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# Single ion magnet behaviour in a two-dimensional network of dicyanamide-bridged cobalt(॥) ions $\dagger$ 

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#### Abstract

A novel two-dimensional coordination polymer of the formula $\left[\mathrm{Co}(\mathrm{dca})_{2}(\text { atz })_{2}\right]_{n}$ (1) resulted from assembling trans-bis(2-amino-1,3,5-triazine)cobalt(II) motifs by dicyanamide spacers. Variabletemperature dc and ac magnetic susceptibility measurements of 1 show that the high-spin cobalt(II) ions act as single ion magnets (SIMs).


In recent years, dicyanamide (dca)-bridged coordination compounds have attracted significant attention for their intriguing architectures and topologies as well as for a large variety of magnetic properties. ${ }^{1,2}$ The binary metal-dicyanamide complexes with a 3D-rutile type architecture $\left(\alpha-\left[M(d c a)_{2}\right]\right)$, formed by six-coordinate metal ions and three-connecting spacers ( $\mu_{1,3,5}$-dca ligands) in the 1:2 metal to dca molar ratio, exhibit a broad diversity long-range magnetic ordering at very low temperatures. ${ }^{3}$ This magnetic behaviour depends on the nature of the metal ion, being weakly ferromagnetic $[\mathrm{Co}(\mathrm{II})$, $\mathrm{Cu}($ II $)$ and $\mathrm{Ni}($ II $)]^{4}$ or spin-canted antiferromagnetic $[\mathrm{Cr}($ III $)$, $\mathrm{Mn}(\mathrm{II})$ and $\mathrm{Fe}(\mathrm{II})] .{ }^{5}$ The modification of the M-dca networks by the incorporation of N -donors as bridges or terminal coligands into the structures resulted in the formation of numerous architectures of the general formula $\left[\mathrm{M}(\mathrm{dca})_{2}(\mathrm{~L})_{n}\right](n=1$ or 2), including linear chains with a single or a double $\mu_{1,5}$-dca connectors, 2D $(4,4)$ nets and herringbone-like lattices with a single $\mu_{1,5^{-}}$dca bridge, triangular lattices with $\mu_{1,3,5^{-}}$and $\mu_{1,5^{-}}$ dca linkers, and also 3D networks. ${ }^{6-8}$ In contrast to the binary $\alpha-\left[\mathrm{M}(\mathrm{dca})_{2}\right]$ systems, the large number of magneto-structurally investigated heteroleptic compounds display very weak anti-

[^0]ferromagnetic couplings ( $J$ less than $1 \mathrm{~cm}^{-1}$ ) because of the poor ability of the extended $\mu_{1,5}$-dca bridges to mediate electronic interactions. ${ }^{1 a, 2 e, 5 c, 7 a}$

Having in mind the recent research results dealing with a few magneto-structural studies of mono- and polynuclear Co (II) compounds where the metal ions exhibit slow relaxation of the magnetization behaving as single ion magnets (SIMs), ${ }^{9-12}$ we have used the dca anion and the 2 -amino-1,3,5triazine (atz) molecule as a linker and an end-cap ligand, respectively. The resulting two-dimensional coordination polymer of the formula $\left[\mathrm{Co}(\mathrm{dca})_{2}(\mathrm{atz})_{2}\right]_{n}(1)$ has been characterized by spectroscopic techniques, X-ray diffraction and magnetic measurements.

Compound 1 was prepared $\ddagger$ by the reaction of an aqueous solution of $\mathrm{Na}(\mathrm{dca})$ with a methanolic solution of $\mathrm{CoCl}_{2} \cdot 2 \mathrm{H}_{2} \mathrm{O}$ and 2-amino-1,3,5-triazine (atz) in a $2: 1: 2$ dca:Co(II): atz molar ratio.$\ddagger$ Intense absorptions associated with the $\nu_{\text {as }}+$ $\nu_{\mathrm{s}}(\mathrm{C} \equiv \mathrm{N}), \nu_{\mathrm{as}}(\mathrm{C} \equiv \mathrm{N})$ and $\nu_{\mathrm{s}}(\mathrm{C} \equiv \mathrm{N})$, stretching modes of dca ligands occur at 2312, 2263 and $2194 \mathrm{~cm}^{-1}$, respectively. Their shift towards higher frequencies as compared to those of the sodium dicyanamide ( 2286,2232 and $2179 \mathrm{~cm}^{-1}$ ) supports the coordination of dca in $\mathbf{1}$. The XRPD pattern measured for the as-synthesized sample was in good agreement with the XRPD patterns simulated from the respective single-crystal X-ray data using the Mercury 2.4 program, ${ }^{13}$ demonstrating that the crystal structure is truly representative of the bulk material (ESI, Fig. S1 $\dagger$ ).

The single crystal X-ray analysis (see Table S1, ESI $\dagger$ ) revealed a polymeric structure consisting of neutral layers of the formula $\left[\mathrm{Co}(\mathrm{dca})_{2}(\mathrm{atz})_{2}\right]$ which are interlinked into a supramolecular 3D structure by hydrogen bonding between the coordinated atz ligands and the central nitrogen atoms of the dca bridges $\left[\mathrm{N}(4)-\mathrm{H}(4 \mathrm{~B}) \cdots \mathrm{N}(97)^{\mathrm{e}}\right.$ with values for $\mathrm{D} \cdots \mathrm{A}$ and D-H $\cdots \mathrm{A}$ of $3.083(3) \AA$ and $164.0^{\circ}$, respectively; symmetry code: (e) $=-1 / 2+x, 1 / 2+y,-1 / 2+z]$ (Fig. S2a $\dagger$ ) and $\pi-\pi$ stacking interactions with centroid-centroid distances of 3.573 and $3.744 \AA$ (Fig. S2b, ESI $\dagger$ ).

Each cobalt(II) ion, located at the crystallographic inversion Wyckoff 4 d center, is six-coordinate with four nitrile nitrogen
atoms at equatorial sites [Co-N distances of 2.0887(17) and $2.0968(16) \AA]$ and two nitrogen atoms of the atz ligand at the axial sites [Co-N distance of $2.1677(15) \AA$ A (see Table S2, ESI $\dagger$ ), building an elongated octahedron (Fig. 1a). The dca ligands adopt the $\mu_{1,5}$ bridging mode and they connect each metal ion to other four neighbouring metal centers leading to distorted rhombus-grid sheets parallel to the crystallographic $b c$ plane (Fig. 1b). The metal centers create 4 -c uninodal net described by a $\left\{4^{2} ; 6^{2}\right\}$ Schläfli symbol and a [4.4.4.4.6(2).6(2)] extended point vertex symbol, which corresponds to the sql topological type. ${ }^{14}$

Within the $[\mathrm{Co}(\mathrm{II})]_{4}$ units, the equatorial coordination planes of the neighboring metal centers are almost perpendicular, exhibiting a dihedral angle of $88.88^{\circ}$. The intralayer Co...Co separation through the dca bridge is $8.041 \AA$, whereas the metal-metal distances through the diagonals are different (13.423 and $8.859 \AA$ ), indicating that metal centers in the $[\mathrm{Co}(\mathrm{II})]_{4}$ units adopt a rhombus arrangement rather than a square disposition. Intralayer hydrogen bonds $[\mathrm{N}(4)-\mathrm{H}(4 \mathrm{~A})$ $\cdots \mathrm{N}(2)^{\mathrm{f}}$ with $\mathrm{D} \cdots \mathrm{A}=3.209(3) \AA$ and $\mathrm{D}-\mathrm{H} \cdots \mathrm{A}=174.0^{\circ}$; symmetry


Fig. 1 (a) Perspective view of the metal environment in 1 together with the atom numbering. Displacement ellipsoids are drawn at $50 \%$ probability level [symmetry code: (a) $=x,-y,-1 / 2+z$; (b) $=x,-y, 1 / 2+z$; (c) $=1 / 2-x,-1 / 2+y, 1 / 2-z$; (d) $=1 / 2-x,-1 / 2-y, 1-z]$. (b) View of a fragment of the 2D structure of 1 formed by $\mu_{1,5}$-dca bridges.
code: $(\mathrm{f})=1 / 2-x, 1 / 2-y, 1-z]$ contribute to the stabilization of the neutral layers. The shortest interlayer $\mathrm{Co} \cdots \mathrm{Co}^{g}$ [symmetry code: $(g)=1 / 2+x, 1 / 2+y, z]$ separation is $8.602 \AA$.

The magnetic properties of $\mathbf{1}$ in the form of a $\chi_{M} T$ versus $T$ plot $\left[\chi_{\mathrm{M}}\right.$ is the magnetic susceptibility per $\mathrm{Co}($ (I) ion $]$ are shown in Fig. 2. At room temperature $\chi_{\mathrm{M}} T$ is equal to $2.77 \mathrm{~cm}^{3} \mathrm{~mol}^{-1}$ $\mathrm{K}\left[\mu_{\text {eff }}=4.71 \mathrm{BM}\right.$ to be compared with $\mu_{\text {eff }}($ spin only $g=2.0)=$ 3.87 BM]. This value is within the range of those observed for six-coordinate high-spin cobalt(II) complexes with an unquenched angular momentum. ${ }^{15}$ Upon cooling, $\chi_{M} T$ remains practically constant in the high temperature range and it decreases sharply below 150 K to reach a value of $1.78 \mathrm{~cm}^{3} \mathrm{~mol}^{-1} \mathrm{~K}$ at 1.9 K . No maximum of the magnetic susceptibility is observed in the $\chi_{\mathrm{M}} v s$. $T$ plot. The decrease of $\chi_{\mathrm{M}} T$ in 1 can be due to an antiferromagnetic interaction between the cobalt(II) ions and/or to the thermal depopulation of the higher energy Kramers doublets of the cobalt(II) centers.

Having in mind the previous magneto-structural results concerning the dca-bridged cobalt(II) compounds (see Table S3, ESI $\dagger$ ) and the large intra [ca. $8.0 \AA$ across the single $\mu_{1,5^{-}}$-dca bridge] and interlayer [shortest value about $8.6 \AA$ i] cobalt-cobalt separation, the magnetic interaction between the local spin quartets, if any, is expected to be very weak. Consequently, the magnetic data of $\mathbf{1}$ were analysed through the Hamiltonian of eqn (1): ${ }^{16}$
$\boldsymbol{H}=-\alpha \lambda \boldsymbol{L}_{\mathrm{Co}} \boldsymbol{S}_{\mathrm{Co}}+\Delta\left[\boldsymbol{L}_{\mathrm{z}, \mathrm{Co}}^{2}-1 / 3 L(L+1)\right]+\beta H\left(-\alpha \boldsymbol{L}_{\mathrm{Co}}+g_{\mathrm{e}} \boldsymbol{S}_{\mathrm{Co}}\right)$
where $\lambda$ is the spin-orbit coupling and $\alpha$ is an orbital reduction factor defined as $\alpha=A \kappa$. The $\kappa$ parameter considers the reduction of the orbital momentum caused by the delocalization of the unpaired electrons and the $A$ parameter contains the admixture of the upper ${ }^{4} \mathrm{~T}_{1 g}\left({ }^{4} \mathrm{P}\right)$ state into the ${ }^{4} \mathrm{~T}_{1 g}\left({ }^{4} \mathrm{~F}\right)$ ground state ( $A=1.5$ and 1 in the weak and strong crystal-field limits, respectively). $\Delta$ is the energy gap between the singlet ${ }^{4} \mathrm{~A}_{2}$ and doublet ${ }^{4} \mathrm{E}$ levels arising from the splitting of the triplet orbital ${ }^{4} \mathrm{~T}_{1 g}$ ground state under an axial distortion of the


Fig. 2 Temperature dependence of the $\chi_{M} T$ product for 1: (o) experimental; (-) best-fit curve through eqn (1) (see text). The inset shows the magnetization vs. H plot 2.0 K .
ideal $O_{\mathrm{h}}$ symmetry of the cobalt(II) ion. The best-fit parameters using the experimental data in the whole temperature range investigated are: $\alpha=1.18(1), \lambda=-125(1) \mathrm{cm}^{-1}$ and $\Delta=$ $-509(10) \mathrm{cm}^{-1}$. Given that the values of $D q\left(955 \mathrm{~cm}^{-1}\right)$ and $B\left(890 \mathrm{~cm}^{-1}\right)$ could be obtained through the analysis of the UV-Vis-NIR electronic spectrum of $\mathbf{1}$ (see Fig. S4, ESI $\dagger$ ), a value of $A=1.41$ can be calculated through eqn (2) and (3): ${ }^{16}$

$$
\begin{equation*}
c=0.75+1.875 B / D q-1.25\left[1+1.8 B / D q+2.25(B / D q)^{2}\right]^{1 / 2} \tag{2}
\end{equation*}
$$

$$
\begin{equation*}
A=\left(3 / 2-c^{2}\right) /\left(1+c^{2}\right) \tag{3}
\end{equation*}
$$

and then $\kappa$ is found to be equal to 0.84 .
The value of the magnetisation vs. $H$ plot for 1 at 2.0 K tends to a quasi-saturation value of 2.08 BM at 6 T (inset of Fig. 2), a value which is as expected, $S_{\text {eff }}=1 / 2$ with $g_{\text {av }} c a .4 .2$. In this respect, it deserves to be noted that the maximum value of the saturation magnetisation in the $M$ against $H$ plots $\left(M_{\text {sat }}=g_{0} \beta M_{\mathrm{S}}\right)$ as a function of the temperature $T<30 \mathrm{~K}$ [a situation where the ground doublet for a six-coordinate high-spin cobalt(II) ion is the only populated state] would be 2.06 BM with $g_{0}=(10+2 \alpha) / 3=4.12$.

The ac susceptibility measurements for 1 under a 0 G static field show no out-of-phase ( $\chi_{\mathrm{M}}{ }^{\prime \prime}$ ) signals suggesting a fast tunnelling of the magnetization (QTM). The application of a dc field of 1000 G removes the QTM and the frequency dependence of $\chi_{\mathrm{M}}{ }^{\prime \prime}$ is observed (Fig. 3a). The relaxation times obtained from the maxima of $\chi_{\mathrm{M}}{ }^{\prime \prime}$ were fitted to an Arrhenius law (Fig. 3b), giving values for the energy barrier $\left(E_{\mathrm{a}}\right)$ and preexponential factor $\left(\tau_{0}\right)$ of $5.1 \mathrm{~cm}^{-1}$ and $1.7 \times 10^{-6} \mathrm{~s}$, respectively.

In conclusion, the X-ray structure of a novel 2D coordination polymer of the formula $\left[\mathrm{Co}(\mathrm{atz})_{2}(\mathrm{dca})_{2}\right]_{n}$ (1) with the cobalt(II) ions bridged by $\mu_{1,5}$-dca bridges has been reported. Combined dc and ac variable-temperature magnetic susceptibility measurements reveal the SIM behaviour of the magneti-


Fig. 3 (a) Frequency dependence of the out-of-phase ac susceptibilities from 2.0 to 9.0 K under a 1000 Oe dc field for 1. (b) The relaxation time of the magnetization $\ln (\tau)$ against $1 / T$ for 1 . The solid line represents the Arrhenius plot.
cally isolated six-coordinate high-spin cobalt(II) ions. Further work will be undertaken to increase the reduced number of magnetic coordination polymers of $\mathrm{Co}(\mathrm{II})$ behaving as SIMs by using the synthetic route described herein aiming to gain deeper insights into these interesting magnetic systems.

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## Notes and references

$\ddagger$ Synthesis: An aqueous solution $(5 \mathrm{~mL})$ of $\mathrm{NaN}(\mathrm{CN})_{2}(0.180 \mathrm{~g}, 2 \mathrm{mmol})$ was slowly poured into a methanolic solution ( 15 mL ) containing $\mathrm{CoCl}_{2} \cdot 2 \mathrm{H}_{2} \mathrm{O}$ ( $0.240 \mathrm{~g}, 1 \mathrm{mmol}$ ) and 3-amino-1,2,4-triazine ( $0.192 \mathrm{~g}, 2 \mathrm{mmol}$ ) under continuous stirring for several hours. The resulting pink solution was filtered to remove any small solid particle and allowed to evaporate in a hood at room temperature. X-ray quality pink prisms of $\mathbf{1}$ were grown from the mother liquor and collected after two weeks. Yield: $70 \%$ on Co. Elemental analysis (\%) for $\mathrm{C}_{10} \mathrm{H}_{8} \mathrm{~N}_{14} \mathrm{Co}$ (1): Calcd: C, 31.34; H, 2.10; N, 51.17; Found: C, 31.52; H. 2.19; N, 51.99. IR (KBr, $\left.\mathrm{cm}^{-1}\right): 3400(\mathrm{~s}), 3319(\mathrm{~s})$ and $3215(\mathrm{~s})[\nu(\mathrm{NH})] ; 2311(\mathrm{~s})\left[\nu_{\mathrm{as}}+\nu_{\mathrm{s}}(\mathrm{C} \equiv \mathrm{N})\right] ; 2262(\mathrm{~s})$ $\left[\nu_{\mathrm{as}}(\mathrm{C} \equiv \mathrm{N})\right]$ and 2194(vs) $\left[\nu_{\mathrm{s}}(\mathrm{C} \equiv \mathrm{N})\right], 1631(\mathrm{~s}), 1586(\mathrm{w}), 1560(\mathrm{~m}), 1542(\mathrm{w})$ and $1527(\mathrm{w})[\nu(\mathrm{C}=\mathrm{N})]$ and $[\nu(\mathrm{C}=\mathrm{C})]$. UV-Vis-NIR (solid, nm): 1020, 427, 355, 302 and 217.

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