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Thermoelectric properties of pulsed current sintered nanocrystalline Al-doped ZnO by chemical vapour synthesis

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ZnO is a promising n-type oxide thermoelectric material, which is stable in air at elevated temperatures. In the present study, we report the bottom-up approach to create Al-doped ZnO nanocomposites from nanopowders, which are prepared by chemical vapour synthesis. With our synthesis route, we are able to create highly doped Al-containing ZnO nanocomposites that exhibit bulk-like electrical conductivity. Moreover, the impact of the microstructure of the nanocomposites on their thermal conductivity is enormous, with a value of $1.0 \text{ W m}^{-1} \text{ K}^{-1}$ for 1% Al-ZnO at room temperature, which is one of the lowest values reported, to date, on ZnO nanocomposites. The optimization of the Al-doping and microstructure with respect to the transport properties of bulk Al-ZnO nanocomposites leads to a zT value of about 0.24 at 950 K, underlining the potential of our technique.

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Introduction

Nanostructured materials based on alloys, namely, tellurides, skutterudites, transition-metal silicides and Si-Ge alloys, exhibit good thermoelectric properties.^{1,2} However, their practical utilization is limited because they cannot be operated in air due to their oxidation and the high cost of the materials. Metal oxides are potential candidates for high temperature thermoelectric applications because they are thermally stable in air and show excellent durability. In particular, zinc oxide (ZnO) has been investigated as an n-type thermoelectric material for high temperature applications because of its better conversion efficiency compared to other oxides, abundance, non-toxicity and low cost.^{3,4}

The thermoelectric conversion efficiency of a material is closely related to the dimensionless figure of merit zT , which is given as follows:

$$zT = \frac{\alpha^2 \sigma}{\kappa} T \quad (1)$$

where α , σ , κ and T denote the Seebeck coefficient, electrical conductivity, thermal conductivity and absolute temperature, respectively.

The thermoelectric efficiency can be enhanced by the maximization of the power factor ($\alpha^2 \sigma$) and reduction of the thermal conductivity. The power factor can be optimized by band gap engineering^{5,6} or doping. However, the most common approach

for the figure of merit enhancement is to reduce the total thermal conductivity. In particular, the lattice contribution can be altered by nanostructuring and placing nanoprecipitates in the matrix;⁷⁻⁹ therefore, creating abundant grain boundaries, which scatter phonons more strongly than electronic charge carriers.

For using nanostructured materials in thermoelectric devices, bulk solids are required. One of the methods to create bulk nanostructured materials is the bottom-up approach, in which nanoparticles are compacted into a bulk material. Chemical vapour synthesis (CVS) is a facile method for the continuous production of nanoparticles that can be up-scaled to industrial levels.¹⁰ The CVS route is advantageous for tailoring the powder properties, such as particle diameter, size distribution, morphology, crystallinity and doping concentration. The sintering of nanoparticles also offers a promising route for the synthesis of bulk samples with nanoscale constituents. However, the challenges during the sintering process are to limit the grain growth and to preserve a high density of grain boundaries. Pulsed electric current sintering (PECS) is a rapid sintering technique in which pulsed electric current and uniaxial pressure are simultaneously applied. It has been successfully demonstrated to manufacture bulk nano-scaled bodies.¹¹

In this work, we report the high temperature electrical and thermal properties of Al-doped ZnO (AZO) nanocomposites. The AZO nanoparticles are prepared by the CVS route, which are then processed into bulk nanocomposites using PECS. The structural and thermoelectric properties of the sintered ceramic pellets are characterized. This study demonstrates the impact of the processing technique on the incorporation of a high

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amount of electrically active Al-dopant, as reflected in the high electrical conductivity value of about 10^5 S m^{-1} and low Seebeck coefficient of about $-50 \mu\text{V K}^{-1}$ at room temperature. In addition, the thermal conductivity is drastically decreased to an extremely low value of $1.0 \text{ W m}^{-1} \text{ K}^{-1}$ at ambient temperature. In regards to the obtained data, we discuss how the doping of nano-microstructure and composite phase-mixture of ZnO and ZnAl_2O_4 in the sintered samples affect the charge transport, and thus the thermoelectric performance of AZO nanocomposites.

Experimental details

Synthesis of nanoparticles

ZnO and AZO nanoparticles are prepared by CVS using a hot-wall reactor. The set-up consists of a precursor delivery unit, a hot-wall reactor and a thermophoretic particle collector. The set-up is described in detail elsewhere.^{12,13} The precursor unit comprises of two bubblers filled with liquid precursors, diethylzinc (DEZn) and triethylaluminium (TEAL), which are placed in a temperature-controlled oil bath to adjust the vapour pressure of the precursors. The precursors DEZn and TEAL are transported into the reactor chamber using helium as the carrier gas. Oxygen acting as a reaction gas with a flow rate of 1000 sccm (standard cubic centimeter per minute) is added to the helium-precursor mixture just before the gas mixture enters the reactor to avoid preliminary oxidation of the precursors. Thermal mass flow controllers are used to control the flow rates of the gases. The reactor consists of a ceramic alumina tube heated by a tube furnace. The synthesis temperature equal to the wall temperature of alumina tube is fixed at 1173 K. The precursors and oxygen react in the reaction zone to form oxide particles, which are subsequently transported to the thermophoretic particle collector by the gas stream. AZO powders are prepared by adjusting the desired mass flow of the DEZn and TEAL vapours, which is achieved by altering the oil bath temperatures and the helium flow rates through the precursors corresponding to nominal doping concentration of 0, 0.5, 1, 2 and 8 at% Al in the particles. In the collector, a thermophoretic force acts on the particles because of the temperature gradient between the hot quartz lamps and the water-cooled wall of the collector, where the particles are finally deposited. The powder is mechanically removed from the wall by a scraper. The process pressure is maintained at 20 mbar by a butterfly valve and measured by a capacitive absolute pressure gauge placed at the outlet of the reactor.

Processing to bulk nanocomposites

The compaction of AZO and ZnO nanopowders into bulk nanocomposites is achieved using the PECS technique. A FCT HP D5/2 (FCT Systeme GmbH, Rauenstein, Germany) instrument is used for sintering. The powder is loaded into a graphite die with an inner diameter of 20 mm fabricated from high strength graphite. A graphite foil coated with boron nitride (BN) is used to avoid the contact of the powder with the inner surface of the die. BN is used as a high temperature electrical insulator to ensure that the current flows through the powder and not

through the dies during sintering. To reduce the radial temperature gradient in the dies and to minimize radiation heat losses, the complete die assembly is surrounded with carbon wool. The nanopowders are pre-pressed at 12 MPa and heated by a pulsed electric current with a heating rate of 100 K min^{-1} to the desired sintering temperature. A pulse pattern of 10 : 5 (10 ms on and 5 ms off) is used for sintering the samples. The samples are sintered at $700 \text{ }^\circ\text{C}$ or $900 \text{ }^\circ\text{C}$ with a holding time of 3 minutes. A uniaxial pressure of 35 MPa is applied to the specimens over the entire heating and cooling cycle. The samples are cooled to $500 \text{ }^\circ\text{C}$ at a cooling rate of 100 K min^{-1} . Furthermore, the samples are furnace cooled to room temperature with a continuous release of applied pressure. The heating power is adjusted using a pyrometer, which monitors the temperature at the bottom of the upper graphite push-punch of the dies. The whole sintering procedure is performed under 1 mbar Ar atmosphere. The graphite foils on the surface of the sintered samples are removed by grinding.

Characterization

The phase composition, crystal structure and isotropic crystallite size of the as-prepared nanopowders and sintered nanocomposites are determined by X-ray diffraction (XRD) using a PANalytical X-ray diffractometer (X-Pert PRO) with Ni-filtered $\text{Cu K}\alpha$ (0.154 nm) radiation detected by an X-Celerator detector. The data are measured in the 2θ range = $20\text{--}120^\circ$ and analyzed by Rietveld refinement using the MAUD software.¹⁴ The geometric density of the sintered samples is determined from the sample mass and geometry. The nano-microstructure and elemental analysis of the as-synthesized powders and sintered pellets are characterized by a high-resolution scanning electron microscope (SEM, JEOL-JSM 7500 F) using a scanning transmission detector (STEM) equipped with an energy dispersive X-ray analyser (EDX, QX200 Bruker Quantax 2000).

The electrical conductivity and the Seebeck coefficient of the sintered nanocomposites are measured under He atmosphere at 100 mbar using a ZEM-3 from Ulvac-Riko, Inc. The thermal conductivity is determined by the laser flash method using a LFA 457 Microflash from Netzsch-Gerätebau GmbH under atmospheric pressure at a flow of 75 sccm N_2 . Because of the limited sample geometry within the available measurement techniques, the thermal transport is measured parallel to the direction of the pressing force and the sintering current, whereas the electrical transport is characterized perpendicular to this direction. This aspect has to be taken into account when both transport processes are considered in combination.

Results and discussion

Microstructure of nanoparticles and nanocomposites

Fig. 1 depicts the XRD patterns of the as-synthesized ZnO nanoparticles doped with 0, 0.5, 1, 2 and 8% of Al. The patterns show that the doped and undoped particles exhibit the hexagonal wurtzite structure and exist in a single phase. It has been reported that the thermodynamic solubility limit of Al in ZnO is below 1%.^{15,16} Even though in the present study the doping



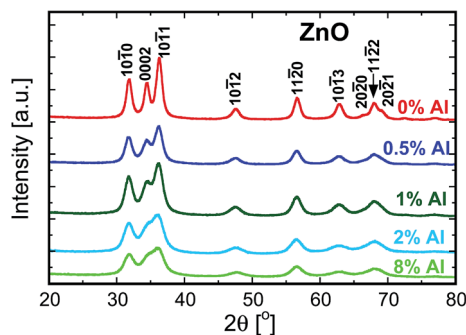


Fig. 1 X-ray diffractograms of the as-synthesized AZO nanopowders by the CVS method.

exceeded the solubility limit, we did not observe a second phase in the as-prepared nanopowders. This could be because of the CVS method, which is a non-equilibrium process, and therefore allows the incorporation of more dopant than permitted by the thermodynamic limit.^{12,13} The continuous increase in the broadening of the XRD signals with increasing Al concentration (see Fig. 1) indicates that the crystallite size of the particles decreases. Both the crystallite size and the microstrain are determined from the Rietveld refinement of the data and are plotted as a function of Al concentration in Fig. 2. The addition of Al as a dopant hinders the growth rate of the nanoparticles because of the impurity drag, which results in smaller crystallites. It can also be seen from Fig. 2 that the lattice microstrain increases with Al concentration, indicating that Al is incorporated into the ZnO lattice. The findings are in agreement with the previous studies.^{17,18} Fig. 3 shows the STEM micrograph of undoped ZnO nanopowders, which consists of nearly spherical crystalline particles with a small degree of agglomeration with a mean diameter of 15 nm, thus substantiating the XRD data.

The XRD patterns of undoped and 8% Al-doped pellets sintered at 900 °C are shown in Fig. 4. As is evident by the narrower Bragg reflections, substantial grain growth occurred during sintering. Most of the Bragg reflections for the AZO nanocomposites belong to the hexagonal wurtzite phase, similar to

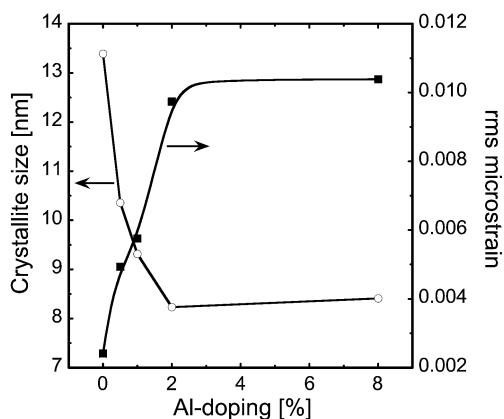


Fig. 2 Evolution of crystallite size and microstrain as a function of the nominal Al concentration for the as-prepared AZO nanoparticles.

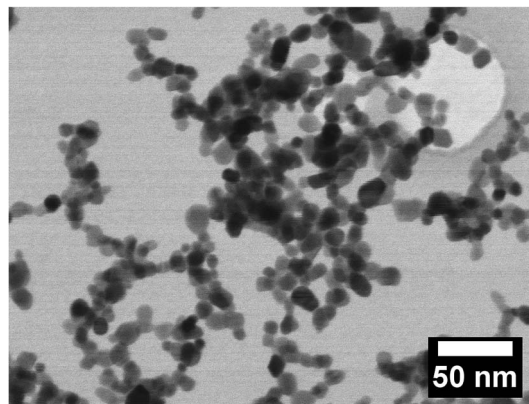


Fig. 3 SEM micrograph (in STEM mode) of the as-synthesized undoped ZnO nanoparticles.

that of undoped ZnO. Moreover, the patterns reveal the presence of a ZnAl_2O_4 spinel phase (gahnite). This additional phase is observed in the patterns when the nanocomposites have an Al concentration $\geq 1\%$ at the sintering temperature of 900 °C. The XRD pattern of undoped ZnO sintered at 900 °C shows significant texturing. The intensity of the (0002) reflection is strongly decreased, while for the (10-10) reflection, the intensity is increased compared to the undoped nanoparticles. This indicates a preferential growth along the c -axis of the crystals,¹⁹ suggesting that the externally applied pressure and the pulsed current during sintering define the texture direction. The addition of Al as a dopant in the nanocomposites enhances the intensity of the (0002) reflection in comparison with the (10-10) reflection, as seen in Fig. 4, indicating a change in the anisotropy of the surface enthalpy of the particles. This modifies the growth direction and dynamics with the presence of Al in the samples.^{20,21}

Fig. 5 shows the SEM micrographs of the fresh fracture surface of AZO nanocomposites. The structure of undoped ZnO sintered at 900 °C shows a grain size of 1–2 μm having a hexagonal cross-section with intense faceting in the material. The faceting substantiates the XRD observation of the preferred growth along the c -axis of the ZnO lattice. The hexagonal cross-sectional structures are created because of the lower energy associated with these facets.²⁰ The micrographs for the AZO

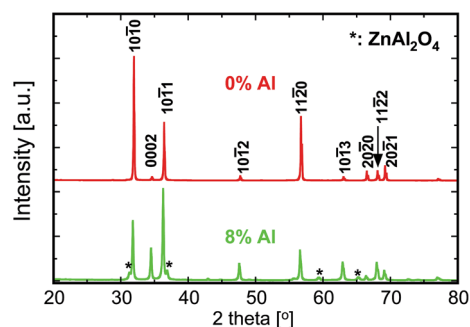


Fig. 4 X-ray diffractograms of the undoped and 8% Al-doped ZnO nanocomposites sintered at 900 °C by PECS.



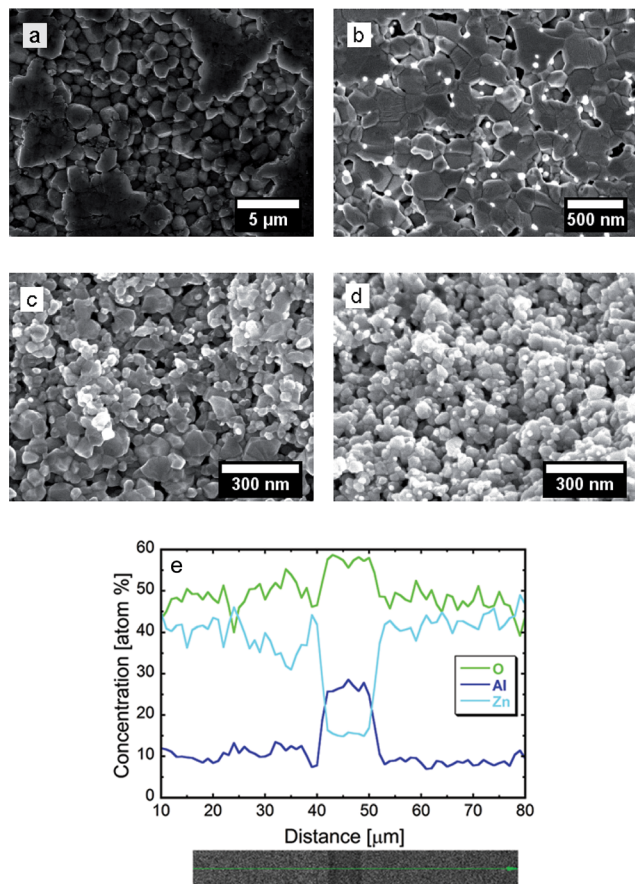


Fig. 5 SEM micrographs of the fractured surfaces of (a) undoped and (b) 0.5% AZO pellets sintered at 900 °C; (c) 0.5%, (d) 1% AZO nanocomposites sintered at 700 °C and (e) EDX line-scan across a region containing ZnAl_2O_4 nanoprecipitates in 8% AZO sample sintered at 900 °C.

nanocomposites reveal that ZnO grains are interspersed with the spinel ZnAl_2O_4 phase, which precipitates at the ZnO grain boundaries. The presence of Al in the nanocomposites inhibits the growth of ZnO grains, and the average size is confined in the range of 80–200 nm. The size of the ZnAl_2O_4 precipitates is in the order of 10–50 nm. The grain growth for both ZnO and ZnAl_2O_4 depends on the sintering temperature, as well as on the Al concentration in the material. Grain/crystallite growth is restricted because of the Al-doping not only during particle synthesis, but also during sintering due to the pinning of the grain boundaries.^{17,19} In addition, the presence of Al changes the surface energy of the (0001) facet, which can give rise to twins and planar defects.²¹ The morphology of the grains changes from hexagonal, as observed for undoped ZnO (see Fig. 5a), to a plate/planar-like structure for AZO nanocomposites, as shown in Fig. 5b–d. The addition of Al suppresses the grain boundary mobility, thus enabling the pores to remain on the moving boundaries during sintering. Thus, most of the porosity is located in the vicinity of the ZnAl_2O_4 precipitates, as revealed in the SEM micrographs (see Fig. 5b–d). These porous materials with fine grains may act as good thermoelectric materials because their thermal

conductivity can be significantly reduced without considerably affecting their electrical conductivity. The EDX line-scan across a region, particularly at a distance of 40–50 μm (see Fig. 5e), in the 8% AZO sample shows an enhancement of Al and O concentration in the nanoprecipitates region, consistent with the formation of ZnAl_2O_4 , as substantiated by the XRD pattern of the sintered pellets (see Fig. 4).

The bulk relative density of the sintered samples is in the range of 65–93%. The density of the specimens is affected both by the sintering temperature and the Al concentration in the samples. Increasing amounts of Al in the material, as well as high sintering temperatures, create more ZnAl_2O_4 precipitates. This makes densification increasingly difficult because of the difference in the physical properties of ZnAl_2O_4 and ZnO. This observation is in agreement with the previous findings in the literature.²²

Dependence of the transport properties of nanocomposites on Al concentration

The room temperature electrical conductivities of the nanocomposites sintered at 700 °C and 900 °C as a function of Al concentration are illustrated in Fig. 6. The undoped ZnO sample sintered at 900 °C has a room temperature electrical conductivity of $3 \times 10^3 \text{ S m}^{-1}$. The addition of 0.5% Al enhances the conductivity by two orders of magnitude to $1.5 \times 10^5 \text{ S m}^{-1}$. The substitution of Zn^{2+} by Al^{3+} increases the free charge carrier concentration to compensate for the electrical charge balance, and in turn the electrical conductivity of the sample.^{17,23} This value is one of the highest for ZnO nanocomposites and is comparable to the bulk value of about $2 \times 10^5 \text{ S m}^{-1}$.^{24,25} In contrast, Nam *et al.* reported an electrical conductivity of about $1.3 \times 10^4 \text{ S m}^{-1}$ for ZnO nanocomposites.²⁶ This indicates that with our processing method we are able to prepare nanocomposites incorporating an active dopant in the material. On further addition of Al, the conductivity slightly decreased for the samples sintered at 900 °C.

The samples sintered at 700 °C exhibit lower electrical conductivity compared to those sintered at 900 °C and their conductivity decreases in a similar pattern with increasing Al

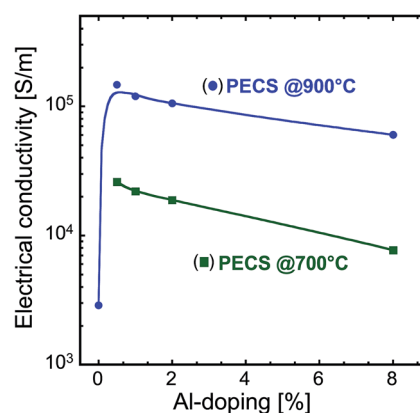


Fig. 6 Room temperature electrical conductivity as a function of Al concentration for the nanocomposites sintered at 700 °C and 900 °C.



concentration (see Fig. 6). First, an increasing fraction of the non-conductive ZnAl_2O_4 phase reduces the volume available for the charge transport. Second, we suggest that smaller AZO grains and more ZnAl_2O_4 precipitate enhances the scattering of the charge carriers in increasingly doped samples, thus decreasing their mobility, and in turn the electrical conductivity. This suggestion is in accordance with the microstructural observations (see Fig. 5a–d). In addition to the mobility, the other main factor for electrical conduction, namely, the charge carrier concentration, originating from the electrically active dopant, cannot be directly obtained from the microstructural data and it is discussed in the subsequent section.

The impact of high electrical conductivity values is also observed in the Seebeck coefficient of the nanocomposites. Fig. 7 illustrates how the Seebeck coefficient varies with Al-doping for nanocomposites sintered at 700 °C and 900 °C. The measured values of the Seebeck coefficients are negative for all the samples, indicating the n-type conduction. The room temperature Seebeck coefficient of the undoped ZnO sample sintered at 900 °C is about $-280 \mu\text{V K}^{-1}$, which is in agreement with the relatively low electrical conductivity of this sample. As anticipated, with increasing Al-doping, the absolute value of the Seebeck coefficient decreases, and reaches a value of about $-50 \mu\text{V K}^{-1}$ at 0.5% Al. The further addition of Al does not have considerable impact on the values of the Seebeck coefficient for the samples sintered at 900 °C. As mentioned above, Al-doping results in the increase in charge carrier concentration by the substitution of Zn by Al at the Zn-site, but this effect is confined by the low solubility limit. Interestingly, even though the electrical conductivities of the samples sintered at 700 °C and 900 °C are remarkably different (see Fig. 6 and 7), their Seebeck coefficient values are very similar, confirming the assumption of Kinemuchi *et al.*⁴ and Nam *et al.*²⁶ that the impact of grain boundaries in the nanocomposite is less significant on the Seebeck coefficient than on the electrical conductivity. However, it is worth noting that we observe a local minimum in the room temperature Seebeck coefficient at 1% Al-doping for both the sintering temperatures. In the case of the sample sintered at 700 °C, the electrical conductivity only slightly decreased

compared to the sample with 0.5 Al%-doping, and the power factor reaches a maximum value at this dopant concentration amongst the samples sintered at 700 °C.

The impact of the nanostructure on the thermal conductivity of sintered samples is enormous. Fig. 8 illustrates how the room temperature thermal conductivity of the samples sintered at 700 °C and 900 °C varies as a function of Al content. Undoped ZnO sintered at 900 °C exhibits a room temperature thermal conductivity of about $40 \text{ W m}^{-1} \text{ K}^{-1}$, similar to the bulk value.^{3,27,28} This phenomenon is expected because the specimen contains grains in the micrometer range (see Fig. 5a), considerably larger than the mean free path of about 30 nm for phonons in ZnO.^{29,30} The addition of Al in the ZnO nanopowders decreases the thermal conductivity of the sintered nanocomposites. The amount of reduction in thermal conductivity depends on the sintering temperature and the Al content. In the case of the samples sintered at 900 °C, the thermal conductivity decreases with increasing Al content, reaching its minimum of $7 \text{ W m}^{-1} \text{ K}^{-1}$ at 8% Al-doping. However, the samples sintered at 700 °C exhibit considerably lower values and have a pronounced minimum at a doping concentration of 1% Al. The thermal conductivity of this sample is reduced by a factor of 40, compared to the bulk and reaches a value of about $1.0 \text{ W m}^{-1} \text{ K}^{-1}$, which is one of the lowest reported for oxide materials. Our findings on the low thermal conductivity of nanocomposites are in agreement with the recently reported observations on AZO thin films by Loureiro *et al.*³¹ Kinemuchi *et al.*^{4,29} reported that the thermal conductivity of ZnO nanocomposites can reach values below $5 \text{ W m}^{-1} \text{ K}^{-1}$ when the grain size is around 50 nm, whereas Nam *et al.*²⁶ found reductions for grain sizes below 200 nm. Because the phonon mean free path in ZnO is approximately 30 nm,^{29,30} our findings of a strong reduction in the thermal conductivity for grain sizes of ZnO in the range of 50–100 nm agree well with these reported values. The smaller grains of ZnAl_2O_4 (10–30 nm, see Fig. 5c, d and 8) also have a major impact on the decrease of thermal conductivity. We attribute this strong reduction in the total thermal conductivity to the reduction in the lattice contribution by the increased phonon scattering from the

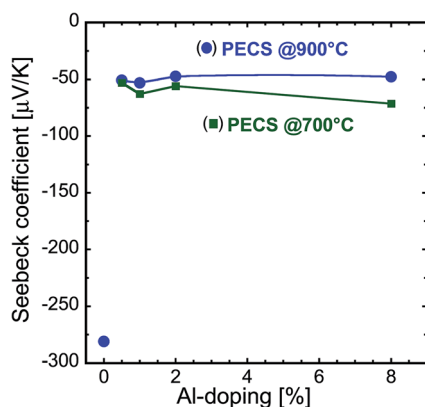


Fig. 7 Room temperature Seebeck coefficient as a function of Al concentration for the nanocomposites sintered at 700 °C and 900 °C.

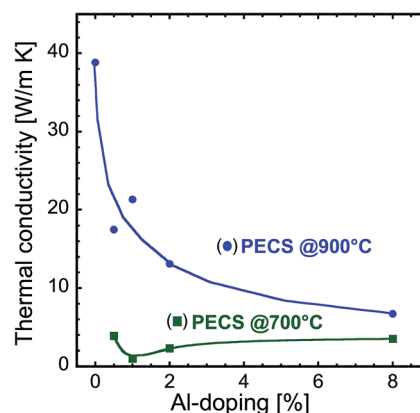


Fig. 8 Variation of room temperature thermal conductivity as a function of Al concentration for the nanocomposites sintered at 700 °C and 900 °C.



ZnAl₂O₄ precipitates, point defects, porosity and refined nano-grains by Al-doping.

Dependence of transport properties of the sintered samples on temperature

As a representative example for illustrating and discussing the charge transport of doped nanocomposites, we used 1% AZO sample and compared it with the undoped specimen. The Seebeck coefficient α of the nanocomposites as a function of measuring temperature is illustrated in Fig. 9 for undoped and 1% AZO nanocomposites sintered at 700 °C and 900 °C. For the undoped ZnO specimen sintered at 900 °C, the Seebeck coefficient has large negative values in the range of -280 to -350 $\mu\text{V K}^{-1}$ indicating a low charge carrier concentration. However, the Al-doped samples exhibit a low absolute value of the Seebeck coefficient over the complete temperature interval, highlighting the enhancement in the concentration of free charge carriers, as previously discussed. As illustrated in Fig. 9, the Seebeck coefficient decreases with temperature, which is in agreement with the findings in the literature for bulk doped-ZnO ceramics.^{3,32–34} Similar results were also reported for doped nanocomposites.^{4,26,29}

The temperature dependence of the Seebeck coefficient of a degenerate semiconductor can be described using the Jonker and Pisarenko relationship,⁸ as expressed in eqn (2)

$$\alpha = \frac{8\pi^2 k_B^2}{3eh^2} m^* \left(\frac{\pi}{3n}\right)^{2/3} T \quad (2)$$

where n is the charge carrier concentration and m^* is the effective mass at the Fermi level. The remaining parameters are natural constants: h is the Planck constant, e is the elementary charge and k_B is the Boltzmann constant. Eqn (2) clearly shows that a higher charge carrier concentration n causes a decrease in α , provided that m^* and T remain constant, which explains the tendency observed in Fig. 7. Moreover, eqn (2) shows that the Seebeck coefficient should linearly vary with the temperature, as illustrated in Fig. 9, provided that m^* and n are constant in the temperature interval. The charge carrier concentration n can be

evaluated by fitting the data in Fig. 9 for doped samples with eqn (2), and using $m^* = 0.33m_e$, where m_e is the electron rest mass.²⁶ In the case of the samples sintered at 700 °C, the low temperature interval (300–670 K) is considered for evaluating the charge carrier concentration. It should be noted that this analysis yields the charge carrier concentration in the volume fraction of the conductive material.

The charge carrier concentration for different samples is plotted in Fig. 10. It can be seen that the carrier concentration in all the sintered samples is considerably high, with values in the order of 10^{20} cm^{-3} , satisfying the criterion of degeneracy. For those samples sintered at 900 °C, we find values that are a factor of 3 to 4 times larger than for the samples sintered at 700 °C, indicating that at higher sintering temperatures more Al is incorporated into the ZnO lattice. Considering the amount of Al-dopants, the carrier concentration is of the same order as the theoretical concentration³⁵ for 0.5% Al, highlighting the effectiveness of our synthesis route to prepare highly doped material (see Fig. 10 and ref. 35). For higher Al contents, the evaluated carrier concentration is different from the theoretical Al content, indicating that only a limited fraction of added Al might be effective as the dopant and the remaining Al is precipitated as ZnAl₂O₄ spinel phase, as demonstrated by the XRD and SEM micrographs of the sintered samples.

Fig. 11 shows the temperature dependence of the electrical conductivity σ of the undoped and 1% AZO nanocomposites. All the samples exhibit decreasing electrical conductivity with increasing temperature, indicating a metallic or strongly degenerate behaviour. It should be noted that even for undoped ZnO, we observe the same tendency, though the absolute values are considerably lower than for the Al-doped samples.

The electrical conductivity can be expressed as follows:

$$\sigma = ne\mu_e \quad (3)$$

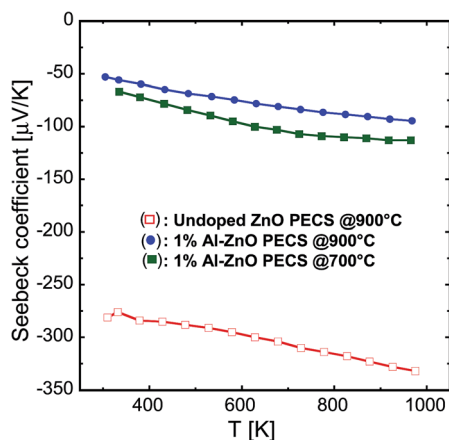


Fig. 9 Temperature dependence of the Seebeck coefficient of the nanocomposites sintered at 700 °C and 900 °C.

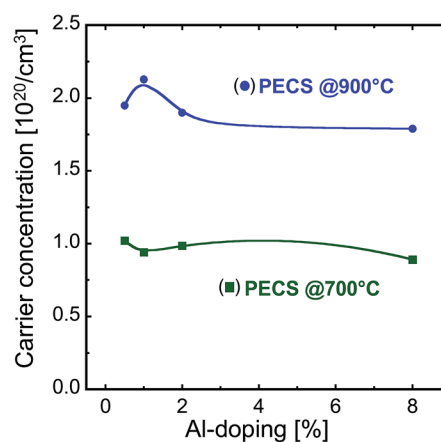


Fig. 10 Room temperature charge carrier concentration as a function of Al concentration for the nanocomposites sintered at 700 °C and 900 °C. The carrier concentration was evaluated from the temperature dependence of the Seebeck coefficient and refers to the electrically active volume fraction.



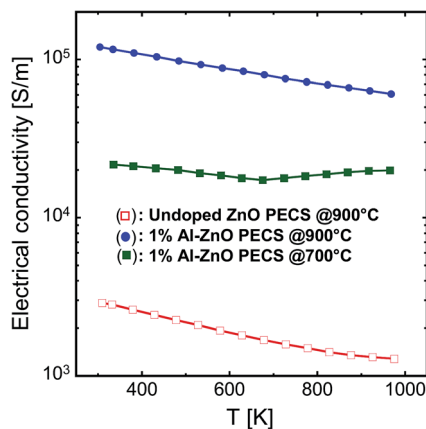


Fig. 11 Temperature dependence of the electrical conductivity of the nanocomposites sintered at 700 °C and 900 °C.

where, μ_e is the average mobility of charge carriers in the total volume of the material.

Taking into account eqn (3) and the charge carrier concentrations from Fig. 10, the room temperature mobility of the charge carriers can be evaluated from the electrical conductivity. By dividing the evaluated average mobility of the charge carriers with the relative density of the samples, we obtain a normalized mobility value for the charge carriers. This normalized mobility takes into account the electrically active volume of the samples and more directly reflects the influence of the nano-microstructure on the transport of the charge carriers. The normalized mobility is illustrated for the room temperature values as a function of Al content in Fig. 12. The mobility of the charge carriers in the sintered nanocomposites decreases with increasing Al-doping in the material, irrespective of the sintering temperature, which is in agreement with the previous reports.^{35–37} At room temperature, the mobility

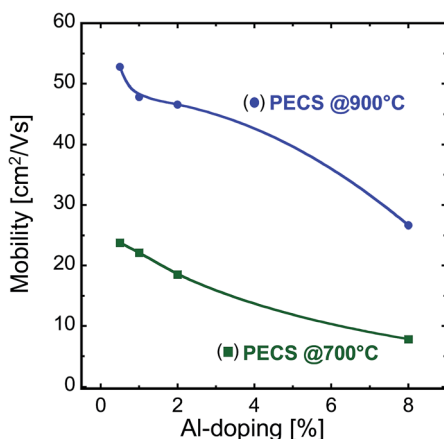


Fig. 12 Normalized room temperature mobility as a function of Al concentration of the nanocomposites sintered at 700 °C and 900 °C. The normalization neglects the contribution of electrically inactive volume fraction (pores), more directly reflecting the influence of the grain boundaries on the transport of the charge carriers. The average mobility in the total volume is lower.

of the charge carriers in most of the samples is well in the range of undoped and Al-doped polycrystalline bulk ZnO ($20\text{--}80\text{ cm}^2\text{ V}^{-1}\text{ s}^{-1}$)³⁵ and pulsed laser deposited thin films ($30\text{--}10\text{ cm}^2\text{ V}^{-1}\text{ s}^{-1}$, decreasing with increasing Al-content).³⁷ However, it should be noted that Baxter *et al.*³⁸ reported higher mobility values for undoped ZnO thin films and nanowires and showed that annealing enhances the mobility of the charge carriers ($140\text{--}230\text{ cm}^2\text{ V}^{-1}\text{ s}^{-1}$). The high mobility found for the sample with 0.5% Al content sintered at 900 °C ($53\text{ cm}^2\text{ V}^{-1}\text{ s}^{-1}$) suggests that even though this sample has grains of about 200 nm and that the spinel ZnAl_2O_4 phase is present as observed from the structural evaluation, the microstructure has little influence on the transport of the charge carriers in this sample. A further increase in the Al content, however, increases the amount of ZnAl_2O_4 nanoprecipitates, which scatter the charge carriers, thus decreasing the mobility. In the case of the samples sintered at 700 °C, the lower mobility values indicate that the microstructure of the specimens contains smaller grains, and therefore an increased fraction of grain boundaries more effectively scatters the charge carriers (see Fig. 5c and d). We cannot directly conclude that the increase in the volume fraction of the spinel ZnAl_2O_4 phase with Al-doping is solely responsible for the mobility reduction and in turn for the electrical conductivity because the amount of the spinel phase is considerably less than the observed diminution in the mobility of the sintered nanocomposites. As already seen from the dependence of the Seebeck coefficient and electrical conductivity on the doping concentration, we also find from their dependence on temperature that the nanostructure and composite nature has an impact on the electrical conductivity, but not on the Seebeck coefficient of the sintered samples, thus supporting the hypothesis by Nam *et al.*²⁶

The thermal conductivity κ of the nanocomposites is shown in Fig. 13. It rapidly decreases with the temperature for undoped ZnO, as also reported by Ohtaki *et al.*³ Doping with Al not only significantly reduces the thermal conductivity of the samples, but also modifies the temperature dependence of κ , as illustrated for 1% Al-doped specimen. It can also be seen that

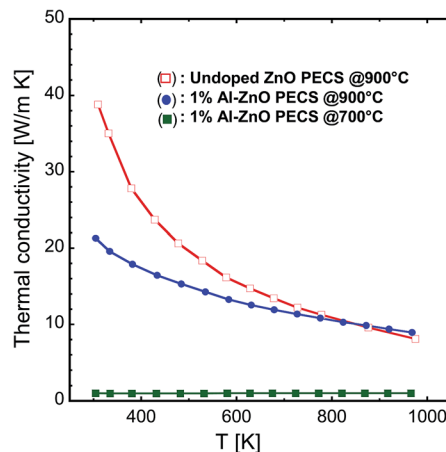


Fig. 13 Temperature dependence of the thermal conductivity of the nanocomposites sintered at 700 °C and 900 °C.



the sample sintered at 700 °C with 1% Al content displays the lowest κ of all the measured specimens and its behaviour is temperature independent. Kinemuchi *et al.*^{4,29} observed a similar dependence of κ on grain size and temperature. We conclude that phonon–phonon scattering is responsible for the decrease in thermal conductivity with increasing temperature in the case of undoped ZnO.^{35,39} In addition, in the case of the Al-doped nanocomposites, the significant reduction in phonon transport occurs because of grain boundary scattering at ZnO nanograins, ZnAl₂O₄ precipitates, point defects and pores present in the sample.^{29,36,39} We estimate the electronic contribution κ_e to the total thermal conductivity using the Wiedemann–Franz law, $\kappa_e = L\sigma T$, where L is the Lorenz number ($2.44 \times 10^{-8} \text{ V}^2 \text{ K}^{-2}$).⁴⁰ The electronic contribution to the thermal conductivity κ_e amounts to about 8–43% at 900 K for all the measured samples, implying that the major contribution to the total thermal conductivity is because of the lattice conductivity κ_L . Therefore, the nanocomposite microstructure of the samples significantly influences the thermal conductivity, and we are able to drastically reduce the lattice contribution to the total thermal conductivity, as reflected in the low κ values and the increasing contribution from the electronic carriers to the thermal conductivity.

Thermoelectric figure of merit

The numerical values for the dimensionless figure of merit zT resulting from the calculation using eqn (1) are shown in Fig. 14. These values have to be evaluated with care because thermal and electrical transports are measured in perpendicular directions. However, the dependence of zT on temperature features a considerably smooth slope, and this behaviour is assumed to be rather insensitive against possible anisotropy within the sample, allowing for a comparison amongst our samples and with temperature dependent data reproduced from Jood *et al.*³⁶ Because of the massive increase in electrical conductivity above 850 K, their zT shows a kink, whereas our

samples exhibit a considerable smoother dependence of zT on temperature. A figure of merit that is high over a wide temperature range is advantageous for power generation in large thermal gradients. The highest value for zT of 0.24 at 965 K is calculated for nanocomposites doped with 1% Al sintered at 700 °C, as illustrated in Fig. 14. This sample exhibits a slightly higher absolute Seebeck coefficient than all other doped samples. Together with an only slight decrease in the electrical conductivity compared to the sample with 0.5% Al-content sintered at 700 °C, this results in the highest power factor amongst the samples sintered at 700 °C, and coupled with a strong reduction in the thermal conductivity, results in the highest zT amongst all the sintered nanocomposites.

We have demonstrated that nanostructuring and the composite nature of the material significantly affect the thermal and electrical transports because of the scattering of phonons and charge carriers at the grain boundaries and nanoprecipitates, while only a minor influence is observed on the Seebeck coefficient because the concentration of active charge carriers does not significantly vary beyond the solubility limit of Al in ZnO. However, the Al content beyond the solubility limit can be used to tailor the nano-microstructure in such a fashion that the lattice contribution to the thermal conductivity is significantly reduced. Thus, we optimized the ratio of σ/κ to enhance zT of the material.

Conclusions

We synthesized Al-doped ZnO nanoparticles by chemical vapour synthesis in a hot-wall reactor. The powders were compacted into bulk nanocomposites by pulsed electric current sintering. The structural characterization shows that the presence of Al-dopant inhibits grain growth, both in the synthesized nanoparticles and in the sintered nanocomposites. The impact of Al-doping, nanograins/grain boundaries and ZnAl₂O₄ nanoprecipitates, *i.e.*, the microstructure of the nanocomposites, on the thermoelectric properties were demonstrated. With our processing methodology, we were able to create efficient Al-doping in ZnO, which produced bulk-like electrical conductivity and a strong reduction in the thermal conductivity of the nanocomposites at room temperature. We observed a strong correlation between the microstructure and the thermoelectric properties of the sintered nanocomposites. However, the detailed mechanism of the impact of the microstructure on the transport of the charge carriers requires further investigations. With this knowledge, a further tuning of the thermoelectric properties would be possible.

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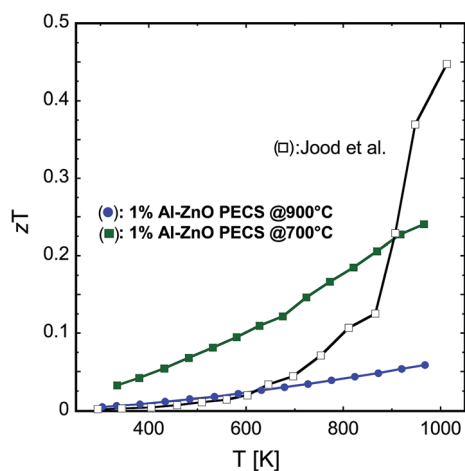


Fig. 14 Temperature dependence of the figure of merit zT for 1% AZO nanocomposites sintered at 700 °C and 900 °C. The present study is compared to the data from Jood *et al.* (see ref. 36).



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References

- M. S. Dresselhaus, G. Chen, M. Tang, R. Yang, H. Lee, D. Wang, Z. Ren, J. P. Fleurial and P. Gogna, *Adv. Mater.*, 2007, **19**, 1043.
- J. P. Heremans, C. M. Thrush and D. T. Morelli, *J. Appl. Phys.*, 2005, **98**, 063703.
- M. Ohtaki, T. Tsubota, K. Eguchi and H. Arai, *J. Appl. Phys.*, 1996, **79**, 1816.
- Y. Kinemuchi, H. Nakano, M. Mikami, K. Kobayashi, K. Watari and Y. Hotta, *J. Appl. Phys.*, 2010, **108**, 053721.
- J. He, M. G. Kanatzidis and V. P. Dravid, *Mater. Today*, 2013, **16**, 166.
- J. P. Heremans, B. Wiendlocha and A. M. Chamoire, *Energy Environ. Sci.*, 2012, **5**, 5510.
- W. Kim, R. Wang and A. Majumdar, *Nano Today*, 2007, **2**, 40.
- G. J. Snyder and E. S. Toberer, *Nat. Mater.*, 2008, **7**, 105.
- W. Liu, X. Yan, G. Chen and Z. Ren, *Nano Energy*, 2012, **1**, 42.
- M. Winterer, *Nanocrystalline Ceramics: Synthesis and Structure*, Springer Verlag, Berlin, 2002.
- U. Anselmi-Tamburini, J. E. Garay, Z. A. Munir, A. Tacca, F. Maglia and G. Spinolo, *J. Mater. Res.*, 2004, **19**, 3255.
- C. Schilling, M. Zähres, C. Mayer and M. Winterer, *J. Nanopart. Res.*, 2014, **16**(2506), 1.
- S. Hartner, M. Ali, C. Schulz, M. Winterer and H. Wiggers, *Nanotechnology*, 2009, **20**, 445701.
- L. Lutterotti and P. Scardi, *J. Appl. Crystallogr.*, 1990, **23**, 246 and MAUD (Material Analysis Using Diffraction) is a free Java program for Rietveld Analysis.
- M. H. Yoon, S. H. Lee, H. L. Park, H. K. Kim and M. S. Jang, *J. Mater. Sci. Lett.*, 2002, **21**, 1703.
- K. Shirouzu, T. Kawamoto, N. Enomoto and J. Hojo, *Jpn. J. Appl. Phys.*, 2010, **49**, 010201.
- B. Ingham, R. Linklater and T. Kemmitt, *J. Phys. Chem. C*, 2011, **115**, 21034.
- J. U. Brehm, M. Winterer and H. Hahn, *J. Appl. Phys.*, 2006, **100**, 064311.
- D. R. Clarke, *J. Am. Ceram. Soc.*, 1999, **82**, 485.
- Z. L. Wang, *J. Phys.: Condens. Matter*, 2004, **16**, R829.
- Y. Ding, X. Y. Kong and Z. L. Wang, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2004, **70**, 235408.
- W. S. Hong, L. C. De-Jonghe, X. Yang and M. N. Rahaman, *J. Am. Ceram. Soc.*, 1995, **78**, 3217.
- K. Hauffe and A. L. Vierk, *Z. Phys. Chem.*, 1950, **196**, 160.
- Y. Kinemuchi, C. Ito, H. Kaga, T. Aoki and K. Watari, *J. Mater. Res.*, 2007, **22**, 1942.
- T. Tsubota, M. Ohtaki, K. Eguchi and H. Arai, *Proceedings of the 16th International Conference on Thermoelectrics (ICT 1997)*, IEEE, Piscataway, 1997, p. 240.
- W. H. Nam, Y. S. Lim, S. M. Choi, W. S. Seo and J. Y. Lee, *J. Mater. Chem.*, 2012, **22**, 14633.
- G. G. Gadzhiev, *High Temp.*, 2003, **41**, 778.
- N. Vogel-Schäuble, R. Dujardin, A. Weidenkaff and M. H. Aguirre, *J. Electron. Mater.*, 2012, **41**, 1606.
- Y. Kinemuchi, M. Mikami, K. Kobayashi, K. Watari and Y. Hotta, *J. Electron. Mater.*, 2010, **39**, 2059.
- S. LeBlanc, S. Phadke, T. Kodama, A. Salles and K. E. Goodson, *Appl. Phys. Lett.*, 2012, **100**, 163105.
- J. Loureiro, N. Neve, R. Barros, T. Mateus, R. Santos, S. Filonovich, S. Reparaz, C. M. Sotomayor-Torres, F. Wyczisk, L. Divay, R. Martins and I. Ferreira, *J. Mater. Chem. A*, 2014, **2**, 6649.
- K. H. Kim, S. H. Shim, K. B. Shim, K. Niihara and J. Hojo, *J. Am. Ceram. Soc.*, 2005, **88**, 628.
- N. Ma, J. F. Li, B. P. Zhang, Y. H. Lin, L. R. Ren and G. F. Chen, *J. Phys. Chem. Solids*, 2012, **71**, 1344.
- Y. Fujishiro, M. Miyata, M. Awano and K. Maeda, *J. Am. Ceram. Soc.*, 2003, **86**, 2063.
- T. Tsubota, M. Ohtaki, K. Eguchi and H. Arai, *J. Mater. Chem.*, 1997, **7**, 85.
- P. Jood, R. Mehta, Y. Zhang, G. Peleckis, X. Wang, R. W. Siegel, T. Borca-Tasciuc, S. X. Dou and G. Ramanath, *Nano Lett.*, 2011, **11**, 4337.
- H. Kim, A. Piqué, J. S. Horwitz, H. Murata, Z. H. Kafafi, C. M. Gilmore and D. B. Chrisey, *Thin Solid Films*, 2000, **377–378**, 798.
- J. B. Baxter and C. A. Schmuttenmaer, *J. Phys. Chem. B*, 2006, **110**, 25229.
- E. S. Toberer, A. Zevkink and G. J. Snyder, *J. Mater. Chem.*, 2011, **21**, 15843.
- M. Ohtaki, K. Araki and K. Yamamoto, *J. Electron. Mater.*, 2009, **38**, 1234.

