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Protonation state of the Cu₄S₂ Cu_Z site in nitrous oxide reductase: redox dependence and insight into reactivity†

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Spectroscopic and computational methods have been used to determine the protonation state of the edge sulfur ligand in the Cu₄S₂ Cu_Z form of the active site of nitrous oxide reductase (N₂OR) in its 3Cu^ICu^{II} (1-hole) and 2Cu^I2Cu^{II} (2-hole) redox states. The EPR, absorption, and MCD spectra of 1-hole Cu_Z indicate that the unpaired spin in this site is evenly delocalized over Cu_I, Cu_{II}, and Cu_{IV}. 1-hole Cu_Z is shown to have a μ₂-thiolate edge ligand from the observation of S–H bending modes in the resonance Raman spectrum at 450 and 492 cm^{−1} that have significant deuterium isotope shifts (−137 cm^{−1}) and are not perturbed up to pH 10. 2-hole Cu_Z is characterized with absorption and resonance Raman spectroscopies as having two Cu–S stretching vibrations that profile differently. DFT models of the 1-hole and 2-hole Cu_Z sites are correlated to these spectroscopic features to determine that 2-hole Cu_Z has a μ₂-sulfide edge ligand at neutral pH. The slow two electron (+1 proton) reduction of N₂O by 1-hole Cu_Z is discussed and the possibility of a reaction between 2-hole Cu_Z and O₂ is considered.

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1. Introduction

The main reductive part of the nitrogen cycle, known as bacterial denitrification, is performed by soil and marine bacteria as a means of anaerobic or microaerobic respiration. Denitrification involves the conversion of nitrate to dinitrogen *via* four successive reductive steps (NO₃[−] → NO₂[−] → NO → N₂O → N₂), each performed by a different metalloenzyme.¹ The terminal product of denitrification can be either N₂O or N₂, depending on the regulatory control of the N₂O reduction process and whether the bacterium involved contains the gene cluster for nitrous oxide reduction (the *nos* cluster; *nosZ* encodes the nitrous oxide reductase enzyme).^{2,3} The N₂O reduction process and its regulation *in vivo* are of significant interest because N₂O is a potent greenhouse gas, with a global

warming potential 300× that of CO₂,^{4,5} and depletes the ozone layer.⁶ Anthropogenic sources of environmental N₂O, the majority of which is due to agricultural activity, is an increasing contribution to the global atmosphere.² Soil studies have indicated that pH,^{7,8} temperature,⁹ acetylene,¹⁰ sulfide,¹¹ and dioxygen¹² all affect the production of N₂O, but the molecular basis of these effects is still not known. A molecular understanding of nitrous oxide reduction and how this process is regulated could enable mitigation of N₂O release from anthropogenic sources.⁵

Nitrous oxide reductase contains two copper sites: a binuclear site known as Cu_A that functions as an electron transfer site, and an unusual tetranuclear copper sulfide cluster active site, where N₂O binds and is reduced (Fig. 1). Two forms of this tetranuclear site have been structurally characterized. One, known as Cu_Z^{*}, has a μ₄ sulfide ligand bridging all four coppers and a solvent derived ligand on an open edge (the Cu_I–Cu_{IV} edge) where N₂O is proposed to bind (Fig. 1A).¹³ This edge ligand has previously been assigned as a bridging hydroxide ligand, due to the presence of a vibration in the resonance Raman spectrum of Cu_Z^{*} that shifts in H₂¹⁸O solvent at high pH and the absence of significant spectroscopic differences between Cu_Z^{*} at high and low pH.¹⁵ The other form of the cluster, known as Cu_Z, has an additional μ₂ sulfur ligand bridging the Cu_I–Cu_{IV} edge (Fig. 1B).¹⁴ Whether the μ₂ edge ligand in Cu_Z is a thiolate (SH[−]) or a sulfide (S^{2−}) and how its protonation depends on the redox state of the cluster are not known. The Cu₄S₂ Cu_Z form of the cluster is dominantly isolated when N₂OR is purified in the absence of

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oxygen¹⁶ or rapidly in the presence of oxygen,¹⁷ while the Cu₄S Cu_Z^{*} form is isolated when the purification is performed aerobically or anaerobically from mutants in the accessory genes;^{16–18} however, all purifications typically yield enzyme with a mixture of the two sites.^{17,19} Which structural form of the cluster is responsible for N₂O reduction *in vivo* is a matter of some debate.^{20–22} As isolated, neither N₂OR containing a high percentage of Cu_Z nor N₂OR containing a high percentage of Cu_Z^{*} shows high enough specific activity in steady-state assays to be consistent with N₂OR activity in whole cells.^{17,23} N₂OR containing Cu_Z can be activated by prolonged dialysis against base,²³ while N₂OR that contains Cu_Z^{*} can be reductively activated by preincubation with methyl viologen, which reduces Cu_Z^{*} to the active fully reduced (4Cu^I) redox state.^{24,25} After activation, both Cu_Z and Cu_Z^{*} show specific activities consistent with whole cell N₂OR activity.^{20,23} However, it has recently been shown that the Cu_Z^{*} site in its fully reduced redox state is the form of the cluster that is responsible for the rapid N₂O reduction in steady state assays with methyl viologen, based on its rapid single turnover reaction with N₂O. Alternatively, in single turnover studies Cu_Z in its 1-hole redox state reduces N₂O but at a rate too slow to be catalytically relevant (10^{−6} that of the fully reduced state of Cu_Z^{*}).²⁶ Thus, the physiological role of the Cu_Z site in nitrous oxide reduction and whether it participates in N₂O reduction *in vivo* is unknown.

The Cu_Z site in nitrous oxide reductase has been extensively studied in N₂OR isolated from *Pseudomonas stutzeri* (PsN₂OR)^{27,28} and *Paracoccus pantotrophus* (PpN₂OR).^{19,29} In the latter enzyme, Cu_Z has been shown to access two redox states, the resting 2Cu^I2Cu^{II} (2-hole) redox state, and a 1 electron reduced 3Cu^I1Cu^{II} (1-hole) redox state (*E*⁰ = +60 mV).¹⁹ Both redox states of Cu_Z have previously been studied using EPR, absorption, MCD, and resonance Raman spectroscopies.^{19,28,30–33} However, these studies were performed before the elucidation of the presence of a second sulfur in the Cu_Z cluster, and so yielded limited direct insight into the cluster and the protonation state of the edge sulfur. Additionally, the previous studies were performed in the presence of background spectroscopic features from ~30% Cu_Z^{*}, which complicates the analysis.^{19,29} These limitations lead to the conclusion that Cu_Z and Cu_Z^{*} were very similar and perhaps differed only in the

second sphere.²⁸ These results are now extended and correlated to the structural insight that Cu_Z contains an additional inorganic sulfur edge ligand.¹⁴ An understanding of the protonation state, electronic structure, and potential reactivity of the Cu_Z site is necessary to gain insight into its reactivity and role *in vivo*.

This study uses EPR, absorption, MCD and resonance Raman spectroscopies coupled with DFT calculations to determine the protonation state of the edge sulfur ligand in the 1-hole and 2-hole redox states of Cu_Z in *Marinobacter hydrocarbonoclasticus* N₂OR (*Mh*N₂OR) and to define the electronic structures of these states. This leads to insight into the nature of the reactivity of the 1-hole and 2-hole states of Cu_Z and the origin of the spectroscopic similarity between 1-hole Cu_Z and 1-hole Cu_Z^{*}, despite significant differences in edge ligation in the two sites.

2. Methodology

2.1 Summary of experimental methodology

Full experimental methodology and computational details can be found in the ESI,[†] while a summary is presented here. Nitrous oxide reductase (N₂OR) was isolated from *Marinobacter hydrocarbonoclasticus* 617 (formerly *Pseudomonas nautica*) grown under microaerobic conditions in the presence of nitrate after two aerobic chromatographic steps without added reductant, as described previously.¹⁷ These purification conditions were shown to maximize the amount of Cu₄S₂ Cu_Z content relative to Cu₄S Cu_Z^{*} in the purified enzyme. Samples containing larger amounts of Cu_Z were purified in parallel with three chromatographic purification steps from a batch of cells grown under anaerobic conditions in the presence of nitrate, and that had been stored at −80 °C for a long period.^{17,26} Both *Mh*N₂OR samples showed copper quantitation results consistent with full occupancy of the Cu_A and Cu_Z/Cu_Z^{*} sites (6.4 ± 0.2 and 6.2 ± 0.7 respectively). The percentage of Cu_Z versus Cu_Z^{*} in the samples used for this study was determined by EPR spin quantitation (Fig. S1[†]). Samples purified with high amounts of Cu_Z contained 60 ± 10% Cu_Z, while samples purified to obtain more Cu_Z^{*} contained 10 ± 10% Cu_Z. Spectroscopic samples of 1-hole and 2-hole Cu_Z were prepared in a glove box under N₂ atmosphere. Samples of 1-hole Cu_Z were prepared from *Mh*N₂OR (60% Cu_Z and 40% Cu_Z^{*}) that had been incubated with 100 equivalents of reduced methyl viologen, with subsequent removal of the methyl viologen using a desalting column. Samples of 2-hole Cu_Z were prepared by reducing *Mh*N₂OR (60 ± 10% Cu_Z, 40 ± 10% Cu_Z^{*}) with 10 equivalents of sodium ascorbate, which reduces the Cu_A site rapidly and the 2-hole Cu_Z site very slowly, and spectra were collected within 1 hour so that minimal reduction of 2-hole Cu_Z was observed. In parallel, *Mh*N₂OR samples containing 90 ± 10% Cu_Z^{*} were reduced with 10 equivalents of sodium ascorbate to obtain the spectral features of 1-hole Cu_Z^{*}. For pH and deuteration studies, samples of 1-hole and 2-hole Cu_Z were buffer exchanged by centrifugation into different pH or pD buffers. Typical *Mh*N₂OR concentrations used for spectroscopic samples were 0.1–0.3 mM for absorption, MCD and EPR, and up to 0.5 mM for resonance Raman.

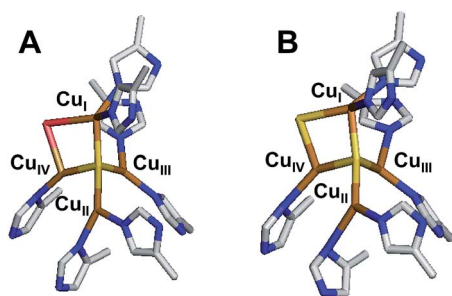


Fig. 1 The two forms of the tetranuclear copper sulfide active site of nitrous oxide reductase. (A) Cu_Z^{*} in PdN₂OR isolated aerobically (PDB ID 1FWX).¹³ (B) Cu_Z in PsN₂OR isolated anaerobically (PDB ID 3SBP).¹⁴



2.2 Computational modeling

A computational model of Cu_Z was built from the atomic coordinates of the crystal structure of *Pseudomonas stutzeri* N_2OR , the only known structure of the Cu_4S_2 cluster (PDB ID 3SBP, resolution 1.7 Å).¹⁴ The model included the Cu_4S_2 core and 7 ligating His residues, where the α carbon and distal nitrogen were constrained at their crystallographic positions. A computational model for Cu_Z^* with a hydroxide bridging ligand and identical α carbon and distal nitrogen constraints was constructed from the crystal structure of *Paracoccus denitrificans* N_2OR (PdN_2OR , PDB ID 1FWX).¹³ Calculations were performed using Gaussian 09 (version d01).³⁴ Geometry optimizations were performed using the B3LYP functional, the TZVP basis set on all core atoms (Cu_4S) and the ligating His nitrogens, and the SV basis set on all remaining atoms, and solvation was modeled with a PCM of 4.0. A larger basis set and different functionals were also explored, as described in the text. The optimized structures were then used for frequency, TD DFT, and single point calculations. To determine the relative energy of deprotonation ($\Delta\Delta E$) of the edge SH^- in the 2-hole *versus* 1-hole redox state, larger models were optimized that included two second sphere carboxylates, Asp127 and Asp240, which hydrogen bond to the His ligands of Cu_I and Cu_{II} . The energy of an internal proton transfer from the edge SH^- to Asp127 was calculated for the 1-hole and 2-hole redox states and compared to obtain the $\Delta\Delta E$.

3 Results and analysis

3.1 Spectroscopy of 1-hole Cu_Z

Previous spectroscopic studies of Cu_Z , undertaken before identification of the presence of a second sulfur, were performed on samples of PpN_2OR and PsN_2OR that contained mixtures of the Cu_Z and Cu_Z^* sites (in a 7 : 3 ratio for PpN_2OR) without a way to resolve the spectral features of the Cu_Z site from the mixture.^{19,28,29} Recently, it has been found that the two-sulfur Cu_Z site cannot be reduced by methyl viologen, which reduces both the Cu_A site and the Cu_Z^* form of the cluster.²⁶ This provides an opportunity to cleanly resolve the spectral features of 1-hole Cu_Z by studying methyl viologen reduced samples after removal of the reductant. This approach allows correlation of the electronic structure of 1-hole Cu_Z , obtained from spectroscopy, with the recently determined Cu_4S_2 structure of the cluster, to determine the nature of the edge sulfur ligand in its 1-hole and resting 2-hole redox states.

EPR. The X-band and Q-band EPR spectra of a methyl viologen reduced sample of 1-hole Cu_Z are given in Fig. 2. The EPR spectrum is axial with $g_{||} > g_{\perp} > 2.0$ and a pattern of five evenly spaced hyperfine lines in the $A_{||}$ region. The axial nature of the spectrum indicates that, while the spin density is delocalized over multiple copper nuclei, it resides in dominantly $d_{x^2-y^2}$ orbitals on each Cu site that contributes to the ground state. The $A_{||}$ hyperfine features can be further resolved in the second derivative of the X-band EPR spectrum, as can hyperfine features in the A_{\perp} region (Fig. 2A inset and S2†). Simulation of the X-band, X-band 2nd derivative, and Q-band EPR spectra

yields the g and A values for 1-hole Cu_Z given in Table 1. The g values for 1-hole Cu_Z are very similar to those previously obtained for 1-hole Cu_Z^* (Table 1) and to those obtained for Cu_Z in PpN_2OR .^{15,29} This is interesting, considering that an edge SH^- (thiolate) or S^{2-} (sulfide) in Cu_Z would be expected to be a more covalent ligand than the hydroxide in Cu_Z^* ¹⁵ and this would lower the g values. However, in 1-hole Cu_Z^* there is a high energy d-d transition that is not present in the 1-hole Cu_Z spectrum (*vide infra*). This transition has previously been assigned as a $d_{xy} \rightarrow d_{x^2-y^2}$ excitation localized on Cu_I .³⁵ The $g_{||}$ value is inversely proportional to the d_{xy} to $d_{x^2-y^2}$ energy splitting, so the presence of a high energy $d_{xy} \rightarrow d_{x^2-y^2}$ transition in Cu_Z^* but not in Cu_Z would lead to a lower $g_{||}$ value for Cu_Z^* than would be expected from covalency alone, which could result in similar $g_{||}$ values between 1-hole Cu_Z^* and the more covalent 1-hole Cu_Z site. The ligand field origin of the lower energy $d_{xy} \rightarrow d_{x^2-y^2}$ transition in 1-hole Cu_Z is considered below.

The $A_{||}$ and A_{\perp} values for 1-hole Cu_Z are similar in magnitude to those for Cu_Z^* , but fitting the hyperfine pattern requires three equivalent contributions rather than the $\sim 5 : 2$ ratio of hyperfine values observed for Cu_Z^* (Table 1).³⁶ This indicates that in the ground state of 1-hole Cu_Z the spin is distributed over three copper centers in dominantly $d_{x^2-y^2}$ orbitals. The three coppers involved are likely Cu_I , Cu_{II} , and Cu_{IV} , since these copper centers are in the same plane as the two sulfur ligands and bonding with the strong donor μ_4 sulfide and μ_2 sulfur ligands should define a common x, y plane for these coppers, with the z axis of the local g tensor of each copper oriented perpendicular to the Cu_3S_2 plane. This is consistent with the

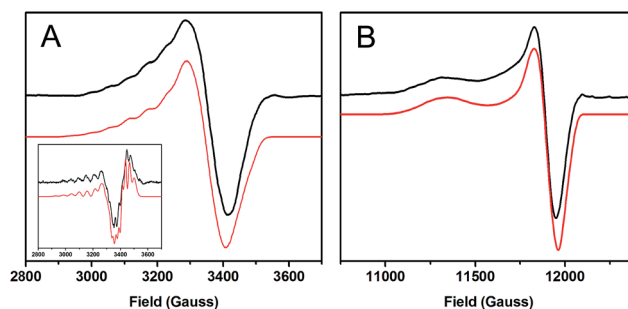


Fig. 2 EPR spectra of 1-hole Cu_Z (black) with simulations (red). (A) X-band at 77 K, 9.6349 GHz. Inset: 2nd derivative of the X-band. (B) Q-band at 77 K, 34.082 GHz.

Table 1 EPR g and A values for 1-hole Cu_Z , obtained from simulations included in Fig. 2 and S2 with values for Cu_Z^* reproduced from ref. 15

	1-hole Cu_Z	1-hole Cu_Z^*
$g_{ }$	2.152	2.160
$A_{ }$	$56 \times 10^{-4} \text{ cm}^{-1}$	$61 \times 10^{-4} \text{ cm}^{-1}$
	$56 \times 10^{-4} \text{ cm}^{-1}$	$23 \times 10^{-4} \text{ cm}^{-1}$
	$56 \times 10^{-4} \text{ cm}^{-1}$	
g_{\perp}	2.042	2.043
A_{\perp}	$20 \times 10^{-4} \text{ cm}^{-1}$	$25 \times 10^{-4} \text{ cm}^{-1}$
	$20 \times 10^{-4} \text{ cm}^{-1}$	$20 \times 10^{-4} \text{ cm}^{-1}$
	$20 \times 10^{-4} \text{ cm}^{-1}$	



axial nature of the g values and with the DFT calculations reported below.

Absorption and MCD. The low temperature absorption and MCD spectra of a methyl viologen reduced sample of 1-hole Cu_Z are presented in Fig. 3A. The absorption maximum of 1-hole Cu_Z occurs at $14\,600\text{ cm}^{-1}$ ($\epsilon \approx 3000\text{ M}^{-1}\text{ cm}^{-1}$), 1000 cm^{-1} lower than the absorption maximum of 1-hole Cu_Z^* (Fig. 3B). There are no additional low energy intense absorption features due to the edge sulfur. The low temperature absorption and MCD spectra can be simultaneously fit to yield a total of 11 transitions, which can be assigned by considering their energies and C_0/D_0 ratios, following ref. 15 (Table S1†). Comparison of the transition assignments and energies of the 1-hole Cu_Z and 1-hole Cu_Z^* sites reveals some key differences. While the absorption maximum of Cu_Z occurs at lower energy than that of Cu_Z^* , from MCD the three $\mu_4\text{S}$ to Cu CT transitions, assigned in Cu_Z^* ,³⁵ occur at very similar energies in the two sites (bands 5, 6, and 7, numbering given in Fig. 3). The shift in the absorption maximum therefore arises from a different intensity pattern for these transitions, where in Cu_Z the lowest energy transition at $14\,600\text{ cm}^{-1}$ is the most intense (band 5) and the transition at $15\,600\text{ cm}^{-1}$ is weaker (band 6), but in Cu_Z^* this is reversed. The $\mu_4\text{S}^{2-}$ to Cu CTs in Cu_Z^* have previously been assigned as transitions from the three different 3p orbitals of the $\mu_4\text{S}^{2-}$ to the β LUMO of the cluster. From our current study of resting 1-hole Cu_Z^* and Cu_Z with different edge ligands,^{15,35} the β LUMO is delocalized in the plane that contains Cu_I , Cu_II , Cu_IV and the $\mu_4\text{S}^{2-}$, with different amounts of spin distributed over Cu_I , Cu_II , and Cu_IV depending on the edge ligation. Two of the $\mu_4\text{S}^{2-}$ p orbitals are in the plane, oriented between Cu_I and Cu_IV ($\text{S } p_x$) and between Cu_IV and Cu_II ($\text{S } p_y$), while the third is perpendicular to the plane ($\text{S } p_z$). Scheme S1† reflects the orientation and simplified composition of these orbitals determined for 1-hole Cu_Z^* from DFT calculations. The CT intensities reflect the overlap of these three S p orbitals with the β LUMO. Since bands 5 and 6 show the highest intensity in the 1-hole forms of Cu_Z and Cu_Z^* , these must reflect charge transfer from the in-plane $\text{S } p_x$ (band 6, dominant in Cu_Z^* due to higher overlap with Cu_I) and $\text{S } p_y$ (band 5) orbitals. Bands 6 and 5 form a pseudo-A feature in the MCD spectrum (*i.e.* derivative-shaped) and thus must arise from two transitions with orthogonal transition moments that spin-orbit couple in a third, mutually perpendicular direction (*i.e.* L_z). Since band 6 arises from a transition to Cu_I (from its dominant intensity in Cu_Z^*), band 5 must reflect a transition to Cu_IV , since the Cu_I -S and Cu_IV -S bonds are close to perpendicular (96° from crystallography) while the Cu_I -S and Cu_II -S bonds are close to parallel (160°).¹³ The change in relative intensities of the $\mu_4\text{S}^{2-}$ to Cu CT transitions in Cu_Z relative to Cu_Z^* , where band 6 decreases in intensity while band 5 increases in intensity, thus indicates that there is less spin on Cu_I and more spin on Cu_IV in 1-hole Cu_Z relative to 1-hole Cu_Z^* . This is consistent with the EPR hyperfine values, which suggest that the spin in Cu_Z is delocalized 1 : 1 : 1 over Cu_I , Cu_II , and Cu_IV , while from ref. 35 in Cu_Z^* the spin is delocalized $\sim 5 : 2$ over Cu_I and Cu_IV . Additionally, in Cu_Z^* , a band at $18\,000\text{ cm}^{-1}$ (band 8) was assigned as a high energy d-d transition due to its high C_0/D_0 ratio; this was assigned as a localized $d_{xy} \rightarrow d_{x^2-y^2}$ transition on Cu_I , where most of the 1-hole

is localized.³⁵ No equivalent high energy d-d transition is observed in the MCD spectrum of Cu_Z . The lower energy of the d-d transitions in Cu_Z relative to Cu_Z^* is likely due to the decreased spin on Cu_I , the only four coordinate site, relative to Cu_II and Cu_IV , which are both 3 coordinate and have a weaker ligand field.

Resonance Raman. The resonance Raman spectrum of 1-hole Cu_Z and the enhancement profiles of the vibrations are presented in Fig. 4A and B, respectively. Seven vibrations are enhanced in the most intense S to Cu CT transition (band 5), including three intense vibrations at 203 , 378 , and 492 cm^{-1} . The vibration at 378 cm^{-1} occurs at the same energy as a Cu-S

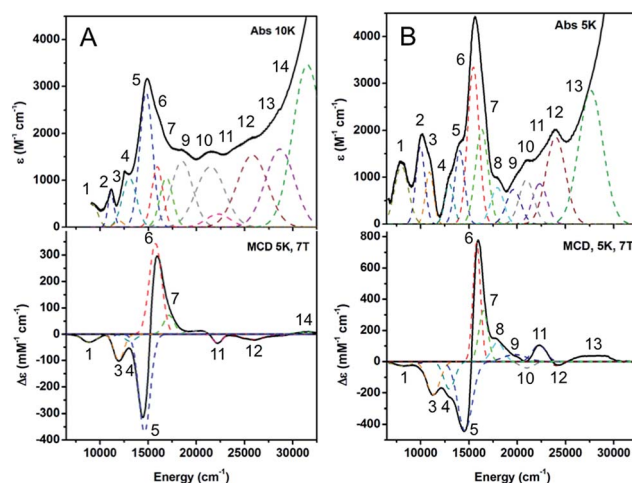


Fig. 3 Low temperature absorption and MCD spectra of (A) 1-hole Cu_Z , 10 K absorption, 5 K and 7 T MCD. (B) 1-hole Cu_Z^* , 5 K absorption, 5 K and 7 T MCD (adapted from ref. 15).

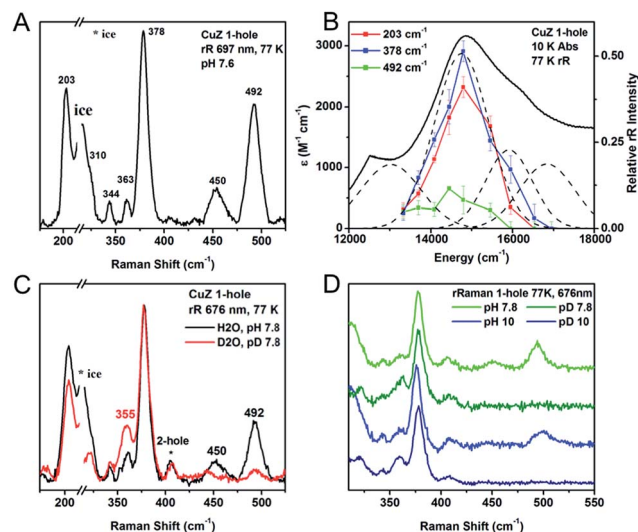


Fig. 4 (A) Resonance Raman spectrum of 1-hole Cu_Z at 77 K, excitation energy 697 nm. (B) Excitation profile of the 203 , 378 , and 492 cm^{-1} vibrations. (C) H/D isotope shift of the vibrations of 1-hole Cu_Z , performed in pH or pD 7.8, 100 mM phosphate, excitation energy 676 nm. (D) Comparison of SH bending vibrations at pH/pD 7.8 (green) and pH/pD 10 (blue).



stretch of the Cu_2^* site (Fig. S3†) and the previously reported ^{34}S isotope sensitivities of both vibrations are similar (-5.8 and -4.7 cm^{-1} , respectively),²⁸ indicating that the 378 cm^{-1} vibration in 1-hole Cu_Z can be assigned as a Cu-S vibration of the μ_4 sulfide. In contrast, the 203 cm^{-1} vibration is significantly lower in energy than the vibrations of Cu_2^* , and thus can be assigned as a Cu-S vibration of the μ_2 sulfur ligand that is only present in Cu_Z . Further, there are two high energy vibrations in Cu_Z at 450 and 492 cm^{-1} that show significant deuterium isotope sensitivity, shifting down in energy by -137 cm^{-1} (for the 492 cm^{-1} vibration) in deuterated buffer (Fig. 4C). This shift requires their assignment as S-H bending modes. Thus, we can definitively identify the edge ligand in 1-hole Cu_Z as a $\mu_2\text{SH}^-$. The S-H bending modes at 492 and 450 cm^{-1} are present at both pH 7.8 and pH 10 (Fig. 4D), indicating that the pK_a of the edge thiolate is ~ 11 or higher. This is further supported by the lack of pH dependence observed in the MCD and EPR spectra of 1-hole Cu_Z between pD 6 and pD 10 (Fig. S4†). Since the second pK_a of free hydrogen sulfide in water is 12, a pK_a range of 11–12 can be estimated for the edge thiolate ligand in 1-hole Cu_Z .

3.2 Spectroscopy of 2-hole Cu_Z

Absorption. The 2-hole redox state has been previously shown to be the resting redox state of Cu_Z . 2-hole Cu_Z is diamagnetic from MCD.¹⁹ The absorption features of 2-hole Cu_Z in as-isolated N_2OR are present with additional spectral contributions from oxidized Cu_A and some amount of 1-hole Cu_Z .²⁶ To remove these contributions, the absorption spectrum of 2-hole Cu_Z (Fig. 5) was obtained after reduction with sodium ascorbate, which reduces the Cu_A site faster than it reduces 2-hole Cu_Z , and subtraction of the spectral contribution of 1-hole Cu_Z , obtained from a separately purified N_2OR sample containing $90 \pm 10\%$ Cu_Z^* . An intense absorption maximum for 2-hole Cu_Z is observed at $18\,300\text{ cm}^{-1}$ ($\epsilon \approx 10\,000\text{ M}^{-1}\text{ cm}^{-1}$) with a weaker low energy shoulder, consistent with absorption spectra previously reported for ascorbate reduced samples containing high amounts of 2-hole Cu_Z .^{19,32} The low temperature absorption spectrum of 2-hole Cu_Z (Fig. S5†) resolves this absorption maximum into three distinct absorption bands. Simulation of the room temperature absorption spectrum with transition energies derived from the low temperature spectrum distinguishes five transitions, all with absorption intensities higher than $1000\text{ M}^{-1}\text{ cm}^{-1}$, sufficiently intense to be S to Cu CT transitions from the $\mu_4\text{S}^{2-}$ or the $\mu_2\text{S}$ ligand (Fig. 5). The two most intense transitions (bands 2 and 3) are either from different ligands ($\mu_2\text{S}$ and $\mu_4\text{S}^{2-}$) or from the same ligand to two different acceptor orbitals (the α and β holes of the broken symmetry singlet ground state). Based on the correlation of resonance Raman excitation profiles of the vibrations of 2-hole Cu_Z to DFT calculations (*vide infra*), the assignment of the two transitions as CT transitions from the $\mu_4\text{S}^{2-}$ to two different holes is preferred.

Resonance Raman. The resonance Raman spectrum of 2-hole Cu_Z was obtained upon excitation into the intense absorption maximum at $18\,300\text{ cm}^{-1}$ (Fig. 6A). Two vibrations are enhanced at 350 and 405 cm^{-1} . The ^{34}S isotope shifts of these vibrations have been previously reported to be -5.6 and -5.8 cm^{-1} , respectively, indicating that they are Cu-S stretches.²⁸ In contrast

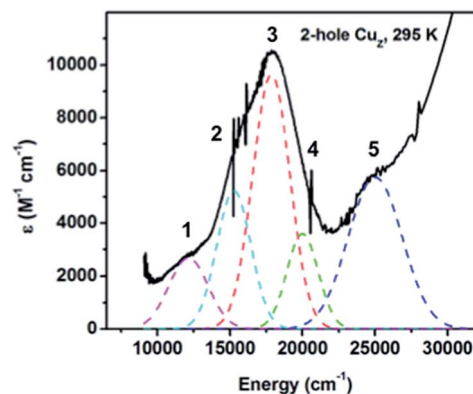


Fig. 5 Absorption spectrum of 2-hole Cu_Z at room temperature, obtained after ascorbate reduction of Cu_A and subtraction of spectral contribution of 1-hole Cu_Z^* .

to 1-hole Cu_Z , no higher energy S-H bending vibration is observed (up to 800 cm^{-1}). The excitation profile of the Cu-S stretching vibrations shows that they are enhanced differently in the most intense absorption bands 2 and 3 (Fig. 6B). The lower energy vibration at 350 cm^{-1} is enhanced in both transitions, while the higher energy vibration at 405 cm^{-1} is dominantly enhanced in the lower energy transition (band 2) and only weakly enhanced in band 3. This difference in profiling behavior is consistent with the Cu-S vibrations obtained computationally for a Cu_4S_2 cluster with a $\mu_2\text{S}^{2-}$ and $\mu_4\text{S}^{2-}$ and with the predicted enhancements of key vibrations in transitions from the $\mu_4\text{S}^{2-}$ to the α and β holes (see 3.3).

The resonance Raman spectrum 2-hole Cu_Z shows no significant shift in the energies of the 350 and 405 cm^{-1} vibrations between pD 6 and pD 10 (Fig. S6†). This suggests that the edge ligand has a pK_a either lower than 5.5 or higher than 10.5. A pK_a higher than 10.5 in the 2-hole redox state is not consistent with observed pK_a of 11–12 for 1-hole Cu_Z , as the increased charge of the 2-hole state will lead to a lower pK_a relative to the 1-hole redox state. The possibility of a pK_a less than 5.5 for 2-hole Cu_Z is evaluated computationally below.

3.3 Calculations

1-hole Cu_Z . A computational model of 1-hole Cu_Z was constructed based on the crystal structure of *Pseudomonas stutzeri*

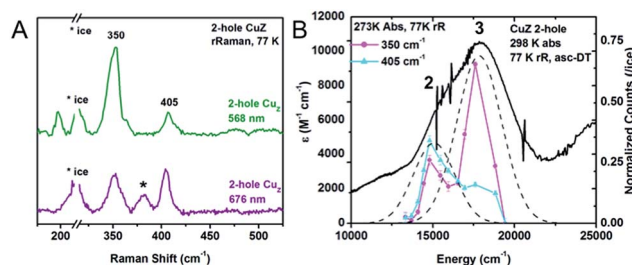


Fig. 6 (A) Resonance Raman spectra of 2-hole Cu_Z at 77 K and two excitation energies, 568 nm (green) and 676 nm (purple). Starred vibration is due to 1-hole Cu_Z^* . (B) Excitation profiles of the 350 and 405 cm^{-1} vibrations overlaid with the room temperature absorption spectrum.



N₂OR (PDB ID 3SBP, resolution 1.7 Å).¹⁴ On the basis of the resonance Raman data, the edge sulfur was modeled as an SH[−] ligand bridging the Cu_I–Cu_{IV} edge (Fig. 7A). This will be compared to an experimentally validated model of the Cu_Z^{*} site, which has a hydroxide ligand bridging the Cu_I–Cu_{IV} edge (Fig. 7B).¹⁵ The optimized structure of the 1-hole SH[−] cluster agrees well with the bond lengths and angles observed in the crystal structure (2.35 Å and 2.48 Å for the Cu_I–μ₂SH[−] and Cu_{IV}–μ₂SH[−] bonds computationally, relative to 2.61 Å and 2.49 Å crystallographically with a resolution of 1.7 Å,¹⁴ Table S2†). Since the crystal was grown from the “purple” resting form of PsN₂OR, containing the resting 2-hole redox state of the Cu_Z site, the Cu_Z site in the crystal may have some photo-reduction due to exposure to X-ray radiation.³⁷ The calculated structures and spin distributions are not significantly perturbed when a triple zeta basis set is used on all His ring atoms (Tables S4 and S5†). Including the second sphere residues Lys397 and Glu435 in the computational model also does not affect the structure or spin distribution, consistent with the small effect on the spectral features of 1-hole Cu_Z^{*} observed experimentally upon deprotonation of Lys397.¹⁵ Thus, the structures including only first sphere ligands were used to model the Cu_Z and Cu_Z^{*} sites in this study.

The 1-hole model with an SH[−] edge ligand reproduces the key spectral features observed for the 1-hole Cu_Z site. The Mulliken atomic spin distribution of the cluster with an SH[−] edge ligand is delocalized over Cu_I, Cu_{II}, and Cu_{IV} in a 2 : 1 : 1 ratio. In going from Cu_Z^{*} to Cu_Z the calculated spin on Cu_I changes from 26% to 17% (Table 2), which is consistent with the decrease in intensity of band 6 observed in the absorption and MCD data for 1-hole Cu_Z and leads to a more equal distribution of spin over Cu_I, Cu_{II}, and Cu_{IV}, consistent with the EPR hyperfine values. The LUMO of the Cu_Z model contains d_{x²−y²} character on Cu_I, Cu_{II}, and Cu_{IV}, which are aligned, consistent with the ground state predicted from the EPR *g* values (Fig. S7†). It also contains significant antibonding μ₄S^{2−} and μ₂SH[−] character, explaining why Cu–S stretching vibrations of both the μ₄S^{2−} and μ₂SH[−] are enhanced in the charge transfer transitions to this acceptor orbital. Additionally, the computational model predicts the Cu_Z site to be more covalent than the Cu_Z^{*} site, with 10% less Cu character in the ground state wavefunction, reflecting delocalization of the

spin from Cu_I onto the edge SH[−] ligand. The low *g*_{||} value for 1-hole Cu_Z is also predicted by the computational model (Table S7†). However, in contrast to experiment, the calculated *g*_{||} values for the Cu_Z and Cu_Z^{*} models differ, with a higher calculated *g*_{||} value for Cu_Z^{*} than that observed experimentally. This suggests that the calculated model of Cu_Z^{*} does not accurately predict the ligand field on Cu_I that leads to the higher energy d_{xy} → d_{x²−y²} transition observed experimentally.

The Cu–S stretching vibrations and S–H bending vibrations for the 1-hole SH[−] model of Cu_Z are given in Table S8 and Fig. S8.† The model predicts two S–H bending modes at 426 and 461 cm^{−1} with H/D isotope shifts of −125 cm^{−1} and −123 cm^{−1}, respectively, similar to the vibrations observed experimentally at 450 and 492 cm^{−1} (with a shift of −137 cm^{−1} for the 492 cm^{−1} vibration; the 450 cm^{−1} vibration cannot be observed after deuteration due to overlap with the ice scattering peak). Equivalent O–H bends are predicted for the OH bridged Cu_Z^{*} model at higher energies, but these are not experimentally observed. The Cu_Z model also predicts the presence of a low energy Cu–S stretching vibration of the μ₂SH[−] (178 cm^{−1}, observed at 203 cm^{−1} experimentally) and both models show similar energies for the Cu–μ₄S stretching vibrations. The absolute energies of the Cu–S stretching vibrations for both the μ₄ sulfide and μ₂ thiolate are underestimated, as has been found for computational models of the Cu_Z^{*} site.^{15,35} The TD DFT calculated absorption spectrum for the Cu_Z model is also very similar to the calculated absorption spectrum for the Cu_Z^{*} model both with B3LYP and with the functional B98, which has been shown to predict the experimental absorption spectrum of a Cu₃S₂ model complex reasonably well.³⁸ Interestingly, neither the experimental absorption spectrum nor the TD-DFT calculation predicts an intense low energy charge transfer transition from the μ₂SH[−] ligand (Fig. S9†). While some weak transitions predicted computationally at lower energy than the μ₄S^{2−} to Cu CT transitions have μ₂SH[−] to Cu CT character, they are predicted to lack intensity and are thus difficult to distinguish from the Cu d to d transitions that are also observed in this energy region.

Thus, a computational model of the tetranuclear copper cluster with an SH[−] edge ligand bridging Cu_I and Cu_{IV} provides a good structural model of 1-hole Cu_Z that reproduces its key spectral features. This spectroscopically calibrated model was then extended to the 2-hole redox state of the Cu_Z site, for which less experimental data are accessible.

2-hole Cu_Z. Two possible computational models were developed for 2-hole Cu_Z, one with an edge thiolate ligand (Cu₄S(SH)) and one with an edge sulfide (Cu₄S₂). These were optimized in both the triplet (*S* = 1) and broken symmetry

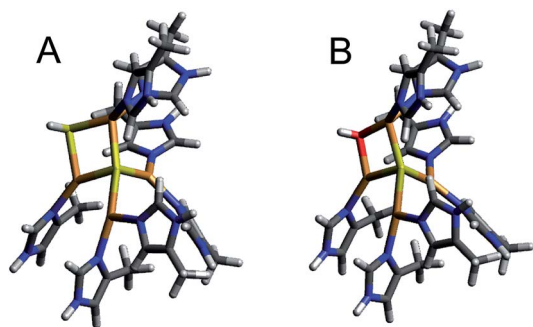


Fig. 7 Computational models of (A) 1-hole Cu_Z and (B) 1-hole Cu_Z^{*} (B3LYP, TZVP on Cu, S, and ligating N atoms, and SV on all remaining atoms, PCM of 4.0).

Table 2 Mulliken atomic spin density of 1-hole computational models with SH[−] and OH[−] bridging ligands on the Cu_I–Cu_{IV} edge, with % Cu and % d orbital character in their ground state wavefunctions

Edge ligand	Mulliken atomic spin density					
	Cu _I	Cu _{II}	Cu _{III}	Cu _{IV}	μ ₄ S ^{2−}	μ ₂ L [−]
SH [−] bridge	0.17	0.11	0.06	0.10	0.34	0.16
OH [−] bridge	0.26	0.09	0.04	0.13	0.31	0.10



singlet ($S = 0$) ground spin states. For both models the singlet is lower in electronic energy, by $-8.0 \text{ kcal mol}^{-1}$ for the sulfide and $-3.4 \text{ kcal mol}^{-1}$ for the thiolate (spin corrected energies using B3LYP). The singlet state was verified to be the ground state using a variety of functionals, including M06L, M06, and TPSSH. Thus, both structures would be consistent with the experimentally determined singlet ground state of 2-hole Cu_2 .¹⁹ The optimized structure of the 2-hole $\text{Cu}_4\text{S}(\text{SH})$ model is similar to that of the 1-hole SH^- model of 1-hole Cu_2 , with slightly shorter $\text{Cu}_\text{I}/\text{Cu}_\text{IV}-\text{SH}^-$ and $\text{Cu}_\text{I}/\text{Cu}_\text{IV}-\mu_4\text{S}$ bonds (Table S9†). The α LUMO is dominantly localized on Cu_I and has equal $\mu_4\text{S}^{2-}$ and $\mu_2\text{SH}^-$ antibonding character (Table 3) while the β LUMO is delocalized equally over Cu_II and Cu_IV and has more $\mu_4\text{S}^{2-}$ antibonding character. Upon deprotonation of the edge SH^- , the 2-hole Cu_4S_2 model has significantly shorter bonds between the edge sulfide and $\text{Cu}_\text{I}/\text{Cu}_\text{IV}$ and similar $\mu_4\text{S}^{2-}-\text{Cu}$ bond lengths to the 2-hole SH^- model (Table S9†). In this model, the α LUMO is localized on Cu_I while the β LUMO is localized on Cu_IV (Fig. 8 and Table 3). Both holes have significant $\mu_2\text{S}^{2-}$ character, indicating that the edge sulfide copper bonds are highly covalent.

The energy of deprotonation calculated from the 2-hole models was compared to the calculated energy of deprotonation of the 1-hole SH^- model, where the experimentally estimated $\text{p}K_\text{a}$ of the edge thiolate is 11–12 (*vide supra*). Examination of the energy required to deprotonate the 1-hole and 2-hole SH^- models shows that deprotonation of the 1-hole is not energetically favored ($\Delta E = 26 \text{ kcal mol}^{-1}$, relative to an energy of $-268 \text{ kcal mol}^{-1}$ for a solvated proton)³⁹ while deprotonation of the 2-hole is favorable ($\Delta E = -9 \text{ kcal mol}^{-1}$). However, the two models have different charges (+2/+1 and +3/+2 for protonated and deprotonated 1-hole and 2-hole models, respectively) and this will significantly affect the relative energies of deprotonation. To minimize the charge effect, the computational models were expanded to include two second sphere Asp residues near the Cu_2 site, such that the 1-hole Cu_2 model with an edge thiolate is neutral and the 2-hole Cu_2 $\text{Cu}_4\text{S}(\text{SH})$ model has a +1 charge. The proton transfer was performed internally to one of the Asp residues, so the total charge of the model does not change upon deprotonation (Fig. S10†). The $\Delta\Delta E$ for deprotonation of the 2-hole edge thiolate relative to the 1-hole species is calculated to be $-25 \text{ kcal mol}^{-1}$ (with a dielectric of 4.0). This value is dependent on the dielectric (Fig. S11†) and, at high dielectric values, converges to a $\Delta\Delta E$ of $-12 \text{ kcal mol}^{-1}$. To

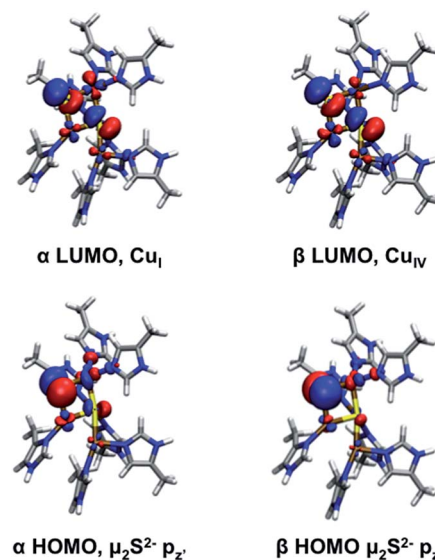


Fig. 8 Frontier molecular orbitals of 2-hole Cu_2 (two sulfide model).

estimate the difference in $\text{p}K_\text{a}$ between 1-hole and 2-hole Cu_2 , the $\Delta\Delta G$ was estimated from the $\Delta\Delta E$ using frequency calculations for structures with identical fixed atom constraints and thus the same number and magnitude of imaginary frequencies (these ΔG corrections vary by only $0.3 \text{ kcal mol}^{-1}$ between the protonated and deprotonated 1-hole structures). This gives a $\Delta\Delta G$ of $-12 \text{ kcal mol}^{-1}$ for deprotonation of 2-hole *versus* 1-hole Cu_2 , which corresponds to a $\Delta\text{p}K_\text{a}$ of -9 . Given the experimental $\text{p}K_\text{a}$ value of 11–12 for the edge thiolate in 1-hole Cu_2 , these calculations predict a $\text{p}K_\text{a}$ for a thiolate in 2-hole Cu_2 of 3 or less, consistent with the absence of a pH effect in resonance Raman of 2-hole Cu_2 at pD 6. This strongly suggests that 2-hole Cu_2 is a two sulfide cluster at neutral pH.

The calculated spectral features for the Cu_4S_2 2-hole Cu_2 model can be compared with those determined experimentally. The TD DFT predicted absorption spectrum (using both B3LYP and B98) is qualitatively similar to the experimental absorption spectrum, showing two intense absorption maxima with a higher energy shoulder (Fig. S12†). The predicted vibrations of the 2-hole Cu_4S_2 model are given in Table S10 and Fig. S13.† All of the calculated vibrations are shifted up in energy in comparison to those of the 1-hole SH^- model, with the most significant energy differences observed for the sulfur edge vibrations, due to the short and highly covalent $\text{Cu}-\mu_2\text{S}^{2-}$ bonds in the 2-hole Cu_4S_2 cluster. In particular, the $\mu_2\text{S}^{2-}-\text{Cu}_\text{I}$ stretch now occurs at a similar energy to and mixes with vibrations of the $\mu_4\text{S}^{2-}$, leading to symmetric and antisymmetric combinations of the $\mu_2\text{S}^{2-}-\text{Cu}_\text{I}$ and $\mu_4\text{S}^{2-}-\text{Cu}_\text{I}$ stretches (predicted at 312 and 309 cm^{-1} , respectively, see ESI†). The symmetric combination is allowed in resonance Raman and will be enhanced in all transitions due to the high amount of $\mu_2\text{S}^{2-}$ character in both the α and β holes. This is a good candidate for the 350 cm^{-1} vibration observed experimentally that profiles in both intense absorption bands (see Fig. 6B). The highest energy core vibration of the 2-hole $\mu_2\text{S}^{2-}$ cluster is a symmetric $\text{Cu}_\text{II}-\mu_4\text{S}^{2-}-\text{Cu}_\text{IV}$ stretch predicted at 344 cm^{-1} which will be

Table 3 Mulliken spin densities on Cu and S atoms in the α and β LUMOs of the broken symmetry singlet 2-hole $\text{Cu}_4\text{S}(\text{SH})$ and Cu_4S_2 models (B3LYP, TZVP on Cu, S, and ligating N atoms, and SV on all remaining atoms, PCM of 4.0)

		Mulliken spin density					
		$\mu_2\text{L}$	Cu_I	Cu_II	Cu_III	Cu_IV	$\mu_4\text{S}^{2-}$
2-Hole SH^- $S = 0$	α LUMO	0.16	0.37	0.06	0.06	0.04	0.20
	β LUMO	0.17	0.03	0.17	0.09	0.16	0.30
2-Hole S^{2-} $S = 0$	α LUMO	0.33	0.22	0.06	0.04	0.07	0.23
	β LUMO	0.37	0.05	0.06	0.05	0.14	0.27



selectively enhanced in a transition to the β LUMO localized on Cu_{IV} (Fig. 8). A symmetric $\text{Cu}_{\text{II}}-\mu_4\text{S}^{2-}-\text{Cu}_{\text{IV}}$ stretch is also computationally predicted in the 1-hole SH^- model at 320 cm^{-1} and the calculated shift in energy of this mode between the 1-hole and 2-hole models ($+24\text{ cm}^{-1}$) is similar to the energy increase of the highest energy Cu-S stretches observed experimentally in 1-hole and 2-hole Cu_Z (378 and 405 cm^{-1} respectively, $\Delta\nu$ of $+27\text{ cm}^{-1}$). Thus, the 2-hole Cu_4S_2 model qualitatively predicts a high energy Cu-S vibration that will be selectively enhanced only in a transition to the β hole and a lower energy Cu-S vibration that will be enhanced in both intense transitions. This is consistent with the enhancement profiles of the two vibrations observed experimentally in Fig. 6B. This establishes that a $\mu_2\text{S}^{2-}$ bridge is energetically favored and consistent with the spectral features of 2-hole Cu_Z .

4. Discussion

A combination of spectroscopic methods and DFT calculations has been used to define the protonation state of the μ_2 sulfur ligand on the $\text{Cu}_\text{I}-\text{Cu}_\text{IV}$ edge in 1-hole and 2-hole Cu_Z . This leads to insight into the spectroscopic similarities between 1-hole Cu_Z and 1-hole Cu_Z^* , the redox reactivity of 1-hole Cu_Z in the slow 2 electron reduction of N_2O , and the interconversion between Cu_Z and Cu_Z^* , the reactive form of the cluster for N_2O reduction *in vitro*.

4.1 Protonation states of 1-hole and 2-hole Cu_Z

The protonation state of the edge ligand in 1-hole Cu_Z has been directly determined by resonance Raman spectroscopy. Two high energy vibrations are enhanced in the most intense $\mu_4\text{S}^{2-}$ to Cu CT transition of 1-hole Cu_Z , at 450 and 492 cm^{-1} , and have large isotope shifts upon solvent deuteration (-137 cm^{-1} for the 492 cm^{-1} mode). This is consistent with S-H bending modes, indicating that the $\mu_2\text{S}$ ligand is a thiolate. The energy and solvent isotope shift of these S-H bending modes are as predicted by DFT calculations for a model with a $\mu_2\text{SH}^-$ bridging the $\text{Cu}_\text{I}-\text{Cu}_\text{IV}$ edge. The EPR spectrum of 1-hole Cu_Z indicates a ground state in which the spin is delocalized over 3 coppers in dominantly $d_{x^2-y^2}$ orbitals. The absorption and MCD spectra show three $\mu_4\text{S}^{2-}$ to Cu charge transfer transitions that have very similar energies to those observed for 1-hole Cu_Z^* (which has a hydroxide bridged $\text{Cu}_\text{I}-\text{Cu}_\text{IV}$ edge) but a different intensity pattern, consistent with a change in spin distribution in the cluster from dominantly on Cu_I in 1-hole Cu_Z^* to more evenly delocalized over Cu_I , Cu_II , and Cu_IV . This ground state spin distribution is consistent with that predicted from DFT calculations for the $\mu_2\text{SH}^-$. Based on the absence of a pH effect in 1-hole Cu_Z up to a pH of 10, the pK_a of the edge thiolate in 1-hole Cu_Z is estimated to be 11–12.

The 2-hole state of Cu_Z was also spectroscopically defined, but no direct spectroscopic evidence for the protonation state of the edge ligand was obtained. DFT calculations of the deprotonation of a $\mu_2\text{SH}^-$ ligand in 2-hole Cu_Z relative to 1-hole Cu_Z were used to determine that 2-hole Cu_Z likely has a sulfide edge ligand. Deprotonation of a $\mu_2\text{SH}^-$ ligand in the 2-hole redox

state is at least 12 kcal mol^{-1} more favorable than in the 1-hole redox state, after accounting for charge and dielectric effects. This yields a calculated pK_a for a $\mu_2\text{SH}^-$ ligand in 2-hole Cu_Z of 3 or less, which strongly suggests that 2-hole Cu_Z has an edge sulfide ligand at physiological pH. The calculated spectroscopic properties of a model of 2-hole Cu_Z with a $\mu_2\text{S}^{2-}$ ligand are also consistent with those observed experimentally.

4.2 Similarities between 1-hole Cu_Z and 1-hole Cu_Z^*

It has previously been observed that the spectral features of 1-hole Cu_Z are rather similar to those of 1-hole Cu_Z^* , despite the change in the nature of the edge ligand from a thiolate to a hydroxide.^{19,28} The spectroscopic similarities between 1-hole Cu_Z and 1-hole Cu_Z^* reflect similar bonding interactions between the $\mu_4\text{S}^{2-}$ and the in plane coppers (Cu_I , Cu_II , and Cu_IV) which are not significantly perturbed by the nature of edge ligand. This results in similar transition energies in the absorption and MCD spectra, as the dominant transitions are due to $\mu_4\text{S}^{2-}$ to Cu charge transfer, and a similar intense core $\text{Cu}-\mu_4\text{S}^{2-}$ stretching mode in the resonance Raman spectrum, observed at 378 cm^{-1} in both sites. Small quantitative differences in the EPR hyperfine values and transition absorption and MCD intensities between the two sites arise from a perturbation of the spin density distribution of the cluster in 1-hole Cu_Z , where the more covalent $\mu_2\text{SH}^-$ leads to delocalization of the spin on Cu_I (dominant in Cu_Z^*) onto the edge SH^- . Despite the higher covalency of the Cu_Z site, the g values in the EPR spectra are similar for Cu_Z and Cu_Z^* , as the localization of spin on the four coordinate Cu_I in 1-hole Cu_Z^* leads to higher energy d-d transitions, opposing the decreased covalency, leading to the net low g values also observed experimentally for Cu_Z . The difference in edge ligation in the two sites is observed primarily in the resonance Raman enhanced vibrations, where a low energy $\text{Cu}-\mu_2\text{SH}^-$ stretch at 203 cm^{-1} and higher energy S-H bending modes at 450 and 492 cm^{-1} are additionally enhanced in the dominant $\mu_4\text{S}^{2-}$ to Cu CT transition in Cu_Z but not Cu_Z^* , due to the more covalent interaction between the coppers and the edge SH^- . Thus, the spectral similarities between 1-hole Cu_Z and 1-hole Cu_Z^* reflect similar bonding with the $\mu_4\text{S}^{2-}$ ligand and the distribution of spin over Cu_I , Cu_II , and Cu_IV . The differences in the vibrational spectra of the two sites reflect the $\mu_2\text{SH}^-$ versus $\mu_2\text{OH}^-$ edge ligation.

4.3 Insights into reactivity of 1-hole and 2-hole Cu_Z

1-hole Cu_Z has been shown to perform a slow 2 electron reduction of N_2O under single turnover conditions, with oxidation of both 1-hole Cu_Z and reduced Cu_A to generate resting 2-hole Cu_Z and 1 electron oxidized Cu_A .²⁶ A structure of PsN_2OR obtained from crystals pressurized with N_2O shows a linear N_2O molecule binding above the $\text{Cu}_\text{IV}-\text{Cu}_\text{II}$ edge of the Cu_Z cluster (Fig. 9).¹⁴ The O of N_2O is thought to be oriented towards a solvent filled cavity between Cu_Z and Cu_A , where there is a hydrogen bonding interaction with a localized solvent molecule, while the N end of the molecule is 2.8 \AA from Cu_IV and 3.5 \AA from the $\mu_2\text{SH}^-$ ligand. The spectroscopically and computationally defined protonation states for 1-hole and



2-hole Cu_Z indicate that the 1-hole Cu_Z site will donate both an electron and a proton upon oxidation, due to the significantly decreased pK_a of the $\mu_2\text{SH}^-$ in the 2-hole redox state. The participation of a proton in the reduction of N_2O by 1-hole Cu_Z avoids the thermodynamically unfavorable 1-electron reduction of N_2O to N_2O^- , which is endergonic by $25.4 \text{ kcal mol}^{-1}$, while the proton-coupled reduction of N_2O to form N_2 and a hydroxyl radical is exergonic by $7.4 \text{ kcal mol}^{-1}$.⁴⁰ However, a substantial barrier exists for this process due to the fact that N_2O is not activated through direct interaction with Cu_Z (the rate of N_2O reduction by 1-hole Cu_Z is $2 \times 10^{-4} \text{ s}^{-1}$).²⁶ Thus the N_2O may alternatively be oriented with the O atom pointed towards Cu_Z , where it can directly accept a proton and an electron from the $\mu_2\text{SH}^-$ to break the N–O bond and generate resting 2-hole Cu_Z , with transfer of the second electron from Cu_A . Since no intermediate is observed in the reduction of N_2O by 1-hole Cu_Z ,²⁶ the hydroxide product that would be formed after N–O bond cleavage would likely be rapidly protonated and released into the nearby solvent-filled cavity, rather than coordinating to the Cu_Z cluster.

The 2-hole resting state of Cu_Z has been defined as having a highly covalent sulfide ligand bridging the Cu_I – Cu_{IV} edge. This resting species is potentially the starting point for the chemical conversion of Cu_Z to Cu_Z^* , the reactive form of the cluster for N_2O reduction.²⁶ *In vitro*, the presence of O_2 is thought to promote the conversion of Cu_Z to Cu_Z^* , as isolation of N_2OR in the presence of O_2 results in samples with a high proportion of resting 1-hole Cu_Z^* , while the resting 2-hole state of Cu_Z is obtained when the purification is performed in the absence of oxygen.¹⁶ DFT calculations on the $\mu_2\text{S}^{2-}$ model of 2-hole Cu_Z suggest that there are frontier molecular orbitals (FMOs) available to interact with O_2 . The α and β HOMOs of 2-hole Cu_Z are occupied $\mu_2\text{S}^{2-}$ orbitals with dominant S p_z character (50% and 66% $\mu_2\text{S}^{2-}$ respectively, Fig. 8). This $\mu_2\text{S}^{2-}$ p_z orbital is oriented perpendicular to the Cu_3S_2 plane, towards the solvent-filled cavity where N_2O , and by analogy O_2 , would access the Cu_Z cluster. Based on these FMOs, reaction of the Cu_Z site with O_2 would proceed *via* oxidation of the edge sulfide, rather than by a Cu-based oxidation. Since this is a four electron process, there will in principle also be electrons available from the sulfide for the reduction of the copper site, dependent on the nature of the oxidized sulfur product. However, it is unlikely that this is the

mechanism involved in interconversion of Cu_Z and Cu_Z^* *in vivo*, since resting Cu_Z^* has been isolated under exclusion of oxygen conditions from anaerobically grown cells in bacterial strains with accessory genes knocked-out.⁴¹ Thus, the *in vivo* mechanism for interconversion of Cu_Z and Cu_Z^* , which is required to maintain N_2OR in the reactive Cu_Z^* form, and the role of accessory proteins in this process, remain to be identified.

5. Conclusions

We have used a combination of spectroscopies and DFT calculations to determine the protonation states of the edge sulfur in the 1-hole and 2-hole redox states of Cu_Z . From resonance Raman spectroscopy, 1-hole Cu_Z has a μ_2 thiolate ligand with a pK_a of 11–12, due to the presence of S–H bending modes that are not perturbed up to pH 10. DFT calculations of a 1-hole cluster with a $\mu_2\text{SH}^-$ ligand reproduce the key spectral features of 1-hole Cu_Z . The computational modeling of the 2-hole Cu_Z site indicates that the edge ligand is a $\mu_2\text{S}^{2-}$ with a pK_a of 3 or less, which is consistent with the absorption and resonance Raman features of 2-hole Cu_Z . The nature of this edge ligand has been used to obtain insight into the slow reduction of N_2O by 1-hole Cu_Z and suggest how 2-hole Cu_Z might react with O_2 , a possible route for the conversion of Cu_Z to Cu_Z^* *in vitro*.

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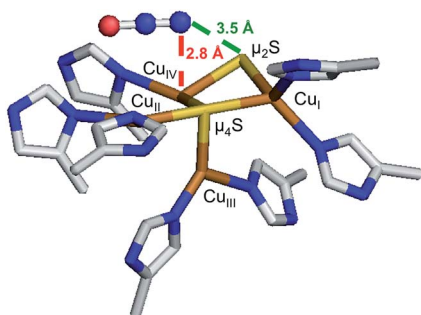


Fig. 9 Crystal structure of N_2O bound to Cu_Z (PDB ID: 3SBR, subunit C, resolution 1.7 Å).¹⁴



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