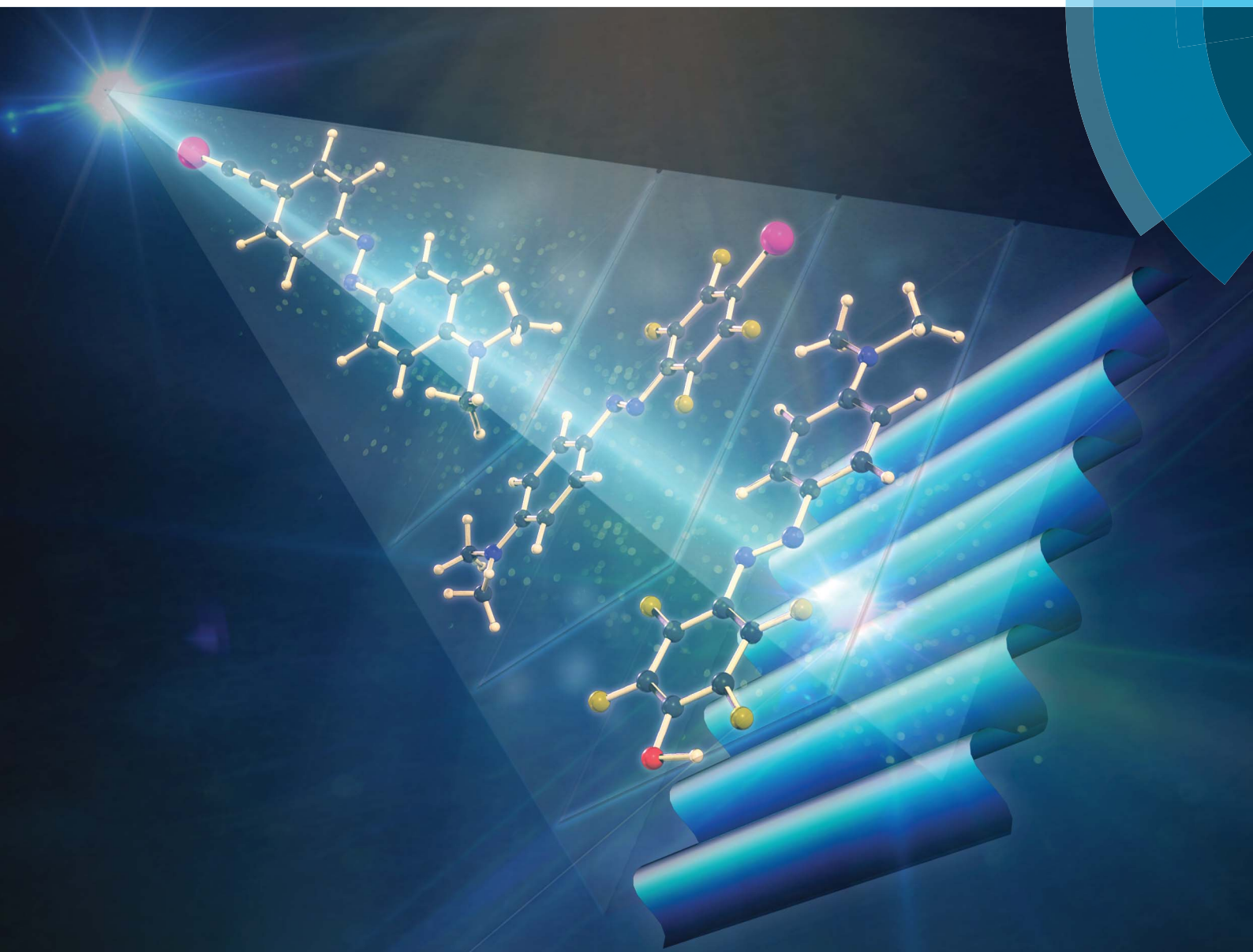


# Journal of Materials Chemistry C

Materials for optical, magnetic and electronic devices

[www.rsc.org/MaterialsC](http://www.rsc.org/MaterialsC)



ISSN 2050-7526



**PAPER**

Alessandra Forni, Giuseppe Resnati, Pierangelo Metrangolo,  
Arri Priimagi *et al.*  
Supramolecular hierarchy among halogen and hydrogen bond donors  
in light-induced surface patterning

CrossMark  
click for updatesCite this: *J. Mater. Chem. C*, 2015, 3, 759

# Supramolecular hierarchy among halogen and hydrogen bond donors in light-induced surface patterning†

Marco Saccone,<sup>a</sup> Valentina Dichiarante,<sup>b</sup> Alessandra Forni,<sup>\*c</sup> Alexis Goulet-Hanssens,<sup>‡d</sup> Gabriella Cavallo,<sup>b</sup> Jaana Vapaavuori,<sup>§a</sup> Giancarlo Terraneo,<sup>b</sup> Christopher J. Barrett,<sup>d</sup> Giuseppe Resnati,<sup>\*b</sup> Pierangelo Metrangolo<sup>\*be</sup> and Arri Priimagi<sup>\*abf</sup>

Halogen bonding, a noncovalent interaction possessing several unique features compared to the more familiar hydrogen bonding, is emerging as a powerful tool in functional materials design. Herein, we unambiguously show that one of these characteristic features, namely high directionality, renders halogen bonding the interaction of choice when developing azobenzene-containing supramolecular polymers for light-induced surface patterning. The study is conducted by using an extensive library of azobenzene molecules that differ only in terms of the bond-donor unit. We introduce a new tetrafluorophenol-containing azobenzene photoswitch capable of forming strong hydrogen bonds, and show that an iodoethynyl-containing azobenzene comes out on top of the supramolecular hierarchy to provide unprecedented photoinduced surface patterning efficiency. Specifically, the iodoethynyl motif seems highly promising in future development of polymeric optical and photoactive materials driven by halogen bonding.

Received 14th October 2014  
Accepted 21st November 2014

DOI: 10.1039/c4tc02315c

[www.rsc.org/MaterialsC](http://www.rsc.org/MaterialsC)

## Introduction

The azobenzene moiety is an important tool in present-day research of molecular-level photoswitching and is extensively used for the development of macroscopic photonic and photo-mechanical materials.<sup>1</sup> The interest in azobenzene lies in its clean, efficient, and reversible photoisomerization between a

rod-like *trans*-form and a bent *cis*-form,<sup>2</sup> which can give rise to a cascade of motions beyond the size scale of an individual molecule. When the azobenzene is incorporated into a liquid-crystalline or a polymeric medium, the photoisomerization can accomplish various tasks, allowing for induction of order-disorder transitions, control over the molecular alignment, photoactuation of the shape and dimensions of the material system, and inscription of surface-relief patterns onto thin azobenzene-containing films.<sup>3</sup> In particular, when an azobenzene-containing thin polymer film is irradiated with a polarization/intensity interference pattern of light, the film surface deforms and creates a replica of the incident interference pattern in the form of a surface topography pattern, often denoted as a surface-relief grating (SRG).<sup>4</sup> The modulation depth of the photoinduced surface patterns can be hundreds of nanometers, creating efficient diffractive optical elements with a first-order diffraction efficiency of up to 30–40%. SRGs have been used as nanostructuring tools, liquid-crystal alignment layers, and periodic light trapping/light outcoupling layers.<sup>5</sup> Despite their application potential, however, the surface pattern formation mechanism is not yet well understood. Therefore, a full understanding of structure–function relationships that govern the SRG process is needed in order to design better performing materials for future advanced applications.

Supramolecular side-chain polymers, in which the photoactive azobenzene units are attached to a passive polymer backbone *via* noncovalent interactions, have in the past few

<sup>a</sup>Department of Applied Physics, Aalto University, P.O. Box 13500, FI-00076 Aalto, Finland

<sup>b</sup>NFMLab-DCMIC “Giulio Natta”, Politecnico di Milano, Via L. Mancinelli 7, IT-20131 Milano, Italy. E-mail: pierangelo.metrangolo@polimi.it; giuseppe.resnati@polimi.it

<sup>c</sup>ISTM-CNR, Institute of Molecular Sciences and Technologies of CNR, Università degli Studi di Milano, Via Golgi 33, IT-20133 Milano, Italy. E-mail: alessandra.forni@istm.cnr.it

<sup>d</sup>Department of Chemistry, McGill University, 801 Sherbrooke Street West, Montreal, Quebec H3A 0B8, Canada

<sup>e</sup>VTT-Technical Research Centre of Finland, P.O. Box 1000, FI-02044 VTT, Finland

<sup>f</sup>Department of Chemistry and Bioengineering, Tampere University of Technology, P. O. Box 541, FI-33101 Tampere, Finland. E-mail: arri.priimagi@tut.fi

† Electronic supplementary information (ESI) available: Synthesis procedures and compound characterization, computational details, description of photochemical experiments, and crystal data. CCDC 1025655 and 1025656. For ESI and crystallographic data in CIF or other electronic format see DOI: 10.1039/c4tc02315c

‡ Present address: Department of Chemistry, Humboldt-Universität zu Berlin, Brook-Taylor-Str. 2, 12489 Berlin, Germany.

§ Present address: Département de Chimie, Université de Montréal, Montréal, QC, Canada H3C 3J7.



years emerged as a new class of SRG-forming materials.<sup>6</sup> The importance of these materials lies in their extreme ease of preparation and modular tunability: the material composition can be easily optimized to meet specific requirements. As a result, high-quality resonance filters, and patterns with record-high modulation depth have been fabricated from azobenzene-containing supramolecular side-chain polymers.<sup>6a,7</sup> We have recently shown that halogen bonding<sup>8</sup> appears particularly promising for the design of supramolecular SRG-forming materials.<sup>9</sup> Halogen bonding is the noncovalent interaction in which a region of positive electrostatic potential on top of the surface of the halogen atom, the  $\sigma$ -hole,<sup>10</sup> interacts with a nucleophilic site. High-level *ab initio* quantum-chemical calculations suggest that dispersion forces may also play an important role in this interaction.<sup>11</sup> The features that distinguish halogen bonding from hydrogen bonding are of particular interest, which include higher directionality, hydrophobicity, size of the interacting atom (halogen *vs.* hydrogen) and the possibility to tune interaction strength by a simple change of the halogen atom.<sup>12</sup> These characteristic properties have allowed for the fine tuning of geometrical features relevant to crystal engineering<sup>13</sup> and to the design of functional materials with unique light-emissive<sup>14</sup> and liquid-crystalline<sup>15</sup> properties.

Recently several examples of halogen-bonded supramolecular structures involving azobenzenes have been reported.<sup>9a,16</sup> For example, we showed that the efficiency of SRG formation increases with the strength of halogen bonding, and that the high performance of halogen-bonded complexes seems to be connected more to the directionality of halogen bonding than to its intrinsic interaction strength. To further develop the practical applications of halogen bonding in the field of photocontrollable materials, we have undertaken a detailed study on the relative surface patterning efficiency that can be imparted by different halogen-bond donors as a result of the presence of electron-withdrawing substituents *versus* the hybridization of the carbon bearing the heavy halogen. Moreover, we aimed to provide a ranking between various halogen and hydrogen bonds in driving SRG formation, since these two interactions show similar chemical behaviour and a detailed knowledge of the performance of competing noncovalent interactions is crucial for the preparation of materials with designed structures and properties.

Herein, we present a comprehensive study on the surface patterning properties of an extensive library of azobenzene derivatives as presented in Fig. 1. We show that halogen-bonded polymer–azobenzene complexes unambiguously surpass their hydrogen-bonded counterparts in terms of SRG formation efficiency, confirming that it is the directionality rather than the strength of the noncovalent interactions that dictates the macroscopic light-induced movements. Most importantly, we introduce in this field a new halogen-bond donor moiety – iodoacetylene – that further boosts the macroscopic light-induced movements and the surface pattern formation. The optical studies are complemented with theoretical calculations and crystallographic structure determinations that further elucidate the different performance of halogen- and hydrogen-bonded systems.

## Results and discussion

### Materials design

The library of azobenzene compounds shown in Fig. 1 allows for (i) establishing an optimal halogen-bond donor in terms of SRG formation capability, and (ii) comparing structurally similar azobenzenes substituted with halogen-bond *versus* hydrogen-bond donor groups. The series of perfluorinated azobenzenes 1–5 comprises a reference molecule incapable of forming halogen bonds (1), two halogen-bond donors (2, 3) with similar spectral and photochemical properties but different interaction strengths ( $\text{Br} < \text{I}$ ), and two hydrogen-bond donors, one weak (4) and one strong perfluorophenol-containing dye (5). Compound 5 serves a dual purpose. Firstly, a close structural relationship with 3 affords an ideal comparison for halogen bonding *versus* hydrogen bonding in driving the SRG formation process. Secondly, the perfluorinated phenol presumably acts as a very powerful hydrogen-bond donor and could therefore give rise to a very strong noncovalent interaction,<sup>17,18</sup> which is interesting in its own right in the field of functional supramolecular materials. This is the rationale for the comparison between 5 and its nonfluorinated counterpart 7, the SRG performance of which is known.<sup>9</sup>

The non-fluorinated dyes 6–9 include the already mentioned phenol-substituted compound 7, an iodo-substituted azobenzene 6 that is expected to act as a weak halogen-bond donor, and ethynyl-substituted dyes 8 and 9. The latter two compounds are particularly interesting, since very recently iodoethynylbenzenes were investigated both theoretically and experimentally with the purpose of establishing a hierarchy among halogen-bond donors in crystal engineering.<sup>19</sup> Based on ATR-FTIR and single crystal X-ray diffraction data, Aakeröy *et al.* proved that the iodoethynyl benzenes form co-crystals with a larger class of halogen-bond acceptors compared to perfluoroiodobenzenes suggesting that this motif might be employed to enhance the performance of photoresponsive polymers. Each of the studied molecules is substituted with an electron-donating dimethylamino group, which red-shifts the absorption spectrum of the dyes, increases the rate of *trans-cis-trans* cycling, and boosts the SRG formation efficiency.

Poly(4-vinyl pyridine) (P4VP) was used as a polymeric bond acceptor as it is a popular choice in constructing supramolecular polymeric complexes based on hydrogen bonding or proton transfer,<sup>20</sup> and it equally functions as a halogen-bond acceptor.<sup>21</sup> P4VP has been used by several groups to prepare light-responsive polymers,<sup>6b,22</sup> making it the polymer of choice for the present study. Low-molecular-weight P4VP ( $M_w$ : 1200 g mol<sup>-1</sup>) was selected since long polymer chains tend to disrupt the photoinduced surface patterning. SRG formation was studied in polymer–azobenzene complexes, denoted as P4VP(*n*), where *n* is the azobenzene unit chosen among 1–9 and *y* is the *n* : P4VP molar ratio. Typically, *y* was set to be 0.1, or, in other words, the studied systems contained one azobenzene moiety per ten polymer repeat units. This low degree of complexation minimizes crystallization and macroscopic phase separation of the azobenzene units, which would compromise



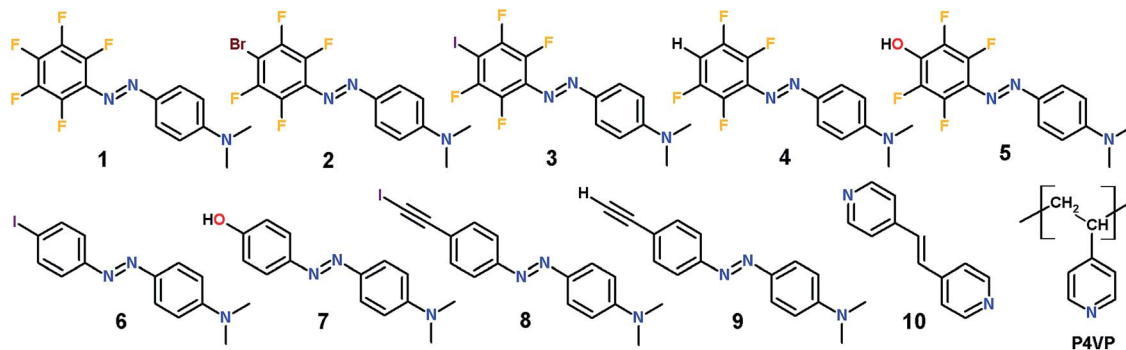


Fig. 1 Chemical structures of fluorinated azobenzene modules 1–5, hydrogenated azobenzene modules 6–9, model dipyrindyl compound 10 used in the co-crystallization studies, and poly(4-vinyl pyridine) (P4VP) that was used as the polymer matrix for the azobenzene dyes.

comparisons between the molecules. Assuming that (*E*)-1,2-di(4-pyridyl)ethylene (**10**) is a good mimic of the minimal binding motif of P4VP, the co-crystals formed by **10** with various azobenzene derivatives were used as models for single-crystal X-ray diffraction analysis of the P4VP–azobenzene systems employed in light-induced surface patterning. 4-Methyl pyridine was used as a model of the minimal binding motif of P4VP in computational studies.

### Computational studies

As mentioned above, 4-methyl pyridine was used as the smallest pyridine model compound mimicking P4VP to compute the binding energy with azobenzenes 1–9. Table 1 presents the PBE0/6-311++G(d,p) interaction energies, both non-corrected ( $\Delta E$ ) and corrected ( $\Delta E_{\text{BSSE}}$ ) for the basis set superposition error, showing a notable energy overestimation (from 6 to 14% depending on the dimer) when neglecting such an error. The  $\Delta E_{\text{BSSE}}$  values computed for the perfluorinated dyes **2** and **3** were  $-3.501$  and  $-5.135$  kcal mol $^{-1}$ ,<sup>9</sup> respectively, in agreement with the consolidated trend of Br < I in halogen-bonded systems,<sup>10</sup> while molecule **1**, lacking the  $\sigma$ -hole on the fluorine atom in the *para* position to the azo group, was confirmed to be

unable to give rise to a stable halogen-bonded complex. For the complex of 4-methyl pyridine with the perfluorinated hydrogen-bond donor **5**, a stabilization energy of  $-11.765$  kcal mol $^{-1}$  was obtained, corresponding to a medium-to-strong hydrogen bond.<sup>23</sup> The corresponding iodinated and phenolic non-fluorinated donors, **6** and **7**, provided stabilization energy values of  $-2.546$  and  $-10.053$  kcal mol $^{-1}$ , respectively, upon binding to the pyridine unit.

The key observation here is that perfluorination of the benzene ring is significantly more effective in strengthening halogen bonding than hydrogen bonding: the interaction strength, in fact, doubles in the first case (**6** vs. **3**), whereas a minor <20% enhancement takes place for the latter (**7** vs. **5**). This difference can be rationalized by examining the changes in charge density distribution around iodine and the phenolic hydrogen in the isolated dyes upon fluorination. While the molecular dipole moment (Table 1,  $\mu_{\text{calc}}$ ) undergoes a greater increase for the phenolic with respect to the iodinated derivative, the electrostatic potential (Fig. 2) is practically unaffected by fluorination near the phenolic hydrogens, though it is strongly positive in both cases (maximum values of  $V_{\text{S,max}}$  are 0.083 and 0.086 au for **7** and **5**, respectively, on the 0.001 a.u. isosurface of electron density). Conversely, the maximum value of the electrostatic potential around the iodine atom, while less positive, doubles upon ring fluorination, from 0.018 a.u. (**6**) to 0.035 au (**3**). This can be explained by polarizability arguments: it is well-known that halogen atoms – apart from fluorine – are highly polarizable, and more so for the heavier halogens. As a result, the strength of halogen bonding significantly increases by introduction of electron-withdrawing groups such as fluorine atoms on the molecular moiety bonded to the halogen-bond donor site.<sup>24</sup>

As recently shown by a systematic investigation of a series of halogen-bond donors,<sup>19</sup> the hybridization of the carbon atom adjacent to the halogen bonded atom heavily influences the strength of the formed halogen bonds. In particular, iodoethynyl-containing moieties have recently emerged as remarkably powerful halogen-bond donor moieties.<sup>19</sup> In agreement with such findings, compound **8** shows, upon binding to the model pyridine unit, a slightly higher stabilization energy ( $-5.170$  kcal mol $^{-1}$ ) than compound **3**. It is also noteworthy that

Table 1 Non-corrected and corrected interaction energies of compounds 1–9 with 4-methyl pyridine ( $\Delta E$  and  $\Delta E_{\text{BSSE}}$ , respectively, kcal mol $^{-1}$ ), local most positive values of electrostatic potential computed on an isosurface (0.001 a.u.) of electron density ( $V_{\text{S,max}}$ , a.u.) and ground-state dipole moments ( $\mu_{\text{calc}}$ , D) of 1–9. Interaction energies and  $V_{\text{S,max}}$  values were computed *in vacuo* and dipole moments in DMF. All properties were evaluated at the PBE0/6-311++G(d,p) level of theory

Compound	$\Delta E$	$\Delta E_{\text{BSSE}}$	$V_{\text{S,max}}$	$\mu_{\text{calc}}$
<b>1</b>	—	—	−0.018	10.67
<b>2</b>	−3.712	−3.501	0.032	10.47
<b>3</b>	−5.565	−5.135	0.035	10.22
<b>4</b>	−4.477	−3.930	0.039	8.94
<b>5</b>	−12.644	−11.765	0.086	7.63
<b>6</b>	−2.828	−2.546	0.018	9.28
<b>7</b>	−10.796	−10.053	0.083	5.65
<b>8</b>	−5.533	−5.170	0.041	8.36
<b>9</b>	−3.832	−3.528	0.043	8.30



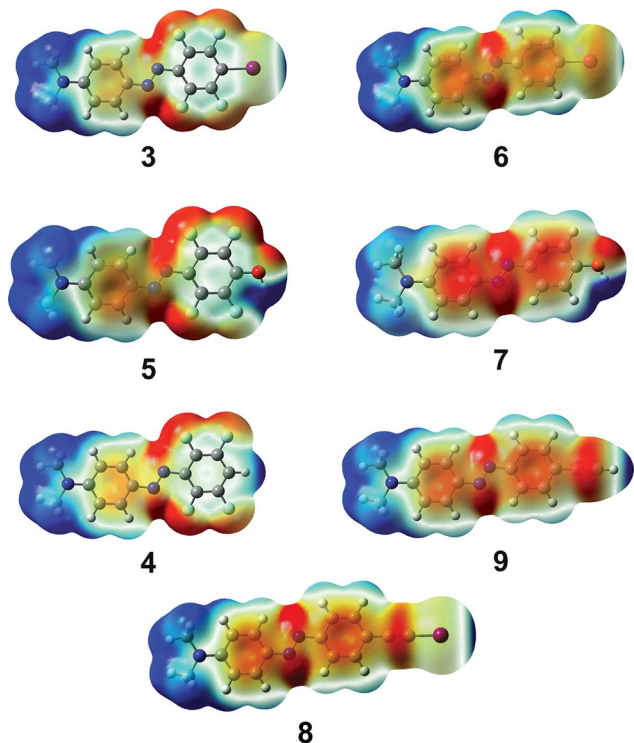


Fig. 2 Plots of the electrostatic potential of compounds 3, 6 (first row), 5, 7 (second row), 4, 9 (third row), and 8 (bottom), computed at the PBE0/6-311++G(d,p) level *in vacuo*. Potentials are mapped on the respective isosurfaces (0.001 a.u.) of electron density. Values of electrostatic potential range from  $-0.03$  (red) to  $0.03$  (blue) a.u. Atom color scheme: C, gray; H, light gray; N, dark blue; O, red; F, sky blue; and I, magenta.

the maximum value of the electrostatic potential around the iodine atom is  $0.041$  a.u., notably larger than the corresponding value obtained for 3 ( $0.035$  a.u.), suggesting that the iodoethynyl moiety should give a more directional halogen bonding.

Such observation prompted us to check the different directionality features of the corresponding hydrogenated derivatives, 9 vs. 4. While the hydrogen bond of the former with 4-methyl pyridine is slightly weaker than that of the latter (stabilization energies are  $-3.528$  and  $-3.930$  kcal mol $^{-1}$ , respectively), the ethynyl derivative is characterized by a highly sharpened positive region in the electrostatic potential (Fig. 2), with a more positive  $V_{S,max}$  value with respect to compound 4 ( $0.043$  and  $0.039$  a.u., respectively). This suggests a potentially much greater directionality for the hydrogen bonding of the ethynyl derivative with respect to the perfluorinated compound, somewhat resembling halogen bonding.

It is thus clear that 4-methyl pyridine is more tightly bound to the phenolic and hydrogen-bonding donor dyes 5 and 7 than to the halogen-bonding donor dyes 3 and 6, but it is also evident from Fig. 2 that halogen bonds are expected to be significantly more directional,<sup>25</sup> as the positive region responsible for the noncovalent interaction is narrowly confined along the extension of the C–I bonds and it is spread out hemispherically around the phenolic hydrogens. On the other hand, the dimeric complexes comprising 4-methyl pyridine and ethynyl

derivatives 8 and 9 show quite similar directionality features but the interaction strength of the iodine derivative is greater. In our hypothesis, a more directional noncovalent interaction between the azobenzene dye and the P4VP provides a more efficient junction in transferring the light-induced motions of the photoactive units to the polymer backbone and enhanced SRG formation results.<sup>9</sup> Molecules 3 and 5 comprise an ideal pair to gain further insight into this hypothesis.

We also computed the UV-visible spectra for 1–9 by time-dependent DFT calculations (Table S1†), which showed that perfluorination of the benzene ring red-shifts the absorption maxima of both iodine and phenolic derivatives, as was also verified experimentally by measuring the absorption spectra in DMF solutions. Moreover, the spectroscopic properties of the tetrafluoroiodo and iodoethynyl derivatives 3 and 8, respectively, are quite similar indicating that the two halogen-bond donors share many parallel features. Such findings are extremely helpful for the design of new functional materials based on halogenated ethynyl moieties. Details on computations are given in the ESI.†

### Structural characterization

(*E*)-1,2-Di(4-pyridyl)ethylene (**10**) was used as a co-crystal forming component that mimics P4VP. The aim was to get specific insight into the binding between the pyridine moiety and the azobenzene derivatives 5 and 8 which are new dyes introduced here. These two molecules were co-crystallized with the bis-pyridine derivative **10** in a 2 : 1 ratio.

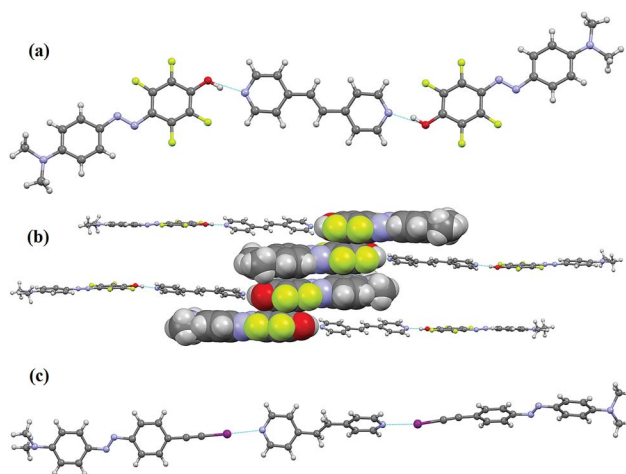


Fig. 3 (a) Ball-and-stick representation of a  $10 \cdot (5)_2$  hydrogen-bonded trimer. The diazo-molecule is disordered by rotation of  $180^\circ$  around the  $O \cdots N$  (amino) axis and only one of the two models is shown. (b) View along the crystallographic *b*-axis of four  $10 \cdot (5)_2$  trimers. Four molecules of compound 5, giving strong quadrupolar interactions, are shown in a spacefill style. (c) Ball-and-stick representation of a  $10 \cdot (8)_2$  halogen-bonded trimer. Both halogen-bond donor and acceptor are heavily disordered over two positions and here only one of the two models is represented. Hydrogen and halogen bonds are pictured as blue dotted lines. Colour code: carbon, grey; hydrogen, white; oxygen, red; nitrogen, blue; fluorine, green and iodine, purple.



Good-quality single crystals of the complex  $10 \cdot (5)_2$  were obtained by slow diffusion of diethyl ether into a methanol solution containing the two starting compounds at room temperature. X-ray diffraction analysis shows that each molecule of **10** binds two different phenol dyes **5** via short N $\cdots$ H hydrogen bonds (Fig. 3a). The O $\cdots$ N distance is 2.557(2) Å, the C–O $\cdots$ N angle is 125.68(3)° and a sigmoidal trimeric unit is formed. The observed hydrogen bond is among the shortest found in the Cambridge Structural Database (CSD) and is comparable to that found in the pentachlorophenol/4-methylpyridine co-crystal wherein the O $\cdots$ N distance is 2.515 Å.<sup>26</sup> The O–H $\cdots$ N hydrogen bond present in the co-crystal  $10 \cdot (5)_2$  is shorter than that found in a similar co-crystal between the nonfluorinated analogue **7** and 1,2-di-pyridylethane (O $\cdots$ N distance 2.703 Å).<sup>9</sup> This difference in hydrogen bond distances of co-crystals formed by the fluorinated and hydrogenated azophenols **5** and **7** parallels very well the difference in interaction energies from theoretical calculations (Table 1), *i.e.*, the enhanced electron density acceptor capability of the phenolic hydrogen in **5**, leading to a shorter and stronger hydrogen bond than in **7**, is consistently proved by experimental and computational studies.

In the three-dimensional crystal packing, the  $10 \cdot (5)_2$  trimers form well-defined anti-parallel dimers due to an off-set stacking of the tetrafluorobenzene rings (the centroid–centroid distance is 3.445 Å) (Fig. 3b).<sup>27</sup> Further  $\pi \cdots \pi$  stacking involves the tetrafluorobenzene rings and the dimethylaminobenzene rings of adjacent dimers or trimers (the centroid–centroid distance is 3.684 Å).<sup>28</sup> Such stacking interactions, while useful in crystal engineering<sup>29</sup> as well as in liquid-crystal self-assembly,<sup>30</sup> may cause dye aggregation within polymer matrices, thus preventing high doping concentrations that would be beneficial in optoelectronic applications.

The co-crystal formed between **8** and **10** is a strong example of the high directionality of halogen bonding. Single crystals of  $10 \cdot (8)_2$  were obtained by slow evaporation at room temperature of a THF solution containing the two starting modules. The single crystal analysis revealed the formation of rod-like halogen-bonded trimers in which a 1,2-di(4-pyridyl)ethylene molecule bridges two iodoethynyl azo-dyes (Fig. 3c) by short and linear halogen bonds with N $\cdots$ I distance of 2.754(3) Å and C–I $\cdots$ N angle of 175.29(10)°. The disorder present in the structure prevents the possible further aggregation of these trimers to be analysed. The structural characterization of the model co-crystals highlights that both of the

new dyes **5** and **8** interact strongly with pyridine-bearing entities; therefore their incorporation into the P4VP matrix may lead to the formation of new supramolecular polymeric materials with efficient photoswitching functionalities.

### Photochemical studies

Table 2 shows the absorption maxima, relative decreases in absorbance upon irradiation (488 nm, 50 mW cm<sup>−2</sup>, circular polarization), and time constant estimates of the *cis*–*trans* thermal isomerization of thin films of the supramolecular complexes between P4VP and dyes **1**–**9**. The absorbance decrease relates to the amount of *cis*-isomers in the photostationary state. The thermal *cis*–*trans* relaxation is a known non-monoexponential process in heterogeneous polymer environments and fitting the relaxation curves as a sum of first-order processes lacks physical justification.<sup>31</sup> Hence the relaxation time constants, which are obtained using reported procedures,<sup>32</sup> are only approximate and are given to facilitate the comparison among compounds of the molecular library of Fig. 1.

Apart from compounds **6** and **7**, the absorption maxima of the molecules studied are very similar. Fluorination of the bond-donor ring red-shifts the  $\lambda_{\text{max}}$  by 25 nm for the iodo-substituted azobenzenes (**6** vs. **3**) and 48 nm upon phenol-substitution (**7** vs. **5**). The ethynyl group plays a similar, though not as pronounced, role (**6** vs. **8**), due to an increase in the Hückel conjugation length. The absorbance decrease upon illumination is very similar, around 30% for compounds **2**–**6**, and around 20% for **1** and **7**. Due to similar *trans*–*cis* isomerization efficiencies, the photoinduced surface patterning may be compared well with each other within our set of molecules. Interestingly, the photostationary state is much more *cis*-rich for the ethynyl-substituted molecules **8** and **9**, and the lifetime of the *cis*-isomer is longer, possibly because of the larger aspect ratio of these molecules as compared to those of **1**–**7**.

To summarize, the compounds shown in Fig. 1 cover the range from weak to strong halogen and hydrogen bond donors and the fact that a central core providing similar absorption properties can be maintained allows for facile and valid comparisons between these dyes in the same polymer matrix.

### Supramolecular hierarchy in SRG formation

The molecular library presented in Fig. 1 allows us to explore various aspects of the surface patterning process, particularly to

**Table 2** Absorption maxima, percentage decrease in absorbance in the photostationary state (pss) at  $\lambda_{\text{max}}$ , estimates for the thermal lifetimes of the *cis*-azobenzenes, diffraction efficiencies (DE) and modulation depths (*d*) of the inscribed SRGs

	P4VP(1)	P4VP(2)	P4VP(3)	P4VP(4)	P4VP(5)	P4VP(6)	P4VP(7)	P4VP(8)	P4VP(9)
$\lambda_{\text{max}}^a$ (nm)	456	467	467	454	463	442	415	454	453
$A_{\text{pss}}/A_0^a$	0.79	0.70	0.69	0.70	0.73	0.69	0.78	0.52	0.53
$\tau_{\text{cis-trans}}^a$ (s)	910	350	620	620	250	370	440	1250	1250
DE <sup>b</sup> (%)	0.6	2.9	6.5	1.8	6.0	2.4	5.0	7.5	3.4
<i>d</i> <sup>b</sup> (nm)	30	105	175	80	170	95	130	180	115

<sup>a</sup> Obtained from drop-cast films of P4VP(*n*)<sub>0.01</sub> with thickness of several microns. <sup>b</sup> Obtained from spin-coated films of P4VP(*n*)<sub>0.1</sub>, thickness: 150 ± 10 nm.



address the role of (i) the interaction strength, (ii) halogen *versus* hydrogen bonding, and (iii) the ethynyl group in the efficiency of the SRG inscription. Such a comprehensive comparison is only valid when using a polymer host of low polydispersity (we used P4VP with  $M_w = 1200 \text{ g mol}^{-1}$ ;  $M_w/M_n = 1.2$ ) containing equal amounts of noncovalently attached azobenzene units (10% mol mol<sup>-1</sup> in this case), and films of equal thickness ( $150 \pm 10 \text{ nm}$ ).

The best way to study the efficiency of SRG formation is to monitor the first-order diffraction generated by periodic modulation of the sample surface due to interference irradiation. Several simultaneously formed and coupled gratings contribute to the overall diffraction efficiency,<sup>33</sup> but in thin amorphous azo-polymer films, by far the largest contribution arises from the SRG formation. Hence, we can relate the efficiency of the SRG formation to the growth dynamics of the diffracted signal. All samples provided high-quality sinusoidal gratings with a well-defined modulation depth ranging from 30 to 180 nm (Table 2).

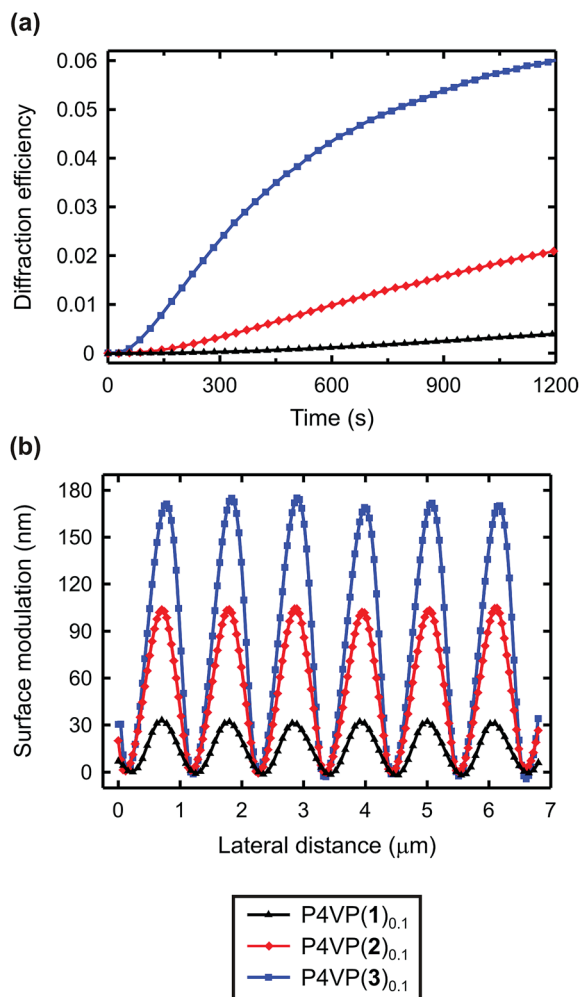


Fig. 4 (a) Kinetics of the first-order diffraction efficiency upon the grating inscription process, and (b) AFM surface profiles of the gratings for P4VP( $n$ )<sub>0.1</sub>, where  $n = 1$  (F is *para* to the azo group), 2 (Br is *para* to the azo group), and 3 (I is *para* to the azo group).

**The role of halogen bonding strength.** The kinetics of the diffraction efficiency of polymer-azobenzene complexes (in which the azobenzene unit is attached to every 10th polymer repeat unit) containing molecules 1–3 is given in Fig. 4a. A clear trend is seen: the non-halogen-bonded reference molecule 1 yields very weak diffraction, whereas the diffraction efficiency of the halogen-bonded complexes based on 2 and 3 develops in the order  $I > Br$ , *i.e.*, a stronger interaction significantly boosts the SRG formation efficiency. The same trend can be seen in the AFM surface profiles shown in Fig. 4b, where the surface modulation develops in the order  $I > Br > F$ . For the iodinated complex P4VP(3)<sub>0.1</sub> the modulation depth is *ca.* 175 nm (see Table 2), exceeding somewhat the thickness of the polymer film (150 nm). This result is in line with our earlier observations<sup>9</sup> and shows that (i) interaction between the photoactive azobenzene units and the passive polymer host is essential for the light-induced mass transport to take place, and (ii) if the azobenzenes are similar in terms of their spectroscopic and photochemical properties (as is the case of 2 and 3), the higher the interaction strength the more efficient the mass transport.

**Halogen bonding *versus* hydrogen bonding.** Molecules 3 and 5, together with their nonfluorinated counterparts 6 and 7, provide a comprehensive set for comparing the SRG formation efficiency in halogen- and hydrogen-bonded complexes, allowing us to assess the role of interaction directionality in the SRG formation. We recall that the interaction strength with pyridine moieties develops in the order 5 ( $11.765 \text{ kcal mol}^{-1}$ )  $>$  7 ( $10.052 \text{ kcal mol}^{-1}$ )  $>$  3 ( $5.135 \text{ kcal mol}^{-1}$ )  $>$  6 ( $2.546 \text{ kcal mol}^{-1}$ ), being significantly higher for the hydrogen-bonded complexes. The absorption maxima of 3 and 5 are practically identical as is the absorbance decrease upon irradiation (see Table 2). The latter relates to the efficiency of *trans-cis* photoisomerization, and to the amount of *cis*-isomers in the photostationary state. The *cis-trans* thermal relaxation is somewhat faster for P4VP(5)<sub>0.1</sub> (lifetime of *ca.* 4 min *vs.* *ca.* 10 min) but as the 488 nm irradiation that is used for SRG inscription triggers both *trans-cis* and *cis-trans* photoisomerization, we believe that the thermal isomerization has a negligible effect on the overall *trans-cis-trans* cycling rate. Therefore, we argue that the differences in thermal isomerization rate do not significantly affect the SRG formation and if they do have a small contribution, it will be in favour of 5.

The diffraction efficiency curves for the complexes between P4VP and 3, 5–7 are presented in Fig. 5. As expected, based on Fig. 4, for the halogen-bonded complexes the diffraction efficiency increases with increasing interaction strength; hence fluorination of the iodobenzene ring of 6 to afford 3 greatly enhances the SRG formation. In the systems based on azophenols 5 and 7, the diffraction kinetics is rather similar but the diffraction efficiency is systematically higher for the fluorinated one P4VP(5)<sub>0.1</sub>. The surface-modulation depths after 30 min inscription are *ca.* 170 nm and *ca.* 130 nm for P4VP(5)<sub>0.1</sub> and P4VP(7)<sub>0.1</sub>, respectively. The difference is likely due to slightly higher interaction strength and more confined electropositive surface of the former.

The most interesting comparison is that between P4VP(3)<sub>0.1</sub> and P4VP(5)<sub>0.1</sub>, that is, between similar halogen-bonded and



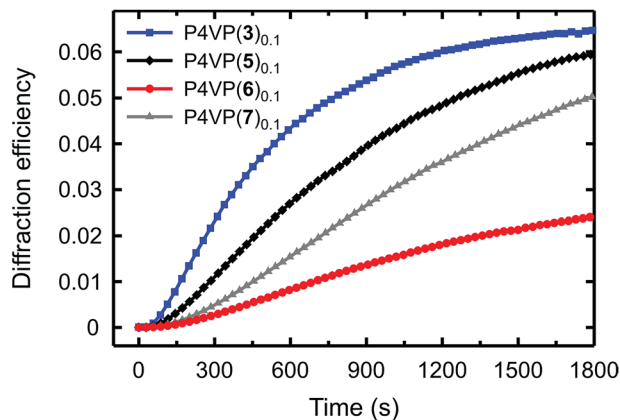


Fig. 5 Kinetics of the first-order diffraction efficiency upon the grating inscription process for P4VP( $n$ )<sub>0.1</sub>, where  $n = 3$  (I is *para* to the azo group), 5 (OH is in the *para* position), 6 (I is in the *para* position on a non-fluorinated ring), and 7 (OH is in the *para* position on a non-fluorinated ring).

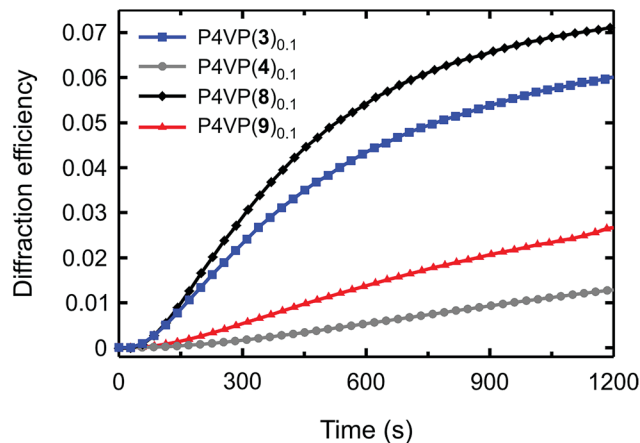


Fig. 6 Time evolution of the first-order diffraction efficiency upon the SRG inscription process for P4VP( $n$ )<sub>0.1</sub>, where  $n = 3$  (I is in the *para* position), 4 (H is in the *para* position), 8 (iodoacetylene is in the *para* position), and 9 (acetylene is in the *para* position).

hydrogen-bonded complexes. Even if the eventual diffraction efficiency and the modulation depth are practically the same, the dynamics of the diffraction growth is significantly faster for the halogen-bonded complex, *i.e.*, the SRG formation efficiency is higher for the iodinated dye despite its much weaker interaction with the polymer. As their chemical structures are identical apart from the bond-donor group, and the spectral and photochemical properties are comparable, this comparison unambiguously shows that, when sufficiently strong, halogen bonding surpasses hydrogen bonding in driving the light-induced mass transport in polymer–azobenzene complexes, due to its higher directionality. The order of the efficiency is reversed only when halogen bonding is significantly weaker than hydrogen bonding.

**The role of the ethynyl group.** Noncovalent interactions mediated by acetylene groups have not been previously used in light-induced surface patterning, and iodoacetylenes have only recently emerged as supramolecular synthons in crystal engineering and topochemical polymerization.<sup>19,34</sup> Hence it is of great interest to compare molecules 3 and 8, to address the question whether or not iodoacetylene is better than iodoperfluorobenzene in driving the SRG formation process. In this respect, it is also interesting to compare the SRG formation in complexes containing molecules 4 and 9. Their interaction energy with pyridine moieties is approximately the same (Table 1), but the molecular shape in the region of the bond donor moiety and the associated charge distribution are completely different, giving rise, in 9, to an electrostatic potential surface around the hydrogen atom that is very similar to that seen in the halogenated derivatives (Fig. 2). The positive area is narrowly confined along the extension of the C–H bond, closely resembling the  $\sigma$ -hole in halogens, surrounded by a belt of negative potential. Hence, this comparison allows us to study whether the more directional (weak) hydrogen bond leads to a more efficient mass transport.

The diffraction curves are shown in Fig. 6, and based on them we can draw two important conclusions. First, among the nine

supramolecular systems studied, the iodoethynyl-based P4VP(8)<sub>0.1</sub> complex is the most efficient for SRG formation. The difference with P4VP(3)<sub>0.1</sub> is not large, but it is systematic and repeatable. We believe that such enhancement results from more favourable photochemical properties, and possibly larger free-volume sweep upon isomerization, which may enhance the SRG formation.<sup>35</sup> Note also that the longer *cis*-lifetime does not seem to hinder the mass transport process as we suggested earlier. The fact that we can maintain efficient optical performance while removing the fluorine atoms from the bond-donor ring may be a great advantage when designing polymeric optical materials based on halogen bonding. The reduced quadrupolar interactions between the chromophores probably allow for using higher chromophore loadings without crystallization and phase separation of the dye, which may be beneficial in applications that require, *e.g.*, high optical anisotropy or nonlinear optical response.

The second important conclusion relates to the role of directionality of the polymer–azobenzene noncovalent interaction in the SRG formation. The P4VP(9)<sub>0.1</sub> complex performs worse than the halogen-bonded complexes (Fig. 6), but better than P4VP(4)<sub>0.1</sub>. This relative efficiency supports the data of Fig. 5 and suggests that more directional interactions of comparable strength lead to more efficient mass transport. We attribute this trend to formation of electrostatics-controlled rigid polymer–azobenzene junctions that “force” the passive polymer chains to follow the movements of the azobenzene molecules upon isomerization. The role of directionality explains the potential of halogen bonding in this application, and also calls for further investigation into the SRG formation in, *e.g.*, complexes driven by strong and directional multiple hydrogen bonds.<sup>36</sup>

## Experimental

Compounds 2, 3 and 4 were synthesized as previously reported.<sup>9</sup> Dye 7 was purchased from TCI Europe. Synthesis of azobenzenes 1, 5, 6, 8 and 9 is described in detail in the ESI.†



The molecular structures of dyes 1–9 were optimized in DMF within the DFT approach, using the PBE0 functional,<sup>37</sup> which has been judged to be well suited for describing the electronic and optical features of a series of organic dyes,<sup>38</sup> and treating the energetic and geometric features of halogen bonding, both *in vacuo*<sup>39</sup> and in a solvent.<sup>40</sup> The 6-311++G(d,p) basis set was used for all atoms. The solvation effects have been included by means of the Polarizable Continuum Model (PCM).<sup>41</sup> To determine the absorption wavelengths, standard vertical time dependent<sup>42</sup> PBE0/6-311++G(d,p) calculations have been carried out. The molecular dimers of compounds 1–9 with 4-methyl pyridine have been optimized at the PBE0/6-311++G\*\* level of theory *in vacuo*. Interaction energies  $\Delta E_{\text{BSSE}}$  have been computed by optimization on the BSSE-free potential energy surface as the difference between the energy of the dimer and the sum of the energies of the single monomers. BSSE correction was made by the standard counterpoise method.<sup>43</sup> All calculations have been performed with the Gaussian suite of programs.<sup>44</sup>

The single-crystal X-ray structures were determined using a Bruker Kappa Apex II CCD diffractometer with Mo K $\alpha$  radiation ( $\lambda = 0.71073 \text{ \AA}$ ) and a Bruker Kryoflex low-temperature device. Crystals were mounted in inert oil on glass fibers. Data collection and reduction were performed by SMART and SAINT and absorption correction, based on the multi-scan procedure, by SADABS. The structures were solved by SIR92 and refined on all independent reflections by full-matrix least-squares based on  $F_o^2$  by using SHELX-97. Crystallographic data are reported in Table S1.† CIFs containing full crystallographic data can be obtained free of charge from the ESI.†

Surface-relief gratings were inscribed on thin films ( $150 \pm 10 \text{ nm}$  as measured with a DEKTAK 6M surface profiler) spin-cast from freshly prepared, filtered DMF solutions of the azobenzene–polymer mixtures on quartz substrates. The inscription was performed using a spatially filtered and collimated Ar<sup>+</sup> laser beam (Coherent Innova 70) with circular polarization at a wavelength of 488 nm with an irradiation intensity of  $150 \text{ mW cm}^{-2}$ . Lloyd's mirror interferometer, with an incidence angle of  $15^\circ$  (spatial period of  $1 \mu\text{m}$ ) was used to create the light-interference pattern. The dynamics of the gratings was monitored by measuring the transmitted first-order diffracted beam from a normally incident 633 nm He–Ne laser. The diffraction efficiency of the gratings was defined as  $\eta = I_1/I_0$ , where  $I_1$  and  $I_0$  are the intensities of the first-order diffracted beam and the transmitted beam prior to irradiation, respectively. The surface profiles of the inscribed gratings were characterized using a Veeco Dimension 5000 atomic-force microscope.

The UV-Vis spectra were collected from both thin films and dilute ( $10^{-5} \text{ M}$ ) DMF solutions with an Ocean Optics USB2000+ fiber-optic spectrometer and a DH-2000-BAL light source, both in the dark and under irradiation ( $488 \text{ nm}$ ,  $50 \text{ mW cm}^{-2}$ ). Thermal *cis-trans* isomerization was studied by exciting the chromophores to the *cis*-state with a circularly polarized pump beam ( $488 \text{ nm}$ ,  $50 \text{ mW cm}^{-2}$ ) and monitoring the transmittance changes after blocking the pump. The probe was a fiber-coupled xenon lamp equipped with proper bandpass filters. The signal was detected with a photodiode and a lock-in amplifier.

## Conclusions

We have investigated the formation of photoinduced surface patterns in halogen-bonded and hydrogen-bonded polymer–azobenzene supramolecular complexes. Using an extensive molecular library we were able to draw several important conclusions on the light-induced mass transport process. First of all, halogen-bonded complexes outperform the corresponding hydrogen-bonded complexes (*cf.* Fig. 1, compounds 3 and 5) in terms of surface patterning efficiency, as a result of the higher directionality of halogen bonding. Greater directionality promotes more efficient mass transport also when hydrogen bonding is responsible for the dye–polymer binding. Furthermore, within the halogen-bonded series (*cf.* molecules 1, 2, 3, and 6) the efficiency increases with the interaction strength. The most efficient halogen-bond donor motif that comes on top of the supramolecular hierarchy in driving SRG formation is the iodoacetylene dye 8. This efficiency is due to its favourable photochemical properties as compared to the tetra-fluoroiodobenzene motif. We believe this motif to widen the future perspectives of the design of halogen-bonded optical and photoresponsive materials by allowing for high chromophore loadings to be used in polymer matrices. Research in this direction is currently ongoing in our laboratories and will be reported elsewhere.

## Acknowledgements

PM acknowledges the MIUR for the PRIN 2010–2011 InfoChem project. GR acknowledges the MIUR for the PRIN 2010–2011 project no. 2010ERFKXL. AF acknowledges the CINECA-LISA Award N. HPL13P3CCB, 2013, for the availability of high-performance computing resources and support. AGH wishes to acknowledge FQRNT for a B3 post-doctoral fellowship. JV wishes to acknowledge Emil Aaltonen Foundation for a post-doctoral grant. AP acknowledges the Politecnico di Milano International Fellowship Program, the Academy of Finland, and the Emil Aaltonen Foundation for financial support.

## Notes and references

- M. M. Russew and S. Hecht, *Adv. Mater.*, 2010, **22**, 3348; H. Yu and T. Ikeda, *Adv. Mater.*, 2011, **23**, 2149; Z. Mahimwalla, K. G. Yager, J. Mamiya, A. Shishido, A. Priimagi and C. J. Barrett, *Polym. Bull.*, 2012, **69**, 967; L. De Sio, N. Tabiryan, T. Bunning, B. R. Kimball and C. Umetsu, *Prog. Opt.*, 2013, **58**, 1; T. Seki, S. Nagano and M. Hara, *Polymer*, 2013, **54**, 6053.
- M. H. Dhammika Bandara and S. C. Burdette, *Chem. Soc. Rev.*, 2012, **41**, 1809.
- T. Ikeda, *J. Mater. Chem.*, 2003, **13**, 2037; A. Priimagi, C. J. Barrett and A. Shishido, *J. Mater. Chem. C*, 2014, **2**, 7155.
- N. K. Viswanathan, D. Y. Kim, S. Bian, J. Williams, W. Liu, L. Li, L. Samuelson, J. Kumar and S. K. Tripathy, *J. Mater. Chem.*, 1999, **9**, 1941; A. Natansohn and P. Rochon, *Chem. Rev.*, 2002, **102**, 4139.



- 5 S. I. Na, S. S. Kim, J. Jo, S. H. Oh, J. Kim and D. Y. Kim, *Adv. Funct. Mater.*, 2008, **18**, 3956; S. Lee, H. S. Kang and J. K. Park, *Adv. Mater.*, 2012, **24**, 2069; L. M. Goldenberg, V. Lisinetskii, Y. Gritsai, J. Stumpe and S. Schrader, *Adv. Mater.*, 2012, **24**, 3339; A. Priimagi and A. Shevchenko, *J. Polym. Sci., Part B: Polym. Phys.*, 2014, **52**, 163.
- 6 (a) O. Kulikovska, L. M. Goldenberg and J. Stumpe, *Chem. Mater.*, 2007, **19**, 3343; (b) J. Gao, Y. N. He, F. Liu, X. Zhang, Q. Z. Wang and X. G. Wang, *Chem. Mater.*, 2007, **19**, 3877; (c) N. Zettsu, T. Ogasawara, N. Mizoshita, S. Nagano and T. Seki, *Adv. Mater.*, 2008, **20**, 516; (d) Q. Zhang, X. Wang, C. J. Barrett and C. G. Bazuin, *Chem. Mater.*, 2009, **21**, 3216; (e) A. Priimagi, K. Lindfors, M. Kaivola and P. Rochon, *ACS Appl. Mater. Interfaces*, 2009, **1**, 1183.
- 7 T. Alasaarela, D. Zheng, L. Huang, A. Priimagi, B. Bai, A. Tervonen, S. Honkanen, M. Kuittinen and J. Turunen, *Opt. Lett.*, 2011, **36**, 2411.
- 8 G. R. Desiraju, P. S. Ho, L. Kloo, A. C. Legon, R. Marquardt, P. Metrangolo, P. Politzer, G. Resnati and K. Rissanen, *Pure Appl. Chem.*, 2013, **85**, 1711.
- 9 A. Priimagi, G. Cavallo, A. Forni, M. Gorynsztejn-Leben, M. Kaivola, P. Metrangolo, R. Milani, A. Shishido, T. Pilati, G. Resnati and G. Terraneo, *Adv. Funct. Mater.*, 2012, **22**, 2572.
- 10 T. Clark, M. Henneman, J. S. Murray and P. Politzer, *J. Mol. Model.*, 2007, **13**, 291; T. Clark, *Wiley Interdiscip. Rev.: Comput. Mol. Sci.*, 2013, **3**, 13.
- 11 K. E. Riley and P. Hobza, *Phys. Chem. Chem. Phys.*, 2013, **15**, 17742; A. Forni, S. Pieraccini, S. Rendine and M. Sironi, *J. Comput. Chem.*, 2014, **35**, 386.
- 12 A. Priimagi, G. Cavallo, P. Metrangolo and G. Resnati, *Acc. Chem. Res.*, 2013, **46**, 2686.
- 13 X. Ding, M. Tuikka and M. Haukka, Halogen Bonding in Crystal Engineering, in *Recent Advances in Crystallography*, ed. J. B. Benedict, InTech, New York, 2012, pp. 143–168; M. Saccone, G. Cavallo, P. Metrangolo, A. Pace, I. Pibiri, T. Pilati, G. Resnati and G. Terraneo, *CrystEngComm*, 2013, **15**, 3102; R. W. Troff, T. Mäkelä, F. Topić, A. Valkonen, K. Raatikainen and K. Rissanen, *Eur. J. Org. Chem.*, 2013, **9**, 1617.
- 14 D. Yan, A. Delori, G. O. Lloyd, T. Friščić, G. M. Dayet, W. Jones, J. Lu, M. Wei, D. G. Evans and X. Duanal, *Angew. Chem., Int. Ed.*, 2011, **50**, 12483; O. Bolton, K. Lee, H. J. Kim, K. Y. Lin and J. Kim, *Nat. Chem.*, 2011, **3**, 205.
- 15 D. W. Bruce, in *Supramolecular Chemistry: From concept to Nanomaterials*, ed. Steed J. W. and Gale P. A., John Wiley & Sons, Ltd, 2012, p. 3493.
- 16 A. Priimagi, M. Saccone, G. Cavallo, A. Shishido, T. Pilati, P. Metrangolo and G. Resnati, *Adv. Mater.*, 2012, **24**, OP345; M. Saccone, G. Terraneo, T. Pilati, G. Cavallo, A. Priimagi, P. Metrangolo and G. Resnati, *Acta Crystallogr., Sect. B: Struct. Sci., Cryst. Eng. Mater.*, 2014, **70**, 149; O. S. Bushuyev, T. C. Corkery, C. J. Barrett and T. Friščić, *Chem. Sci.*, 2014, **5**, 3158; Y. Chen, H. Yu, L. Zhang, H. Yang and Y. Lu, *Chem. Commun.*, 2014, **50**, 9647.
- 17 J. P. W. Wong, A. C. Whitwood and D. W. Bruce, *Chem.–Eur. J.*, 2012, **18**, 16073; A. Takemura, L. J. McAllister, S. Hart, N. E. Pridmore, P. B. Karadakov, A. C. Whitwood and D. W. Bruce, *Chem.–Eur. J.*, 2014, **20**, 6721.
- 18 T. Steiner, *Angew. Chem., Int. Ed.*, 2002, **41**, 48.
- 19 C. B. Aakeröy, M. Baldrighi, J. Desper, P. Metrangolo and G. Resnati, *Chem.–Eur. J.*, 2013, **19**, 16240.
- 20 O. Ikkala and G. ten Brinke, *Chem. Commun.*, 2004, 2131.
- 21 R. Bertani, P. Metrangolo, A. Moiana, E. Perez, T. Pilati, G. Resnati, I. Rico-Lattes and A. Sassi, *Adv. Mater.*, 2002, **14**, 1197.
- 22 S. Xiao, X. Lu and Q. Lu, *Macromolecules*, 2007, **40**, 7944; A. Priimagi, J. Vapaavuori, F. J. Rodriguez, C. F. J. Faul, M. T. Heino, O. Ikkala, M. Kauranen and M. Kaivola, *Chem. Mater.*, 2008, **20**, 6358; Q. Zhang, C. G. Bazuin and C. J. Barrett, *Chem. Mater.*, 2008, **20**, 29; J. Vapaavuori, V. Valtavirta, T. Alasaarela, J. Mamiya, A. Priimagi, A. Shishido and M. Kaivola, *J. Mater. Chem.*, 2011, **21**, 15437.
- 23 G. A. Jeffrey, *An Introduction to Hydrogen Bond*, Oxford University Press, New York, 1997.
- 24 K. E. Riley, J. S. Murray, J. Fanfrlik, J. Rezac, R. J. Sola, M. C. Concha, F. M. Ramos and P. Politzer, *J. Mol. Model.*, 2011, **17**, 3309.
- 25 A. C. Legon, *Angew. Chem., Int. Ed.*, 1999, **38**, 2686; Z. P. Shields, J. S. Murray and P. Politzer, *Int. J. Quantum Chem.*, 2010, **110**, 2823; P. Politzer and J. S. Murray, *ChemPhysChem*, 2013, **14**, 278.
- 26 (a) I. Majerz, Z. Malarski and L. Sobczyk, *Chem. Phys. Lett.*, 1997, **274**, 361; (b) T. Steiner, I. Majerz and C. C. Wilson, *Angew. Chem., Int. Ed.*, 2001, **40**, 2651.
- 27 J. H. Williams, J. K. Cockcroft and N. Fitch, *Angew. Chem., Int. Ed.*, 1992, **31**, 1655.
- 28 E. A. Meyer, R. K. Castellano and F. Diederich, *Angew. Chem., Int. Ed.*, 2003, **42**, 1210.
- 29 M. H. Yoon, A. Facchetti, E. S. Charlotte and T. J. Marks, *J. Am. Chem. Soc.*, 2006, **128**, 5792.
- 30 K. Kishikawa, *Isr. J. Chem.*, 2012, **52**, 800.
- 31 C. S. Paik and H. Morawetz, *Macromolecules*, 1972, **5**, 171.
- 32 C. Barrett, A. Natansohn and P. Rochon, *Macromolecules*, 1994, **27**, 4781; C. Barrett, A. Natansohn and P. Rochon, *Chem. Mater.*, 1995, **7**, 899.
- 33 F. Lagugné Labarthe, P. Rochon and A. Natansohn, *Appl. Phys. Lett.*, 1999, **75**, 1377A; A. Sobolewska and S. Bartkiewicz, *Appl. Phys. Lett.*, 2008, **92**, 253305; H. Audorff, R. Walker, L. Kador and H. W. Schmidt, *J. Phys. Chem. B*, 2009, **113**, 3379; A. Sobolewska, S. Bartkiewicz, A. Miniewicz and E. Schab-Balcerzak, *J. Phys. Chem. B*, 2010, **114**, 9751.
- 34 J. W. Lauher, F. W. Fowler and N. S. Goroff, *Acc. Chem. Res.*, 2008, **41**, 1215; C. Perkins, S. Libri, H. Adams and L. Brammer, *CrystEngComm*, 2012, **14**, 3033; O. Dumele, D. Wu, N. Trapp, N. Goroff and F. Diederich, *Org. Lett.*, 2014, **16**, 4722.
- 35 A. Goulet-Hanssens, T. C. Corkery, A. Priimagi and C. J. Barrett, *J. Mater. Chem. C*, 2014, **2**, 7505.



- 36 L. J. Prins, D. N. Reinhoudt and P. Timmerman, *Angew. Chem., Int. Ed.*, 2001, **40**, 2382; J. V. Barth, *Annu. Rev. Phys. Chem.*, 2007, **58**, 375.
- 37 M. Ernzerhof and G. E. Scuseria, *J. Chem. Phys.*, 1999, **110**, 5029; C. Adamo and V. Barone, *J. Chem. Phys.*, 1999, **110**, 6158.
- 38 D. Jacquemin, V. Wathelet, E. A. Perpète and C. Adamo, *J. Chem. Theory Comput.*, 2009, **5**, 2420; D. Jacquemin, A. Planchat, C. Adamo and B. Mennucci, *J. Chem. Theory Comput.*, 2012, **8**, 2359.
- 39 S. Kozuch and J. M. L. Martin, *J. Chem. Theory Comput.*, 2013, **9**, 1918.
- 40 A. Forni, S. Rendine, S. Pieraccini and M. Sironi, *J. Mol. Graphics Modell.*, 2012, **38**, 31.
- 41 M. Cossi and V. Barone, *J. Chem. Phys.*, 2001, **115**, 4708.
- 42 E. Runge and E. K. U. Gross, *Phys. Rev. Lett.*, 1984, **52**, 997; R. E. Stratmann, G. E. Scuseria and M. J. Frisch, *J. Chem. Phys.*, 1998, **109**, 8218; M. E. Casida, *J. Mol. Struct.: THEOCHEM*, 2009, **914**, 3.
- 43 S. F. Boys and F. Bernardi, *Mol. Phys.*, 1970, **19**, 553.
- 44 M. Frisch, *et al.*, *Gaussian 09, Revision D.01.*, Gaussian Inc., Wallingford CT, 2013.

