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## Formation of hexagonal-molybdenum trioxide (h-MoO<sub>3</sub>) nanostructures and their pseudocapacitive behavior†

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The crystallographic structure and morphology of redox active transition metal oxides have a pronounced effect on their electrochemical properties. In this work, h-MoO<sub>3</sub> nanostructures with three distinct morphologies, *i.e.*, pyramidal nanorod, prismatic nanorod and hexagonal nanoplate, were synthesized by a facile solvothermal method. The morphologies of h-MoO<sub>3</sub> nanostructures were tailored by a controlled amount of hexamethylenetetramine. An enhanced specific capacitance about 230 F g<sup>-1</sup> at an applied current density of 0.25 A g<sup>-1</sup> was achieved in h-MoO<sub>3</sub> pyramidal nanorods. Electrochemical studies confirmed that the h-MoO<sub>3</sub> pyramidal nanorods exhibit superior charge-storage ability. This improved performance can be ascribed to the coexistence of its well exposed crystallographic planes with abundant active sites, *i.e.*, hexagonal window (HW), trigonal cavity (TC) and four-coordinated square window (SW). The mechanism of charge-storage is likely facilitated by the vehicle mechanism of proton transportation due to the availability of the vehicles, *i.e.*, NH<sub>4</sub><sup>+</sup> and H<sub>2</sub>O. The promising, distinct and unexploited features of h-MoO<sub>3</sub> nanostructures reveal a strong candidate for pseudocapacitive electrode materials.

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### 1. Introduction

The physicochemical properties and electronic structure of transition metal oxides have been a subject of intense fundamental research.<sup>1,2</sup> Transition metal oxides are considered promising candidates for electrochemical energy-storage (supercapacitors and batteries) applications due to their excellent reduction-oxidation (redox) properties (M<sup>n+</sup> + e<sup>-</sup> → M<sup>(n-1)+</sup>).<sup>3,4</sup> Recently, supercapacitors have emerged as promising electrochemical energy-storage devices due to their spectacular features such as higher power density compared to batteries (~10<sup>4</sup> W kg<sup>-1</sup> for supercapacitors and ~100 W kg<sup>-1</sup> for batteries), quick charging/discharging (ms–min) and excellent cycling stability (10<sup>4</sup>–10<sup>6</sup>) in aqueous as well as non-aqueous

electrolytes.<sup>5</sup> Supercapacitors have attracted considerable attention because of their potential use in vehicle start/stop/start operation and micro-electro-mechanical and portable electronic devices.<sup>6,7</sup> To date, binary transition metal oxides (MnO<sub>2</sub>, CoO<sub>3</sub>, V<sub>2</sub>O<sub>5</sub>, *etc.*)<sup>8,9</sup> have gained enormous attention as prospective electrode materials for supercapacitor applications (based on intercalation pseudocapacitance and surface pseudocapacitance). The capacitive behavior can be improved by tuning the state of crystallinity of the materials, *e.g.*, amorphous or crystalline. However, this approach is found to be strongly materials system dependent, for example Zheng *et al.*<sup>10,11</sup> and Farsi *et al.*<sup>12</sup> showed that the amorphous RuO<sub>2</sub>·xH<sub>2</sub>O and amorphous MoO<sub>3</sub> nanostructures store more charge than their crystalline counterparts, respectively, while recently Brezesinski *et al.*<sup>13</sup> and Kim *et al.*<sup>14</sup> found that the crystalline MoO<sub>3</sub> and Nb<sub>2</sub>O<sub>5</sub> are superior to their amorphous counterparts as charge-storage carriers.

Nanostructuring of the binary transition metal oxides in various crystallographic structures has also been considered an effective and feasible approach to improve the electrochemical performance of the electrode materials.<sup>15</sup> To date, layered binary transition metal oxides have been studied frequently due to their two dimensional crystallographic structure, which provide effective paths for the conduction of electrolyte ions. However, binary transition metal oxides with three dimensional (3D) crystallographic structures (open-structure), which can provide 3D conduction paths for electrolyte ions have not been explored. The 3D crystallographic structure

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†Electronic supplementary information (ESI) available: Transmission electron microscopy (TEM) and high resolution transmission electron microscopy (HRTEM) of pyramidal nanorods; comparison of reaction times to prepared h-MoO<sub>3</sub> nanostructures, FTIR spectrum of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates; atomic representation of hexagonal framework; Cyclic voltammograms (CVs) of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates at higher scan rates; BET isotherms of all three samples; charge/discharge profiles of pyramidal nanorods; prismatic nanorods and hexagonal nanoplates at the various current densities; FESEM images of pyramidal nanorods before and after cycling; Bode plot of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates; electrochemical performance comparison table. See DOI: 10.1039/c5nr01505g



primarily exists in the metastable phase of binary transition metal oxides such as spinel-MnO<sub>2</sub>, beta-V<sub>2</sub>O<sub>5</sub>, MoO<sub>3</sub>, etc.<sup>9,10,13,15</sup>

Among them MoO<sub>3</sub> is considered an interesting material due to the high electrochemical activity of its stable as well as metastable polymorphs.<sup>16–18</sup>

The thermodynamic stable phase of MoO<sub>3</sub>, *i.e.*,  $\alpha$ -MoO<sub>3</sub> possesses a layered (ABA) framework structure (*i.e.*, 2D), while its metastable polymorphs possess an open-structure (*i.e.*, 3D) in the hexagonal phase (h-MoO<sub>3</sub>), a layered (AAA) framework structure in the MoO<sub>3</sub>-II phase and rutile or distorted rutile structure in the  $\beta$ -MoO<sub>3</sub> phase.<sup>19</sup> To date, the electrochemical properties of  $\alpha$ -MoO<sub>3</sub> have been investigated. Various strategies have been employed to enhance the electrochemical performance of  $\alpha$ -MoO<sub>3</sub>, for example nanowire arrays<sup>20</sup> core/shell architecture<sup>21</sup> decorations of conducting supports (polymers)<sup>22</sup> and incorporation of CNT and nanocomposites with other oxides.<sup>23,24</sup> However, these efforts did not improve the specific capacitance significantly. Recently, a kinetically favoured intercalation mechanism was proposed, benefiting from the mesoporous and iso-oriented nanocrystalline domain of the layered  $\alpha$ -MoO<sub>3</sub>,<sup>13</sup> for improved Li<sup>+</sup> ions storage (605 C g<sup>-1</sup>). In addition to all the efforts made to date, several computational and experimental studies have shown that the crystal plane structure of the electrode materials has a substantial influence on the electrochemical performance of materials.<sup>25–27</sup> Therefore, size and shape controlled nanomaterials are essential in unleashing and understanding their intrinsic properties, which are critical in electrochemical applications. More recently, the highest achievable specific capacitance of 326 F g<sup>-1</sup> (at 0.25 A g<sup>-1</sup>) using  $\alpha$ -MoO<sub>3</sub> nanobelts has been reported, but the loading mass tends to be low (~0.51 mg).<sup>28</sup> Also, the specific capacitance of  $\alpha$ -MoO<sub>3</sub> nanoplates as high as 280 F g<sup>-1</sup> (calculated from non-rectangular CV curves) has been achieved, benefiting from the small diffusion length of the ions into the nanosheets.<sup>29</sup> Most recently, Chang *et al.* introduced the concept of the work function difference to enhance the working potential window of an asymmetric supercapacitor device based on rGO- $\alpha$ -MoO<sub>3</sub>, in an aqueous electrolyte solution.<sup>30</sup> Due to an inter-layer spacing (6 Å) between the two layers of  $\alpha$ -MoO<sub>3</sub>, it can accommodate a massive amount of electrolyte ions (0.28 Å for H<sup>+</sup>), but its poor conductivity limits the diffusion of the ions. Compared to  $\alpha$ -MoO<sub>3</sub> that has intercalation sites along one dimension only, the open-structure of h-MoO<sub>3</sub> offers an increase in the degree of freedom (from 1 to 3) for electrolyte ions to intercalate in all possible directions. The hexagonal phase of MoO<sub>3</sub> (h-MoO<sub>3</sub>) can be considered a promising host material for electrolyte ions because of the availability of the various intercalation sites, *i.e.*, hexagonal window (HW), trigonal cavity (TC) and four-coordinated square window (SW). The open-structure of h-MoO<sub>3</sub> could facilitate intercalation as well as diffusion of ions to realize improved supercapacitor performance.

Taking into account all these factors, we have synthesized h-MoO<sub>3</sub> nanostructures using a facile solvothermal method. The synthesis of h-MoO<sub>3</sub> nanostructures with diverse mor-

phologies has been rarely found; this is related to the difficulty in phase formation due to thermodynamic (formation enthalpy and entropy of constituents) barriers.<sup>31–33</sup> We have previously attempted to tune the morphology of h-MoO<sub>3</sub> using thiourea as the source of ammonium ions. With the increase in the loading amount of the ammonium ions, a predominant formation of a flower-like morphology was observed. In this work, we are able to circumvent the challenges limiting pseudo-capacitive behaviour by judicious control of the synthesis in tuning the physical structure of the h-MoO<sub>3</sub> nanostructures.

## 2. Experimental

### 2.1. Chemical used

All chemicals, Mo powder 99.9%, H<sub>2</sub>O<sub>2</sub> (30%), HNO<sub>3</sub> (conc.), potassium hydroxide (KOH) flakes and hexamethylenetetramine (HMTA) were used as received from Alpha Aesar and Sigma Aldrich, respectively, without further purification.

### 2.2. Preparation of h-MoO<sub>3</sub> nanostructures

To synthesize h-MoO<sub>3</sub> with pyramidal nanorod, prismatic nanorod and hexagonal nanoplate-like morphology, a simple yet effective solvothermal approach was employed, modified from our previous report<sup>32</sup> with the use of hexamethylenetetramine (HMTA) as a structure directing agent instead of thiourea (CS(NH<sub>2</sub>)<sub>2</sub>). In brief, peroxomolybdic solution was prepared by the addition of hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>) into a 20 ml glass vial containing 0.50 g of Mo powder. In a typical synthesis procedure, 100 mg, 150 mg and 250 mg of HMTA in 10 ml of DI water were used. The solution of HMTA was slowly added into peroxomolybdic solution and left for stirring. The reaction times to prepare h-MoO<sub>3</sub> are compared with the other synthesis routes, and presented in the ESI, Table S1.† The solvothermal reaction was carried out at 100 °C for 18 h. Finally, the precipitates were collected and washed thoroughly with ethanol and DI water to remove loosely bound or bulk particles. The as-obtained precipitates were dried at 70 °C to obtain the final product.

### 2.3. Electrode preparation for electrochemical testing

The working electrodes were prepared by mixing the as-prepared h-MoO<sub>3</sub> powder (85%), carbon black (10%) and polyvinylidene fluoride (PVDF) (5%) into a 5 ml glass vial containing 500  $\mu$ l of *N*-methyl-2-pyrrolidone (NMP), under magnetic stirring to make a slurry. 20  $\mu$ l of this slurry was coated onto a current collector (graphite paper, 1 cm<sup>2</sup>) and dried in a laboratory oven at 100 °C. The loading mass was measured to be 0.95  $\pm$  0.05 mg, for all three samples. Electrochemical testing was carried out in an aqueous solution of H<sub>2</sub>SO<sub>4</sub> (1 M) in a potential window of 0.05 V to 0.65 V vs. Ag/AgCl.

### 2.4. Structural characterization

The samples were characterized using powder X-ray diffractometer (Shimadzu XRD - 6000, Cu K $\alpha$  radiation  $\lambda$  = 1.54 Å; power 2 kW) at a scan rate of 1° min<sup>-1</sup> in the 2 $\theta$  range of



10°–60°. Field emission scanning electron microscopy (FESEM; JEOL, JSM 7600F thermal FEG, JSM 6340F cold cathode FEG) and transmission electron microscopy (TEM; JEOL, JEM 2010 and JEM 2100F) was used to evaluate the morphology of the samples. Thermal analysis was done using thermal gravimetric analysis (TGA) (Q 500, max. temperature 900 °C), and IR spectroscopy (Perkin Elmer, Model Spectrum GX) was used to obtain information about the state of the bonds present. X-ray photoelectron measurements (XPS) were carried out using a VG ESCALAB 220I-XL spectrometer with a twin monochromated Al K $\alpha$  X-ray source (1486.6 eV) at a constant dwell time of 50 ms and a pass-energy of 50 eV. CasaXPS peak fitting software was used to perform curve fitting. The physical surface area (Brunauer–Emmett–Teller) and pore size distribution were measured using a Tristar-II surface area and porosity analyser. Electrochemical measurements were carried out using Solartron, S1470E electrochemical interface and AutoLab PGSTAT 30 potentiostat. A three electrode cell was used for electrochemical testing, in which graphite paper coated with active materials, platinum sheet and saturated Ag/AgCl were employed as the working, counter and reference electrodes, respectively, in an aqueous electrolyte (1 M H<sub>2</sub>SO<sub>4</sub>). Electrochemical impedance spectroscopy (EIS) measurements were carried out by applying an AC voltage with 1 mV amplitude in a frequency range from 100 mHz to 100 kHz at open circuit potential (OCP).

### 3. Results and discussion

#### 3.1. Structural characterization

The crystallographic phases of the as-prepared samples were evaluated using X-ray diffraction (XRD), as shown in Fig. 1. All

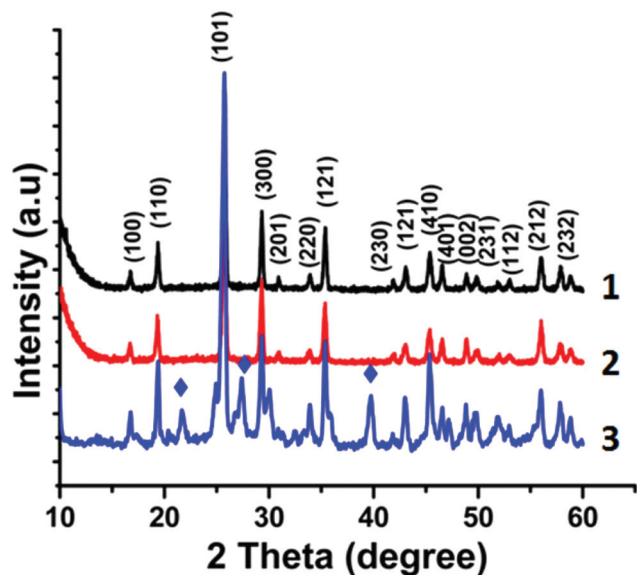


Fig. 1 X-Ray diffraction patterns of pyramidal nanorods (1), prismatic nanorods (2) and hexagonal nanoplates (3). The nanoplates show some impurity peaks labelled in blue diamond (◆).

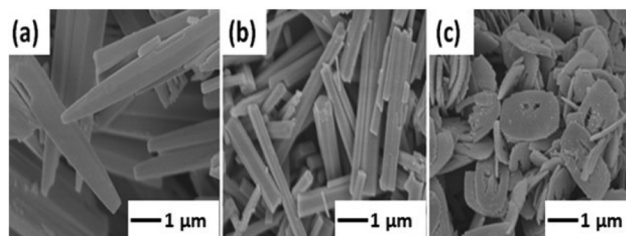


Fig. 2 FESEM micrographs of the samples prepared using (a) 100 mg, (b) 150 mg and (c) 250 mg of HMTA in peroxomolybdic acid solution.

the XRD peaks in the patterns were matched exclusively with the hexagonal phase of MoO<sub>3</sub> (ICDD # 15-5332) with lattice parameters  $a = 10.54 \text{ \AA}$ ,  $c = 3.72 \text{ \AA}$  and cell volume  $V = 359.37 \text{ \AA}^3$ .<sup>34</sup> A few impurity peaks (◆) were observed in the pattern of hexagonal nanoplates (probably due to the residue of the by-product formed by Mo<sup>5+</sup> ions, which made complex ammonium hydroxides of Mo, *i.e.*, (NH<sub>4</sub>)<sub>8</sub>Mo<sub>10</sub>O<sub>34</sub>). The major peaks positioned at  $2\theta \sim 25.5^\circ$  and  $19.2^\circ$  are characteristic of the hexagonal phase of MoO<sub>3</sub>.

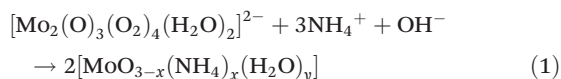
The morphology of the as-prepared h-MoO<sub>3</sub> nanostructures was evaluated using FESEM, as shown in Fig. 2. Fig. 2(a)–(c) depict pyramidal nanorods, prismatic nanorods and the hexagonal nanoplate-like morphology of the samples prepared using 100 mg, 150 mg and 250 mg of HMTA, respectively. The typical diameter of pyramidal nanorods was about 400 and 100 nm at the base and tip, respectively, while the diameter of prismatic nanorods was about 350 nm. The length and thickness of the hexagonal nanoplates were about 1 μm and 80 nm, respectively. The effect of the concentration of ammonium ions on the nucleation and growth is critical in controlling the morphology of the as prepared nanostructures, as explained in our previous report using thiourea as the source of the ammonium ions.<sup>32</sup>

The concentration of HMTA or ammonium (NH<sub>4</sub><sup>+</sup>) ions is essential to control the morphology of the final products. HMTA is a non-ionic, hetero-cyclic organic compound with the chemical formula (CH<sub>2</sub>)<sub>6</sub>N<sub>4</sub>, which eventually dissociates into NH<sub>4</sub><sup>+</sup> and OH<sup>−</sup> ions upon hydrolysis under the ambient conditions.<sup>35</sup> The reaction of molybdenum (Mo) powder with the hydrogen peroxide (H<sub>2</sub>O<sub>2</sub>) generates a variety of species attached to 1 to 6 peroxide groups.<sup>36,37</sup> The attachment of the peroxide groups to the molybdic units strongly depends on the pH of the medium. A yellow peroxomolybdic solution was obtained in the acidic medium (pH  $\sim$  1–1.5), which mainly consists of monomer- and dimer-peroxo species, *i.e.*, [MoO(O<sub>2</sub>)<sub>2</sub>(H<sub>2</sub>O)<sub>2</sub>], [O{MoO(O<sub>2</sub>)<sub>2</sub>(H<sub>2</sub>O)<sub>2</sub>}]<sup>2−</sup> and other even more complex poly-peroxo species.<sup>38</sup>

The chemistry of the HMTA with the peroxomolybdic acid solution generates the nuclei and forms the lattice of molybdenum trioxide (MoO<sub>3</sub>) in the hexagonal phase. The interaction of NH<sub>4</sub><sup>+</sup> and OH<sup>−</sup> ions (hydrolysed product of



HMTA) with the peroxomolybdic species can be expressed according to the following reaction (1):



where  $x$  and  $y$  cause the alteration of the morphology of the h-MoO<sub>3</sub> nanostructure. The NH<sub>4</sub><sup>+</sup> ions reside in the tunnel of the hexagonal framework and are responsible for maintaining the stability of the framework.<sup>39</sup>

When a small amount (100 mg) of HMTA was incorporated into the peroxomolybdic acid solution, a pyramidal or pencil-like morphology of h-MoO<sub>3</sub> predominates. This morphology is due to the depletion of the NH<sub>4</sub><sup>+</sup> ions during the synthesis reaction. The NH<sub>4</sub><sup>+</sup> ions continuously alter the interaction of the growth species [Mo<sub>8</sub>O<sub>26</sub>]<sup>4-</sup> on the crystal surfaces. At the instant NH<sub>4</sub><sup>+</sup> ions are depleted, the growth of the low-energy planes at the tip of the pyramidal nanorods leads to crystal structure stability. As the loading amount of HMTA was increased (150 mg), a perfect prismatic or hexagonal morphology could be realized. It is interesting to note that the hexagonal nanoplates were obtained with a relatively high amount (250 mg) of HMTA. The as-obtained NH<sub>4</sub><sup>+</sup> ions are pivotal in the formation of the hexagonal framework and simultaneously form partially-reduced molybdenum cations (Mo<sup>6+</sup> to Mo<sup>6-x</sup>). Further growth of the hexagonal framework is hindered by the dissolution of the compound formed by Mo<sup>6-x</sup> cations in the vicinity (due to excessive ammonium ions), as the course of reaction proceeds. The surfaces of the as-prepared hexagonal nanoplates were found to be extremely rough or defective, which could probably due to the dissolution of the product formed by Mo<sup>6-x</sup> ions. This phenomenon was not prevalent during the formation of pyramidal and prismatic nanorods due to a relatively low amount of NH<sub>4</sub><sup>+</sup> ions in the reaction.

The morphologies of the as-obtained nanostructures were further characterized by TEM. Fig. 3(a)–(c), unambiguously depict the pyramidal nanorods, prismatic nanorods and hexagonal nanoplate-like morphology, respectively, which are consistent with FESEM analysis.

Fig. 3(d)–(f) show the HRTEM micrographs of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates, respectively. The spacings between two neighbouring fringes were found to be 0.21 nm and 0.18 nm, as can be seen in Fig. 3(d) and (e), corresponding to the  $d$ -spacing of the (221) plane and (002) plane of h-MoO<sub>3</sub>, respectively. The growth direction of the bulk of pyramidal nanorods was found similar to the prismatic nanorods (ESI, Fig. S1†). The interplanar spacings in hexagonal nanoplates were related to the (101) and (210) planes of h-MoO<sub>3</sub>, as shown in Fig. 3(f). Furthermore, the selected area electron diffraction (SAED) and fast Fourier transformation (FFT) patterns (Fig. 3(g)–(i)) show the diffraction spots along the [110], [100] and [1–2–1] zone axis for the tip of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates, respectively.

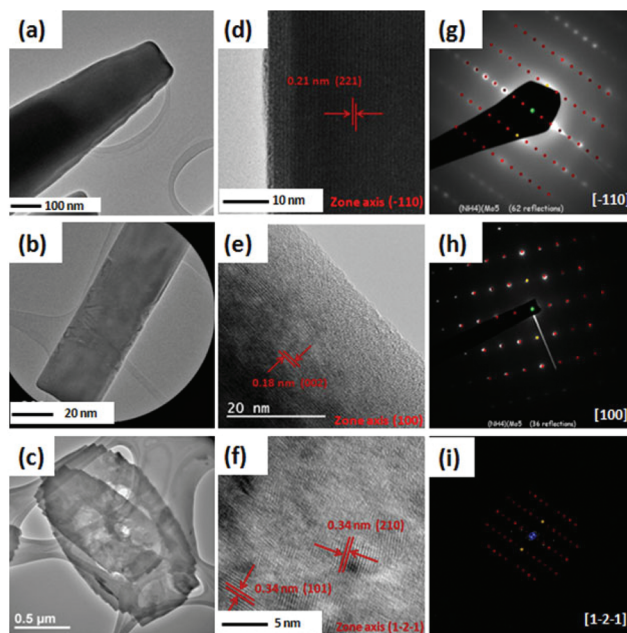


Fig. 3 TEM micrographs (a)–(c), HRTEM micrographs (d)–(f) and SAED and FFT patterns (g)–(i) of the samples prepared using 100 mg, 150 mg and 250 mg of HMTA, respectively, are shown. The red and yellow spots in SAED images show standard patterns of h-MoO<sub>3</sub>.

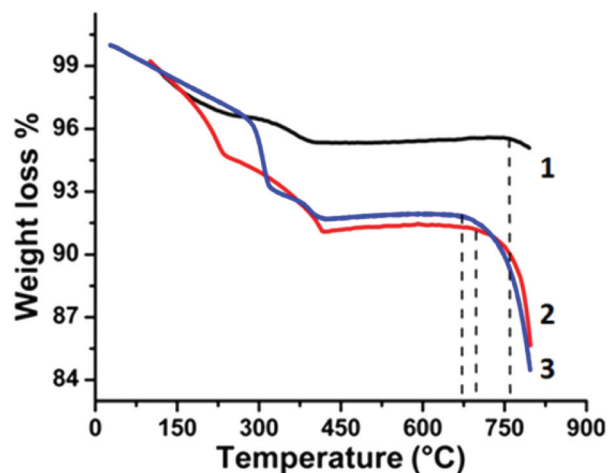


Fig. 4 TGA analyses of (1) pyramidal nanorods, (2) prismatic nanorods and (3) hexagonal nanoplates. The step-wise weight loss of the materials corresponds to removal of water and ammonium molecules from the lattice of h-MoO<sub>3</sub>.

The chemical purity and thermal stability of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates were evaluated using TGA. Fig. 4 shows the typical characteristics of the ammonium ion-assisted formation of h-MoO<sub>3</sub>.<sup>32</sup> The amount of water and ammonia molecules was calculated to be 3.4%, 1.3% for pyramidal nanorods, 4.5%, 2.5% for prismatic nanorods and 3.35%, 4.95% for hexagonal nanoplates, respectively, taking their weight loss characteristics into account. The



weight losses in the temperature range of 150 °C–300 °C and 300 °C–450 °C for all three samples correspond to the removal of water and ammonium molecules, respectively.

The major weight losses of the material occur at 760 °C, 698 °C and 670 °C for pyramidal nanorods, prismatic nanorods and hexagonal nanoplates, respectively. The weight loss at different temperatures suggests that the stability of the MoO<sub>6</sub> octahedra decreases from pyramidal nanorods to prismatic nanorods to hexagonal nanoplates. Approximate formula units (NH<sub>4</sub>)<sub>0.11</sub>(H<sub>2</sub>O)<sub>0.28</sub>MoO<sub>2.89</sub>, (NH<sub>4</sub>)<sub>0.23</sub>(H<sub>2</sub>O)<sub>0.39</sub>MoO<sub>2.77</sub> and (NH<sub>4</sub>)<sub>0.44</sub>(H<sub>2</sub>O)<sub>0.30</sub>MoO<sub>2.56</sub> are derived for pyramidal nanorods, prismatic nanorods and hexagonal nanoplates, respectively, taking thermal analysis into consideration. Effects of the concentration of NH<sub>4</sub><sup>+</sup> ions on the bonding between Mo and O was further evaluated using IR spectroscopy, as shown in the ESI, Fig. S2 and Table S2.†

XPS measurements were performed to analyse the valance state and chemical shift of Mo and N atoms in all three types of samples. Fig. 5(a)–(c) show the XPS spectra of Mo atoms in pyramidal nanorods, prismatic nanorods and hexagonal nanoplates, respectively. Fig. 5(a)–(c) consist of two peaks which correspond to spin orbit doublets of Mo<sup>6+</sup> (3d<sub>5/2</sub>) and Mo<sup>6+</sup> (3d<sub>3/2</sub>).<sup>40</sup> The Mo<sup>6+</sup> (3d<sub>5/2</sub>) and Mo<sup>6+</sup> (3d<sub>3/2</sub>) peaks are centred at 235.4 eV and 238.5 eV for pyramidal nanorods, 234.9 eV and 238.2 eV for prismatic nanorods, and 234.2 eV and 238 eV for hexagonal nanoplates. With respect to the binding energies of the deconvoluted Mo<sup>6+</sup> peaks in pyramidal nanorods, a negative shift of about 0.5 eV and 0.7 eV in the binding energies is identified for prismatic nanorods and hexagonal nanoplates, respectively. This shift in the binding energies indicates that the valance state of Mo decreases from pyramidal nanorods to prismatic nanorods to hexagonal nanoplates. Besides a shift in binding energies, the integrated intensity of the deconvoluted Mo<sup>6+</sup> peaks decreases from pyramidal nanorods to hexagonal nanoplates, which suggests that the pyramidal nanorods possess a higher content of Mo<sup>6+</sup>. The change in the binding

energies can be ascribed to the change in the NH<sub>4</sub><sup>+</sup> content in all three samples. To verify this change in NH<sub>4</sub><sup>+</sup> content, XPS analyses of N 1s atoms were performed, as depicted in Fig. 5(d–f). The N 1s peak for all three samples centred at 401 eV originates from NH<sub>4</sub><sup>+</sup> ions.<sup>41</sup> The integrated intensity of N 1s increases comparing the spectra from pyramidal nanorods to prismatic nanorods to hexagonal nanoplates, which indicates that the hexagonal nanoplates possess a higher loading amount of NH<sub>4</sub><sup>+</sup> ions. Due to the relatively high loading amount of NH<sub>4</sub><sup>+</sup> ions in hexagonal nanoplates, the binding energies of deconvoluted Mo<sup>6+</sup> peaks are found less positive than that of the pyramidal nanorods. A shake-up satellite peak for all three samples is identified at about 404 eV that could be due to nonstoichiometric NH<sub>4</sub><sup>+</sup> ions generated by N–O complexes, but not much is known about them.<sup>42</sup> The relative amount of these oxidized nitrogen complexes is found to be much smaller than that of the N–H complexes. XPS analyses show the same trend in the amount of NH<sub>4</sub><sup>+</sup> ions as revealed by TGA and IR analyses.

### 3.2. Electrochemical testing of the samples

It is widely accepted that the electrochemical properties of the materials can be altered by tuning the morphology and size of the materials.<sup>43–45</sup> The electrochemical measurements of h-MoO<sub>3</sub> with three different morphologies, *i.e.*, pyramidal nanorods, prismatic nanorods and hexagonal nanoplates, for supercapacitor application, were carried out in a three electrode cell configuration using 1 M H<sub>2</sub>SO<sub>4</sub> solution. Cyclic voltammograms (CVs) of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates were carried out at scan rates of 1–50 mV s<sup>-1</sup>, as shown in Fig. 6(a) and (b), respectively. The strong redox peaks in CVs of h-MoO<sub>3</sub> with different morphologies indicate that the charges are stored in association with the faradic redox reactions.<sup>46,47</sup> A slight shift in the peak potential signifies the effect of morphologies on the electrochemical activities of the electrode materials.<sup>43,48</sup> The origin of this polarization lies in the fact that the reaction potential (over potential) is extremely sensitive towards the activity of the surfaces. From the electrochemical measurements, we found that the pyramidal nanorods exhibit superior charge-

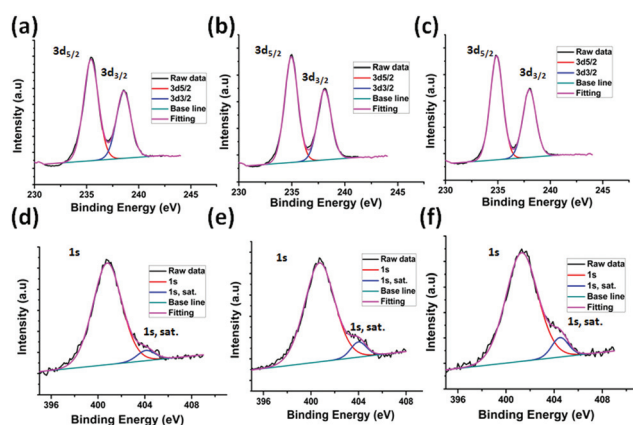


Fig. 5 XPS analysis of Mo components in (a) pyramidal nanorods, (b) prismatic nanorods and (c) hexagonal nanoplates, and analysis of N components in (d) pyramidal nanorods, (e) prismatic nanorods and (f) hexagonal nanoplates.

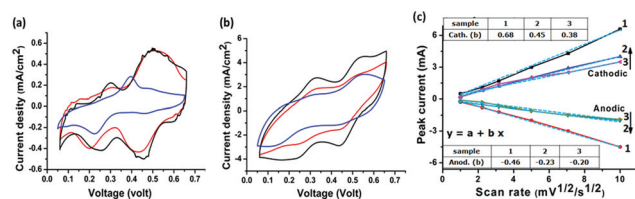
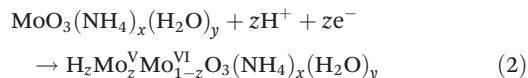


Fig. 6 Cyclic voltammograms (CV) of (1) pyramidal nanorods, (2) prismatic nanorods and (3) hexagonal nanoplates are carried out at scan rate of (a) 1 mV s<sup>-1</sup> and (b) 50 mV s<sup>-1</sup>. The redox peaks in CV analysis indicate reversible ingress and degrees of electrolyte ions, and (c) peak current response of the as-prepared electrodes with respect to the square root of the scan rates indicates diffusion controlled charge transfer process. The dotted lines in (c) indicate fitting profiles according to equation  $y = a + bx$ .



storage capacity compared to prismatic nanorods or hexagonal nanoplates. In the case of hexagonal nanoplates, there is one set of strong redox peaks at 0.22 V/0.39 V ( $\Delta E = 170$  mV) and another set of faint peaks appearing at 0.45/0.49 ( $\Delta E = 40$  mV). For the pyramidal nanorods there are two sets of strong redox peaks positioned at 0.24 V/0.29 V ( $\Delta E = 50$  mV) and 0.46 V/0.50 V ( $\Delta E = 40$  mV), respectively. These peaks correspond to the reversible ingress/digress of  $H^+$  ions into/out of the framework of h-MoO<sub>3</sub>, according to reaction (2):



The intercalation of  $H^+$  ions into the h-MoO<sub>3</sub> is favoured by the redox centres ( $\text{Mo}^{6+}$  ions), which are located at the various intercalation sites, *i.e.*, hexagonal window (HW), hexagonal cavity (TC) and square window (SW), in the framework of h-MoO<sub>3</sub>. These intercalation sites are available in the range of potentials, *i.e.*,  $0.1 \leq V \leq 0.3$  (SW),  $0.3 \leq V \leq 0.6$  (TC) and  $V \geq 0.7$  (HW).

It is commonly believed that the proton conduction in hydrated compounds occur *via* the Grotthuss ( $\text{H}_2\text{O}$ ,  $\text{H}_5\text{O}_2^+$ ,  $\text{H}_3\text{O}^+$ ,  $\text{OH}^-$  *etc.*) or Vehicle ( $\text{H}_2\text{O}$ ,  $\text{NH}_4^+$ ,  $\text{N}_2\text{H}_5^+$ ,  $\text{CH}_3\text{NH}_3^+$ , *etc.*) type mechanism facilitated by the water molecules or proton carriers.<sup>10,49</sup> In our case, due to the availability of proton carriers (Vehicle), *i.e.*,  $\text{NH}_4^+$  and  $\text{H}_2\text{O}$  molecules, the charge-storage mechanism may favour the Vehicle mechanism of proton transportation. Vehicle mechanism is readily encountered in aqueous solutions, liquid/melts and solids. In solids, Vehicle mechanism is usually limited to the materials with layered or open-structures.<sup>50</sup> It is believed that the compounds with a smaller amount of water molecules (<2  $\text{H}_2\text{O}$ ) conduct protons by Vehicle mechanism.<sup>51</sup> Here, we considered that the charge-storage mechanism in h-MoO<sub>3</sub> does not rely on the Grotthuss mechanism of proton transportation due to the open-structure of h-MoO<sub>3</sub>, a small amount of water molecules and the availability of proton carriers ( $\text{NH}_4^+$  and  $\text{H}_2\text{O}$ ). Additionally, in the crystal structure of h-MoO<sub>3</sub> (shown in the ESI, Fig. S3†), water molecules are coordinatively attached with the octahedral units, which do not favour the formation of Zundel ( $\text{H}_5\text{O}_2^+$ ) cations (essential for the Grotthuss mechanism).<sup>52</sup> Therefore, the charge-storage mechanism in h-MoO<sub>3</sub> is more likely facilitated by Vehicle mechanism of proton transportation.

The scan rate response of the as-prepared electrodes was also carried out at different scan rates (1  $\text{mV s}^{-1}$  to 100  $\text{mV s}^{-1}$ ), as depicted in Fig. 6(c). The anodic and cathodic peak currents (in a potential window of 0.4 V to 0.6 V) increase linearly with the increase in scan rates (CVs are shown in the ESI, Fig. S4†). Additionally, linear dependence of the peak currents ( $I_{\text{cathodic}}$  and  $I_{\text{anodic}}$ ) on the square root of the scan rate ( $\nu^{1/2}$ ) describes the process as a diffusion controlled electron transfer process.<sup>46,47</sup> The slope of curves, shown in Fig. 6(c), can be used to estimate the diffusion coefficient ( $D$ ) of the redox-species ( $H^+$  ions) using the Randle Sevcik equation (3),<sup>46</sup> and

this relationship can also be used to evaluate the electrochemical surface area of the electrode:<sup>54</sup>

$$I_p = 2.69 \times 10^5 n^{3/2} ACD^{1/2} \nu^{1/2} \quad (3)$$

where  $I_p$  is the peak current,  $n$  is the number of the electrons involved in the reaction ( $n = 1$ ),  $A$  is the area of the electrode ( $\text{cm}^2$ ),  $C$  is the concentration of the redox-species (1 M),  $D$  is the diffusion coefficient of  $H^+$  ions in MoO<sub>3</sub> ( $1.55 \times 10^{-10} \text{ cm}^2 \text{ s}^{-1}$ ),<sup>53</sup> and  $\nu$  is the scan rate. When the peak currents are plotted as a function of scan rates, then the electrochemically active surface area of the electrodes will be directly proportional to the slope of the cathodic or anodic peak currents *vs.* the scan rate curves, as shown in Fig. 6(c). Thus, the electrochemically active surface area of the pyramidal nanorods is found to be about 1.5 times the area of prismatic nanorods and about 2 times the area of hexagonal nanoplates. It is noteworthy that the BET surface area (ESI, Fig. S5†) of pyramidal nanorods ( $2.2 \text{ m}^2 \text{ g}^{-1}$ ) and prismatic nanorods ( $1.92 \text{ m}^2 \text{ g}^{-1}$ ) is commensurate with the electrochemically active surface area (with similar pore size of 15 nm and 20 nm for pyramidal nanorods and prismatic nanorods samples, respectively). However, the hexagonal nanoplate sample with a higher BET surface area ( $7.86 \text{ m}^2 \text{ g}^{-1}$ ) and larger pore size ( $\sim 40$  nm) does not translate to an electrochemically active surface area, likely due to the ineffective bonding of electrolyte ions at the pore walls with reduced electric field that leads to sluggish redox reactions.<sup>55</sup> Therefore, it is evident from the peak current *vs.* scan rate analysis that the pyramidal nanorods exhibit a larger electrochemical active area, which results in improved electrochemical performance.

Fig. 7(a) shows the galvanic charge/discharge curves of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates at an applied current density of  $0.25 \text{ A g}^{-1}$  (charge/discharge curves at higher current densities are shown in the ESI, Fig. S6†). The non-linearity associated with the charge/discharge curves further confirms the pseudocapacitive response of the as-prepared h-MoO<sub>3</sub> nanostructures. The specific capacitance ( $C_{\text{sp}}$ ) of the as-prepared electrodes can be calculated from charge/discharge curves using eqn (4):

$$C_{\text{sp}} = 2I \int \nu dt / A \Delta V^2 \quad (4)$$

where  $I$ ,  $A$ , and  $\Delta V$  are the current density at which charges/discharges are obtained, area of the electrode and voltage interval of the discharge, respectively. The change in slope of the discharge curves is attributable to the redox reaction, as can be seen in CVs as well (Fig. 6(a) and (b)).<sup>56</sup> The specific capacitance of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates is calculated to be  $230 \text{ F g}^{-1}$ ,  $160 \text{ F g}^{-1}$  and  $103 \text{ F g}^{-1}$ , respectively, at an applied current density of  $0.25 \text{ A g}^{-1}$ . Due to relatively high loading amount of  $\text{Mo}^{6+}$  ions as well as a larger electrochemically active surface area, an improvement in the specific capacitance of pyramidal nanorods is realized. The attainable charge-storage capacity of pyramidal nanorods is distinctly better than that of  $\alpha$ -MoO<sub>3</sub> nanorods, nanowires and



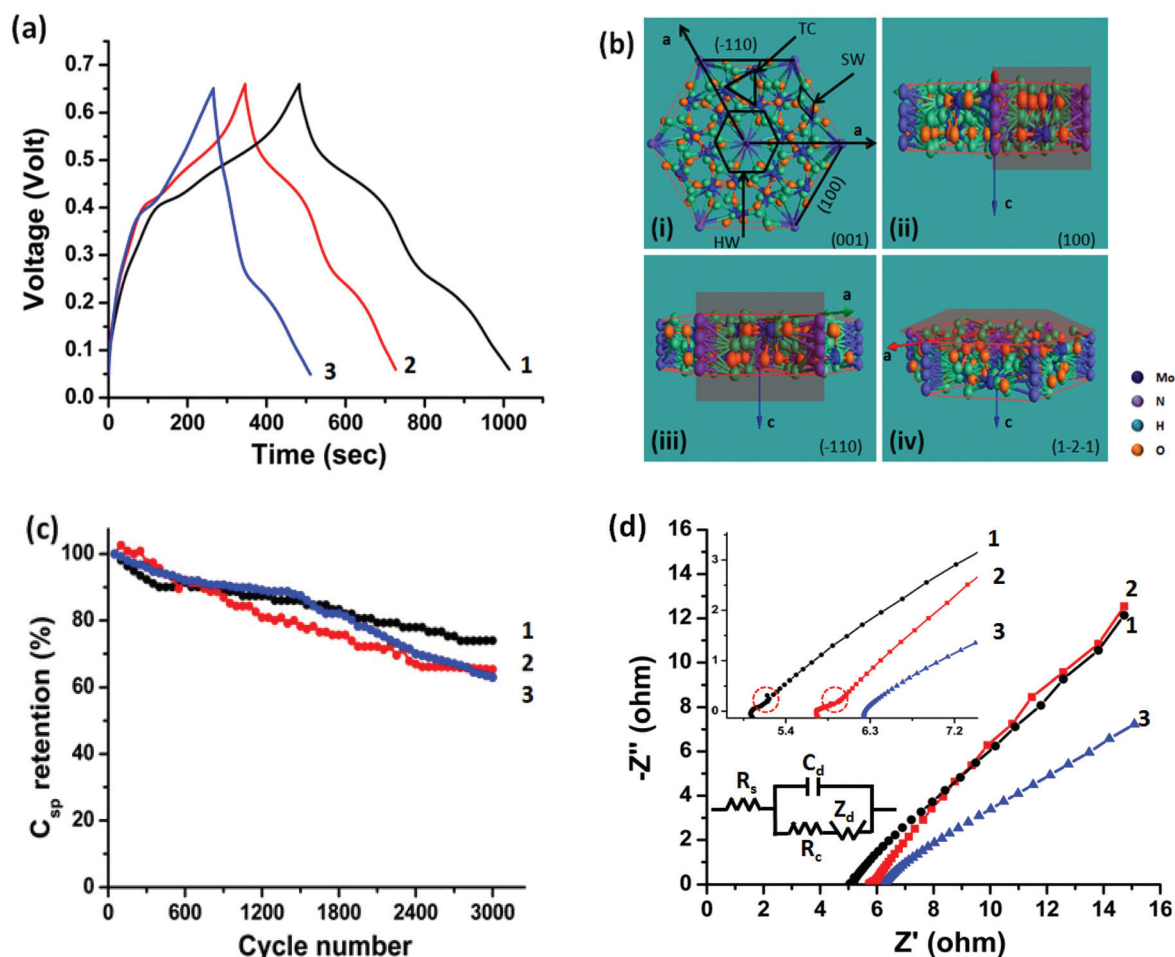


Fig. 7 Charge/discharge profiles of (1) pyramidal nanorods, (2) prismatic nanorods and (3) hexagonal nanoplates are performed at an applied current density of (a)  $0.25 \text{ A g}^{-1}$ , (b) atomic representation of hexagonal framework in the various plane configurations, (c) cycling test of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates using cyclic voltammetry (CV) at a scan rate of  $50 \text{ mV s}^{-1}$  and (d) Nyquist plots of pyramidal nanorods (1), prismatic nanorods (2) and hexagonal nanoplates (3) in a frequency bandwidth of 100 mHz to 100 kHz at open circuit potential (OCP). The inset shows the response of the as-prepared electrodes (pyramidal and prismatic nanorods) in the high frequency region (100 kHz–10 mHz). In the inset of figure (d) (Randles circuit),  $R_s$ ,  $Z_d$ ,  $R_c$  and  $C_{dl}$  indicates solution resistance, diffusion resistance or Warburg resistance, charge transfer resistance and double layer capacitance, respectively.

nanoribbons.<sup>57,58</sup> A comparison of the electrochemical performance of our electrode material with the  $\alpha\text{-MoO}_3$  is shown in the ESI, Table S3.†

We assert that the exposure of desirable crystallographic planes is of paramount importance. The hexagonal framework of MoO<sub>3</sub> resembles the hexagonal framework of WO<sub>3</sub>.<sup>59</sup> The amount of the stabilizing ions, *i.e.*, ammonium ions, occupancies in the hexagonal framework affects the available intercalation sites and therefore the electrochemical performance. The h-MoO<sub>3</sub> framework is constructed by MoO<sub>6</sub> octahedra sharing the equatorial oxygen in the *ab* plane and stacks along the *c* axis by sharing MoO<sub>6</sub> octahedra. The hexagonal framework contains three different intercalation sites, *i.e.*, hexagonal window (HW), trigonal cavity (TC) and four-coordinated square window (SW) running along the *c* axis, and these intercalation sites are available for the guest ions,<sup>60</sup> as shown in Fig. 7(b)-(i). Among the various available intercalation sites, most of the

HW sites were partially occupied by the NH<sub>4</sub><sup>+</sup> ions to stabilize the hexagonal framework, and exposed along the [001] direction of h-MoO<sub>3</sub>. On the other hand, TC and SW are positioned along the [001] and [100] directions, respectively. It is believed that the TC sites can accommodate a large number of electrolyte ions.<sup>60</sup> In our case, pyramidal nanorods are constructed by the (100) and ( $\bar{1}10$ ) planes (shadow plane in the atomic representation, Fig. 7(b)-(ii) and (iii)), representing their base and tip plane, respectively, while prismatic nanorods possess only (100) planes (shadow plane in the representation, Fig. 7(b)-(ii)). The coexistence of the ( $\bar{1}10$ ) plane and the (100) plane in the case of pyramidal nanorods offers an additional degree of freedom to electrolyte ions for intercalation into the TC and SW sites.<sup>60</sup> The existence of two distinct redox peaks in the CV analysis signifies the intercalation into the SW ( $0.1 \leq V \leq 0.3$ ) and TC ( $0.3 \leq V \leq 0.6$ ) sites; similar behavior was also identified in structurally identical Na doped



$V_2O_5$ .<sup>9</sup> The improved current with distinct redox peaks in the case of pyramidal nanorods can be considered a consequence of this process, as can be seen in Fig. 6(a). The persistence of redox peaks even at a higher scan rate ( $50 \text{ mV s}^{-1}$ ), Fig. 6(b), indicates the ease of intercalation into pyramidal nanorods due to the reasons discussed above. A pitiable capacitive charge-storage performance of hexagonal nanoplates originates from not only a higher content of ammonium ions within the framework but also the plane exposed by hexagonal nanoplates. The  $(1\bar{2}\bar{1})$  plane of hexagonal nanoplates predominantly exposes HW and TC sites, which are partially occupied by  $\text{NH}_4^+$  ions. The shadow plane in Fig. 7(b)-(iv) shows the  $(1\bar{2}\bar{1})$  plane of h-MoO<sub>3</sub>. The availability of SW sites in hexagonal nanoplates gives rise to the strong redox peaks (0.22 V/0.39 V) and is attributed to the lack of appropriate TC sites, a set of faint or depressed redox peaks appeared at 0.45 V/0.49 V, as can be seen in Fig. 6(a).

The cycling stability of h-MoO<sub>3</sub> with various morphologies (pyramidal nanorods, prismatic nanorods and hexagonal nanoplates) is evaluated by the cyclic voltammetry (CV) test, at a scan rate of  $50 \text{ mV s}^{-1}$  in 1 M H<sub>2</sub>SO<sub>4</sub> electrolyte solution, as shown in Fig. 7(c). A capacitance of 74%, 65% and 62% of the initial value is retained for pyramidal nanorods, prismatic nanorods and hexagonal nanoplates after 3000 cycles of continuous charge/discharge. This is attributable to the increase in the contact resistance that comes from poor contact (between the current collector and active materials) upon expansion/contraction of the active material (due to ingress/digress of electrolyte ions).<sup>54,56</sup> Despite achieving good cycling stability, an early stage ( $\sim 300$  cycles) degradation (12%) of capacitance is observed in pyramidal nanorods. To address the early stage degradation in the capacitance of pyramidal nanorods, an h-MoO<sub>3</sub> electrode without addition of carbon black and PVDF was subjected to 300 cycles. It is evident from the FESEM micrographs (ESI, Fig. S7(a) and (b)†) that the initial stage degradation is caused by breakage of pyramidal nanorods. The framework of metal oxides may collapse during ingress/digress of electrolyte ions.<sup>57</sup> Due to a relatively high intake of electrolyte ions in pyramidal nanorods, a small number of  $\text{Mo}^{5+}$  ions comes out from the framework (caused by the mobile  $\text{Mo}^{5+}$  ions). The excess negative charge of the system (produced by the removal of  $\text{Mo}^{5+}$  ions) is compensated by electrolyte ions ( $\text{H}^+$ ) and results in the framework being subjected to stress. Eventually, the framework of the nanorods breaks at the low energy site (neck of the nanorods) and isolates the tip and base of pyramidal nanorods. Due to formation of the abundant active sites at the tip as well as at the base of the nanorods (broken pyramidal nanorods), a long term cycling stability of the sample was realized, as shown in Fig. 7(c).

The benefit of h-MoO<sub>3</sub> with pyramidal shape becomes more apparent when Electrochemical Impedance Spectroscopy (EIS) is used to evaluate the charge transfer and electrode kinetics of h-MoO<sub>3</sub>. The Nyquist plot in Fig. 7(d) shows the charge transfer characteristics of pyramidal nanorods, prismatic nanorods and hexagonal nanoplates in the higher and lower

**Table 1** Calculated values of EIS elements from the Randles circuit

Sample	$R_s/\text{ohm}$	$R_c/\text{ohm}$	$Z_d$	$C_{dl}/\text{mF}$
1-Pyramidal nanorods	5.03	0.10	9.71	2.31
2-Prismatic nanorods	5.79	0.16	10.2	2.73
3-Hexagonal nanoplates	6.31	0.45	11.71	3.01

frequency regions. The semicircle in the high frequency region (as shown in the inset of Fig. 7(d)) signifies the charge transfer resistance associated with the faradic reactions. To determine the parameters ( $R_s$ ,  $C_{dl}$ ,  $R_c$  and  $Z_d$ ) of EIS, a Randles circuit diagram was used, as presented in the inset of Fig. 7(d). The values of the parameters are listed in Table 1. Pyramidal nanorods give a low intercept value ( $5.03 \Omega$ ) on the real axis ( $X$ -axis), which indicates its low internal resistance (solution resistance and uncompensated resistance) in comparison to prismatic nanorods ( $5.79 \Omega$ ) and hexagonal nanoplates ( $6.31 \Omega$ ). The slight kinks (dotted circle in the inset of Fig. 7(d)) in the high frequency region of the Nyquist plot (for pyramidal and prismatic nanorods) are likely due to the adsorption of electrolyte ions with the surface of the nanorods. Due to the adsorption of electrolyte ions with the surface of the nanorods, generation of a new resistive element in the series with the charge transfer resistance takes place. In the low frequency region, the sloping lines are related to the diffusion resistance (Warburg resistance), when the ions enter into the interior of the host material. However, in the very low frequency region ( $f < 10 \text{ Hz}$ ) the resistance of h-MoO<sub>3</sub> in pyramidal morphology increased slightly (due to sluggish redox reactions), as evaluated by Bode plot (shown in the ESI, Fig. S8†).

## 4. Conclusions

In summary, we have demonstrated a facile approach to control the desirable morphology of h-MoO<sub>3</sub> nanostructures. By rationally controlling the amount of HMTA, h-MoO<sub>3</sub> with pyramidal nanorod, prismatic nanorod and hexagonal nanoplate-like morphologies are achieved, with an approximate molecular formula unit of  $(\text{NH}_4)_{0.11}(\text{H}_2\text{O})_{0.28}\text{MoO}_{2.89}$ ,  $(\text{NH}_4)_{0.23}(\text{H}_2\text{O})_{0.39}\text{MoO}_{2.77}$  and  $(\text{NH}_4)_{0.44}(\text{H}_2\text{O})_{0.30}\text{MoO}_{2.56}$ , respectively. The as-synthesized h-MoO<sub>3</sub> in pyramidal morphology exhibited high specific capacitance ( $230 \text{ F g}^{-1}$ ) compared to prismatic nanorods ( $160 \text{ F g}^{-1}$ ) or hexagonal nanoplate morphology ( $103 \text{ F g}^{-1}$ ) at an applied current density of  $0.25 \text{ A g}^{-1}$ . The coexistence of the (100) and  $(\bar{1}10)$  planes with pyramidal morphology exposes various intercalation sites (HW, SW and TC) for the intercalation of electrolyte ions. The ammonium ions not only stabilize the hexagonal framework, but also facilitate the mechanism of charge-storage (by the Vehicle mechanism of proton transportation). This unique strategy to enhance the electro-kinetics leads to increase in the charge-storage capacity of MoO<sub>3</sub>. This represents the first attempt to prepare h-MoO<sub>3</sub> for electrochemical





capacitors, offering possibilities in catalytic, sensing and solar applications.

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## Notes and references

- 1 S. Li and D. A. Dixon, *J. Phys. Chem. A*, 2006, **110**, 6231–44.
- 2 F. Zhou, M. Cococcioni, C. Marianetti, D. Morgan and G. Ceder, *Phys. Rev. B: Condens. Matter*, 2004, **70**, 235121.
- 3 P. S. Directorate and F. Monmouth, *J. Electrochem. Soc.*, 1995, **142**, 6–8.
- 4 H. Y. Lee and J. B. Goodenough, *J. Solid State Chem.*, 1999, **144**, 220–223.
- 5 B. E. Conway, *Electrochemical Supercapacitors: Scientific Fundamentals and Technological Applications*, Kluwer Academic/Plenum, New York, 1999.
- 6 A. C. Arias, J. D. MacKenzie, I. McCulloch, J. Rivnay and A. Salleo, *Chem. Rev.*, 2010, **110**, 3–24.
- 7 M. S. Gordon, D. Dabdub and R. B. Gerber, *Proc. Natl. Acad. Sci. U. S. A.*, 2009, **106**, 16889–16889.
- 8 X. Wang, B. D. Myers, J. Yan, G. Shekhawat, V. Dravid and P. S. Lee, *Nanoscale*, 2013, **5**, 4119–22.
- 9 E. Khoo, J. Wang, J. Ma and P. S. Lee, *J. Mater. Chem.*, 2010, **20**, 8368.
- 10 J. P. Zheng, P. J. Cygan and T. R. Jow, *J. Electrochem. Soc.*, 1995, **142**, 2699.
- 11 J. P. Zheng and T. R. Jow, *J. Electrochem. Soc.*, 1995, **142**, L6.
- 12 H. Farsi, F. Gobal, H. Raissi and S. Moghiminia, *J. Solid State Electrochem.*, 2009, **14**, 643–650.
- 13 T. Brezesinski, J. Wang, S. H. Tolbert and B. Dunn, *Nat. Mater.*, 2010, **9**, 146–51.
- 14 J. W. Kim, V. Augustyn and B. Dunn, *Adv. Energy Mater.*, 2012, **2**, 141–148.
- 15 O. Ghodbane, J. L. Pascal and F. Favier, *ACS Appl. Mater. Interfaces*, 2009, **1**, 1130–1139.
- 16 R. L. Smith and G. S. Rohrer, *J. Solid State Chem.*, 1996, **115**, 104–115.
- 17 E. M. McCarron III, *J. Chem. Soc., Chem. Commun.*, 1986, 336–338.
- 18 J. C. Calabrese and E. M. McCarron, *J. Solid State Chem.*, 1991, **125**, 121–125.
- 19 V. Kumar, A. Sumboja, J. Wang, V. Bhavanasi, V. C. Nguyen and P. S. Lee, *Chem. Mater.*, 2014, **26**, 5533–5539.
- 20 T. Tao, Q. Chen, H. Hu and Y. Chen, *Mater. Lett.*, 2012, **66**, 102–105.
- 21 G. R. Li, Z. L. Wang, F. L. Zheng, Y. N. Ou and Y. X. Tong, *J. Mater. Chem.*, 2011, **21**, 4217.
- 22 V. Kumar and P. S. Lee, *J. Phys. Chem. C*, 2015, **119**, 9041–9049.
- 23 I. Shakir, M. Shahid, S. Cherevko, C.-H. Chung and D. J. Kang, *Electrochim. Acta*, 2011, **58**, 76–80.
- 24 I. Shakir, M. Shahid, M. Nadeem and D. J. Kang, *Electrochim. Acta*, 2012, **72**, 134–137.
- 25 M. S. Islam, D. J. Driscoll, C. A. J. Fisher, P. R. Slater, M. C. Group, C. Di V, V. Uni, G. Gu and U. Kingdom, *Chem. Mater.*, 2005, 5085–5092.
- 26 C. Nan, J. Lu, C. Chen, Q. Peng and Y. Li, *J. Mater. Chem.*, 2011, **21**, 9994.
- 27 G.-Z. Wei, X. Lu, F.-S. Ke, L. Huang, J.-T. Li, Z.-X. Wang, Z.-Y. Zhou and S.-G. Sun, *Adv. Mater.*, 2010, **22**, 4364–7.
- 28 J. Jiang, J. Liu, S. Peng, D. Qian, D. Luo, Q. Wang, Z. Tian and Y. Liu, *J. Mater. Chem. A*, 2013, **1**, 2588.
- 29 W. Tang, L. Liu, S. Tian, L. Li, Y. Yue, Y. Wu and K. Zhu, *Chem. Commun.*, 2011, **47**, 10058–60.
- 30 J. Chang, M. Jin, F. Yao, T. H. Kim, V. T. Le, H. Yue, F. Gunes, B. Li, A. Ghosh, S. Xie and Y. H. Lee, *Adv. Funct. Mater.*, 2013, **23**, 5074–5083.
- 31 L. V. Gurvich, I. V. Veyts and C. B. Alcock, *Thermodynamic Properties of Individual Substances*, CRC Press, Boca Raton, 4th edn, 1994.
- 32 V. Kumar, X. Wang and P. S. Lee, *CrystEngComm*, 2013, **15**, 7663.
- 33 H. J. Lunk, H. Hartl, M. A. Hartl, M. J. G. Fait, I. G. Shenderovich, M. Feist, T. A. Frisk, L. L. Daemen, D. Mauder, R. Eckelt and A. A. Gurinov, *Inorg. Chem.*, 2010, **49**, 9400–8.
- 34 B. Rachid, *Ann. Chim.*, 2007, **32**, 277–282.
- 35 K. M. McPeak, T. P. Le, N. G. Britton, Z. S. Nickolov, Y. A. Elabd and J. B. Baxter, *Langmuir*, 2011, **27**, 3672–7.
- 36 F. Taube, M. Hashimoto, I. Andersson and L. Pettersson, *J. Chem. Soc., Dalton Trans.*, 2002, 1002–1008.
- 37 V. Nardello, J. Marko, G. Vermeerch and J. M. Aubry, *Inorg. Chem.*, 1995, **34**, 4950–4957.
- 38 M. Cindric, Z. Veksli and B. Kamenar, *Croat. Chem. Acta*, 2009, **82**, 345–367.
- 39 I. M. Szilagyi, J. Madrasaz, G. Pokol, P. Kiraly, G. Tarkanyi, S. Saukko, J. Mizsei, A. L. Toth, A. Szabo and K. V. Josepovits, *Chem. Mater.*, 2008, **20**, 4116–4125.
- 40 Y. L. Leung, P. C. Wong, K. A. R. Mitchell and K. J. Smith, *Appl. Surf. Sci.*, 1998, **136**, 4–6.
- 41 R. Bacsá, J. Kiwi, T. Ohno, P. Albers and V. Nadtochenko, *J. Phys. Chem. B*, 2005, **109**, 5994–6003.
- 42 P. Serp and J. L. Figueiredo, *Carbon Materials For Catalysis*, John Wiley & Sons, Inc., Hoboken, New Jersey, USA, 2009.
- 43 H. Jiang, T. Zhao, C. Li and J. Ma, *J. Mater. Chem.*, 2011, **21**, 3818.
- 44 J. S. Chen, M. F. Ng, H. Bin Wu, L. Zhang and X. W. (David) Lou, *CrystEngComm*, 2012, **14**, 5133.
- 45 D. Cai, D. Wang, B. Liu, Y. Wang, Y. Liu, L. Wang, H. Li, H. Huang, Q. Li and T. Wang, *ACS Appl. Mater. Interfaces*, 2013, **5**, 12905–10.
- 46 X. Wang, W. S. Liu, X. Lu and P. S. Lee, *J. Mater. Chem.*, 2012, **22**, 23114.



- 47 J. Wang, J. Polleux, J. Lim and B. Dunn, *J. Phys. Chem. C*, 2007, **2**, 14925–14931.
- 48 S. M. Majhi, P. Rai, S. Raj, B. Chon, K. Park and Y. Yu, *ACS Appl. Mater. Interfaces*, 2014, **6**, 7491–7497.
- 49 K. D. Kreuer, A. Rabenau and W. Weppner, *Angew. Chem., Int. Ed. Engl.*, 1982, **21**, 3.
- 50 K. D. Kreuer, A. Rabenau and R. Messer, *Appl. Phys. A: Solids Surf.*, 1983, **32**, 45–53.
- 51 R. Thakkar and U. Chudasama, *J. Sci. Ind. Res.*, 2009, **68**, 312–318.
- 52 H. Lin, F. Zhou, C. P. Liu and V. Ozolins, *J. Mater. Chem. A*, 2014, **2**, 12280–12288.
- 53 S. S. Mahajan, S. H. Mujawar, P. S. Shinde, A. I. Inamdar and P. S. Patil, *Int. J. Electrochem. Sci.*, 2008, **3**, 953–960.
- 54 X. Wang, A. Sumboja, M. Lin, J. Yan and P. S. Lee, *Nanoscale*, 2012, **4**, 7266–72.
- 55 J. Chmiola, G. Yushin, R. Dash and Y. Gogotsi, *J. Power Sources*, 2006, **158**, 765.
- 56 Q. Mahmood, W. S. Kim and H. S. Park, *Nanoscale*, 2012, **4**, 7855–60.
- 57 I. Shakir, M. Shahid, H. W. Yang and D. J. Kang, *Electrochim. Acta*, 2010, **56**, 376–380.
- 58 R. Liang, H. Cao and D. Qian, *Chem. Commun.*, 2011, **47**, 10305–7.
- 59 M. Hibino, W. Han and T. Kudo, *Solid State Ionics*, 2000, **135**, 61–69.
- 60 S. Balaji, Y. Djaoued, R. Z. Ferguson and R. Bru, *Chem. Mater.*, 2009, **21**, 1381–1389.

