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Title: Origin of exotic ferromagnetic behavior in exfoliated layered transition metal dichalcogenides MoS₂ and WS₂

The room temperature ferromagnetism is closely related to the structural disorder on the sheet edges and the presence of a conductive metallic phase. The origin of room temperature ferromagnetism in transition metal dichalcogenides (TMDCH: MoS₂ and WS₂) was investigated in detail. The ferromagnetism is strongly dependent on the concentration of defects predominantly located on the TMDCH sheets edges as well as on the formation of the metallic 1T phase. Significant differences between the ferromagnetism of TMDCH exfoliated using butyllithium and sodium naphtalenide originate from differences in the exfoliation degree.
Origin of exotic ferromagnetic behavior in exfoliated layered transition metal dichalcogenides MoS$_2$ and WS$_2$

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Bulk layered transition metal dichalcogenides (TMDs) show diamagnetic properties. When exfoliated, the materials’ band gap increases and changes from an indirect band gap to a direct one. During the exfoliation, the TMDs may undergo a phase transition from 2H to 1T polymorph, which is likely electronically driven and accompanied by a metal–insulator transition. A significantly higher efficiency of the exfoliation was observed using sodium naphthalenide compared to butyllithium. Moreover we demonstrate that the exfoliation has a dramatic influence on the magnetic properties of two TMDs, MoS$_2$ and WS$_2$. These materials become partly ferromagnetic upon exfoliation, which is a highly unexpected behavior. Exotic ferromagnetism is generally observed on samples with a high degree of exfoliation, which indicates the association of this effect with defects formed on the edges of dichalcogenide sheets. Such an exotic ferromagnetic behavior, if properly understood and brought under material engineering control, shall open the door to new applications of these materials.

Introduction

Layered transition metal dichalcogenides (TMDs) have received revived interest in the past few years being dubbed as the “next graphene” due to their possible applications in energy generation and their highly anisotropic material properties. In contrast to graphene, where the layer is single atom thick and truly 2-dimensional, MoS$_2$ (or WS$_2$) sheets consist of transition metals sandwiched between two layers of chalcogen atoms. These triple layers are bonded to each other by van der Waals interactions. Such a “triple-layer” can exhibit three different crystallographic arrangements, the most common semiconducting 2H (trigonal prism), metallic 1T (octahedral) and semiconducting 3R (trigonal prism). The broad selection of metal elements and chalcogens in TMD as well as the large variety of different crystal structures of the materials yield a large number of combinations for tailoring the TMD properties. One of them is the magnetic behaviour. Similarly to graphite, MoS$_2$ or WS$_2$ in their bulk forms do not exhibit ferromagnetic or paramagnetic behavior. The magnetic properties of TMDs have been studied both theoretically and experimentally. Theoretically it has been predicted that the introduction of sulphur vacancies onto the S-edges of MoS$_2$ nanoribbons would lead to a non-zero net magnetic moment. Furthermore zigzag and armchair terminated edges of MoS$_2$ nanoribbons were studied by density functional theory (DFT) calculations. It was predicted that the zigzag terminated nanoribbons are metallic and show ferromagnetic properties while armchair terminated nanoribbons are semiconducting and nonmagnetic. A magnetic ordering has also been theoretically described for WS$_2$ nanosheets. The results of DFT calculations have shown that the adsorption of non-metallic elements such as H, B, O or F onto the edges of TMDs induces magnetic properties. More recently experimental studies have been emerging. While the synthesis of magnetic TMD nanoribbons has not been reported, probably because of the difficulties in their synthesis, only a few reports on TMD magnetic nanosheets exist. The first report on magnetic MoS$_2$ nanosheets was based on thermal evaporation synthesis onto silicon or tantalum foil substrates. However the role of the substrate and the interface between MoS$_2$ and the substrate remains unclear and may have influenced the magnetic properties of the prepared samples. Further studies have shown the preparation of MoS$_2$ by a two-step hydrothermal synthesis, which introduces sulphur vacancies responsible for magnetic behaviour. Magnetic WS$_2$ nanosheets have also
been prepared by exfoliation in DMP. However, the frequently used organometallic intercalation followed by exfoliation has not been employed for the synthesis of the above mentioned TMD nanomaterials. These methods are easily scalable and relatively undemanding. Here we wish to show that exfoliated MoS₂ and WS₂ exhibit exotic ferromagnetic properties, which we consider as an unexpected behavior in view of the fact that bulk TMDs exhibit diamagnetic properties.

**Experimental**

**Materials**

Molybdenum disulphide (99.8%) and tungsten disulphide (99.8%) were obtained from Alfa Aesar. n-Butyllithium (1.6 M in hexane), benzophenone and sodium were obtained from Sigma-Aldrich, Czech Republic. Tetrahydrofuran (THF), hexane and naphthalene were delivered by Lach-Ner, Czech Republic. Hexane was dried with sodium before use and THF was dried by distillation from a sodium/benzophenone mixture under an argon atmosphere before use.

**Synthesis procedure**

The exfoliation of MoS₂ and WS₂ with butyllithium was performed at room temperature. 3 g of MoS₂ was stirred with butyllithium solution for 72 hours. Subsequently the intercalated MoS₂ and WS₂, respectively, were filtered out and washed with hexane. All these operations were performed under an argon atmosphere in a glovebox. Intercalated MoS₂ and WS₂ were dispersed in water and repeatedly centrifuged and dispersed in water. Finally the exfoliated material was dried in a vacuum oven for 48 hours at 50 °C before further use.

Sodium naphthalenide used as an intercalation compound was prepared by stirring sodium metal (0.35 g) with naphthalene (1.8 g) for 2 h in 50 ml of THF under an argon atmosphere. After the addition of transition metal dichalcogenide (TMD) powders (3 g), intercalation was carried out by stirring the solutions at room temperature for 24 h under an argon atmosphere. The prepared intercalates were then suction filtered and washed several times with THF. Separated samples were dispersed in water and centrifuged several times. Finally the exfoliated material was dried in a vacuum oven for 48 hours at 50 °C before further use.

**Analytical techniques**

Morphology studies were performed by using a scanning electron microscope (SEM) equipped with a FEG source of electrons (Tescan Lyra Dual Beam Microscope). The sample composition was determined by means of an energy dispersive spectroscopy (EDS) analyser (X-MaxN) with a 20 mm² SDD detector (Oxford Instruments). Data were evaluated using AZtecEnergy software. Before the measurement, samples were placed onto a conductive carbon tape to avoid charging effects. All the measurements were carried out with 15 kV acceleration voltage.

High-resolution transmission electron microscopy (HRTEM) was carried out on a transmission electron microscope JEOL JEM 2200FS operated at 200 kV (autoemission Schottky gun, point resolution 0.19 nm) with an in-column energy Ω-filter for EELS/EFTEM analyses, a STEM unit, and an Energy Dispersive X-ray (EDX) SDD detector (Oxford Instruments X-Max attached). Images were recorded on a Gatan CCD camera with a resolution of 2048 × 2048 pixels using the Digital Micrograph software package. EDX analyses were performed and treated with the INCA software package.

High resolution X-ray photoelectron spectroscopy (XPS) was performed with an ESCAProbeP (Omicron Nanotechnology Ltd, Germany) spectrometer using a monochromatic aluminium X-ray radiation source (1486.7 eV). A wide scan survey of all elements was performed with consequent high-resolution scans of the Mo, W and S core levels. Relative sensitivity factors were used to determine the chalcogen-to-metal ratio and the quantitative distribution of 2H and 1T phases. The samples were placed onto a carbon conductive tape prior to the measurement.

Characterization by Atomic Force Microscopy (AFM) was performed on the NT-MDT Ntegra Spectra from NT-MDT in tapping mode. The samples were dispersed in isopropyl alcohol (0.5 mg mL⁻¹), sonicated for 10 minutes and then dropped onto a freshly cleaved mica substrate before the measurement.

An InVia Raman microscope (Renishaw, England) was used for Raman spectroscopy measurements. The spectrometer operates in backscattering geometry with a CCD detector. A Nd-YAG laser (532 nm, 50 mW) was used with 50× magnification objective. The instrument was calibrated on a silicon reference which gives a peak position at 520 cm⁻¹. The laser power did not exceed 5% of the maximum 50 mW power in order to avoid damage to the sample.

The magnetic measurements were performed using a SQUID magnetometer MPMS-XL (Quantum Design). The hysteresis loops were recorded between −10 and 10 kOe at T = 300 K. The magnetization measurements at T = 2 K were conducted between 0 and 70 kOe. The surface areas were measured using adsorption of methylene blue. About 0.05 g of each TMD material (weighed exactly to within ±0.01 mg) was dispersed in 70 mL of deionized water by ultrasonication (400 W, 30 min). A defined amount of 1 mM of methylene blue aqueous solution was added to the suspension and then filled to 100 mL. The concentration of unabsorbed methylene blue was determined by UV-Vis spectroscopy at λ_max = 665 nm (Cary 50, Agilent). For the calculation of surface areas, the value of 1.05 nm² per methylene blue molecule was used.

**Results and discussion**

Here we have investigated the effect of the exfoliation process on the magnetic and structural properties of molybdenum and tungsten disulfide. We have used two agents for chemical exo-
liaction of bulk MoS$_2$ and WS$_2$: butyllithium (BuLi) and sodium naphthalenide (NAPH). While butyllithium is commonly used for the chemical exfoliation of transition metal dichalcogenides (TMDs), only little has been reported on TMD exfoliation using sodium naphthalenide. Further on, we will show that sodium naphthalenide is not only more efficient in TMD exfoliation than butyllithium for both MoS$_2$ and WS$_2$ but also induces magnetic ordering. Before the discussion of exotic magnetic properties of exfoliated TMDs we thoroughly examined the exfoliated TMDs by scanning electron microscopy (SEM), atomic force microscopy (AFM), X-ray photoelectron spectroscopy (XPS), Raman spectroscopy, high-resolution transmission electron microscopy (HR-TEM), selective area electron diffraction (SAED) and surface area measurement by methylene blue adsorption. We will show that the magnetic properties measured by using a superconducting quantum interference device (SQUID) strongly depend on the character and concentration of the induced defects. Simultaneously, the phase transition from the 2H to 1T form, which sets in as a result of intercalation (likely due to charge carrier doping), has to be correlated with the peculiar magnetic behavior.

The morphology of the exfoliated and bulk TMDs was examined by SEM as shown in the ESI (Fig. SI1†). One can observe that for BuLi exfoliated WS$_2$, most of the sheets remained stacked with sharp edges, indicating an incomplete intercalation and the subsequent exfoliation has been incomplete. It is known that intercalation of Li to WS$_2$ is incomplete at room temperature and becomes more pronounced at higher temperatures. NAPH exfoliation of WS$_2$, on the other hand, leads to sheets with smaller lateral sizes and slight edge wrinkling, which points to a significantly higher degree of exfoliation than that in the case of WS$_2$-BuLi. A similar trend can be observed for MoS$_2$, however, with more substantial edge wrinkling, especially for NAPH exfoliated MoS$_2$.

XRD measurements were further conducted (Fig. SI2†) to confirm the presence of pure TMDs. We have not detected any diffraction lines originating from molybdenum or tungsten oxides. This indicates a low degree of oxidation of MS$_2$ to MO$_x$/MO$_y$. Even though some increase in MO$_3$ was observed by XPS measurements (see XPS discussion), it was only on the surface and was only detectable by XPS. It is also obvious from the data that NAPH exfoliated TMDs show less intense and broadened reflections. This is especially apparent for NAPH-MoS$_2$ which only shows weak and very broad reflections indicating a high degree of exfoliation. As for NAPH-WS$_2$ we have detected much higher tendency for restacking (even though AFM and Raman spectroscopy indicate a similar level of exfoliation as for MoS$_2$ and all XRD measurements were performed within a couple of weeks after preparation) which is obvious from its quite sharp XRD pattern. Also the BuLi exfoliated TMDs show much lower degree of exfoliation which is in good agreement with the data discussed later on.

We used AFM to determine the sheet thickness and to further investigate the morphology of the exfoliated MoS$_2$ and WS$_2$. Fig. 1 shows the AFM images documenting that the use of BuLi for exfoliation of MoS$_2$ and WS$_2$ leads to multilayered materials with sheet thicknesses of ∼20–25 nm. While WS$_2$ retains its sharp edged structure observed by SEM, MoS$_2$ consists of sub-micron sized flakes. The application of NAPH led to a higher degree of exfoliation of both MoS$_2$ and WS$_2$ manifested by a typical sheet thickness (∼0.7 nm) corresponding to a single layer. This result is in good agreement with the previously reported values of the monolayer thickness of MoS$_2$ and WS$_2$, respectively.

To obtain insight into the structural properties of exfoliated TMDs, we performed Raman spectroscopy. It is well known that Raman spectroscopy is a very useful tool for identifying structural properties as well as the number of layers as has been shown several times for various TMDs including MoS$_2$ and WS$_2$. WS$_2$ shows two first-order modes denoted as E$_{2g}$ and A$_{1g}$ at 336 cm$^{-1}$ and 418 cm$^{-1}$, respectively and a disorder activated mode denoted as LA at 352 cm$^{-1}$, which dominates the WS$_2$ Raman spectrum. These modes can be effectively used to determine the number of layers. Raman spectra show no apparent difference between BuLi exfoliated WS$_2$ and the spectrum of bulk materials (Fig. 2). This would suggest that the material was delaminated only to a small extent, which is in good agreement with the results previously observed by SEM and AFM. The spectrum of NAPH exfoliated WS$_2$ shows major differences. We can note the decrease in the I(A$_{1g}$)/I(E$_{2g}$) ratio from ∼2 in the bulk material to ∼0.35 in

![Fig. 1 AFM images of exfoliated transition metal dichalcogenides.](image)

![Fig. 2 Raman spectra of exfoliated transition metal dichalcogenides in their bulk state (top), after exfoliation with butyllithium (middle) and after exfoliation with sodium naphthalenide (bottom).](image)
NAPH exfoliated WS₂. This has been previously reported as a characteristic feature of a single atom layered WS₂.²⁴ It should be noticed that the band centered at around 354 cm⁻¹ is broadened, which is associated with an activation of LA mode due to disorder introduced during the exfoliation (e.g. edge wrinkling observed by SEM and HR-TEM). In addition we have located low intensity peaks in the low frequency region (see Fig. S13†). These peaks correspond to the 1T phase and are denoted as J₁, J₂ and J₃.²⁷ Even though these modes confirm the presence of the 1T form since they are not allowed in the 2H phase structure, it is more accurate to use other methods, such as XPS, to calculate the 1T concentrations.²⁷ It is noticeable that these modes are also visible even in the spectrum of BuLi-WS₂ where AFM and XRD showed a very low, if any, degree of exfoliation. This suggests that even though the intercalation of BuLi into WS₂ and its subsequent exfoliation are highly inefficient, some charge carrier doping takes place, leading to a slight increase in the 1T phase concentration (see XPS for a detailed discussion).

The modes denoted as E₁₂g at 383.5 cm⁻¹ and A₁g at 408.6 cm⁻¹ in the spectra of MoS₂ are of analytical value. The peak separation Δ is used in the case of MoS₂ for the determination of the number of layers. We have obtained Δ = 26.1 cm⁻¹ for bulk MoS₂ and Δ = 26 cm⁻¹ for BuLi exfoliated MoS₂. These values are nearly identical and they signify that the thickness of BuLi exfoliated MoS₂ is more than four layers.²⁸ By contrast, we have shown by AFM measurements that the average sheet thickness is lower than that of the bulk material. More importantly, NAPH exfoliated MoS₂ shows the peak separation of Δ = 21.4 cm⁻¹. This value is close to that of monolayer MoS₂ documenting that the high efficiency of NAPH was highly efficient in MoS₂ exfoliation.²⁸,²⁹

In addition to E₁₂g and A₁g modes, several additional features can be seen in the low wavenumber region of MoS₂ and WS₂ spectra (see Fig. S13†). Firstly, modes denoted as J₁, J₂ and J₃ (pink regions) are clearly visible in the spectra of MoS₂·BuLi and MoS₂·NAPH. These modes are not allowed in 2H modification but become allowed in 1T MoS₂ polymorph.³⁰,³¹ This indicates a successful 2H to 1T phase transition in BuLi·MoS₂ and NAPH·MoS₂ samples. In the spectrum of bulk MoS₂ none of these modes (J₁, J₂, J₃) can be detected. Even though the presence of J₁, J₂ and J₃ phonon modes prove the presence of the 1T phase, other methods such as XPS are more suitable for determining the relative 1T phase concentrations.³⁰ The peak at ~280 cm⁻¹ is related to the presence of MoO₃.³² The presence of MoO₃ on the surface has been proven by XPS measurements (see XPS discussion). The increase in the intensity of the peak located at ~280 cm⁻¹, compared to the bulk material, for MoS₂·BuLi and MoS₂·NAPH is likely caused by slightly higher amounts of MoO₃ on the surface of these samples, also observed by XPS.

So far we have shown by means of SEM, AFM and Raman measurements that sodium naphthalenide is more efficient in exfoliation of bulk MoS₂ and WS₂ than the usually used butyllithium. In order to obtain better information on the chemical changes taking place within the prepared materials, we conducted XPS measurements. Detailed high-resolution spectra of core-level Mo 3d and W 4f after Shirley background subtraction are displayed in Fig. 3.

It is well known that bulk materials consist mainly of the semiconducting 2H phase. During the exfoliation varying amounts of metallic 1T phase emerge as a consequence of electron doping introduced during intercalation.³³ The relative changes of the transition from the 2H to 1T phase are summarized in Table 1. For WS₂, the spectra consist of W⁴⁺ 4f⁵/₂ and W⁴⁺ 4f⁷/₂ peaks at approximately 32 eV and 35.1 eV, respectively, corresponding to the 2H phase. A contribution of the 1T phase can be also found in the spectra at around 31.1 eV and 34.2 eV as well as the pair corresponding to W⁶⁺ at ~36.8 eV and ~38.5 eV. The deconvolution of those components is essential when trying to identify the chemical changes occurring during the exfoliation.

For BuLi exfoliated WS₂ it is clearly visible that there is only a small contribution from the metallic 1T phase of about 10 at%. This reflects the low degree of exfoliation (and the low degree of electron doping) also supported by electron microscopy and Raman measurements. Then for NAPH exfoliated WS₂ a much more pronounced increase in the 1T phase to ~37% was observed. The presence of W⁶⁺ is also notable from the spectra for both samples.

Let us note that even in bulk materials there is a certain amount of tungsten atoms (as well as Mo in MoS₂) in the

Table 1 Contents of 1T, 2H phase and M⁶⁺ as obtained from the high-resolution core-level spectra of Mo 3d and W 4f

<table>
<thead>
<tr>
<th>Sample</th>
<th>1T (at%)</th>
<th>2H (at%)</th>
<th>M⁶⁺ (at%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>MoS₂·BuLi</td>
<td>15.3</td>
<td>73.3</td>
<td>11.4</td>
</tr>
<tr>
<td>MoS₂·NAPH</td>
<td>38.9</td>
<td>48.2</td>
<td>12.9</td>
</tr>
<tr>
<td>WS₂·BuLi</td>
<td>10.2</td>
<td>77.8</td>
<td>12.0</td>
</tr>
<tr>
<td>WS₂·NAPH</td>
<td>37.1</td>
<td>48.1</td>
<td>14.7</td>
</tr>
</tbody>
</table>

Fig. 3 High-resolution XPS spectra of W 4f in WS₂ and Mo 3d in MoS₂ exfoliated with butyllithium and sodium naphthalenide.
oxidation state of M⁶⁺ (see Table SI1†). The slight increase of W⁶⁺ concentration in NAPH-WS₂ is likely caused by a higher surface area available in the more exfoliated material. Now for MoS₂ the detailed high-resolution core-level spectra can be evaluated in a similar manner as for WS₂. In this case these are the Mo⁴⁺ 3d₃/₂ at ~229.1 eV and Mo⁶⁺ 3d₅/₂ at ~232.5 eV peaks that correspond to the 2H phase. The peaks corresponding to the 1T phase appear at about 0.9 eV lower binding energies (~228.2 eV and ~231.4 eV). Finally Mo⁶⁺ species are located at approximately 232.7 eV and 235.5 eV. For BuLi exfoliated MoS₂ the 1T phase content is ~15 at% while for NAPH exfoliated MoS₂ it is as high as ~39 at%. The higher amount of Mo⁶⁺ in NAPH-MoS₂ can be explained in the same way as for WS₂. It is also possible to track chemical changes in molybdenum and tungsten disulphides by deconvolution of the S 2p core-level, but it is rather difficult to determine the changes with satisfying accuracy because of the low proximity and overlap of individual components. The S 2p high-resolution core-level spectra are shown in Fig. SI4.† In addition the survey spectra are shown in the ESI (Fig. SI5†) together with the individual element concentrations (Table SI2†).

The HR-TEM measurement was performed in order to further investigate the structure of exfoliated MoS₂ and WS₂ at the atomic level. The SAED in combination with EDS was applied to confirm the structure and elemental composition of exfoliated TMDs. Both materials, exfoliated with butyllithium and sodium naphthalenide, are composed of both single and few layer sheets. All exfoliated TMDs with the exception of the WS₂-BuLi sample exhibit wrinkling close to the rim of the sheets (Fig. SI6†). The EDS analysis yielded the S/Mo ratio close to the theoretical value of 2 (2.09 for MoS₂-BuLi and 1.93 for MoS₂-NAPFT). The EDS analysis of WS₂ exfoliated with butyllithium and sodium naphthalenide gave the S/W ratio of 0.41 and 1.23, respectively. Such a large difference in the S/W ratio can be explained by the large difference in tungsten and sulfur molar mass. The selective area electron diffraction (SAED) was performed on selected sheets of TMDs which were further investigated by HR-TEM microscopy. The diffraction shows the presence of both single and few layer TMDs with the hexagonal symmetry of the structure. The results of SAED are shown in the ESI (Fig. SI7†). The HR-TEM can be used to distinguish between 2H and 1T phases. In order to identify different phases, detailed images were obtained. In comparison with BuLi exfoliated TMDs we can observe a high degree of disorder on the rims of TMD sheets exfoliated with sodium naphthalenide. This can be seen in detail in Fig. 4 for exfoliated MoS₂ and WS₂. The disorder in the MoS₂ and WS₂ structures was theoretically predicted to be responsible for room temperature ferromagnetism3⁴ and also experimentally reported for MoS₂ layers with a high concentration of edges.3⁵ Since the samples exfoliated with sodium naphthalenide exhibit a significantly larger degree of exfoliation with a higher concentration of edges and a high degree of disorder, the relationship with the observed room temperature ferromagnetism appears to be quite straightforward.

In order to recognize the presence of the 2H and 1T phases from HR-TEM images, we used the method reported in the literature.3⁰ The MoS₂-NAPFT showed the dominant presence of 1T modifications while for the sample MoS₂-BuLi, we can see both the 1T and 2H structures. This is in qualitative agreement with XPS results, where the higher concentration of the 1T phase was observed for samples exfoliated with sodium naphthalenide. The HR-TEM images of exfoliated MoS₂ and WS₂ are shown in Fig. 5.
In order to better quantify the degree of MoS$_2$ and WS$_2$ exfoliation we performed the surface area measurement using methylene blue adsorption. This method was reported to be useful for surface area measurement of TMDs and other layered compounds such as clay silicates.\(^{36,37}\) The obtained surface area of MoS$_2$ and WS$_2$ exfoliated with butyllithium was 62.1 m$^2$ g$^{-1}$ and 36.1 m$^2$ g$^{-1}$, respectively. A significantly larger area was obtained for MoS$_2$ and WS$_2$ exfoliated using sodium naphthalenide, 88.3 m$^2$ g$^{-1}$ for MoS$_2$·NAPH and 190.2 m$^2$ g$^{-1}$ for WS$_2$·NAPH. The observed differences in the surface area of MoS$_2$ and WS$_2$ exfoliated with sodium naphthalenide and butyllithium are in good agreement with the results obtained by microscopic techniques, where a substantially higher degree of exfoliation was observed for NAPH samples.

The magnetic measurements of butyllithium and sodium naphthalenide exfoliated TMDs were carried out using a SQUID magnetometer. At room temperature ($T = 300$ K) the field dependence of the magnetization was found to be given by a superposition of a diamagnetic moment proportional to the magnetic field (the corresponding susceptibilities $\chi_D$ are listed in Table 2) and a small ferromagnetic (FM) contribution $m_{\text{FM}}$, for which the hysteresis loops are plotted in Fig. 6a.

The hysteresis curve is characterized by a very small coercivity of the order of 100 Oe and by the saturation magnetization $m_S$. The FM contribution is roughly independent of temperature. For samples WS$_2$·BuLi and MoS$_2$·BuLi the FM magnetization, if any, lies under a limit of $2.5 \times 10^{-5}$ emu g$^{-1}$ given by the accuracy of the measurement. At $T = 2$ K all samples exhibit a paramagnetic (PM) behaviour (Fig. 6b) which can be approximately described by the sum of the Brillouin contribution for spin $S$ with the saturated value $m_{\text{PM0}}$ and a small linear contribution with a susceptibility $\chi_1$.

The value $m_{\text{PM0}}$, the estimated spin $S$ as well as the susceptibility $\chi_1$ are given in Table 2. In addition we performed a careful magnetic characterization of the starting bulk MoS$_2$ and WS$_2$. For both bulk materials the room temperature hysteresis curves corrected for a diamagnetic contribution with the susceptibility $\chi_D$ are shown in the ESI (Fig. S18†). The ferromagnetic (FM) magnetization is observed only for the sample MoS$_2$, for which the value of the saturation magnetization $m_S$ (see Table 2) reliably exceeds the error level $2.5 \times 10^{-5}$ emu g$^{-1}$. In this case the hysteresis curve exhibits a finite remanent magnetization of $8.1 \times 10^{-5}$ emu g$^{-1}$ and a coercivity of 240 Oe. However this value is one order of magnitude lower compared to the exfoliated MoS$_2$. For the WS$_2$ sample we can also formally speak about the FM contribution but its value is lower than the error level. At the lowest temperature $T = 2$ K we measured the magnetization curves providing information about the paramagnetic (PM) contribution to the magnetic moment (Fig. 6b). The $m(H)$ curves can be well approximated by the sum of the PM Brillouin contribution for spin $S = 2$, $g = 2$ and a linear susceptibility $\chi_1$ (Table 2). First we shall discuss the values of diamagnetic susceptibility contributions. The theoretical diamagnetic susceptibility for MoS$_2$ and WS$_2$ can be calculated using the data for Mo$^{4+}$, W$^{4+}$ and S$^{2-}$ ions.\(^{38}\) This yields $\chi_D (\text{MoS}_2) = -2.46.10^{-7}$ emu (g Oe)$^{-1}$ and $\chi_D (\text{WS}_2) = -2.46.10^{-7}$ emu (g Oe)$^{-1}$. Whereas for the bulk materials MoS$_2$ and WS$_2$ we could speak about a rough agreement with the measured values of $\chi_D$, a clear disagreement occurs for the sample MoS$_2$·NAPH (having the largest FM magnetization) where the measured value of $\chi_D$ is of the order of $10^{-8}$ emu (g Oe)$^{-1}$. The observed FM magnetization in the bulk is significantly lower compared to the material exfoliated with sodium naphthalenide. As far as the PM magnetization is concerned it is interesting that for the MoS$_2$ bulk sample the Brillouin approximation yields $S = 2$ in contrast to $S = 1.5$ for MoS$_2$·NAPH. This difference originates from the high degree of exfoliation and the 2H–1T phase transition observed by XPS. In the case of so small FM contributions, it is necessary to take into account the possible influence of impurities.

**Table 2** Summary of magnetization measurements on exfoliated as well as bulk MoS$_2$ and WS$_2$

<table>
<thead>
<tr>
<th>Sample</th>
<th>$x_D$ (emu (g Oe)$^{-1}$)</th>
<th>$m_S$ (emu g$^{-1}$)</th>
<th>$m_{\text{PM0}}$ (emu g$^{-1}$)</th>
<th>$\chi_1$ (emu (g Oe)$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>MoS$_2$·NAPH</td>
<td>$-6.818 \times 10^{-8}$</td>
<td>0.0019</td>
<td>0.156</td>
<td>$5.311 \times 10^{-7}$</td>
</tr>
<tr>
<td>WS$_2$·NAPH</td>
<td>$-2.161 \times 10^{-7}$</td>
<td>0.0007</td>
<td>0.164</td>
<td>$6.1 \times 10^{-9}$</td>
</tr>
<tr>
<td>WS$_2$·BuLi</td>
<td>$-3.188 \times 10^{-7}$</td>
<td>$&lt;2.0 \times 10^{-5}$</td>
<td>0.045 ($S = 2$)</td>
<td>$2.1 \times 10^{-8}$</td>
</tr>
<tr>
<td>MoS$_2$·BuLi</td>
<td>$-3.551 \times 10^{-7}$</td>
<td>$&lt;2.0 \times 10^{-5}$</td>
<td>0.01 ($S = 1$)</td>
<td>$4.23 \times 10^{-7}$</td>
</tr>
<tr>
<td>MoS$_2$</td>
<td>$-3.98 \times 10^{-7}$</td>
<td>8.1 $\times 10^{-5}$</td>
<td>0.0473 ($S = 2$)</td>
<td>$2.78 \times 10^{-7}$</td>
</tr>
<tr>
<td>WS$_2$</td>
<td>$-3.461 \times 10^{-7}$</td>
<td>$&lt;2.0 \times 10^{-5}$</td>
<td>0.0266 ($S = 2$)</td>
<td>$2.49 \times 10^{-8}$</td>
</tr>
</tbody>
</table>
mostly represented by iron. Starting from the data reported in the literature, we can infer that the FM saturation magnetizations \( m_s = 0.0019 \text{ emu g}^{-1} \) and \( m_n = 0.0007 \text{ emu g}^{-1} \) detected for MoS\(_2\) and WS\(_2\) correspond to the concentration of Fe impurities 8.6 and 3.18 ppm, respectively.\(^{19}\)

For bulk MoS\(_2\) and WS\(_2\) no FM ordering has been observed so far to the best of our knowledge. However, detailed magnetic measurement of starting MoS\(_2\) and WS\(_2\) showed no or extremely weak magnetization (MoS\(_3\)) at room temperature and also the ferromagnetism of samples MoS\(_2\)-BuLi and WS\(_2\)-BuLi was below the detection limit of the SQUID device. This leads us to the conclusion that the room temperature ferromagnetism observed in sodium naphthenalene exfoliated MoS\(_2\) and WS\(_2\) is associated with several effects which take place during the TMD exfoliation. First we have to take into account the 2H–1T phase transition which is significantly larger compared to the samples exfoliated with butyllithium. The transformation of MoS\(_2\) and WS\(_2\) from 2H to 1T modification was observed by HR-TEM and well documented by XPS spectroscopy from deconvolution of Mo 3d and W 4f HR-XPS spectra, respectively. Moreover, a larger surface area imposes a higher concentration of edges, where a large amount of defects can be observed in HR-TEM images. Since NAPH samples possess a high degree of exfoliation, this will be one of the main factors necessary for the room temperature ferromagnetic ordering caused by the high concentration of edges related to a high degree of exfoliation and possibly also connected to the defects present on the edges. These defects may also lead to a PM contribution which superimposes on the intrinsic diamagnetic component. This manifests itself in the reduction of the measured susceptibility \( \chi_0 \) as was observed, e.g., in the case of the sample MoS\(_2\)-NAPH. Let us mention that the FM state induced by defects was predicted for various types of defects in MoS\(_2\) and also experimentally observed on nanostructured MoS\(_2\) layers with a high concentration of defects.\(^{15,40}\)

We also have to note one possible origin of ferromagnetism, which has been recently reported on another transition metal dichalcogenide – ReS\(_2\). The room temperature ferromagnetism was observed on the wrinkled area of ReS\(_2\).\(^{31}\) Such a wrinkling of the layers was also observed in the TEM images of MoS\(_2\) and WS\(_2\) sheets exfoliated with sodium naphthenalide. Although this feature was also observed in the MoS\(_2\) exfoliated with butyllithium, a significantly lower degree of exfoliation led only to a minor influence on the magnetic behavior. Another interesting finding is the presence of a small PM contribution corresponding to the concentration of \( 10^{-3}–10^{-4} \) of PM centers with a spin \( S \). The spin \( S = 1 \) could be connected with the occurrence of a spin-pairing of d\(^2\) centers mentioned by other authors.\(^{7}\) Our results suggest larger values of spins for the exfoliated WS\(_2\) samples.

**Conclusions**

The detailed investigation of MoS\(_2\) and WS\(_2\) exfoliated using various alkali metal intercalating compounds like butyllithium and sodium naphthenalide yields materials with significant differences in the degree of exfoliation. The exfoliation using butyllithium predominantly resulted in a few layer materials, while the exfoliation using sodium naphthenalide led to almost complete exfoliation and formation of single layer TMDs. Such an intensive process of intercalation and subsequent exfoliation is associated with the 2H to 1T phase transition and formation of defects on the edges of TMD sheets. The presence of a metallic phase (1T) and formation of defects in the TMD structure can be responsible for the exotic ferromagnetic behavior which was observed up to room temperature. Even though the room temperature ferromagnetism is of the order of \( 1 \times 10^{-3} \text{ emu g}^{-1} \), this effect was not observed in bulk MoS\(_2\) and WS\(_2\) or even in the butyllithium exfoliated material. Butyllithium exfoliated samples exhibit a significantly lower degree of exfoliation and subsequently a lower concentration of defects on the edges of individual sheets. This led us to draw a conclusion that the room temperature ferromagnetism is closely related to the structural disorder on the sheet edges and the presence of a conductive metallic phase. Our finding can have a strong impact on the future applications of transition metal dichalcogenides in the construction of novel spintronic devices.

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**Notes and references**