The first near-linear bis(amide) f-block complex: a blueprint for a high temperature single molecule magnet†

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We report the first near-linear bis(amide) 4f-block compound and show that this novel structure, if implemented with dysprosium(III), would have unprecedented single molecule magnet (SMM) properties with an energy barrier, $U_{\text{eff}}$, for reorientation of magnetization of 1800 cm$^{-1}$.

Since their initial discovery, single molecule magnets (SMMs) have been lauded as candidates for high density data storage devices. A major breakthrough in the field occurred in 2003 with the observation of SMM behavior in a monometallic {TbPc$_2$} complex with an energy barrier, $U_{\text{eff}} = 230$ cm$^{-1}$. The ensuing decade saw rapid growth in lanthanide SMMs with the $U_{\text{eff}}$ barrier to magnetization reversal increased to 652 cm$^{-1}$ for another derivative of {TbPc$_2$}, and 585 cm$^{-1}$ for a polymeric Dy$_5$ complex. The highest blocking temperature $T_B$ (i.e. the temperature at which hysteresis is observed) was also increased to 14 K, via an N$_2$ bridge in a {Tb$_2$N$_3$} complex.

Although three of these milestones employ the Tb(III) ion, by far the most utilized lanthanide ion in SMMs is Dy(III) by virtue of its unique electronic structure. Apart from a radical-bridged {[Dy$_2$N$_2$]} complex,

Very low coordination numbers for 4f-ions are difficult to achieve as these are large, electropositive ions, which require a sterically demanding ligand. Such a pro-ligand HN(SiiPr$_3$)$_2$ was designed, and synthesised from ClSiiPr$_3$ and LiHN(SiiPr$_3$)$_2$, and this was converted to the group 1 transfer agent [KN(SiPr$_3$)$_2$] with KH.

An electrostatic model for the design of ideal ligand environments to exploit the maximal anisotropy of Dy(III) has been postulated, and shown to be in good agreement with multi-configurational complete active space Self consistent field (CASSCF) $ab$ initio calculations that are often employed to examine 4f complexes, pioneered by Chibotaru. Electrostatic approaches suggest that the optimal ligand environment to exploit the oblate spheroidal electron density of Dy(III) is axial, where rigorously axial systems have the benefit of maintaining a single, unique quantization axis for the total angular momentum $m_I$ states. A set of unadulterated $m_I$ states implies that the probability of quantum tunnelling of the magnetization (QTM) is reduced, therefore increasing magnetic relaxation times.

The simplest axial ligand environment is a linear two-coordinate complex with donor atoms exclusively on a single Cartesian axis; the $U_{\text{eff}}$ barrier is so large for the {Dy$_5$} and {Dy$_4$K$_2$} alkoxide complexes because of the strongly axially repulsive crystal field potentials along the local $z$-direction of each Dy$^{III}$. Other compounds such as {[C$_6$H$_6$]Ln} or Cloke’s bis(arene) lanthanide complexes are sometimes described as linear, but lack donor atoms directly on the axis. Linear 3d-metal compounds also show remarkable magnetic behaviour with very high $U_{\text{eff}}$ values. A one coordinate lanthanide complex [Dy(O)]$^+$ has been considered theoretically with a very large $U_{\text{eff}}$ predicted, however such an entity is not chemically feasible.

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Reacting two equivalents of [KN(SiPr$_3$)$_2$] with samarium(II) diiodide yields the mononuclear homoleptic bis(amide) complex, [(Pr$_3$Si)$_2$N–Sm–N(SiPr$_3$)$_2$] (Fig. 1, see ESI† for details).

Complex 1 is the first near-linear f-element complex, with an N–Sm–N angle of 175.52(18)° in the solid state (Fig. 2, see ESI† for details); this near-linearity contrasts with the bent C–Ln–C angles of [Ln$^{III}$C(SiMe$_3$)$_2$] complexes (Ln = Sm, Yb, Eu). The

$$[\text{Sm(THF)}_2 + 2 [\text{K[N(SiPr}_3]_2])] \rightarrow \text{toluene} \rightarrow [\text{Pr}_3\text{Si} \rightarrow \text{Pr}_3\text{Si}]_2$$

Fig. 1 Synthetic route to 1.
bulky iPr groups are vital for the isolation of a homoleptic complex, as [Sm[N(SiMe3)2]2(THF)2] exhibits additional O-donors.21 The Sm–N distances in 1 [2.483(6) Å] are longer than those observed in [Sm[N(SiMe3)2]2(THF)2] [mean Sm–N 2.433(9) Å] but this is compensated by 1 exhibiting four short Sm···Cmethine distances [Sm···C 3.082(7)–3.224(7) Å] that are closer than the analogous Sm···Cmethol contacts observed in [Sm[N(SiMe3)2]2(THF)2] [Sm···C 3.32(1)–3.46(1) Å].21 The approximately planar SmNSi2 fragments in 1 are staggered with respect to each other (twist angle of 44.42°), with the deviation from 90° attributed to agostic Sm···Cmethine interactions.

Formally each nitrogen atom carries a single negative charge and the SmIII ion is divalent, with an [Xe]4f6 configuration. The f6 configuration leads to a formally diamagnetic 7F0 ground state, with close lying excited states that provide a non-zero magnetic moment at room temperature. Magnetic measurements on 1 give a room temperature magnetic moment of 3.62 μB that falls towards zero at low temperature (Fig. S2 and S3, ESI†). This is clearly incompatible with interesting low temperature magnetic behaviour. However, the structure of 1 is close to the ideal linear arrangement to stabilize the large angular momentum states of DyIII and produce monstrous uniaxial magnetic anisotropy.

Such a DyIII compound is challenging to make; we believe a route via the heteroleptic [Dy[N(SiPr3)2]2] treated with the potassium salt of a large anion might work through precipitation of a potassium iodide. Other routes can be imagined, and here we present predictions of the magnetic properties of such a complex, intending to inspire synthetic work towards the linear DyIII complex, and, more ambitiously, the iso-electronic TbII analogue.

The properties of [[Pr6Si8–N–DyIII–N(SiPr3)2]]2 are predicted by CASSCF/RASSI/SINGLE_ANISO and ab initio calculations (see ESI† for details) employing the structure of 1, where SmIII has been replaced by DyIII. The validity of the method was tested by calculating the variable temperature magnetic behavior of 1, where the agreement is excellent (Fig. S2 and S3, ESI†). DyIII has a 6H15/2 ground multiplet, which is split by the field into eight Kramer’s doublets with total angular momentum projections mJ = ±1/2, ±3/2, . . . ±15/2. The ab initio calculations show that the lowest six Kramer’s doublets are strongly mixed; a characteristic of low symmetry complexes due to the lack of a rigorous molecular C∞ axis.14 Along the main magnetic axis these two states can be expressed as ψab = 64%|−1/2⟩ + 26%|+1/2⟩ and ψcd = 68%|−1/2⟩ + 31%|+1/2⟩ and (Table S2, ESI†), giving the most energetic Kramers doublet a large gJ value of ~17.5 perpendicular to the main magnetic axis.

Magnetic relaxation in lanthanides follows three possible routes: (1) QTM within the ground doublet (e.g. |−15/2⟩ → |+15/2⟩ in Fig. 3), (2) thermally assisted QTM (TA-QTM) via excited states (e.g. |−15/2⟩ → |−13/2⟩ → |+13/2⟩ → |+15/2⟩), or (3) an Orbach process composed of direct and/or Raman mechanisms (e.g. |−15/2⟩ → |−13/2⟩ → |+15/2⟩). The most probable pathway depends on the composition of the states involved and their interactions with phonons. For example, the slow magnetic relaxation for DyIII coupled with its large g-factor suggests that TA-QTM will occur via the excited states which transverse...
g-factors above a certain threshold or where its main magnetic axis is non-collinear with that of the ground state. All non-QTM transitions are induced by the vibrational modes of the lattice (phonons) which create local oscillating magnetic fields through modulation of dipolar fields as well as an oscillating crystal field potential. To a first approximation, we can associate the probability of a phonon induced transition with the average magnetic and crystal field perturbation matrix elements (see ESI† for details).

Compared to all known DyIII complexes the calculated properties for 2 are unique with very small transverse g-factors and a common principal axis for the lowest six Kramers doublets. This suggests that both the probability of QTM within the ground doublet and TAZ-QTM is vanishingly small until the two most energetic doublets. Orbach relaxation is also strongly disfavoured in the low lying states as magnetic transition probabilities due to phonons are miniscule (Fig. 3). Efficient magnetic relaxation will only occur via the highest energy doublets (Fig. 3, Fig. S4 and Tables S4 and S5, ESI†). Therefore the ab initio calculation predicts $U_{\text{eff}} \approx 1800 \, \text{cm}^{-1}$ for 2 – far greater than any complex to date. Whilst such calculations may over-estimate the energies like to thank J. P. S. Walsh for assistance with graphics.

**Notes and references**