

Cite this: *J. Mater. Chem. A*, 2015, 3, 10949

# Nickel oxide interlayer films from nickel formate–ethylenediamine precursor: influence of annealing on thin film properties and photovoltaic device performance†

K. X. Steirer,<sup>\*acd</sup> R. E. Richards,<sup>a</sup> A. K. Sigdel,<sup>‡bc</sup> A. Garcia,<sup>§c</sup> P. F. Ndione,<sup>c</sup> S. Hammond,<sup>c</sup> D. Baker,<sup>¶d</sup> E. L. Ratcliff,<sup>e</sup> C. Curtis,<sup>c</sup> T. Furtak,<sup>d</sup> D. S. Ginley,<sup>c</sup> D. C. Olson,<sup>c</sup> N. R. Armstrong<sup>a</sup> and J. J. Berry<sup>\*c</sup>

An organometallic ink based on the nickel formate–ethylenediamine ( $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$ ) complex forms high performance  $\text{NiO}_x$  thin film hole transport layers (HTL) in organic photovoltaic (OPV) devices. Improved understanding of these HTLs functionality can be gained from temperature-dependent decomposition/oxidation chemistries during film formation and corresponding chemical structure–function relationships for energetics, charge selectivity, and transport in photovoltaic platforms. Investigations of as-cast films annealed in air (at 150 °C–350 °C), with and without subsequent  $\text{O}_2$ -plasma treatment, were performed using thermogravimetric analysis, Fourier transform infrared spectroscopy, ultraviolet and X-ray photoelectron spectroscopy, and spectroscopic ellipsometry to elucidate the decomposition and oxidation of the complex to  $\text{NiO}_x$ . Regardless of the anneal temperature, after exposure to  $\text{O}_2$ -plasma, these HTLs exhibit work functions greater than the ionization potential of a prototype donor polymer poly(*N*-9'-heptadecanyl-2,7-carbazole-*alt*-5,5-(4',7'-di-2-thienyl-2',1',3'-benzothiadiazole) (PCDTBT), thereby meeting a primary requirement of energy level alignment. Thus, bulk-heterojunction (BHJ), OPV solar cells made on this series of  $\text{NiO}_x$  HTLs all exhibit similar open circuit voltages ( $V_{\text{oc}}$ ). In contrast, the short circuit currents increase significantly from 1.7 to 11.2  $\text{mA cm}^{-2}$  upon increasing the anneal temperature from 150 °C to 250 °C. Concomitantly, increased conductivity and electrical homogeneity of  $\text{NiO}_x$  thin films are observed at the nanoscale using conductive tip-AFM. Similar  $V_{\text{oc}}$  observed for all the  $\text{O}_2$ -plasma treated  $\text{NiO}_x$  interlayers and variations to nanoscale conductivity suggest that the HTLs all form charge selective contacts and that their carrier extraction efficiency is determined by the amount of precursor conversion to  $\text{NiO}_x$ . The separation of these two properties: selectivity and conductivity, sheds further light on charge selective interlayer functionality.

Received 21st February 2015  
Accepted 15th April 2015

DOI: 10.1039/c5ta01379h

www.rsc.org/MaterialsA

## Introduction

High efficiency bulk heterojunction (BHJ) organic photovoltaic (OPV) devices often require contacts modified with hole or

electron charge-transport interlayers in order to increase the charge carrier collection efficiency above that of the unmodified transparent conducting oxide or metal contact.<sup>1,2</sup> The efficiency of charge collection interlayers relies upon their thin film conductivity,<sup>3</sup> work function ( $\Phi_w$ ),<sup>4–7</sup> alignments of interlayer valence and conduction band edges with the highest occupied molecular orbital (HOMO) and lowest unoccupied molecular orbital (LUMO) energies across the heterojunction,<sup>8–12</sup> as well as the degree of heterogeneity for the contact surface.<sup>13</sup> The combination of these mechanisms provide for selective charge collection in competition with bimolecular and surface recombination under low internal electric fields (*i.e.* near open-circuit conditions).<sup>14–16</sup> In particular, the importance of selectivity and contact-extraction efficiency becomes increasingly important in solution-processed photovoltaic platforms as free carrier mobilities of photo-active layers increase. Metal oxides formed with Mo, V, W, Ni or NiCo<sup>6</sup> have been shown to exhibit favorable attributes as hole transport layers (HTLs). Photovoltaic

<sup>a</sup>Department of Chemistry and Biochemistry, University of Arizona, Tucson, Arizona, USA

<sup>b</sup>Department of Physics and Astronomy, University of Denver, Denver, Colorado, USA

<sup>c</sup>National Renewable Energy Laboratory, Golden, Colorado, USA. E-mail: Kenneth.Steirer@NREL.GOV; Joseph.Berry@NREL.GOV

<sup>d</sup>Applied Physics Department, Colorado School of Mines, Golden, Colorado, USA

<sup>e</sup>Department of Materials Science and Engineering, University of Arizona, Tucson, Arizona, USA

† Electronic supplementary information (ESI) available: XPS of N 1s, UPS and additional FTIR for  $\text{NiO}_x$  films. See DOI: 10.1039/c5ta01379h

‡ Present address: Intel Corporation, Portland, OR, USA.

§ Present address: Next Energy Technologies Inc. Santa Barbara, CA, USA.

¶ Present address: Washington Nanofabrication Facility at the University of Washington, Seattle WA, USA.



applications in organic-,<sup>17–22</sup> colloidal quantum structure-<sup>23,24</sup> and methyl ammonium lead halide-based platforms<sup>25–27</sup> utilize HTLs with high transmission at operational wavelengths and work function ( $\Phi_w$ ) values equal to or in excess of the donor IP value, although presumably due to different mechanisms for n- and p-type oxides. The grand challenge in contact design remains to effectively create solution-processed deposition methods for high performing thin film devices while still maintaining material and interfacial functionalities outlined above.

NiO<sub>x</sub> is one of few p-type metal oxides that has transversed numerous energy relevant technologies such as catalysis, batteries, fuel cells and photovoltaics. Hence, it is of fundamental interest and several organometallic precursor formulations compatible with solution processing have been identified for thin film formation. Examples of these are: nickel acetate tetrahydrate complexed with methanolamine (275 °C),<sup>28</sup> nickel nitrate hexahydrate with monoethanolamine (500 °C)<sup>29</sup> and nickel formate dihydrate with ethylenediamine (250 °C).<sup>30</sup> Lowering the processing temperature required to convert these precursors to the oxide allows use of plastic substrates, which in general cannot tolerate prolonged processing above 150 °C.<sup>31</sup> There is considerable literature precedent for the decomposition of nickel formate to form Ni and NiO.<sup>32–38</sup> Diamine complexation with nickel formate lowers the thermal requirement for decomposition and thus enables formation of NiO<sub>x</sub> at lower temperature. Solutions made with the complexed organometallic precursor in ethylene glycol and water allow fabrication of NiO<sub>x</sub> thin films by spin coating the nickel formate-ethylenediamine-ethylene glycol-water (Ni(O<sub>2</sub>CH)<sub>2</sub>-en-eg-water) ink followed by thermal annealing in air. Formation of NiO<sub>x</sub> by this method is unique as it produces conformal, high performance thin films with few processing steps.

A detailed understanding of the interconnected decomposition chemistry with the material and interface functionality can drive metal oxide ink development beyond empirical approaches. For example, exposure to reactive oxygen during annealing may further reduce thermal post-treatments. Zhai *et al.* demonstrated this relaxation in processing conditions for the acetate precursor, below 150 °C.<sup>39</sup> As a direct result of the film growth and processing, NiO<sub>x</sub> interlayers strongly affect the OPV device performance.<sup>28,30,40</sup> After annealing in air and treating with an O<sub>2</sub>-plasma, NiO<sub>x</sub> outperforms a benchmark HTL of poly(ethylene dioxythiophene):poly(styrene sulfonate) (PEDOT:PSS) in prototypical OPV devices using the BHJ poly[*N*-9'-heptadecanyl-2,7-carbazole-*alt*-5,5-(4',7'-di-2-thienyl-2',1',3'-benzothiadiazole)] (PCDTBT):[6,6]-phenyl-C71 butyric acid methyl ester (PC<sub>70</sub>BM).<sup>20</sup> When NiO<sub>x</sub> interlayers are included in OPV devices, the surface chemistry, band edge energies and mid-gap defect states determine the surface electrical properties and charge selectivity towards holes. Detailed spectroscopic analyses of these solution-deposited NiO<sub>x</sub> thin films have shown that these are complex NiO<sub>x</sub> surfaces, with a wide range of possible oxide stoichiometries that influence their optoelectronic properties, and their interactions with semiconductors such as those found in organic and hybrid photovoltaic platforms.<sup>8,25–27,41</sup> Previous UPS and XPS measurements on these

films correlated surface hydroxyl species and their dipolar character with an increased band gap energy and improved band edge alignment with BHJ films.<sup>8,41</sup> More specifically, the NiO<sub>x</sub> surface formed from decomposition of these solution precursors is comprised predominantly of a mixture of NiO<sub>x</sub>, Ni(OH)<sub>2</sub> and NiOOH, as revealed by XPS characterization.<sup>41</sup> The dipolar character of this modified surface leads to a high  $\Phi_w$  and favorable energetic matching to the highest occupied molecular orbital (HOMO<sub>D</sub>) hole-transport energy level of PCDTBT, while the wide band gap, and an apparent lack of mid-gap states, functions to block reverse electron transfer from the lowest unoccupied molecular orbital (LUMO<sub>A</sub>) of the fullerene electron acceptor.<sup>8,20,41</sup> Furthermore, as these processing conditions for the NiO<sub>x</sub> interlayers led to variations in the measured local density of states observed in UPS, this resulted in higher hole selectivity and lower leakage currents in hole only devices.<sup>41</sup> Through improved charge selectivity and limiting carrier injection from the contact, these NiO<sub>x</sub> interlayers lower leakage current and increase shunt resistance in OPV devices.<sup>14,42</sup> However, systematic investigation of precursor decomposition in relation to device performance has yet to be addressed and hence, is the focus of this paper.

Here, we study the effects of varying the annealing temperature between 150 °C and 350 °C for thin films spin coated from the Ni(O<sub>2</sub>CH)<sub>2</sub>-en-eg-water formulation. The effects of incomplete precursor decomposition are important to understanding their influence on the interlayer optoelectronic properties and the ability to collect photocurrent in OPV devices. We observe changes to both the chemical and electronic properties of the resulting NiO<sub>x</sub> thin films that correlate with large changes in short-circuit photocurrent ( $J_{sc}$ ) and little to no changes in open-circuit photovoltage ( $V_{oc}$ ) in PCDTBT:PC<sub>70</sub>BM OPV devices. Decomposition/oxidation reactions for the films were investigated by thermal gravimetric analysis (TGA), differential scanning calorimetry (DSC), Fourier transform infrared absorption spectroscopy (FTIR) and X-ray photoelectron spectroscopy (XPS). By increasing the anneal temperature for thin films spin-coated from the Ni(O<sub>2</sub>CH)<sub>2</sub>-en-eg-water ink, from 150 °C to above 250 °C, and subjecting the films to an O<sub>2</sub>-plasma, amorphous thin films are formed with: (i) increased conductivity as measured by conductive AFM; (ii) increased surface oxygen content (O/Ni ratio revealed by XPS); (iii) increase of the NiO<sub>x</sub> band gap; (iv) high performance in OPV devices, as revealed by analysis of their series resistance and  $J_{sc}$ .  $V_{oc}$  is shown to be affected primarily by the surface oxidation chemistry of NiOOH even if the precursor decomposition/oxidation is incomplete, while losses observed in  $J_{sc}$  depend primarily upon the nano-scale conductivity threshold reached upon decomposition of the Ni-formate-diamine complex. These studies decouple the underlying oxide formation from the surface effects by O<sub>2</sub>-plasma treating for photovoltaic device applications.

## Experimental

### Ink preparation

Preparation of the Ni(O<sub>2</sub>CH)<sub>2</sub>-en-eg precursor formulation for NiO<sub>x</sub> films has been reported earlier.<sup>30</sup> To summarize, nickel



formate (1 g) was combined with ethylene glycol (10 ml) followed by ethylenediamine (0.87 ml). The mixture was heated and shaken multiple times, and then filtered at near ambient temperature through a 0.45  $\mu\text{m}$  pore filter. The ink was a deep purple color, consistent with the violet color reported for  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$ .<sup>43</sup> For spin coating, the ink was diluted 1 : 1 by volume with water (nanopure).

### Thin films & devices

Patterned ITO substrates were first solvent cleaned in acetone and isopropyl alcohol followed by an  $\text{O}_2$ -plasma treatment.  $\text{NiO}_x$  films were deposited by spin coating the  $\text{Ni}(\text{O}_2\text{CH})_2$ -en-eg-water ink at 4000 rpm onto the ITO substrates and immediately annealing the films in air at 150–400  $^\circ\text{C}$  for one hour. After annealing, all  $\text{NiO}_x$  layers were exposed to  $\text{O}_2$ -plasma treatment for 2 minutes at 155 W and 800 mTorr. The 1 : 4 ratio PCDTBT:PC<sub>70</sub>BM solution was prepared in 1,2-dichlorobenzene under an inert atmosphere at a total concentration of 35  $\text{mg ml}^{-1}$ . The solution was stirred at 90  $^\circ\text{C}$  for 8 hours before cooling to 60  $^\circ\text{C}$  followed by immediate use, which is a variation on a previously reported procedure.<sup>20</sup> Spin coated active layers were deposited on top of the  $\text{NiO}_x$  HTL films at a spin rate of 2000 rpm for 120 seconds. The coated substrates were annealed at 70  $^\circ\text{C}$  on a hot plate for one hour. Top electrodes composed of Ca/Al (20 nm/100 nm) were thermally evaporated using an Angstrom Engineering thermal evaporator with a base pressure below  $1 \times 10^{-7}$  Torr to produce 0.11  $\text{cm}^2$  devices. Films of  $\text{NiO}_x$  were prepared on freshly  $\text{O}_2$ -plasma cleaned Au substrates for AFM and C-AFM studies.

### AFM & C-AFM

Scanning probe measurements employed an Asylum Research MFP-3D Atomic Force Microscope in conductive mode (c-AFM) using a Pt/Cr coated conductive tip (ElectricMulti75-G by Budget Sensors Inc.) with a radius less than 25 nm. Both height topography and c-AFM were obtained simultaneously. To obtain a good c-AFM signal,  $\text{NiO}_x$  films were deposited on top of Au-coated glass substrates rather than ITO, since ITO has non-uniform conductive regions<sup>44</sup> and thermally-induced electrical degradation at small length scales.<sup>13</sup> The Au substrate used for calibration had a highly uniform c-AFM profile at very small sample-to-tip bias (VST). All c-AFM measurements on  $\text{NiO}_x$  films were at ambient conditions with identical scan parameters such as scan speed and drive amplitude, using a VST of 40 mV.

### Photoelectron spectroscopy

Experiments were performed on a Kratos Axis Ultra X-ray photoelectron spectrometer equipped with a monochromatic Al K-alpha X-ray source ( $h\nu$  1486.6 eV) and a He UV source ( $h\nu$  21.22 eV). Linear calibration of the binding energy scale for the detector was performed following the procedure outlined by M. P. Seah.<sup>45</sup> A bias of  $-10.00$  V was applied to the sample during UPS experiments to spectrally separate the lowest kinetic energy electrons and secondary electrons from the local environment. An Ar sputter-etched, atomically-clean gold sample was measured before characterization of the  $\text{NiO}_x$  samples to establish the Fermi edge of the spectrometer.

### TGA & DSC

The ink (nickel formate-ethylenediamine-ethylene glycol-water) was placed in a Pt pan at 120  $^\circ\text{C}$  to evaporate the bulk of the water and ethylene glycol solvent with minimal disruption of the nickel formate-ethylenediamine complex. This procedure was repeated twice to obtain an initial mass of 13.87 mg of the  $\text{NiO}_x$  film precursor (less ethylene glycol). The pan temperature was increased at 10  $^\circ\text{C min}^{-1}$  under dry synthetic air (20%  $\text{O}_2$ , 80%  $\text{N}_2$ ) in a TA Instruments SDT Q600 operated in TGA/DSC mode.

### FTIR

Transmission spectra were measured using a Thermo-Nicolet 6700 FTIR. A liquid  $\text{N}_2$ -cooled mercury-cadmium-telluride (MCT) detector and a KBr beamsplitter were used. Scans were collected to provide data with a resolution of 2  $\text{cm}^{-1}$ . For each measurement 100 scans were averaged for both the sample and the background. Absorbance spectra were calculated from the sample and background using the Beer-Lambert equation.

### Ellipsometry

Thin film optical properties and thicknesses ( $9.5 \pm 0.5$  nm) were characterized using an M-2000 spectroscopic ellipsometer (J.A. Woolam Co Inc.) at wavelengths of 250–1000 nm and angle of 65–75 degrees. Spectroscopic ellipsometry data were processed with the aid of WVASE software. The  $\text{NiO}_x$  complex refractive index constants ( $N = n - ik$ ) were obtained using a Lorentz parameterized model, which is consistent with the Kramers-Kronig relations. Thicknesses were verified using a Dektak 8 stylus profilometer.

## Results and discussion

### 1. Decomposition and oxidation of $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$

Decomposition and oxidation processes for the metal organic precursor depends on the nickel complex formed in solution and spin cast into films. We follow the conversion to  $\text{NiO}_x$  with TGA and DSC shown in Fig. 1. Previous literature isothermal studies have reported the complete decomposition/oxidation of  $\text{Ni}(\text{O}_2\text{CH})_2 \cdot 2\text{H}_2\text{O}$  (at 240  $^\circ\text{C}$ –280  $^\circ\text{C}$  in air)<sup>38</sup> and  $\text{Ni}(\text{O}_2\text{CH})_2$  (at 215–250  $^\circ\text{C}$  in oxygen and 242–262  $^\circ\text{C}$  in air).<sup>35,36</sup> For example,  $\text{Ni}(\text{O}_2\text{CH})_2 \cdot 2\text{H}_2\text{O}$  was converted to NiO in less than an hour at 240  $^\circ\text{C}$ –280  $^\circ\text{C}$  in air.<sup>38</sup> This is consistent with the expectation that the films annealed in this study at 250  $^\circ\text{C}$  or 300  $^\circ\text{C}$  in air should be comprised of  $\text{NiO}_x$ , which is also supported by the decomposition/oxidation temperature for the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  complex observed by TGA/DSC in Fig. 1. Narain reported the synthesis of  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$ , which should be robust at 120  $^\circ\text{C}$  without any loss of the ethylenediamine ligand.<sup>43</sup> Therefore, the DSC/TGA measurements (Fig. 1) were performed in a Pt pan on an ink sample after two rounds of evaporation of the bulk of the ethylene glycol and water solvents at 120  $^\circ\text{C}$ , leaving a sufficient mass of primarily the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  complex. On heating in synthetic air (80%  $\text{N}_2$ , 20%  $\text{O}_2$ ), a small endothermic peak is evident near 125  $^\circ\text{C}$ , with a corresponding mass loss of ca. 5%, which is interpreted as a loss of residual solvent (water, ethylene



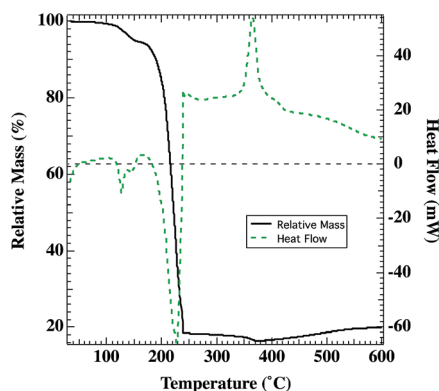


Fig. 1 Thermal gravimetric analysis of the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  precursor after evaporation of ethylene glycol and water (black). Differential scanning calorimetry data for the same thermal scan (green-dotted).

glycol, excess ethylenediamine). Between 180–240 °C, an additional 77% of the initial mass (81% of mass at 150 °C) is lost in an endothermic process, which is comparable to the expected 72% mass change for conversion of  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  to  $\text{NiO}$ . Mass loss is due to a combination of evaporation of residual solvent (likely ethylene glycol, boiling point 197 °C), loss of the ethylenediamine ligands, which have been shown to leave stepwise,<sup>46–48</sup> and the decomposition/oxidation of the  $\text{Ni}(\text{O}_2\text{CH})_2$  complex (above 225 °C) to  $\text{NiO}_x$ .<sup>34,36,38</sup> In comparison, studies with  $[\text{Ni}(\text{saccharinate})_2(\text{en})_2]\text{en}$ <sup>47</sup>  $[\text{Ni}(\text{en})_3]\text{oxalate}$ ,<sup>48</sup>  $\text{Ni}(\text{succinate})(\text{en})$ ,<sup>49</sup> and  $\text{Ni}(\text{en})_3\text{SO}_4$  (ref. 50) reported temperatures for the final stage decomposition/oxidation to  $\text{NiO}$  of 300 °C, 325 °C, 410 °C, and 466 °C respectively. On further heating, there is some additional mass lost (ca. 2% of the initial mass) until 375 °C, when an exothermic transition is seen and correlates with a gradual mass increase. Hence, the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  complex is to date, preferred for lower temperature formation of solution-deposited  $\text{NiO}_x$  but photovoltaic applications require subsequent surface treatments to increase the work-function.<sup>7,30</sup>

## 2. Chemical composition characterization

Surface elemental composition determined from XPS measurements confirmed predominant  $\text{NiO}_x$  composition as well as the presence of C and N due to the  $\text{O}_2$ -plasma treatment and ambient exposure prior to measurement (see Table 1).  $\text{O}_2$ -plasma is known to remove adventitious organic compounds. However, the N content in the films likely originates from the ethylenediamine, which decreases as the anneal temperature increases from 150 °C to 250 °C. Yet, after correcting for the C 1s binding energy (BE) the BE values for N 1s peak centroids are located at 406.6 eV (see S1†) and are too high to be explained by the presence of unreacted amine groups. A more likely assignment consistent with the  $\text{O}_2$ -plasma treatment is one or more forms of near surface N–O species such as  $-\text{NO}_3$ , which can be identified using vibrational spectroscopy.

Exposure of the  $\text{NiO}_x$  thin films to  $\text{O}_2$ -plasma predominantly affects the exposed surface creating similar structures for all the films while providing less effect on the subsurface

Table 1 Surface atomic composition of spin-coated ( $\text{NiO}_x$ ) films calculated from XPS peaks normalized to photoemission sensitivity factors: Ni 2p, C 1s, O 1s and N 1s. Work function ( $\Phi_w$ ) values were calculated from UPS data. All films were  $\text{O}_2$ -plasma treated. Satellites for Ni 2p photoelectrons were omitted from area calculations. Error is  $\pm 10\%$

Anneal $T$ (°C)	Ni (%)	O (%)	C (%)	N (%)	O/Ni	$\Phi_w$ (eV)
150	33	55	8	4	1.7	5.4
200	29	56	12	3	2.0	5.4
250	27	60	11	2	2.2	5.5
300	26	62	10	2	2.4	5.5

material. We analyzed FTIR spectra taken for  $\text{NiO}_x$  thin films to understand the  $\text{O}_2$ -plasma effects for the whole system. FTIR spectra are shown in Fig. 2 for films spin-coated from the  $\text{Ni}(\text{O}_2\text{CH})_2\text{-en-eg-water}$  ink, comparing the as-deposited film ('no anneal') to films annealed for one hour in air at 150 °C, 200 °C, 250 °C or 300 °C. Band assignments for chemical constituents of the film precursor are listed in Table 2. The major band assignments reported in the literature for the fundamental vibrations of the formate group in nickel formate (dihydrate) are the  $\nu_1 \nu(\text{CH})$  mode at ca. 2900  $\text{cm}^{-1}$ , the intense  $\nu_4 \nu_{\text{as}}(\text{COO})$  mode at ca. 1570  $\text{cm}^{-1}$ , and asymmetric deformation ( $\nu_5 \delta(\text{C-H})$ ) and symmetric stretch ( $\nu_2 \nu_{\text{s}}(\text{COO})$ ) modes between 1400–1350  $\text{cm}^{-1}$ .<sup>34,51–53</sup> For liquid ethylene glycol, the major band assignments reported in the literature are the strong  $\nu(\text{OH})$  stretching mode at 3400–3150

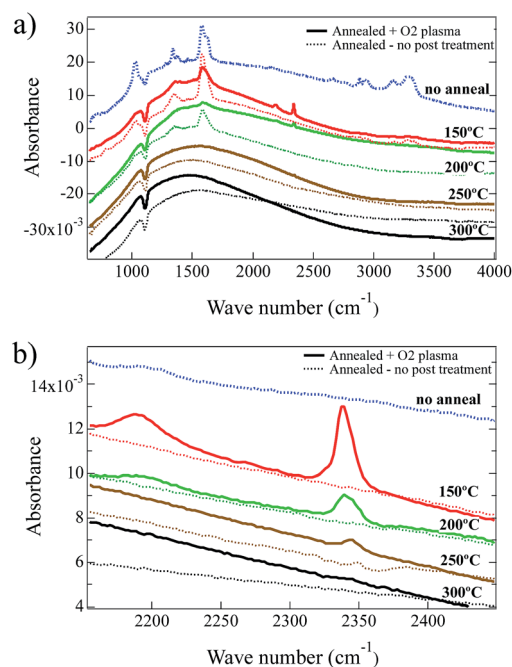


Fig. 2 FTIR results for spin-coated ( $\text{NiO}_x$ ) thin films on Si wafers, comparing without anneal and annealed conditions (150 °C, 200 °C, 250 °C and 300 °C). Films that received  $\text{O}_2$ -plasma treatment are indicated. Spectra are offset in order to clearly show the spectral changes.



Table 2 Peak assignments for FTIR spectra shown in Fig. 2

Wavenumber (cm <sup>-1</sup> )	Assignment	Source of vibration	Reference
1018	$\nu(\text{CN})$	Ethylenediamine (ligand)	1–3
1027	$\nu(\text{CO})$	Ethylene glycol	4 and 5
1338	$\omega(\text{NH}_2)$	Ethylenediamine (ligand)	2 and 3
1371	$\nu_5 \delta(\text{CH}) \nu_2 \nu_s(\text{COO})$	Formate	6–9
1458	$\delta(\text{CH}_2)$	Ethylene glycol	4 and 5
1587	$\nu_4 \nu_{\text{as}}(\text{COO})$	Formate	6–9
1627	$\delta(\text{NH}_2)$	Ethylenediamine (ligand)	1–3
2190	$\nu(\text{CN})$	Cyanate/isocyanate	10–13
2340	$\nu_3 \nu_{\text{as}}(\text{CO}_2)$	CO <sub>2</sub>	14–16
2884	$\nu_1 \nu(\text{C-H})$	Formate	6–9
	$\nu_s(\text{CH})$	Ethylene glycol	4 and 5
2936	$\nu_{\text{as}}(\text{CH})$	Ethylene glycol	4 and 5
	$\nu(\text{CH}_2)$	Ethylenediamine	1 and 2
3300–3160	$\nu(\text{OH})$	Ethylene glycol	4 and 5
3313	$\nu_s(\text{NH}), \nu_{\text{as}}(\text{NH})$	Ethylenediamine (ligand)	1–3 and 17
3570	$\nu(\text{OH})$	surface hydroxyls, Ni(OH) <sub>2</sub>	18–20

cm<sup>-1</sup>, the strong asymmetric ( $\nu_{\text{as}}(\text{CH})$ ) and symmetric ( $\nu_s(\text{CH})$ ) stretch modes at 2935 cm<sup>-1</sup> and 2875 cm<sup>-1</sup> respectively, the strong  $\delta(\text{CH}_2)$  mode at *ca.* 1450 cm<sup>-1</sup>, and the very strong  $\nu(\text{CO})$  and  $\nu(\text{CC})$  modes at 1100–1050 cm<sup>-1</sup>.<sup>54,55</sup> IR spectra have not been reported for the complex Ni(O<sub>2</sub>CH)<sub>2</sub>(en)<sub>2</sub>. However, literature reports indicate the major band assignments for ethylenediamine ligands in Ni(acetylacetonates)<sub>2</sub>en<sup>56</sup> and [Ni(succinate)(en)]·2H<sub>2</sub>O<sup>49</sup> complexes to be the asymmetric ( $\nu_{\text{as}}(\text{NH})$ ) and symmetric ( $\nu_s(\text{NH})$ ) modes at 3350–3150 cm<sup>-1</sup>, the strong  $\delta(\text{NH}_2)$  mode at 1613 cm<sup>-1</sup>, the  $\omega(\text{NH}_2)$  mode at 1318 cm<sup>-1</sup>, and the strong  $\nu(\text{CN})$  mode at 1025 cm<sup>-1</sup>. These ethylenediamine ligand band assignments for Ni complexes are consistent with values reported for other transition metal organometallic complexes<sup>56–58</sup> and also liquid ethylenediamine.<sup>59</sup>

The bands at *ca.* 1020–1030 cm<sup>-1</sup> ( $\nu(\text{CN})$  and  $\nu(\text{CO})$ ) and *ca.* 3200–3350 cm<sup>-1</sup> ( $\nu(\text{OH})$ ) indicate ethylene glycol and/or ethylenediamine,<sup>54,55,57,59,60</sup> and are discernible only in the no-anneal film and the film annealed at 150 °C as shown in Fig. 2a. The FTIR spectra indicate that ethylene glycol and ethylenediamine are virtually eliminated by a one hour anneal in air at 200 °C. Bands at 1627 cm<sup>-1</sup> ( $\delta(\text{NH}_2)$ ), and 1338 cm<sup>-1</sup> ( $\nu_5 \delta(\text{C-H})$ ) and ( $\nu_2 \nu_s(\text{COO})$ ) and 1587 cm<sup>-1</sup> ( $\nu_4 \nu_{\text{as}}(\text{COO})$ ), also seen in Fig. 2a, indicate ethylenediamine and formate respectively.<sup>34,51–53,57,59,60</sup> The formate bands are present in the FTIR spectra for the as spun film and the films annealed at 150 °C or 200 °C without an O<sub>2</sub>-plasma treatment. A one hour anneal in air at 250 °C and 300 °C eliminates the formate from the films resulting in a near featureless spectra consistent with NiO<sub>x</sub> except for broad bands at *ca.* 3570 cm<sup>-1</sup> that are surface hydroxyls.<sup>61–63</sup> This result is consistent with the TGA/DSC data described above.

A comparison of the impact of the O<sub>2</sub>-plasma treatment, typically used for NiO HTLs, is also included in Fig. 2a and b for films annealed for one hour in air at 150 °C, 200 °C, 250 °C or 300 °C. The intensities of all the ethylenediamine, ethylene glycol, and formate bands were lowered significantly after treatment with O<sub>2</sub>-plasma, consistent with an O<sub>2</sub>-plasma being highly efficient at removing organic compounds from materials

and surfaces. After O<sub>2</sub>-plasma treatment, two new bands emerge in the FTIR spectra located at 2190 cm<sup>-1</sup> and 2340 cm<sup>-1</sup>. The 2190 cm<sup>-1</sup> band is present after O<sub>2</sub>-plasma treatment in the film annealed at 150 °C (*i.e.*, before the ethylenediamine is eliminated), and is very weak in the film annealed at 200 °C. The 2340 cm<sup>-1</sup> band is present in the films annealed at 150 °C, 200 °C and 250 °C after O<sub>2</sub>-plasma treatment, although the intensity decreases significantly with increasing temperature consistent with greater conversion decomposition/oxidation of the precursor to NiO<sub>x</sub>. Given the oxidizing environment in the O<sub>2</sub>-plasma and the presence of C and N in the partially decomposed/oxidized films annealed at 150 °C and 200 °C, the 2190 cm<sup>-1</sup> band is tentatively assigned to the  $\nu(\text{C-N})$  modes for oxygen-bonded cyanate (OCN) groups or nitrogen-bonded isocyanate (NCO) groups to Ni<sup>2+</sup>; rather than the stretch modes of C=N in a carbon nitride film.<sup>64</sup> For example,  $\nu(\text{CN})$  modes have been reported at *ca.* 2200 cm<sup>-1</sup> for the nickel isocyanate complex [Et<sub>4</sub>N]<sub>2</sub>[Ni(NCO)<sub>4</sub>],<sup>65</sup> CNO<sup>-</sup> intercalated in  $\alpha$ -Ni(OH)<sub>2</sub> (ref. 66) and a theoretical study of the adsorption of cyanate and isocyanate on a Ni(100) surface;<sup>67</sup> and at 2262 cm<sup>-1</sup> and 2200 cm<sup>-1</sup> for Ni(NCO)<sub>2</sub>·H<sub>2</sub>O,<sup>68</sup> In contrast, a theoretical study of the adsorption of cyanide on a Ni(100) surface reported the  $\nu(\text{CN})$  mode at only *ca.* 2000 cm<sup>-1</sup>,<sup>69</sup> and experimentally the  $\nu(\text{CN})$  mode for Ni(CN)<sub>2</sub>·2H<sub>2</sub>O was reported at 2172 cm<sup>-1</sup>.<sup>70</sup> The 2340 cm<sup>-1</sup> band is tentatively assigned to the  $\nu_3 \nu_{\text{as}}(\text{CO}_2)$  mode of CO<sub>2</sub> trapped in the films annealed at 150 °C–200 °C after O<sub>2</sub>-plasma treatment of the partially decomposed/oxidized Ni(O<sub>2</sub>CH)<sub>2</sub>(en)<sub>2</sub> complex. Similar IR bands have been reported for free CO<sub>2</sub> trapped during the thermal decomposition in air of hexahydrated nickel iron citrate to form ultrafine NiFe<sub>2</sub>O<sub>4</sub> particles (2320 cm<sup>-1</sup>),<sup>71</sup> propanol/TaCl<sub>5</sub> gel to form Ta<sub>2</sub>O<sub>5</sub> thin films (at 2345 cm<sup>-1</sup> and 2333 cm<sup>-1</sup>),<sup>72</sup> and zinc acetate dihydrate/sodium hydrogen carbonate mixtures in argon to form ZnO nanoparticles (at *ca.* 2340 cm<sup>-1</sup>).<sup>73</sup> Both of these modes appear to be eliminated after O<sub>2</sub>-plasma treatment for the film annealed at 300 °C since decomposition/oxidation of the precursor to NiO<sub>x</sub> is complete. However, as described above a small percentage of N is still observed in the XPS spectra with high BE values for



O<sub>2</sub>-plasma treated films, most likely due to near surface N–O species such as –NO<sub>3</sub>. For nickel nitrate the strong –NO<sub>3</sub> ion stretching band would appear at *ca.* 1300–1400 cm<sup>-1</sup>.<sup>74–76</sup> This is also the region where medium strength formate and ethylenediamine bands are anticipated. Confirmation of nitrate cannot be provided by the FTIR spectra after 150 °C anneal and O<sub>2</sub>-plasma treatment. However, the peak position of the ethylenediamine ω(NH<sub>2</sub>) mode shifts from 1338 cm<sup>-1</sup> to 1358 cm<sup>-1</sup> for the as-deposited, and 150 °C anneal plus O<sub>2</sub>-plasma respectively, and may indicate possible spectral contribution from –NO<sub>3</sub> ions (see S2†). After the 200 °C anneal and O<sub>2</sub>-plasma treatment, the 1300–1400 cm<sup>-1</sup> region is nearly featureless.

FTIR analysis suggests the trapping of CO<sub>2</sub> and the formation of N-based anions in the films annealed at the lower temperatures and after an O<sub>2</sub>-plasma treatment. The confinement of CO<sub>2</sub> in solid-state NiO<sub>x</sub> films implies that a dense surface barrier forms during the O<sub>2</sub>-plasma treatment. The cyanate species assigned in the FTIR spectra were not identified in the more surface-sensitive XPS measurements. Moreover, nitrates observed by XPS could not be unambiguously identified with FTIR. These complementary surface and through-film measurements lead to the tentative hypothesis that the low concentration of nitrates are most likely surface confined.

### 3. Characterization using spectroscopic ellipsometry

In order to study the effect of annealing on the optical properties we employed angular dependent UV-vis reflection using spectroscopic ellipsometry. Fig. 3 displays the square of absorbance *versus* incident photon energy. Absorbance was calculated using the extinction coefficient from Lorentz oscillator model fits to the ellipsometric data. The onset of absorption was extrapolated after a linear fit to the square of the absorption coefficient between 4.1 and 4.4 eV.<sup>77</sup> For the lowest temperature anneal at 150 °C the NiO<sub>x</sub> absorbance is low and barely resembles the absorption edge of a semiconductor, which is

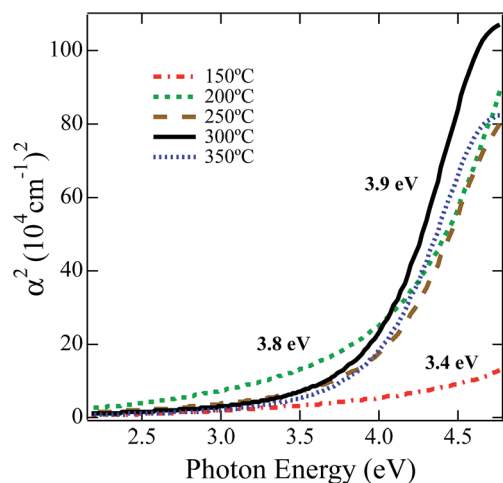


Fig. 3 Square of absorption coefficient plot calculated with *n* and *k* measured from spectroscopic ellipsometry *versus* photon energy. Linear fits were obtained between 4.1 and 4.4 eV for (NiO<sub>x</sub>) films at each anneal temperature.

consistent with the minimal decomposition/oxidation of the Ni(O<sub>2</sub>CH)<sub>2</sub>(en)<sub>2</sub> complex at this temperature. Nevertheless, fitting the onset region results in an estimate for the optical gap of 3.4 eV. For the film annealed at 200 °C, the absorbance increases as the precursor has partially decomposed/oxidized resulting in a 3.8 eV estimate of the optical band gap. For films annealed at 250 °C and 300 °C, because the decomposition of the Ni(O<sub>2</sub>CH)<sub>2</sub>(en)<sub>2</sub> complex to NiO<sub>x</sub> is well advanced the absorption edge is more clearly defined and results in an optical gap estimate of 3.9 eV, which is in good agreement with the accepted band gap of 4.0 eV for NiO.<sup>78,79</sup>

### 4. OPV device performance

To investigate the impact of the NiO<sub>x</sub> film composition as a function of annealing temperature, on OPV performance we utilized these films as HTLs, annealed at different temperatures on ITO and integrated them into ITO/NiO<sub>x</sub>/PCDTBT:PC<sub>70</sub>BM/Ca/Al OPV devices. Current–voltage (*JV*) measurements were performed under one-sun illumination. The data from these devices are shown in Fig. 4a with calculated performance metrics in Tables 3 and 4.

Power conversion efficiencies (PCE) are shown normalized in Fig. 4b and increased from 0.5% to 5.7% with increasing anneal temperature between 150 °C and 250 °C. These PCE values trend directly with short-circuit current density and as a function of thermal annealing temperature. Likewise, the PCE and *J*<sub>sc</sub> inversely trend with *R*<sub>s</sub> as a function of annealing temperature. For the lowest annealing temperatures (150 °C and 200 °C) the devices suffer from large resistive losses, poor current extraction and low fill factors. At 250 °C and above the *R*<sub>s</sub> drops substantially and the device performance improves with gains in *J*<sub>sc</sub> and FF. This drop in series resistance within the device is commensurate with the decomposition/oxidation of the NiO<sub>x</sub> layer.

It is important to note that the open-circuit voltages do not appear to trend with annealing temperature. Ultraviolet photoelectron spectroscopy (UPS) measurements (see S3†) for these solution-deposited NiO<sub>x</sub> films after annealing between 150 °C and 300 °C all produce films with very similar work function values that range from 5.4 to 5.5 eV and IP values of 5.7–5.8 ± 0.1 eV, in agreement with earlier reports.<sup>15</sup> This is consistent with the relatively uniform *V*<sub>oc</sub> found across the devices when one considers work-function and the interface electronic structure of the contact determining factors of *V*<sub>oc</sub>. For the devices annealed at 250 °C and above there is very little statistically significant difference in the device data. As shown in Table 3, there is a modest increase just above the statistical noise from 250 °C to 300 °C. It is clear that the compositional changes from annealing the NiO<sub>x</sub> HTL to 250 °C significantly alter the electronic properties and result in enhanced hole-collection from the BHJ.

### 5. Conductive tip-AFM

To examine the charge transport across the NiO<sub>x</sub> films processed at different temperatures and to better understand hole-collecting efficiency, samples were prepared (in an identical way



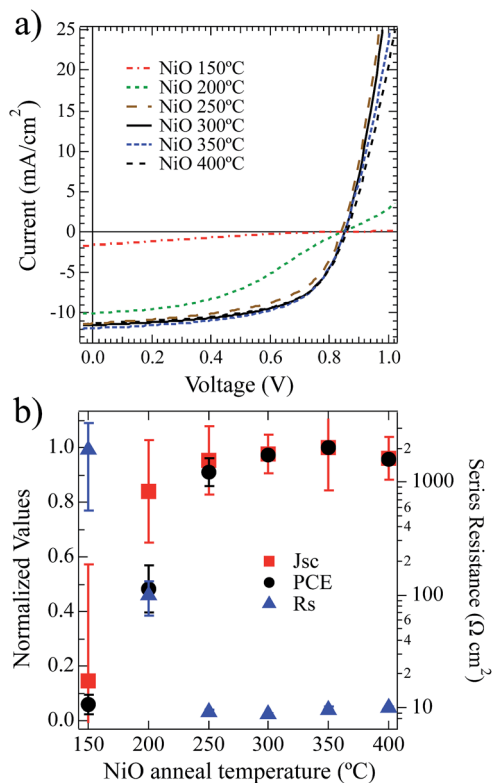


Fig. 4 (a) Current–voltage response of PCDTBT:PC<sub>70</sub>BM solar cells with increasing anneal temperatures for the (NiO<sub>x</sub>) films spin-coated from the Ni(O<sub>2</sub>CH)<sub>2</sub>-en-eg-water solution. (b) Short circuit current ( $J_{sc}$  – squares) and power conversion efficiency (PCE – circles) normalized and plotted for the current–voltage curves in (a).  $R_s$  taken as the inverse slope at +1.5 V is plotted on the right axis. Error bars represent one standard deviation.

to the films used in devices) on Au as opposed to ITO substrates. Gold was used to eliminate underlying effects of electrical heterogeneity of the ITO, compared to the uniform and stable background provided on the gold surface. The nanoscale topography and electrical transport of these films were measured with conductive tip-AFM using a Pt/Ir tip held at the same tip bias for all films. For these experiments the tip forms the top contact on the Au-NiO<sub>x</sub>-Pt/Ir tip junction. Fig. 5 shows AFM and c-AFM data for five different anneal temperatures between 150 °C and 350 °C. The white pixels in the bottom row of Fig. 5 indicate a good electrical transport or an Ohmic junction, while dark pixels indicate poor electrical conductivity

Table 4 Normalized device metrics from Table 3

Anneal temp	$V_{oc}$ norm	$J_{sc}$ norm	FF norm	PCE norm	$R_s$ norm
150	0.98	0.14	0.38	0.06	1.0000
200	0.97	0.84	0.59	0.48	0.0504
250	0.99	0.96	0.96	0.91	0.0047
300	0.99	0.98	1.00	0.97	0.0045
350	0.99	1.00	0.99	1.00	0.0049
400	1.00	0.96	1.00	0.96	0.0051

or diode like behavior. Since all the NiO<sub>x</sub> films have similar work functions as seen in Table 1, we can assume minimal changes in tip-surface injection/extraction barriers. Hence, we conclude that the measured c-AFM data is indicative of changes in the conductivity of the NiO<sub>x</sub>. Conformal, thin-film coatings were observed for each annealing condition on the Au substrates. Moreover, XRD analysis of a series of NiO<sub>x</sub> thin-films heated from 150 °C to 350 °C showed no sign of diffraction peaks. The 150 °C anneal resulted in an undulating surface topology (1.46 nm RMS) and a highly insulating film with no current transfer between tip and the Au ground. In contrast, the film annealed at 200 °C has a much flatter surface morphology (0.46 nm RMS) and a small but measurable current, which indicates improved charge transfer with <10% area being electrically active. As the annealing temperature is increased from 200 °C to 350 °C, increases in the NiO<sub>x</sub> film roughness and electrical property are observed. After the 250 °C anneal, the film roughness is 0.65 nm RMS and conductive regions occupy most of the surface. The 300 °C anneal results in a NiO<sub>x</sub> film with larger protrusions (1.35 nm RMS) and a sharp increase in the measured current that saturated the 20 nA c-AFM detection limit. The NiO<sub>x</sub> films resulting from a 350 °C anneal exhibits slightly lower conductivity compared to the 300 °C film, and a further increase in particle/grain size. The improved through-film conductivity, as assessed by the area of the saturated pixels, correlate well with lowering of the series resistance and improved current collection for the OPV devices. This suggests that high series resistance for the lower anneal temperatures is a result of the poor conductivity of the HTL, which is a direct result of incomplete precursor decomposition/oxidation.

The large volume fraction of incompletely decomposed precursor results in resistive properties at the nanoscale and macroscale. However, separating improvements to carrier concentration and mobility remains elusive. One would expect

Table 3 Device performance metrics for current–voltage responses reported in Fig. 4.  $R_s$  corresponds to the inverse slope at 1.5 V

Anneal temp (°C)	$V_{oc}$ (mV)	±	$J_{sc}$ (mA cm <sup>-2</sup> )	±	FF	±	PCE (%)	±	$R_s$ – dark (Ω cm <sup>2</sup> )	±
150	842	6	1.7	0.7	0.23	0.04	0.3	0.2	1955.1	1397.9
200	834	11	9.8	0.3	0.35	0.05	2.8	0.5	98.6	33.7
250	850	14	11.2	0.2	0.57	0.04	5.3	0.3	9.3	0.4
300	852	10	11.5	0.1	0.59	0.01	5.6	0.1	8.7	0.2
350	852	7	11.7	0.3	0.58	0.00	5.7	0.1	9.7	0.6
400	860	8	11.2	0.1	0.59	0.00	5.5	0.1	10.1	0.3



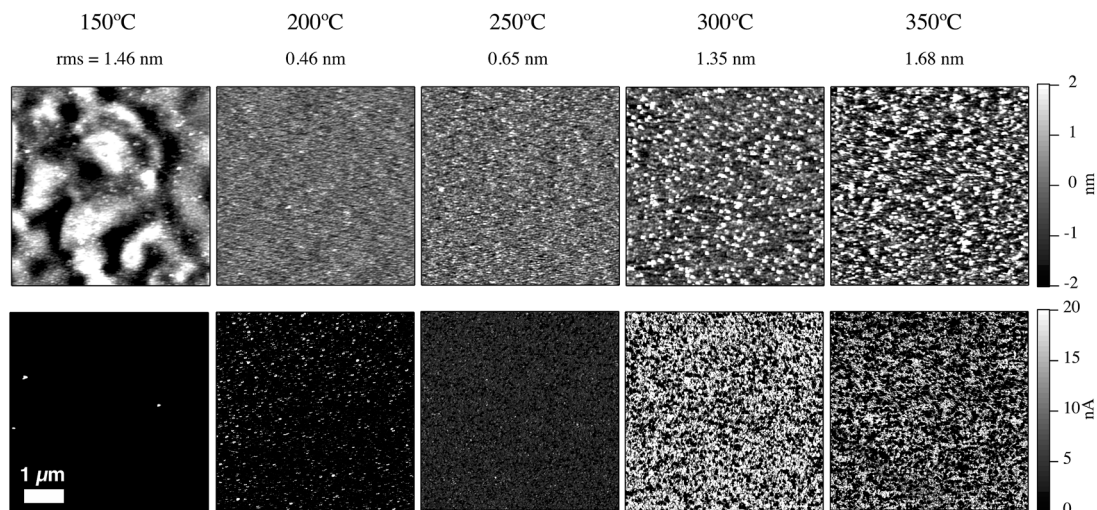


Fig. 5 AFM (top row) and cAFM (bottom row) of  $\text{NiO}_x$  films deposited on Au. From left-to-right, anneals at 150 °C, 200 °C, 250 °C, 300 °C and 350 °C for the spin-coated films. RMS roughness is stated below the corresponding anneal temperature.

that the improved conductivity with annealing temperature occurs due to increases of both mobility and carrier concentrations as the precursor decomposes and formation of NiO proceeds. Note that the formation of new species (cyanate, isocyanate, or nitrate) after exposing the low temperature (150 °C–200 °C) annealed films to  $\text{O}_2$ -plasma, may also contribute to the poor semiconductor properties of these films. However, solution-deposited  $\text{NiO}_x$  HTLs annealed at 250 °C–400 °C, and hence above the decomposition/oxidation temperature of the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  complex, exhibit compositional changes that improve the nanoscale conductivity. These  $\text{NiO}_x$  thin films can then effectively collect photogenerated holes from the BHJ with minimal resistive loss. The sub-stoichiometric nature of the  $\text{NiO}_x$  produced at low annealing temperatures suggests that in addition to establishing a high work-function, the  $\text{O}_2$ -plasma treatment required for optimal device performance oxidizes the  $\text{NiO}_x$  but to a large extent is limited to the surface of the HTL. If the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  is not completely decomposed (anneal temperature <240 °C), then the  $\text{O}_2$ -plasma treatment does not significantly improve the NiO stoichiometry and electrical transport, and instead affects only the near surface of the thin film.

## Conclusions

Thermal annealing in air and subsequent  $\text{O}_2$ -plasma treatment of solution-deposited  $\text{NiO}_x$  films spin-coated from a  $\text{Ni}(\text{O}_2\text{CH})_2$ -en-eg-water ink have been investigated to gain an understanding of their composition and properties in relation to hole collection in PCDTBT:PC<sub>70</sub>BM solar cells. The impacts of incomplete precursor decomposition on OPV device performance are investigated with the main effect arising as increased series resistance and concomitant lowering of short circuit current. For annealing temperatures of 150 °C–200 °C, below the decomposition temperature of the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  complex, BHJ solar cells incorporating these 'NiO<sub>x</sub>' HTLs have low  $J_{sc}$  and high  $R_s$ . TGA identified a lower limit of 240 °C anneal

temperature that is required to ensure adequate decomposition/oxidation of the  $\text{Ni}(\text{O}_2\text{CH})_2(\text{en})_2$  complex. The resulting formation of  $\text{NiO}_x$  HTLs with  $\sim 3.9$  eV optical gap and adequate conductivity permits efficient hole collection in BHJ devices. TGA/DSC analysis suggests such films are initially sub-stoichiometric  $\text{NiO}_x$ , which helps to explain the improved performance for subsequent  $\text{O}_2$ -plasma treatment to obtain a high work function and improved film stoichiometry. Results are consistent with recent reports of other  $\text{NiO}_x$  precursor formulations that produce  $\text{NiO}_x$  HTLs with similar device properties at low decomposition temperature.<sup>28</sup> These studies show increased series resistance as a direct result of incomplete decomposition of the  $\text{NiO}_x$  precursor complex and how these resistive HTL films lead to lower  $J_{sc}$  in devices with concomitant performance loss.

However, for interlayers where conversion to  $\text{NiO}_x$  is performed below the decomposition temperature and is incomplete,  $V_{oc}$  remains at high values. The surface chemistry and work function of the  $\text{NiO}_x$  interlayer determines the  $V_{oc}$  in OPV devices whether or not the organometallic precursor has fully decomposed to form  $\text{NiO}_x$ . In comparison, non-selective self-assembled molecular interlayer contacts provide paths to electron transfer from the BHJ LUMO levels and reduce quasi- $E_F$  splitting, which lead to lower  $V_{oc}$ .<sup>7</sup> From the results presented here, we conclude that the nanoscale electrical changes as a function of converted precursor observed seem to strongly affect the ability of charge selective  $\text{NiO}_x$  interlayers to extract holes from the adjacent BHJ and transport those to the external circuit. Recently, the surface polarity of  $\text{NiO}_x$  interlayers was investigated and shown to dominate the interface properties when compared to the interlayer surface roughness and crystal structure.<sup>40</sup> Hence, post treatment and formation of a dipolar surface with low defects is related to the increased polar component of the total surface energy. Data presented here shows that surface composition for these films are similar. However, differences in their nanoscale conductivity do not strongly affect the  $V_{oc}$ . Hence, surface recombination velocity is



not significantly enhanced as quasi- $E_F$  splitting seems nominally equivalent for these devices at  $V_{oc}$ , which is consistent with steady-state and transient photocurrent studies on similar systems.<sup>42</sup> We hypothesize that the majority of  $J_{sc}$  loss observed as the processing temperature is lowered below the precursor decomposition threshold proceeds *via* recombination in the BHJ and is not mediated by  $NiO_x$  surface states. If this postulate holds, then charge selectivity and efficient carrier transport are functionally separate and proceed by different mechanisms for this particular active layer. Moreover, these properties are also spatially separate as the selectivity is determined by the surface composition and local density of states that provide a low defect interface and low surface recombination while the interlayer subsurface enables charge delocalization and carrier transport to the transparent electrode. Implications for the separation of selectivity and transport mechanisms could result in designs for bilayer selective contacts and indeed examples exist in literature.<sup>80</sup> This can also help to decouple surface and subsurface effects of decomposition temperature, organometallic precursor formulations and subsequent surface modifications for efficient interlayer contacts in photovoltaic technologies. However in more demanding photovoltaic systems with higher carrier mobilities and photogenerated charge densities, it may be necessary to increase the  $NiO_x$  thickness in order to effectively passivate high carrier density electrodes such as TCOs and metals.

## Acknowledgements

Thank you Professor Reuben T. Collins at the Colorado School of Mines for thoughtful discussions. Research supported as part of the Center for Interface Science: Solar Electric Materials, an Energy Frontier Research Center funded the U.S. Department of Energy, Office of Science, Basic Energy Sciences, under Award Number DE-SC0001084 (NRA, ELR, RER, KXS), and NREL LDRD Number ARE21000.

## Notes and references

- E. L. Ratcliff, B. Zacher and N. R. Armstrong, *J. Phys. Chem. Lett.*, 2011, **2**, 1337–1350.
- C. Brabec, S. Gowrisanker, J. Halls, D. Laird, S. Jia and S. Williams, *Adv. Mater.*, 2010, **22**, 3839–3856.
- J. Chen, H. Wei and K. Ho, *Sol. Energy Mater. Sol. Cells*, 2007, **91**, 1472–1477.
- H. Frohne, S. Shaheen, C. J. Brabec, D. Muller, N. Sariciftci and K. Meerholz, *ChemPhysChem*, 2002, **9**, 795–799.
- S. Khodabakhsh, B. Sanderson, J. Nelson and T. Jones, *Adv. Funct. Mater.*, 2006, **16**, 95–100.
- P. Ndione, A. Garcia, E. Widjonarko, A. K. Sigdel, K. X. Steirer, D. C. Olson, P. A. Parilla, D. S. Ginley, N. R. Armstrong, R. E. Richards, E. L. Ratcliff and J. J. Berry, *Adv. Energy Mater.*, 2013, **3**, 524–531.
- E. L. Ratcliff, A. Garcia, S. A. Paniagua, S. R. Cowan, A. J. Giordano, D. S. Ginley, S. R. Marder, J. J. Berry and D. C. Olson, *Adv. Energy Mater.*, 2013, **3**, 647–656.
- E. L. Ratcliff, J. Meyer, K. Steirer, N. R. Armstrong, D. C. Olson and A. Kahn, *Org. Electron.*, 2012, **13**, 744–749.
- J. Meyer and A. Kahn, *Proc. SPIE*, 2011, **1**, 1–7.
- J. Meyer, A. K. Zilberberg, T. Riedl and A. Kahn, *J. Appl. Phys.*, 2011, **110**, 033710.
- J. H. Seo, S. Cho, M. Leclerc and A. J. Heeger, *Chem. Phys. Lett.*, 2011, **503**, 101–104.
- H. Ishii, K. Sugiyama, E. Ito and K. Seki, *Adv. Mater.*, 1999, **11**, 605–625.
- N. R. Armstrong, P. A. Veneman, E. L. Ratcliff, D. Placencia and M. Brumbach, *Acc. Chem. Res.*, 2009, **42**, 1748–1757.
- T. Ripolles-Sanchis, A. Guerrero, E. Azaceta, R. Tena-Zaera and G. Garcia-Belmonte, *Sol. Energy Mater. Sol. Cells*, 2013, **117**, 564–568.
- J. Reinhardt, M. Grein, C. Bühler, M. Schubert and U. Würfel, *Adv. Energy Mater.*, 2014, DOI: 10.1002/aenm.201400081.
- M. Bolognesi, A. Sánchez-Díaz, J. Ajuria, R. Pacios and E. Palomares, *Phys. Chem. Chem. Phys.*, 2011, **13**, 6105–6109.
- M. Irwin, D. Buchholz, A. Hains, R. Chang and T. Marks, *Proc. Natl. Acad. Sci. U. S. A.*, 2008, **105**, 2783–2787.
- F. Liu, S. Shao, X. Guo, Y. Zhao and Z. Xie, *Sol. Energy Mater. Sol. Cells*, 2010, **94**, 842–845.
- V. Shrotriya, G. Li, Y. Yao, C.-W. Chu and Y. Yang, *Appl. Phys. Lett.*, 2006, **88**, 073508.
- K. X. Steirer, P. F. Ndione, N. E. Widjonarko, M. T. Lloyd, J. Meyer, E. L. Ratcliff, A. Kahn, N. R. Armstrong, C. J. Curtis, D. S. Ginley, J. J. Berry and D. C. Olson, *Adv. Energy Mater.*, 2011, **1**, 813–820.
- M. D. Irwin, J. D. Servaites, D. Buchholz, B. Leever, J. Liu, J. Emery, M. Zhang, J. Song, M. Durstock, A. Freeman, M. Bedzyk, M. Hersam, R. Chang, M. Ratner and T. Marks, *Chem. Mater.*, 2011, **23**, 2218–2226.
- J. Zhang, J. Wang, Y. Fu, B. Zhang and Z. Xie, *J. Mater. Chem. C*, 2014, **2**, 8295–8302.
- V. Wood, M. J. Panzer, J. E. Halpert, J. M. Caruge, M. G. Bawendi and V. Bulović, *ACS Nano*, 2009, **vol. 3**, 3581–3586.
- B. S. Mashford, T.-L. Nguyen, G. J. Wilson and P. Mulvaney, *J. Mater. Chem.*, 2009, **20**, 167.
- J.-Y. Jeng, K.-C. Chen, T.-Y. Chiang, P.-Y. Lin, T.-D. Tsai, Y.-C. Chang, T.-F. Guo, P. Chen, T.-C. Wen and Y.-J. Hsu, *Adv. Mater.*, 2014, **26**, 4107–4113.
- W. Chen, Y. Wu, J. Liu, C. Qin, X. Yang, A. Islam, Y.-B. Cheng and L. Han, *Energy Environ. Sci.*, 2015, **8**, 629–640.
- L. Hu, J. Peng, W. Wang, Z. Xia, J. Yuan, J. Lu, X. Huang, W. Ma, H. Song, W. Chen, Y.-B. Cheng and J. Tang, *ACS Photon.*, 2014, **1**, 547–553.
- J. R. Manders, S.-W. Tsang, M. J. Hartel, T.-H. Lai, S. Chen, C. M. Amb, J. R. Reynolds and F. So, *Adv. Funct. Mater.*, 2013, **23**, 2993–3001.
- J. Jung, D. L. Kim, S. H. Oh and H. J. Kim, *Sol. Energy Mater. Sol. Cells*, 2012, **102**, 103–108.
- K. X. Steirer, J. P. Chesin, N. E. Widjonarko, J. J. Berry, A. Miedaner, D. S. Ginley and D. C. Olson, *Org. Electron.*, 2010, **11**, 1–5.



- 31 V. Zardetto, T. M. Brown, A. Reale and A. Di Carlo, *J. Polym. Sci., Part B: Polym. Phys.*, 2011, **49**, 638–648.
- 32 A. A. Balandin, E. S. Grigoryan and Z. S. Yanysheva, *Acta Physicochim. URSS*, 1940, **12**, 737–758.
- 33 L. L. Bircumshaw and J. Edwards, *J. Chem. Soc.*, 1950, 1800–1809.
- 34 G. R. Rao, K. C. Patil and C. N. R. Rao, *Inorg. Chim. Acta*, 1970, **4**, 215–218.
- 35 B. R. Wheeler and A. K. Galwey, *J. Chem. Soc., Faraday Trans. 1*, 1974, **70**, 661–670.
- 36 A. K. Galwey and M. E. Brown, *Proc. R. Ir. Acad., Sect. B*, 1977, **77**, 465–471.
- 37 A. B. Edwards, C. D. Garner and K. J. Roberts, *J. Phys. Chem. B*, 1997, **101**, 20–26.
- 38 A. H. Qusti, A. A. Samarkandy, S. Al-Thabaiti and E.-H. M. Diefallah, *J. King Abdulaziz Univ.*, 1997, **9**, 73–81.
- 39 Z. Zhai, X. Huang, M. Xu, J. Yuan and J. Peng, *Adv. Energy Mater.*, 2013, **3**, 1614–1622.
- 40 N. E. Widjonarko, P. Schulz, P. A. Parilla, C. L. Perkins, P. F. Ndione, A. K. Sigdel, D. C. Olson, D. S. Ginley, A. Kahn, M. F. Toney and J. J. Berry, *Adv. Energy Mater.*, 2014, **4**, 1301879.
- 41 E. L. Ratcliff, J. Meyer, K. X. Steirer, A. Garcia, J. J. Berry, D. S. Ginley, D. C. Olson, A. Kahn and N. R. Armstrong, *Chem. Mater.*, 2011, **23**, 4988–2000.
- 42 S. R. Cowan, J. V. Li, D. C. Olson and E. L. Ratcliff, *Adv. Energy Mater.*, 2014, **5**, 1400549.
- 43 G. Narain, *Bull. Chem. Soc. Jpn.*, 1966, **39**, 2298–2300.
- 44 N. Armstrong, C. Carter, C. Donley, A. Simmonds, P. Lee, M. Brumbach, B. Kippelen, B. Domercq and S. Yoo, *Thin Solid Films*, 2003, **445**, 342–352.
- 45 M. P. Seah, *Surf. Interface Anal.*, 1989, **14**, 488–488.
- 46 M. Ghaemy, A. Omrani and A. Rostami, *J. Appl. Polym. Sci.*, 2005, **97**, 265–271.
- 47 H. Icbudak, T. K. Yazicilar and W. T. Yilmaz, *Thermochim. Acta*, 1999, **335**, 93–98.
- 48 K. S. Rejitha and S. Mathew, *J. Therm. Anal. Calorim.*, 2010, **102**, 931–939.
- 49 V. T. Yilmaz, S. Demir, O. Andac and W. T. A. Harrison, *J. Coord. Chem.*, 2002, **55**, 863–872.
- 50 K. S. Rejitha and S. Mathew, *J. Therm. Anal. Calorim.*, 2008, **93**, 213–217.
- 51 H. G. Edwards and A. Knowles, *J. Mol. Struct.*, 1992, **268**, 13–22.
- 52 F. J. A. Garcia and M. M. P. Manrique, *Ion*, 1976, **36**, 422–425.
- 53 D. Stoilova and V. Koleva, *Spectrochim. Acta, Part A*, 2001, **57**, 2629–2636.
- 54 P. D. Buckley and P. A. Giguere, *Can. J. Chem.*, 1967, **45**, 397–407.
- 55 W. Sawdony, K. Niedenzu and J. W. Dawson, *Spectrochim. Acta, Part A*, 1967, **23**, 799–806.
- 56 J. A. Broomhead and L. A. P. Kane-Maguire, *J. Chem. Soc. A*, 1967, **4**, 546–553.
- 57 T. Iwamoto and D. F. Shriver, *Inorg. Chem.*, 1971, **10**, 2428–2432.
- 58 J. J. A. McLean, A. F. Schreiner and A. F. Laethem, *J. Inorg. Nucl. Chem.*, 1964, **26**, 1245–1248.
- 59 A. Sabatini and S. Califano, *Spectrochim. Acta*, 1960, **16**, 677–688.
- 60 M. L. Niven and G. C. Percy, *Transition Met. Chem. (Dordrecht, Neth.)*, 1978, **3**, 267–271.
- 61 J. L. Bantignies, S. Deabate, A. Righi, S. Rols, P. Hermet, J. L. Sauvajol and F. Henn, *J. Phys. Chem. C*, 2008, **112**, 2193–2201.
- 62 N. Bellakhal and J. L. Brisset, *J. Chem. Res., Synop.*, 2001, 238–239.
- 63 D. S. Hall, D. J. Lockwood, S. Poirier, C. Bock and B. R. MacDougall, *J. Phys. Chem. A*, 2012, **116**, 6771–6784.
- 64 L. Escobar-Alarcon, A. Arrieta, E. Camps, S. Romero and M. A. Camacho-Lopez, *Superficies Vacio*, 2005, **18**, 9–12.
- 65 D. Forster and D. M. L. Goodgame, *J. Chem. Soc.*, Jan., 1965, 262–267.
- 66 B. Mavis and M. Akinc, *Chem. Mater.*, 2006, **18**, 5317–5325.
- 67 H. Yang and J. L. Whitten, *Surf. Sci.*, 1998, **401**, 313–321.
- 68 R. Micu-Semeniuc and C. G. Macarovici, *Rev. Roum. Chim.*, 1972, **17**, 981–992.
- 69 H. Yang and J. L. Whitten, *Journal of Chem. Phys.*, 1997, **107**, 8518–8524.
- 70 A. Ludi and R. Huegi, *Helv. Chim. Acta*, 1968, **51**, 349–355.
- 71 N. S. Gajbhiye and S. Prasad, *Thermochim. Acta*, 1996, **285**, 325–336.
- 72 J. Kristof, A. D. Battisti, G. Keresztury, E. Horvath and T. Sziligyi, *Langmuir*, 2001, **17**, 1637–1640.
- 73 W. M. H. Oo, M. D. McCluskey, A. D. Lalonde and M. G. Norton, *Appl. Phys. Lett.*, 2005, **86**, 073111.
- 74 E. L. Williams, R. A. Bailey and S. E. Wberley, *J. Inorg. Nucl. Chem.*, 1970, **32**, 2197–2203.
- 75 T. J. Taylor, D. Dollimore, G. A. Gamlen, A. J. Barnes and M. A. Stuckey, *Thermochim. Acta*, 1986, **101**, 291–304.
- 76 S. Srinivasan and S. Ganguly, *Catal. Lett.*, 1991, **10**, 279–288.
- 77 F. Urbach, *Phys. Rev.*, 1953, 1324.
- 78 A. Fujimori and F. Minami, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1984, **30**, 957–971.
- 79 S. Hüfner, *Adv. Phys.*, 1994, **43**, 183–356.
- 80 J. D. Zimmerman, B. Song, O. Griffith and S. R. Forrest, *Appl. Phys. Lett.*, 2013, **103**, 243905.

