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Quantum chemical insights into the dependence of porphyrin basicity on the *meso*-aryl substituents: thermodynamics, buckling, reaction sites and molecular flexibility

Martin Presselt,^{*abc} Wim Dehaen,^d Wouter Maes,^e Andreas Klamt,^{fg}
Todd Martinez,^{bh} Wichard J. D. Beenken^c and Mikalai Krukⁱ

The chemical and sensing properties of porphyrins are frequently tuned via the introduction of peripheral substituents. In the context of the exceptionally fast second protonation step in the case of 5,10,15,20-tetraphenylporphyrin (TPP), as compared to porphyrin and 5,10,15,20-tetramesitylporphyrin (TMesP), we investigated the macrocycle-substituent interactions of these three porphyrin derivatives in detail. Using quantum chemical thermodynamics calculations, the analysis of geometric structures, torsional profiles, electrostatic potential distributions, and particularly the analysis of molecular flexibilities via *ab initio* molecular dynamics simulations, we obtained a comprehensive picture of the reactivities of the studied porphyrins and how these are influenced by the *meso*-substituents. As compared to porphyrin and TMesP the second protonation of TPP is energetically more favorable and is particularly energetically comparable to its first protonation, instead of being significantly less favorable like in the case of porphyrin and TMesP. Additionally, the second TPP protonation is facilitated by an interplay between out-of-plane (oop) distortion of the protonation site and a pronounced electrostatic binding spot at the protonation site. Furthermore, the second protonation is particularly facilitated in the case of TPP by the large oop-flexibility of the diprotonated species as unraveled by *ab initio* molecular dynamics simulations.

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Introduction

The photophysical and chemical properties of porphyrins and related tetrapyrrolic macrocycles can be tuned over a wide range via peripheral functionalization, control of the acid–base equilibria of the porphyrin core, and the core size.^{1–20} Thus, a large variety of reports on tetrapyrrolic compounds can be found in the literature, like the occurrence of Fe-protoporphyrin IX as the prosthetic group of the hemoproteins in nature,²¹ the application as photosensitizers in photodynamic therapy,²² including the fast developing

field of two-photon excitation,^{23–25} as active materials for organic solar cells^{26,27} or light emitting diodes,²⁸ and in chemical sensors. Sensing applications can utilize the chelating core of the macrocycles or their acid–base equilibria, thus enabling *e.g.* optical pH-sensing with porphyrins.^{29–34}

To optimize porphyrin structures for the above-mentioned applications, particularly for sensing, control over the chemical properties of the porphyrin core, such as the acid–base equilibria shown in Scheme 1, is necessary. Therefore, possibilities to tune binding and reactivity properties of the macrocycle core via electronic and steric influences of the peripheral substituents have been extensively studied.^{35–42} Interestingly, the relative rates of the two possible subsequent protonations of the porphyrin core can be controlled by the *meso*-aryl substituents.^{43,44} In addition to the effect of protonation, the *meso*-substituents might cause nonplanar distortions of the macrocycle and change the molecular flexibility, as studied in-depth for diprotonated porphyrins.^{21,32,36,37,41,45} The type and extent of nonplanar distortions of the macrocycle have been extensively discussed in terms of molecular flexibility,^{32,37,41} the peripheral substitution pattern,^{21,32,36,37,41} and the strength of intermolecular interactions with acid residues or other anionic species in solution.^{10,36,43,46} The relationship between the molecular structure and the macrocycle acid–base equilibrium has also been

^a Institute of Physical Chemistry, Friedrich Schiller University Jena, Helmholtzweg 4, 07743 Jena, Germany. E-mail: martin.presselt@uni-jena.de

^b Department of Chemistry and PULSE Institute, Stanford University, Stanford, California 94305, USA

^c Ilmenau University of Technology, Institute for Physics, 98684 Ilmenau, Germany

^d Molecular Design and Synthesis, Department of Chemistry, KU Leuven, Celestijnenlaan 200F, 3001 Leuven, Belgium

^e Design & Synthesis of Organic Semiconductors (DSOS), Institute for Materials Research (IMO-IMOMEC), Hasselt University, 3590 Diepenbeek, Belgium

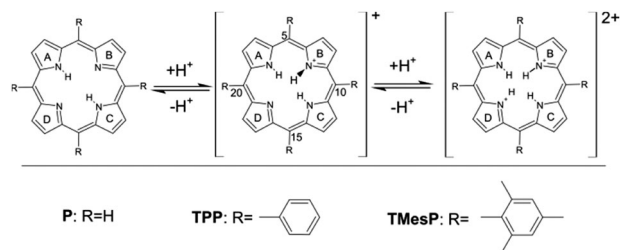
^f COSMOlogic GmbH&CO KG, Imbacher Weg 46, 51379 Leverkusen, Germany

^g University of Regensburg, 93040 Regensburg, Germany

^h SLAC National Accelerator Laboratory, Menlo Park, California 94309, USA

ⁱ Belarusian State Technological University, Physics Department, Sverdlova str. 13a, Minsk 220050, Belarus



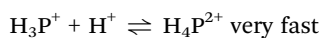
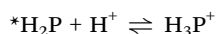
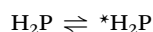


Scheme 1 Acid–base equilibria of different porphyrin derivatives. Carbon atoms 5, 10, 15, and 20 are the *meso*-carbons.

investigated.^{36,46–52} The monoprotonated forms of porphyrins are, however, much less studied with respect to molecular conformation, stability, and optical properties.^{47,51,53–55} For the porphine (H_2P) and alkyl-substituted porphyrins, both mono- and diprotonated forms were reported (see Fig. 1).^{53,55}

Particularly, comparing the evolution of the UV-vis absorption spectra in the course of acid–base titration of tetra-*meso*-methylporphyrin and tetra-*meso*-phenylporphyrin, TMesP and TPP, respectively (Fig. 1), demonstrates the dramatic influence of the type of *meso*-aryl substituents on the basicity of the porphyrin core. While the titration of TMesP yields three well distinguishable UV-vis absorption spectra in different pH ranges, that are attributed, respectively, to the free base (H_2TMesP), monoprotonated (H_3TMesP^+) and diprotonated ($\text{H}_4\text{TMesP}^{2+}$) species⁵² (as known from porphine and alkyl-substituted porphyrins), no spectral signature of the monoprotonated form (H_3TPP^+) was found for TPP and also not for the structurally very similar tetra-*meso*-(4-*N*-methylpyridyl)porphyrin ($\text{H}_2\text{TMPyP}^{4+}$) in the course of titration.^{43,44} To obtain stable monoprotonated derivatives of TPP, structural modifications are necessary, such as introducing hydrocarbon-capped moieties that face the macrocyclic plane⁵⁶ or complexation of the diprotonated form with bulky and poorly coordinating anions followed by dilution to decrease the acidity.⁵⁴

It was reported by Pasternack *et al.* as well as by Stone and Fleischer that the experimental titration data of TPP and (TMPyP)⁴⁺ are consistent with the following set of equilibria:^{43,44}



where ${}^*\text{H}_2\text{P}$ represents an activated form of the free base. The authors assumed that this activated form was buckled and that it dominates protonation, while similar mechanisms were not discussed for the mono-protonated form, as the second protonation is “very fast”. However, since we observed a significant difference in the second protonation between P, TMesP and TPP,⁵² the second protonation needs to be analyzed in more detail to understand how it is influenced by the *meso*-aryl substituents.

Therefore, we have first treated the thermodynamics of the acid–base equilibria of P, TMesP and TPP quantum chemically within the present paper. Since the base strength $\text{p}K_{\text{B}}$ depends linearly on the Gibbs free energy ΔG at a given temperature T

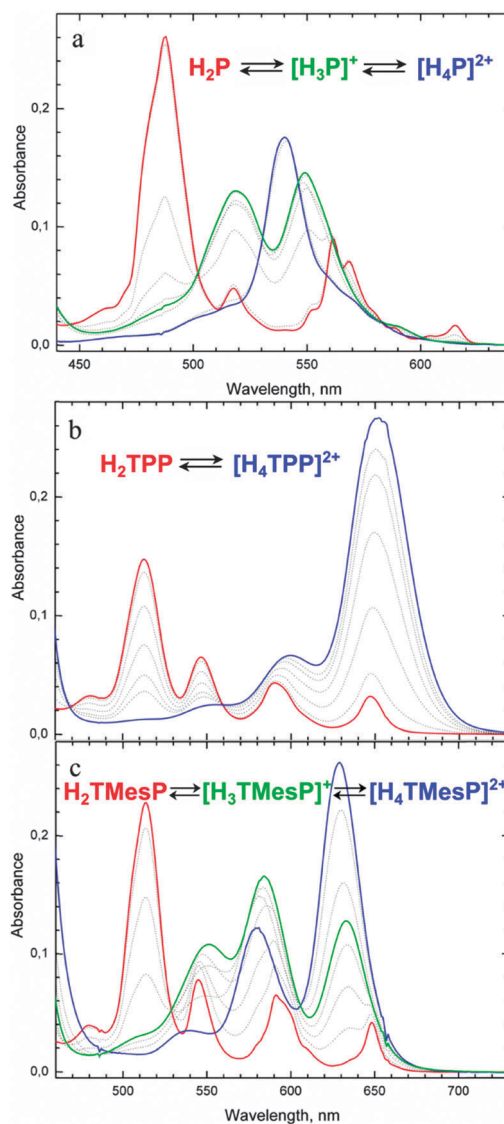


Fig. 1 Ground-state absorption spectra of the free base, mono- and diprotonated species (red, green and blue colors, respectively). Panels from top to bottom: porphine (H_2P , $[\text{H}_3\text{P}]^+$, $[\text{H}_4\text{P}]^{2+}$), tetra-*meso*-phenylporphyrin (H_2TPP , $[\text{H}_4\text{TPP}]^{2+}$), and tetra-*meso*-methylporphyrin (H_2TMesP , $[\text{H}_3\text{TMesP}]^+$, $[\text{H}_4\text{TMesP}]^{2+}$).

according to $\Delta G = -RT \ln K = -RT \ln 10 \lg K = -RT \ln 10 \lg(K_{\text{A}}/[\text{H}_2\text{O} = 55.5]) = RT \ln 10(\text{p}K_{\text{A}} + \lg 55.5)$, thus $\text{p}K_{\text{A}} = \Delta G/(RT \ln 10) - 1.74$ (R is the gas constant and K the equilibrium constant), ΔG needs to be calculated carefully under consideration of the solvent environment.⁵⁷ We compared the energies of isolated molecules and those enclosed by a dielectric continuum solvent environment (COSMO) in the conductor limit and additionally involved the interaction of surface elements to better account for intermolecular interactions, *i.e.* we applied a statistical thermodynamics treatment for more realistic solvation (RS) simulations.^{58–60}

If buckling or any kind of out-of-plane distortion determines the kinetics of the first protonation,^{43,44} it might also be important for the second protonation. Therefore, we analyzed the pyrrole out-of-plane distortions along with the *meso*-aryl tilts. The buckling

might accelerate protonation, either because it increases structural similarity between the mono- and diprotonated forms, or because it increases accessibility of the protonable nitrogen. The latter was analyzed by means of reactivity volumes. In any case, the *meso*-aryl substituents must play an essential role, possibly in terms of the porphyrin flexibility, as studied by means of *ab initio* molecular dynamics and analysis of the torsional profiles of the *meso*-substituents.

Methods

Quantum chemical structure optimizations, polarizable continuum solvation, and calculations of electrostatic potential distributions were performed using density functional theory (DFT) implemented in Turbomole^{61,62} and applying the GGA (generalized gradient approximation) B-P86 exchange–correlation functional, the def2-TZVP triple- ζ basis set^{63,64} and the MARI-J approximation.⁶⁵ This combination has been shown to give reliable geometries, electron density distributions and spectroscopic properties in many cases at a very reasonable computational cost.^{66–72} For simulation of solvent environments, the conductor like screening model (COSMO, performed in the conductor limit, *i.e.* $\epsilon = \infty$) and additionally COSMO for realistic solvation (COSMO-RS) implemented in COSMOtherm were used to model the dissolution of porphyrins in water and tetrahydrofuran (THF) that was used in the experiments.^{57–59} *Ab initio* molecular dynamics (MD) were run at 298 K to determine the range of variation of internal molecular coordinates as estimates for molecular flexibility. Therefore the highly efficient program TeraChem was used.^{73,74} For analyzing electrostatic potential distributions and MD trajectories, Mathematica⁷⁵ Version 8 was used.

Results and discussion

The experimentally observed differences in spectroscopic responses upon acid addition between porphine and both differently *meso*-substituted porphyrins might originate from thermodynamics and/or kinetics. First, we compared DFT-calculated energies of the porphyrin derivatives involved in the acid–base equilibria with the experimental basicity relations. In the next step, we analyzed the steady state geometries of the porphyrin derivatives involved in the acid–base equilibria and estimated molecular flexibility from the distribution of internal coordinate values as sampled by MD runs. Thus, we will ascertain if the exceptionally fast second protonation in

the case of tetra-*meso*-phenylporphyrin is due to a high structural similarity between the mono- and diprotonated species. Additionally, a high molecular out-of-plane (oop) flexibility would promote the formation of stable saddle-type geometries of diprotonated species and would facilitate exceptional oop-exposure of the reaction site in $[\text{H}_3\text{TPP}]^+$.

Protonation energies and $\text{p}K_{\text{B}}$ values

The Gibbs free energy difference of dissociation, ΔG_{diss} , that drives the protonation in the acid–base equilibria $[\text{H}_{2+n}\text{P}]^{n+} + [\text{H}_3\text{O}]^+ \rightleftharpoons [\text{H}_{3+n}\text{P}]^{(n+1)+} + \text{H}_2\text{O}$, exemplarily shown for porphyrin (P) with $n = 0$ or 1 for the first or second protonation, is defined according to eqn (1), with B being the base.

$$\Delta G_{\text{diss}} = G(\text{B}) - G([\text{HB}]^+) + [G([\text{H}_3\text{O}]^+) - G(\text{H}_2\text{O})] \quad (1)$$

The differences in ΔG_{diss} between the porphyrin derivatives P, TMesP, and TPP dissolved in the same solvent equal the differences in corresponding ΔG values defined as $\Delta G = G(\text{B}) - G([\text{HB}]^+)$. To better understand the origin of these ΔG differences between the protonation of P, TMesP and TPP, we discuss the independent contributions from the isolated molecules E_{QC}^* , the dielectric energies, E_{diel} , of protonated and non-protonated species according to COSMO at $\epsilon = \infty$, *i.e.* the energy contribution that is caused by the dielectric screening of charges at a solvent-excluded surface (SAS) in the conductor limit, and the chemical potentials obtained by COSMO-RS, μ .⁵⁸

The pseudo gas-phase energy difference, ΔE_{QC}^* , is defined as $\Delta E_{\text{QC}}^* = (E_{\text{COSMO}}^{\text{B}} - E_{\text{diel}}^{\text{B}}) - (E_{\text{COSMO}}^{[\text{BH}]^+} - E_{\text{diel}}^{[\text{BH}]^+})$.⁵⁸ As shown in Table 1, ΔE_{QC}^* for the first protonation step is similar for all the three derivatives H_2P , H_2TMesP , and H_2TPP ($\Delta E_{\text{QC}}^* = -1034, -1070, -1073 \text{ kJ mol}^{-1}$; differences in this series: 36 and 3 kJ mol^{-1}), with a particular small difference between the tetra-*meso*-aryl-substituted porphyrins. The ΔE_{QC}^* differences for the second protonation step between $[\text{H}_3\text{P}]^+$, $[\text{H}_3\text{TMesP}]^+$, and $[\text{H}_3\text{TPP}]^+$ are larger ($\Delta E_{\text{QC}}^* = -680, -763, -771 \text{ kJ mol}^{-1}$; differences in this series: 83 and 8 kJ mol^{-1}) but their relations are similar. Consequently, in the gas phase, *meso*-aryl-substitution leads to a significantly increased stabilization of the protonated species (first/second protonation: $\Delta E_{\text{QC}}^*(\text{TMesP}) - \Delta E_{\text{QC}}^*(\text{P}) = 36/83 \text{ kJ mol}^{-1}$, $\Delta E_{\text{QC}}^*(\text{TPP}) - \Delta E_{\text{QC}}^*(\text{P}) = 39/91 \text{ kJ mol}^{-1}$), but the second protonation step for TPP is energetically only slightly

Table 1 Gas-phase, conductor, COSMO and COSMO-RS energies, chemical potentials and Gibbs free energies for the free base and protonated states of porphine and the tetra-*meso*-aryl-substituted derivatives

	Pseudo-gas	COSMO		COSMO-RS H ₂ O		
	ΔE_{QC}^*	ΔE_{diel}	ΔE_{COSMO}	$\Delta \mu$	$\Delta \Delta G$	p <i>K</i> _B
	[kJ mol ^{−1}]	[kJ mol ^{−1}]		[kJ mol ^{−1}]		
H ₂ P → [H ₃ P] ⁺	−1034	−127.9	−1162	9	−1153	11
[H ₃ P] ⁺ → [H ₄ P] ²⁺	−680	−453.9	−1134	4	−1130	14
H ₂ TMesP → [H ₃ TMesP] ⁺	−1070	−96.9	−1167	12	−1155	11
[H ₃ TMesP] ⁺ → [H ₄ TMesP] ²⁺	−763	−381.9	−1145	5	−1140	13
H ₂ TPP → [H ₃ TPP] ⁺	−1073	−104.2	−1177	13	−1164	10
[H ₃ TPP] ⁺ → [H ₄ TPP] ²⁺	−771	−395.3	−1166	5	−1162	10



favorable as compared to TMesP and the second proton affinities are generally smaller than the first ones.

The energy gains due to dielectric screening E_{diel} increase approximately linearly with the molecular size of the free bases. However, in the case of the protonated species, E_{diel} of the pristine porphyrin significantly exceeds the E_{diel} stabilization of the tetra-*meso*-aryl-substituted derivatives, while the E_{diel} stabilization in the case of TPP is slightly larger than in the case of TMesP. Thus, ΔE_{COSMO} for the first protonation is similar for P and TMesP (-1162 and -1167 kJ mol $^{-1}$), while it is largest for TPP (-1177 kJ mol $^{-1}$) and in the case of the second protonation ΔE_{COSMO} increases from P to TMesP to TPP (-1134 , -1145 and -1166 kJ mol $^{-1}$), as shown in Table 1. Consequently, the highest basicity for TPP as compared to P and TMesP observed experimentally is reproduced by the COSMO-energies for protonation, but the ΔE_{COSMO} values are again smaller for the second protonation as compared to the first one. However, this $\Delta \Delta E_{\text{COSMO}}$ difference is the smallest, *i.e.* the energy gains due to the first and second protonation are almost similar, for TPP ($\Delta \Delta E_{\text{COSMO}}$ between the second and first protonation for P, TMesP, TPP: -28 , -22 , -10 kJ mol $^{-1}$).

Additional accounting of interactions of molecular surface elements in real solutions using COSMO-RS slightly alters the protonation energies, which now even better match the experimental findings. The ΔE_{COSMO} relations are approximately preserved, but particularly ΔG values for the first protonation are reduced due to the consideration of the chemical potentials of the solutes in water determined *via* COSMO-RS. Thus, the COSMO-RS derived ΔG values for the first and second protonation of TPP are almost identical (-1164 and -1162 kJ mol $^{-1}$). The $\Delta \Delta G$ value between the first and second protonation of TPP equates to -2 kJ mol $^{-1}$, while $\Delta \Delta G$ -values for P and TMesP are both significantly higher, -23 and -15 kJ mol $^{-1}$, respectively. The same relation is mirrored by the corresponding pK_{B} values reported in Table 1.

As detailed in Table 2, the mean interaction energies decrease with subsequent protonation, which is mainly driven by hydrogen-bond-interactions that are most energetically beneficial for the

diprotated species. Opposingly, the surface charges are best compensated in the case of the free-bases (misfit energy E_{mf}), while van der Waals interactions E_{vdw} are energetically most favorable for the mono-protonated species. However, both contributions to the mean interaction energy, as well as their changes, are small as compared to hydrogen bonds. The energetic benefit upon improved hydrogen bonding of the diprotated species is similar for TMesP and TPP, thus their change in the mean interaction energy at the second protonation is virtually identical. Just small differences in the mean interaction energy between TMesP and TPP at the first protonation contribute to the larger similarity between the Gibbs free energies of the first and the second protonation of TPP as compared to TMesP.

The energetic contributions due to hydrogen bonds with tetrahydrofuran (THF) are smaller for the free bases but larger for the diprotated species as compared to water as shown in Table 3. Their changes upon protonation are slightly larger in the case of THF as solvent as compared to water. van der Waals interactions of the solutes with THF are generally larger than in the case of water, hence causing similar energetic changes upon protonation. The compensation of surface-charges depends more strongly on the protonation state and its energetic change upon protonation is generally more unfavorable than in the case of water. However, the relations between the changes in the mean interaction energies upon protonation are very similar to the ones discussed above for water. Finally, THF yields porphyrins with higher basicity when dissolved in water, but the pK_{B} -difference between the first and second protonation is again smaller for TPP ($\Delta pK_{\text{B}} = 1.3$) as compared to TMesP ($\Delta pK_{\text{B}} = 3.2$).

In conclusion, the experimentally observed thermodynamic relations between the successive protonation of P, TMesP and TPP could be reproduced by using COSMO-RS and are predominantly determined by the interplay between gas-phase and dielectric screening energies. These gas-phase energies themselves depend on the interplay between steric interactions and molecular π -conjugation, while the dielectric screening energies are

Table 2 Solute–water interactions as calculated *via* COSMO-RS. E_{i} is the mean interaction energy, E_{mf} , E_{H} , and E_{vdw} refer to the misfit, hydrogen bonding, and van der Waals energy contributions, respectively. All energies and potentials are given in kJ mol $^{-1}$

	ΔE_{i}	ΔE_{mf}	ΔE_{H}	ΔE_{vdw}	$\Delta \mu$	ΔG	pK_{B}
$\text{H}_2\text{P} \rightarrow [\text{H}_3\text{P}]^+$	−6.4	1.1	−7.0	−0.5	8.7	−1153.4	11.0
$[\text{H}_3\text{P}]^+ \rightarrow [\text{H}_4\text{P}]^{2+}$	−29.5	0.1	−30.7	1.2	3.6	−1130.1	14.3
$\text{H}_2\text{TMesP} \rightarrow [\text{H}_3\text{TMesP}]^+$	−3.6	1.8	−4.9	−0.6	12.2	−1154.9	10.8
$[\text{H}_3\text{TMesP}]^+ \rightarrow [\text{H}_4\text{TMesP}]^{2+}$	−22.2	0.3	−24.6	2.1	4.8	−1139.8	12.9
$\text{H}_2\text{TPP} \rightarrow [\text{H}_3\text{TPP}]^+$	−5.1	1.9	−7.2	0.2	13.0	−1163.9	9.5
$[\text{H}_3\text{TPP}]^+ \rightarrow [\text{H}_4\text{TPP}]^{2+}$	−22.2	−0.2	−24.0	2.0	4.8	−1161.7	9.8

Table 3 Solute–tetrahydrofuran (THF) interactions as calculated *via* COSMO-RS. E_{i} is the mean interaction energy, E_{mf} , E_{H} , and E_{vdw} refer to the misfit, hydrogen bonding, and van der Waals energy contributions, respectively. All energies and potentials are given in kJ mol $^{-1}$

	ΔE_{i}	ΔE_{mf}	ΔE_{H}	ΔE_{vdw}	$\Delta \mu$	ΔG	pK_{B}
$\text{H}_2\text{P} \rightarrow [\text{H}_3\text{P}]^+$	−9.3	5.4	−14.6	−0.1	−2.7	−1164.8	11.0
$[\text{H}_3\text{P}]^+ \rightarrow [\text{H}_4\text{P}]^{2+}$	−26.9	6.3	−35.1	1.9	−3.8	−1137.5	12.5
$\text{H}_2\text{TMesP} \rightarrow [\text{H}_3\text{TMesP}]^+$	−5.8	6.7	−12.1	−0.4	0.1	−1166.9	8.1
$[\text{H}_3\text{TMesP}]^+ \rightarrow [\text{H}_4\text{TMesP}]^{2+}$	−17.2	7.8	−27.4	2.4	−0.4	−1145.0	11.4
$\text{H}_2\text{TPP} \rightarrow [\text{H}_3\text{TPP}]^+$	−6.7	6.7	−13.6	0.2	1.1	−1175.8	6.8
$[\text{H}_3\text{TPP}]^+ \rightarrow [\text{H}_4\text{TPP}]^{2+}$	−17.2	7.1	−27.0	2.7	−0.8	−1167.3	8.1



determined by the charge distribution at the molecular surface, hence depending also on the molecular geometry and electron delocalization. To identify the molecular origin of the outstanding protonation thermodynamics and kinetics of TPP, we subsequently focused on the geometric and electronic properties of the tetra-*meso*-aryl-substituted porphyrins.

Structural features and geometric similarity between different protonation states

One contribution to the fast second protonation of TPP might be a larger structural similarity between its mono- and diprotonated form as compared to P and TMesP, what is discussed after presentation of the basic structural features of each investigated derivative.

The oop-exposures of all pyrroles are quantified by the dihedral angles defined by α -, *meso*-, α - and either β -carbons or nitrogen as the fourth point (θ^{BC} or θ^{N}). As shown by these angles, given in Table 4, H₂P and H₂TMesP are planar, while H₂TPP is slightly distorted ($\theta < 1^\circ$). In [H₃P]⁺ and [H₃TMesP]⁺ the protonated pyrroles B are tilted out of plane to a virtually identical extent ($\theta^{\text{N}} \approx -9^\circ$). Pyrroles A and C (for assignment see Scheme 1) are both tilted oop less and in the opposite direction than pyrrole B ([H₃P]⁺: $\theta^{\text{N}} = 3^\circ$, [H₃TMesP]⁺: $\theta^{\text{N}} = 5^\circ$). The protonated pyrrole B is bent slightly more ($\Delta\theta = |\theta^{\text{BC}}| - |\theta^{\text{N}}| = 3$ to 4°) in all mono-protonated porphyrins than all other pyrroles A, C and D ($\Delta\theta < 2^\circ$). In both mono-protonated derivatives [H₃P]⁺ and [H₃TMesP]⁺, the non-protonated pyrrole D is almost in plane, but with a very small oop-distortion towards a saddle-type geometry ([H₃P]⁺: $\theta^{\text{N}} = 0^\circ$, [H₃TMesP]⁺: $\theta^{\text{N}} = -2^\circ$). In contrast to [H₃P]⁺ and [H₃TMesP]⁺, the non-protonated pyrrole D in [H₃TPP]⁺ is tilted oop considerably ($\theta^{\text{N}} = -8^\circ$), thus possessing a oop-distorted structure similar to the saddle-type geometry of the diprotonated

pristine porphyrin [H₄P]²⁺. For all diprotonated derivatives the saddle-type geometry is energetically most favourable, see Table 4 ($|\theta|$ for [H₄P]²⁺, [H₄TMesP]²⁺, [H₄TPP]²⁺: 10, 15, 20°).

The change in the macrocycle oop-geometry upon protonation is concisely reflected in the root-mean-square-deviations (RMSD) of the above mentioned dihedral angles from the planar geometry. These RMSDs clearly show that the total macrocycle oop-deformations increase from P to TMesP to TPP for each protonation step, respectively. The changes in these RMSDs upon the first protonation are small for P and TMesP (5.5° and 6.5°), but significantly larger for TPP (11.5°). The second protonation causes further oop-deformation, but the oop-geometry-change ΔRMSD is actually smaller for P and TPP (3.8° and 8.3°), but slightly larger for TMesP (8.2°) than in the case of the first protonation. Thus, TPP undergoes an exceptional structural rearrangement during the first protonation, which is significantly smaller for the second protonation. This relation might contribute to the fact that the second protonation is significantly faster than the first one in the case of TPP, hence leading to a domination of H₂TPP and [H₄TPP]²⁺ related absorption features in the course of the titration experiment. However, since the ΔRMSD -values for the second protonation are almost identical for TPP and TMesP, the macrocycle geometry of mono-protonated TPP is not closer to its diprotonated form than in the case of TMesP.

Macrocycle flexibility

After discussing the similarities between the energetically most favorable geometries of the mono- and diprotonated species and the accessibilities of the non-protonated nitrogen atoms, the conclusions drawn were reappraised in view of molecular flexibilities as derived from *ab initio* molecular dynamics (AIMD) simulations.

Table 4 Dihedral angles θ (in degrees) within the porphyrin macrocycle (pyrrole tilt) and between the tetra-*meso*-aryl substituents and the porphyrin macrocycle (aryl twist). The values of the tilting angle for nonprotonated pyrrole D and twisting angles of adjacent aryls in monoprotonated species are given in bold

		H ₂ P	H ₃ P ⁺	H ₄ P ²⁺	H ₂ TMesP	H ₃ TMesP ⁺	H ₄ TMesP ²⁺	H ₂ TPP	H ₃ TPP ⁺	H ₄ TPP ²⁺
Pyrrole A ^a	C _{2,3}	0.0	−1.9	−9.6	0.0	−3.6	−14.2	0.0	−9.5	−19.5
	N ₂₁	0.0	2.7	9.6	0.0	4.6	14.8	0.4	10.3	20.5
Pyrrole B ^a	C _{7,8}	0.0	12.2	9.5	0.0	12.7	14.5	0.7	17.9	19.0
	N ₂₂	0.0	−8.5	−9.2	0.0	−9.7	−15.0	−0.2	−15.2	−20.1
Pyrrole C ^a	C _{12,13}	0.0	−2.1	−9.3	0.0	−4.1	−14.4	0.5	−9.6	−19.4
	N ₂₃	0.0	2.9	9.5	0.0	5.1	15.0	0.3	10.5	20.1
Pyrrole D ^a	C _{17,18}	0.0	0.8	9.3	0.0	1.3	14.5	0.6	5.5	19.1
	N ₂₄	0.0	− 0.3	−9.0	0.0	− 1.6	−15.1	−0.7	− 8.3	−20.1
RMSD ^b		0.0	5.5	9.4	0.0	6.5	14.7	0.5	11.5	19.7
ΔRMSD^c			5.5	3.8		6.5	8.2		11.0	8.3
5-Aryl ^d					−90.0	−82.4	−71.4	−71.2	−55.6	−45.6
10-Aryl ^d					90.0	81.8	70.4	70.5	55.1	45.6
15-Aryl ^d					−90.0	− 88.5	−71.8	−70.8	− 60.1	−45.6
20-Aryl ^d					90.0	88.1	71.2	69.7	60.3	45.2

^a Numbering is according to Scheme 1. The pyrrole tilts are measured by dihedral angles separately for nitrogen (N_{21–24}) and C _{β} carbon atoms (C_{2,3,7,8,12,13,17,18}) on the basis of the bonds between the respective C _{α} (C_{1,4,6,9,11,14,16,19}) and C_{*m*} carbons (C_{5,10,15,20}) averaged for the left and right side of each pyrrole. Pyrroles A and C are protonated in the free bases (H₂P, H₂TPP, and H₂TMesP). Pyrrole B is additionally protonated in the monoprotonated forms (H₃P⁺, H₃TPP⁺, and H₃TMesP⁺) and all four pyrroles are protonated in the diprotonated forms (H₄P²⁺, H₄TPP²⁺, and H₄TMesP²⁺). ^b Root mean square deviation from planar, i.e. $\theta = 0^\circ$, geometries: $\text{RMSD} = \sqrt{(1/n) \sum \theta^2}$. ^c RMSD difference between two differently protonated species. ^d The aryl substituent is labelled by the *meso*-position (C_{5,10,15,20}) to which it is attached. The basis for these dihedral angles is again the bond between the respective carbon atoms in *alpha*- (C_{1,4,6,9,11,14,16,19}) and *meso*-positions (C_{5,10,15,20}) and the C₁- and C₂-carbon atoms in the aryl ring. Thereby, for counterclockwise order the sign of the angle has been taken as negative.



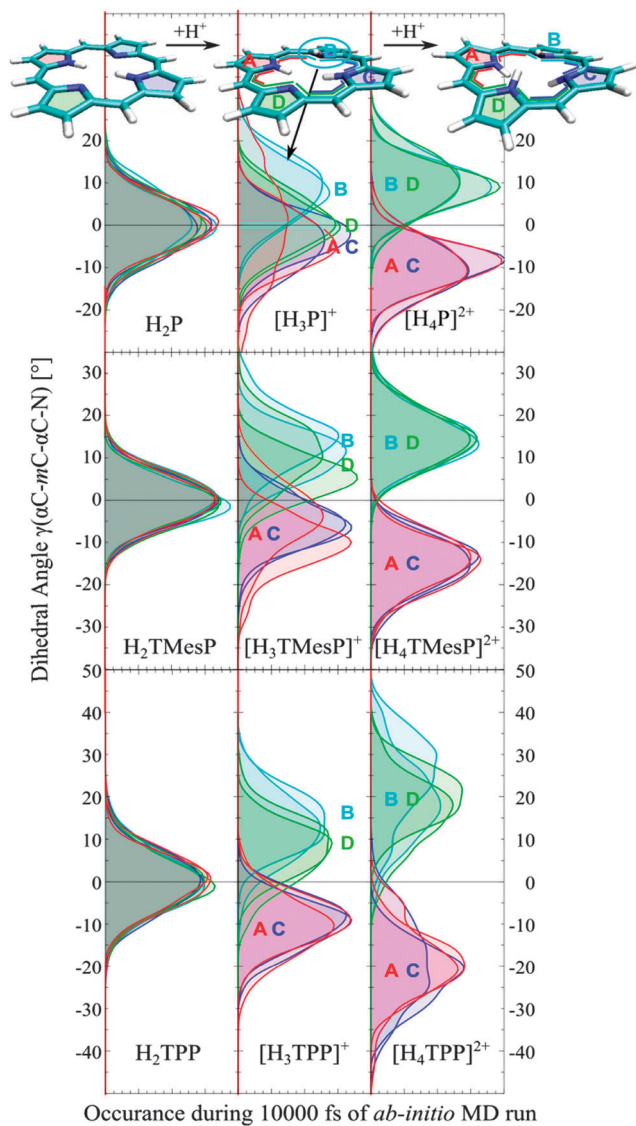


Fig. 2 Histograms of dihedral angles θ^N , which quantify the oop-distortion and are determined by the flexibility of the macrocycle. For each pyrrole two dihedral angles involving the same nitrogen, but α -, *meso*-, and next α -carbons in opposing directions can be defined, as indicated in the protonated structures at the top of the figure. Thus, two identically colored lines represent the smoothed histograms of these two dihedral angles.

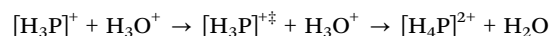
To analyze the geometries sampled during the AIMD runs, histograms of θ^N angles for each geometry are plotted in Fig. 2.

In the case of all investigated free base porphyrins, the histograms are almost fully superposed and show maxima at 0° and a full-width-at-half maximum of roughly 16° , *i.e.* predominantly planar macrocycles with similar flexibilities are present for all investigated derivatives. Upon subsequent protonation, saddle-type structures are formed, while in the case of the mono-protonated porphyrin $[H_3P]^+$, the non-protonated pyrrole D stays planar. In contrast, in $[H_3TMesP]^+$ and $[H_3TPP]^+$ the *meso*-aryl groups transmit the tilts of pyrroles A and C to D so that the non-protonated pyrrole D is tilted oop by about 10° . Significant differences between the oop-tilt of the non-protonated pyrrole D between $[H_3TMesP]^+$ and $[H_3TPP]^+$ are not present. In contrast,

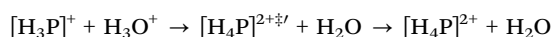
the histograms of $[H_4TMesP]^{2+}$ and $[H_4TPP]^{2+}$ are vastly different. While the histograms of $[H_4TMesP]^{2+}$ are smooth Gaussian profiles with virtually full superposition for B, D and A, C and maxima at *ca.* $\pm 15^\circ$, respectively, the histograms of $[H_4TPP]^{2+}$ are significantly broader with larger contributions at small θ^N tilt angles. Thus, we conclude that $[H_4TPP]^{2+}$ has a significantly larger oop-flexibility than $[H_4TMesP]^{2+}$ and can more easily adopt to the geometry of its mono-protonated form than the mesityl-substituted derivative. In other words, the weaker restriction in aryl-torsions in the case of TPP as compared to TMesP causes an exceptionally high oop-flexibility of diprotonated TPP,³⁶ what is likely to be an important reason for the exceptionally fast second protonation step in TPP.

The energies that are necessary to distort the mono-protonated forms to the geometry of the diprotonated form, without the second proton, and, opposing, those energies that are necessary to distort the diprotonated form to match the geometry of the mono-protonated form are shown in Fig. 3. The given energies on the left side of each panel, respectively, refer to summed educt energies while those on the right side of each panel refer to the products, respectively, according to the chemical reactions shown below.

Oop-distortions of the mono-protonated forms so that they match the geometries of the diprotonated species $[H_3P]^{2+}$ according to the reaction shown below, necessitate energies of about 50 kJ mol^{-1} in the gas phase and about 100 to 170 kJ mol^{-1} if dielectric screening is considered.[§] This hypothetical reaction path is indicated by dashed lines in Fig. 3.



In contrast, if the geometry of the mono-protonated form is preserved and a proton is attached according to the reaction below[¶] (see dotted lines in Fig. 3),



an energetic downhill reaction results for the tetra-*meso*-aryl-substituted porphyrins in case of a gas-phase calculation. In contrast, dielectric screening in the conductor limit strongly increases the energies, relative to the monoprotated species, of the hypothetical transient diprotonated species at the geometries of their mono-protonated forms. Essentially, these COSMO-energies show that the energies of the $[H_4P]^{2+}$ states and those of the tetra-*meso*-aryl derivatives are more sensitive towards their solvent environment than the energies of the $[H_3P]^{2+}$ state and those of the tetra-*meso*-aryl derivatives. It appears reasonable to assume that in a realistic picture, the transient species are not perfectly screened as in the COSMO-calculation. Thus, we obtain a picture in between the gas-phase and the COSMO results, shown in Fig. 3, where the protonation kinetics are substantially influenced by the energy of $[H_4P]^{2+}$ related states, *i.e.* by the flexibility of the diprotonated species. This is in excellent agreement with both the AIMD result, that

[§] Theoretical details: Energy calculated by a single-point calculation on a geometry-optimized diprotonated structure with one proton removed.

[¶] Theoretical details: Energy derived in a constrained optimization that just allows modification of the second proton coordinates.



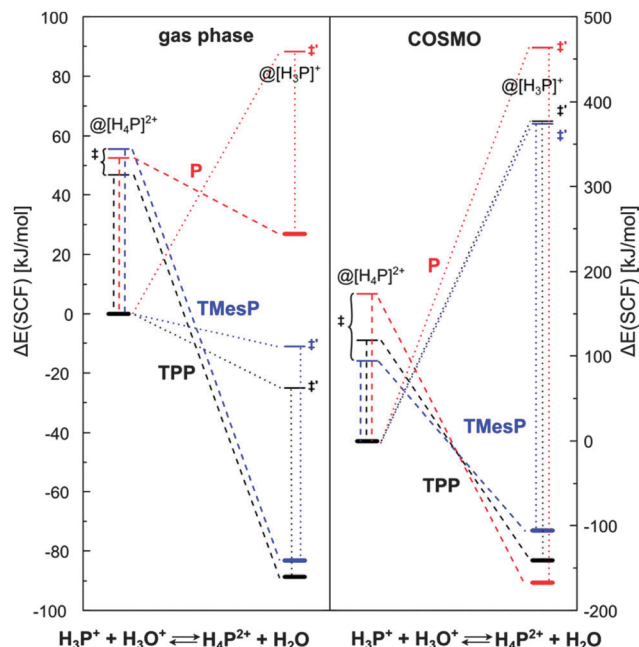


Fig. 3 Energies of geometry-optimized mono- and diprotonated species as well as states with geometries matching those of the neighboring protonation state. †labels: mono-protonated species that are distorted to match the geometries of their diprotonated forms (via removing one proton from the diprotonated forms). ‡labels: diprotonated species that are distorted to match the geometries of their mono-protonated forms (via constraint geometry optimization after adding a proton to the mono-protonated forms).

particularly $[\text{H}_4\text{TPP}]^{2+}$ has an outstandingly large flexibility, and the experimental findings of an exceptional fast second protonation in the case of TPP as compared to P and TMesP.

Steric interactions between the *meso*-aryl substituents and the macrocycle

To unravel the influence of the *meso*-aryl substituents on the above discussed geometric and energetic features, as well as their interplay and the impact of electronic changes on the porphyrin induced by the substituents, the impact of their torsions against the porphyrins on molecular properties are analyzed in the following.

In the case of H_2TMesP , the mesityl groups are perpendicular to the porphyrin and get successively twisted reaching the saddle-type geometry in $[\text{H}_4\text{TMesP}]^{2+}$, as shown in Table 4. The two mesityls (*meso*-C5 and *meso*-C10) next to pyrrole B, which is protonated first, tilt constantly with protonation (H_2TMesP , $[\text{H}_3\text{TMesP}]^+$, $[\text{H}_4\text{TMesP}]^{2+}$: $|\delta| \approx 90, 82, 71^\circ$), while the mesityls next to the non-protonated pyrrole D in $[\text{H}_3\text{TMesP}]^+$ (*meso*-C15 and *meso*-C20) change their tilt only little compared to the free base structure (H_2TMesP , $[\text{H}_3\text{TMesP}]^+$: $|\delta| \approx 90, 88^\circ$). This change in the mesityl tilt upon the first protonation is small because the neighboring pyrroles A and C do not get tilted significantly oop upon protonation of H_2TMesP and the non-protonated pyrrole D stays approximately in plane in $[\text{H}_3\text{TMesP}]^+$, thus transmitting little torsion to the mesityls neighboring the non-protonated pyrrole D. In contrast, the *meso*-phenyl groups in H_2TPP are tilted significantly already

in the free base form ($|\delta| \approx 70^\circ$) and the *meso*-C5 and *meso*-C10 phenyls next to the first protonated pyrrole B constantly tilt in slightly larger steps upon protonation (H_2TPP , $[\text{H}_3\text{TPP}]^+$, $[\text{H}_4\text{TPP}]^{2+}$: $|\delta| \approx 70, 55, 46^\circ$) as compared to the mesityl-substituted derivatives discussed above. The two phenyls next to the non-protonated pyrrole D in $[\text{H}_3\text{TPP}]^+$ (at *meso*-C15 and *meso*-C20) significantly tilt already in the mono-protonated species $[\text{H}_3\text{TPP}]^+$, thus both sterically interacting with the non-protonated pyrrole D, which tilts oop exceptionally strong in $[\text{H}_3\text{TPP}]^+$ ($\theta^N = -8^\circ$) as compared to $[\text{H}_3\text{P}]^+$ and $[\text{H}_3\text{TMesP}]^+$ ($\theta^N = 0, -2^\circ$).

In the case of TMesP, protonated pyrrole tilts are transmitted via the *meso*-mesityls to the neighboring pyrroles, thus significantly disturbing the porphyrin geometry and π -electron delocalization already at small tilt angles at which the mesityl groups do not yet extend the porphyrin π -system. Because of the lacking *ortho*-methyl groups in the *meso*-phenyl substituents, their tilts and the porphyrin distortion are significantly smaller than in the case of the mesityl groups.

From calculations of torsional profiles of the model substances 5-mono-phenyl- and 5-mono-mesityl-porphyrin (H_2monoPP , $\text{H}_2\text{monoMesP}$), which are shown in the upper panel of Fig. 4, it can be concluded that at room temperature the mesityl tilts are approximately $90 \pm 10^\circ$. In contrast, phenyl tilts are possible between approximately 40° and 140° , with two energetic minima at $\sim 60^\circ$ and $\sim 120^\circ$ in H_2monoPP . At these tilt angles, considerable π -conjugation between the porphyrin and the *meso*-phenyls is present⁷⁶ and improves further with co-planarization between both moieties. The energy gain because of this increased π -electron delocalization compensates the energy demand caused by the accompanied oop-tilt of the non-protonated porphyrin according to quantum chemical geometry optimization.

Upon the first protonation, both monoPP and monoMesP torsional profiles get steeper, which is attributed to the higher steric strain in the protonated macrocycle core. Thus, just one energetic minimum at approximately 50° is obtained for $[\text{H}_3\text{monoPP}]^+$, as shown in the middle panel of Fig. 4. At room temperature ($\Delta E = kT = 2.4 \text{ kJ mol}^{-1}$), the torsional angle in $[\text{H}_3\text{monoPP}]^+$ varies approximately between 45° and 72° ($\Delta\delta \approx 27^\circ$). This variation is similar for $[\text{H}_3\text{monoPP}]^+$ and $[\text{H}_3\text{monoMesP}]^+$ ($63^\circ < \delta < 95^\circ$, $\Delta\delta \approx 32^\circ$).

Electronic interactions between the *meso*-aryl substituents and the macrocycle

While the proton affinities (PAs), which neglect solvent interactions with the porphyrins, for TPP and TMesP are almost identical (cf. differences between pseudo-gas phase energies in Table 1, $\Delta E_{\text{QC}}^* = 1070$ and $-1073 \text{ kJ mol}^{-1}$), mono-*meso*-MesP and mono-*meso*-PP have smaller and significantly different PAs (~ 990 and $\sim 1050 \text{ kJ mol}^{-1}$). Thus, the more *meso*-aryl substituents are attached to the porphyrin, the better a proton gets stabilized. In the case of the mono-*meso*-aryl derivatives, the electron-pushing methyl substituents of the *meso*-mesityl group lead to a significantly larger stabilization (by $\sim 60 \text{ kJ mol}^{-1}$) of the protonated species than in the case of unsubstituted *meso*-phenyl functionalization. As shown in Fig. 4, this energy relation holds true for a very broad



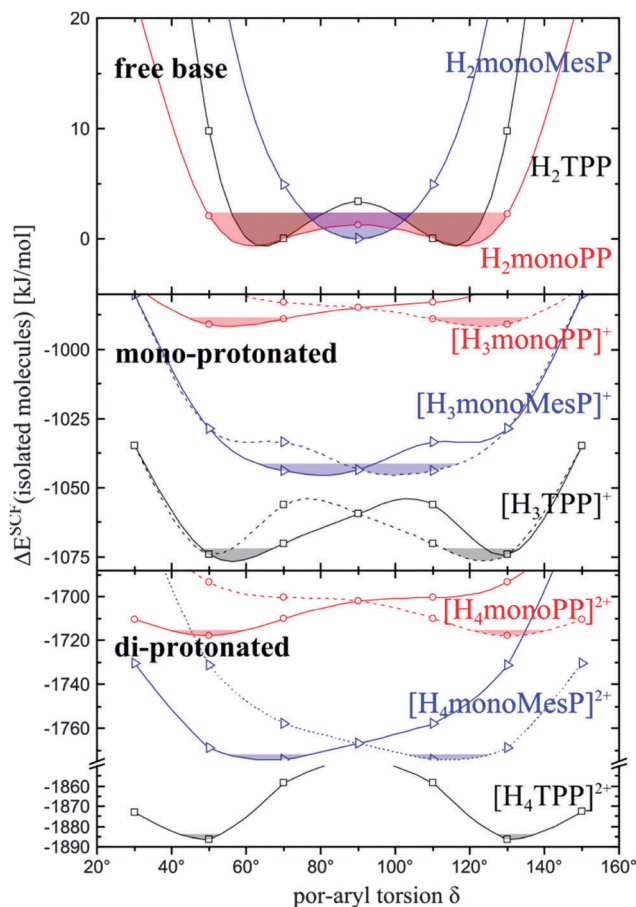


Fig. 4 Torsional profiles of mono-*meso*-substituted porphyrins. The influence of the different *meso*-substituents (phenyl, P, and mesityl, Mes) on the SCF energy upon torsion against the porphyrin in the model compounds H_2 monoPP and H_2 monoMesP and their mono- and diprotonated forms is compared. The torsional profiles of the mono-protonated species refer to protonation of the pyrrole that is closest to the *meso*-substituent. The dashed graphs are mirrored representatives of the continuous line graphs. Both are asymmetric for the protonated species since the pyrrole oop-twist is different for the two pyrroles neighboring the *meso*-substituent. The ΔE^{SCF} values refer to the absolute SCF energies with respect to the energetic minima for H_2 monoPP and H_2 monoMesP. The gray background (darker: H_2 monoPP, lighter: H_2 monoMesP) represents torsional angles at which the $\Delta E^{\text{SCF}} < RT$, i.e. which should be the dominant contribution at room temperature.

range of por-aryl torsions, for perpendicular geometries as well as for those around $\delta = 60^\circ$, where considerable π -conjugation between the porphyrin and the *meso*-aryl groups is present. Thus, improved π -conjugation due to co-planarization between the porphyrin and *meso*-aryl groups plays a minor role for proton stabilization as compared to the electron pushing influence of the *meso*-aryl groups.

However, addition of more *meso*-phenyl groups yields a larger additional cumulative stabilization of the protonated species than addition of *meso*-mesityl groups, since $[H_3\text{TPP}]^+$ and $[H_3\text{TMesP}]^+$ show virtually identical proton affinities ΔE_{QC}^* . The reason for the similar stabilization by the four *meso*-phenyls and the four *meso*-mesityls is the interplay between electronic and geometric stabilization. While the mesityls push more electron density to the core, the porphyrin, with its core

crowded with three hydrogens, can sterically relax better towards a saddle-like structure in the case of phenyl- than in the case of mesityl-substitution, as shown by the geometry analysis and the difference in the shape of the torsional potentials discussed above. The difference in sterical relaxation just gets pronounced at multiple *meso*-aryls attached to the porphyrin core and is negligible for the mono-mesityl-substituted derivatives, where the electronic influence of the *meso*-substituents dominates the difference in proton affinities.

For the diprotonated species, the torsional profiles get steeper and, in accordance to the geometric features of the energetically most favorable structures discussed above, the positions of the minima slightly shift to larger porphyrin-aryl torsions. Thus, the torsional profiles of monoPP and monoMesP confirm that smaller porphyrin-aryl torsion angles are energetically favorable in monoPP as compared to monoMesP and that these torsional angles generally shift to smaller values on protonation. In the case of TPP, the steric demands of the four *meso*-phenyls accumulate and cause significantly steeper torsional profiles than those of monoPP, while the energetic minima are approximately at the same torsional angles. Since the minima-positions in the torsional profile of TPP are retained from the one of monoPP, it is assumed that the quantitative comparison between the torsional profile of monoPP and monoMesP can be transferred qualitatively to the comparison of porphyrin-aryl torsions between TPP and TMesP.

Attraction of and accessibility for protons

As briefly mentioned above, the θ^{N} -oop exposure of the non-protonated pyrrole D in $[H_3\text{TPP}]^+$ is exceptionally large (-8° , cf. Table 4) as compared to the angles of $[H_3\text{P}]^+$ and $[H_3\text{TMesP}]^+$. Thus, it appears probable that the nitrogen atom is sterically and electrically shielded significantly less against protons by the mono-protonated and positively charged macrocycle in the case of TPP as compared to the other derivatives, hence probably contributing to its instantaneous (on the time scale of the titration experiments) protonation. However, in a deeper analysis of accessibility, the molecular surface that can be reached by solvent molecules (the “solvent excluded surface” SES⁷⁷) needs to be considered. This SES is one boundary of the reactivity volume V_{react} ,⁷⁶ while a sphere around the reaction center with a radius of 2 Å that accounts for intermolecular distances is the outer boundary of this reactivity volume. Thus, the V_{react} accounts for the geometric accessibility of a reaction center for reactants. These reactivity volumes V_{react} are represented by small 3D-voxels in the molecular representations in Fig. 5 and are divided into one cap above and one beneath the macrocycles.

In the free base form, the non-protonated pyrrole D is not tilted oop in $H_2\text{P}$ and $H_2\text{TMesP}$, as listed in Table 1. Accordingly, the reactivity caps above and beneath the macrocycle show identical volumes ($V_{\text{react}}(H_2\text{P}) = 1.88 \text{ \AA}^3$, $V_{\text{react}}(H_2\text{TMesP}) = 1.87 \text{ \AA}^3$, see gray lines in the bottom panel of Fig. 5). Since the porphyrin in TPP is already distorted in the free base form, the reactivity cap above the macrocycle is larger than the one beneath ($V_{\text{react}}(\text{above/beneath})$: 2.14/1.67 \AA^3). This asymmetry



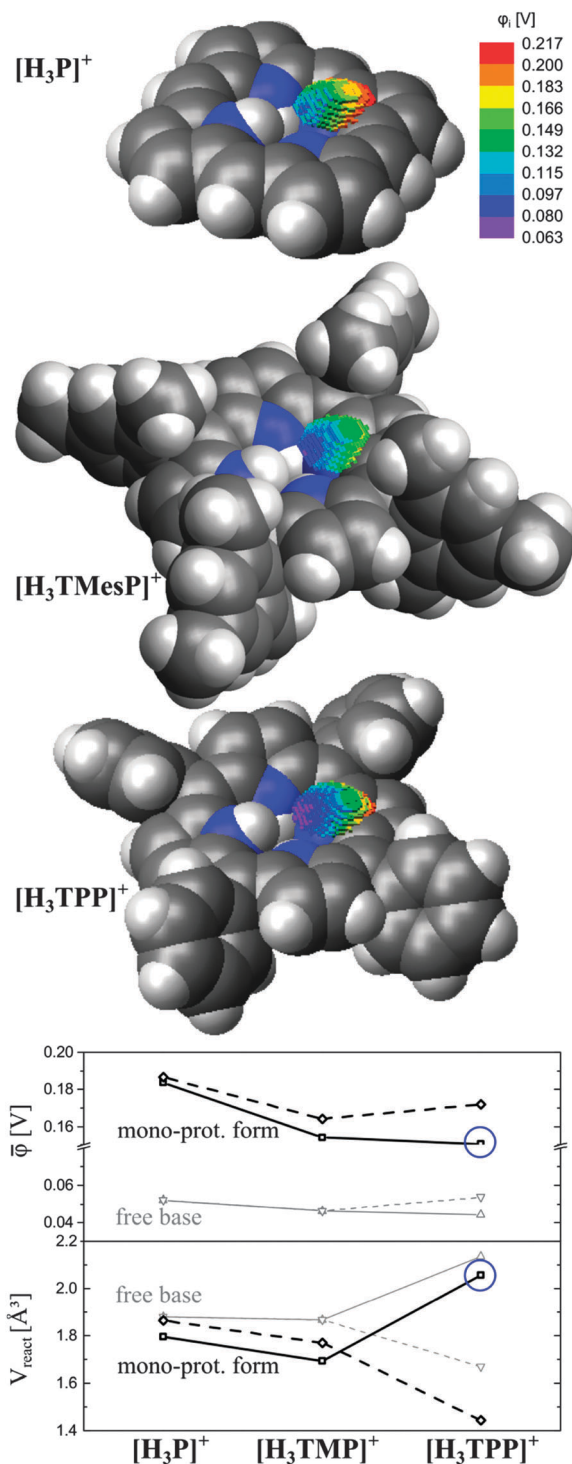


Fig. 5 van der Waals surface representations of the mono-protonated species $[H_3P]^+$, $[H_3TMesP]^+$, $[H_3TPP]^+$ together with voxels that represent the reactivity volumes V_{react} . The color coding of the voxels corresponds to the electrostatic potential in the center of each voxel. Color coding of atoms: white, gray, and blue – H, C, and N. Reactivity volumes V_{react} with their associated mean electrostatic potentials $\bar{\phi}$ are found above and beneath each of the macrocycles and are labelled by solid and dashed lines, respectively, in the bottom panel.

in the V_{react} of the different caps is present for the mono-protonated species of all considered porphyrin derivatives because of their

pyrrole oop-torsion (V_{react} (above/beneath): $[H_3P]^+$: 1.80/1.87 \AA^3 , $[H_3TMesP]^+$: 1.69/1.77, $[H_3TPP]^+$: 2.06/1.44). In accordance with the geometric analysis discussed above (see Table 4), $[H_3TPP]^+$ shows an exceptionally large reactivity volume of $V_{\text{react}} = 2.06 \text{ \AA}^3$. In contrast to the discussion on dihedral angles, *i.e.* despite the large dihedral angle of pyrrole D in the mono-protonated form of TPP as compared to the free base, the accessibility volume in $[H_3TPP]^+$ is actually slightly smaller than in H_2TPP due to steric shielding in the buckled macrocycle.

For proton transfer not just the geometric accessibility, but also the electron distribution around the reaction center is important. The corresponding mean electrostatic potentials $\bar{\phi}$ within the reactivity volumes of the free bases are slightly positive, $\bar{\phi} \approx 0.05 \text{ V}$, and increase to 1.5–1.9 V for the mono-protonated species. $[H_3P]^+$ shows higher $\bar{\phi}$ -values, *i.e.* stronger proton repulsion, than $[H_3TMesP]^+$ and $[H_3TPP]^+$. Even if this mean electrostatic potential within the reactivity volumes of these tetra-*meso*-aryl-substituted derivatives are similar, their $\phi(r)$ -distributions differ, as shown in Fig. 5. The purple colored spot of low positive, and for protons weakly repulsive, potentials at the nitrogen lone-pair is larger and has a lower minimum- ϕ of 0.066 V in $[H_3TPP]^+$ as compared to $[H_3TMesP]^+$ with $\phi_{\text{min}} = 0.078 \text{ V}$.

Consequently, the considered nitrogen atom is exceptionally exposed in the free base as well as in the mono-protonated form of TPP as compared to P and TMesP. Additionally, for protons repulsive electrostatic potential shows a large spot with low ϕ -values, *i.e.* weak repulsion, at the position of the nitrogen lone pair in $[H_3TPP]^+$ in comparison to $[H_3TMesP]^+$ and $[H_3TPP]^+$. It is expected that both effects contribute to the fact that the second protonation of TPP is faster than for the other derivatives and similarly fast or even faster than the first protonation of TPP, thus explaining why the experimental absorption spectra detected during titration are dominated by H_2TPP and $[H_4TPP]^{2+}$ related absorption features.

Conclusions

In conclusion, the ease of oop-buckling, *i.e.* molecular flexibility, appears to be the key for understanding the exceptionally fast second protonation of TPP as compared to P and TMesP. Our *ab initio* molecular dynamics simulations show that particularly $[H_4TPP]^{2+}$ is significantly more flexible than $[H_4TMesP]^{2+}$, while just small differences are found between the mono-protonated species. Thus, the accessible geometric conformational space of $[H_4TPP]^{2+}$ is significantly larger than for $[H_4TMesP]^{2+}$. Accordingly, transitions between geometries of mono- and diprotonated forms are strongly facilitated for TPP as compared to TMesP or P, which corresponds to a lower activation barrier and a faster kinetics of the second protonation. Together with the energetic relations and the pronounced binding spot in $[H_3TPP]^+$ mentioned above, the exceptional molecular flexibility of $[H_4TPP]^{2+}$ explains why spectroscopic signatures of just the free base and diprotonated species of TPP are found in the titration experiments, but not $[H_3TPP]^+$. Thus, the very similar experimental findings on the exceptional basicities of the



monoprotonated form of sterically less demanding *meso*-aryl porphyrins of Pasternack *et al.* and of Stone and Fleischer,^{43,44} which were discussed focusing on buckling of the free-base forms, need to be reappraised in view of our quantum chemical results focusing on the diprotonated species.

The general conclusions that can be drawn from this work for the development of improved molecular designs are the following. Despite the low π -conjugation between the *meso*-mesityl substituents and the porphyrin center due to their approximately perpendicular arrangement, the mesityls still significantly push electron density to the porphyrin's core, as shown by quantification of proton stabilization in the section on torsional profiles, *i.e.* approximately perpendicular geometries between the porphyrin and its substituents do not dramatically disturb their electronic interactions. Nevertheless, those perpendicular geometries are usually caused by steric strain that significantly affects reaction kinetics, as shown by the relative speed of the second protonation between the *meso*-mesityl and *meso*-phenyl-porphyrins in this work.

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