



Cite this: *Chem. Soc. Rev.*, 2021, **50**, 2173

Received 23rd August 2020

DOI: 10.1039/d0cs00357c

[rsc.li/chem-soc-rev](http://rsc.li/chem-soc-rev)

## Fundamentals and applications of photo-thermal catalysis†

Diego Mateo, Jose Luis Cerrillo, Sara Durini and Jorge Gascon \*

Photo-thermal catalysis has recently emerged as an alternative route to drive chemical reactions using light as an energy source. Through the synergistic combination of photo- and thermo-chemical contributions of sunlight, photo-thermal catalysis has the potential to enhance reaction rates and to change selectivity patterns, even under moderate operation conditions. This review provides the fundamentals of localized surface plasmon resonance (LSPR) that explain the photo-thermal effect in plasmonic structures, describes the different mechanistic pathways underlying photo-thermal catalysis, suggests methodologies to disentangle the reaction mechanisms and proposes material design strategies to improve photo-thermal performance. Ultimately, the goal is to pave the way for the wide implementation of this promising technology in the production of synthetic fuels and chemicals.

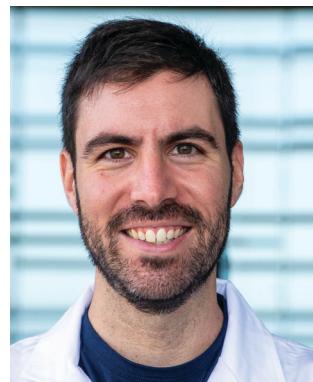
### 1. Introduction

Catalysis is an integral part of chemistry with large implications for the efficient production of everyday goods. Indeed, *circa* 95% of all chemical products are manufactured through catalytic routes. At the same time, many of these processes require an energy input that has, traditionally, been supplied by non-renewable sources.<sup>1,2</sup>

King Abdullah University of Science and Technology, KAUST Catalysis Center (KCC), Advanced Catalytic Materials, Thuwal 23955-6900, Saudi Arabia.

E-mail: [Jorge.gascon@kaust.edu.sa](mailto:Jorge.gascon@kaust.edu.sa)

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d0cs00357c



**Diego Mateo**

Diego Mateo received his bachelor's degree in Chemistry in 2009 and his bachelor's degree in Food Technology in 2011, both from the University of Valencia. After working for two years as a research technician at Complutense University of Madrid, in 2015 he started his PhD in the development of new graphene-based photocatalysts for the production of solar fuels in the group of Prof. Hermenegildo Garcia at the Institute of Chemical Technology (UPV-CSIC). After obtaining his PhD in 2019, he moved to the group of Prof. Jorge Gascon at King Abdullah University of Science and Technology (KAUST) as a postdoctoral fellow. His current research is focused on novel photo- and electro-catalytic routes for the production of fuels and chemicals.



**Jose Luis Cerrillo**

Jose Luis Cerrillo (1986) obtained his BSc in Chemistry in 2011 from the Universitat de València (Spain). After two years working as a R&D technician in a chemical company, he came back to academia and obtained his MSc in 2014 and his PhD in Sustainable Chemistry in 2019, both from Universitat Politècnica de València (Spain). His PhD research was funded by a Severo-Ochoa National scholarship and held at the Instituto de Tecnología Química (CSIC-UPV), focusing on the catalytic removal of water pollutants and on the study of novel biocide materials based on Ag-zeolites. Since October 2019, he has been doing postdoctoral research in Prof. Gascon's research group (Advanced Catalytic Materials) in King Abdullah University of Science and Technology (Saudi Arabia) where he investigates diverse heterogeneous catalytic processes related to the reutilization of CO<sub>2</sub> to produce valued fuels and chemicals.



In the last few decades, with the discovery of heterogeneous semiconductor-based photocatalysis, the possibility of catalyzing chemical transformations using light as a source of energy has become a reality.

In a semiconductor, upon absorption of a photon with equal or higher energy than the band gap, electron–hole pairs are generated. Eventually these charge carriers can migrate to the semiconductor surface and be transferred to adsorbed molecules, thereby initiating reduction or oxidation processes. Since the seminal paper of Fujishima and Honda in 1972, tremendous advances in the field of photocatalysis have been reported. In spite of this, the efficiency of most photocatalytic processes remains insufficiently low (typically in the range of hundreds of  $\mu\text{mol g}^{-1} \text{h}^{-1}$ ) mainly due to fast charge carrier recombination and low absorption and utilization of the solar spectrum by traditional wide band gap semiconductors.<sup>3–8</sup>

Next to charge separation upon light excitation, plasmon-mediated processes have attracted a great interest. Hot charge carriers generated by means of the decay of localized surface plasmon resonance (LSPR) possess higher energies than those induced by direct photoexcitation.<sup>9,10</sup> Furthermore, these energetic carriers can relax internally and dissipate their energy by local heating of the surroundings, causing a thermal effect on the material.<sup>11–13</sup> This photo-thermal effect has been extensively applied in a large number of fields, including cancer therapy, degradation of pollutants, seawater desalination and water vaporization.<sup>14–19</sup> It was therefore only a matter of time until similar concepts were applied to speed up chemical reactions. Indeed, photo-thermal catalysis combines photochemical and thermochemical contributions of sunlight and has emerged as a rapidly growing and exciting new field of research.<sup>20–31</sup> Photo-thermal catalysis allows for a more effective harvesting of the solar spectrum, including low-energy visible and infrared photons that would be insufficient to promote photocatalytic reactions.<sup>30,32–34</sup> Furthermore, due to

the increase in the temperature of the catalytic active sites, photo-thermal catalysis renders outstanding production rates while operating under moderate conditions.<sup>35–37</sup>

Because of the rapidly growing nature of the field, we believe now is the time to write a holistic review on the field of photo-thermal catalysis focusing on the production of fuels and chemicals. As relatively newcomers, we make special emphasis on those topics that should facilitate the entrance of other newcomers to this exciting field. We therefore pay special attention to the fundamentals of localized surface plasmon resonance and the underlying mechanisms behind photochemical and thermochemical processes that define the photo-thermal effect. Based on this analysis, we suggest characterization techniques and methodologies that should help assess the dominating pathway of reaction for a given catalytic system, and therefore gain a better insight. Next, we review in detail reactions driven by means of photo-thermal catalysis, gathering the most relevant information available in terms of efficiency, selectivity targets and mechanistic knowledge. Last but not least, we share our opinion on design strategies to further optimize catalytic performance and highlight potential future directions and challenges ahead.

## 2. Localized surface plasmon resonance and the photo-thermal effect

The localized surface plasmon resonance (LSPR) band in metallic nanoparticles (NPs) is an intense and broad absorption band along the UV-visible-NIR region of the electromagnetic spectrum. The physics behind this phenomenon has been traditionally explained by means of two theories: the Drude–Maxwell model and the theory developed by Gustav Mie in 1908.<sup>38,39</sup> Although it is beyond the scope of this review to



Sara Durini

University of Science and Technology (KAUST, Saudi Arabia) as a postdoctoral fellow, where she investigates porous materials for application in heterogeneous catalysis.

Sara Durini obtained her bachelor's degree in Chemistry and Industrial Chemistry in 2009, her master's degree in Chemical Sciences in 2011, and her PhD in Chemistry in 2015 from the University of Insubria in Como (Italy). In 2016 she continued her work as a postdoctoral researcher at the University of Leipzig, where she worked on the synthesis of new luminescent metal–organic frameworks. In 2019 she joined Prof Gascon's group at King Abdullah University of Science and Technology (KAUST, Saudi Arabia) as a postdoctoral fellow, where she investigates porous materials for application in heterogeneous catalysis.



Jorge Gascon

co-authored over 275 publications, several patents and books. His research group focuses on the design, development and demonstration of new heterogeneous catalysts and reactor engineering concepts.

Jorge Gascon (1977) received his MSc in Chemistry in 2002 and his PhD in Chemical Engineering in 2006, both from the University of Zaragoza (Spain). After his PhD he spent eleven years at TU Delft, the last four as Antonie van Leeuwenhoek Professor of Catalysis Engineering. In 2018 he moved to King Abdullah University of Science and Technology (KAUST), where he is Professor of Chemical Engineering and the Director of the KAUST Catalysis Center. He has



contribute to a deep and extensive discussion of these theories, these models can provide insightful information to describe the LSPR phenomenon and its application to catalysis, so a brief overview of them will be addressed in the ESI<sup>†</sup> of the present work.

As per these theories, inside metal NPs conducting free electrons can move when guided by external incident irradiation. This motion is damped by electron inelastic collisions and the restoring force on the electron cloud created by the accumulation of surface charges (Fig. 1).

Under resonant conditions, both the incident electromagnetic wave and the resonant frequency of conduction electrons are in phase, thus maximizing the electric field in so-called “hot spots” on the surface of plasmonic NPs. Apart from the enhanced electric field, other relevant properties of LSPR arise from the different relaxation processes that occur within plasmonic structures. After excitation, the energy stored in surface plasmons can decay through several pathways, either radiative (as re-emitted photons) or non-radiative (electron–hole pair excitations and electron–electron collisions).<sup>40</sup> Although for some specific purposes non-radiative losses are considered as a drawback for plasmonic performance, actually they can show potential scientific and technological applications.<sup>13,41</sup> For instance, non-radiative decay may produce local heating as the electronic kinetic energy can be transmitted to metal lattice phonons (see Section 3.2). In spite of the fact that photo-thermal cancer therapy has traditionally benefited from this phenomenon, in the recent years the plasmon-induced heating effect has been also applied in a variety of research fields including environmental remediation, solar steam generation and, of more relevance to the scope of this review, catalytic processes.<sup>17,42,43</sup>

Besides, plasmon energy can be non-radiatively dissipated by absorption within the metal NPs and, eventually, generate energetic hot charge carriers in the plasmonic structure.

This contribution to plasmon dephasing comes from electronic collisions with the surface of the plasmonic structure (surface scattering, also known as “Landau damping”) and is a pure quantum mechanical process in which the energy from a plasmon quantum is transferred on the timescale of femtoseconds (1–100 fs) into a single electron–hole pair excitation.<sup>40</sup> This phenomenon is due to the non-conservation of the linear momentum of electrons near the surface and in hot spots, and enables an electron to absorb a photon quantum with energy  $\hbar\omega$ .<sup>44</sup> As we will describe in detail in the next section, hot carriers are able to escape from the plasmonic NPs and induce further chemical reactions.

However, before we delve into the mechanisms underlying photo-thermal catalysis, we would like to clarify some terms commonly used in the literature that, due to the lack of a strict definition, may cause misinterpretations in this field. Plasmon-enhanced catalysis derives from the synergistic combination of three properties arising from LSPR: hot carrier generation, local heating effect and optical near-field enhancement.<sup>45</sup> When it comes to photo-thermal catalysis, Ozin *et al.* determined that hot carriers and the thermal effect contribute synergistically to the enhanced overall reaction rate, with both mechanisms playing a role to various degrees depending on the operating system.<sup>30,46</sup> Therefore, in the case of plasmonic materials the photo-thermal effect arises from the combination of both thermal and photochemical contributions of non-radiative plasmon decay, as we will illustrate in subsequent sections. It should be emphasized, however, that the photo-thermal effect is not an exclusive phenomenon of plasmonic NPs. In fact, semiconductor oxides, chalcogenides, metal organic frameworks and other carbon-based materials such as graphene or carbon nanotubes (CNTs) can display photo-thermal performance as well, although in these cases the operating mechanisms have been less studied.<sup>29,47</sup> For these reasons, in the next

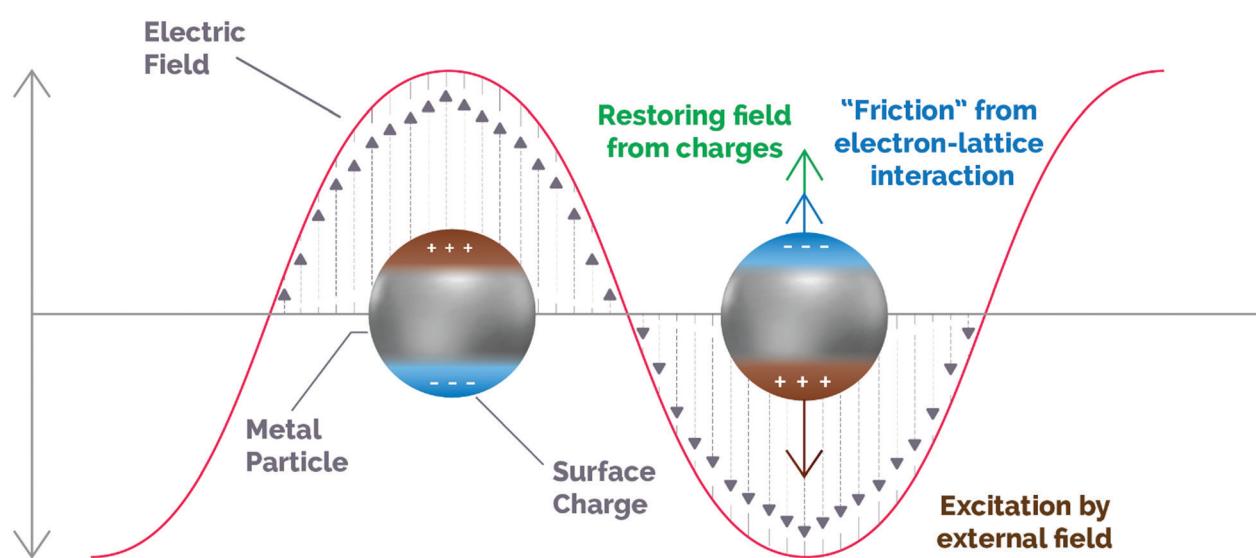


Fig. 1 Schematic illustration of the dynamics of an excited plasmonic nanoparticle. The motion of carriers follows the changes induced by the incoming electromagnetic field (brown), while a restoring force is generated by the out-of-equilibrium surface charges (green) and the ionic network produces damping due to electronic collisions (blue). Adapted with permission from ref. 45. Copyright (2019) American Chemical Society.



section we will also give a brief overview of these non-plasmonic approaches to photo-thermal catalysis.

### 3. Photo-thermal enhanced catalysis

#### 3.1. Photochemical enhancement in plasmonic structures

As mentioned above, the excitation of plasmonic NPs can generate hot carriers (*i.e.* electrons and holes) *via* electronic intraband and/or interband transitions through non-radiative Landau damping.<sup>48</sup> The absorption of a photon with energy  $h\nu$  promotes hot electrons with energies above the Fermi level ( $E_F + h\nu$ ) that, upon interaction with other species possessing electron-accepting orbitals, may be injected from the metal nanoparticle to these species.<sup>48,49</sup> The kinetic energy of hot electrons excited during plasmon decay is transferred to the adsorbates, that are chemically activated *via* vibrational or electronic transitions. Ultimately, the promotion of high-energy electrons to antibonding orbitals of adsorbed molecules can result in the cleavage of molecular bonds, triggering subsequent chemical transformations.

These above-mentioned hot electrons can be exploited for catalytic applications in four main pathways, depending on the existence of a single-component or a heterostructured plasmonic photocatalyst. The first two pathways make use of hot electrons of single plasmonic NPs interacting with adsorbates *via* indirect or direct electron transfer. The last two pathways are related to interactions of supported plasmonic metal nanostructures with semiconductors *via* an indirect injection mechanism to an acceptor or through direct promotion of the carrier into the conduction band (CB).<sup>50,51</sup>

##### 3.1.1. Indirect hot electron transfer into the adsorbate.

According to the indirect electron transfer, hot electrons are first generated within the metal nanoparticle after plasmon excitation and subsequently transferred to the lowest unoccupied molecular orbital (LUMO) of the adsorbed species on a metal surface.<sup>40,52,53</sup> Since in this case the electron transfer into the adsorbate takes place after the generation of hot carriers, it is restrained by energy losses due to electron-electron scattering. The efficiency of this indirect hot electron transfer shows a positive correlation with incident photon energy, as a higher energy will generate more electrons with sufficient potential to be injected into the LUMO of adsorbates.<sup>54</sup>

This indirect mechanism for electron transfer has been evidenced in the plasmon-mediated activation of small molecules such as molecular hydrogen ( $H_2$ ) and oxygen ( $O_2$ ).<sup>49,55-58</sup> Recently, Halas *et al.* using both DFT calculations and H/D exchange experiments proved that hot electrons generated on Au NPs could transfer *via* an indirect pathway to the antibonding orbital of molecular  $H_2$ , thereby creating a negatively charged  $H_2^{\delta-}$  species. Eventually, electrons transferred back to Au and  $H_2$  returned to their electronic ground state with accumulated vibrational energy that led to ultimate  $H_2$  dissociation.<sup>55</sup> Linic and co-workers also reported the plasmon-induced partial oxidation of ethylene using Ag nanocubes. In this study, hot electrons generated on the surface of Ag nanocubes were transferred to antibonding states of  $O_2$  by an indirect mechanism, leading to the generation of  $O_2^-$ . In the same line as in the case of  $H_2$  activation, the accumulated vibrational

energy overcomes the activation energy barrier and causes  $O_2$  dissociation.<sup>49</sup>

**3.1.2. Direct hot electron transfer into the adsorbate.** The excitation of LSPR can also lead to the direct injection of hot electrons from metal NPs to adsorbates.<sup>48,53,59</sup> This plasmon dephasing pathway originates from chemical interface damping (CID), which arises from the coupling between unoccupied adsorbate states and excited surface plasmons.<sup>50</sup> In the direct transfer, hot electrons are directly transferred into these hybridized states between the metal nanoparticle and adsorbed molecules.<sup>60</sup> Therefore, as opposed to the indirect mechanism, which takes place after the generation of hot electrons, direct electron transfer occurs simultaneously during the dephasing of the plasmon excitation.<sup>45</sup> For this reason, the direct electron transfer mechanism is expected to show a higher electron-transfer efficiency and minor energy losses, as it circumvents electron-electron scattering.<sup>61</sup> Nevertheless, this alternative pathway for direct electron promotion has a lower probability to occur, as it requires an intense adsorbate/metal interaction for surface orbital hybridization that is not usual in plasmonic photocatalysts.<sup>50</sup>

Interestingly, Linic's group observed that the photo-degradation of methylene blue by Ag nanocubes was driven *via* direct electron transfer.<sup>53,60</sup> Although Ag nanocubes showed two plasmon absorption bands centered at 532 and 785 nm, the authors reported a higher degradation yield of methylene blue under irradiation at 785 nm. These results indicated that the longest wavelength, and so the lowest in energy, could transfer hot electrons more efficiently, thus suggesting a one-step process (direct electron transfer) rather than a two-step pathway (indirect mechanism). Based on their hypothesis, a strong hybridization between methylene blue and the Ag interface was established, generating electron-acceptor orbitals (LUMO) centered on methylene blue and electron-donor orbitals (HOMO) centered on Ag.<sup>53</sup> However, for the LSPR decay of Ag to directly generated hot electrons in the electron-acceptor states, the energy gap of the HOMO-LUMO transition had to be resonant with the Ag plasmon band at 785 nm. Under this condition, the Ag LSPR decay allowed the direct electron promotion to the excited hybridized states.<sup>53,60</sup> These results demonstrated that the direct electron transfer pathway can be a powerful strategy to selectively activate reactants only by choosing the wavelength that is resonant with the transition between the hybridized states.<sup>50</sup>

##### 3.1.3. Indirect hot electron transfer into the semiconductor.

The use of supported plasmonic NPs as light absorbers has been an extensively described methodology in the field of photocatalysis since it provides a better harvesting of the solar spectrum than traditional wide band gap semiconductors.<sup>62</sup> Metallic NPs with LSPR bands along the visible and IR region allow the complete exploitation of solar radiation, especially from low-energy wavelengths that would be inadequate to excite electron-hole pairs in UV-active semiconductor photocatalysts. Furthermore, the use of hybrid plasmonic metal-semiconductor structures displays the additional advantage of favoring the spatial separation of photo-induced electron-hole pairs once electrons are promoted to the semiconductor, preventing charge carrier recombination within the metal and thus extending their lifetime.<sup>45,63,64</sup>



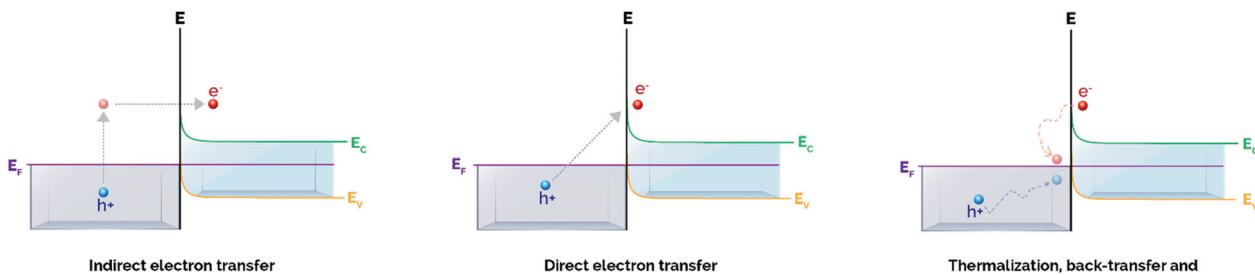


Fig. 2 Plasmon-induced hot carrier generation and hot electron transfer/back-transfer processes in metal/semiconductor heterostructures. Adapted with permission from ref. 50. Copyright (2017) American Chemical Society.

In a similar way to the indirect mechanism for non-supported plasmonic NPs, the most common mechanism for hot electron transfer in heterostructures takes place *via* a two-step pathway that involves hot carrier generation within the plasmonic nanoparticle in the time scale of femtoseconds followed by electron transfer through the metal/semiconductor interface (Fig. 2).<sup>50</sup>

Upon the formation of a junction between a semiconductor and a metal nanoparticle, the Fermi values of both components align to favor the redistribution of charge and to ensure equilibrium in the system (Fig. 3).<sup>65</sup> This gives rise to bending of both the valence and conduction bands of the semiconductor, thus creating a Schottky barrier at the semiconductor/metal interface whose energy ( $\varphi_{SB}$ ) corresponds to the difference between the Fermi level of the metal and the interfacial CB edge.<sup>52,66</sup>

Importantly for catalytic applications, only hot electrons with energies above the energy threshold settled by the Schottky barrier will inject into the CB and will be available for subsequent electron-induced reactions at the semiconductor surface. In this regard, a low  $\varphi_{SB}$  will allow the promotion of a larger number of carriers at energies sufficiently high to traverse the barrier, and hence available for participating in chemical transformations. However, it is also of pivotal importance that the conduction band bending at the interface is high

enough to prevent the back-transfer of electrons to the metal nanoparticle (Fig. 2), thereby guaranteeing the spatial separation of electron–hole pairs and increasing the average lifetime of charge carriers.<sup>45</sup> For these reasons, an adequate balance between these two effects has to be taken into account when designing metal/semiconductor heterojunctions.

Typically, the magnitude of the Schottky barrier ranges between 0.5 and 1.5 eV and clearly determines the efficiency of hot electron injection from the metal nanoparticle to the CB of the semiconductor.<sup>52</sup> Most relevant, the height of the Schottky barrier is smaller than the band gap of many semiconductors, and this features one of the key advantages of these heterostructures, as excited electrons are not required to possess greater energy than the semiconductor band gap to be extracted.<sup>67</sup> As a matter of fact, many studies have described so far this metal–semiconductor indirect electron transfer mechanism and its practical application in a vast number of photocatalytic reactions including water splitting, CO<sub>2</sub> reduction or the degradation of pollutants.<sup>21,68–81</sup>

### 3.1.4. Direct hot electron transfer into the semiconductor.

Besides the indirect electron transfer pathway previously described, some studies have reported a possible direct electron transfer mechanism within metal/semiconductor heterostructures (Fig. 2). It is believed that in the indirect pathway the time scale for electron transfer should be in the order of picoseconds.<sup>52</sup> However, according to transient absorption spectroscopy measurements in Au/TiO<sub>2</sub> nanostructures, Furube *et al.* determined a much faster electron transfer time scale in the order of hundreds of femtoseconds.<sup>82</sup> In addition to this, the overall quantum yield (QY) of the electron transfer process was as high as 40%, compared to the maximum theoretical value for indirect hot electron transfer established below 8%.<sup>83</sup> In line with this work, several studies have also described ultrafast hot electron injection in plasmonic metal–semiconductor hybrid materials, some of them showing surprisingly high QYs.<sup>84–87</sup> All these data suggested the existence of a single-step electron transfer based on the direct injection of hot electrons to acceptor orbitals in the semiconductor upon plasmon dephasing in metal/semiconductor heterojunctions. In accordance with the direct electron transfer from a plasmonic metal nanoparticle to an adsorbate (see Section 3.1.2), this pathway is assumed to be more efficient as it avoids the energy loss of hot carriers associated to electron–electron and electron–phonon scattering within the metal nanoparticle.<sup>88</sup>

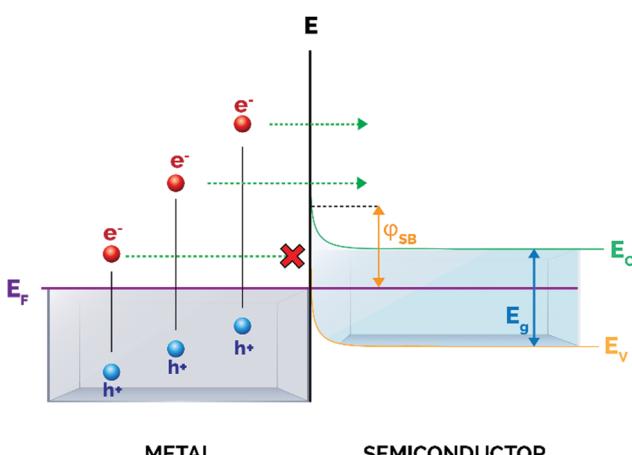


Fig. 3 Schematic illustration of the metal/semiconductor Schottky barrier. Upon excitation, only hot electrons with enough energy can surpass the Schottky barrier ( $\varphi_{SB}$ ) and transfer into the conduction band of the semiconductor.



In a recent work, Lian *et al.* observed a broadened absorption peak in the UV-visible spectrum for a Au/CdSe nanorod composite, suggesting a strong orbital hybridization between both Au and CdSe components.<sup>89</sup> Transient absorption spectroscopy data also revealed an electron transfer time scale in the order of tens of femtoseconds with an overall efficiency exceeding 24%.<sup>89</sup> In 2017, Ratchford and co-workers also determined an outstanding hot electron injection efficiency up to 45% of Au NPs embedded in TiO<sub>2</sub> films.<sup>85</sup> All these results were inconsistent with the traditional indirect pathway, hence indicating the existence of a direct electron transfer mechanism.

The assessment of the dominating electron transfer mechanism in metal/semiconductor hybrid materials can be tricky, but the study of QY can be a useful tool to distinguish the reaction pathway. For example, recent works have reported a very low QY of electron transfer for Au/graphene (~10%) and Au/CdS (~2.75%) heterostructures, thus revealing the existence of an indirect electron transfer mechanism.<sup>90,91</sup> Nevertheless, most of the current studies still have not developed a solid distinction of electron transfer mechanisms in metal/semiconductor heterojunctions.<sup>50</sup>

### 3.2. Thermal enhancement in plasmonic structures

In the previous sections we have described the relevant physical mechanisms underlying photo-thermal catalysis of ultrafast processes (below 100 fs) occurring upon plasmon decay (*i.e.* electron-hole pair excitations *via* Landau damping). At longer relaxation times, hot carriers decay and convert their energy into heat through inelastic coulombic electron-electron scattering over a period from 100 fs to 1 ps. In parallel with this process, low energy

electrons couple with phonon modes *via* electron-lattice collisions and so increasing the temperature of the metal lattice with a timescale of several to some hundreds of ps. At a later stage, this heat is transferred from the metal lattice to the surroundings in the time scale from 100 ps to 10 ns through phonon-phonon scattering (Fig. 4).<sup>59,92,93</sup> Therefore, plasmon-induced heat can result in energy transfer from the metal nanoparticle to surface adsorbates and subsequently trigger further chemical transformations following an Arrhenius dependence of the reaction rate. Ultimately, this implies that the use of plasmon heating to perform chemical reactions is similar to externally heating the system and does not offer a single pathway to control product selectivity.

In the particular case of a single spherical metal nanoparticle, the local temperature increase ( $\Delta T$ ) caused by the photo-thermal effect can be described by eqn (1):<sup>94</sup>

$$\Delta T(r) = \frac{V_{NP}Q}{4\pi k_0 r} \quad (1)$$

where  $r$  is the distance between the center of the NP and the boundary,  $V_{NP}$  the volume of the nanoparticle,  $Q$  the heat generation and  $k_0$  the thermal conductivity of the medium.

If we consider that the incident light wavelength is much larger than the nanoparticle radius, then the heat generation  $Q$  can be expressed as follows (eqn (2)):

$$Q = \frac{\omega}{8\pi} E_0^2 \left| \frac{3\epsilon_0}{2\epsilon_0 + \epsilon_{NP}} \right|^2 \text{Im}\epsilon_{NP} \quad (2)$$

where  $E_0$  is the amplitude of the incident electromagnetic radiation and  $\epsilon_0$  and  $\epsilon_{NP}$  are the dielectric constants of the

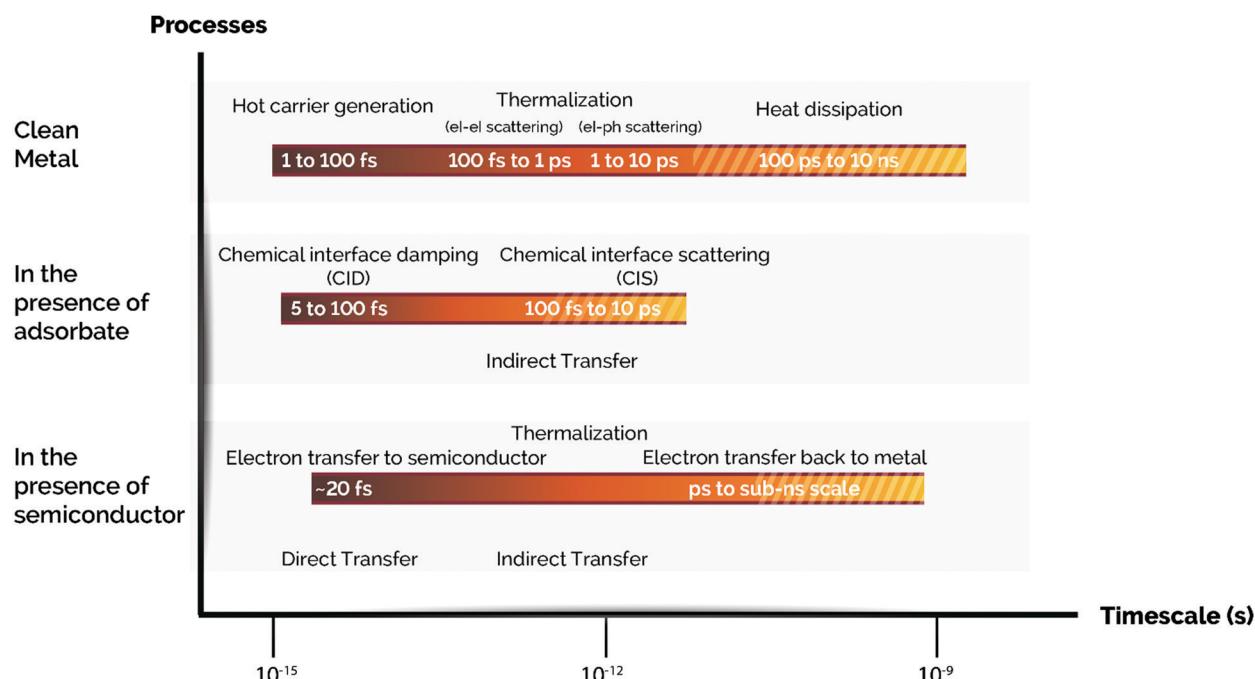


Fig. 4 Time scales of plasmon-induced hot carrier generation, hot electron transfer, and thermalization processes with/without an adsorbate or a semiconductor. From top to bottom: A series of time scales corresponding to the fate of hot carriers in a clean metal NP, additional events in a metal NP capped with an adsorbate, and the processes involved in a metal NP loaded on a semiconductor support *via* a Schottky contact are shown. Note: el stands for electron and ph stands for phonon. Adapted with permission from ref. 50. Copyright (2017) American Chemical Society.



surrounding medium and the nanoparticle, respectively. By means of derivation whose details can be found elsewhere,<sup>95</sup> it is possible to express the maximum temperature increase at  $r = R_{\text{NP}}$  (eqn (3)):

$$\Delta T(I_0) = \frac{R_{\text{NP}}^2 \omega}{3k_0} \left| \frac{3\epsilon_0}{2\epsilon_0 + \epsilon_{\text{NP}}} \right|^2 \text{Im} \epsilon_{\text{NP}} \frac{8\pi I_0}{c\sqrt{\epsilon_0}} \quad (3)$$

where  $I_0$  is the light intensity inside the matrix.

It is of great importance that from this equation it is possible to establish the dependence of temperature on the nanoparticle size:

$$\Delta T \propto R_{\text{NP}}^{-2}$$

Based on this, the temperature increase is proportional to the square of the nanoparticle radius. Essentially, this implies that as the nanoparticle size increases, the thermal contribution to the photocatalytic rate is amplified.

Nevertheless, several experimental and theoretical studies suggest that the thermal contribution to photo-thermal catalysis represents a minority weight in the overall catalytic rate.<sup>49,94</sup> An interesting work from Govorov *et al.* theorized that for light intensities similar to continuous-wave solar irradiance (around 100 mW cm<sup>-2</sup>) the temperature increment due to irradiation of a single gold NP became negligible ( $\sim 10^{-2}$  K). Only when the photon flux was about 10 000 times more intense a significant temperature rise (in the order of a few K) was achieved.<sup>95</sup> If we take into account that for commercial catalytic processes an increase of 10 K approximately doubles the reaction rate, it is difficult to imagine a scenario in which the light-induced heating plays a crucial role in the overall catalytic performance, at least under common solar irradiance operation conditions.<sup>96</sup> It is worth reminding, however, that according to Govorov *et al.* the light-induced heating effect can be strongly enhanced in the presence of a large number of NPs owing to collective effects, so this opens up the possibility of achieving further temperature increases with appropriate material engineering.<sup>95</sup> These observations are in line with those from Baffou *et al.* regarding significant collective temperature increases even in the presence of small amounts of plasmonic NPs.<sup>97</sup>

In fact, several groups have reported so far the use of plasmon-mediated heating in order to perform catalytic reactions. Boyd *et al.* described the catalytic steam reforming of ethanol using a flow-type microfluidic channel reactor with embedded 20 nm diameter Au NPs.<sup>98</sup> Upon irradiation of a mixture of ethanol–water with a low-power laser at the resonant frequency of Au NPs, the authors could observe the formation of gaseous products corresponding to CO, CO<sub>2</sub> and H<sub>2</sub>. These results demonstrated the possibility to drive endothermic reactions by plasmon-assisted heat generation.

More recently, Jiang and co-workers reported the first application in catalysis of the synergistic combination of the photo-thermal effect of plasmonic NPs with the advantageous properties of metal–organic frameworks (MOFs).<sup>99</sup> In this work, the authors developed a Pd@ZIF-8 composite for the light-assisted selective hydrogenation of olefins at room temperature. Interestingly, the authors found that while Pd nanocube cores acted as nano-heater

active sites, the ZIF-8 shell played multiple roles in the reaction mechanism. On the one hand, it accelerated the reaction rate due to its ability to act as a H<sub>2</sub> gas reservoir. On the other hand, it favored Pd nanocube stabilization preventing them from migration/aggregation and acted as a molecular sieve for selective catalysis. As a result, the Pd@ZIF-8 composite effectively catalyzed the hydrogenation of olefins under UV-visible irradiation at room temperature, achieving a 66% conversion. Noticeably, the reaction yield was about 50% by heating at 50 °C under dark conditions, hence indicating the superior catalytic performance of the Pd@ZIF-8 composite upon light irradiation.

Similarly, Xu *et al.* also described a hybrid core–shell hierarchical nanostructure based on plasmonic NPs and MOFs for the photo-thermally catalyzed cyclocondensation reaction.<sup>35</sup> Upon NIR irradiation, the plasmonic photo-thermal core of Cu<sub>7</sub>S<sub>4</sub> generated a high temperature and efficiently transferred heat to the catalytic shell of ZIF-8 containing the acid–base Lewis active sites. This resulted in a 4-fold enhancement of the catalytic activity for cyclocondensation reaction compared to dark conditions.

It should be mentioned, however, that there is a lack of consensus across the literature when evaluating the thermal contribution to the overall photocatalytic rate in plasmon-mediated reactions.<sup>100</sup> Consequently, it is still tricky to successfully distinguish the thermochemical pathway from the pure photochemical mechanism both underlying the photo-thermal effect. In Section 3.4, we will describe different methodologies that can help to identify which of the aforementioned mechanisms (*i.e.* photochemical and/or thermal) is preeminent in the photo-thermal reactions under study.

### 3.3. Photo-thermal catalysis in non-plasmonic structures

In a similar way to plasmonic structures, non-plasmonic elements can display the photo-thermal effect through direct intraband and/or interband electronic transitions. For instance, Sarina *et al.* demonstrated that non-plasmonic metal NPs supported on ZrO<sub>2</sub> could catalyze cross-coupling reactions at low temperatures under visible light.<sup>101</sup> According to the authors, upon irradiation with UV light electrons could shift to high-energy levels through interband transitions, and only those with enough energy could transfer to the LUMO of adsorbed molecules, just like in the case of plasmonic metal NPs. When excited with low-energy visible–IR light, electrons were not energetic enough to be injected into adsorbate states, thus contributing to the enhancement of the reaction rate by means of thermal effects.

In the case of metal oxide non-plasmonic semiconductors, an effective charge separation is of pivotal importance to maximize the efficiency of the photochemical process. However, in these systems non-radiative relaxation processes compete with charge separation and occur in the form of either Auger or Shockley–Read–Hall recombination.<sup>30</sup> It is of great importance that ultimately both processes are responsible for the emission of energy as lattice vibrations, hence providing an active mechanism for heat generation in this type of structures. Auger recombination takes place when an electron and hole recombine but, instead of primarily releasing the excess of energy in the form of light or heat, the energy is transferred to a third charge carrier that



thermalizes back to the conduction band edge by emitting phonon vibrations. An alternative pathway for non-radiative relaxation is Shockley–Read–Hall recombination (also known as “trap-assisted recombination”). This mechanism usually takes place in semiconductors displaying a high degree of defects or vacancies that create mid-gap energy states. After excitation, electrons quickly move to these states releasing energy as phonons that eventually increase the temperature of the semiconductor. Altogether, photochemical and thermal effects contribute to the photo-thermal performance in non-plasmonic semiconductors. In 2019, Ye *et al.* reported an example of this approach for the photo-thermal reduction of  $\text{CO}_2$  using 2D  $\text{In}_2\text{O}_{3-x}$  as a catalyst.<sup>102</sup> Compared with commercial  $\text{In}_2\text{O}_3$ , defective 2D  $\text{In}_2\text{O}_{3-x}$  displayed an outstanding photo-thermal CO production of  $103.2 \text{ mmol g}^{-1} \text{ h}^{-1}$  by increasing the temperature of the system to  $280^\circ\text{C}$  within 10 min. The authors rationalized that in this case oxygen vacancies in defective 2D  $\text{In}_2\text{O}_{3-x}$  had a dual function; in other words, they provided efficient local heat generation and active sites for  $\text{CO}_2$  adsorption that basically enhanced the catalytic performance.

Besides non-plasmonic metallic NPs and semiconductors, the photo-thermal effect has been also reported in a myriad of materials including MOFs, chalcogenides, MXenes and other carbon-based materials. In spite of the fact that in these cases the mechanisms underlying the photo-thermal effect are not completely understood, all of them offer promising applications not only in the field of catalysis, but also in photovoltaics and solar thermal energy.<sup>47</sup>

### 3.4. Methodologies to identify the reaction mechanisms in photo-thermal processes

The photo-thermal effect results from the combination of both photochemical and thermochemical contributions. Taking into account that the two mechanisms can coexist simultaneously, it is very challenging and not always evident to distinguish which pathway drives predominantly the system. For this reason, in this section we have summarized and listed different experimental and computational methodologies that can help to

distinguish between hot-carrier and thermal effects in photo-thermal processes.

**3.4.1. Experimental procedures.** The study of the influence of light intensity on the photocatalytic rate can be a powerful and straightforward analytic tool to comprehend the reaction mechanisms in photo-thermal processes. In 2014, Christopher *et al.* established four different kinetic categories for the light intensity ( $I$ ) dependence on the photocatalytic reaction rate: (1) sublinear (rate  $\propto I^n$ ,  $n < 1$ ), linear (rate  $\propto I$ ), superlinear (rate  $\propto I^n$ ,  $n > 1$ ) and exponential (rate  $\propto e^{f(I)}$ ).<sup>48</sup> Plasmon-induced photocatalysis has been demonstrated to exhibit each of these regimes depending on the operating conditions, except for the sublinear dependence that is characteristic of non-plasmonic semiconductor catalysis when charge carrier recombination is dominant.<sup>45</sup>

Previous experimental and theoretical studies have shown that there is a linear relationship between the photon flux and the surface temperature increase of plasmonic NPs.<sup>94</sup> However, if this linear relationship is combined with the Arrhenius equation for thermally driven reactions, the reaction rates arising from plasmon-induced heating exhibit an exponential dependence on light intensity.<sup>49</sup> An example of this behavior was reported by Wang *et al.* when they used Au NPs supported on ZnO for the plasmon-assisted thermal hydrogenation of  $\text{CO}_2$  (Fig. 5a).<sup>103</sup> The authors not only found an exponential dependence of the CO conversion rate on laser intensity, but they could also tune product distribution to either CO or  $\text{CH}_4$  by just varying the intensity of the laser. Therefore, an exponential dependence of the reaction rate on illumination intensity is a characteristic feature of thermally driven transformations.

A linear dependence of the reaction rate on light intensity is, however, a signature of an electron-driven process and it is the most commonly reported mechanism for plasmonic reactions.<sup>49,104–106</sup> A first order relationship between the photon flux and reaction rate implies that a single photon absorption induces the chemical reaction. It is worth mentioning that both linear (rate  $\propto I$ ) and superlinear (rate  $\propto I^n$ ,  $n > 1$ ) regimes are distinctive features of electron-driven chemical transformations on metal surfaces, but the

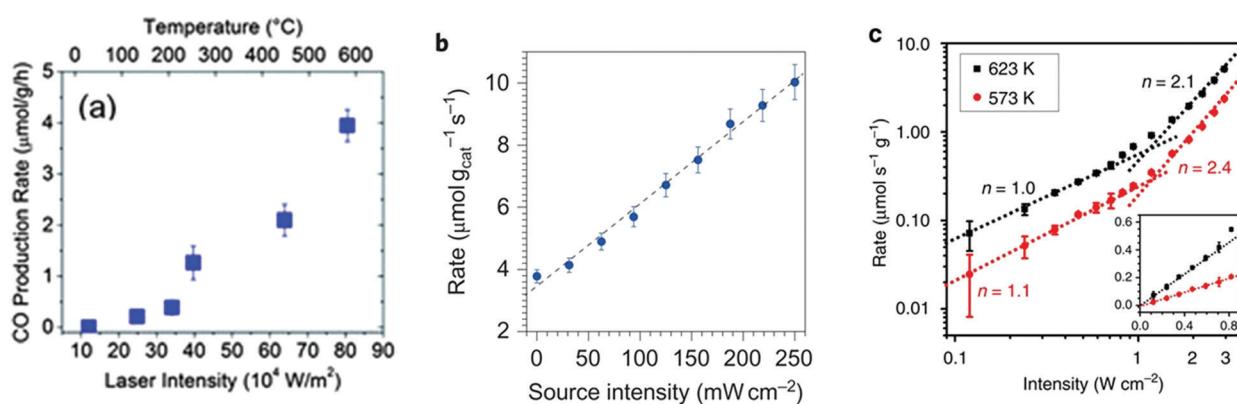


Fig. 5 (a) CO production rate as a function of continuous wave laser intensity at 532 nm using ZnO-supported Au NPs. Reproduced with permission from ref. 103. Copyright (2013) Royal Chemical Society. (b) The rate of photo-thermal ethylene epoxidation as a function of light source intensity. Reproduced with permission from ref. 49. Copyright (2011) Springer Nature. (c) Rate of  $\text{CH}_4$  photo-production as a function of ultraviolet light intensity at 623 (black squares) and 573 K (red circles). The intensity-dependent reaction rate shows a linear to super-linear transition with increasing light intensity. The inset shows the intensity-dependent reaction rate in the linear region. Reproduced with permission from ref. 108. Copyright (2017) Springer Nature.



operation conditions are the key factor determining the prevalence of each of them. For instance, Linic *et al.* observed a linear dependence of the photocatalytic rate of ethylene epoxidation on light intensity up to 300 mW cm<sup>-2</sup>, while at higher source intensities the dependence turned into the superlinear regime (Fig. 5b).<sup>49,107</sup> In a similar way, Zhang *et al.* described a linear dependence of the photocatalytic CO<sub>2</sub> hydrogenation on Rh/Al<sub>2</sub>O<sub>3</sub> heterostructures that became superlinear at the highest light intensities (Fig. 5c).<sup>108</sup> In this case, the authors attributed the transition to superlinearity to multiple excitations of the vibrational modes of adsorbed species by plasmon-induced hot electrons.

The reaction rate dependence on the photon wavelength typically correlates absorption properties of the photocatalyst with its catalytic performance and is of great importance for this section; it can provide further evidence for electron-driven processes. In this respect, an agreement between the absorption spectrum of the plasmonic photocatalyst and the wavelength dependent reaction rates corroborates an electron-driven photochemical transformation.<sup>54</sup> For instance, Ye and co-workers recently observed that the quantum efficiency for the photocatalytic steam reforming of CH<sub>4</sub> using TiO<sub>2</sub>-supported Rh NPs was consistent with the optical absorption spectrum of the catalyst, thus suggesting that the hot electrons from interband transition in Rh NPs were responsible for the catalytic activity.<sup>109</sup> In 2014, the same group reported the photo-thermal conversion of CO<sub>2</sub> with H<sub>2</sub> to produce methane (CH<sub>4</sub>) with group VIII nanocatalysts.<sup>110</sup> However, in this case the electron-driven reaction pathway was completely excluded since the photo-action spectrum using monochromatic light showed no dependence on the irradiation wavelength. In view of these findings, the authors concluded that the photo-methanation of CO<sub>2</sub> was mediated by a pure thermal effect. In the early 2020, Su and Xiao studied the visible light  $\alpha$ -alkylation of ketones with alcohols using a titanium nitride (TiN) photocatalyst.<sup>111</sup> The authors not only found that both action and optical absorption spectra were in good agreement, but also reported a positive relationship between reaction rate and light intensity. Altogether, the results certainly indicated that the reaction proceeded through an electron-driven mechanism operated by energetic charge carriers on TiN. The irradiation wavelength dependence of the reaction rate was also evaluated for O<sub>2</sub> activation by different plasmonic Au-nanorods@Pd superstructures.<sup>112</sup> In this work, the authors observed a clear matching between the photocatalytic performance and the visible-NIR spectrum of the photocatalyst, suggesting that plasmon-induced hot electrons drove the reaction mechanism by electron injection into the anti-bonding orbital of molecular oxygen. Furthermore, on the basis of theoretical calculations the thermal effect was found to be negligible, so the high photocatalytic activity was mainly ascribed to the large field-enhancement that maximized plasmon absorption and boosted hot electron generation in Au-nanorods@Pd heterostructures. In spite of the fact that all these studies are insufficient to distinguish between the direct and indirect electron transfer pathways, the combination of wavelength dependent experiments with a linear relationship of the reaction rate with light intensity provides convincing evidence of the hot electron-induced reaction mechanism.

In a recent review, Baffou and co-workers proposed a series of simple experimental procedures to discriminate and quantify the different thermal and non-thermal contributions of plasmon-driven chemical reactions.<sup>113</sup> One of these methodologies is based on the variation of the light beam diameter using two approaches: a constant irradiance or a constant power. In the first case, the photon flux is proportional to the beam diameter, whereas in the second case the amount of photons remains constant. Both thermal and non-thermal contributions typically present a proportional dependence between reaction rate and the area of light beam under constant irradiance conditions. However, in the constant-power regime, the rate of electron-driven processes is independent of the beam diameter, while thermal-driven reactions show an inverse relationship between rate and beam radius. The study of these different behaviors by varying the illumination diameter appears to be an efficient means to elucidate the underlying reaction mechanism.

The activation energy ( $E_a$ ) is a key parameter in catalysis that provides insightful information about the catalytic performance of a material and the mechanistic pathway driving the reaction.<sup>114</sup> Because of that, a straightforward method to ascertain if photon-induced charge carriers contribute to the enhancement of the reaction rate consists in comparing activation barriers under both light and dark conditions at constant temperature. The reduction in the value of  $E_a$  under light irradiation compared to dark conditions is indeed typical for a hot carrier-driven mechanism.<sup>108,109,115,116</sup>

Another distinctive feature of electron-driven processes can arise from isotopic labelling experiments. Traditionally, the isotopic labelling of reactants has been a useful analytical probe to elucidate the reaction mechanisms in thermal catalysis through the measurement of the kinetic isotope effect (KIE). The KIE is obtained by measuring the change in the reaction rate due to the introduction of labelled reactants, as heavy isotopes require a higher energy input to reach the transition state. Interestingly, an electron-driven process displays a greater KIE under light irradiation than the thermally driven counterpart under dark conditions, so the comparison of both values can be a convenient methodology to differentiate between thermally and electron-driven reactions. For example, Ye *et al.* measured the KIE for photocatalytic CH<sub>4</sub> steam reforming over a Rh/TiO<sub>2</sub> catalyst by comparison of the H<sub>2</sub> production rates using either CH<sub>4</sub> or CD<sub>4</sub> as a reactant.<sup>109</sup> The authors found that the KIE under visible light irradiation (KIE<sub>light</sub> = 1.69) was larger than the one obtained under dark conditions (KIE<sub>dark</sub> = 1.55). Consequently, it was concluded that photo-induced hot electrons from interband transitions of Rh were responsible for the enhanced catalytic activity. In a similar way, for the photocatalytic ethylene epoxidation Linic and co-workers used <sup>16</sup>O<sub>2</sub> and isotopically labelled <sup>18</sup>O<sub>2</sub> to measure the KIE under both illumination and dark conditions.<sup>49</sup> Experiments confirmed that the activation of O<sub>2</sub> determined the overall reaction rate and the considerably larger KIE reported for the photocatalytic process suggested an electron-driven reaction mechanism.

Alternatively, the analysis of reaction selectivity can be another useful technique to assess the dominating reaction pathway. In thermally driven reactions, energy is provided to all available degrees of freedom including vibrational, rotational



and translational states, leading to the simultaneous activation of a plethora of reaction pathways.<sup>117</sup> In contrast, electron-driven processes offer the opportunity to modulate reaction yield towards the desired product by selectively injecting hot carriers to specific adsorbate–metal bonds.<sup>108</sup> In this respect, Christopher and co-workers investigated the photocatalytic CO oxidation in a H<sub>2</sub> rich stream over Al<sub>2</sub>O<sub>3</sub>-supported Pt NPs.<sup>118</sup> The results showed that the resonant photoexcitation of Pt–CO bonds arising from the hybridization of adsorbate–metal states enhanced the selective oxidation of CO to produce CO<sub>2</sub> over the competing H<sub>2</sub> oxidation reaction. Accordingly, the authors proposed that the specific photoactivation of metal–adsorbate bonds could be a good methodology to control reaction selectivity in electron-driven plasmonic reactions. Cronin *et al.* studied the photoreduction of CO<sub>2</sub> with H<sub>2</sub>O vapor using Au/TiO<sub>2</sub> heterostructures and found that the product distribution was strongly dependent on the irradiation wavelength.<sup>74</sup> While the excitation at 532 nm mostly produced CH<sub>4</sub>, the use of UV light of 254 nm yielded the generation of other products including ethane, methanol and formaldehyde. It was rationalized that depending on the irradiation wavelength it was possible to trigger different reaction pathways to drive the chemical reaction towards certain product distribution. Interestingly, any product was observed when the reaction was performed at 400 °C under dark conditions, thus excluding any thermal contribution to the reaction mechanism.

Transient absorption spectroscopy (TAS) is an effective methodology to study the dynamics of electron transfer processes. Consequently, many groups have used TAS to elucidate the mechanism of light-mediated reactions. For instance, Ye and co-workers determined an ultrafast electron transfer from Rh NPs to TiO<sub>2</sub> under photoexcitation at 450 nm, attributable to the injection of hot electrons across the Schottky barrier at the metal–oxide interface (Fig. 6).<sup>109</sup> Importantly, the authors found that the electron transfer event occurred in the time scale of hundreds of

femtoseconds, while the transient signal showed a decay in the order of thousands of nanoseconds, thus indicating an effective charge carrier separation that prevented undesirable electron–hole recombination. In view of this, the authors concluded that the observed photo-enhancement of methane reforming using a Rh/TiO<sub>2</sub> photocatalyst was due to excited hot carriers, rather than thermal effects. In the same way, Garcia and Corma performed transient spectroscopy measurements to study the reaction mechanism of photocatalytic CO<sub>2</sub> hydrogenation to CH<sub>4</sub> by Ni/SiO<sub>2</sub>–Al<sub>2</sub>O<sub>3</sub>.<sup>119</sup> Experiments allowed the detection of transient signals under excitation at 355 and 532 nm that confirmed the presence of charge separation states. Since the reaction temperature was monitored at 150 °C, therefore insufficiently low to perform CO<sub>2</sub> methanation, it was assumed that the reaction was driven by H<sub>2</sub> activation through the formation of Ni–H species. A more recent work from the same group also used TAS measurements to demonstrate that NiO/Ni NPs supported on graphene could hydrogenate CO<sub>2</sub> to CH<sub>4</sub> by photo-excitation of electron–hole pairs.<sup>120</sup>

### 3.4.2. Nanoscale thermal measurements

**3.4.2.1. Non-luminescence methods.** Another approach to differentiate between the photochemical and the thermochemical contribution to photo-thermal reactions implies the direct temperature measurement at the surface active sites. Although carrying out this type of measurements is still very challenging, recent advances in metrology have allowed great improvements in nanoscale thermometry. For instance, scanning thermal microscopy equipped with nanoscale probe tips renders surface temperature measurements with spatial resolutions in the range of 10 nm and ~10–50 mK precision.<sup>121,122</sup> It is also possible to obtain local temperature maps with nanometer-scale resolution by means of the analysis of the temperature-dependent energy shift of plasmon peaks using scanning transmission electron microscopy (STEM).<sup>123</sup> More recently in 2018, Idrobo *et al.* performed local temperature assessments in boron nitride nanoflakes by calculating the ratio between the gain and loss phonon peaks in the

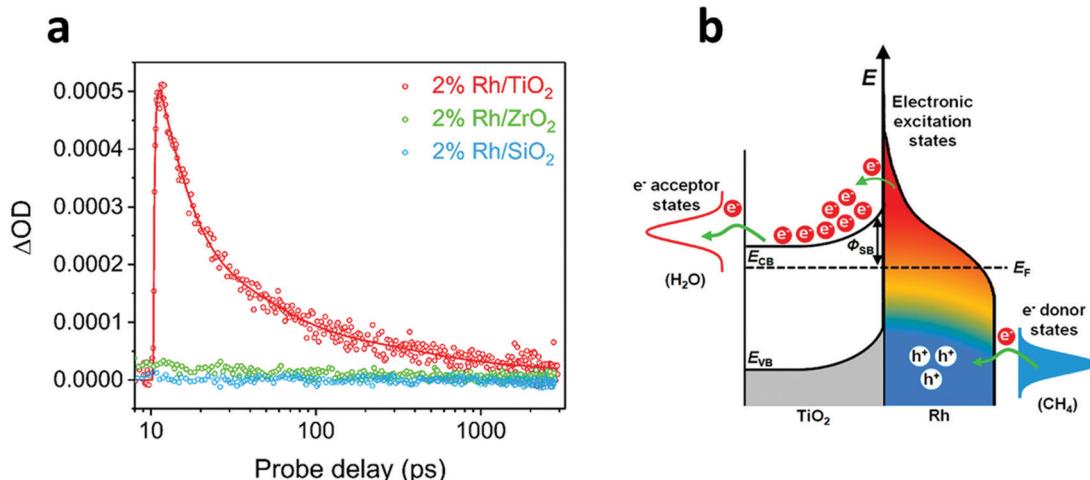


Fig. 6 (a) Transient absorption kinetics at 5000 nm of Rh/TiO<sub>2</sub>, Rh/ZrO<sub>2</sub> and Rh/SiO<sub>2</sub> under laser excitation of 450 nm. (b) Mechanism of hot carrier-enhanced steam methane reforming. Visible light-excited hot carriers can be separated quickly at the Rh/TiO<sub>2</sub> interface; the positively charged Rh surface favours C–H bond activation, whereas electron-rich states at the interface promote the reduction process. Reproduced with permission from ref. 109. Copyright (2018) American Chemical Society.



electron energy spectrum using monochromated aberration-corrected STEM.<sup>124</sup>

**3.4.2.2. Luminescence methods.** IR thermography is based on the Planck blackbody emission principle and estimates the temperature profile of a body in relation to its emitted energy. Although local and global temperatures can be obtained by thermal mapping with IR cameras in combination with bulk thermocouple measurements, the accuracy of these methodologies is still controversial due to different interpretations of the emissivity of catalysts and the parameter choice of thermal imaging cameras.<sup>100,125,126</sup> It is also worth mentioning that under irradiation of thick catalyst beds, thermal gradients may exist due to the limited penetration ability of light and non-uniform heat diffusion. In these cases, the measurement from IR cameras only represents the surface temperature of the catalyst, that is commonly higher than the temperature inside the catalyst bed.<sup>125</sup> For these reasons, IR camera measurements must be precisely recorded and attention has to be paid in order to avoid misinterpretations in the thermal contribution to photo-thermal systems.

Thermoreflectance and optical interferometry are luminescence techniques that use the reflection of light to obtain thermal measurements. Thermoreflectance relies on the relationship between the refractive index of a material and its temperature, giving rise to temperature profiles in the micrometer/submicrometer scale with high thermal and temporal

resolutions typically in the order of  $10^{-2}$  K and  $10^{-1}$   $\mu$ s, respectively.<sup>127</sup> In spite of the fact that this technique has been widely implemented in microelectronics to study electrical self-heating, Wang *et al.* recently used different thermoreflectance techniques to evaluate the thermal enhancement of Au/Al<sub>2</sub>O<sub>3</sub> plasmonic structures with outstanding spatial and temporal resolutions in the order of 100 nm and several nanoseconds.<sup>128</sup> As for optical interferometry, it renders local temperature and thermal expansion measurements with very high temperature resolutions of  $\sim 10$   $\mu$ K, but in this case the spatial resolutions are below the nanoscale ( $\sim 1$   $\mu$ m).<sup>127</sup>

Apart from the previously described techniques, tip-enhanced Raman spectroscopy (TERS) also enables mapping local temperatures at the nanometer scale by virtue of calculation of the ratio of the intensities of the anti-Stokes and Stokes Raman signals.<sup>129</sup> In fact, the analysis of the observed linear shift in the position of Raman bands as a function of temperature can be a distinctive feature to characterize thermally driven photo-thermal reactions as well.<sup>130,131</sup> Ozin and co-workers recently reported an appealing example of this effect using Pd NPs supported on Nb<sub>2</sub>O<sub>5</sub> in the photocatalytic reduction of CO<sub>2</sub> (Fig. 7).<sup>132</sup> The authors observed a shift in the position of the niobium–oxygen phonon modes as a function of the excitation laser power, hence revealing a photo-thermal phenomenon. This heating effect was probed through a study of the Stokes and anti-Stokes Raman bands and ultimately allowed extracting the local temperature of the catalyst as high as 470 °C. According to the available data, Pd NPs acted as

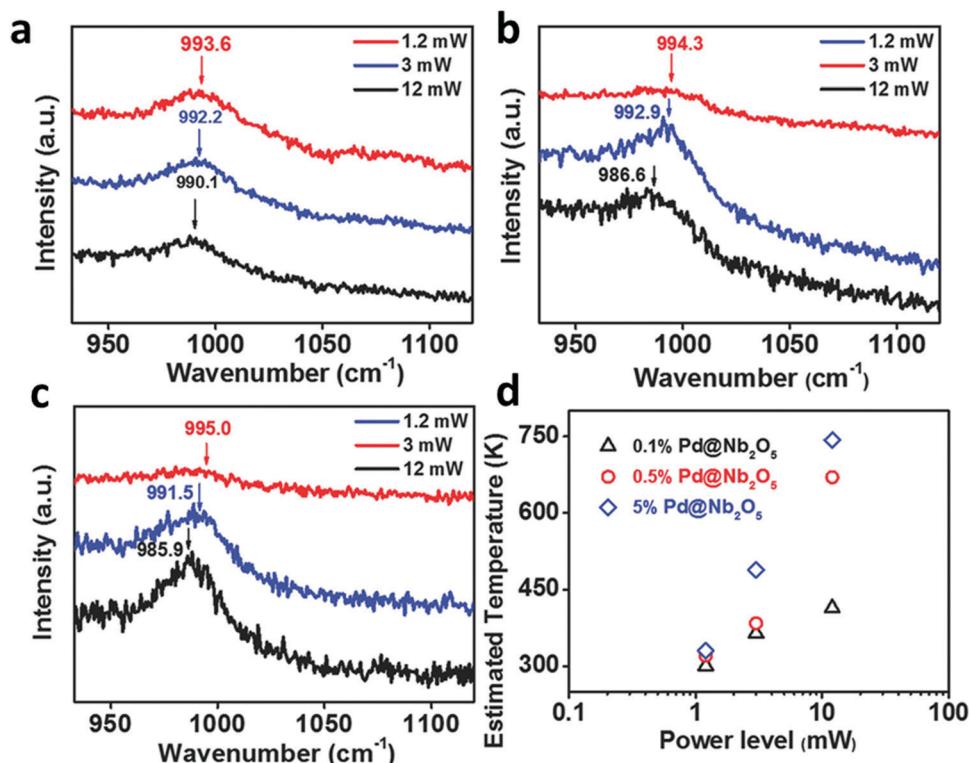


Fig. 7 Dependence of the Raman frequency for Nb–O stretching vibrations of (a) 0.1% Pd@Nb<sub>2</sub>O<sub>5</sub>, (b) 0.5% Pd@Nb<sub>2</sub>O<sub>5</sub>, and (c) 5% Pd@Nb<sub>2</sub>O<sub>5</sub> at different incident light power levels. (d) Estimated temperatures of 0.1%, 0.5%, and 5% Pd loaded onto Nb<sub>2</sub>O<sub>5</sub> nanorods at different incident light power levels. Reproduced with permission from ref. 132. Copyright (2017) John Wiley and Sons.



photo-thermal “nano-heaters” that effectively elevated the local temperature of the  $\text{Nb}_2\text{O}_5$  support, which drove the  $\text{CO}_2$  hydrogenation reaction.

It is well established that the fluorescence signal strongly depends on temperature, since its intensity decreases and the peaks shift to longer wavelengths when the temperature increases. By taking advantage of this effect, nanoscale fluorescence thermometry can provide temperature measurements with nanometric spatial resolution using quantum dots (QDs) or organic dyes as temperature probes.<sup>133–136</sup> For instance, Garcia *et al.* demonstrated the potential of this strategy to assess the reaction mechanism in the photocatalytic  $\text{CO}_2$  reduction by  $\text{Cu}_2\text{O}$  NPs supported on graphene (Fig. 8).<sup>137</sup> In this work, the authors used commercial core–shell  $\text{CdSe}@\text{ZnS}$  QDs as local thermometers and measured their emission lifetime as a function of temperature. Interestingly, the authors found that the irradiation of samples containing these QDs in contact with  $\text{Cu}_2\text{O}$  NPs did not lead to significant changes in the QDs’ emission lifetime, thereby excluding a light-induced heating effect. In view of these findings, the authors concluded that the most likely reaction mechanism was the photo-generation of electron–hole pairs rather than a thermal activation. For further details regarding nanoscale thermal probing methodologies, the readers are referred to an excellent review by Wang *et al.*<sup>138</sup> Carlos and Palacio’s groups also performed a comprehensive revision of thermal measurement techniques at the nanoscale.<sup>127</sup>

**3.4.2.3. Computational methods.** Finally, besides the aforementioned experimental methodologies, computational methods can aid elucidation of the mechanisms underlying photo-thermal reactions. In this regard, the density functional theory (DFT) framework has proven to be a reliable model to study electronic structures in metal–semiconductor interfaces or adsorbed species in surface active sites.<sup>86,108,109,139</sup> A remarkable example of the use of this model to assess the reaction pathway in  $\text{Rh}/\text{TiO}_2$  plasmonic nanostructures for  $\text{CH}_4$  activation was recently reported by Ye *et al.*<sup>109</sup> Theoretical calculations suggested that under visible light excitation, hot electrons in Rh 4d orbitals could transfer to unoccupied 3d orbitals in Ti, so preventing them from recombination within the

metal and favoring the generation of electron-deficient sites ( $\text{Rh}^{\delta+}$ ) in the Rh surface. Consequently, these  $\text{Rh}^{\delta+}$  sites were more prone to accept electrons from  $\text{CH}_4$ , thus facilitating C–H bond cleavage and promoting  $\text{CH}_4$  activation. Liu and Everitt used DFT calculations as well to elucidate how hot electrons interacted with intermediate species in the photocatalytic  $\text{CO}_2$  hydrogenation by Rh nanocubes supported on  $\text{Al}_2\text{O}_3$ .<sup>108</sup> This study revealed that anti-bonding orbitals from CHO intermediates could accept photo-generated electrons from Rh, thus weakening C–O bonds and assisting  $\text{CH}_4$  formation. Curiously, it was also found that the photo-injection of hot electrons into anti-bonding orbitals of the CO intermediate was far unlikely, somehow explaining the reaction selectivity towards  $\text{CH}_4$  production rather than CO formation.

It is worth reminding that all these methodologies can help to distinguish between photochemical and thermal contribution to photo-thermal conversion and, in the best-case scenario, determine which one is predominant in the system under study. Nonetheless, it is possible and very common indeed that both mechanisms coexist together and operate synergistically to enhance the overall photocatalytic rate. For example, Willets *et al.* quantified the relative contribution of hot carriers and heating to plasmon-mediated photoelectrochemical reactions based on scanning electrochemical microscopy (SECM) measurements.<sup>140</sup> The authors determined that both mechanisms contributed to the photoelectrochemical performance, with a different dependence on excitation intensity. Another example of successful exploration and discrimination of thermal and non-thermal activities working synergistically in the photocatalytic activity of Rh NPs supported on  $\text{TiO}_2$  was recently provided by Zhang and co-workers.<sup>141</sup> Based on both experimental measurements and theoretical models, it was found that heat and light worked not only simultaneously, but also cooperatively, and that a temperature increase gave rise to an improvement of the plasmon-enhanced catalysis. In a related precedent from the same group, the authors introduced a novel indirect illumination technique to distinguish between thermal and non-thermal effects on the photocatalytic  $\text{CO}_2$  methanation by  $\text{Rh}/\text{TiO}_2$ .<sup>142</sup> This technique makes use of an inactive layer of  $\text{Ti}_2\text{O}_3$  (a well-known photo-thermal heater) on the top of the  $\text{Rh}/\text{TiO}_2$  catalyst. This  $\text{Ti}_2\text{O}_3$  layer effectively absorbs

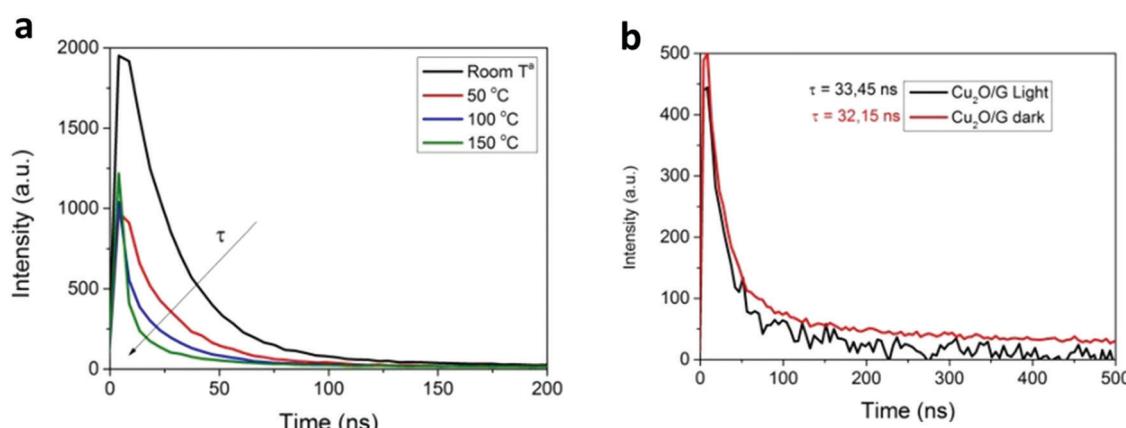


Fig. 8 (a) Time correlated emission decay of commercial  $\text{CdSe}@\text{ZnS}$  QDs as a function of temperature. (b) Emission lifetime of thin films of  $\text{CdSe}@\text{ZnS}$  QDs supported on a  $\text{Cu}_2\text{O}/\text{graphene}$  film under light and dark conditions. Reproduced with permission from ref. 137. Copyright (2017) Royal Society of Chemistry.



light and produces heat, displaying a similar temperature profile to the one achieved from the direct illumination of the Rh/TiO<sub>2</sub> catalyst. Under these indirect illumination conditions, the hot carrier contribution derived from light is completely extracted, hence allowing a solid distinction between thermal and non-thermal effects on the photocatalytic rate. In 2019, Zhan *et al.* studied the photoelectrochemical response of plasmonic Au electrodes and quantitatively disentangled the thermal from the hot carrier effect.<sup>143</sup> To do this, the authors separated the overall plasmonic photocurrent into a rapid response current (RRC) and a slow response current (SRC) attributable to photo-electronic and photo-thermal effects, respectively. The results suggested that both effects influenced the reaction and the different timescales for hot carrier generation and heating effect served to assess each of the contributions to the overall reaction. For a deeper discussion of the synergy between thermal and non-thermal effects on plasmon-mediated photocatalysis the readers are referred to an excellent review by Liu and Everitt.<sup>144</sup>

## 4. Photo-thermal applications in catalysis

### 4.1. CO<sub>2</sub> conversion

The possibility of producing synthetic fuels or chemicals through the reaction between CO<sub>2</sub> and H<sub>2</sub>O (or indirectly between CO<sub>2</sub> and H<sub>2</sub> derived from H<sub>2</sub>O using green energy sources) is among the most promising alternatives to achieve CO<sub>2</sub> neutrality in transportation and the chemical industry. In general, the reduction of CO<sub>2</sub> can produce different chemicals and fuels, CO, CHOOH, CH<sub>3</sub>OH or CH<sub>4</sub> among others, depending on the number of electrons that come into play. Over the last decade, the use of photothermal catalysis to speed up these processes has grown exponentially.

**4.1.1. Artificial photosynthesis.** Artificial photosynthesis is described as the most direct approach to reduce CO<sub>2</sub> and it tries to replicate the natural photosynthesis, one of the most important reactions on Earth.<sup>145</sup> The perfect route will employ CO<sub>2</sub> and water as reagents in a significant carbon recycling, producing hydrogen or carbon-based fuels by using ideally sunlight as a sustainable energy source.<sup>146,147</sup> This process is divided into two main reactions, the water splitting reaction that produces H<sub>2</sub>, and CO<sub>2</sub> reduction, which produces carbon-based products such as CO, CH<sub>3</sub>OH or CH<sub>4</sub>.<sup>148</sup> The second reaction is much more challenging than the first one due to its requirement of the conduction band edge and multi-electron processes, depending on the product. On this basis, the main limitations of the photocatalytic methodology are the relatively low charge separation efficiency, light harvesting and reduction selectivity.<sup>149,150</sup> However, the cooperation between photo and thermocatalytic processes could be a solution.<sup>151</sup>

Some of the most used semiconductor oxides for the photocatalytic CO<sub>2</sub> reduction are TiO<sub>2</sub>, ZnO, CdS, Fe<sub>x</sub>O<sub>y</sub> or WO<sub>3</sub> among others. All of them share properties such as low or non-toxicity, low cost and powerful redox properties. For the application of these photocatalysts in CO<sub>2</sub> reduction processes, these materials should present high light harvesting, low charge recombination rates and easy activation of CO<sub>2</sub>. In order to improve some of these

properties, the creation of superficial oxygen vacancies is one of the most suitable and used methods, usually by means of H<sub>2</sub>-treatments. The incorporation of these O-defects improves the adsorption of reagents and can also enhance the separation of photo-generated charges.<sup>152</sup> Additionally, most of these semiconductor oxides are stable at high temperatures, permitting their adaptation to photo-thermal reaction conditions. To improve the photo-thermocatalysis, the same semiconductors have been functionalized by different metals or metallic oxide NPs. The use of these composites improves the yield of the process by combining the individual roles of each component in the overall reaction mechanism (*e.g.* photo-generated charges at TiO<sub>2</sub> under UV light + H<sub>2</sub> adsorption in metallic NPs such as Pd) or even, by means of cooperation between them (*e.g.* metallic NPs can reduce the recombination of photo-generated charges of semiconductors, by transfer of e<sup>-</sup>/h<sup>+</sup>). Moreover, most of the composites are designed in order to optimize the absorption of sunlight and exploit the full solar spectrum as a sustainable energy source.

Since the 70s, TiO<sub>2</sub> has been one of the most used semiconductors in photocatalysis.<sup>153,154</sup> Pure photocatalytic studies using this semiconductor with oxygen vacancies as a catalyst under artificial photosynthesis conditions have demonstrated a high CH<sub>4</sub> selectivity. For example, Yin *et al.* produced 16.2 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup> and 2.7 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>, with a CH<sub>4</sub> selectivity of 79 and 73% under solar light and visible light, respectively.<sup>155</sup> However, the introduction of thermal energy in the photocatalytic process modifies the production rates and the main product. One of the first studies evaluating the use of TiO<sub>2</sub> in photo-thermal catalysis was published in 1997.<sup>156</sup> TiO<sub>2</sub> was evaluated as an “artificial leaf” for the photocatalytic reduction of CO<sub>2</sub> with H<sub>2</sub>O. The principal results showed CO, CH<sub>4</sub> and C<sub>2</sub>H<sub>6</sub> as the main products. However, the most relevant result was the increase of the CH<sub>4</sub> rate when the temperature of the system was increased from 25 to 200 °C. The authors attributed this trend to the different physical adsorption/desorption capacities of the catalyst at each temperature and not to a clear cooperation between photo and thermocatalytic mechanisms. Anyway, this study can be considered as one of the first examples of photo-thermocatalysis.

DFT studies using supported MnO<sub>x</sub> NPs on TiO<sub>2</sub> revealed that thermocatalysis aids photocatalysis by creating oxygen vacancies and modifying CO<sub>2</sub> surface bindings, which results in an increased CO yield.<sup>157</sup> Wang *et al.* employed TiO<sub>2</sub> nanosheets alone and TiO<sub>2</sub> nanosheets doped with CuS, for improving the light utilization efficiency, even in the IR, due to the narrow band gap energy of CuS. The catalytic results at room temperature using different light ranges revealed that the 2% CuS/TiO<sub>2</sub> catalyst reduced CO<sub>2</sub> to CO as the main product and the use of UV-vis-IR light achieved the best production rate (26 µmol CO g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>) increasing the temperature to 138 °C.<sup>158</sup> The use of full spectrum light and cooling to 102 °C resulted in a reduction in the CO rate by almost 4 times, evidencing the photo-thermal cooperation. Similarly, metallic NPs (Ru and Au 1 wt%) were supported over TiO<sub>2</sub> by an impregnation-reduction method.<sup>159</sup> The UV-visible diffuse reflectance spectra uncovered that the incorporation of Au and Ru did not modify the absorption related to the band gap of the semiconductor, although Au-incorporation enhanced the absorption



profile due to the LSPR at 570 nm. This LSPR absorption was shifted to 580 nm when Ru was incorporated, revealing a Au–Ru interaction. Photo-thermal experiments were carried out in a fixed-bed reactor at different temperatures (50–150 °C). The results demonstrated that, under temperature and light illumination, H<sub>2</sub> and CH<sub>4</sub> were formed simultaneously by water splitting coupled to CO<sub>2</sub> reduction, obtaining the best results at 85 °C (0.3 μmol H<sub>2</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup> and 27.1 μmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>). Control experiments without the catalyst or without irradiation did not produce any reduced product, concluding that the reaction is principally driven by photocatalysis. Briefly, the photocatalysis of water splitting reaction generates H<sub>2</sub> that is able to react with CO<sub>2</sub> in a thermocatalytic process over Au–Ru/TiO<sub>2</sub>. Besides, the rapid consumption of this generated H<sub>2</sub> facilitates the migration of photo-generated electrons, minimizing the recombination of photo-induced carriers (Fig. 9).

As a way to create more sustainable processes, sunlight must be considered as the main energy source (photo and thermocatalytically). Solar light is the largest sustainable source of energy: renewable, clean, non-toxic and non-polluting, but also, almost 50% of the total sunlight is included in the range of visible light.<sup>160,161</sup> This fact requires finding active photocatalysts that can efficiently harvest visible light.<sup>160</sup> In this spirit, plasmonic metal NPs are one of the best options to enhance visible light absorption through plasmon resonance.<sup>162,163</sup> Kumar *et al.* optimized a catalyst based on Au and Pt NPs on a SiO<sub>2</sub> shell (Pt–AuNPs@SiO<sub>2</sub>) for the reduction of CO<sub>2</sub> with H<sub>2</sub>O.<sup>164</sup> The outcomes revealed an efficient, stable and recyclable photo-catalyst under visible-IR light at room temperature, with HCOOH as the main product along with traces of CH<sub>3</sub>OH. When the tests were carried out without temperature control, the temperature increased after 30 min from 25 to 85 °C, together with an increase of quantum yield up to 4.32%. These values confirmed the influence of the photo-thermal effect due to the irradiation of plasmonic nanomaterials using only visible and IR light. However, when comparing the formic acid yield at

different irradiation wavelengths, the obtained profile was similar to the absorption spectrum of the catalyst. This equivalence is clear evidence of a mechanism directed mainly by the photo approach.

In the same direction, other studies focused on the incorporation of metallic NPs with LSPR into semiconductors, in order to generate hot charge carriers for CO<sub>2</sub> photo-thermoreduction *via* plasmonic effects.<sup>74</sup> This functionalization permits light absorption in a wide range of the spectrum, principally visible light, and also modifies the band gap of the semiconductor through the creation of a Schottky barrier. Recently, Hongjia Wang *et al.* studied the catalytic activity of Au/TiO<sub>2</sub>.<sup>165</sup> This Au incorporation improved the harvest of light (LSPR absorption from Au at 550 nm) and increased about two times the CO rate (~20 μmol g<sub>cat</sub><sup>-1</sup>) and CH<sub>4</sub> rate (~4 μmol g<sub>cat</sub><sup>-1</sup>) after 4 hours compared to the use of bare rutile-TiO<sub>2</sub> under the full light spectrum. Moreover, the catalytic activity of Au/TiO<sub>2</sub> under different irradiation conditions was proportional to the reached temperature: full spectrum (60 °C), IR cut-off filter (30 °C) and the one cooling down (15 °C). This similar trend suggests a temperature dependence of the catalytic process. Finally, two tests with Au/TiO<sub>2</sub> at the same temperature were carried out using full spectra and IR cut-off irradiation + heating bath. The outcomes revealed a higher CO rate and a slightly increased CH<sub>4</sub> rate in the full spectra experiment demonstrating that the reaction is boosted by IR (sunlight full spectra) and photo-thermal effects produced by LSPR of Au NPs.

Despite these results, the incorporation of temperature in a photocatalytic system produces some general alterations in the reaction: one benefit, the temperature enhances water splitting due to the endothermicity of this reaction; and two drawbacks, the recombination of hot carriers increases due to the thermal motion of photo-generated species and CO<sub>2</sub> hydrogenation, being an exothermic reaction, is not thermodynamically favored. In order to minimize these drawbacks, disordered TiO<sub>2</sub> was employed by Yu *et al.*, reducing the charge carrier recombination and increasing the superficial oxygen defects. The disordered semiconductor was doped with Pt NPs, improving the visible light harvesting and increasing the dissociative adsorption of H<sub>2</sub>, aiding the reduction of CO<sub>2</sub>.<sup>166</sup> The study always gave better results in the photo-thermocatalytic experiments (sunlight + 120 °C) than in photocatalytic tests (sunlight + 25 °C), giving the first evidence of the synergy between photo and thermal processes. The highest CH<sub>4</sub> rate was achieved using Pt/disordered-TiO<sub>2</sub> (17.1 μmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>) with a selectivity of 87.5%. Thermal H<sub>2</sub> decomposition studies were carried out in dark conditions observing the highest H<sub>2</sub>-decomposition rate with the Pt/disordered-TiO<sub>2</sub> catalyst, and then the best promotion of H<sub>2</sub> splitting aiding the formation of CH<sub>4</sub>. In short, the incorporation of oxygen vacancies and disorder in the semiconductor, linked to the deposition of Pt NPs, are the most important variables to promote CH<sub>4</sub> production, at least in the photo-thermocatalytic process.

Although TiO<sub>2</sub> is the most studied, other semiconductors have also been evaluated for CO<sub>2</sub> photo-thermocatalytic reduction using H<sub>2</sub>O. One of these examples is ZnO. Guo *et al.* evaluated initially the photocatalytic activity of CO<sub>2</sub> artificial photosynthesis under UV and visible light at 25 and 200 °C using two different shapes of

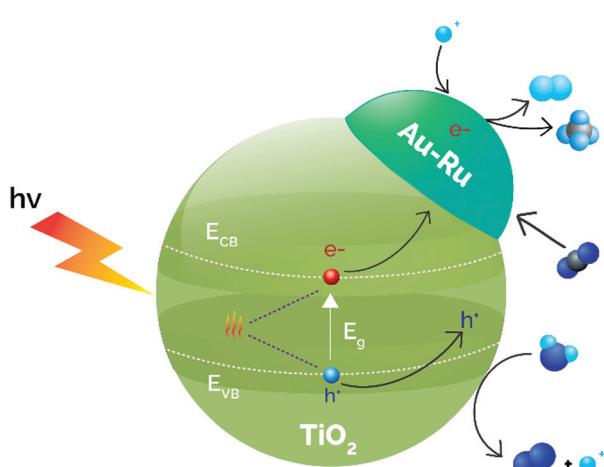


Fig. 9 Synergistic effects between photocatalysis and thermocatalysis in the direct thermophotocatalytic reduction of CO<sub>2</sub>–H<sub>2</sub>O over Au–Ru/TiO<sub>2</sub>. Adapted with permission from ref. 159. Copyright (2017) John Wiley and Sons.



ZnO, rods and belts, alone or in layered-double-hydroxides (LDHs).<sup>167</sup> The effects of the shape of the photocatalysts were observed and related to the electronic structure and surface defects, but the most important was the enhancement of production rates in all the evaluated photocatalysts when the temperature was increased to 200 °C. Besides, core–shell ZnO@LDHs (LDHs are 2D layered Cu–Zn–Al materials) were evaluated in the same photo-thermal conditions obtaining the best results with the belt shaped ZnO@LDHs (11.4 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>). The composite or hetero-junction permitted an efficient transfer of photo-generated charges that minimized the recombination and also increased the photocatalyst surface area, improving the CO<sub>2</sub> reduction. Again, the cooperation between light and heat increased the catalytic performance although the authors did not study the contribution of each of the energy sources to the overall reaction mechanism.

Other choice of a photo-thermocatalyst is the versatile, low cost, and non-toxic MoO<sub>3</sub>. This oxide was used by Li *et al.*, tailoring it (MoO<sub>3-x</sub>) with extra superficial defects (O-vacancies).<sup>168</sup> The introduction of O-defects extended the absorption of light permitting one to employ the whole spectrum of sunlight by means of the LSPR effect that narrowed the band gap of the material. Moreover, the increase of superficial defects improved the surface area and the available active sites, but also oxygen vacancies are responsible for decreasing the recombination of charge carriers. Catalytic tests using specific range of lights without any external heating system modified the activity as UV-vis-IR > IR > UV-vis, and also the reached temperature followed the same order from 160 to 105 °C, evidencing a high influence of thermal catalysis in the process. However, pure thermal reactions without irradiation were less active than experiments at the same temperature using light, and thus, photocatalysis played a role in the reaction too. Finally, synergistic photo-thermocatalytic artificial photosynthesis using MoO<sub>3-x</sub> produced the highest reaction rates, obtaining 41.2 µmol CO g<sub>cat</sub><sup>-1</sup> and 8.3 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> using the full sunlight spectrum (UV-visible-IR).

Another example is WO<sub>3</sub>, a stable and non-toxic semiconductor able to harvest 12% of sunlight while offering a low conduction band that facilitates recombination. Wang *et al.* synthesized this semiconductor over a mesoporous template (KIT-6) and this was employed as a catalyst to reduce CO<sub>2</sub> with H<sub>2</sub>O under visible light and/or temperature (250 °C).<sup>169</sup> The mesoporous catalysts were pre-reduced under different temperatures controlling the amount of oxygen vacancies (related to the presence of W<sup>5+</sup>) and increasing lattice expansion and surface area compared to commercial WO<sub>3</sub>. The characterization of the diverse mesoporous WO<sub>3</sub> (m-WO<sub>3</sub>) revealed that oxygen vacancies improved the visible-light absorption as well as narrowed the band gap. Accordingly, the higher the amount of oxygen vacancies, the higher the catalytic performance. The principal product was CH<sub>4</sub> regardless of the used conditions; meanwhile CH<sub>3</sub>OH was formed only under thermal and photo-thermal catalysis conditions (<3 µmol CH<sub>3</sub>OH g<sub>cat</sub><sup>-1</sup>). Oppositely, the photocatalytic route did not produce CH<sub>3</sub>OH due to the insufficient energy of hot electrons to reduce CO<sub>2</sub> to the simplest carbon alcohol. Moreover, the different amount of oxygen defects at the catalysts using different reduction temperatures drove to different product distribution. The m-WO<sub>3</sub> reduced at 550 °C

presented the highest number of O-defects and also presented the highest CH<sub>4</sub> production rate, specifically 0.15, 21.42 and 25.77 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> after 12 h under photo, thermo and photo-thermal catalysis respectively. Examining these data, the extra 4.35 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> after 12 h in the photo-thermal catalysis process compared to the thermocatalytic process was not just an addition of the other individual photo-processes. There was a clear cooperation between the two processes. The authors assigned this improvement to the fact that the temperature of the catalytic surface increased faster when the system was heated and irradiated. This intensified the photocatalytic performance because the electron excitation and relaxation were easier, increasing the production rate. Moreover, the m-WO<sub>3</sub> reduced at 250 °C presented less oxygen vacancies and produced 6.0 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> and almost 10 µmol CH<sub>3</sub>OH g<sub>cat</sub><sup>-1</sup> after 12 h under photo-thermal catalytic conditions. These outcomes disclosed that oxygen vacancies have an influence on the reaction mechanism and depending on the amount of these defects, the product distribution varies.

The same research group incorporated Mo NPs into the mesoporous WO<sub>3</sub> (m-Mo/WO<sub>3</sub>), as a strategy to improve the photo-thermocatalytic performance by enhancing light harvesting, reducing the recombination of photo-generated charges and/or incorporating new catalytic active sites.<sup>170</sup> Under the same experimental conditions, CH<sub>4</sub> and CH<sub>3</sub>OH were the only products and the production rates were significant only when photo-thermal catalysis was carried out using m-Mo/WO<sub>3</sub> after 12 h (5.96 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> and 13.80 µmol CH<sub>3</sub>OH g<sub>cat</sub><sup>-1</sup>). The authors reported a photo-thermal effect due to the better production rate when heating and irradiation were combined. The results highlighted how mesoporosity and Mo-doping increase the photocatalytic behavior by the reduction of charge carrier recombination due to the high crystallinity and ordered porosity together with the increase of the separation of photo-generated charges when Mo acts as cocatalyst.<sup>169</sup> Anyway, more experiments at different temperatures or with different lights should be performed to analyze the effect of each energetic source.

Bai *et al.* have evaluated Bi NPs as a co-catalyst (substituting the usual noble metals) supported on the Bi<sub>4</sub>O<sub>5</sub>I<sub>2</sub> semiconductor under sunlight irradiation.<sup>171</sup> The study demonstrated photo-thermocatalytic reduction of CO<sub>2</sub> without any external heating. Experiments were carried out with different lights, showing how the full spectrum (UV-vis-IR) yielded the best rates, 40.02 µmol CO g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup> and 7.19 µmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>. The temperature of the system using the full spectrum or just IR was close to 75 °C, although the catalytic performance was relevant only in the case of full irradiation. Finally, other tests were performed only under UV-vis light and external heating up to 70 °C, showing a higher catalytic behavior than the same experiment at room temperature. Based on these outcomes, the semiconductor was excited by UV-vis light and electrons can be transferred to the Bi surface (co-catalyst). Simultaneously, Bi NPs absorbed IR light increasing the temperature of the active sites and facilitating the reduction of CO<sub>2</sub>. In conclusion, the reduction mechanism is a photo-thermocatalytic process without external heating and using the whole range of sunlight and it is described in Fig. 10.



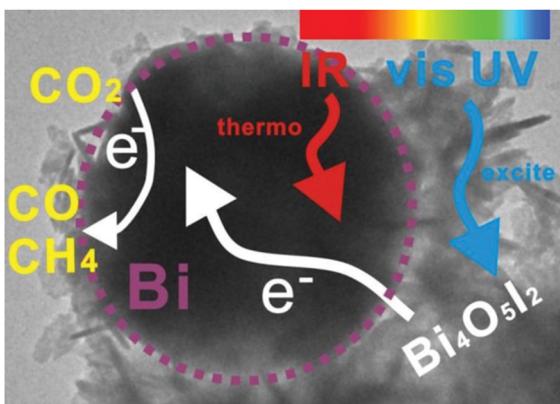


Fig. 10 Scheme of the photo-thermocatalysis mechanism of semimetal Bi decorated  $\text{Bi}_4\text{O}_5\text{I}_2$  nanosheets under UV-vis-IR light. Reproduced with permission from ref. 171. Copyright (2018) Elsevier.

Different research groups have studied the use of materials based on carbon or including carbon in their composition as photocatalysts for the  $\text{CO}_2$  reduction. Under UV irradiation a triple-catalyst based on carbon dots (0.9 wt%),  $\text{TiO}_2$  and Cu (0.9 wt%) was employed in the artificial photosynthesis reaction ( $\text{Cu}/\text{TiO}_2\text{-C}$ ) using a fixed bed flow reactor under atmospheric pressure.<sup>172</sup> Characterization analysis determined that Cu species were dispersed homogeneously in the semiconductor, as well as carbon dots. XPS analyses highlighted carbon dots as electron savers, permitting the delivery of electrons during the reactions due to the excellent separation of charges. Regarding Cu, XPS revealed Cu(i) as a stable species during several treatments, with the stability of this species being one responsible factor for the higher activity of this triple catalyst. The catalytic tests were performed with the triple catalyst and compared to the same catalyst without carbon dots.  $\text{Cu}/\text{TiO}_2$  catalysts did not produce  $\text{CH}_4$  in dark conditions at any temperature, although with UV-irradiation the best results were  $8 \mu\text{mol CH}_4 \text{ g}_{\text{cat}}^{-1} \text{ h}^{-1}$  at  $250^\circ\text{C}$ . On the other hand, the  $\text{Cu}/\text{TiO}_2\text{-C}$  catalyst showed significant thermocatalytic results at temperatures higher than  $150^\circ\text{C}$  in dark conditions which demonstrates that carbon dots promoted the thermocatalysis of the reaction. Similarly, when UV light was involved, the  $\text{CH}_4$  production rate increased significantly, giving the best results of  $60 \mu\text{mol CH}_4 \text{ g}_{\text{cat}}^{-1} \text{ h}^{-1}$  at  $250^\circ\text{C}$ . Moreover, the results with UV light at low temperatures and highest temperatures in dark conditions were very similar, proving that the activity is not ruled by only the photo or thermal effect, revealing a synergistic photo-thermal effect. Finally, a specific mechanism under UV irradiation was described which included the reduction of  $\text{CO}_2$  to  $\text{CO}$  followed by hydrogenation to  $\text{CH}_4$  via the  $^*\text{HCOO}$  intermediate, along with a WGS reaction for the production of  $\text{H}_2$  from  $\text{H}_2\text{O}$  which occurred at the same time. This reaction pathway is based on the reduction of  $\text{CuO}$  to  $\text{Cu}_2\text{O}$  by electrons from  $\text{TiO}_2$  and the *in situ* regeneration of  $\text{Cu}_2\text{O}$  using more electrons from  $\text{TiO}_2$  or the stored electrons in carbon dots, creating an excellent  $\text{CuO}/\text{Cu}_2\text{O}$  cycle to ensure a continuous reaction.

In a similar approach, Xu *et al.* combined a traditional photocatalyst ( $\text{TiO}_2$ ) with graphene achieving a  $\text{CH}_4$  productivity of  $26.7 \mu\text{mol CH}_4 \text{ g}_{\text{cat}}^{-1} \text{ h}^{-1}$  at  $117^\circ\text{C}$  without any external

heating.<sup>173</sup> Further investigations should be carried out to elucidate the action of photo and thermocatalysis and to determine which one is the dominant pathway in the reaction mechanism.

Other research groups have also studied photonic crystals as photo-thermal catalysts. Low *et al.* prepared photonic crystals based on  $\text{TiO}_2$  via an anodization-calcination method.<sup>174</sup> The light-absorption characteristics of the catalyst showed a peak at 378 nm related to the band gap of anatase (increasing light utilization) and an absorption range between 400 and 800 nm due to the layered photonic crystal structure, allowing the use of IR light. The absorption in the IR region can produce localized surface photo-thermal (LSPT) effects raising the temperature of the catalyst. The increase of temperature should enhance the photo-conversion as well as catalytic efficiency, favoring the reactant adsorption and/or desorption. The  $\text{CO}_2$  reduction using chemically produced  $\text{CO}_2$  and  $\text{H}_2\text{O}$  as reagents and irradiation with a Xe lamp produced  $35 \mu\text{mol CH}_4 \text{ h}^{-1} \text{ m}^{-2}$ , which is 15 times higher than that achieved with  $\text{TiO}_2\text{-P25}$ , and some  $\text{CH}_3\text{OH}$  as a secondary product. The authors base this enhancement on the optical abilities of  $\text{TiO}_2$  in combination with the LSPT effect of the photonic crystalline structure, which permits absorption of heat radiation of IR light and accelerates the transfer and generation of photo-generated electron-hole pairs, in brief, improving the carbon dioxide reduction by means of photo-thermal effects. Moreover, the authors describe the concept of slow photon as another reason for the best catalytic performance. This concept is related to the position of the photonic band gap, which matches with the absorption range of  $\text{TiO}_2$  and lowers the propagation speed of light. Consequently, photons can be easily absorbed and employed for the photo-generation of hot carriers.

Perovskite-type materials are semiconductors with good stability and flexibility.<sup>175</sup> Many of these compounds have been widely studied as photocatalysts for the production of  $\text{H}_2$  through water splitting.<sup>176–178</sup>

In 2017, Xu *et al.* reported a modified  $\text{LaFeO}_3$  perovskite oxide doped with Co as a catalyst for the  $\text{CO}_2$  reduction with  $\text{H}_2\text{O}$ , using visible-IR light ( $\lambda > 420 \text{ nm}$ ).<sup>179</sup> The photo-thermal catalytic activities were evaluated at  $350^\circ\text{C}$  after 6 h of irradiation obtaining  $\text{CH}_4$  ( $437.28 \mu\text{mol g}_{\text{cat}}^{-1}$ ) and  $\text{CH}_3\text{OH}$  ( $13.75 \mu\text{mol g}_{\text{cat}}^{-1}$ ) as unique products, with  $\text{CH}_4$  selectivity being higher than 96% in all cases. The results revealed that the modified-perovskite,  $\text{LaCo}_{0.6}\text{Fe}_{0.4}\text{O}_3$ , generated 3.2 and 4.0 times more  $\text{CH}_4$  and  $\text{CH}_3\text{OH}$  respectively than the unmodified perovskites ( $\text{LaFeO}_3$  or  $\text{LaCoO}_3$ ). This improvement is partially due to the narrower band gap of this perovskite (1.68 eV), making easier the formation of charge carriers ( $e^-/h^+$ ). Besides, XPS analysis revealed that the substitution of Fe by Co created more oxygen vacancies, and the promotion of these defects, also adsorption sites, enhanced the catalytic behavior. Other works have evaluated similar perovskites.<sup>180</sup> Specifically, using  $\text{LaNi}_{0.4}\text{Co}_{0.6}\text{O}_3$  under visible light at  $350^\circ\text{C}$  it is possible to obtain  $678.57 \mu\text{mol CH}_4 \text{ g}_{\text{cat}}^{-1}$  and  $20.83 \mu\text{mol CH}_3\text{OH} \text{ g}_{\text{cat}}^{-1}$  after 6 h, the best results at artificial photosynthesis based on perovskites. As well as the previous work, the substitution of Ni by Co ( $x = 0.4$ ) produced a higher amount of oxygen vacancies (67.90%), hence facilitating the adsorption and dissociation of  $\text{CO}_2$  and  $\text{H}_2\text{O}$ .



This substitution also modulated the positions of conduction and valence bands and the energy band gap (1.42 eV) of this perovskite to better promote both reactions involved in artificial photosynthesis. Finally, both studies shared that CH<sub>4</sub> should be produced mainly by photocatalytic reduction of CO<sub>2</sub>, as well as water splitting. Meanwhile, CH<sub>3</sub>OH should be produced through a thermo-catalytic route because the conduction band of the modified-perovskite is not negative enough to accomplish the reduction of CO<sub>2</sub> to CH<sub>3</sub>OH. Nevertheless, the mechanisms are not completely clear and there is a lack of evidence to evaluate the synergy of photo and thermal catalysis to clear up which is the main pathway or how thermal and/or photo effects act in each step of the reaction.

Double perovskites prepared by combustion with the formula La<sub>1-x</sub>Sr<sub>x</sub>Co<sub>1-y</sub>Fe<sub>y</sub>O<sub>3-δ</sub> have also been studied alone (LSCF) or synthesized over a mesoporous template (3DOM-LSCF) for the artificial photosynthesis combined reaction.<sup>181</sup> The photo-thermal reaction (UV-vis light + 350 °C) exhibited better results using both types of double perovskites than only thermocatalytic experiments, obtaining CH<sub>4</sub> as the only product in all tests. Light absorption experiments demonstrated that these perovskites can absorb UV and visible light, photo-generating electron–hole pairs that improve the reduction of CO<sub>2</sub> with H<sub>2</sub>O by means of photocatalysis. In accordance with the above results, the catalytic performance depends on the presence of oxygen vacancies and the smallest crystal size that increase the accessibility to active sites, with both properties being more related to thermocatalysis. On the other hand, although the band gaps of both double perovskites were very similar, the incorporation of the mesoporous structure in the 3DOM-LSCF catalyst influences the position of the conduction and valence bands, being more suitable in the case of the mesostructured material. Moreover, this specific structure also avoids the recombination of the photo-generated e<sup>-</sup>/h<sup>+</sup>, and thus, the photocatalytic process was improved. This was evidenced in the production rates which were 1.6 times higher when using 3DOM-LSCF (560 μmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>) than LSCF (350 μmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>). Concluding, this study demonstrates a photo-thermal coupling effect and both photo and thermo-catalytic mechanisms are shown in the study. Nevertheless, more studies must be performed to elucidate which is the dominant mechanism.

Additionally, the use of perovskites as supports for different metallic NPs has drawn special attention. A triple catalytic system based on the CaTiO<sub>3</sub> perovskite (70 wt%) as a semiconductor and support for Ni NPs (30 wt% NiO) and Pt NPs (1.0 wt%) has been applied as an “artificial leaf”.<sup>151</sup> Each one of the components holds specific functions: perovskites are not only good light-harvesters but, according to CO<sub>2</sub>-TPD experiments, good CO<sub>2</sub> adsorbents; Ni is the best metal to achieve the direct hydrogenation of CO<sub>2</sub> to CH<sub>4</sub>, and also interacts with the perovskite generating oxygen vacancies demonstrated by XPS characterization; and finally small amounts of Pt lend efficient photo and thermal H<sub>2</sub> activation, avoid the recombination of photo-generated charges and improve the formation of hot charges and local heating by means of LSPR. This work reveals unique reaction conditions because the authors employed both H<sub>2</sub> and H<sub>2</sub>O as an effective H<sub>2</sub> source or reducing agent (CO<sub>2</sub>:H<sub>2</sub>:H<sub>2</sub>O ratio of 1:2:2). Moreover, the catalysts were irradiated with a UV lamp in the temperature range of 40–240 °C in

a fixed bed quartz reactor, and four catalytic conditions were tested modifying the use of the reduced or not-reduced catalyst and the energy sources (photo, thermo and photo-thermo). The photo-thermal results of the reduced catalyst showed the best catalytic performance at 180 °C reaching 46.48% CO<sub>2</sub> conversion and 99.5% CH<sub>4</sub> selectivity (35 μL CH<sub>4</sub>). Compared with thermal-methanation in the literature for which at least 200 °C is necessary, the photo-thermal system mixing H<sub>2</sub>O and H<sub>2</sub> boosts the reaction at lower temperatures. Otherwise, the photo-results after 10 hours (not fixed reactor) displayed very low CO<sub>2</sub> conversion although again the catalysts with Ni and Pt raised CH<sub>4</sub> selectivity close to 100%. With respect to the thermal system without any irradiation, it showed similar trends to the photo-thermal system but with a lower production rate (26 μL CH<sub>4</sub>) and similar conversions (44.5%). Finally, the photo-thermal study was repeated with the non-reduced catalysts presenting the lowest results when heating is involved. This unlikeness is related to the better H<sub>2</sub> adsorption in metallic components or alloys and the lesser amount of oxygen vacancies in perovskites that are usually promoted in the presence of metallic NPs. Summarizing, cooperation between photo and thermal approaches allows a successful catalytic reduction of CO<sub>2</sub> using a mixture of H<sub>2</sub> and H<sub>2</sub>O as a reducing agent.

As it has already been mentioned in the previous sections of this review, to optimize photo-thermal performance, photocatalysts should display a strong light absorption across the solar spectrum, low charge-carrier recombination rates and a good capability for heat transfer and/or generation. In the particular case of CO<sub>2</sub> reduction, a high adsorption of CO<sub>2</sub> is also desirable. A huge variety of catalysts have been studied with very diverse properties that permit tuning the products of the reaction or even working at different temperatures and with different ranges of light. In the specific case of artificial photosynthesis, two reactions must be carried out simultaneously, complicating the search of an active, stable, inexpensive and environmentally friendly catalyst.<sup>182</sup> It seems that perovskite supports combined with plasmonic nanostructures could be especially convenient composite materials to perform efficient photo-thermal CO<sub>2</sub> conversion.

**4.1.2. Reverse water gas shift (RWGS).** This endothermic reaction can convert CO<sub>2</sub> into CO by a 2-electron reduction using H<sub>2</sub> as a reducing agent (eqn (4)) and is considered as a sustainable strategy for CO<sub>2</sub> capture and utilization (CUU). Together with a sustainable and economically acceptable hydrogen source, the produced CO can form syngas (CO + H<sub>2</sub>), which *via* Fischer-Tropsch synthesis allows for the synthesis of important hydrocarbons, such as alkanes, olefins, alcohols, dimethyl ether or gasoline.<sup>183,184</sup>



Experiments using Pd on TiO<sub>2</sub> at different times of irradiation (Hg lamp) and a fixed temperature (electric furnace) and *vice versa* demonstrated the cooperation between photo and thermocatalysis, improving the CO production by a photo-thermal coupling factor.<sup>185</sup> The maximum CO production rate was more than 8 times higher than for pristine TiO<sub>2</sub>-P25 (11.05 μmol CO g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>). The incorporation of Pd not only improved the photo-generated species



and the number of oxygen vacancies, but also enhanced  $\text{CO}_2$  adsorption.

Ozin *et al.* described hydride-terminated silicon nanocrystals as potential photo-thermal catalysts for the reduction of  $\text{CO}_2$  to CO, according to the thermal and optical properties of silicon materials. Only- $\text{CO}_2$  experiments using a metal halide lamp as a unique energy source at fixed 150 °C produced 4.5  $\mu\text{mol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ , but the reaction rate decreased after each run.<sup>186</sup>

Many more examples can be found in the literature on the production of CO using  $\text{H}_2$  as a reducing agent. The most interesting works involve catalysts based on the incorporation of metals on photoactive supports or semiconductors. Many works have studied the efficiency of Au NPs of different shapes and sizes in different photocatalytic reactions.<sup>187–189</sup> Using a Au/ZnO catalyst with a characteristic plasmon resonance around 538 nm, the authors were able to produce CO (4.22  $\mu\text{mol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ ) and small amounts of  $\text{CH}_4$  after 5 hours of irradiation with a continuous wave 532 nm laser, reaching 600 °C in a gas-tight photocatalysis cell with a  $\text{H}_2 : \text{CO}_2$  ratio of 3:1.<sup>103</sup> Thermal catalytic tests in dark conditions revealed a different temperature profile than that under irradiation, offering evidence of photo-generated heat that improved the catalytic performance. Other catalysts based on Au NPs supported on  $\text{TiO}_2$ ,  $\text{CeO}_2$  and  $\text{Al}_2\text{O}_3$  by deposition–precipitation have been evaluated in photo-thermal  $\text{CO}_2$  reduction.<sup>190</sup> The use of Au/TiO<sub>2</sub> with a metal loading of 1.4 wt% resulted in a  $\text{CO}_2$  consumption rate of 2.7  $\text{mmol g}_{\text{cat}}^{-1} \text{h}^{-1}$  at 400 °C under visible light irradiation. Lu *et al.* described a photocatalytic deposition method to produce Au/CeO<sub>2</sub> nanorods.<sup>191</sup>  $\text{CO}_2$  reduction under thermal conditions was not significant (<4%  $\text{CO}_2$  conversion). Nevertheless, the photo-thermal experiments achieving the same temperature as the thermal reactions just with irradiation (varying intensity) were much more active.

Another example is the use of Pd NPs below 10 nm size supported on the nanorods of  $\text{Nb}_2\text{O}_5$  semiconductor.<sup>34</sup> This catalyst was studied to elucidate the photo-thermocatalytic reduction of  $\text{CO}_2$ , using a batch reactor pressurized up to 1.8 atm with a  $\text{CO}_2 : \text{H}_2$  ratio of 1, during 3 h under three different situations: no light and no heating, just irradiation with a Xe lamp and 160 °C without irradiation. In this way, the best results were obtained under light conditions (4.9  $\text{mmol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ ) followed by the externally heated system (1  $\text{mmol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ ); meanwhile no significant  $\text{CO}_2$  conversion was observed when neither light nor temperature was used. Focusing on the photocatalytic conditions, several experiments were carried out cutting-off the light by filters. The results demonstrated a clear influence of the range of irradiation; for instance, the CO reaction rate decreased until 1.8  $\text{mmol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ , when UV-light was cut-off (only visible and IR light). According to the energy band gap of the semiconductor and the LSPR of Pd NPs, photoexcitation was not possible without UV-light. Therefore, the photons from visible and IR, not energetic enough to produce charge carriers, provided just thermal energy to the system by means of intra or interband transitions and non-radiative relaxation. In other words, when using visible and IR light,  $\text{CO}_2$  reduction to CO was possible by means of a photo-induced thermocatalysis. The most important conclusion was obtained when the authors compared the influence of both

wavelength and light intensity. Both sets of experiments resulted in approximately the same exponential profile (Fig. 11), evidencing that the photocatalytic pathway was not the dominating one.

Using the same catalyst and conditions, the selectivity towards RWGS (CO) or Sabatier reaction ( $\text{CH}_4$ ) was tuned by varying the Pd loading.<sup>132</sup> The 0.1% Pd/ $\text{Nb}_2\text{O}_5$  catalyst produced the highest CO production rate (18.8  $\text{mol CO g}_{\text{Pd}}^{-1} \text{h}^{-1}$ ) with a CO selectivity of 99.5%, meanwhile 10% Pd/ $\text{Nb}_2\text{O}_5$  significantly increased the  $\text{CH}_4$  selectivity by 24 times, decreasing the CO production rate ( $\sim 1 \text{ mol CO g}_{\text{Pd}}^{-1} \text{h}^{-1}$ ). The results normalized by Pd content pointed out an influence of the Pd-NP size on the activity and selectivity of the studied reaction; higher NP sizes produced higher  $\text{CH}_4$ . Furthermore, different amounts of 10% Pd/ $\text{Nb}_2\text{O}_5$  were evaluated obtaining a dependence between the mass of the catalyst and the  $\text{CH}_4$  production, and thus it was confirmed that selectivity was not related to just photocatalysis, but photo-thermal effects seemed to drive the reaction (light-to-heat). Ultimately, the characterization of the catalysts also demonstrated that oxygen vacancies and surface reduced Nb species are formed during the photo-thermal process by light irradiation, and this could improve the visible light absorption, avoid the recombination of  $e^-/h^+$  and besides, increase the  $\text{CO}_2$  affinity.

Other metals with well-known hydrogenation properties have been analyzed. Anatase- $\text{TiO}_2$  and  $\gamma\text{-Al}_2\text{O}_3$  were doped with Pt (2 wt%) to obtain two similar catalysts for the reduction of  $\text{CO}_2$  with  $\text{H}_2$  using a ratio of 1:6.<sup>192</sup> The experiments were carried out at different temperatures from 80 to 400 °C, under dark and irradiation conditions (Schott KL 2500 LCD visible-light). For both catalysts, the addition of light irradiation (photo-thermal process) promoted  $\text{CO}_2$  conversion and increased the CO production rate at every temperature, proving the photo-thermal effect. Specifically, for Pt/ $\text{Al}_2\text{O}_3$ , the increase was 9 times higher at 80 °C when the

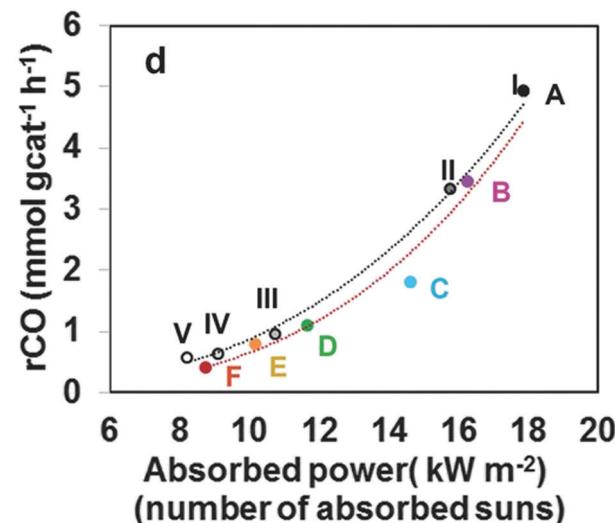


Fig. 11 Photo-thermal catalytic performance of the nanostructured Pd@ $\text{Nb}_2\text{O}_5$  samples. The RWGS reaction rates plotted as a function of absorbed power for the series of batch reactions with different cut-off filters A–F (red line) and diverse intensities I–V (black line). Reproduced with permission from ref. 34. Copyright (2016) John Wiley and Sons.



system was illuminated. However, better results were obtained with Pt/TiO<sub>2</sub> (110  $\mu\text{mol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$  at 400 °C) probably due to the strong interaction in Pt-TiO<sub>2</sub>. According to the presence of other products such as CH<sub>3</sub>OH, CHOOH or CH<sub>4</sub>, thermodynamic studies and operando DRIFT analysis, a possible mechanism was described based on thermocatalysis using Pt/Al<sub>2</sub>O<sub>3</sub>. Consequently, CO desorption was found to be the limiting step, which was enhanced by both heating and light irradiation, weakening the Pt–CO interaction.

Ru catalysts have also been described as efficient photo-thermal materials for CO<sub>2</sub> reduction. Through an ion-exchange method Ru was incorporated in a Mg(OH)<sub>2</sub> support with a three-dimensional nanoflower structure followed by 150 °C H<sub>2</sub>-reduction to create the metallic Ru.<sup>193</sup> Using a batch reactor at atmospheric pressure irradiated with a Xe arc lamp (300–2500 nm and 18 suns) and a mixture of CO<sub>2</sub> and H<sub>2</sub> (1:1), the best CO<sub>2</sub> conversion rate was 545  $\text{mmol g}_{\text{Ru}}^{-1} \text{h}^{-1}$  using the 8.3% Ru/Mg(OH)<sub>2</sub> catalyst. A clear exponential dependence between light intensity and photocatalytic activity was demonstrated (1413  $\text{mmol g}_{\text{Ru}}^{-1} \text{h}^{-1}$  at 20 suns), therefore suggesting that thermocatalysis rules the photo-thermal reaction. Anyway, deeper investigations must be done to explain the photo-thermal effects and mechanisms but it is obvious that a suitable amount and dispersion of Ru increase the light absorption and can modify the product distribution, defining these variables as key factors in photo-thermal CO<sub>2</sub> reduction.

The use of noble metals entails a significant cost. In order to reduce the price of the catalysts, non-noble transition metals have also been studied as part of catalysts in photo-thermocatalytic reactions as a more economical option. Ning *et al.* synthesized a hybrid carbon-based catalyst with embedded Co NPs together with dispersed Co–N species, denoted as Co@CoN & C.<sup>194</sup> The reactions were carried out at 55 kPa, using an equal amount of CO<sub>2</sub> and H<sub>2</sub> (1:1) and a Xe lamp in a batch reactor. Initially, the photoinduced thermal-effect was demonstrated as evidenced by the increase of temperature of the catalyst surface during the reaction, which reached 518 °C, while the plain Co NP system increased only up to 300 °C. After 30 min of irradiation, the optimized catalyst (ratio of CN:Co(OH)<sub>2</sub> of 1) showed 41% CO<sub>2</sub> conversion and a CO selectivity of 91% (132  $\text{mmol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ ), meanwhile the unsupported Co NPs showed 43% CO selectivity. The high reaction rate of the hybrid catalyst was attributed to the significant photo-thermal effect, and the CO selectivity was improved by the special characteristics provided by the Co–N shell in comparison with plain Co NPs.

The utilization of metal organic frameworks (MOFs) as a sacrificing template for the manufacture of catalysts has also been applied in photo-thermal catalysis. A core–shell catalyst (Fe@C) was prepared by a double step pyrolysis of MIL-101(Fe).<sup>195</sup> This catalyst was irradiated in a batch reactor using a Xe lamp and a CO<sub>2</sub>:H<sub>2</sub> molar ratio of 1:1, presenting a CO<sub>2</sub> reaction rate of 26.12  $\text{mmol g}_{\text{cat}}^{-1} \text{h}^{-1}$  with almost 100% CO selectivity. Meanwhile, other catalysts based on naked Fe NPs (Fe/SiO<sub>2</sub> and Fe/CNT) also presented catalytic activity but never as active as Fe@C. Moreover, the selectivity to CO was a bit lower for the naked Fe-catalysts and for longer reaction times total selectivity was reduced, improving the selectivity to CH<sub>4</sub> and even hydrocarbons

such as C<sub>2</sub>H<sub>4</sub>, C<sub>2</sub>H<sub>6</sub>, C<sub>3</sub>H<sub>6</sub> or C<sub>3</sub>H<sub>8</sub>. The same catalysts were employed under identical conditions in a fixed-bed reactor, producing 55.75  $\text{mmol CO min}^{-1}$  as the highest reaction rate using Fe@C. These photo-thermal experiments presented the best results compared to traditional photocatalytic and thermocatalytic experiments, an evidence of photo-thermal effects. In addition, catalytic tests at diverse wavelengths showed optimum results using UV light, probably due to the plasmonic activity of Fe NPs. According to the characteristics of Fe@C and diverse catalytic outcomes, the absorption of light by Fe (LSPR) created local heating which was able to activate CO<sub>2</sub> at high temperatures. It should be remarked that when carbon materials are used as supports, it is very important to demonstrate that all carbon products are derived from the fed CO<sub>2</sub> and not from the direct hydrogenation of the support.

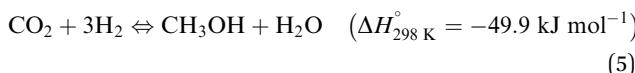
The hydrogenation of In<sub>2</sub>O<sub>3</sub> at 400 °C results in the formation of oxygen vacancies (In<sub>2</sub>O<sub>3-x</sub>/In<sub>2</sub>O<sub>3</sub>),<sup>196</sup> and in a heterostructure with a lower bandgap (2.36 eV) that improves solar light harvesting and absorption. This photo-enhancement (In<sub>2</sub>O<sub>3-x</sub>/In<sub>2</sub>O<sub>3</sub> catalyst) was correlated to the best catalytic performance *via* RWGS, obtaining CO as the only product (19.64  $\text{μmol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ ) in a batch reactor system irradiated with a Xe lamp (H<sub>2</sub>:CO<sub>2</sub> of 1). The increase of the temperature up to 160 °C under light irradiation evidenced a photo-thermal effect. Moreover, the authors evaluated the catalysts using a gas-phase flow reactor irradiated with the same lamp and with external heating. Compared to dark tests, the photo-experiments presented a 417% greater CO rate at 300 °C. Other researchers have studied the same semiconductor but using novel synthesis methods to obtain ultrathin 2D black In<sub>2</sub>O<sub>3-x</sub> nanosheets with photoinduced oxygen vacancies.<sup>102</sup> This catalyst was compared to traditional In systems (bulk, hydroxide, defective oxide). The 2D structure improved the photoactivity (more electrons were trapped at oxygen defects) and the oxygen vacancies activated CO<sub>2</sub> molecules providing more active sites for the photo-thermal CO<sub>2</sub> reduction.

The good results of modified In<sub>2</sub>O<sub>3</sub> as a semiconductor in photocatalytic studies motivated Hoch *et al.* to impregnate silicon nanowires (NWs) with In<sub>2</sub>O<sub>3-x</sub>(OH)<sub>x</sub> NPs. This catalyst was employed to reduce CO<sub>2</sub> to CO at a rate of 22  $\text{μmol g}_{\text{cat}}^{-1} \text{h}^{-1}$  under irradiation with a Xe lamp (20 suns) without external heating.<sup>197</sup> The same experiments using only  $\lambda > 495$  nm (photons without the possibility to create charge separation in In<sub>2</sub>O<sub>3-x</sub>(OH)<sub>x</sub>) presented a CO rate 20 times lower. Most of the photons with wavelengths higher than 495 nm did not participate in the reaction in a radiative way. Instead, these photons were trapped by SiNWs and increased the temperature, facilitating the release of produced H<sub>2</sub>O and improving the thermocatalytic reaction. On the other hand, more energetic photons (<495 nm) contributed to the photocatalytic reduction of CO<sub>2</sub> to CO.

Hydride-terminated silicon nanocrystals were also studied as catalysts for the reduction of CO<sub>2</sub> using H<sub>2</sub>.<sup>186</sup> Comparing experiments under dark and light conditions at the same temperature, the CO production rate is higher when the catalytic system is heated and irradiated ( $\sim 240 \text{ μmol CO g}_{\text{cat}}^{-1} \text{h}^{-1}$ ).

**4.1.3. Methanol production.** CH<sub>3</sub>OH (formed *via* eqn (5)) is an important raw material that can be either directly used or

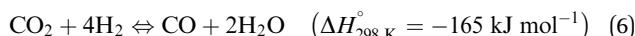
transformed to synthetic hydrocarbons (*i.e.* methanol to olefins process, MTO) and other valuable compounds.<sup>198,199</sup> In addition to being traditionally produced from syngas, methanol can also be directly formed from CO<sub>2</sub> and H<sub>2</sub> *via* traditional thermocatalysis.<sup>200</sup> As it could be expected, photo-thermal catalysis has also been tried as an alternative.



Hybrid structures formed by Pt nanocubes (6 nm) and Au nanocages (31 nm) inside of the well-known ZIF-8 MOF were reported for the photothermal transformation of 1:3 CO<sub>2</sub>:H<sub>2</sub> mixtures at 30 bar under irradiation (Xe lamp) in batch mode.<sup>36</sup> While the Au nanocages did not show any conversion, Pt nanocubes after 3 hours produced small amounts of CH<sub>3</sub>OH at 150 °C with high selectivity (96%), both under dark and under illumination. The incorporation of Pt-nanocubes in ZIF-8 increased slightly the methanol production up to 0.7 mmol CH<sub>3</sub>OH, with no difference between dark and illumination conditions. In contrast, the 3-component catalyst consisting of Pt@Au@ZIF-8 showed an enhanced methanol productivity under light irradiation. It was concluded that (i) Pt nanocubes act as active sites for the transformation of CO<sub>2</sub> to CH<sub>3</sub>OH; (ii) the addition of Au results in localized heating of the active sites by plasmonic photo-thermal effects (IR adsorption), and (iii) ZIF-8 acts as a heat insulator.

Pd-Zn catalysts have also been tested under photo-thermal conditions.<sup>201</sup> Experiments at different temperatures (190–270 °C) using a 3:1 H<sub>2</sub>:CO<sub>2</sub> ratio presented the best CO<sub>2</sub> conversion and CH<sub>3</sub>OH production rates when the system was irradiated with UV light (Hg lamp). More importantly, the use of only visible light also resulted in an important enhancement in catalytic activity compared to experiments performed in the dark. Diffuse reflectance spectroscopy confirmed a shift to the visible range of the LSPR of Pd when alloyed with Zn. In view of these results, the authors speculate that hot-electrons generated upon light irradiation migrate to the active sites (Pd-Zn alloy), increasing in this way the electron density and therefore catalytic activity *via* electron injection to the antibonding orbital of CO<sub>2</sub>. We find that these results are very encouraging.

**4.1.4. Methane production.** The well-known Sabatier reaction (eqn (6))<sup>26,202,203</sup> is one of the most studied processes for the recycling of CO<sub>2</sub>.



In general, Ni-supported catalysts are widely used due to the high activity and CH<sub>4</sub> selectivity, in addition to the fact that Ni is a highly available and cheap metal.<sup>204</sup> In contrast, only a few studies in the photo-thermal literature on the topic focus on this metal. Ni-CeO<sub>2</sub>-TiO<sub>2</sub> catalysts prepared by sol-gel chemistry under UV irradiation (Xe lamp) are among the most active photo-thermal catalysts reported to date.<sup>205,206</sup> The presence of CeO<sub>2</sub> was claimed to enhance CO<sub>2</sub> adsorption and to improve surface heating, while TiO<sub>2</sub> was responsible for the stabilization of oxygen vacancies. Analysis of CH<sub>4</sub> evolution *versus* temperature showed an initial increase of the temperature required to reduce the passivated

catalyst followed by CH<sub>4</sub> production and an extra heating coming from the exothermicity of the reaction.

Incorporation of Ni in Y<sub>2</sub>O<sub>3</sub> nanosheets allowed CO<sub>2</sub> methanation under sunlight.<sup>207</sup> Under 1 sun irradiation the catalyst with Ni NPs did not present good catalytic results (11% CO<sub>2</sub> conversion and total CH<sub>4</sub> selectivity), probably because the catalytic system reached only 78 °C temperature, not enough for the activation of the Sabatier reaction. To improve this catalytic activity, Ni was incorporated as single Ni atom species instead of Ni NPs, achieving 87% CO<sub>2</sub> conversion and a temperature up to 240 °C. Finally, in order to obtain a photo-thermal catalytic system, the single atom Ni/Y<sub>2</sub>O<sub>3</sub> catalyst was covered by a selective light absorber (AlN<sub>x</sub>/Al). This new system permitted elevating the temperature of the catalysts to 288 °C under just irradiation of 1 sun describing a reaction mechanism mainly driven by thermocatalytic effects. Using natural sunlight, a CH<sub>4</sub> production rate of 7.5 L m<sup>-2</sup> h<sup>-1</sup> could be achieved (Fig. 12).

Since 1987, different researchers have employed Ru-based catalysts for photo and thermocatalytic CO<sub>2</sub> reduction.<sup>32,208–210</sup> Incorporation of Ru in TiO<sub>2</sub>, according to Wang *et al.*, can narrow the energy band gap of the semiconductor (2.69 eV).<sup>211</sup> Diverse catalytic tests were carrying out in a fixed bed reactor with a CO<sub>2</sub>:H<sub>2</sub> ratio of 1:3. However, catalytic tests using sunlight (equivalent to 1 sun) at 150 °C did not show any difference with experiments performed in the dark at the same temperature. On the contrary, when increasing the temperature to 300 °C, the catalytic activity was three times higher under light illumination (69 mmol CH<sub>4</sub> g<sub>cat</sub><sup>-1</sup> h<sup>-1</sup>) than in dark conditions. Nitrogen doping of TiO<sub>2</sub> in similar systems can further enhance the photo-thermal behavior.<sup>212</sup>

Multimetallic systems such as Ru and Au supported on TiO<sub>2</sub> have also been studied,<sup>159</sup> once again confirming that the addition of Au greatly improves heat effects.<sup>108</sup>

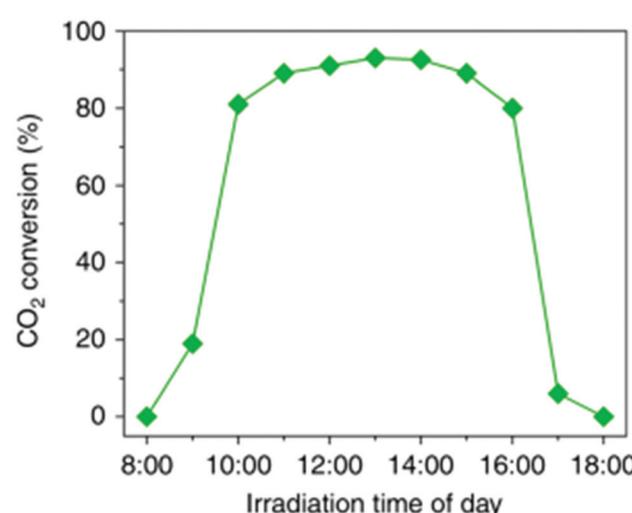


Fig. 12 Photo-thermal CO<sub>2</sub> conversion using SA Ni/Y<sub>2</sub>O<sub>3</sub> nanosheets with a selective light absorber-assisted photo-thermal system as a function of time. Experiment carried out on June 30, 2018 from 8:00 to 18:00 in Baoding, Hebei, China. Reproduced with permission from ref. 207. Copyright (2019) Springer Nature.



When it comes to layered double hydroxides (LDHs) as catalytic supports,<sup>213,214</sup> Ren *et al.* supported Ru NPs on ultrathin Mg-Al LDHs and tested the resulting catalysts under flow and UV illumination (Xe lamp).<sup>215</sup> Under these conditions, the photo-thermal catalysts could reach up to 350 °C and showed similar activities to those obtained when the reaction is performed in the dark at the same temperature, highlighting the important thermal and the negligible photo effect for these specific catalysts. In the same spirit, Ru has been supported on many other semi-conductors such as silicon nanowires,<sup>32</sup> silica opal and inverted silicon opal photonic crystals,<sup>203</sup> planar silicon wafers<sup>26</sup> or strontium titanate.<sup>216</sup>

Similar photo-thermal effects have been found when using Rh as the active phase and different semiconductors as supports,<sup>217</sup> with very interesting studies on the effect of light intensity on catalyst activity (Fig. 13).<sup>141</sup>

As it can be concluded from the above, group VIII metals have attracted a great deal of attention. This is mostly due to an effective absorption in the full range of the solar spectrum (covering UV, visible and IR light) together with a high H<sub>2</sub> activation ability. In a comparative study, Ye *et al.* impregnated an inert Al<sub>2</sub>O<sub>3</sub> with 2–3 wt% group VIII metals (Ru, Rh, Ni, Co, Pd, Pt, Ir and Fe) and studied the catalytic properties of the resulting catalysts in a batch reactor with a gas mixture of CO<sub>2</sub> and H<sub>2</sub> (1 : 4) under Xe lamp irradiation.<sup>110</sup> As a result of the strong increase in temperature (going up to 300–400 °C), CH<sub>4</sub> production rates measured under these conditions were six orders of magnitude higher than under only photocatalytic conditions, highlighting, once again, the importance of photo-thermal processes in the final outcome of the reaction. Further studies on similar metals supported on SiO<sub>2</sub><sup>218</sup> revealed a linear relation between light intensity and CO<sub>2</sub> conversion rate, demonstrating that the reaction is driven by hot electrons from the metal.

**4.1.5 Towards C<sub>2+</sub> products.** The creation of C–C bonds during CO<sub>2</sub> activation has proved challenging, specially under electro or photo-catalytic conditions. The main reason is that,

usually, a high partial pressure of the intermediate CO is required and this is difficult to achieve under regular electro and photo-catalytic conditions. Indeed, Fischer–Tropsch synthesis, which transforms syngas (usually with CO : H<sub>2</sub> ratios close to 1) into longer hydrocarbons, is carried out at 20–30 bar and moderate temperature.

Surprisingly, following a photo-thermal approach, several groups have been able to demonstrate the formation of C<sub>2+</sub> hydrocarbons from CO<sub>2</sub> : H<sub>2</sub> mixtures at atmospheric pressure. In 2018, Zhang and co-workers, using an alumina-supported CoFe alloy reduced at 650 °C, were able to achieve 37% C<sub>2+</sub> selectivity and a total CO<sub>2</sub> conversion of 78.6% under irradiation with a Xe lamp at 5.2 W cm<sup>-2</sup> irradiance and with no external heating.<sup>219</sup> In order to elucidate if pure photochemical or thermal effects drove the CO<sub>2</sub> hydrogenation, the authors monitored the temperature of the catalyst using a thermocouple in intimate contact with the catalyst bed. According to these experiments, under UV-light irradiation, the catalyst achieved temperatures as high as 310 °C, while in the absence of catalyst the reactor temperature was only 120 °C, proving the existence of a strong photo-thermal effect. To gain further evidence of this photo-thermal effect, CO<sub>2</sub> conversion was also monitored as a function of temperature under both light and dark conditions. Interestingly, CO<sub>2</sub> conversion profiles were almost identical in both cases, proving that the reaction was driven photo-thermally rather than photochemically. Computational calculations showed that the CoFe bimetallic alloy promoted C–C coupling reactions during CO<sub>2</sub> hydrogenation, thus explaining the higher selectivity towards C<sub>2+</sub> products.

More recently, Garcia and Corma reported a sodium-promoted Co@C (Na–Co@C) nanocomposite for the photo-thermal hydrogenation of CO<sub>2</sub> to ethanol and C<sub>2+</sub> hydrocarbons.<sup>220</sup> Using this catalyst under photo-thermal conditions (235 °C and 24 kW m<sup>-2</sup> power irradiance), the selectivity to ethanol and C<sub>2+</sub> hydrocarbons was 7 and 36%, respectively, at a CO<sub>2</sub> conversion of 37%. A blank experiment under dark conditions at the same temperature and CO<sub>2</sub> conversion led to markedly lower selectivity, thus indicating that light irradiation had a strong influence on product selectivity. In order to further demonstrate the role of light in the catalytic performance, the authors conducted wavelength-dependent studies using monochromatic light. The results showed that UV light was mainly responsible for CO<sub>2</sub> hydrogenation, and photo-action experiments had a clear match with the absorption spectrum of the material. Ultimately, this was clear evidence that the reaction mechanism was driven predominantly by the photo-generation of charge carriers. Additional near-ambient pressure X-ray photoelectron spectroscopy (AP-XPS) and *in situ* Raman mechanistic studies not only demonstrated that photo-induced electrons on Na–Co@C promoted CO<sub>2</sub> dissociation to CO, but also that light irradiation stabilized CO on the catalyst surface, therefore contributing to a higher selectivity towards ethanol by a CO insertion mechanism. Furthermore, the authors also observed that UV light had an evident effect on the chain growth process according to the different surface intermediates compared to dark experiments, a fact that explains the difference in the reported selectivity. Altogether, these results evidenced the potential of light to selectively activate different reaction pathways in CO<sub>2</sub> hydrogenation.

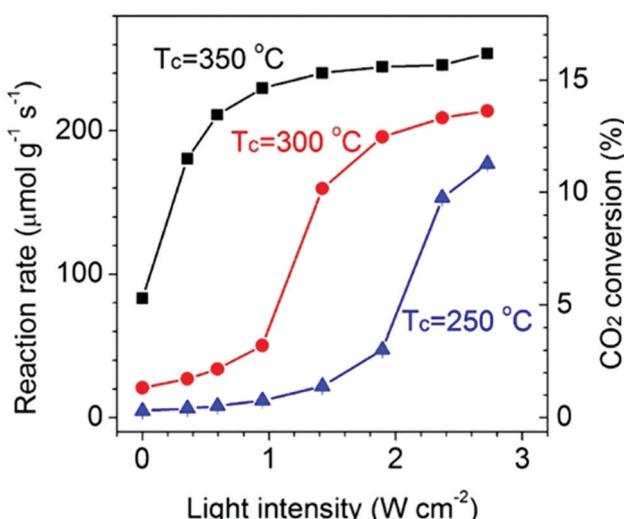


Fig. 13 Total reaction rate and CO<sub>2</sub> conversion at 350 °C (black squares), 300 °C (red circles) and 250 °C (blue triangles) as a function of light intensity. Reproduced with permission from ref. 141. Copyright (2018) American Chemical Society.



#### 4.2. Fischer-Tropsch chemistry

In 2015, two research groups reported interesting results in the photo-thermally catalyzed Fischer-Tropsch (PFT) process.<sup>221,222</sup> Guo and co-workers, using a worm-like ruthenium nanostructure dispersed on graphene sheets, reported the efficient formation of  $C_{20+}$  hydrocarbons from  $CO:H_2$  (1:1) mixtures at 20 bar, under irradiation from 400 nm to 800 nm using a 300 W Xe lamp.<sup>221</sup> The catalytic activity in these conditions is  $14.4 \text{ mol CO mol}_{Ru}^{-1} \text{ h}^{-1}$ , while the dark reaction, *i.e.* conducted without light irradiation, shows a halved activity value,  $7.8 \text{ mol CO mol}_{Ru}^{-1} \text{ h}^{-1}$ . Moreover, this last value is comparable with the one reported by Kou and co-workers for another Ru nanocluster-based catalyst applied in FT synthesis, and this confirms that light irradiation plays indeed a key role.<sup>223</sup> Guo's group attributes these improved FT reaction features to the formation of hot electrons in graphene; these high energy electrons can transfer to Ru, which also generates energetic electrons due to light absorption and initiates the process by activating CO. The authors confirmed this hypothesis proving distinctive features of the hot-electron mechanism: a linear relationship between reaction rate and light intensity, as shown in Fig. 14a, a lower activation energy under light irradiation compared to dark conditions, as shown in Fig. 14b, and a correlation between catalytic activity and irradiation wavelength, as shown in Fig. 14c and d.

Yu *et al.* presented a series of iron-based catalysts supported on mesoporous titania ( $TiO_2$ ) at different Fe/Ti ratios, and they

tested the catalytic performance of these materials in the photo-thermal Fischer-Tropsch process to convert coal to FT products.<sup>222</sup> After a thorough characterization of five catalytic systems having 5%, 10%, 15%, 20%, and 30% Fe/Ti ratio, the researchers performed PFT tests at atmospheric pressure, at a temperature of  $220^\circ\text{C}$  and with a  $H_2/CO$  ratio of 1:2. Although the selectivity results did not follow a clear trend in correspondence with the amount of Fe in the catalyst, and the authors did not report a clear reaction mechanism, the study shows that the investigated catalytic systems are active in PFT and that light irradiation allows one to carry out PFT synthesis using mild conditions.

In recent years, more research groups have put their efforts to investigate deeper this topic; specifically they have tried to address the problem of increasing the selectivity towards light olefins. Su and co-workers designed a catalytic system combining a thermo-active part and a photosensitive support.<sup>224</sup> They prepared *via* impregnation two  $TiO_2$  nanotube supported Co catalysts (20% and 30% Co/ $TiO_2$ ), and studied them for the FT reaction ( $CO:H_2 = 1:2$ , pressure 20 bar) under photo-thermal conditions using UV light, thermal only condition and UV irradiation only condition. The results were compared after 40 hours on stream, and the data clearly showed that UV irradiation alone is not enough to overcome the energy barrier to catalyze CO hydrogenation. The 20% Co/ $TiO_2$  catalyst shows the most interesting result in terms of CO conversion; in fact,

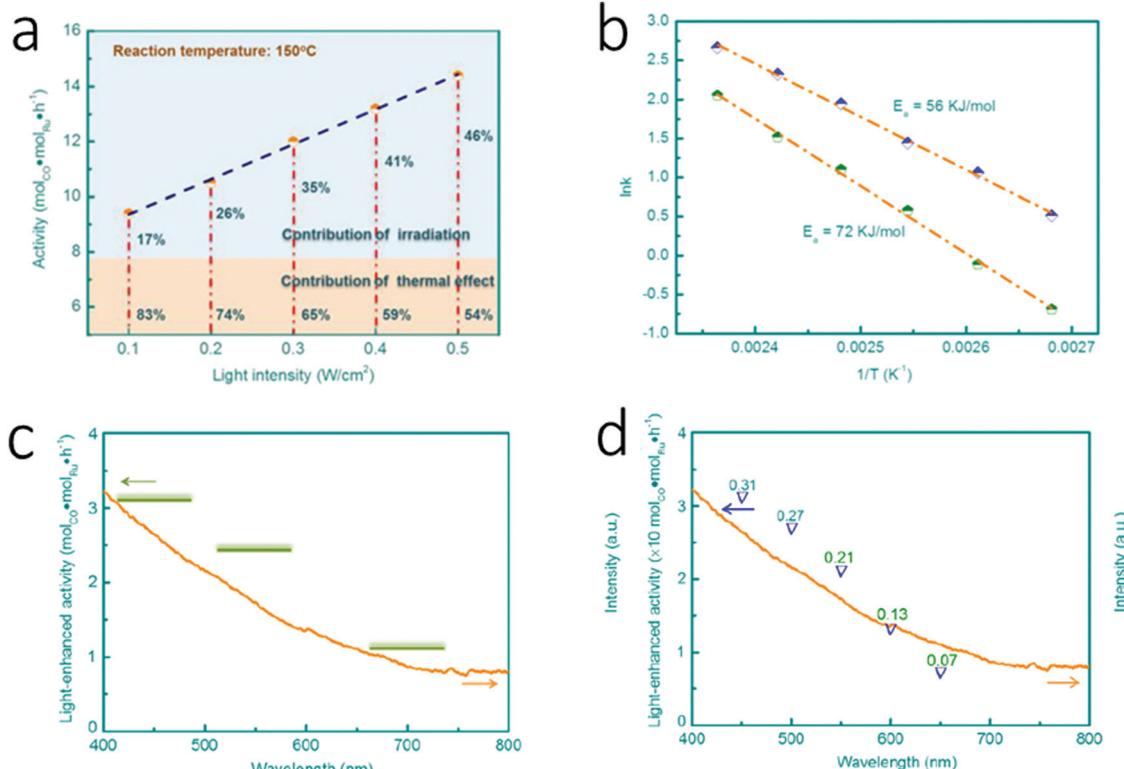


Fig. 14 (a) Dependence of the photocatalytic activity of Ru/graphene on light intensity. (b) Arrhenius plots of photo-thermal Fischer-Tropsch reactions under light and dark conditions. (c) Dependence of the photocatalytic activity of Ru/graphene on wavelength regions. (d) Dependence of the photocatalytic activity of Ru/graphene on specific wavelength. Reproduced with permission from ref. 221. Copyright (2015) American Chemical Society.



although the CO conversion value for the thermocatalytic process is 9.2%, it increases significantly to 63.9% in the presence of UV irradiation under the same temperature and pressure conditions. As Guo *et al.* already proposed, the support, graphene in their work and  $\text{TiO}_2$  in Su and co-workers' paper, plays a key role as a photococatalyst.<sup>221</sup> As a consequence of UV light irradiation, hot carriers (electron-hole pairs) are generated in  $\text{TiO}_2$ : electrons can be collected on Co surface sites and transferred to CO to initiate the dissociation of C–O bonds; holes can be intercepted by the available reduced species. To prove that UV irradiation is the reason for the higher CO conversion because it activates CO dissociation, Su's group performed *in situ* Raman spectroscopy. This technique confirmed the adsorption of CO on the metal center and consequently the weakening of the C–O bond. Furthermore, for both catalysts the product distribution showed that the  $\text{CO}_2$  selectivity decreases under photo-thermal conditions. As far as the hydrocarbon selectivity is concerned, the quantity of  $\text{CH}_4$  is enhanced under UV light, due to the photohydrogenolysis of hydrocarbons with longer chains and the formation of hydrogen atoms from  $\text{H}_2$  dissociation. At the same time, in the  $\text{C}_2$ – $\text{C}_4$  fraction the photo-thermal process shifts the selectivity towards alkanes with a remarkable value of 98.6% *versus* 89.0% for the dark reaction. Overall, Su and co-workers also showed that the same CO conversion obtained with conventional thermo catalysis at 300 °C could be achieved at 220 °C using photo-thermal catalysis, when performing a more environmentally friendly Fischer–Tropsch synthesis.

In the same year Zhang and co-workers published a paper in which they investigated the photo-thermal Fischer–Tropsch process using a cobalt based catalytic system.<sup>225</sup> They fabricated a series of catalysts consisting of Co and  $\text{Co}_3\text{O}_4$  NPs supported on zinc oxide–alumina ( $\text{ZnO}–\text{Al}_2\text{O}_3$ ), and they used them for the photo-thermal Fischer–Tropsch to olefins (FTO) synthesis. A detailed characterization of the obtained materials showed that, depending on the reduction temperature at which the catalysts are prepared, the presence at the surface of cobalt oxides and metallic cobalt can be tuned. The different catalysts were tested under the same reaction conditions: no external heating, 1.8 bar, 300 W Xe light (200–800 nm), 1 hour irradiation time, and the outcomes showed that the different surface composition strongly influences the CO conversion as well as the product selectivity. In particular, the most notable results were achieved using the catalyst prepared when the reduction temperature was 450 °C (Co-450), yielding 15.4% CO conversion and good olefins selectivity (36%) with a olefins to paraffins ratio of 6.1. In order to have a deeper insight into the process occurring, Zhang's group performed catalytic tests over Co-450 in different conditions: by monitoring the temperature under UV-irradiation; by using electric heating only; and by controlling light irradiation with monochromatic filters. Overall, they proved that the photo-thermal effect is the driving phenomenon for CO hydrogenation over this new cobalt-based catalyst rather than a pure photochemical route.

Another research group in 2018 investigated photo-thermal Fischer–Tropsch as a way to convert syngas into lower olefins. Ma and co-workers designed an iron-based catalyst,  $\text{Fe}_5\text{C}_2$ , whose

surface was spontaneously decorated with oxygen atoms *in situ* when irradiated with UV light.<sup>226</sup> After preparation, the catalyst was analyzed *via* energy dispersive X-ray (EDX) spectroscopy and powder X-ray diffraction (PXRD) to determine the elemental distribution and the phase purity. The researchers tested  $\text{Fe}_5\text{C}_2$  as a catalyst for PFT performed at atmospheric pressure, molar ratio of  $\text{CO}/\text{H}_2 = 1:2$ , 300 W Xe light (200–1100 nm), 0.5 hour irradiation time, and the results were quite surprising: CO conversion was 49.5%, selectivity toward  $\text{CO}_2$  was 18.9%, selectivity toward  $\text{C}_{2-4}$  olefins was 55.5 and selectivity toward  $\text{C}_{2-4}$  paraffins was 5.1% (o/p ratio 10.9). The temperature of the reaction was also monitored, and the catalyst bed reached 200 °C after 60 seconds under irradiation, while after 28 minutes the temperature rose to 490 °C. This behavior, supported by 3D excitation emission photoluminescence analysis, shows that the catalytic system undergoes photo-induced heat conversion, corroborating the hypothesis of the photo-thermal catalytic route. Additional investigations were performed to understand the reason behind the promoted production of olefins. X-ray photoelectron spectroscopy performed on fresh and used catalysts showed that photoirradiation leads to the formation of iron oxide on the  $\text{Fe}_5\text{C}_2$  surface, in a significantly minor amount compared to the formation of iron oxide observed in the thermal FT process. This finding, integrated with density functional theory (DFT) calculations on  $\text{Fe}_5\text{C}_2$  and oxygen-decorated  $\text{Fe}_5\text{C}_2$ , allowed Gao *et al.* to correlate the presence of a suitable amount of oxygen on the catalyst surface with the favorable production of olefins.

Zhang's research group expanded the investigation field on the PFT process and reported a new class of Co-based catalysts prepared *via* hydrogen reduction performed at different temperatures (from 300 °C until 700 °C) of CoAl layered-double-hydroxide (LDH) nanosheets.<sup>227</sup> These materials perform better than those containing the same metal that the research group presented in a previous paper discussed above.<sup>225</sup> The different catalysts were tested for CO hydrogenation at 1.8 bar and irradiation with a 300 W Xe light (200–800 nm) for 1 hour, without external heating, and an improving trend was observed with the increase of the reduction temperature used for the preparation of the catalyst. The material prepared at 700 °C, namely Co-700, gave the best catalytic performance yielding a CO conversion of 35.4% and a selectivity to higher hydrocarbons of 65%; therefore it was selected for additional catalytic tests using different conditions in order to optimize the procedure. Finally, Li *et al.* conducted a series of tests modifying the light source while monitoring the temperature, and they were able to prove that the investigated CO hydrogenation is a process driven by a photo-thermal route rather than a non-thermal photoexcitation process.

Very recently another paper from Zhang and co-workers was published, in which they reported the use of an unusual metal for the Fischer–Tropsch process: nickel.<sup>228,229</sup> Wang *et al.* fabricated a novel class of catalysts, consisting of Ni NPs supported on manganese oxide ( $\text{MnO}$ ), labelled Ni-*x*, where “*x*” represents the temperature at which the Ni–Mn mixed metal oxide was reduced to obtain the target materials. This metal–support combination was a winning strategy to optimize the hydrogenation properties of Ni to make it suitable for light olefins



production *via* PFT synthesis, performed at 1.8 bar and irradiation with a 300 W Xe light (200–800 nm) for 1 hour, without external heating. Particularly, the performance of Ni-500 was the best compared to the other catalysts of the series. Ni-500 showed the highest value of selectivity to light olefins (33.0%) and a really low CO<sub>2</sub> selectivity (0.2%). The Zhang group performed catalytic tests monitoring the temperature under UV irradiation, as well as dark tests with direct heating, and they confirmed that the studied Fischer–Tropsch process follows a photo-thermal route. Furthermore, they were able to clarify the role of MnO and Ni in the process. Ni K-edge X-ray-absorption near edge spectroscopy (XANES) showed that the properties of the support change the electronic features of metallic Ni NPs, and X-ray photoelectron spectroscopy (XPS) proved that charge transfer from MnO to Ni NPs occurs. These phenomena reduce the hydrogenation activity of Ni, and clarify the unexpectedly high selectivity of Ni-500 towards olefins.

#### 4.3. Methane activation

In the last years, scientists have made significant efforts to develop new catalytic approaches for one of the biggest challenges in the chemical community: performing methane conversion to upgraded fuels and commodity chemicals under milder conditions. Photo-thermal-driven methane activation is a recent strategy that several research groups are exploring, because it allows exploring new routes, which require lower activation energies and, consequently, a more environmentally friendly approach.

In 2019, two interesting reviews were published about sustainable conversion of methane driven by solar energy, and about the role of solar energy in obtaining fuels from catalytic C<sub>1</sub> chemistry.<sup>230,231</sup> Both publications contain a substantial section that revises a series of photo-thermal catalytic systems, which were tested for dry reforming of methane (DRM) and steam reforming of methane (SRM), reported until then.<sup>109,116,232–243</sup>

In addition to SRM and DRM, the partial oxidation of methane (POM) in the presence of oxygen has also been explored by photo-thermal means. In 2019, Miyauchi and co-workers used Pd NPs supported on strontium tantalate (Sr<sub>2</sub>Ta<sub>2</sub>O<sub>7</sub>) to perform the reaction under UV irradiation at temperatures lower than those conventionally required.<sup>244</sup> They observed that, in contrast to any other thermal process reported, UV light initiates the partial oxidation at 150 °C, while in dark conditions the conversion of CH<sub>4</sub> and O<sub>2</sub> begins only above 350 °C. In addition, at this temperature, the conversion of CH<sub>4</sub> goes from 21% in the dark to 60% under irradiation, and the yield of CO and H<sub>2</sub> goes from 2% and 2% to 47% and 53%. The Arrhenius plot (logarithm of formation rate of CO *versus* the inverse temperature) allows calculating the activation energy ( $E_a$ ) for the reaction under UV light (11.4 kJ mol<sup>−1</sup>) and in dark conditions (93.6 kJ mol<sup>−1</sup>), and the difference between the two shows that the light-induced energy plays a major role in the high performance at lower temperatures. The authors investigated the reaction mechanism and they proved that hot carriers are generated by the d to s interband excitation of Pd NPs *via* UV light irradiation. These hot carriers are transferred into the conductive band (CB) of Sr<sub>2</sub>Ta<sub>2</sub>O<sub>7</sub>, and reduce O<sub>2</sub> to radical species. In parallel, hot holes in the Pd d-band activate

CH<sub>4</sub> for cracking. Finally, the authors found that the ultrawide-bandgap (UWBG) of the semiconductor, Sr<sub>2</sub>Ta<sub>2</sub>O<sub>7</sub>, is crucial to drive the photo-thermal reaction towards the desired product.

Very recently, the research group of Miyauchi expanded the investigation in photo-thermal DRM with two papers, in which the researchers studied the role of Rh as a plasmonic metal with two different supports, namely SrTiO<sub>3</sub> and TaON.<sup>245,246</sup> Surprisingly, in both cases the catalytic system tested exceeded the theoretical thermal limit calculated by using thermodynamic simulation software. Shoji *et al.* observed that the Rh based catalytic system supported on strontium titanate (Rh/STO) irradiated with a 150 W Hg–Xe lamp converts CH<sub>4</sub> and CO<sub>2</sub> giving H<sub>2</sub> with a yield higher than 50% even at a temperature of 200 °C, while the performance in dark conditions was less than half even at 450 °C.<sup>245</sup> In addition, they tested six different semiconductors as a support for Rh NPs and Ni/Al<sub>2</sub>O<sub>3</sub> as catalytic systems only with photoirradiation, and they found that the band gap and the conduction band position of STO are optimal to drive CH<sub>4</sub> and CO<sub>2</sub> conversion. To understand thermal and photophysical contribution in this process, the authors performed the reaction under light irradiation while monitoring the working temperature, *i.e.* the temperature at the surface of the catalyst. They observed indeed that the temperature reached 300 °C; however this thermal energy is quite low to justify the high yield. *In situ* electron spin resonance (ESR) measurements conducted under light irradiation and at −173 °C proved that trapped electrons and trapped holes are generated in the valence band, that Rh NPs accept electrons in the DRM process, and that CH<sub>4</sub> interacts with the holes while CO<sub>2</sub> does not. In addition, Miyauchi's research group performed several experiments to elucidate the mechanism involved in this process: Kelvin probe force microscopy (KPFM) shows that the charge transfer process occurs at the semiconductor–metal interface; the correlation between light intensity and hydrogen generation rate shows that photoexcited charge carriers play the main role to drive the reaction; the results of the action spectrum obtained below 380 nm, corresponding to the bandgap energy of STO (3.2 eV), shows that the promotion of electrons from the valence band to the conductive band is also a driving force of the reaction. Catalytic tests conducted in the presence of a single gas, and the use of isotopic labeling allowed one to understand in detail the formation of CO. All this evidence allowed the authors to present a detailed mechanism for this photo-thermal DRM process (Fig. 15).

Cho *et al.* studied the catalytic activity of rhodium-loaded tantalum oxynitride (Rh/TaON).<sup>246</sup> In this case, the H<sub>2</sub> yield at 450 °C under visible-light irradiation is 50.3%, a value that is significantly higher than the one obtained in dark conditions (15.6%), and higher than the theoretical thermal limit calculated by using thermodynamic simulation software (36.3%). The authors extended their studies to different combinations of metal and semiconductor support, and they confirmed that the interaction between these two components is pivotal to achieve the best results. Finally, they performed different experiments, in which they could monitor the activity of Rh/TaON with respect to the temperature generated by either light irradiation or external heating, and they proved that the high performance is not only due to the photogenerated thermal process, but it



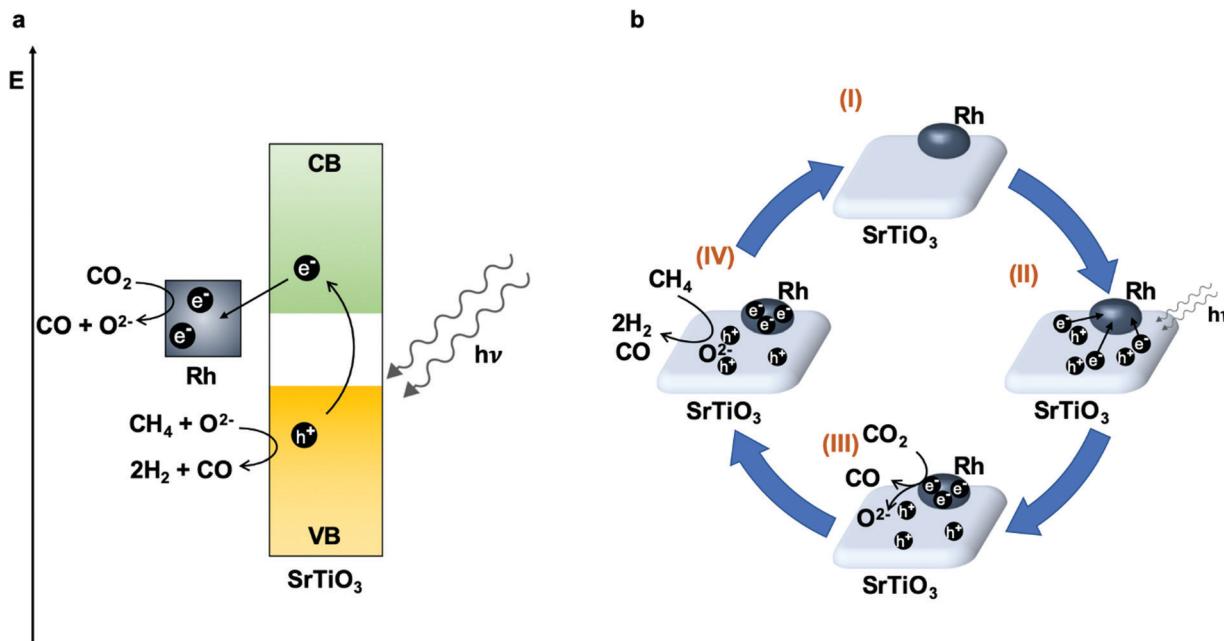


Fig. 15 Mechanism of the photocatalytic DRM by Rh/STO. (a) Band diagram showing generation of electron–hole pairs in STO and expected redox reactions. (b) Schematic dynamics of charge carriers and oxygen ions. Adapted with permission from ref. 245. Copyright (2020) Springer Nature.

also involves the contribution of the plasmonic effect on Rh NPs and of hot electron–hole pairs in the support.

Inspired by the works of Ye's group on bimetallic-alloy plasmonic photocatalysts, Halas and co-workers prepared a series of Cu nanoparticle–Ru single-atom catalysts supported onto a MgO–Al<sub>2</sub>O<sub>3</sub> composite (Cu<sub>x</sub>Ru<sub>y</sub>) for photo-thermal DRM.<sup>116,235,247</sup> Cu<sub>19.9</sub>Ru<sub>0.1</sub> and Cu<sub>19.8</sub>Ru<sub>0.2</sub> showed the best stability due to coking resistance, with a minimal decrease in activity only after 50 hours, and Cu<sub>19.5</sub>Ru<sub>0.5</sub> instead showed higher catalytic performance, but with a decrease in activity after 5 hours. Cu<sub>19.9</sub>Ru<sub>0.1</sub> and Cu<sub>19.8</sub>Ru<sub>0.2</sub> also showed higher selectivity compared to Cu<sub>19.5</sub>Ru<sub>0.5</sub>. The authors were able to prove that Ru atoms can be monodispersed on the Cu NPs in Cu<sub>19.9</sub>Ru<sub>0.1</sub> and Cu<sub>19.8</sub>Ru<sub>0.2</sub>, and they attributed their superior performances to the presence of these atomically dispersed Ru sites. The researchers investigated the mechanism comparing the activity of the reaction under light irradiation and in dark conditions adding external heating, and they found that the DRM process is mainly driven by the generation of hot carriers, which promote the C–H activation on Ru sites and the release of H<sub>2</sub>. Finally, the study of the energy efficiency, expressed as the percentage of methane converted, proved even in this case that the photo-thermal induced process performs better than the thermal process having the same surface temperature.

#### 4.4. NH<sub>3</sub> synthesis/decomposition

The Haber–Bosch catalytic process is considered as one of the greatest game changer discoveries of the 20th century.<sup>248</sup> When this synthesis to produce ammonia starting from its elements, hydrogen and nitrogen, was performed in BASF's laboratories, the operative conditions were really extreme. Since then, this century old process has been optimized in many aspects and is still applied worldwide to provide ammonia, a key chemical for the

world's agricultural industry.<sup>249</sup> However, this reaction is severely hampered by thermodynamics, and to achieve competitive industrial performance, the reactors need to operate at high temperatures (400 °C to 500 °C) and pressures (100 bar to 200 bar).

In 2018, Zhang and co-workers presented a new strategy to achieve a Haber–Bosch comparable process based on sunlight only as the energy source.<sup>250</sup> They prepared a catalytic system where K promoted Ru is supported on disordered TiO<sub>2-x</sub>H<sub>x</sub> (K/Ru/TiO<sub>2-x</sub>H<sub>x</sub>). This catalyst could generate ammonia simply by irradiation with sunlight at atmospheric pressure, without the need of external heating. K/Ru/TiO<sub>2-x</sub>H<sub>x</sub> is able to absorb in the UV-vis-NIR region of sunlight, and it can convert part of it to thermal energy; in fact, after 13 minutes of irradiation with a 300 W xenon lamp, the reactor reached a temperature of 360 °C. As a result, the process yields an ammonia generation rate of 112.6 μmol g<sup>-1</sup> h<sup>-1</sup> (gas flow rate of 6 mL min<sup>-1</sup>, TOF  $3.9 \times 10^{-4}$  s<sup>-1</sup>), which is surprisingly high when compared with the same and similar Ru based catalyst used under thermal conditions. In their work, Mao *et al.* highlight the importance of the surface disorder of TiO<sub>2-x</sub>H<sub>x</sub> to avoid H<sub>2</sub> poisoning of Ru. In addition, they propose a mechanism, which involves the localized surface plasmon resonance (LSPR) effect of Ru facilitated by coupling with the support, TiO<sub>2-x</sub>H<sub>x</sub>, and which lowers activation barriers.

In the same year, Halas *et al.* presented an interesting study distinguishing the role of thermal and photochemical effects in reducing the activation barrier for ammonia decomposition.<sup>115</sup> Particularly, the researchers tested a plasmonic Cu–Ru antenna-reactor (Cu–Ru-AR) as a photocatalyst, and they found out that the reaction rate was directly proportional to light irradiation. They were able to differentiate the plasmon-induced hot carriers and the photo-thermal heating contribution, and to mainly attribute to plasmon-induced hot carriers the ability of catalyzing NH<sub>3</sub>

decomposition. In addition, they calculated the activation energy ( $E_a$ ) as a function of the irradiation wavelength and intensity, and they plotted the different values, obtained using the Arrhenius equation, obtaining a three-dimensional map, which shows that under optimized illumination conditions  $E_a$  can be reduced to 0.27 eV. This mapping proved to be a useful tool that provides information to predict and optimize catalytic performance.

One year later, Halas and co-workers published another paper in the same field, and they demonstrated an innovative route to produce ammonia using titanium nitride NPs as plasmonic antennas, which can heat magnesium-based nanomaterials to undergo chemical looping as active materials.<sup>251</sup> The researchers manufactured a customized low-pressure reaction chamber to perform the three-step chemical looping process and to monitor the product formation *in situ* via isotopic labelling and rotational spectroscopy. During the experiment, the catalytic system was irradiated with a Ti:sapphire laser (808 nm wavelength, 150 fs pulses, 80 MHz repetition rate) with different timing for each step. Sweareer *et al.* studied each stage of the looping process, and they proved that TiN plasmonic NPs play a crucial role in photo-thermal heating, which led to the increase of the surface temperature of the catalytic system above 550 °C. Furthermore, they succeeded in detecting ammonia evolution in the final step of the photo-thermal-promoted chemical looping. As the authors highlight in the conclusion, this process needs further improvements; however its advantages compared to classical strategies are evident, making this approach a promising green option for ammonia production.

Another innovative approach to produce ammonia without using thermal energy or elevated pressures was proposed by Liu and co-workers in 2019.<sup>252</sup> Their strategy exploits photo-thermal heating in order to create and control a thermal gradient, which allows achieving high reaction rates and good conversion yields. For their investigation, Li *et al.* built a custom-made fixed-bed reactor equipped with a quartz window and two embedded thermocouples to monitor the temperature of the catalyst on the top-surface ( $T_1$ ) and on the bottom-surface ( $T_2$ ). The catalytic system consists of cesium-promoted ruthenium NPs supported on magnesium oxide (Ru–Cs/MgO); the powdered catalyst is loosely packed in the reactor, and it has a thickness of 3 mm. After reduction of the Ru catalyst, N<sub>2</sub> and H<sub>2</sub> were introduced in the reactor with a 1:3 ratio and a flow rate of 75 mL min<sup>-1</sup>. NH<sub>3</sub> was produced giving the best results under irradiation from the top using UV and visible LEDs, at atmospheric pressure without adding external heating. The researchers observed that in these conditions photo-thermal heating forms a negative thermal gradient, meaning that  $T_1$  is higher than  $T_2$ , within the vertical temperature profile of the packed catalyst. They demonstrated that this negative thermal gradient plays a crucial role in improving the reaction rate and the product yield, because it acts as a thermodynamic pump, which is able to shift the equilibrium towards the products. The hotter temperature on the top of the catalyst accelerates nitrogen scission, the rate-limiting step, while the product is pushed to the coldest region due to thermophoretic forces, and NH<sub>3</sub> does not undergo decomposition reaction.<sup>253</sup>

Concluding this section, we consider worth mentioning three publications in the field of plasmonic assisted ammonia synthesis,

although the authors do not discuss thermal contribution.<sup>254–256</sup> The work presented in 2016 by Misawa and co-workers describes the production of ammonia and oxygen starting from nitrogen in water under visible light irradiation promoted by plasmonic induced charge separation.<sup>254</sup> The researchers coated a niobium-doped strontium titanate photoelectrode with a thin layer composed of gold NPs, as a plasmonic metal, and with zirconium/zirconium oxide as a cocatalyst (AuNP/Nb–SrTiO<sub>3</sub>/Zr). The catalytic reaction was performed under ambient conditions, and the photoelectrode was irradiated using a xenon lamp using arbitrary light intensity and wavelength. Then the quantity of ammonia formed was quantitatively analyzed by GC-MS, and expressed as apparent quantum yield (AQY, %), *i.e.* the percentage of the ratio between reaction rate of NH<sub>3</sub> formation and the flux of incident photon. Although the efficiency of this method is still not competitive for industrial applications, it proved to be highly selective and highly responsive to visible light. These aspects give high expectations for further development of sunlight-driven ammonia synthesis.

Gong and co-workers published in 2018 a work, where they described surface oxygen vacancies (O<sub>vac</sub>) together with plasmon generation as a winning strategy for N<sub>2</sub> fixation based on environmentally friendly methods.<sup>255</sup> As in the previously reported example, the plasmonic metal of interest is gold, because it can broaden the light absorption range of the metal oxide support, TiO<sub>2</sub>. The optimized catalytic system consists of Au nanorods (NR) supported on rutile TiO<sub>2</sub> modified on the surface by atomic layer deposition of amorphous TiO<sub>2</sub> (TiO<sub>2</sub>/Au/a-TiO<sub>2</sub>). This modification introduces more defects, or O<sub>vac</sub>, that promote the adsorption of N<sub>2</sub> and its activation *via* hot electrons, without modifying the bulk properties. This was proved by testing the catalytic reaction by immersing different electrodes in water with continuous nitrogen bubbling under UV-visible light irradiation. The NH<sub>3</sub> production rate increased with the enhancing of the surface modification of the electrodes: bare TiO<sub>2</sub> < TiO<sub>2</sub>/AuNRs < TiO<sub>2</sub>/a-TiO<sub>2</sub> < TiO<sub>2</sub>/Au/a-TiO<sub>2</sub>. With this paper Li *et al.* suggest a promising method to combine electrochemical ammonia synthesis with photovoltaic energy, exploiting a plasmon-enhanced reaction pathway.

Xiong and co-workers published very recently a paper about nitrogen fixation, where LSPR promotes nitrogen fixation through a N≡N dissociative mechanism.<sup>256</sup> This reaction pathway is normally unusual, since the nitrogen triple bond has a high dissociation energy. The series of catalysts that Hu *et al.* prepared, AuRu<sub>x</sub> ( $x = 0.14, 0.23, 0.31, 0.39$ ; it represents the percentage of Ru incorporated into Au NPs), contain Au in the core, because it absorbs efficiently a wide region of light, and Ru on the surface, because it facilitates N<sub>2</sub> dissociation. AuRu<sub>0.31</sub> was the catalyst that gave the best performance achieving a production rate of 101.4 μmol g<sup>-1</sup> h<sup>-1</sup> at room temperature and 2 bar of N<sub>2</sub> pressure, under full-spectrum light irradiation (300 W Xe lamp, 400 mW cm<sup>-2</sup>). In addition, they performed electron spin resonance (ESR) to detect the reactive oxygen species, while catalytic N<sub>2</sub> reduction by altering photon flux and wavelength-dependent apparent quantum efficiencies (AQEs) proved the plasmonic-driven mechanism. The researchers employed several techniques to follow carefully the different steps that N<sub>2</sub> undergoes to generate NH<sub>3</sub>.



They used  $^{15}\text{N}_2$  to verify the nitrogen source, they performed low-temperature Fourier-transform infrared (LT-FTIR) spectroscopy to identify the configuration of the adsorbed  $\text{N}_2$ , and they brought evidence of the surface plasmon-driven  $\text{N}_2$  dissociation mechanism owing to *in situ* diffuse reflectance infrared Fourier-transform spectroscopy (DRIFTS) and *in situ* near ambient pressure X-ray photoelectron spectroscopy (NAP-XPS).

#### 4.5. Water gas shift reaction

Water gas shift reaction (WGS) is a key process for  $\text{H}_2$  production, as well as for ammonia or methanol synthesis, Fischer-Tropsch process, reforming systems for fuel cells or even one of the key steps in the automobile exhaust processes, producing  $\text{H}_2$  for the removal of toxic  $\text{NO}_x$  gases.<sup>257,258</sup> This reaction was discovered by Felice Fontana in 1780 and it is considered as one of the most important reactions for industry and in the history of Chemistry.<sup>259</sup>

The reaction is a reversible and moderately exothermic process that permits, simultaneously, CO conversion into  $\text{CO}_2$  and  $\text{H}_2$  production using water as the H source (eqn (7)):



Liu and co-workers described several Cu-based photo-thermal catalysts to perform WGS reaction.<sup>260</sup> The study was carried out with 50 kPa of CO and 3 mL of water in a glass reactor with a closed gas circulation system, using a Xe lamp as the light source (300 W and 1 sun). The catalysts were based on Cu supported on  $\text{Al}_2\text{O}_3$ ,  $\text{ZnO}$  and a mixture of them ( $\text{Al}_2\text{O}_3$ – $\text{ZnO}$ ). The characterization of the catalysts revealed that all of them presented the same Cu loading. Besides, XRD and XPS demonstrated that  $\text{Cu}/\text{Al}_2\text{O}_3$  presented amorphous alumina and  $\text{CuO}$  phases without any metallic Cu. Meanwhile in the other two cases, metallic Cu NPs and crystalline  $\text{ZnO}$  can be observed. According to the UV-visible absorption spectra of the samples,  $\text{Cu}/\text{Al}_2\text{O}_3$  presented the most intense light absorption in spite of the fact that the characteristic Cu LSPR band was not detected. In contrast, this plasmon band centered at 560 nm appeared in the case of Cu/ $\text{ZnO}$  and  $\text{Cu}/\text{Al}_2\text{O}_3$ – $\text{ZnO}$ . The catalytic activity of different molar ratios of Cu, Al and Zn was studied obtaining the best results with the catalysts having 1:0.6:0.4 Cu:Zn:Al (4.0 mmol  $\text{H}_2 \text{ h}^{-1}$ ). However, these results did not improve the catalytic performance of the catalysts with pure  $\text{ZnO}$  as a support (Cu/ $\text{ZnO}$ ) that yielded a hydrogen production of 4.75 mmol  $\text{H}_2 \text{ h}^{-1}$  and a CO conversion of 85.6%. On the other hand, the pure  $\text{Al}_2\text{O}_3$  catalyst (Cu/ $\text{Al}_2\text{O}_3$ ) was less active (2.1 mmol  $\text{H}_2 \text{ h}^{-1}$ ). Furthermore, the study evaluated the time-dependent temperature increase revealing higher temperatures than 300 °C after three minutes of irradiation when using Cu/ $\text{ZnO}$ , Cu/ $\text{Zn}_{0.6}\text{Al}_{0.4}$ , Cu/ $\text{Al}_2\text{O}_3$  and pure Cu NPs; meanwhile bare supports did not present higher temperatures than 150 °C, probably due to the absence of metallic Cu NPs. It is worth mentioning that Cu/ $\text{Al}_2\text{O}_3$  reached the highest temperature (367 °C) with the highest photon absorption together with the plasmonic effect of metallic Cu NPs. The study, however, needs a deeper analysis of the photo and thermal approaches due to the existence of controversial results. The temperature study displayed the LSPR effect from plasmonic Cu, but the temperature was not indeed described

as the dominant variable in the catalytic performance as the three catalysts presented higher temperatures than 250 °C (usual for WGS reaction at low temperatures), and the catalyst which reached higher temperatures was not the most active one. As more Cu is detected in Cu/ $\text{ZnO}$  and Cu/ $\text{Zn}_{0.6}\text{Al}_{0.4}$  catalysts than in Cu/ $\text{Al}_2\text{O}_3$ , it seems that the presence of  $\text{ZnO}$  improves the dispersion of Cu, offering a higher amount of active plasmonic sites. Finally, the authors also studied different reducing pretreatments for several samples obtaining as a better option the photo-reduction method, using also light for the reduction of the catalysts prior to their use in the photo-thermal reaction.

In the same line, Cu supported on  $\text{Al}_2\text{O}_3$  ( $\text{CuO}_x/\text{Al}_2\text{O}_3$ ) was compared with Pd and Au supported on the same material as catalysts for WGS reaction in a closed-circulated system under light irradiation (UV-visible-IR).<sup>261</sup> The catalysts with 2 wt.% loading of each metal performed different, ordered as Cu > Pt > Au (24.6, 6.3 and 3.5  $\mu\text{mol H}_2 \text{ g}_{\text{cat}}^{-1} \text{ s}^{-1}$  respectively). Focusing on the most active and also cheapest metal, the loading of Cu was optimized up to 19 wt% revealing four times more catalytic activity, an improvement up to 122  $\mu\text{mol H}_2 \text{ g}_{\text{cat}}^{-1} \text{ s}^{-1}$ . The characterization of this Cu-based catalyst (XRD and HRTEM) showed mostly metallic Cu NPs (well-dispersed) and some  $\text{Cu}_2\text{O}$  after the WGS reaction. The temperature reached during photoreaction was 285 °C offering twice the activity of usual thermal catalysis which occurs at 300 °C. This comparison demonstrated additional effects apart from the just heating of the catalyst (not only thermocatalysis). To elucidate the additional effects, several catalytic tests were carried out using light-filters: those filtering IR radiation and reducing the light-to-thermal conversion (only up to 60 °C) presented 54.6% less catalytic activity after 30 min; the other filters removed UV contribution and the catalytic activity was decreased to 48.4% achieving the same temperature as in full-spectrum experiments. Summarizing, these results revealed the photo-thermal oxidation of CO and reduction of  $\text{H}_2\text{O}$  as a coupled and synergistic process, a cooperative action between photo-generated carriers and photo-induced heat.

#### 4.6. $\text{H}_2$ evolution and other processes driven by photo-thermal catalysis

The approach of using solar energy to improve reactions that normally require high temperature and pressure has been exploited successfully in several catalytic processes. Hydrogen evolution generated *via* water splitting reaction is indeed among these processes, and the first paper in this field was published in 2008.<sup>262</sup> Li and Lu reported the formation of  $\text{H}_2$  from water steam splitting over titanium disilicide ( $\text{TiSi}_2$ ) and Pt supported on titanium oxide ( $\text{Pt}/\text{TiO}_2$ ) under visible light irradiation, and they studied the influence of pressure and temperature on the catalytic activity. The researchers carried out the reaction in a Pyrex flask, which was irradiated with a 300 W tungsten halogen lamp. They observed that for both catalysts the increase of pressure is essential to generate  $\text{H}_2$  at 190 °C with good yields.

This topic gained attention again last year; in fact three different research groups studied water splitting reaction for hydrogen generation catalyzed by nickel-based materials.<sup>263–265</sup>



Zhang and co-workers combined the photo-thermal effect with electrocatalysis for hydrogen evolution reaction (HER) and oxygen evolution reaction (OER).<sup>263</sup> They prepared Ni NPs supported on reduced graphene oxide nanosheets (Ni/RGO) *via* hydrothermal synthesis and reduction in a H<sub>2</sub>/Ar atmosphere at 500 °C. The researchers compared the performance of the catalytic system under irradiation and in dark conditions, and they observed that light improves both the thermodynamics and the kinetics of the reaction, yielding better HER and OER activity. They explained this behavior with the ability of the catalyst to generate hot active species after light absorption, which increase the system's temperature to 50 °C in less than 5 minutes. Guo and co-workers observed a similar behavior, when they investigated nickel–tungsten–boride nanoneedles (Ni–W–B) supported on carbon cloth (CC) for photo-thermal catalysis in electrochemical water splitting.<sup>264</sup> They tested the catalytic system for HER with and without irradiation, and they also compared the results with those obtained from bare CC and Pt–C deposited on CC. As expected, HER activity is facilitated by light irradiation, without an increase of the electrolyte temperature. In addition, the authors performed the catalytic process at different temperatures, and they observed that the performance of the photo-thermal HER was similar to the performance achieved at 48 °C. As far as the OER is concerned, the activity of Ni–W–B/CC was compared with bare CC and IrO<sub>2</sub> supported on CC, and even in this case the Ni-based photo-thermal catalyst showed the best results. Finally, the researchers assembled an electrolyzer using Ni–W–B/CC as the cathode and anode, and they demonstrated its superior overall photo-thermal-assisted water splitting performance among the reported bifunctional catalysts. Even Liu and co-workers used abundant and cheap nickel to assemble a novel photo-thermal-enhanced electrochemical device for water splitting.<sup>265</sup> They prepared a Ni nanosheets array supported on an Al<sub>2</sub>O<sub>3</sub> ceramic chip, and they characterized thoroughly its electrochemical properties by performing electrocatalytic HER activity, electrochemical impedance spectroscopy (EIS), cyclic voltammetry (CV) and durability tests. Then, the researchers assembled a two-electrode device to investigate the overall water splitting during exposure to simulated sunlight, using the supported Ni nanosheets as the cathode, for HER, and NiFe hydroxide film on CC as the anode, for OER. The good production rates of hydrogen and oxygen make this device promising, although the authors suggest that they aim to improve the process with modifications in order to avoid thermal diffusion between the two ends of the device.

An alternative process to generate hydrogen proposed in this current year by Dai and by Naldoni is the photo-thermally enhanced hydrolysis of ammonia borane, whose hydrogen content is 19.6%.<sup>266–268</sup> Wu *et al.* prepared a series of urchin-like sodium titanate microspheres supported on the reduced graphene oxide (RGO/Na<sub>2</sub>Ti<sub>3</sub>O<sub>7</sub>) photo-thermal catalyst with different graphene oxide content (0, 33, 50, 60, 66 wt%).<sup>266</sup> Particularly, they focused on the challenging aspect of synthesis optimization, elucidating the role of graphene oxide sheets in the growth mechanism of Na<sub>2</sub>Ti<sub>3</sub>O<sub>7</sub> microspheres. Also, they found that the photoelectrochemical

performance of the prepared materials was related to the structural organization of the catalyst and the surface features of the Na<sub>2</sub>Ti<sub>3</sub>O<sub>7</sub> microspheres. In fact, the branches of the urchin-like microspheres are crucial because they provide plenty of active sites, and they facilitate electron transfer at the interfaces. Then the authors studied the photo-thermal effect and its influence on the catalytic hydrolysis of ammonia borane. They observed that the high results in H<sub>2</sub> evolution are a combination of photo-thermal properties of RGO and the ability of Na<sub>2</sub>Ti<sub>3</sub>O<sub>7</sub> to generate hot spots. Finally, the researchers proved that the morphology of Na<sub>2</sub>Ti<sub>3</sub>O<sub>7</sub> is also important, since they observed that RGO/Na<sub>2</sub>Ti<sub>3</sub>O<sub>7</sub> microspheres have almost 2 times higher activity than RGO/Na<sub>2</sub>Ti<sub>3</sub>O<sub>7</sub> nanowires. The synergistic effect between hot carriers and photo-thermal properties is also the driving mechanism which justifies the results reported by Naldoni and co-workers.<sup>267</sup> They developed a procedure to decorate TiN nanocubes with Pt nanocrystals (TiN–Pt nanohybrids); they fully characterized the obtained materials, and they evaluated their catalytic performance for photo-thermal hydrogen evolution from ammonia borane dehydrogenation. As expected, visible-NIR light irradiation increases hydrogen production, compared to the results obtained in dark conditions, due to plasmonic generation of hot electrons. When the irradiation wavelength matches the LSPR, *i.e.* 700 nm, the value of apparent quantum yield (AQY) reaches 120%. This high value is explained with the high absorption ability of the catalysts and the low energy barrier of the dehydrogenation reaction, aspects that allow both hot electrons and low energy hot electrons (those which underwent inelastic scattering) to react efficiently with ammonia borane. In addition, the authors were able to detect with good approximation the temperature at the surface of the nanoparticles in the liquid medium, and to evaluate the photo-thermal heating accordingly. Finally, they demonstrated by performing theoretical COMSOL simulations the contribution of photo-thermal heating to the reaction rate, and they elucidated in detail the mechanism for the cleavage of the O–H bond, the rate determining step of the reaction, by hot electrons.

To conclude this part of the review, we believe that there are other noteworthy works, which did not fit in the previous sections. In fact, in the last few decades photo-thermal catalysis has contributed to the development of catalysis field bringing a different perspective to many catalytic reactions. One of the first examples was reported in 1998 by Kennedy and Datye, who described how the combination of a typical thermal oxidation catalyst, Pt, with a photocatalyst, TiO<sub>2</sub>, can improve the efficiency of ethanol oxidation.<sup>269</sup> Almost twenty years later, Amal and co-workers investigated the mechanism of the same reaction using a TiO<sub>2</sub>-supported Au catalyst, and they reported an almost doubled catalytic activity in the photo-thermal regime.<sup>270</sup> In the same year, Yang and co-workers reported the synthesis of AgPt bimetallic hollow nanoparticles, and they tested this material for LSPR-enhanced methanol oxidation.<sup>271</sup> Linic and co-workers prepared another photo-thermal catalyst consisting of silver nanocubes supported on Al<sub>2</sub>O<sub>3</sub>, and they tested it for commercially relevant catalytic oxidations, which are normally performed at high temperatures and over expensive platinum-group metals.<sup>49</sup> Zhang and co-workers investigated the effect and the mechanism of the photo-thermal enhanced catalytic oxidation of formaldehyde



promoted by a graphene-based  $\text{MnO}_2$  composite.<sup>272</sup> Another important application of photo-thermal catalysis is removal of VOCs (volatile organic compounds) *via* catalytic oxidation, as reported by Jia and coworkers.<sup>273</sup> They tested four  $\text{ACo}_2\text{O}_4$  (A = Ni, Cu, Fe, Mn) type spinels for photo-thermal degradation of toluene, and they observed the superior activity of  $\text{NiCo}_2\text{O}_4$  due to better structural and photophysical features. Xu and co-workers reported the synthesis of a  $\text{SmCoO}_3/\text{SBA-15}$  composite *via* an impregnation method, which confers to the material an outstanding surface area.<sup>37</sup> This feature, together with the good photo-thermal conversion, makes  $\text{SmCoO}_3/\text{SBA-15}$  a promising candidate for photo-thermal degradation of propane.

Finally, to have a more comprehensive overview on the advantages of plasmon-driven processes, we suggest three papers, which focus on plasmon induced crystal growth and material transformation.<sup>274–276</sup> Xu and co-workers give a closer look at the different transformation processes that can occur among molecules when they are in the proximity of a plasmonic electromagnetic field, hot carriers, and heat.<sup>274</sup> In their work, they explain the different paths that allow the energy transfer that induce the chemical transformation, as well as a significant contribution to plan efficient and tunable plasmonic catalytic systems. Zheng and his research group presented a nice example of an *in situ* transformation mediated by Au NPs.<sup>275</sup> A polycrystalline material,  $\text{NaYF}_4:\text{Eu}^{3+}$ , can be converted into a single crystal luminescent material,  $\text{Y}_2\text{O}_3:\text{Eu}^{3+}$ , in a rapid and controlled way. Sasaki and co-workers demonstrated that they were able to control the synthesis of a  $\text{ZnO}$  based plasmonic hybrid nanostructure in a targeted position.<sup>276</sup> They realized a nanobutterfly gold structure, which directs plasmonic hotspots in a precise location, according to the direction of the incident light polarization. Thanks to this procedure,  $\text{ZnO}$  layers can selectively grow *via* hydrothermal synthesis where required, leaving the surrounding gold particles uncoated, and yielding the desired plasmonic hybrid nanostructure.

## 5. Design strategies for photo-thermal catalysts

For an optimal exploitation of the photo-thermal effect in catalytic processes, photocatalysts must meet a series of requirements, including an intense light absorption across the solar spectrum, an efficient charge carrier generation and a high capacity for heat production and/or heat transfer.<sup>30</sup> In the subsequent sections, we outline different strategies for materials' design and engineering in order to maximize photo-thermal effects and therefore efficiency for a given mechanism, *i.e.* hot electron or local heating effects.

### 5.1. Size and shape effects

In plasmonic materials, size and shape are key factors determining the LSPR phenomenon (*i.e.* ratio of absorption to scattering, number of LSPR modes and position of the LSPR peak) and can be adjusted appropriately in order to match the plasmon band with the solar spectrum and thereby improve catalytic performance.<sup>277,278</sup> In fact, both parameters have shown a strong

influence not only on the capability for charge carrier generation, but also on the thermal properties of a material.

Theoretical calculations have demonstrated that larger plasmonic NPs provide a larger amount of hot carriers; however, they display energies close to the Fermi level, since most of the absorbed energy is dissipated either by scattering or by heating.<sup>9</sup> In contrast, smaller NPs render less electron-hole pairs than their counterparts, but they exhibit higher energies but showing very short lifetimes in the order of femtoseconds.<sup>279</sup> This is of crucial importance in photo-thermal catalysis, as the higher the energy of hot electrons, the higher the probability for them to populate acceptor states or overcome Schottky barriers. Therefore, it seems that smaller NPs are especially convenient for hot electron induced catalysis. However, smaller NPs do not have to necessarily enhance the overall photo-thermal performance as the reduction of particle size can lead to inefficient charge carrier separation due to spatial confinement.<sup>280</sup> In addition, plasmonic NPs have shown a red shift in their resonant frequency as the particle size increases, thus allowing a more efficient harvesting of the visible and IR regions of sunlight compared to their smaller equivalents.<sup>281,282</sup> Given these circumstances, a precise balance between these effects should be considered when developing photo-thermal catalysts.

Morphology is another key parameter determining the optical properties of nanostructures, especially those featuring the LSPR effect. In this regard, certain shapes such as bars, bipyramids or branched structures have shown a red-shift in the position of their main LSPR absorption peaks, therefore being more appropriate to take advantage of the low energy regions of the solar spectrum.<sup>283–285</sup> Theoretical studies from Goriely and co-workers pointed out that the ability for the generation of hot carriers was also closely related to the shape of plasmonic structures. In these works, the authors established that geometries displaying strong and inhomogeneous electric fields together with important confinement effects could generate high-energy hot electrons and holes. In contrast, those geometries without defects or featuring uniform electric fields like thin films or nanospheres afforded a less efficient production of hot carriers.<sup>286,287</sup>

The thermal conductivity of nanostructures can also be optimized by modifying phonon transport properties, by either size or shape engineering. For instance, it has been demonstrated that the thermal conductivity of a material can be drastically reduced by introducing a high density of grain boundaries which intensifies phonon scattering.<sup>288</sup> In line with this, increasing the surface roughness has also proven to slow down phonon velocity.<sup>289</sup> Then, decreasing the size of certain components of the photo-thermal system and maximizing its surface roughness contributes to the reduction of the thermal conductivity of the material, consequently favoring heat generation close to catalytic active sites.<sup>30</sup>

It should be also emphasized that size has direct consequences on the heating capability of plasmonic structures. In fact, as mentioned in the previous sections, the plasmon-heating effect in metallic nanospheres scales quadratically with the NP radius. However, larger plasmonic NPs are more prone to decay radiatively, thus reducing the amount of locally dissipated heat.<sup>290</sup> For these reasons, an accurate balance between these two opposite effects has



to be pondered when tailoring new photo-thermal catalysts. In addition, according to theoretical works developed by Baffou and co-workers, morphology also seems to play a key role when evaluating the heat production in plasmonic nanostructures.<sup>291</sup> For example, the authors found that sharp, flat or elongated geometries exhibited a more intense heat generation than spheres. In the case of nanospheres, only the outermost part of the structure effectively faced the incoming light, while in the rest of the cases the whole volume of the structure participated in the heating process. Apart from this, heat production also benefited from the presence of corners and edges that enhanced electric fields locally.<sup>292</sup>

## 5.2. Hybrid materials

A very common and straightforward strategy to improve photo-thermal performance is the design of hybrid structures. In these composite materials, the combination of two or more components with different optical, electronic and thermal properties works synergistically to enhance catalytic rates. Typically, photo-thermal systems are comprised of an inorganic support or host (in most of the cases metal oxides) functionalized by either plasmonic or non-plasmonic metallic NPs. It is worth mentioning that each of the components is not restricted to a single function inside the photo-thermal structure and that, depending on the nature of each system, only one or more elements can act as catalytic active sites in the reaction.

**5.2.1. Metal/semiconductor structures.** As stated above, metal/semiconductor hybrid materials are one of the most studied systems. An ideal semiconductor support for photo-thermal applications should display a high surface area together with a strong broadband optical absorption that enables the generation of charge carriers and/or high local temperatures. These types of structures generally include narrow band gap metal oxide semiconductors with a high degree of defects and mid-bandgap states that ultimately enhance the absorption of light in the low energy regions of the solar spectrum. Examples of this type of structures are non-stoichiometric materials with oxygen/metal vacancies, mixed metal valences and/or conduction electrons.<sup>30,293–295</sup> An appealing precedent of this approach was recently reported by Ozin and co-workers in the photo-thermal hydrogenation of CO<sub>2</sub> using the non-stoichiometric black In<sub>2</sub>O<sub>3-x</sub>/In<sub>2</sub>O<sub>3</sub> catalyst.<sup>46</sup> The authors found that the hydrogenation of pale yellow In<sub>2</sub>O<sub>3</sub> at different temperatures led to the generation of black non-stoichiometric heterostructures composed of In<sub>2</sub>O<sub>3-x</sub> domains in crystalline stoichiometric In<sub>2</sub>O<sub>3</sub>. These heterostructures displayed a broad absorption across the entire solar spectrum that provided an excellent photo-thermal performance. On the one hand, light irradiation caused a strong local heating in the In<sub>2</sub>O<sub>3-x</sub>/In<sub>2</sub>O<sub>3</sub> composite that catalyzed the thermochemical hydrogenation of CO<sub>2</sub>. On the other hand, photo-generated electrons could also promote photochemical reduction of CO<sub>2</sub> due to the presence of mid-gap states. The effective combination of both reaction pathways gave rise to an incredibly high activity and 100% selectivity towards CO production under long-term conditions of more than 70 h. Therefore, the use of this type of broadband absorption semiconductors, alone or in combination with metallic active sites, offers the possibility for excellent photo-thermal systems exhibiting noteworthy reaction rates.

**5.2.2. Metal/porous support structures.** Materials with high porosity and high surface area such as MOFs and zeolites have been widely utilized as supports or hosts for metal NPs in a vast variety of catalytic applications. However, photo-thermal catalysis over this type of materials is a much less explored discipline and it appears to be an emerging and promising field of research.

Recently, Jiang *et al.* described for the first time the preparation of a composite catalyst based on a porphyrinic MOF displaying the photo-thermal effect.<sup>296</sup> The authors stabilized Pt nanocrystals on the MOF structure and the composite photocatalyst was tested for the oxidation of benzyl alcohol with molecular O<sub>2</sub>. Interestingly, according to the catalytic results a synergy between the photo-thermal effect of both the MOF and supported Pt nanocrystals was established. This cooperative work enhanced the photocatalytic rate of the composite compared to the individual contributions of each component separately.

More recently in 2018, Maspoch and co-workers studied the photo-thermal effect of a series of MOFs representing the major subfamilies of this type of porous materials and calculated the photo-thermal transduction efficiency for each of the MOFs under study.<sup>297</sup> For example, UiO-66-NH<sub>2</sub> and CPO-27-Ni reached 149 and 167 °C respectively under 30 min of irradiation, both showing photo-thermal efficiencies above 55%. In contrast, UiO-66 and ZIF-8 under the same experimental conditions were heated to 57 and 70 °C respectively, with photo-thermal efficiencies far below 10%. These results indicated that the observed photo-thermal effect was highly dependent on the optical absorption band of the MOFs. In other words, MOFs displaying absorption bands in the range of 300–600 nm featured a more pronounced photo-thermal effect than those that did not exhibit intense absorption in this region.

Some groups, however, have taken advantage of the heat insulating properties of the MOFs to create hybrid systems by combining them with active photo-thermal metallic species. In a work referenced in the previous section of this review (see Section 4.1.3), Zeng *et al.* designed a three-component heterostructure for the photocatalytic hydrogenation of CO<sub>2</sub> based on Pt nanocubes and Au nanocages encapsulated in a zeolitic imidazole framework (ZIF).<sup>36</sup> In this system, Pt nanocubes acted as catalytic sites facilitating CO<sub>2</sub> reduction, while plasmonic Au nanocages effectively transformed light into heat. Particularly, ZIF served as a thermal insulator preventing heat dissipation, thus confining heat inside the structure and enhancing the catalytic activity of Pt nanocubes. This approach offers obvious advantages, as photo-thermal catalysis benefits from low thermal conductivity to enhance local heat generation, and opens up the possibility for further advances by proper material engineering in hybrid architectures.

Although much less studied, other porous materials such as zeolites have also shown potential applications as active components in photo-thermal catalysis. Zhu *et al.* evaluated the photocatalytic activity for the visible light-mediated oxidation of benzyl alcohol and derivatives by Au NPs supported on various zeolites.<sup>298</sup> On the basis of experimental data, the authors proposed a reaction mechanism in which zeolitic supports effectively adsorbed reagent alcohols while Au NPs activated O<sub>2</sub> molecules by transferring plasmon-induced hot electrons. The resulting O<sub>2</sub><sup>–</sup> species then

triggered the oxidation reaction towards the corresponding aldehydes. From the kinetic study of the reaction, it was also possible to calculate the activation energy under both dark and visible light irradiation conditions. Surprisingly, it was found that the plasmon contribution decreased the activation barrier by 40%.

**5.2.3. Core-shell structures.** Core-shell structured hybrid materials present the advantages of tunable optical, electronic and thermal properties because of their bifunctional nature and exemplify a very attractive strategy to develop new materials exhibiting good photo-thermal performances.<sup>299</sup>

For instance, a recent work from Kumar *et al.* described the preparation of Au plasmonic NPs coated with TiO<sub>2</sub> and Pt NPs creating a core-shell structure (Pt@TiO<sub>2</sub>-AuNPs) and its application in CO<sub>2</sub> photoreduction.<sup>300</sup> It was demonstrated that the presence of a thin layer of TiO<sub>2</sub> increased the quantum yield of the reaction compared to bare Au NPs due to a combination of positive effects including an enlargement of the surface area, a red-shift broadening of the LSPR peak and an increase in the light absorption. Furthermore, the deposition of Pt NPs assisted the trapping of hot electrons from the CB of TiO<sub>2</sub>, thus preventing electron-hole recombination and enhancing the overall reaction rate. The successful combination of these three components rendered a high efficiency (up to 70%) and selectivity for the photocatalytic reduction of CO<sub>2</sub> to formic acid under NIR laser irradiation.

In a related precedent, Xiong and co-workers demonstrated that a precise tuning of the Pd shell thickness in the Au nanorods-Pd core-shell nanostructure could control not only the photo-thermal local heating but also the hot electron injection process in the photocatalytic styrene hydrogenation.<sup>301</sup> By virtue of ultrafast TAS measurements, the authors could obtain the time constants of electron-phonon scattering (attributable to the photo-thermal effect) and the charge recombination process (attributable to hot electron lifetime). Interestingly, the Pd shell thickness appeared to be a key parameter to determine the preeminent surface plasmon decay pathway, either local heating or hot electron transfer.

Lastly, Yang *et al.* developed AuCu-CuS core-shell heterostructures immobilized on TiO<sub>2</sub>.<sup>302</sup> This alloy-chalcogenide combination led to a synergistic effect between LSPR and photo-thermal heating that enhanced the photocatalytic glycerol oxidation at room temperature. Upon light irradiation, hot electrons from the plasmonic AuCu nanoalloy effectively transfer into the CB of TiO<sub>2</sub>, where they react with molecular oxygen to produce H<sub>2</sub>O<sub>2</sub> or reactive superoxide ions (O<sub>2</sub><sup>·</sup>) that drive the oxidation of glycerol. Meanwhile, the non-radiative recombination of electron-hole pairs in the plasmonic CuS shell induces thermal heating that provides an additional driving force for the overall catalytic process. In this work, the rational design of an effective bifunctional core-shell architecture proved to be a successful approach to perform photo-oxidation reactions under mild conditions.

## 6. Conclusions

In this review, we have succinctly explained the physics behind the LSPR phenomenon and its direct influence on the photo-thermal

effect displayed by plasmonic NPs. Upon light irradiation, non-radiative plasmon decay leads to the local heat generation and/or hot carrier formation in plasmonic structures. Individually or synergistically working together, thermal and photochemical contributions of the photo-thermal effect in plasmonic or non-plasmonic nanostructures can be exploited to promote catalytic chemical transformations. A comprehensive analysis of the contribution of both approaches to the overall reaction rate is, however, crucial to further understand the reaction mechanisms, so an extensive revision of straightforward methodologies to successfully distinguish the dominant reaction pathway has been included in this review. Indeed, photo-thermal catalysis has proven to be an efficient strategy to perform a wide variety of high energy-demanding catalytic processes including CO<sub>2</sub> hydrogenation, Fischer-Tropsch reaction, ammonia synthesis, methane activation or H<sub>2</sub> production. In order to improve catalytic performances and selectivity, a suitable material engineering is also advisable. In the last section of this review, we have included strategies for catalyst design in order to maximize the efficiency for a given photo-thermal-mediated pathway. The vast number of catalytic applications together with the possibility of tailoring the thermal, optical and electronic properties of photocatalysts demonstrates the huge versatility and potential of photo-thermal catalysis.

## 7. Outlook

Despite its afore-mentioned promising applications, the wide implementation of photo-thermal catalysis has still to compete with traditional thermal processes. In this regard, one of the main advantages of the photo-thermal approach is the possibility of increasing selectivity to target products by choosing adequate excitation wavelengths or light intensities. Hot-carrier chemistry has proved to be particularly efficient to trigger certain reaction mechanisms by selectively activating molecular bonds of adsorbed species, hence resulting in better selectivity than the thermochemical pathway. Photo-thermal catalysis also provides direct local heating at the nanoscale, thus concentrating heat at the surface of the active sites and avoiding unnecessary heating of the whole reactor system. Compared to traditional thermal catalysis, where fossil fuels are commonly used to achieve harsh reaction conditions, the photo-thermal approach renders a high productivity (in the order of mmol g<sub>cat</sub><sup>-1</sup>) under mild conditions employing sunlight as a unique energy source. This offers obvious benefits not only in terms of sustainability, but also in terms of recyclability and stability of catalysts. Furthermore, the photo-thermal effect allows a much faster cooling and this opens up the possibility to apply this strategy to discontinuous operation conditions.

Regardless of its outstanding performances presented here, there is still plenty of room for the improvement of catalytic efficiency in the field of photo-thermal catalysis. Taking into account that the photo-thermal effect derives from the combination of both photochemical and thermochemical pathways, different approaches can be followed depending on the dominating mechanism. For instance, when it comes to the photochemical process, photocatalysts with a strong light absorption, high



population of charge carriers, long charge carrier lifetimes and low radiative emission are advisable.<sup>30</sup> As we have previously commented in Section 5, size, morphology or surface defects clearly determine the efficiency of hot electron generation in plasmonic nanostructures. However, besides hot carrier generation efficiency, charge transfer efficiency is another key parameter that defines the photochemical contribution to the overall photo-thermal process. For example, in metal/semiconductor junctions only high-energy electrons are able to surpass the Schottky barrier at the metal/semiconductor interface, so the efficiency of the charge transfer will be dictated not only by the energy of hot carriers, but also the height of the Schottky barrier. Thus, a possible strategy to enhance charge transfer efficiency in this type of systems is to tailor the energy levels of the semiconductor by doping or formation of vacancies/defects in the structure. The spatial organization of the heterostructure or metal loading also determines the hot carrier transfer efficiency through the interface and, for that reason, the synergistic combination of all these parameters has to be pondered when designing photo-thermal nanostructures. In addition, it has been shown that small NPs shorten the pathway of hot electrons, hence favoring their injection to adsorbates. Last but not least, increasing the adsorption of intermediates is also a useful approach to extend their residence time on the surface of the active sites, which eventually enhances the hot electron transfer efficiency to these species. When it comes to the thermochemical process, materials with a high optical absorbance, low radiative emission and low thermal transport tend to amplify the temperature of the active sites, thus enhancing the catalytic performance. In this regard, non-stoichiometric black materials or materials with a high degree of defects in their structure are promising examples of photo-thermal catalysts. Overall, an adequate material engineering to improve both photochemical and thermochemical mechanisms can result in an extraordinary cooperation to boost photo-thermal performance in a vast number of chemical processes, as we have shown in this review.

Nevertheless, if photo-thermal catalysis aims to become a genuine source of solar fuels and chemicals it must overcome a series of limitations:

(a) In spite of the fact that photo-thermal catalysis does not require scaling considerations due to its close relationship with traditional catalytic heterogeneous systems, its main challenge is the incorporation of light into the processes. In contrast to other technologies, the photo-thermal approach does not demand extremely high temperatures but adequate photoreactors capable of harvesting and utilizing sunlight. Most of the works reported so far have been performed under batch conditions; however, the use of continuous flow systems appears to be more convenient for industrial-scale applications. For these reasons, research and development in reactor design and configuration are necessary for the wide implementation of this technique.

(b) Only a few studies among the vast number of works regarding photo-thermal catalysis have clarified the dominating pathway that rules the reaction mechanism. More fundamental studies in this direction are necessary.

(c) Although in the last few years a higher number of species apart from noble plasmonic metals (*i.e.* Au, Ag, Pd, Pt) have

been studied, it is still advisable to move towards abundant and cost-efficient materials. Chalcogenides, nitrides or pnictogenides have recently appeared as promising candidates for photo-thermal catalysis owing to their capability to transform light into heat.<sup>303–305</sup>

(d) The investigation of long-term stability of photocatalysts remains insufficient, as most of the stability tests are limited to several hours. For an adequate scale-up to commercial plants, feasible and stable photocatalysts are required.

(e) Particularly in the case of CO<sub>2</sub> conversion, CO and CH<sub>4</sub> are the major products in most of the reported works. However, it would be more convenient to obtain high-value-added C<sub>2+</sub> products, namely ethylene, propylene and other olefins. Selectivity control is one of the most attractive features of photo-thermal catalysis, so more work should be dedicated to the design of catalytic systems with remarkable selectivity. Furthermore, and especially in the case of carbon-based photocatalysts, experiments with isotopically labelled CO<sub>2</sub> are necessary in order to demonstrate that all carbon based products from these processes are actually derived from CO<sub>2</sub> and not from the support itself.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

Funding for this work was provided by King Abdullah University of Science and Technology (KAUST). The authors would like to thank Mrs Sandra Ramirez for her contribution to the graphical abstract and to the rest of the artwork in this manuscript.

## References

- 1 R. Schlögl, *ChemSusChem*, 2010, **3**, 209–222.
- 2 Y. Zhao, W. Gao, S. Li, G. R. Williams, A. H. Mahadi and D. Ma, *Joule*, 2019, **3**, 920–937.
- 3 A. Fujishima and K. Honda, *Nature*, 1972, **238**, 37–38.
- 4 T. Paik, M. Cargnello, T. R. Gordon, S. Zhang, H. Yun, J. D. Lee, H. Y. Woo, S. J. Oh, C. R. Kagan and P. Fornasiero, *ACS Energy Lett.*, 2018, **3**, 1904–1910.
- 5 H. Liu, P. Wu, H. Li, Z. Chen, L. Wang, X. Zeng, Y. Zhu, Y. Jiang, X. Liao and B. S. Haynes, *Appl. Catal., B*, 2019, **259**, 118026.
- 6 C. Xiao, H. Wang, L. Zhang, S. Sun and W. Wang, *Chem-CatChem*, 2019, **11**, 6467–6472.
- 7 D. Li, S. Ouyang, H. Xu, D. Lu, M. Zhao, X. Zhang and J. Ye, *Chem. Commun.*, 2016, **52**, 5989–5992.
- 8 Z.-C. Kong, J.-F. Liao, Y.-J. Dong, Y.-F. Xu, H.-Y. Chen, D.-B. Kuang and C.-Y. Su, *ACS Energy Lett.*, 2018, **3**, 2656–2662.
- 9 A. Manjavacas, J. G. Liu, V. Kulkarni and P. Nordlander, *ACS Nano*, 2014, **8**, 7630–7638.

10 B. Y. Zheng, H. Zhao, A. Manjavacas, M. McClain, P. Nordlander and N. J. Halas, *Nat. Commun.*, 2015, **6**, 1–7.

11 H. H. Richardson, Z. N. Hickman, A. O. Govorov, A. C. Thomas, W. Zhang and M. E. Kordesch, *Nano Lett.*, 2006, **6**, 783–788.

12 H. H. Richardson, M. T. Carlson, P. J. Tandler, P. Hernandez and A. O. Govorov, *Nano Lett.*, 2009, **9**, 1139–1146.

13 G. Baffou and R. Quidant, *Laser Photonics Rev.*, 2013, **7**, 171–187.

14 X. Huang, I. H. El-Sayed, W. Qian and M. A. El-Sayed, *J. Am. Chem. Soc.*, 2006, **128**, 2115–2120.

15 W. Tao, X. Ji, X. Xu, M. A. Islam, Z. Li, S. Chen, P. E. Saw, H. Zhang, Z. Bharwani and Z. Guo, *Angew. Chem., Int. Ed.*, 2017, **56**, 11896–11900.

16 S. Lu, F. Liu, P. Qiu, M. Qiao, Y. Li, Z. Cheng, N. Xue, X. Hou, C. Xu and Y. Xiang, *Chem. Eng. J.*, 2020, **379**, 122382.

17 K.-K. Liu, Q. Jiang, S. Tadepalli, R. Raliya, P. Biswas, R. R. Naik and S. Singamaneni, *ACS Appl. Mater. Interfaces*, 2017, **9**, 7675–7681.

18 A. Politano, P. Argurio, G. Di Profio, V. Sanna, A. Cupolillo, S. Chakraborty, H. A. Arafat and E. Curcio, *Adv. Mater.*, 2017, **29**, 1603504.

19 Y. Shao, Z. Jiang, Y. Zhang, T. Wang, P. Zhao, Z. Zhang, J. Yuan and H. Wang, *ACS Nano*, 2018, **12**, 11704–11710.

20 S. Linic, P. Christopher and D. B. Ingram, *Nat. Mater.*, 2011, **10**, 911–921.

21 J. Lee, S. Mubeen, X. Ji, G. D. Stucky and M. Moskovits, *Nano Lett.*, 2012, **12**, 5014–5019.

22 M. Xiao, R. Jiang, F. Wang, C. Fang, J. Wang and C. Y. Jimmy, *J. Mater. Chem. A*, 2013, **1**, 5790–5805.

23 C. Wang and D. Astruc, *Chem. Soc. Rev.*, 2014, **43**, 7188–7216.

24 X. Meng, L. Liu, S. Ouyang, H. Xu, D. Wang, N. Zhao and J. Ye, *Adv. Mater.*, 2016, **28**, 6781–6803.

25 S. Li, P. Miao, Y. Zhang, J. Wu, B. Zhang, Y. Du, X. Han, J. Sun and P. Xu, *Adv. Mater.*, 2020, 2000086.

26 A. A. Jelle, K. K. Ghuman, P. G. O'Brien, M. Hmadedh, A. Sandhel, D. D. Perovic, C. V. Singh, C. A. Mims and G. A. Ozin, *Adv. Energy Mater.*, 2018, **8**, 1702277.

27 S. Tang, J. Sun, H. Hong and Q. Liu, *Front. Energy*, 2017, **11**, 437–451.

28 L. Zhu, M. Gao, C. K. N. Peh and G. W. Ho, *Mater. Horiz.*, 2018, **5**, 323–343.

29 J.-D. Xiao and H.-L. Jiang, *Acc. Chem. Res.*, 2018, **52**, 356–366.

30 M. Ghoussoub, M. Xia, P. N. Duchesne, D. Segal and G. Ozin, *Energy Environ. Sci.*, 2019, **12**, 1122–1142.

31 Z. j. Wang, H. Song, H. Liu and J. Ye, *Angew. Chem., Int. Ed.*, 2020, **59**, 8016–8035.

32 P. G. O'Brien, A. Sandhel, T. E. Wood, A. A. Jelle, L. B. Hoch, D. D. Perovic, C. A. Mims and G. A. Ozin, *Adv. Sci.*, 2014, **1**, 1400001.

33 S. I. Nikitenko, T. Chave, C. Cau, H.-P. Brau and V. R. Flaud, *ACS Catal.*, 2015, **5**, 4790–4795.

34 J. Jia, P. G. O'Brien, L. He, Q. Qiao, T. Fei, L. M. Reyes, T. E. Burrow, Y. Dong, K. Liao and M. Varela, *Adv. Sci.*, 2016, **3**, 1600189.

35 F. Wang, Y. Huang, Z. Chai, M. Zeng, Q. Li, Y. Wang and D. Xu, *Chem. Sci.*, 2016, **7**, 6887–6893.

36 W. Zhang, L. Wang, K. Wang, M. U. Khan, M. Wang, H. Li and J. Zeng, *Small*, 2017, **13**, 1602583.

37 R. Zhang, J. Xu, C. Lu and Z. Xu, *Mater. Lett.*, 2018, **228**, 199–202.

38 P. Drude, *Ann. Phys.*, 1900, **306**, 566–613.

39 G. Mie, *Ann. Phys.*, 1908, **330**, 377–445.

40 M. L. Brongersma, N. J. Halas and P. Nordlander, *Nat. Nanotechnol.*, 2015, **10**, 25.

41 A. Gellé and A. Moores, *Curr. Opin. Green Sustainable Chem.*, 2019, **15**, 60–66.

42 L. Hu, S. Gao, X. Ding, D. Wang, J. Jiang, J. Jin and L. Jiang, *ACS Nano*, 2015, **9**, 4835–4842.

43 J. Qiu and W. D. Wei, *J. Phys. Chem. C*, 2014, **118**, 20735–20749.

44 L. V. Besteiro, X.-T. Kong, Z. Wang, G. Hartland and A. O. Govorov, *ACS Photonics*, 2017, **4**, 2759–2781.

45 A. Gellé, T. Jin, L. de la Garza, G. D. Price, L. V. Besteiro and A. Moores, *Chem. Rev.*, 2019, **120**, 986–1041.

46 L. Wang, Y. Dong, T. Yan, Z. Hu, A. A. Jelle, D. M. Meira, P. N. Duchesne, J. Y. Y. Loh, C. Qiu and E. E. Storey, *Nat. Commun.*, 2020, **11**, 1–8.

47 Y. Zhao, A. Dunn, J. Lin and D. Shi, *Novel Nanomaterials for Biomedical, Environmental and Energy Applications*, Elsevier, 2019, pp. 415–434.

48 M. J. Kale, T. Avanesian and P. Christopher, *ACS Catal.*, 2014, **4**, 116–128.

49 P. Christopher, H. Xin and S. Linic, *Nat. Chem.*, 2011, **3**, 467.

50 Y. Zhang, S. He, W. Guo, Y. Hu, J. Huang, J. R. Mulcahy and W. D. Wei, *Chem. Rev.*, 2017, **118**, 2927–2954.

51 G. V. Hartland, L. V. Besteiro, P. Johns and A. O. Govorov, *ACS Energy Lett.*, 2017, **2**, 1641–1653.

52 P. Christopher and M. Moskovits, *Annu. Rev. Phys. Chem.*, 2017, **68**, 379–398.

53 C. Boerigter, R. Campana, M. Morabito and S. Linic, *Nat. Commun.*, 2016, **7**, 1–9.

54 S. Sarina, E. Jaatinen, Q. Xiao, Y. M. Huang, P. Christopher, J. C. Zhao and H. Y. Zhu, *J. Phys. Chem. Lett.*, 2017, **8**, 2526–2534.

55 S. Mukherjee, F. Libisch, N. Large, O. Neumann, L. V. Brown, J. Cheng, J. B. Lassiter, E. A. Carter, P. Nordlander and N. J. Halas, *Nano Lett.*, 2013, **13**, 240–247.

56 S. Mukherjee, L. Zhou, A. M. Goodman, N. Large, C. Ayala-Orozco, Y. Zhang, P. Nordlander and N. J. Halas, *J. Am. Chem. Soc.*, 2014, **136**, 64–67.

57 L. Zhou, C. Zhang, M. J. McClain, A. Manjavacas, C. M. Krauter, S. Tian, F. Berg, H. O. Everitt, E. A. Carter and P. Nordlander, *Nano Lett.*, 2016, **16**, 1478–1484.

58 A. Marimuthu, J. Zhang and S. Linic, *Science*, 2013, **339**, 1590–1593.

59 C. Bauer, J.-P. Abid, D. Fermin and H. H. Girault, *J. Chem. Phys.*, 2004, **120**, 9302–9315.

60 C. Boerigter, U. Aslam and S. Linic, *ACS Nano*, 2016, **10**, 6108–6115.

61 B. Foerster, A. Joplin, K. Kaefer, S. Celiksoy, S. Link and C. Sönnichsen, *ACS Nano*, 2017, **11**, 2886–2893.

62 M. Wang, M. Ye, J. Iocozzia, C. Lin and Z. Lin, *Adv. Sci.*, 2016, **3**, 1600024.

63 H. G. Baldoví, S. T. Neațu, A. Khan, A. M. Asiri, S. A. Kosa and H. Garcia, *J. Phys. Chem. C*, 2015, **119**, 6819–6827.

64 J. S. DuChene, B. C. Sweeny, A. C. Johnston-Peck, D. Su, E. A. Stach and W. D. Wei, *Angew. Chem., Int. Ed.*, 2014, **53**, 7887–7891.

65 M. R. Khan, T. W. Chuan, A. Yousuf, M. Chowdhury and C. K. Cheng, *Catal. Sci. Technol.*, 2015, **5**, 2522–2531.

66 K. Marchuk and K. A. Willets, *Chem. Phys.*, 2014, **445**, 95–104.

67 P. J. Schuck, *Nat. Nanotechnol.*, 2013, **8**, 799.

68 Y. Tian and T. Tatsuma, *J. Am. Chem. Soc.*, 2005, **127**, 7632–7637.

69 D. B. Ingram and S. Linic, *J. Am. Chem. Soc.*, 2011, **133**, 5202–5205.

70 C. Gomes Silva, R. Juárez, T. Marino, R. Molinari and H. García, *J. Am. Chem. Soc.*, 2011, **133**, 595–602.

71 Z. Zhang, L. Zhang, M. N. Hedhili, H. Zhang and P. Wang, *Nano Lett.*, 2013, **13**, 14–20.

72 Y. Zhu, Z. Xu, W. Jiang, W. Yin, S. Zhong, P. Gong, R. Qiao, Z. Li and S. Bai, *RSC Adv.*, 2016, **6**, 56800–56806.

73 M.-Z. Ge, C.-Y. Cao, S.-H. Li, Y.-X. Tang, L.-N. Wang, N. Qi, J.-Y. Huang, K.-Q. Zhang, S. Al-Deyab and Y.-K. Lai, *Nanoscale*, 2016, **8**, 5226–5234.

74 W. Hou, W. H. Hung, P. Pavaskar, A. Goeppert, M. Aykol and S. B. Cronin, *ACS Catal.*, 2011, **1**, 929–936.

75 S. Bera, J. E. Lee, S. B. Rawal and W. I. Lee, *Appl. Catal., B*, 2016, **199**, 55–63.

76 X. Cheng, P. Dong, Z. Huang, Y. Zhang, Y. Chen, X. Nie and X. Zhang, *J. CO<sub>2</sub> Util.*, 2017, **20**, 200–207.

77 J. Y. Do, R. K. Chava, K. K. Mandari, N.-K. Park, H.-J. Ryu, M. W. Seo, D. Lee, T. Senthil and M. Kang, *Appl. Catal., B*, 2018, **237**, 895–910.

78 P. Christopher, D. B. Ingram and S. Linic, *J. Phys. Chem. C*, 2010, **114**, 9173–9177.

79 X. Chen, Z. Zheng, X. Ke, E. Jaatinen, T. Xie, D. Wang, C. Guo, J. Zhao and H. Zhu, *Green Chem.*, 2010, **12**, 414–419.

80 H. Zhang, X. Fan, X. Quan, S. Chen and H. Yu, *Environ. Sci. Technol.*, 2011, **45**, 5731–5736.

81 M. Misra, N. Singh and R. K. Gupta, *Catal. Sci. Technol.*, 2017, **7**, 570–580.

82 A. Furube, L. Du, K. Hara, R. Katoh and M. Tachiya, *J. Am. Chem. Soc.*, 2007, **129**, 14852–14853.

83 T. P. White and K. R. Catchpole, *Appl. Phys. Lett.*, 2012, **101**, 073905.

84 A. Naldoni, T. Montini, F. Malara, M. M. Mróz, A. Beltram, T. Virgili, C. L. Boldrini, M. Marelli, I. Romero-Ocaña and J. J. Delgado, *ACS Catal.*, 2017, **7**, 1270–1278.

85 D. C. Ratchford, A. D. Dunkelberger, I. Vurgaftman, J. C. Owrutsky and P. E. Pehrsson, *Nano Lett.*, 2017, **17**, 6047–6055.

86 S. Tan, A. Argondizzo, J. Ren, L. Liu, J. Zhao and H. Petek, *Nat. Photonics*, 2017, **11**, 806–812.

87 Y. Yu, Z. Ji, S. Zu, B. Du, Y. Kang, Z. Li, Z. Zhou, K. Shi and Z. Fang, *Adv. Funct. Mater.*, 2016, **26**, 6394–6401.

88 J.-Y. Bigot, V. Halté, J.-C. Merle and A. Daunois, *Chem. Phys.*, 2000, **251**, 181–203.

89 K. Wu, J. Chen, J. R. McBride and T. Lian, *Science*, 2015, **349**, 632–635.

90 A. Hoggard, L.-Y. Wang, L. Ma, Y. Fang, G. You, J. Olson, Z. Liu, W.-S. Chang, P. M. Ajayan and S. Link, *ACS Nano*, 2013, **7**, 11209–11217.

91 K. Wu, W. E. Rodríguez-Córdoba, Y. Yang and T. Lian, *Nano Lett.*, 2013, **13**, 5255–5263.

92 K. Watanabe, D. Menzel, N. Nilius and H.-J. Freund, *Chem. Rev.*, 2006, **106**, 4301–4320.

93 S. Yu, A. J. Wilson, G. Kumari, X. Zhang and P. K. Jain, *ACS Energy Lett.*, 2017, **2**, 2058–2070.

94 A. O. Govorov and H. H. Richardson, *Nano Today*, 2007, **2**, 30–38.

95 A. O. Govorov, W. Zhang, T. Skeini, H. Richardson, J. Lee and N. A. Kotov, *Nanoscale Res. Lett.*, 2006, **1**, 84.

96 H. Fogler, *Essentials of Chemical Reaction Engineering*, 2011.

97 G. Baffou, P. Berto, E. Bermúdez Ureña, R. Quidant, S. Monneret, J. Polleux and H. Rigneault, *ACS Nano*, 2013, **7**, 6478–6488.

98 J. R. Adleman, D. A. Boyd, D. G. Goodwin and D. Psaltis, *Nano Lett.*, 2009, **9**, 4417–4423.

99 Q. Yang, Q. Xu, S. H. Yu and H. L. Jiang, *Angew. Chem., Int. Ed.*, 2016, **55**, 3685–3689.

100 Y. Sivan, J. H. Baraban and Y. Dubi, *OSA Continuum*, 2020, **3**, 483–497.

101 S. Sarina, H. Y. Zhu, Q. Xiao, E. Jaatinen, J. Jia, Y. Huang, Z. Zheng and H. Wu, *Angew. Chem.*, 2014, **126**, 2979–2984.

102 Y. Qi, L. Song, S. Ouyang, X. Liang, S. Ning, Q. Zhang and J. Ye, *Adv. Mater.*, 2019, 1903915.

103 C. Wang, O. Ranasingha, S. Natesakhawat, P. R. 2Ohodnicki, M. Andio, J. P. Lewis and C. Matranga, *Nanoscale*, 2013, **5**, 6968–6974.

104 K. H. Kim, K. Watanabe, D. Mulugeta, H.-J. Freund and D. Menzel, *Phys. Rev. Lett.*, 2011, **107**, 047401.

105 S. Navalon, M. de Miguel, R. Martin, M. Alvaro and H. Garcia, *J. Am. Chem. Soc.*, 2011, **133**, 2218–2226.

106 M. J. Landry, A. Gellé, B. Y. Meng, C. J. Barrett and A. Moores, *ACS Catal.*, 2017, **7**, 6128–6133.

107 P. Christopher, H. Xin, A. Marimuthu and S. Linic, *Nat. Mater.*, 2012, **11**, 1044–1050.

108 X. Zhang, X. Li, D. Zhang, N. Q. Su, W. Yang, H. O. Everitt and J. Liu, *Nat. Commun.*, 2017, **8**, 1–9.

109 H. Song, X. Meng, Z.-J. Wang, Z. Wang, H. Chen, Y. Weng, F. Ichihara, M. Oshikiri, T. Kako and J. Ye, *ACS Catal.*, 2018, **8**, 7556–7565.

110 X. Meng, T. Wang, L. Liu, S. Ouyang, P. Li, H. Hu, T. Kako, H. Iwai, A. Tanaka and J. Ye, *Angew. Chem., Int. Ed.*, 2014, **53**, 11478–11482.

111 P. Li, G. Xiao, Y. Zhao and H. Su, *ACS Catal.*, 2020, **10**, 3640–3649.

112 J. Guo, Y. Zhang, L. Shi, Y. Zhu, M. F. Mideksa, K. Hou, W. Zhao, D. Wang, M. Zhao and X. Zhang, *J. Am. Chem. Soc.*, 2017, **139**, 17964–17972.

113 G. Baffou, I. Bordacchini, A. Baldi and R. Quidant, *Light: Sci. Appl.*, 2020, **9**, 1–16.

114 Y. Kim, D. Dumett Torres and P. K. Jain, *Nano Lett.*, 2016, **16**, 3399–3407.

115 L. Zhou, D. F. Swearer, C. Zhang, H. Robatjazi, H. Zhao, L. Henderson, L. Dong, P. Christopher, E. A. Carter and P. Nordlander, *Science*, 2018, **362**, 69–72.

116 H. Song, X. Meng, T. D. Dao, W. Zhou, H. Liu, L. Shi, H. Zhang, T. Nagao, T. Kako and J. Ye, *ACS Appl. Mater. Interfaces*, 2018, **10**, 408–416.

117 U. Aslam, V. G. Rao, S. Chavez and S. Linic, *Nat. Catal.*, 2018, **1**, 656–665.

118 M. J. Kale, T. Avanesian, H. Xin, J. Yan and P. Christopher, *Nano Lett.*, 2014, **14**, 5405–5412.

119 F. Sastre, A. V. Puga, L. Liu, A. Corma and H. Garcia, *J. Am. Chem. Soc.*, 2014, **136**, 6798–6801.

120 D. Mateo, J. Albero and H. García, *Appl. Catal., B*, 2018, **224**, 563–571.

121 D. G. Cahill, P. V. Braun, G. Chen, D. R. Clarke, S. Fan, K. E. Goodson, P. Keblinski, W. P. King, G. D. Mahan and A. Majumdar, *Appl. Phys. Rev.*, 2014, **1**, 011305.

122 F. Menges, P. Mensch, H. Schmid, H. Riel, A. Stemmer and B. Gotsmann, *Nat. Commun.*, 2016, **7**, 1–6.

123 M. Mecklenburg, W. A. Hubbard, E. White, R. Dhall, S. B. Cronin, S. Aloni and B. Regan, *Science*, 2015, **347**, 629–632.

124 J. C. Idrobo, A. R. Lupini, T. Feng, R. R. Unocic, F. S. Walden, D. S. Gardiner, T. C. Lovejoy, N. Dellby, S. T. Pantelides and O. L. Krivanek, *Phys. Rev. Lett.*, 2018, **120**, 095901.

125 L. Zhou, D. F. Swearer, H. Robatjazi, A. Alabastri, P. Christopher, E. A. Carter, P. Nordlander and N. J. Halas, *Science*, 2019, **364**, eaaw9545.

126 Y. Sivan, J. Baraban, I. W. Un and Y. Dubi, *Science*, 2019, **364**, eaaw9367.

127 C. D. Brites, P. P. Lima, N. J. Silva, A. Millán, V. S. Amaral, F. Palacio and L. D. Carlos, *Nanoscale*, 2012, **4**, 4799–4829.

128 D. Wang, Y. R. Koh, Z. A. Kudyshev, K. Maize, A. V. Kildishev, A. Boltasseva, V. M. Shalaev and A. Shakouri, *Nano Lett.*, 2019, **19**, 3796–3803.

129 M. V. Balois, N. Hayazawa, F. C. Catalan, S. Kawata, T.-A. Yano and T. Hayashi, *Anal. Bioanal. Chem.*, 2015, **407**, 8205–8213.

130 S. Xie, E. Iglesia and A. T. Bell, *J. Phys. Chem. B*, 2001, **105**, 5144–5152.

131 W. Sun, G. Zhong, C. Kübel, A. A. Jelle, C. Qian, L. Wang, M. Ebrahimi, L. M. Reyes, A. S. Helmy and G. A. Ozin, *Angew. Chem., Int. Ed.*, 2017, **56**, 6329–6334.

132 J. Jia, H. Wang, Z. Lu, P. G. O'Brien, M. Ghoussoub, P. Duchesne, Z. Zheng, P. Li, Q. Qiao and L. Wang, *Adv. Sci.*, 2017, **4**, 1700252.

133 G. W. Walker, V. C. Sundar, C. M. Rudzinski, A. W. Wun, M. G. Bawendi and D. G. Nocera, *Appl. Phys. Lett.*, 2003, **83**, 3555–3557.

134 S. Li, K. Zhang, J.-M. Yang, L. Lin and H. Yang, *Nano Lett.*, 2007, **7**, 3102–3105.

135 U. Resch-Genger, M. Grabolle, S. Cavaliere-Jaricot, R. Nitschke and T. Nann, *Nat. Methods*, 2008, **5**, 763.

136 P. Löw, B. Kim, N. Takama and C. Bergaud, *Small*, 2008, **4**, 908–914.

137 D. Mateo, J. Albero and H. García, *Energy Environ. Sci.*, 2017, **10**, 2392–2400.

138 Y. Yue and X. Wang, *Nano Rev.*, 2012, **3**, 11586.

139 V. A. Spata and E. A. Carter, *ACS Nano*, 2018, **12**, 3512–3522.

140 Y. Yu, V. Sundaresan and K. A. Willets, *J. Phys. Chem. C*, 2018, **122**, 5040–5048.

141 X. Zhang, X. Li, M. E. Reish, D. Zhang, N. Q. Su, Y. Gutiérrez, F. Moreno, W. Yang, H. O. Everitt and J. Liu, *Nano Lett.*, 2018, **18**, 1714–1723.

142 X. Li, H. O. Everitt and J. Liu, *Nano Res.*, 2019, **12**, 1906–1911.

143 C. Zhan, B.-W. Liu, Y.-F. Huang, S. Hu, B. Ren, M. Moskovits and Z.-Q. Tian, *Nat. Commun.*, 2019, **10**, 1–8.

144 X. Li, H. O. Everitt and J. Liu, *Nano Res.*, 2020, 1–13.

145 M. Najafpour, *Artificial photosynthesis*, BoD-Books on Demand, 2012.

146 B. AlOtaibi, S. Fan, D. Wang, J. Ye and Z. Mi, *ACS Catal.*, 2015, **5**, 5342–5348.

147 I. Shown, H.-C. Hsu, Y.-C. Chang, C.-H. Lin, P. K. Roy, A. Ganguly, C.-H. Wang, J.-K. Chang, C.-I. Wu and L.-C. Chen, *Nano Lett.*, 2014, **14**, 6097–6103.

148 A. Nawaz, A. Kuila, A. Rani, N. S. Mishra, L. C. Sim, K. H. Leong and P. Saravanan, *Industrial Applications of Nanomaterials*, Elsevier, 2019, pp. 151–179.

149 I. McConnell, G. Li and G. W. Brudvig, *Chem. Biol.*, 2010, **17**, 434–447.

150 R. Poudyal, I. Tiwari, A. Koirala, H. Masukawa, K. Inoue, T. Tomo, M. Najafpour, S. Allakhverdiev and T. Veziroğlu, *Compendium of Hydrogen Energy*, Elsevier, 2015, pp. 289–317.

151 J. H. Lee, J. Y. Do, N.-K. Park, M. W. Seo, H.-J. Ryu, J.-P. Hong, Y. S. Kim, S. K. Kim and M. Kang, *J. Photochem. Photobiol., A*, 2018, **364**, 219–232.

152 Y. Li, C. Wang, M. Song, D. Li, X. Zhang and Y. Liu, *Appl. Catal., B*, 2019, **243**, 760–770.

153 T. Inoue, A. Fujishima, S. Konishi and K. Honda, *Nature*, 1979, **277**, 637–638.

154 J. Low, B. Cheng and J. Yu, *Appl. Surf. Sci.*, 2017, **392**, 658–686.

155 G. Yin, X. Huang, T. Chen, W. Zhao, Q. Bi, J. Xu, Y. Han and F. Huang, *ACS Catal.*, 2018, **8**, 1009–1017.

156 F. Saladin and I. Alxneit, *J. Chem. Soc., Faraday Trans.*, 1997, **93**, 4159–4163.

157 K. Schwartzenberg, J. Hamilton, A. K. Lucid, E. Weitz, J. Notestein, M. Nolan, J. A. Byrne and K. A. Gray, *Catal. Today*, 2017, **280**, 65–73.

158 L. Wang, X. Liu, Y. Dang, H. Xie, Q. Zhao and L. Ye, *Solid State Sci.*, 2019, **89**, 67–73.

159 L. Zhang, G. Kong, Y. Meng, J. Tian, L. Zhang, S. Wan, J. Lin and Y. Wang, *ChemSusChem*, 2017, **10**, 4709–4714.

160 M. Sun, W. Hao, C. Wang and T. Wang, *Chem. Phys. Lett.*, 2007, **443**, 342–346.

161 A. Tang, Y. Jia, S. Zhang, Q. Yu and X. Zhang, *Catal. Commun.*, 2014, **50**, 1–4.

162 L. J. Sherry, R. Jin, C. A. Mirkin, G. C. Schatz and R. P. Van Duyne, *Nano Lett.*, 2006, **6**, 2060–2065.

163 Z. Zhang, S. Wang, T. Yu and T. Wu, *J. Phys. Chem. C*, 2007, **111**, 17500–17505.

164 D. Kumar, C. H. Park and C. S. Kim, *J. Mater. Chem. A*, 2020, **8**, 5734–5743.

165 H. Wang, Y. Wang, L. Guo, X. Zhang, C. Ribeiro and T. He, *Chin. J. Catal.*, 2020, **41**, 131–139.

166 F. Yu, C. Wang, H. Ma, M. Song, D. Li, Y. Li, S. Li, X. Zhang and Y. Liu, *Nanoscale*, 2020, **12**, 7000–7010.

167 Q. Guo, Q. Zhang, H. Wang, Z. Liu and Z. Zhao, *Catal. Commun.*, 2016, **77**, 118–122.

168 J. Li, Y. Ye, L. Ye, F. Su, Z. Ma, J. Huang, H. Xie, D. E. Doronkin, A. Zimina and J.-D. Grunwaldt, *J. Mater. Chem. A*, 2019, **7**, 2821–2830.

169 L. Wang, Y. Wang, Y. Cheng, Z. Liu, Q. Guo, M. N. Ha and Z. Zhao, *J. Mater. Chem. A*, 2016, **4**, 5314–5322.

170 L. Wang, M. N. Ha, Z. Liu and Z. Zhao, *Integr. Ferroelectr.*, 2016, **172**, 97–108.

171 Y. Bai, P. Yang, P. Wang, H. Xie, H. Dang and L. Ye, *J. CO2 Util.*, 2018, **23**, 51–60.

172 K. Wang, R. Jiang, T. Peng, X. Chen, W. Dai and X. Fu, *Appl. Catal., B*, 2019, **256**, 117780.

173 M. Xu, X. Hu, S. Wang, J. Yu, D. Zhu and J. Wang, *J. Catal.*, 2019, **377**, 652–661.

174 J. Low, L. Zhang, B. Zhu, Z. Liu and J. Yu, *ACS Sustainable Chem. Eng.*, 2018, **6**, 15653–15661.

175 G. Zhang, G. Liu, L. Wang and J. T. Irvine, *Chem. Soc. Rev.*, 2016, **45**, 5951–5984.

176 D. W. Hwang, H. G. Kim, J. Kim, K. Y. Cha, Y. G. Kim and J. S. Lee, *J. Catal.*, 2000, **193**, 40–48.

177 H. Fujito, H. Kunioku, D. Kato, H. Suzuki, M. Higashi, H. Kageyama and R. Abe, *J. Am. Chem. Soc.*, 2016, **138**, 2082–2085.

178 X. Sun, Y. Mi, F. Jiao and X. Xu, *ACS Catal.*, 2018, **8**, 3209–3221.

179 L. Xu, M. N. Ha, Q. Guo, L. Wang, Y. Ren, N. Sha and Z. Zhao, *RSC Adv.*, 2017, **7**, 45949–45959.

180 G. Wei, D. Zheng, L. Xu, Q. Guo, J. Hu, N. Sha and Z. Zhao, *Mater. Res. Express*, 2019, **6**, 086221.

181 M. N. Ha, G. Lu, Z. Liu, L. Wang and Z. Zhao, *J. Mater. Chem. A*, 2016, **4**, 13155–13165.

182 U. Ulmer, T. Dingle, P. N. Duchesne, R. H. Morris, A. Tavasoli, T. Wood and G. A. Ozin, *Nat. Commun.*, 2019, **10**, 1–12.

183 L. Pastor-Pérez, F. Baibars, E. Le Sache, H. Arellano-García, S. Gu and T. R. Reina, *J. CO2 Util.*, 2017, **21**, 423–428.

184 X. Yang, X. Su, X. Chen, H. Duan, B. Liang, Q. Liu, X. Liu, Y. Ren, Y. Huang and T. Zhang, *Appl. Catal., B*, 2017, **216**, 95–105.

185 C. Xu, W. Huang, Z. Li, B. Deng, Y. Zhang, M. Ni and K. Cen, *ACS Catal.*, 2018, **8**, 6582–6593.

186 W. Sun, C. Qian, L. He, K. K. Ghuman, A. P. Wong, J. Jia, A. A. Jelle, P. G. O'Brien, L. M. Reyes and T. E. Wood, *Nat. Commun.*, 2016, **7**, 1–9.

187 R. Costi, A. E. Saunders, E. Elmalem, A. Salant and U. Banin, *Nano Lett.*, 2008, **8**, 637–641.

188 R. Kaur and B. Pal, *J. Mol. Catal. A: Chem.*, 2012, **355**, 39–43.

189 I. Tanabe, T. Ryoki and Y. Ozaki, *RSC Adv.*, 2015, **5**, 13648–13652.

190 A. A. Upadhye, I. Ro, X. Zeng, H. J. Kim, I. Tejedor, M. A. Anderson, J. A. Dumesic and G. W. Huber, *Catal. Sci. Technol.*, 2015, **5**, 2590–2601.

191 B. Lu, F. Quan, Z. Sun, F. Jia and L. Zhang, *Catal. Commun.*, 2019, **129**, 105724.

192 Z. Zhao, D. E. Doronkin, Y. Ye, J.-D. Grunwaldt, Z. Huang and Y. Zhou, *Chin. J. Catal.*, 2020, **41**, 286–293.

193 N. Kong, B. Han, Z. Li, Y. Fang, K. Feng, Z. Wu, S. Wang, A.-B. Xu, Y. Yu, C. Li, Z. Lin and L. He, *ACS Appl. Nano Mater.*, 2020, **3**, 3028–3033.

194 S. Ning, H. Xu, Y. Qi, L. Song, Q. Zhang, S. Ouyang and J. Ye, *ACS Catal.*, 2020, **10**, 4726–4736.

195 H. Zhang, T. Wang, J. Wang, H. Liu, T. D. Dao, M. Li, G. Liu, X. Meng, K. Chang, L. Shi, T. Nagao and J. Ye, *Adv. Mater.*, 2016, **28**, 3703–3710.

196 L. Wang, Y. Dong, T. Yan, Z. Hu, A. A. Jelle, D. M. Meira, P. N. Duchesne, J. Y. Y. Loh, C. Qiu, E. E. Storey, Y. Xu, W. Sun, M. Ghoussoub, N. P. Kherani, A. S. Helmy and G. A. Ozin, *Nat. Commun.*, 2020, **11**, 2432.

197 L. B. Hoch, P. G. O'Brien, A. Jelle, A. Sandhel, D. D. Perovic, C. A. Mims and G. A. Ozin, *ACS Nano*, 2016, **10**, 9017–9025.

198 G. A. Olah, *Angew. Chem., Int. Ed.*, 2005, **44**, 2636–2639.

199 J. F. Haw and D. M. Marcus, *Top. Catal.*, 2005, **34**, 41–48.

200 F. Arena, K. Barbera, G. Italiano, G. Bonura, L. Spadaro and F. Frusteri, *J. Catal.*, 2007, **249**, 185–194.

201 D. Wu, K. Deng, B. Hu, Q. Lu, G. Liu and X. Hong, *ChemCatChem*, 2019, **11**, 1598–1601.

202 E. T. Kho, T. H. Tan, E. Lovell, R. J. Wong, J. Scott and R. Amal, *Green Energy Environ.*, 2017, **2**, 204–217.

203 P. G. O'Brien, K. K. Ghuman, A. A. Jelle, A. Sandhel, T. E. Wood, J. Y. Y. Loh, J. Jia, D. Perovic, C. V. Singh, N. P. Kherani, C. A. Mims and G. A. Ozin, *Energy Environ. Sci.*, 2018, **11**, 3443–3451.

204 S. Tada, T. Shimizu, H. Kameyama, T. Haneda and R. Kikuchi, *Int. J. Hydrogen Energy*, 2012, **37**, 5527–5531.

205 E. T. Kho, S. Jantarang, Z. Zheng, J. Scott and R. Amal, *Engineering*, 2017, **3**, 393–401.

206 S. Jantarang, E. C. Lovell, T. H. Tan, J. Scott and R. Amal, *Prog. Nat. Sci.: Mater. Int.*, 2018, **28**, 168–177.

207 Y. Li, J. Hao, H. Song, F. Zhang, X. Bai, X. Meng, H. Zhang, S. Wang, Y. Hu and J. Ye, *Nat. Commun.*, 2019, **10**, 1–9.

208 K. R. Thampi, J. Kiwi and M. Graetzel, *Nature*, 1987, **327**, 506–508.

209 K. K. Bando, H. Arakawa and N. Ichikuni, *Catal. Lett.*, 1999, **60**, 125–132.

210 J. H. Kwak, L. Kovarik and J. Szanyi, *ACS Catal.*, 2013, **3**, 2449–2455.

211 C. Wang, S. Fang, S. Xie, Y. Zheng and Y. H. Hu, *J. Mater. Chem. A*, 2020, **8**, 7390–7394.

212 L. Lin, K. Wang, K. Yang, X. Chen, X. Fu and W. Dai, *Appl. Catal., B*, 2017, **204**, 440–455.

213 K. Teramura, S. Iguchi, Y. Mizuno, T. Shishido and T. Tanaka, *Angew. Chem., Int. Ed.*, 2012, **51**, 8008–8011.

214 Y. Zhao, X. Jia, G. I. N. Waterhouse, L.-Z. Wu, C.-H. Tung, D. O'Hare and T. Zhang, *Adv. Energy Mater.*, 2016, **6**, 1501974.

215 J. Ren, S. Ouyang, H. Xu, X. Meng, T. Wang, D. Wang and J. Ye, *Adv. Energy Mater.*, 2017, **7**, 1601657.

216 D. Mateo, J. Albero and H. García, *Joule*, 2019, **3**, 1949–1962.

217 X. Li, H. O. Everitt and J. Liu, *Nano Res.*, 2019, **12**, 1906–1911.

218 C. Kim, S. Hyeon, J. Lee, W. D. Kim, D. C. Lee, J. Kim and H. Lee, *Nat. Commun.*, 2018, **9**, 3027.

219 G. Chen, R. Gao, Y. Zhao, Z. Li, G. I. N. Waterhouse, R. Shi, J. Zhao, M. Zhang, L. Shang, G. Sheng, X. Zhang, X. Wen, L.-Z. Wu, C.-H. Tung and T. Zhang, *Adv. Mater.*, 2018, **30**, 1704663.

220 L. Liu, A. V. Puga, J. Cored, P. Concepcion, V. Perez-Dieste, H. Garcia and A. Corma, *Appl. Catal., B*, 2018, **235**, 186–196.

221 X.-N. Guo, Z.-F. Jiao, G.-Q. Jin and X.-Y. Guo, *ACS Catal.*, 2015, **5**, 3836–3840.

222 S. Yu, T. Zhang, Y. Xie, Q. Wang, X. Gao, R. Zhang, Y. Zhang and H. Su, *Int. J. Hydrogen Energy*, 2015, **40**, 870–877.

223 C.-x. Xiao, Z.-p. Cai, T. Wang, Y. Kou and N. Yan, *Angew. Chem., Int. Ed.*, 2008, **47**, 746–749.

224 L. Wang, Y. Zhang, X. Gu, Y. Zhang and H. Su, *Catal. Sci. Technol.*, 2018, **8**, 601–610.

225 Z. Li, J. Liu, Y. Zhao, G. I. Waterhouse, G. Chen, R. Shi, X. Zhang, X. Liu, Y. Wei and X. D. Wen, *Adv. Mater.*, 2018, **30**, 1800527.

226 W. Gao, R. Gao, Y. Zhao, M. Peng, C. Song, M. Li, S. Li, J. Liu, W. Li, Y. Deng, M. Zhang, J. Xie, G. Hu, Z. Zhang, R. Long, X.-D. Wen and D. Ma, *Chem.*, 2018, **4**, 2917–2928.

227 Z. Li, J. Liu, Y. Zhao, R. Shi, G. I. Waterhouse, Y. Wang, L.-Z. Wu, C.-H. Tung and T. Zhang, *Nano Energy*, 2019, **60**, 467–475.

228 Y. Wang, Y. Zhao, J. Liu, Z. Li, G. I. Waterhouse, R. Shi, X. Wen and T. Zhang, *Adv. Energy Mater.*, 2020, **10**, 1902860.

229 V. P. Ananikov, *ACS Catal.*, 2015, **5**, 1964–1971.

230 H. Song, X. Meng, Z.-J. Wang, H. Liu and J. Ye, *Joule*, 2019, **3**, 1606–1636.

231 G. Chen, G. I. N. Waterhouse, R. Shi, J. Zhao, Z. Li, L.-Z. Wu, C.-H. Tung and T. Zhang, *Angew. Chem., Int. Ed.*, 2019, **58**, 17528–17551.

232 H. Liu, X. Meng, T. D. Dao, H. Zhang, P. Li, K. Chang, T. Wang, M. Li, T. Nagao and J. Ye, *Angew. Chem., Int. Ed.*, 2015, **54**, 11545–11549.

233 B. László, K. Baán, E. Varga, A. Oszkó, A. Erdőhelyi, Z. Kónya and J. Kiss, *Appl. Catal., B*, 2016, **199**, 473–484.

234 B. Han, W. Wei, L. Chang, P. Cheng and Y. H. Hu, *ACS Catal.*, 2016, **6**, 494–497.

235 H. Liu, M. Li, T. D. Dao, Y. Liu, W. Zhou, L. Liu, X. Meng, T. Nagao and J. Ye, *Nano Energy*, 2016, **26**, 398–404.

236 H. Liu, T. D. Dao, L. Liu, X. Meng, T. Nagao and J. Ye, *Appl. Catal., B*, 2017, **209**, 183–189.

237 H. Liu, X. Meng, T. D. Dao, L. Liu, P. Li, G. Zhao, T. Nagao, L. Yang and J. Ye, *J. Mater. Chem. A*, 2017, **5**, 10567–10573.

238 Q. Zhang, M. Mao, Y. Li, Y. Yang, H. Huang, Z. Jiang, Q. Hu, S. Wu and X. Zhao, *Appl. Catal., B*, 2018, **239**, 555–564.

239 H. Huang, M. Mao, Q. Zhang, Y. Li, J. Bai, Y. Yang, M. Zeng and X. Zhao, *Adv. Energy Mater.*, 2018, **8**, 1702472.

240 H. Liu, H. Song, W. Zhou, X. Meng and J. Ye, *Angew. Chem., Int. Ed.*, 2018, **57**, 16781–16784.

241 M. Mao, Q. Zhang, Y. Yang, Y. Li, H. Huang, Z. Jiang, Q. Hu and X. Zhao, *Green Chem.*, 2018, **20**, 2857–2869.

242 F. Pan, X. Xiang, W. Deng, H. Zhao, X. Feng and Y. Li, *ChemCatChem*, 2018, **10**, 940–945.

243 Z. Jiang, Y. Li, Q. Zhang, Y. Yang, S. Wu, J. Wu and X. Zhao, *J. Mater. Chem. A*, 2019, **7**, 4881–4892.

244 H. Jiang, X. Peng, A. Yamaguchi, S. Ueda, T. Fujita, H. Abe and M. Miyauchi, *Sol. RRL*, 2019, **3**, 1900076.

245 S. Shoji, X. Peng, A. Yamaguchi, R. Watanabe, C. Fukuhara, Y. Cho, T. Yamamoto, S. Matsumura, M.-W. Yu and S. Ishii, *Nat. Catal.*, 2020, **3**, 148–153.

246 Y. Cho, S. Shoji, A. Yamaguchi, T. Hoshina, T. Fujita, H. Abe and M. Miyauchi, *Chem. Commun.*, 2020, **56**, 4611–4614.

247 L. Zhou, J. M. P. Martirez, J. Finzel, C. Zhang, D. F. Swearer, S. Tian, H. Robatjazi, M. Lou, L. Dong, L. Henderson, P. Christopher, E. A. Carter, P. Nordlander and N. J. Halas, *Nat. Energy*, 2020, **5**, 61–70.

248 V. Smil, *Nature*, 1999, **400**, 415.

249 J. W. Erisman, M. A. Sutton, J. Galloway, Z. Klimont and W. Winiarwter, *Nat. Geosci.*, 2008, **1**, 636–639.

250 C. Mao, L. Yu, J. Li, J. Zhao and L. Zhang, *Appl. Catal., B*, 2018, **224**, 612–620.

251 D. F. Swearer, N. R. Knowles, H. O. Everitt and N. J. Halas, *ACS Energy Lett.*, 2019, **4**, 1505–1512.

252 X. Li, X. Zhang, H. O. Everitt and J. Liu, *Nano Lett.*, 2019, **19**, 1706–1711.

253 S. Duhr and D. Braun, *Proc. Natl. Acad. Sci. U. S. A.*, 2006, **103**, 19678.

254 T. Oshikiri, K. Ueno and H. Misawa, *Angew. Chem., Int. Ed.*, 2016, **55**, 3942–3946.

255 C. Li, T. Wang, Z.-J. Zhao, W. Yang, J.-F. Li, A. Li, Z. Yang, G. A. Ozin and J. Gong, *Angew. Chem., Int. Ed.*, 2018, **57**, 5278–5282.

256 C. Hu, X. Chen, J. Jin, Y. Han, S. Chen, H. Ju, J. Cai, Y. Qiu, C. Gao, C. Wang, Z. Qi, R. Long, L. Song, Z. Liu and Y. Xiong, *J. Am. Chem. Soc.*, 2019, **141**, 7807–7814.

257 T. Bunluesin, R. J. Gorte and G. W. Graham, *Appl. Catal., B*, 1998, **15**, 107–114.

258 W.-H. Chen and C.-Y. Chen, *Appl. Energy*, 2020, **258**, 114078.

259 Y. Choi and H. G. Stenger, *J. Power Sources*, 2003, **124**, 432–439.

260 F. Liu, L. Song, S. Ouyang and H. Xu, *Catal. Sci. Technol.*, 2019, **9**, 2125–2131.

261 L. Zhao, Y. Qi, L. Song, S. Ning, S. Ouyang, H. Xu and J. Ye, *Angew. Chem., Int. Ed.*, 2019, **58**, 7708–7712.

262 Q. Li and G. Lu, *Catal. Lett.*, 2008, **125**, 376–379.

263 L. Gu, C. Zhang, Y. Guo, J. Gao, Y. Yu and B. Zhang, *ACS Sustainable Chem. Eng.*, 2019, **7**, 3710–3714.

264 W. Hao, R. Wu, H. Yang and Y. Guo, *J. Mater. Chem. A*, 2019, **7**, 12440–12445.

265 L. Zhao, Z. Yang, Q. Cao, L. Yang, X. Zhang, J. Jia, Y. Sang, H.-J. Wu, W. Zhou and H. Liu, *Nano Energy*, 2019, **56**, 563–570.

266 Y. Wu, Y. Sun, W. Fu, X. Meng, M. Zhu, S. Ramakrishna and Y. Dai, *ACS Appl. Nano Mater.*, 2020, **3**, 2713–2722.

267 S. Rej, L. Mascaretti, E. Y. Santiago, O. Tomanec, S. Kment, Z. Wang, R. Zbořil, P. Fornasiero, A. O. Govorov and A. Naldoni, *ACS Catal.*, 2020, **10**, 5261–5271.

268 A. Staubitz, A. P. M. Robertson and I. Manners, *Chem. Rev.*, 2010, **110**, 4079–4124.

269 J. C. Kennedy III and A. K. Datye, *J. Catal.*, 1998, **179**, 375–389.

270 T. H. Tan, J. Scott, Y. H. Ng, R. A. Taylor, K.-F. Aguey-Zinsou and R. Amal, *ACS Catal.*, 2016, **6**, 1870–1879.

271 J. Bi, H. Cai, B. Wang, C. Kong and S. Yang, *Chem. Commun.*, 2019, **55**, 3943–3946.

272 J. Wang, G. Zhang and P. Zhang, *Appl. Catal., B*, 2018, **239**, 77–85.

273 X. Chen, S. Cai, E. Yu, J. Li, J. Chen and H. Jia, *Appl. Surf. Sci.*, 2019, **484**, 479–488.

274 Z. Zhang, C. Zhang, H. Zheng and H. Xu, *Acc. Chem. Res.*, 2019, **52**, 2506–2515.

275 C. Zhang, J. Lu, N. Jin, L. Dong, Z. Fu, Z. Zhang and H. Zheng, *Small*, 2019, **15**, 1901286.

276 H. Fujiwara, T. Suzuki, C. Pin and K. Sasaki, *Nano Lett.*, 2019, **20**, 389–394.

277 J. Mock, M. Barbic, D. Smith, D. Schultz and S. Schultz, *J. Chem. Phys.*, 2002, **116**, 6755–6759.

278 J. Zhang, H. Liu, Z. Wang and N. Ming, *Adv. Funct. Mater.*, 2007, **17**, 3295–3303.

279 S. Dal Forno, L. Ranno and J. Lischner, *J. Phys. Chem. C*, 2018, **122**, 8517–8527.

280 M. Rycenga, C. M. Cobley, J. Zeng, W. Li, C. H. Moran, Q. Zhang, D. Qin and Y. Xia, *Chem. Rev.*, 2011, **111**, 3669–3712.

281 S. Link and M. A. El-Sayed, *J. Phys. Chem. B*, 1999, **103**, 4212–4217.

282 Q. Zhang, W. Li, C. Moran, J. Zeng, J. Chen, L.-P. Wen and Y. Xia, *J. Am. Chem. Soc.*, 2010, **132**, 11372–11378.

283 B. J. Wiley, Y. Chen, J. M. McLellan, Y. Xiong, Z.-Y. Li, D. Ginger and Y. Xia, *Nano Lett.*, 2007, **7**, 1032–1036.

284 B. J. Wiley, Y. Xiong, Z.-Y. Li, Y. Yin and Y. Xia, *Nano Lett.*, 2006, **6**, 765–768.

285 M. J. Mulvihill, X. Y. Ling, J. Henzie and P. Yang, *J. Am. Chem. Soc.*, 2010, **132**, 268–274.

286 H. Zhang and A. O. Govorov, *J. Phys. Chem. C*, 2014, **118**, 7606–7614.

287 L. V. Besteiro and A. O. Govorov, *J. Phys. Chem. C*, 2016, **120**, 19329–19339.

288 H. Wu, J. Carrete, Z. Zhang, Y. Qu, X. Shen, Z. Wang, L.-D. Zhao and J. He, *NPG Asia Mater.*, 2014, **6**, e108–e108.

289 J. Lim, K. Hippalgaonkar, S. C. Andrews, A. Majumdar and P. Yang, *Nano Lett.*, 2012, **12**, 2475–2482.

290 S. Link and M. A. El-Sayed, *Int. Rev. Phys. Chem.*, 2000, **19**, 409–453.

291 G. Baffou, R. Quidant and F. J. García de Abajo, *ACS Nano*, 2010, **4**, 709–716.

292 G. Baffou, R. Quidant and C. Girard, *Appl. Phys. Lett.*, 2009, **94**, 153109.

293 K. Manthiram and A. P. Alivisatos, *J. Am. Chem. Soc.*, 2012, **134**, 3995–3998.

294 Z. Fang, S. Jiao, Y. Kang, G. Pang and S. Feng, *ChemistryOpen*, 2017, **6**, 261–265.

295 R. Du, W. Liu, H. Bai, H. Wang and G. Xi, *RSC Adv.*, 2020, **10**, 2075–2084.

296 Y.-Z. Chen, Z. U. Wang, H. Wang, J. Lu, S.-H. Yu and H.-L. Jiang, *J. Am. Chem. Soc.*, 2017, **139**, 2035–2044.

297 J. Espín, L. Garzón-Tovar, A. Carné-Sánchez, I. Imaz and D. Maspoch, *ACS Appl. Mater. Interfaces*, 2018, **10**, 9555–9562.

298 X. Zhang, X. Ke and H. Zhu, *Chem. – Eur. J.*, 2012, **18**, 8048–8056.

299 S. Das, J. Pérez-Ramírez, J. Gong, N. Dewangan, K. Hidajat, B. C. Gates and S. Kawi, *Chem. Soc. Rev.*, 2020, **49**, 2937–3004.

300 D. Kumar, C. H. Park and C. S. Kim, *ACS Sustainable Chem. Eng.*, 2018, **6**, 8604–8614.

301 H. Huang, L. Zhang, Z. Lv, R. Long, C. Zhang, Y. Lin, K. Wei, C. Wang, L. Chen and Z.-Y. Li, *J. Am. Chem. Soc.*, 2016, **138**, 6822–6828.

302 L. Guo, Q. Sun, K. Marcus, Y. Hao, J. Deng, K. Bi and Y. Yang, *J. Mater. Chem. A*, 2018, **6**, 22005–22012.

303 S. Okeil, S. Yadav, M. Bruns, A. Zintler, L. Molina-Luna and J. J. Schneider, *Dalton Trans.*, 2020, **49**, 1032–1047.

304 P. Ren and X. Yang, *Sol. RRL*, 2018, **2**, 1700233.

305 X. Niu, Y. Li, Y. Zhang, Z. Zhou and J. Wang, *ACS Appl. Mater. Interfaces*, 2019, **11**, 17987–17993.

