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Pure-shift IMPRESS EXSIDE – Easy measurement of ^1H – ^{13}C scalar coupling constants with increased sensitivity and resolution†

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The easy-to-interpret and rapid SelEXSIDE NMR experiment for measuring long-range ^1H – ^{13}C scalar coupling constants $^nJ_{\text{CH}}$ is improved with increased sensitivity and resolution by the incorporation of ‘pure-shift’ homonuclear decoupling and IMPRESS-Hadamard encoding, with typical increases in signal-to-noise of over 300% and 400% demonstrated for small molecule examples, and narrower linewidths enabling measurement of smaller coupling constants.

NMR analyses of the 3D-structure of conformationally flexible small molecules in solution are often under-determined *i.e.* the stereochemistry or conformation cannot be definitively assigned because the number and accuracy of experimental parameters is too small to account for all of the possible configurations and conformations of the molecule. In organic molecules, ^1H – ^1H scalar coupling constants ($^3J_{\text{HH}}$) and NOE measurements are the common NMR parameters used for this and are routinely combined with computational methods to improve structure elucidations or semi-quantitative methods exemplified by ‘ J -based configurational analysis’ to acyclic systems.^{1–3} We have recently demonstrated that increased accuracy in quantitative NOE-distance analysis of small molecules⁴ allows finer details of dynamic 3D molecular structures to be resolved – for example identifying very small populations (~2%) of molecular conformers,⁵ accurate determination of the rotational populations of a flexible chain⁶ and stereochemical elucidation of contiguous quaternary centres.⁷

Adding quantitative analysis of 2- and 3-bond ^1H – ^{13}C scalar coupling constants ($^nJ_{\text{CH}}$), of which there are many more than $^3J_{\text{HH}}$ couplings in most organic molecules, enables even greater levels of discrimination such that we have successfully

described the strong helical or linear conformational preference of iteratively homologated alkane diastereomers, despite thousands of potential conformers.⁸ However, if chemists are to routinely use the potentially powerful $^nJ_{\text{CH}}$ couplings it is crucial that simple, reliable and accurate experimental and analysis methods must be available which are accessible to non-spectroscopists. The practical challenges of extracting $^nJ_{\text{CH}}$ values from NMR data include sensitivity, inability to measure values for quaternary carbons, and analysis of convoluted lineshapes.^{9,10} Most existing 2-dimensional methods rely on measuring ^1H – ^{13}C coupling constants in the direct (horizontal) F2 dimension and suffer from simultaneous overlap of ^1H – ^1H and ^1H – ^{13}C couplings with severe lineshape distortions arising from this. Castanar *et al.* recently reported the ‘pure in-phase’ (PIP) HSQMBC¹¹ which ensures clean ^1H – ^1H and ^1H – ^{13}C coupling co-evolution resulting in substantially improved F2 lineshapes, although complex ^1H – ^1H multiplets still make it challenging to accurately measure small coupling constants (<2 Hz) with this approach.

An alternative approach to dealing with lineshape distortion is to use J -resolved approaches, exemplified by EXSIDE.¹² Here the desired heteronuclear coupling constant is isolated and resolved as clean in-phase doublets in the indirect (vertical) F1 dimension of the spectrum (Fig. 1a). The F1-splitting of the peaks is the multiple of $^nJ_{\text{CH}}$ and a user-selected J -scaling factor, N . This provides an easy accurate means to measure even small $^nJ_{\text{CH}}$ values, down to ~1 Hz in most cases. Such F1-based approaches, however, are relatively slow and weak and it generally requires several hours of measurement for each proton, while intensities are also reduced by T_2 relaxation in the extended J -scaled INEPT preparation period. The long experiment time arises because a large number of t_1 increments are needed to resolve small couplings in F1. Our recently reported ^{13}C band-selective EXSIDE¹³ (SelEXSIDE) method substantially reduces the experiment time to just a few minutes by measuring each $^nJ_{\text{CH}}$ separately with a much narrower ^{13}C window, thus needing fewer t_1 increments. However the short experiment time also reduces signal averaging and thus sensitivity, so in

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† Electronic supplementary information (ESI) available: Experimental details, including pulse sequence diagrams, additional data acquired for menthol, and alternative acquisition modes for IMPRESS data and a tar file containing VNMRJ 4 compatible pulse sequence and acquisition/processing macros. See DOI: 10.1039/c5ra24926k



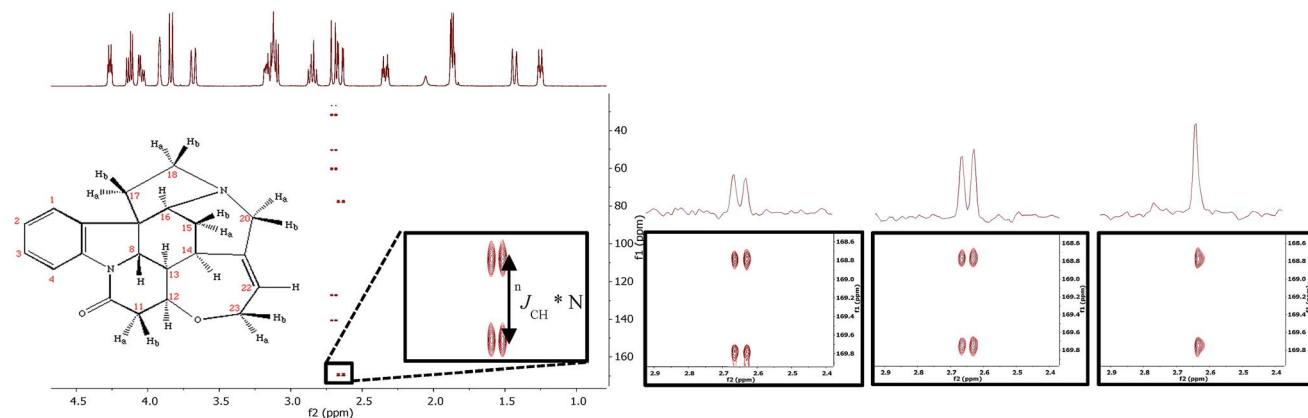


Fig. 1 Strychnine structure and H11b–C10 correlations from (left-to-right): H20b/H11b EXSIDE spectrum (~7.5 hours) with H11b–C10 inset; SeLEXSIDE spectrum (50 minutes); ImpEXSIDE spectrum (50 minutes, as one of four correlations measured simultaneously); ImpEXSIDE_PS spectrum (50 minutes, as one of four correlations measured simultaneously).

practice it requires >10 mmol samples to obtain reasonable signal intensities using a 500 MHz room temperature NMR probehead.

Herein we combine two methods to substantially improve the sensitivity of EXSIDE-based experiments: greater signal-averaging by spectrum aliasing and F1-Hadamard encoding; and increasing signal intensity (height) through F2-homonuclear decoupling.

Hadamard encoding of the F1 dimension of EXSIDE in 2D NMR spectra allows simultaneous measurement of multiple ^{13}C regions of the EXSIDE spectrum – effectively simultaneous measurement of several SeLEXSIDE¹³ experiments, providing greater signal averaging for each peak. Hadamard encoding is achieved by the IMPRESS technique described by Krishnamurthy,¹⁴ replacing SeEXSIDEs selective ^{13}C inversion pulses with a series of n symmetrically shifted laminar pulses¹⁵ which are phase-encoded to match Hadamard matrices, H_n .^{14,16–18} The IMPRESS EXSIDE (ImpEXSIDE) pulse sequence (Fig. S1†)¹⁹ incorporates the ^{13}C symmetrically shifted pulses (S_1 and S_2) during the DPGFSE refocussing period. S_1 and S_2 combine to phase encode the ^{13}C chemical shift regions of interest. For a 4-region IMPRESS EXSIDE, in all scans S_1 will always have an x phase at all four chemical shifts (S_{xxxx}), however the phase of S_2 at each chemical shifts (A, B, C, D) will correspond to the rows of the H_4 matrix shown below where ‘+’ and ‘-’ represent x and y phase pulses respectively at each chemical shift (that is S_2 will be S_{xxxx} , S_{xyxy} , S_{xxyy} and S_{xyyx} for scans 1, 2, 3 and 4 respectively).

$$H_4 = \begin{bmatrix} & A & B & C & D \\ \text{Scan 1} & + & + & + & + \\ \text{Scan 2} & + & - & + & - \\ \text{Scan 3} & + & + & - & - \\ \text{Scan 4} & + & - & - & + \end{bmatrix}$$

The desired sub-spectra for each ^{13}C chemical shift region can then be readily extracted by the appropriate sum/difference of the four scans following the appropriate columns of the Hadamard matrix. For example, the sum of all four scans (Scan 1 + Scan 2 + Scan 3 + Scan 4) will give sub-spectrum A (while

cancelling contributions from B, C and D) while Scan 1 + Scan 2 – Scan 3 – Scan 4 will give sub-spectrum C, etc (See ESI† for full details, along with pulse sequence and processing templates/macros for Varian/Agilent NMR spectrometers). The IMPRESS data for H11b–C10 of strychnine is shown in Fig. 1c alongside data for EXSIDE¹² and SeLEXSIDE¹³ experiments (Fig. 1a and b respectively).

ImpEXSIDE and SeLEXSIDE data were obtained for all ^{13}C couplings to H15a/b, H11b and H20b of strychnine. For comparability both techniques measured four $^nJ_{\text{CH}}$ values in 50 minutes (one 50 minute experiment for ImpEXSIDE, four 12.5 minute experiments for SeLEXSIDE) and in all cases the $^nJ_{\text{CH}}$ values measured from these experiments matched to within 0.2 Hz, and were in line with literature experimental and computed values.¹³ Fig. 2 illustrates the systematic increase in signal-to-noise ratio obtained with ImpEXSIDE (green) compared to the parent SeLEXSIDE (blue). In fact in a number of cases correlations too weak to observe by SeLEXSIDE under these conditions were clear in ImpEXSIDE experiments (H15a–C14, H15a–C16, H15b–C16, H15b–C12, H15b–C21, H20b–C18 and H20b–C22). In theory, the IMPRESS method should increase sensitivity for all correlations by around 2-fold compared to SeLEXSIDE as signals are averaged for four-times longer but in practice a mean 2.5-fold increase in sensitivity was found, with some variation of responses from ~1.5-fold (H20b–C21) to ~4.0-fold (H11b–C12 and H15b–C13). This variability is ascribed to the differing placement of the selective ^{13}C inversion pulses compared to SeLEXSIDE (these are in the initial DPGFSE INEPT period for SeLEXSIDE) and differential sum/difference decoding of the acquired FIDs, for example the all-plus column of a Hadamard matrix is more sensitive to instrumental errors.¹⁸ In any case, there is no doubt that substantial sensitivity improvements arise from the ImpEXSIDE as compared to the original SeLEXSIDE when more than one $^nJ_{\text{CH}}$ value is required to be measured. Fig. 1 also highlights a reduction in the F1 half-height line-width moving from EXSIDE (20.2 Hz) to the faster selective variants ImpEXSIDE (13.4 Hz) and SeLEXSIDE (11.0



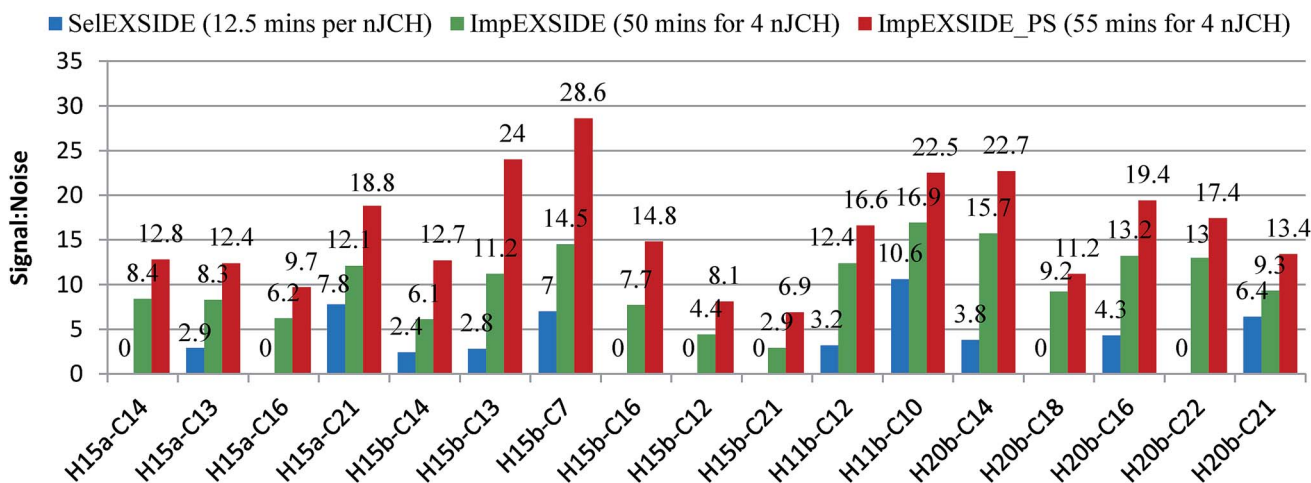


Fig. 2 The signal to noise ratio of peaks arising from heteronuclear coupling of H15a/b, H11b and H20b of strychnine from SelEXSIDE (blue), ImpEXSIDE (green) and Pure-shift ImpEXSIDE (red).

Hz), allowing even better discrimination of smaller J -scaled couplings.

Two practical limitations of ImpEXSIDE relate to the ^1H and ^{13}C selection frequencies. As with all EXSIDE techniques, the selected ^1H region must not contain two protons that are mutually scalar coupled to ensure that only heteronuclear coupling is J -scaled in F1 (although multiple uncoupled protons can be selected if desired). Further, the selected ^{13}C regions (A, B, C and D) must not overlap to any great extent as the phase-encoded pulse profiles for each region are not rectangular and imperfect phase behaviour will be observed in overlap regions – although there is no reason why multiple carbon peaks cannot be present in each selected ^{13}C region.

Further improvement to the sensitivity of EXSIDE-based experiments can be made by eliminating ^1H – ^1H coupling in F2 through gated band-selective refocusing of the homonuclear coupling during t_2 acquisition (known variously as “pure-shift”, “HOBs” (HOMonuclear Band Selective) or “bash” (BAND-Selective Homonuclear) decoupling methods).^{20–22} Broadly classified as ‘pure-shift’ spectra, these methods have raised a lot of interest recently²³ with significant improvements in sensitivity where the ^1H of interest (‘active’) is decoupled by selective inversion of their coupled partner (‘passive’) spins, by either band-selection²² or isotopic selection.²⁴

Fig. 3 illustrates the band-selective J -refocussing element which is incorporated into the standard acquisition period of the ImpEXSIDE.¹⁹ The J -refocussing element (consisting of a hard 180° ^1H pulse on all spins followed by a selective 180° ^1H pulse on the active spin only) refocuses the t_2 evolution of chemical shift, heteronuclear scalar coupling as well as homonuclear scalar coupling of the selected (active) ^1H . Acquisition is systematically interrupted with this J -refocussing element after each acquired chunk of data. The length of each chunk, at/n (where a is acquisition time and n is number of repeats) should be $<1/(3^{*n}J_{\text{HH}})$,²² to minimize the breakthrough of homonuclear coupling. In practice, the choice of chunk length is largely determined by a compromise between resolution achieved and

spectrum artifacts and in some cases can be as small as $1/(5^{*n}J_{\text{HH}})$. When applied to ImpEXSIDE, the resulting correlations (Fig. 1d) are thus collapsed to simple singlets in F2. This decoupling method is also readily implemented in the parent EXSIDE and SelEXSIDE sequences with essentially the same improvements and limitations (see ESI† for details).

Importantly, this band-selective homonuclear decoupling can be incorporated without substantial loss of information or sensitivity as the EXSIDE-based experiments are already inherently F2(^1H)-selective. Hence collapsing the F2-multiplets to narrow singlets should result in a substantial improvement in signal:noise (ideally up to 2-fold for a doublet or triplet, 4-fold for a resolved doublet of doublets, *etc.*). On the other hand the F2-width of the final decoupled ‘singlet’ is limited to ~ 4 – 6 Hz by B1 inhomogeneity of the pulses and the artificially increased T_2 relaxation in the FID (resulting from non-acquisition during the J -refocusing period which, in some cases, contains relatively long ^1H selective inversion pulses). Small F2-sideband artifacts appear at regular intervals equivalent to $1/(\text{chunk length})$ either side of the decoupled resonance (visible in Fig. 1), arising from the regular discontinuities in the intensity of sequential data

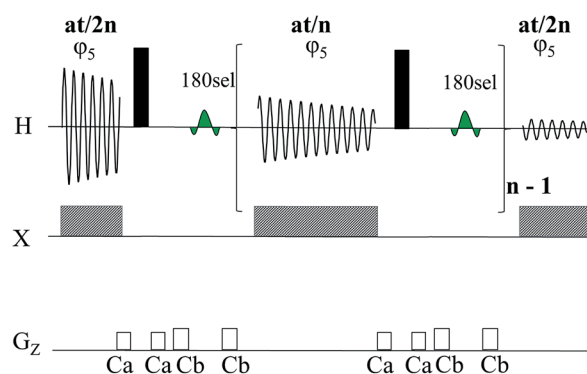


Fig. 3 Band-selective J -refocussing element incorporated into the t_2 acquisition period for band-selective homodecoupling. The 180sel pulse (green) is a refocussing 180° RSNOB.



points where the FID is 'stitched' back together around the refocussing periods. Crucially, the presence of these low intensity F2 artefacts do not affect the measurement of the desired $^nJ_{\text{CH}}$ value in F1.

Fig. 2 illustrates the resulting 1.25–2.3 fold improvement in sensitivity of the Pure-shift modification (ImpEXSIDE_PS, red) on correlations for strychnine compared to the parent ImpEXSIDE (green). As expected, the largest increases were observed from the most highly coupled signal (H15a, doublet of triplets) while the smallest increases arise for the least coupled (H20b, doublet). It should be noted that strychnine is nearly a worst-case substrate for sensitivity improvement by homonuclear decoupling as its ^1H spectrum consists mostly of sharp doublets/triplets. With more highly coupled substrates greater sensitivity enhancement can be achieved. For instance, application of ImpEXSIDE_PS to menthol (see ESI†) gave over 2-fold sensitivity enhancement compared to ImpEXSIDE in every case for the highly coupled signals of this more complex ^1H NMR spectrum.

In summary, the sensitivity of EXSIDE and SeEXSIDE experiments can be substantially improved by the incorporation of IMPRESS (F1-spectral aliasing and Hadamard encoding) and F2-homonuclear decoupling either independently or in combination. In addition to improved resolution in both dimensions, there is no loss of signal or correlations from incorporating these modifications and >300% sensitivity increases are achieved for most resonances of strychnine, and substantially higher for more complex multiplets. With this sensitivity improvement, data acquisition should be an order of magnitude faster for sample-limited experiments and $^nJ_{\text{CH}}$ values should be measurable for single-digit mmol concentrations (corresponding to ~ 1 mg of a 500 g mol^{-1} compound) in accessibly short experiment times.

Notes and references

- 1 N. Matsumori, D. Kaneno, M. Murata, H. Nakamura and K. Tachibana, *J. Org. Chem.*, 1999, **64**, 866.
- 2 A. Ardá, M. I. Nieto, M. Blanco, C. Jiménez and J. Rodríguez, *J. Org. Chem.*, 2010, **75**, 7227.
- 3 P. Cimino, G. Bifulco, A. Evidente, M. Abouzeid, E. Riccio and L. Gomez-Paloma, *Org. Lett.*, 2002, **4**, 2779.

- 4 C. P. Butts, C. R. Jones, E. C. Towers, J. L. Flynn, L. Appleby and N. J. Barron, *Org. Biomol. Chem.*, 2011, **9**, 177.
- 5 C. P. Butts, C. R. Jones and J. N. Harvey, *Chem. Commun.*, 2011, **47**, 1193.
- 6 C. R. Jones, C. P. Butts and J. N. Harvey, *Beilstein J. Org. Chem.*, 2011, **7**, 145.
- 7 M. G. Chini, C. R. Jones, A. D. Zampella, M. V. D'Auria, B. Renga, S. Fiorucci, C. P. Butts and G. Bifulco, *J. Org. Chem.*, 2011, **77**, 1489.
- 8 M. Burns, S. Essafi, J. R. Bame, S. P. Bull, M. P. Webster, S. Balieu, J. W. Dale, C. P. Butts, J. N. Harvey and V. K. Aggarwal, *Nature*, 2014, **513**, 183.
- 9 B. L. Marquez, W. H. Gerwick and R. T. Williamson, *Magn. Reson. Chem.*, 2001, **39**, 499.
- 10 T. Parella and J. Felix Espinosa, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2013, **73**, 17.
- 11 L. Castanar, J. Sauri, R. T. Williamson, A. Virgili and T. Parella, *Angew. Chem., Int. Ed.*, 2014, **53**, 8379.
- 12 V. V. Krishnamurthy, *J. Magn. Reson., Ser. A*, 1996, **121**, 33.
- 13 C. P. Butts, B. Heise and G. Tatolo, *Org. Lett.*, 2012, **14**, 3256.
- 14 K. Krishnamurthy, *J. Magn. Reson.*, 2001, **153**, 124.
- 15 S. L. Patt, *J. Magn. Reson.*, 1992, **96**, 94.
- 16 R. Crouch, R. D. Boyer, R. Johnson and K. Krishnamurthy, *Magn. Reson. Chem.*, 2004, **42**, 301.
- 17 K. Krishnamurthy, *J. Magn. Reson.*, 2001, **153**, 144.
- 18 E. Kupče, T. Nishida and R. Freeman, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2003, **42**, 95.
- 19 The pulse sequence diagram and a detailed walk-through can be found in the ESI†
- 20 R. W. Adams, L. Byrne, P. Kiraly, M. Foroozandeh, L. Paudel, M. Nilsson, J. Clayden and G. A. Morris, *Chem. Commun.*, 2014, **50**, 2512.
- 21 L. Castañar, J. Sauri, P. Nolis, A. Virgili and T. Parella, *J. Magn. Reson.*, 2014, **238**, 63.
- 22 N. H. Meyer and K. Zangger, *Angew. Chem., Int. Ed.*, 2013, **52**, 7143.
- 23 (a) K. Zangger, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2015, **86–87**, 1; (b) L. Castanar, T. Parella *Magn. Reson. Chem.*, 2015, **53**, 399.
- 24 L. Paudel, R. W. Adams, P. Király, J. A. Aguilar, M. Foroozandeh, M. J. Cliff, M. Nilsson, P. Sándor, J. P. Waltho and G. A. Morris, *Angew. Chem., Int. Ed.*, 2013, **52**, 11616.

