{total} is the intensity of light (power) from a solar simulator. Since the photocurrent developed at low applied bias values play an important role in characterizing any photo material, we have taken the true activity of the photomaterial by measuring the applied bias to photo current conversion efficiency (ABPE), which signifies the minimum applied bias (below thermodynamic potential of water *i.e.* 1.23 V), at which the maximum photon to current conversion is achieved.^{42–44} The measured ABPE values of unhydrogenated and hydrogenated samples prepared at different temperatures and time periods are compared at 0.63 V and tabulated in Table S3 (ESI†). Among all the prepared electrodes, the best ABPE value is found for the hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ sample with a value of $\sim 0.38\%$, at 0.63 V applied bias (Fig. 6f). Therefore, the sample with the best photo activity performance in this work is hydrogenated N-S doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ and is chosen based on the true ABPE value. We conclude from above values that hydrogenated undoped/doped TiO_2 nanostructures show better PEC performance than unhydrogenated nanostructures and we can attribute the enhanced performance to reduction in charge carrier recombination as well as an improvement in optical absorption with the proposed design strategy. From Table S3 (ESI†), it is found that both N-doping and S-doping yield an enhancement of ABPE which can be ascribed to the enhancement of the photon absorption cross-section. Co-doping leads to a better ABPE value than any singular N- or S-doping (Fig. S16, ESI†). The order of ABPE among co-doped hydrogenated TiO_2 samples is found to be N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h} >$ hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/72\text{ h} >$ hydrogenated N-S co-doped $\text{TiO}_2@650^\circ\text{C}/24\text{ h} >$ hydrogenated N-S co-doped $\text{TiO}_2@500^\circ\text{C}/24\text{ h}$. The PEC activity in the disordered mixed phase is higher in comparison to the anatase phase.^{40,41,45} It is evident from the above results that hydrogenation at higher temperature enhances

the efficiency. It is to be noted that the 24 h hydrogenated sample with an optical bandgap of $\sim 2.95\text{ eV}$ shows enhanced efficiency in comparison to the 72 h hydrogenated sample having a bandgap of $\sim 2.94\text{ eV}$. The reduction in the PEC performance of the 72 h hydrogenated sample over the 24 h hydrogenated sample can be attributed to (1) the decrease in the electrochemical surface area (ECSA) (Table S6b, ESI†) as explained in Section 3.3, and (2) high charge carrier recombination that persists due to defect states/oxygen vacancies developed in bulk.

The PEC performance of the electrodes is also benchmarked by comparing the intrinsic solar to chemical conversion (ISTC) efficiency. This efficiency is calculated using three-electrode cell potentials (Fig. 7) which measure the maximum intrinsic power of the photo-anode to convert the solar light into chemical energy as per below relation,

$$\text{ISTC} = \frac{\eta_F \times 1.23(\text{V})}{V_{\text{dark}}(\text{V})} \times \frac{|j_{ph}(\text{mA cm}^{-2})| \times V_{\text{photo}}(\text{V})}{P_{\text{solar}}(\text{mW cm}^{-2})} \quad (4)$$

where photocurrent (j_{photo} , difference in current density obtained with and without light at the same potential) is given in mA cm^{-2} , on the other hand, photovoltage *i.e.* V_{photo} is the difference in potentials under light (V_{light}) and dark (V_{dark}) conditions at same current density. Furthermore, the maximum internal photovoltaic power of TiO_2 nanostructures is calculated from the plot between j_{photo} and V_{photo} as shown in Fig. 7b. For instance, under light illumination hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ shows the photocurrent of $\sim 0.58\text{ mA cm}^{-2}$ at 1.45 V (*vs.* RHE) whereas to obtain same current under dark conditions, this sample requires 1.48 V (*vs.* RHE) which means that solar light provides energy that corresponds to the potential of 0.036 V or power of $\sim 0.021\text{ mW cm}^{-2}$. As mentioned by Grätzel and co-workers, external power is saved as a result of the internal photovoltaic power of this sample which is calculated by the reduction in the water electrolysis conversion efficiency *i.e.* $1.23(V_{\text{RHE}})/V_{\text{dark,RHE}} = \sim 82\%$. Thus, the ISTC efficiency of this sample is $\sim 0.82 \times 0.021 = 0.0172$ or 1.72% (Fig. 7a). The maximum ISTC efficiency of $\sim 3.1\%$ is obtained for hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ compared to other hydrogenated samples. The significance of ISTC can be further understood as follows; the maximum current density obtained at 1.45 V (*vs.* RHE) with hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ is 0.775 mA cm^{-2} which means that this electrode can produce the maximum chemical power of $\sim 0.953\text{ mW cm}^{-2}$ and in this total power output, the internal solar conversion to the chemical power of hydrogenated N-S co-doped $\text{TiO}_2@800^\circ\text{C}/24\text{ h}$ contributes ~ 0.0172 *i.e.* 1.8% and the remaining power could be supplied from the external electrical energy. ISTC provides only half cell efficiency (as it is measured in three-electrode configuration) and does not include losses from other components of the PEC cell *i.e.* the counter-electrode, electrolyte, wires *etc.* Furthermore, both the efficiencies *i.e.* ABPE and ISTC are mainly used to evaluate the performance of either PEC device or photoanode based on the photocurrent and photovoltage. However, none of the above efficiencies provides the performance evaluation based on the total power output generated from the PEC device which is based on the total current



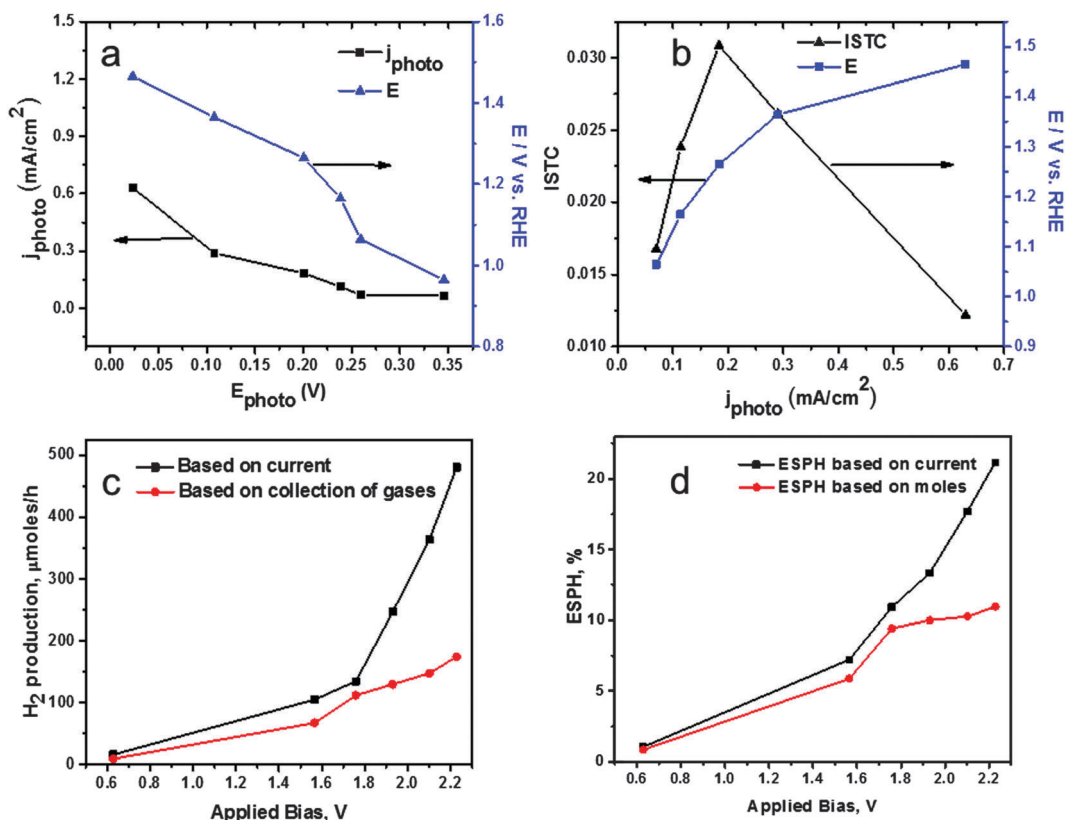


Fig. 7 (a) Intrinsic photovoltaic power characteristics; (b) intrinsic solar to chemical conversion (ISTC) efficiency; (c) micromoles of hydrogen production; and (d) the electrical and solar power-to-hydrogen (ESPH) conversion efficiency based on current and moles of hydrogen collected for the hydrogenated N-S co-doped TiO_2 @800 °C/24 h sample.

density obtained due to contribution from applied bias and light. Hence, in order to do so, we propose a new efficiency metric called “electrical and solar power-to-hydrogen (ESPH)” which is based on the definition of efficiency *i.e.* total power output divided by total power input.

We measured ESPH for the hydrogenated N-S co-doped TiO_2 @800 °C/24 h electrode, which includes the external applied electrical power, along with solar power as the total power input.

ESPH (based on collected H_2 gas)

$$= \left[\frac{(\text{mmoles of } \text{H}_2/\text{s}) \times (237 \text{ kJ mol}^{-1})}{(V_b \times I_{\text{dark}} + P_{\text{total}})(\text{mW cm}^{-2}) \times \text{Area}(\text{cm}^2)} \right]_{\text{AM1.5G}} \times 100 \quad (5)$$

ESPH (based on current)

$$= \left[\frac{|j_{\text{PEC}}(\text{mA cm}^{-2})| \times 1.23 \text{ (V)}}{(V_b \times I_{\text{dark}} + P_{\text{total}})(\text{mW cm}^{-2}) \times \text{Area}(\text{cm}^2)} \right]_{\text{AM1.5G}} \quad (6)$$

During the water splitting reaction, oxygen and hydrogen gases were collected at working and counter-electrodes respectively. Evolved gases were observed to be in a stoichiometric ratio which suggests that there is no side product formation. The sample shows ESPH efficiencies of ~1.18% and 7.2% at 0.63 and 1.56 V applied bias, respectively, which is the highest ESPH efficiency

among all other samples (Fig. 7d). The possible reasons for lower ESPH efficiency are: (1) it is measured in a two-electrode configuration and gives full cell efficiency and (2) it includes all the parasitic losses that arise from different components of the PEC cell. Fig. 7c shows a plot of hydrogen collected during reaction *vs.* applied bias for the same sample under light illumination conditions. Furthermore, ESPH efficiency (based on moles of H_2 production) under illumination (@ ~1.56 V bias) for this sample is ~5.83%. Additionally, we calculated the amount of hydrogen generated based on current density as well as based on collection as shown in Fig. 7c. The difference in ESPH as well as production of H_2 based on current and moles of hydrogen collection can be attributed to the ohmic drop (IR drop) and other associated resistances during the collection of evolved gases using the “inverted beaker method” (Fig. S19, ESI†).

3.3 Electrochemical impedance spectra (EIS) analysis

The electrode/electrolyte interface kinetics and dynamics of electron transport of fabricated nanostructures can be analyzed using EIS analysis.^{46,47} We performed EIS measurements to investigate the conductivity and impedance values of the best sample “hydrogenated N-S co-doped TiO_2 @800 °C/24 h” on SS substrate under dark and light conditions (Fig. S17, ESI†); The experimental and simulated patterns match with the $R_s(R_1C)(R_2Q)$ circuit (inset image of Fig. S17a, ESI†), where R_s represents the



solution resistance of the electrolyte (*i.e.* ohmic resistance including the resistance of the electrolyte, separator and electrical contacts), C and Q being capacitance and constant phase element, and R_1 and R_2 being the resistances associated with the charge transfer process. The two semicircles correspond to the oxide/hydroxide covered mixed surface.⁴⁸ The double layer charge capacitance (C_{dl}) vs. applied potential data for this sample is given in Table S6a (ESI[†]). Furthermore, a decrease in charge transfer resistance upon illumination is observed from a decrease in the diameter of the semicircle as shown in Fig. S17b (ESI[†]), which signifies the enhancement in the charge transport rate due to an increase in the photon absorption cross-section. This is because the illuminated light source provides additional electrochemical potential which helps in overcoming the activation energy for electrochemical reaction. Furthermore, EIS measurements were performed on three control samples (anatase TiO₂, hydrogenated N-S co-doped TiO₂@800 °C/24 h, and hydrogenated N-S co-doped TiO₂@800 °C/72 h) in Faradaic zone (0.5 V, 0.6 V, 0.7 V, and 0.8 V vs. Ag/AgCl) under light conditions (Fig. 8) to trace the charge transfer kinetics which is informative of the charge carrier recombination rate (films on fluorine doped tin oxide FTO substrate). The inset image in Fig. 8(d) shows the equivalent circuit (($R_s(R_f(R_{ct}C_{\mu}))C_{dl}$)) which is chosen for fitting the Nyquist data proposed by Bisquert *et al.*^{49,50} for three control samples. The reason for choosing this circuit is, after illumination of light on the photoanode, the electron-hole pair gets generated. Some of the excited electrons in the conduction band

form the double layer capacitance (C_{dl}) with ions in the electrolyte and the rest of the electrons recombine with holes in the valence band (which is denoted film resistance (R_f)). Since C_{dl} and R_f are mutually inclusive events, these two are kept in parallel. Furthermore, the fraction of holes which are left after recombination with electrons is transported from the surface to species to oxidize them (R_{ct} , which is the measure of charge transport from the valence band of the semiconductor to the electrolyte (*i.e.* less the charge transfer resistance, more the activity towards oxidation/reduction of the species)) and the rest of the fraction forms the chemical capacitance (C_{μ} , which denotes the concentration of holes in the valence band). Since R_{ct} and C_{μ} are inclusive events and series events with charge carrier recombination (R_f), R_{ct} and C_{μ} are in parallel and these are in series with R_f . After the inclusion of solution resistance (R_s) into the circuit, which is a series event with all of the above-mentioned systems, the overall EIS circuit is given by ($R_s(R_f(R_{ct}C_{\mu}))C_{dl}$).

From the fitted data of the EIS circuits, we infer that there is a decrease in the resistance of film (which is a measure of charge carrier recombination) for the hydrogenated N-S co-doped TiO₂@800 °C/24 h sample over anatase TiO₂ and hydrogenated N-S co-doped TiO₂@800 °C/72 h samples at 0.5, 0.6, 0.7 and 0.8 V potentials (Fig. 8 and Table S6b, ESI[†]). This result confirms the hydrogenated N-S co-doped TiO₂@800 °C/24 h sample having faster charge carrier separation in comparison to other two samples.⁵¹ Therefore, improvement in optical absorption does not lead to an enhancement of the photo performance

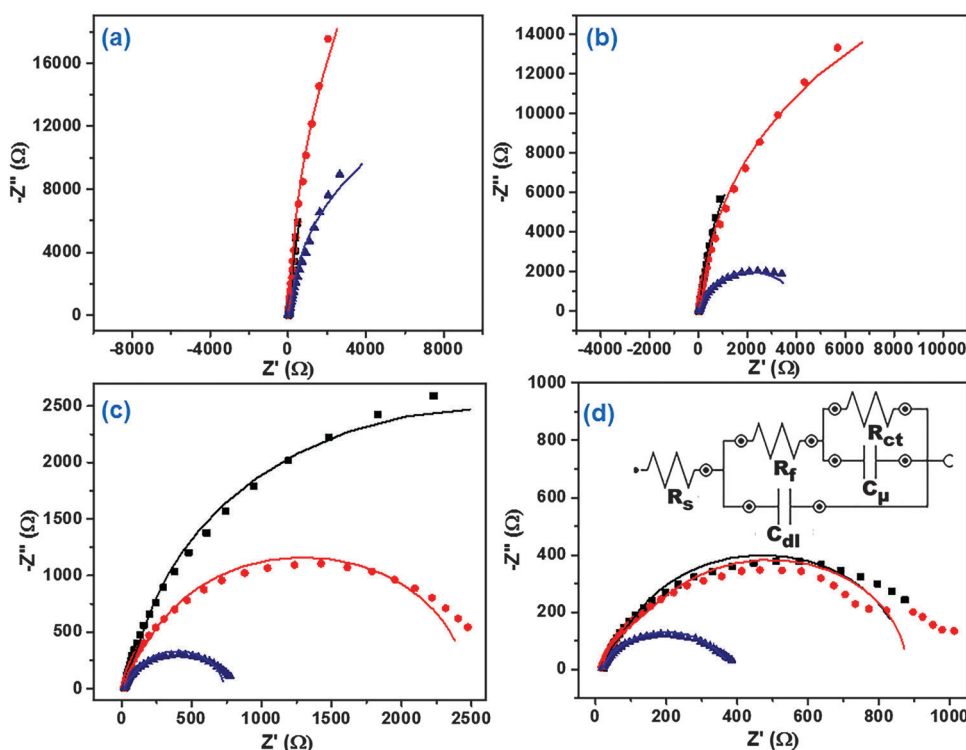


Fig. 8 Electrochemical impedance spectra (EIS) of anatase TiO₂ (black square), hydrogenated N-S co-doped TiO₂@800 °C/24 h (blue triangle) and hydrogenated N-S co-doped TiO₂@800 °C/72 h (red circle) under light conditions at different potentials (vs. Ag/AgCl): (a) 0.5 V, (b) 0.6 V, (c) 0.7 V, and (d) 0.8 V (symbols represent the experimental data, lines represent the fitted data, and the inset in figure d represents the EIS circuit).



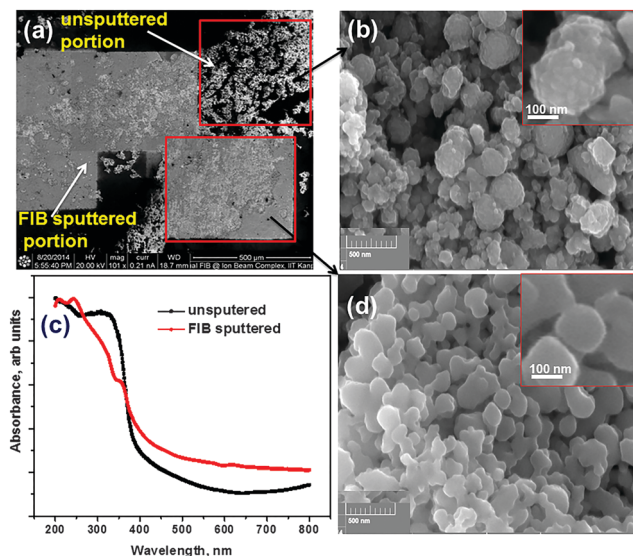


Fig. 9 SEM images of hydrogenated N doped $\text{TiO}_2@500^\circ\text{C}/24\text{ h}$: (a) at normal resolution, (b) the unspun portion, (d) the FIB sputtered portion; and (c) UV-Vis spectra of FIB sputtered and unspun portions of the hydrogenated N doped $\text{TiO}_2@500^\circ\text{C}/24\text{ h}$ sample.

as the charge carrier recombination step is dominating. It is to be noted that the measure of double layer capacitance is an indication of the active surface available for the electrochemical reaction (ECSA) if the specific capacitance C_s (capacitance of an

atomically smooth planar surface per unit area) of the material under identical electrolytic conditions is known.⁵² Therefore, the double layer capacitance (C_{dl}) calculated from EIS circuits (Table S6b, ESI†) represents the electrochemical surface area (ECSA) available for the electrochemical reaction. From Table S6b (ESI†), it can be concluded that the 24 h hydrogenated sample possesses better ECSA than the 72 h hydrogenated sample. Additionally, the low charge transfer resistance R_{ct} (which is a measure of Faradaic activity at the electrode/electrolyte interface) of the 24 h hydrogenated sample also reflects the same. Therefore, the surface electrochemical parameter cannot be ruled out in true activity measurements of electrodes used in the PEC water splitting reaction. These experimental data assist in inferring the interplay of the coupled parameters (optical absorption, charge carrier separation, and surface electrochemistry) in enhancing the PEC water splitting reaction.

3.4 Role of amorphous layer/surface disorder

Sputtering *via* focused ion beam (FIB) was employed to remove few atomic layers from surface disordered/hydrogenated TiO_2 . For this purpose, we have taken the N-doped $\text{TiO}_2@500^\circ\text{C}/24\text{ h}$ sample for study. SEM images of the FIB sputtered sample show a smoother surface (Fig. 9b) in comparison to the unspun sample (Fig. 9d). The UV-Vis spectrum of the sputtered sample (Fig. 9c) shows a blue shift in absorption with respect to the unspun sample. This confirms that defects contributing to lowering of the band gap and they are concentrated in the first

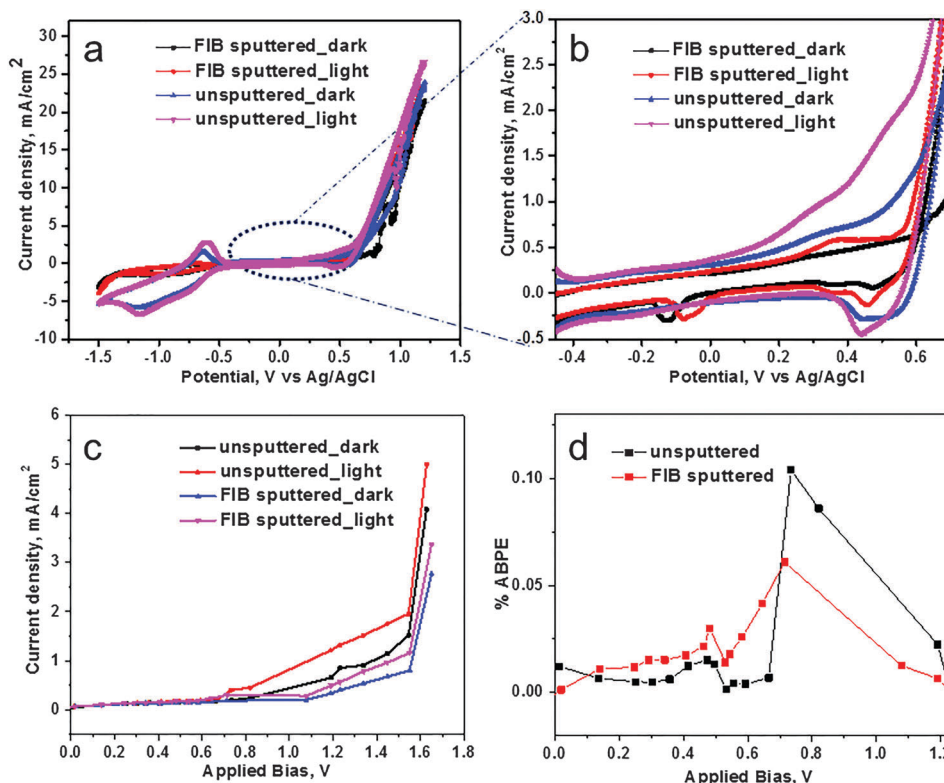


Fig. 10 EC/PEC data for the hydrogenated N doped $\text{TiO}_2@500^\circ\text{C}/24\text{ h}$ sample before and after FIB sputtering: (a) cyclic voltammetry data (3-probe), (b) the enlarged portion of fig. a, (c) current vs. applied bias plots (2-probe); and (d) the % applied bias photon to current conversion efficiency (ABPE).



few surface and subsurface layers in comparison to bulk. From Raman spectra, a decrease in the peak intensity ratio (B_{1g}/E_g or D/G) is observed for the sputtered portion of the sample over the unsputtered portion (Fig. S18a and b, ESI†). This observation further confirms the removal of surface defects present on the disordered/amorphous surface upon FIB sputtering. Furthermore, the XPS spectral peaks of sputtered and un-sputtered portions (Fig. S12d, ESI†) of the sample are tabulated in Table S7 (ESI†). There is no noticeable change from FTIR and XRD analyses before and after FIB sputtering (Fig. S18c and d, ESI†).

The sputtered sample shows less EC activity in comparison to the unsputtered sample (Fig. 10a). Furthermore, the change in the surface structure of the sample upon sputtering is observed from the additional redox peaks in the non-Faradiac region of CV (Fig. 10b). In addition, the decrease in PEC performance is observed for the sputtered sample over the unsputtered part (Fig. 10c). The calculated ABPE efficiencies for the sputtered and unsputtered portions of the sample are ~ 0.06 and 0.1% , respectively (Fig. 10d). Furthermore, a reduction in differential capacitance values (at 0.1 V and 0.2 V) is observed after sputtering/removing few atomic layers from the sample (Table S8, ESI†) which suggests a reduction in electrocatalytic active sites. The surface disorder/amorphous layer is found to be beneficial for enhancement of electrocatalytic activity, photon absorption cross-section, and charge carrier transport which is in agreement with recent literature reports.^{53,54}

4. Conclusions

In the present work, we explore the issue of enhancement in optical absorption cross-section and EC activity without compromising electron-hole pair separation in PEC water splitting reaction *via* hydrogenated (N-S) doped TiO_2 nanostructures. The maximum EC/PEC activity is found in the hydrogenated N-S co-doped $TiO_2@800^\circ C/24$ h sample with an applied bias photon to current conversion efficiency (ABPE) of $\sim 0.38\% @ 0.628$ V applied bias, an intrinsic solar to chemical conversion efficiency (ISTC) of $\sim 3.1\% @ 1.27$ V (*vs.* RHE) and an electrical and solar power-to-hydrogen (ESPH) conversion efficiency of $\sim 6.6\% @ 1.56$ V applied bias. The enhancement in PEC performance in the co-doped hydrogenated TiO_2 samples over doped/unhydrogenated TiO_2 samples is attributed to the combined effect of dopants, oxygen vacancy states, and Ti^{3+} states. The photon absorption cross-section is increased by shifting the VB upward through introduction of less electronegative dopants (N and S) and lowering of the CB through the formation of Ti^{3+} upon hydrogenation. The presence of the Ti^{3+} state and oxygen vacancies is confirmed through Raman, EPR, and XPS studies. Furthermore, double layer capacitance (C_{dl}), charge transfer resistance (R_{ct}), and film resistance (R_f) data obtained from electrochemical impedance spectroscopy (EIS) analysis also support the view point that optical absorption, charge carrier separation, and surface electrochemistry are not mutually exclusive pathways in enhancing the PEC performance of surface disordered/hydrogenated TiO_2 . The formation of the Ti^{3+} state and oxygen

vacancies due to hydrogenation leads to a few layers of the disordered surface/subsurface region. The removal of surface layers through focused ion beam sputtering results in a blue-shift of the absorption spectrum and also in the decrease of both electrochemical (EC) and photoelectrochemical (PEC) activities. This signifies that the surface disordered layer plays a vital role in enhancing both electrocatalytic activity and photon absorption. However, increase in defect states (Ti^{3+} and oxygen vacancy) beyond a point (extending towards bulk), while increasing the photon absorption cross-section, reduces the charge-carrier separation/transport and EC activity which in turn reduces the PEC activity. The understanding from the present work helps in the partial decoupling of the effects of photon absorption cross-section, surface electrochemistry, and charge-carrier transport and paves the way for the rational material design of water oxidation PEC catalysts.

Acknowledgements

We gratefully acknowledge the support of the Technology System Development program of the Department of Science and Technology (DST), Government of India *via* project DST/TSG/SH/2011/106.

References

- 1 X. B. Chen, L. Liu, P. Y. Yu and S. S. Mao, *Science*, 2011, **331**, 746–750.
- 2 A. Naldoni, M. Allieta, S. Santangelo, M. Marelli, F. Fabbri, S. Cappelli, C. L. Bianchi, R. Psaro and V. Dal Santo, *J. Am. Chem. Soc.*, 2012, **134**, 7600–7603.
- 3 N. Roy, Y. Sohn, K. T. Leung and D. Pradhan, *J. Phys. Chem. C*, 2014, **118**, 29499–29506.
- 4 T. Lin, C. Yang, Z. Wang, H. Yin, X. Lu, F. Huang, J. Lin, X. Xie and M. Jiang, *Energy Environ. Sci.*, 2013, **7**, 967–972.
- 5 D. M. Jang, I. H. Kwak, E. L. Kwon, C. S. Jung, H. S. Im, K. Park and J. Park, *J. Phys. Chem. C*, 2015, **119**, 1921–1927.
- 6 Q. Xiang, J. Yu, W. Wang and M. Jaroniec, *Chem. Commun.*, 2011, **47**, 6906–6908.
- 7 R. Su, R. Bechstein, J. Kibsgaard, R. T. Vang and F. Besenbacher, *J. Mater. Chem.*, 2012, **22**, 23755–23758.
- 8 L. Cai, I. S. Cho, M. Logar, A. Mehta, J. He, C. H. Lee, P. M. Rao, Y. Feng, J. Wilcox, F. B. Prinz and X. Zheng, *Phys. Chem. Chem. Phys.*, 2014, **16**, 12299–12306.
- 9 J. Buha, *J. Phys. D: Appl. Phys.*, 2012, **45**, 385305.
- 10 W. D. Chemelewski, H.-C. Lee, J.-F. Lin, A. J. Bard and C. B. Mullins, *J. Am. Chem. Soc.*, 2014, **136**, 2843–2850.
- 11 R. D. L. Smith, M. S. Prellivot, R. D. Fagan, S. Trudel and C. P. Berlinguette, *J. Am. Chem. Soc.*, 2013, **135**, 11580–11586.
- 12 R. D. L. Smith, B. Sporinova, R. D. Fagan, S. Trudel and C. P. Berlinguette, *Chem. Mater.*, 2014, **26**, 1654–1659.
- 13 E. Fabbri, A. Habereder, K. Waltar, R. Kotz and T. J. Schmidt, *Catal. Sci. Technol.*, 2014, **4**, 3800–3821.
- 14 S. Hoang, S. P. Berglund, N. T. Hahn, A. J. Bard and C. B. Mullins, *J. Am. Chem. Soc.*, 2012, **134**, 3659–3662.



- 15 L. Mi, Y. Zhang and P.-N. Wang, *Chem. Phys. Lett.*, 2008, **458**, 341–345.
- 16 S. P. Russo, I. E. Grey and N. C. Wilson, *J. Phys. Chem. C*, 2008, **112**, 7653–7664.
- 17 O. Diwald, T. L. Thompson, T. Zubkov, E. G. Goralski, S. D. Walck and J. T. Yates, *J. Phys. Chem. B*, 2004, **108**, 6004–6008.
- 18 R. W. Shunhang Wei, J. Jian, F. Chen and Y. Sun, *Dalton Trans.*, 2015, **44**, 1534–1538.
- 19 X. Yang, C. Cao, L. Erickson, K. Hohn, R. Maghirang and K. Klabunde, *Appl. Catal., B*, 2009, **91**, 657–662.
- 20 H. Liu and L. Gao, *J. Am. Ceram. Soc.*, 2004, **87**, 1582–1584.
- 21 X. Pan, M.-Q. Yang, X. Fu, N. Zhang and Y.-J. Xu, *Nanoscale*, 2013, **5**, 3601–3614.
- 22 Z. Wang, C. Yang, T. Lin, H. Yin, P. Chen, D. Wan, F. Xu, F. Huang, J. Lin and X. Xie, *Adv. Funct. Mater.*, 2013, **23**, 5444–5450.
- 23 A. P. Upadhyay, D. K. Behara, G. P. Sharma, A. Bajpai, N. Sharac, R. Ragan, R. G. S. Pala and S. Sivakumar, *ACS Appl. Mater. Interfaces*, 2013, **5**, 9554–9562.
- 24 J. Lynch, C. Giannini, J. K. Cooper, A. Loiudice, I. D. Sharp and R. Buonsanti, *J. Phys. Chem. C*, 2015, **119**, 7443–7452.
- 25 C. L. Haynes and R. P. Van Duyne, *J. Phys. Chem. B*, 2001, **105**, 5599–5611.
- 26 J. M. Coronado, A. J. Maira, J. C. Conesa, K. L. Yeung, V. Augugliaro and J. Soria, *Langmuir*, 2001, **17**, 5368–5374.
- 27 H. Liu, H. T. Ma, X. Z. Li, W. Z. Li, M. Wu and X. H. Bao, *Chemosphere*, 2003, **50**, 39–46.
- 28 N. Serpone, *J. Phys. Chem. B*, 2006, **110**, 24287–24293.
- 29 T. Berger, M. Sterrer, O. Diwald, E. Knözinger, D. Panayotov, T. L. Thompson and J. T. Yates, *J. Phys. Chem. B*, 2005, **109**, 6061–6068.
- 30 A. Fujishima, T. N. Rao and D. A. Tryk, *J. Photochem. Photobiol., C*, 2000, **1**, 1–21.
- 31 T. Su, Y. Yang, Y. Na, R. Fan, L. Li, L. Wei, B. Yang and W. Cao, *ACS Appl. Mater. Interfaces*, 2015, **7**, 3754–3763.
- 32 J. Cai, K. Lv, J. Sun and K. Deng, *RSC Adv.*, 2014, **4**, 19588–19593.
- 33 L. Xiong, M. Ouyang, L. Yan, J. Li, M. Qiu and Y. Yu, *Chem. Lett.*, 2009, **38**, 1154–1155.
- 34 X. Lu, G. Wang, T. Zhai, M. Yu, J. Gan, Y. Tong and Y. Li, *Nano Lett.*, 2012, **12**, 1690–1696.
- 35 J. M. Buriak, P. V. Kamat and K. S. Schanze, *ACS Appl. Mater. Interfaces*, 2014, **6**, 11815–11816.
- 36 C. H. Choi, S. H. Park and S. I. Woo, *ACS Nano*, 2012, **6**, 7084–7091.
- 37 P. Chen, T.-Y. Xiao, H.-H. Li, J.-J. Yang, Z. Wang, H.-B. Yao and S.-H. Yu, *ACS Nano*, 2012, **6**, 712–719.
- 38 E. L. Miller and R. E. Rocheleau, *J. Electrochem. Soc.*, 1997, **144**, 3072–3077.
- 39 H. Zhang, X. Lv, Y. Li, Y. Wang and J. Li, *ACS Nano*, 2009, **4**, 380–386.
- 40 D. C. Hurum, A. G. Agrios, K. A. Gray, T. Rajh and M. C. Thurnauer, *J. Phys. Chem. B*, 2003, **107**, 4545–4549.
- 41 A. Scalfani and J. M. Herrmann, *J. Phys. Chem.*, 1996, **100**, 13655–13661.
- 42 C. Yang, Z. Wang, T. Lin, H. Yin, X. Lü, D. Wan, T. Xu, C. Zheng, J. Lin and F. Huang, *J. Am. Chem. Soc.*, 2013, **135**, 17831–17838.
- 43 G. Wang, H. Wang, Y. Ling, Y. Tang, X. Yang, R. C. Fitzmorris, C. Wang, J. Z. Zhang and Y. Li, *Nano Lett.*, 2011, **11**, 3026–3033.
- 44 H. Cui, W. Zhao, C. Yang, H. Yin, T. Lin, Y. Shan, Y. Xie, H. Gu and F. Huang, *J. Mater. Chem. A*, 2014, **2**, 8612–8616.
- 45 Y. Yao, G. Li, S. Ciston, R. M. Lueptow and K. A. Gray, *Environ. Sci. Technol.*, 2008, **42**, 4952–4957.
- 46 B. Klahr, S. Gimenez, F. Fabregat-Santiago, J. Bisquert and T. W. Hamann, *Energy Environ. Sci.*, 2012, **5**, 7626–7636.
- 47 P. L. Varo, L. G. Bertoluzzi, J. A. J. Tejada and J. Bisquert, *J. Mater. Chem. A*, 2016, DOI: 10.1039/C5TA03210E.
- 48 Y.-H. Fang and Z.-P. Liu, *J. Am. Chem. Soc.*, 2010, **132**, 18214–18222.
- 49 L. Bertoluzzi and J. Bisquert, *J. Phys. Chem. Lett.*, 2012, **3**, 2517–2522.
- 50 B. Klahr, S. Gimenez, F. Fabregat-Santiago, J. Bisquert and T. W. Hamann, *J. Am. Chem. Soc.*, 2012, **134**, 16693–16700.
- 51 T. Lopes, L. Andrade, F. Le Formal, M. Gratzel, K. Sivula and A. Mendes, *Phys. Chem. Chem. Phys.*, 2014, **16**, 16515–16523.
- 52 C. C. McCrory, S. Jung, J. C. Peters and T. F. Jaramillo, *J. Am. Chem. Soc.*, 2013, **135**, 16977–16987.
- 53 P. Yan, G. Liu, C. Ding, H. Han, J. Shi, Y. Gan and C. Li, *ACS Appl. Mater. Interfaces*, 2015, **7**, 3791–3796.
- 54 T. Xia, W. Zhang, W. Li, N. A. Oyler, G. Liu and X. Chen, *Nano Energy*, 2013, **2**, 826–835.

