Central-transition double-quantum sideband NMR spectroscopy of half-integer quadrupolar nuclei: estimating internuclear distances and probing clusters within multi-spin networks†

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We introduce a strategy to estimate the size of clusters of recoupled homonuclear half-integer quadrupolar nuclei under magic-angle spinning (MAS) conditions, by combining double-quantum (2Q) sideband NMR experiments with an approximate numerical analysis based on the summation of all spin-pairs present over a given radius of the structure. The experiment relies solely on the evolution of homonuclear 2Q coherences (2QC) among the central-transitions (CT) of half-integer spins and is suitable for probing clusters in network structures, such as those encountered in large groups of oxide-based materials. Experimental 11B, 23Na and 27Al NMR results are presented on bis(catecholato)diboron, Na2SO4 and Al2O3, respectively; in each case, the growth of the spin-cluster size was monitored from a series of experiments that employed progressively lengthened 2QC excitation intervals. Our new approach is the first option for probing larger constellations of half-integer spins; it provides similar information as the "multiple-quantum spin counting" experiment, which is well-established for spin-1/2 applications but has hitherto not been demonstrated for half-integer spins undergoing MAS. We also discuss various options for determining the internuclear distance within a (nearly) isolated pair of half-integer spins by comparing the experimental 2Q sideband NMR spectra with results from numerical simulations involving various degrees of approximation.

1 Introduction

Precise interatomic-distance information constitutes the core of a detailed understanding of a given molecular structure, and is generally best determined by diffraction methods. However, whenever the possibility of obtaining large and well-ordered crystals is not feasible, either because the structure lacks long-range order or feature dynamics within its sub-units, solid state nuclear magnetic resonance (NMR) spectroscopy offers an attractive alternative. The power of high-resolution magic-angle spinning (MAS) NMR that utilizes distance-dependent through-space dipolar interactions restored by “dipolar recoupling” techniques is well documented for providing highly accurate information about selected internuclear distances and dihedral angles in well-defined clusters of spin-1/2 nuclei—primarily involving 1H, 13C and 15N in bio-molecules. This is generally arranged by site-specific isotopic labeling, which is feasible thanks to the combination of well-established synthetic routes and the low natural abundance of 13C and 15N. Recent progress in dipolar recoupling NMR methodology is gradually also permitting the simultaneous accurate extraction of increased numbers of internuclear distances in uniformly (13C, 15N)-enriched biological macromolecules.

Yet, comparatively few distance-determinations involving homonuclear (i.e., of the same species) spins-1/2 are demonstrated in the large groups of inorganic network structures, e.g., minerals, ceramics, glasses and porous materials, where the nature of the structures implies very large networks of mutually coupled spins. Further, the state-of-affairs is much worse when considering even semi-quantitative internuclear distance-estimates between half-integer spin quadrupolar nuclei despite a strong impetus for exploiting them for structural studies by NMR, stemming from their dominating abundance in inorganic materials. Owing to their multi Zeeman-level character and the presence of strong quadrupolar interactions, the development of homonuclear dipolar recoupling methodology targeting half-integer spins is a complex task, as reviewed in ref. 14 and 15. Yet, during the past decade, several alternatives have been proposed. However, these recoupling techniques,
including dipolar self-recoupling under MAS\textsuperscript{32,33} and experimentation on non-spinning samples,\textsuperscript{34–36} have predominantly been applied for gaining qualitative information about connectivities, proximities, and distributions of homonuclear spins, or relative orientations of NMR interaction tensors,\textsuperscript{15,17–19,22–28,30,31,37–44} whereas present reports on direct internuclear distance-estimates remains very sparse.\textsuperscript{20,21,45,46} Furthermore, recent work exploits homonuclear through-bond ("p") interactions among half-integer spins to gain information about site-connectivities in inorganic structures.\textsuperscript{37–49}

All existing strategies to directly measure internuclear distances via their associated dipolar interactions within homonuclear pairs of half-integer spins undergoing MAS, have relied on double-quantum (2Q) dipolar recoupling radio-frequency (rf) pulse sequences. The magnetization may either (i) be allowed to evolve under their resulting effective dipolar Hamiltonian,\textsuperscript{26,21} or be converted into 2Q coherences (2QC) among the central transitions (CTs) of the two spins, onwards referred to as "2QCCT". The latter option has involved either (ii) monitoring of the 2QCCT generation (which is directly dependent on the dipolar interaction-strength) for prolonged excitation intervals (\(t_{\text{exc}}\))\textsuperscript{45} or (iii) using 2Q–1Q 2D correlation spectroscopy, which results in an internuclear distance-sensitive spinning sideband manifold if the \(t_{\text{exc}}\)-evolution is sampled faster than at completed rotational periods.\textsuperscript{46} Strategy (iii) was introduced by Spiess and co-workers as a general means of determining NMR interaction-parameters in systems involving at least one spin-1/2,\textsuperscript{6,50–54} All three approaches to interatomic distance-measurements have been applied extensively in the area of spins-1/2, which possesses a very substantial arsenal of 2Q-recoupling schemes (reviewed in ref. 27, 29 and 30) are beneficial for cases relatively low resonance-frequency dispersions,\textsuperscript{27,29} the CT 2Q sideband experiment may utilize any 2QCCT generator for half-integer spins,\textsuperscript{23,24,27–31} or supercyclets thereof. The pulse trains described in ref. 27, 29 and 30 are beneficial for cases featuring large frequency-spreads among the recoupled spin sites. Next follows a Hahn spin-echo of duration \(T\), timed such

### 2 Double-quantum sideband NMR spectroscopy

#### 2.1 Experimental protocol

Fig. 1 displays the rf pulse scheme for the CT 2Q sideband NMR experiment. It starts by enhancing the spectral signal-to-noise ratio by increasing the population-difference across the CT.\textsuperscript{39} Here, we employ a smoothly attenuated single frequency sweep,\textsuperscript{46,65} but any alternative technique is equally applicable, e.g., those of ref. 66–69. The longitudinal magnetization is subsequently converted into 2QCCT, in this work by repeating a \(R_2^1R_2^{-1}\) 2Q-recoupling sequence\textsuperscript{26} \(m\) times to arrange a 2QC excitation interval of \(t_{\text{exc}} = 4m\tau_c\), where the sample rotational period \(\tau_c\) relates to the angular spinning frequency \(\omega_s\) as \(\tau_c = 2\pi/\omega_s\). This pulse scheme conforms to the R-symmetry class,\textsuperscript{2,3,70} and requires that the spin-I CT nutation frequency \(\omega_{\text{CT}} = I(I + 1)/2\gamma B_1/2\) equals \(\omega_{\text{res}} = \omega_s/2\), where \(\gamma\) is the spin gyromagnetic ratio and \(B_1\) the rf amplitude. All recoupling pulses are sandwiched between two CT-selective π/2-pulses, each of duration \(\tau_{\text{rec}}\), as discussed in detail in ref. 26, 27 and 29. While we recommend using the \(R_2^1R_2^{-1}\) scheme for 2QCCT-excitation in samples displaying relatively low resonance-frequency dispersions,\textsuperscript{27,29} the CT 2Q sideband experiment may utilize any 2QCCT generator for half-integer spins,\textsuperscript{23,24,27–31} or supercyclets thereof. The pulse trains described in ref. 27, 29 and 30 are beneficial for cases featuring large frequency-spreads among the recoupled spin sites.
that $T + 2\pi \delta \tau = p \tau$, where $p$ is any even integer; the phase-cycling\textsuperscript{27} of the CT-selective $\pi$-pulse should retain solely the $\pm 2Q_{CT} \rightarrow \mp 2Q_{CT}$ pathways to block $2QC$ involving the satellite transitions (STs),\textsuperscript{24} as indicated in Fig. 1c.

In the absence of $t_1$-evolution, the $2QC_{CT}$ coherences are transferred directly into detectable CT single-quantum coherences (1QC) by repeating the $\pi/2$-sandwiched $R_2^C R_2^C$\textsuperscript{1} sequence for a reconversion interval $\tau_{\text{rec}}$, followed by a CT-selective $\pi/2$-pulse (Fig. 1a). While one may arrange that $\tau_{\text{exc}} \neq \tau_{\text{rec}}$\textsuperscript{72} herein we employ equal excitation and reconversion segments throughout all our experiments. However, when $t_1 > 0$, the incrementation of the evolution interval has strong bearings on the $2Q$-1Q 2D NMR spectrum. When internatural proximity-information is desired, $t_1$ is normally sampled at integer multiples of $\tau_{1/4}$,\textsuperscript{24,26,28,30,39,41,56} which ensures that the rotor position is identical at the start of each $\tau_{\text{exc}}$ and $\tau_{\text{rec}}$ event, thereby forcing all spinning sidebands to fold onto each respective centerband (“isotropic”) frequency: the 2D spectral coordinates then reveals all unique $2QC_{CT}$ frequencies $\omega_{jk}^C$ along the indirect (vertical) dimension, correlated with their respective CT 1QC frequencies $\omega_1$ and $\omega_2$ that appear along the direct (horizontal) spectral dimension.

However, if the $t_1$-incrementation is performed in steps smaller than $\tau_r$, a spinning sideband formation occurs along the indirect spectral dimension. It typically originates from two distinct mechanisms: (i) Rotor-modulation, stemming from $2QC_{CT}$-evolution under anisotropic interactions, produces sidebands at positions shifted by $p\omega_1$ ($p$ is any integer) from the centerband frequency. They are well-known and become significant in all MAS NMR experimentation when the magnitude(s) of the anisotropic interaction(s) exceed the spinning frequency. (ii) Rotor-encoded spinning sidebands, on the other hand, only appear at odd integer multiples of $\omega_1$ and arise if the spatial sample position is different at the start of each 2QC excitation and reconversion segment.\textsuperscript{6,50–54} These sideband-intensities are essentially independent on the relative sizes of the MAS frequency and the chemical shift and quadrupolar anisotropies active during the evolution interval $t_1$, provided that rotor-modulated contributions are eliminated by employing fast MAS. However, the precise appearance of the rotor-encoded sideband-envelope depends on the product $b_{jk}^\tau_{\text{exc}} = b_{jk}^\tau_{\text{rec}}$, which provides a route to measure the dipolar coupling constant $b_{jk}\textsuperscript{6,50–54}$ related as $r_{jk}^{-3}$ to the internuclear separation $r_{jk}$ between spins $j$ and $k$. We exploit this feature to estimate internuclear distances between half-integer spins,\textsuperscript{46} as explored further herein. There are two options for implementing the $t_1$-sampling in Fig. 1a, which have bearings both on the experimental time and the precise numerical protocol to extract information from the rotor-encoded sidebands:

(1) The protocol in Fig. 1a is used to record a complete $2Q$-1Q 2D NMR spectrum by sampling the evolution period over a ms-range, i.e., over many rotational periods of the sample. This is required to achieve sufficient signal-resolution in the presence of several distinct spin-pairs in the structure.

(2) The evolution interval is only sampled over the first rotational period, i.e., at $N$ time-points $t_1 = j\tau_1/N$, with $j = 0, 1, \ldots, N - 1$, which in practice is implemented as an array of $N$ 1D NMR experiments. This approach is only beneficial when the sample produces one unique $2QC_{CT}$ frequency, but is then highly advantageous for reducing the experimental time (as well as that for the numerical simulations discussed below), as the $2QC_{CT}$ spinning sideband manifold results directly on Fourier transformation (with respect to $t_1$) of the “pseudo-2D” data-set $\{S(t_1, \tau_2 = 0)\}$; these values are in practice obtained by integrating the total intensity along the $\omega_2$ spectral dimension for each consecutive $t_1$-value. Similar approaches have been utilized both for numerical simulations of MAS NMR experiments (e.g., see ref. 73), and in several experimental NMR protocols, such as for spin-counting\textsuperscript{57–64} and molecular torsion-angle determinations.\textsuperscript{7}

Implementation (1) was used for the bis(catecholato)-diboron sample and protocol (2) was employed for the samples $\text{Na}_2\text{SO}_4$ and $\text{Al}_2\text{O}_3$. The $2QC_{CT}$ sideband manifold along the $\omega_1$ dimension will be depicted by a “stick-spectrum”, where the height of each stick reflects the respective integrated sideband intensity, plotted against the sideband order $k$.

2.2 Strategies for extracting inter-atomic distances and spin-cluster sizes

2.2.1 Overview. As the number of spinning sidebands of significant amplitude in the CT 2Q NMR spectrum is proportional to the product $b_{jk}^\tau_{\text{exc}}$, whenever an experimental result is available from one $\{b_{jk}^\tau_{\text{exc}}\}$ pair of parameters, a qualitative estimation of an unknown dipolar coupling constant is straightforward by counting the number of odd-order sidebands in the spectrum. For low values of $b_{jk}^\tau_{\text{exc}}$, only the first-order ($\pm 1$) sidebands appear, whereas the presence of a significant dipolar interaction between two spins is signified by the growth of the higher-order ($\pm 3, \pm 5, \pm 7, \ldots$) sidebands for increasing $2QC_{CT}$ excitation intervals.\textsuperscript{46} Yet, accurate quantitative internuclear-distance estimates must involve explicit numerical simulations of the spin dynamics in one way or the other. We then distinguish between two scenarios:

(A) The rare case when the internuclear separation within an isolated pair of half-integer spins is to be determined. Such spin systems may feature one internuclear distance $r_{jk}$ that is much shorter than all others; a typical scenario is a spin-pair within a molecular unit where the crystal structure involves an ordered arrangement of such units. This case may be handled by directly fitting the experimental data to numerically exact spin-pair simulations.\textsuperscript{46} Even for structures built from separate molecular units, truly isolated spin-pairs are extremely rare due to intermolecular dipolar interactions. This will be expanded on in Section 4.1, where we explore the consequences of including longer-range intermolecular couplings, and hence the deviation from a truly “isolated” $^{11}\text{B}$-$^{11}\text{B}$ pair, for determining the B–B interatomic distance in the bis(catecholato)-diboron molecule.

In scenario (B) of a large continuous spin network, for which the CT 2Q sideband NMR experiment will create a significant number of recoupled spin-pairs, the size of such a spin-$I$ cluster grows for increasing $\tau_{\text{exc}}$. Then a large number of distinct spin-pairs associated with a spread in $r_{jk}$-values contributes to the
sideband pattern of the CT 2Q experiment. The measurement of a “single distance” is then not a well-defined concept for such cases, and one must simultaneously determine the large number of distances associated with the entire set of individual spin-pairs. Accurate numerical simulations cannot address this task due to the substantial time required to calculate the NMR spectrum from even a few dipolar coupled $I = 3/2$ spins, and the prohibitive scaling of the computation time for higher-spin numbers. Numerically exact simulations accounting for more than four coupled spins-3/2 appears presently intractable by using standard computer resources, at least if employed for the purpose of iterative fitting to experimental data. Hence, one must resort to approximating the complex spin dynamics, as well as lower the ambition level from extracting multiple inter-nuclear distances to just defining the spin cluster size. Here we employ the strategy of ref. 56, which amounts to first calculating the NMR spectrum from each unique spin-pair in the structure over a given radius $R$ onwards denoted the “cut-off distance”—followed by summation over all contributing spectra. Note that whereas each spin-pair is treated numerically exact, the procedure ignores all cross-couplings between spins of distinct pairs and their mutual interferences. Despite invoking a crude approximation that may not in general be justified rigorously, this strategy allowed for the successful determination of the absolute orientations of all $^{23}$Na quadrupolar tensors in Na$_2$SO$_4$ further validation of the approach is gained from the results presented in Section 4.1.

2.2.2 Spin-pair summation for calculating NMR spectra from multi-spin systems. Ref. 56 describes the numerical modeling procedure in detail; here we limit ourselves to recapitulating its key stages, as well as outlining the relationship between the number of dipolar interactions and number of contributing spins in the cluster.

For each pair $jk$ of unique crystallographic sites, all dipolar coupling constants and their accompanying orientations are extracted from the (known) crystal structure: firstly, the primitive unit cell is built from its corresponding asymmetric unit. Secondly, the central unit cell is expanded into an odd order supercell, e.g. $3 \times 3 \times 3$ or $5 \times 5 \times 5$. Finally, all dipolar interactions from crystallographic sites in the central unit cell sites in the same, as well as to surrounding, unit cells are located and sorted according to the relative absolute values of their dipolar coupling constants, i.e., by the corresponding internuclear distances. In the analysis below, dipolar interactions are considered up to the cut-off distance $R$. The corresponding minimum dipolar coupling constant is denoted by $b(R) = -\mu_0/4\pi I I^2 h R^{-3}$. Dipolar interactions that are related by translational symmetry are considered equivalent. Each of the inequivalent dipolar couplings is a member of one set of $M(b_{jk})$ symmetry-related couplings, where $M(b_{jk})$ is referred to as the dipolar multiplicity: all dipolar interactions within this set are associated with the same magnitude of $b_{jk}$ but differ in their orientations in the crystal structure. Concrete examples of this approach are provided in Sections 4.1–4.3 for the samples used in this study.

Next, the NMR spectrum for each dipolar vector is calculated: it may represent either the entire 2Q–1Q 2D correlation NMR spectrum or its $\omega_1$-projection, depending on which underlying $t_1$-sampling approach is used (Section 2.1). This computation also involves all relevant chemical shift and quadrupolar interaction parameters, whose values must be a priori known. In general, all $M(b_{jk})$ members that feature a fixed value of the dipolar coupling constant $b_{jk}$ may be divided into $N(b_{jk})$ subsets. The NMR spectrum generated from each subset is distinct from that of any other subset, meaning that there are $N(b_{jk})$ unique (i.e., distinct) spectra. In contrast, all dipolar coupling orientations within a given subset produce identical NMR spectra, owing to additional symmetries in the NMR responses. Then explicit simulations must only be performed for one sole dipolar vector orientation from each of the $N(b_{jk})$ subsets. The resulting set of dipolar coupling parameters is denoted $[(b_{jk}, \Omega_{DD}^{(jk)})]$, with the index $i$ running over all distinct subsets; $i = 1, \ldots, N(b_{jk})$. The corresponding individual simulated spectra are labeled $S_{jk}(\omega_1, \omega_2; \tau_{exc}, b_{jk}, \Omega_{DD}^{(jk)})$. With this approach the computation time is dictated by the factor $M(b_{jk})/N(b_{jk})$.

The NMR spectrum $S_{\text{sum}}^R(\omega_1, \omega_2; \tau_{exc})$ that is representative of the entire spin system—the latter being defined by the $\tau_{exc}$-determined radius $R$—is constructed by summing all individual contributions $S_{jk}(\omega_1, \omega_2; \tau_{exc}, b_{jk}, \Omega_{DD}^{(jk)})$, while considering their respective dipolar multiplicities:

$$S_{\text{sum}}^R(\omega_1, \omega_2; \tau_{exc}) = \sum_{jk} M(b_{jk}) N(b_{jk}) \sum_{i=1}^{N(b_{jk})} S_{jk}(\omega_1, \omega_2; \tau_{exc}, b_{jk}, \Omega_{DD}^{(jk)}).$$

The number $N_{\text{spins}}$ of interacting (re)coupled half-integer spins may be assessed by comparing the experimental 2Q–1Q sideband manifold with those calculated by eqn (1) for progressively expanded radii of the (potentially) infinitely large structure, with the aim of locating the $R$-value that minimizes the root mean square (RMS) deviation between experimental and calculated NMR spectra. Once the value of $R$ is determined, the resulting set of contributing dipolar interactions may be translated into the number of individual half-integer spins $N_{\text{spins}}$ involved in the corresponding coupling network, as follows: all distinct crystallographic sites in the constructed supercell are indexed consecutively. Next, the set of contributing dipolar interactions is translated into a set of index pairs, whereupon $N_{\text{spins}}$ is obtained from the number of unique spin indices found within this set.

3 Materials and methods

3.1 Solid-state NMR experiments

All experiments were performed at a static magnetic field of $B_0 = 11.75$ T using a Bruker Avance-III spectrometer. The sample of bis(catecholato)diboron was studied at a $^{11}$B Larmor frequency ($-\gamma_B B_0$) of $-160.5$ MHz in a Varian 3.2 mm double-resonance standard-bore Varian MAS probehead. The powder of bis(catecholato)diboron was purchased from Sigma-Aldrich.
and was spun at 25 kHz. This corresponds to an rf field of 6.25 kHz during dipolar recoupling, i.e., to a CT nutation frequency equal to half the MAS frequency, as for the HORROR condition. The CT selective 90° and 180° pulses operated at 9.0 kHz rf field amplitude. Complete 2D nutation experiments were performed to calibrate all rf amplitudes.

To enhance the CT population difference, a single-frequency sweep (SFS) of 2 ms duration was performed at the same rf field-strength, by using an rf frequency offset and a sweep bandwidth of 500 kHz and 800 kHz, respectively, where the beginning and end of the sweep-pulse shape were attenuated by a $\sin^2$ and $\cos^2$ function, respectively. The States-TPPI scheme was employed to obtain 2D pure absorption lineshapes and for distinguishing positive and negative 2Q coherences.

The experiments on the Na$_2$SO$_4$ sample were performed on a 3.2 mm double-resonance standard-bore Varian MAS probehead at a 27Al Larmor frequency of $-132.3$ MHz and a spinning frequency of 20 kHz. The rf field during the 2QC$_{CT}$-recoupling was set to 5 kHz. The CT selective 90° and 180° pulses operated at 7.151 kHz. The CT population was enhanced by a SFS of 2 ms duration at the same rf field strength, 550 kHz rf offset and 900 kHz bandwidth, where the beginning and the end of the sweep where attenuated, as described above.

Prior to the NMR experiments, the sample of Al$_2$O$_3$ was heated to 1360 °C for 96 h to ensure a phase-pure $\alpha$-Al$_2$O$_3$ specimen, which was confirmed by a directly excited 27Al MAS NMR spectrum. The NMR experiments were performed on a 4 mm double-resonance standard-bore Varian MAS probehead at a 27Al Larmor frequency of $-130.3$ MHz and a spinning frequency of 10 kHz. The rf field during the 2QC$_{CT}$-recoupling was set to 1.674 kHz. The CT selective 90° and 180° pulses operated at 980 Hz. The CT population was enhanced by a SFS of 2 ms duration at 2.577 kHz rf field strength, 500 kHz rf offset and 800 kHz bandwidth, as described above.

### 3.2 Numerical spin dynamics simulations

The numerical simulations of the 11B, 27Al and 23Na 2QC$_{CT}$ sideband patterns were performed by using either the SIMPSON package (version 3.0.1) or software developed in Stockholm. The simulations invoked the experimental parameters listed in Section 3.1. All pulses were explicitly accounted for, except for the CT population enhancement and the 90° read pulse.

Both the initial density operator and the detection operator represented CT 2Q-magnetization. The CT 2Q$_C$ were selected by nulling all other elements in the density matrix before and after the 180° spin-echo pulse. The CT 2Q sideband amplitudes were simulated by incrementing $t_1$ in $N$ steps of duration $t_2/N$, with $N$ equal to 35, 35 and 16 for the simulations involving 11B, 27Al and 23Na, respectively. The results of the separately calculated time-dependent signal amplitudes of the +2Q and −2Q were combined to form “cosine” and “sine” NMR signal components required by the procedure of States et al. The resulting complex $t_1$-signal was subjected to a complex Fourier transformation to obtain the integrated CT 2Q spinning sideband amplitudes. Powder averaging was accomplished using a set of 6044 triplets of Euler angles selected according to the ZCW scheme.

In all cases, we used the respective experimentally estimated values of the quadrupolar coupling constant and asymmetry parameter listed in Table 1. In case of bis(catecholato)diboron, the parameters of the CSA tensor and the orientation of the quadrupolar tensor as calculated by GIPAW-DFT were employed in the numerical simulations. Both for $\alpha$-Al$_2$O$_3$ and Na$_2$SO$_4$ the experimentally estimated value of the anisotropic chemical shift and asymmetry parameter were used together with the corresponding tensor orientations obtained by GIPAW-DFT calculations.

### 3.3 Quantum chemical calculations

For the GIPAW-DFT calculations of the quadrupolar and chemical shift tensors CASTEP and CASTEP-NMR (version 4.5) were used together with Accelrys’ Materials Studio. The generalized gradient approximation with the Perdew–Burke–Ernzerhof exchange correlation functional and “on-the-fly” pseudo-potentials were chosen. In all cases the plane wave cutoff energy was selected to be 550 eV. In case of bis(catecholato)diboron the $k$-point grid was set to $5 \times 2 \times 4$. Starting from the structure determined by single-crystal XRD, only the proton positions were optimized, whereupon the NMR parameters were calculated. In case of Na$_2$SO$_4$ the conventional orthorhombic unit cell of the single-crystal structure was converted to the primitive unit cell prior to the DFT calculations employing a $5 \times 5 \times 5$ $k$-point grid. For $\alpha$-Al$_2$O$_3$ the conventional hexagonal unit cell of the structure determined by single-crystal XRD was converted to the

<table>
<thead>
<tr>
<th>Sample</th>
<th>C$_O$ [MHz]</th>
<th>$\eta_0$</th>
<th>$\delta_{\text{iso}}$ [ppm]</th>
<th>$\delta_{\text{aniso}}$ [ppm]</th>
<th>$\eta$</th>
<th>Ref.</th>
</tr>
</thead>
<tbody>
<tr>
<td>C$_2$H$_5$B$_2$O$_4$</td>
<td>2.71</td>
<td>0.745</td>
<td>28.6</td>
<td>-</td>
<td>-</td>
<td>46</td>
</tr>
<tr>
<td>Na$_2$SO$_4$</td>
<td>2.6</td>
<td>0.6</td>
<td>-8.5</td>
<td>-18.9</td>
<td>0.0</td>
<td>37 and 98</td>
</tr>
<tr>
<td>$\alpha$-Al$_2$O$_3$</td>
<td>2.40</td>
<td>0.0</td>
<td>18.8</td>
<td>17.3</td>
<td>0.0</td>
<td>99</td>
</tr>
<tr>
<td>C$_{12}$H$_8$B$_2$O$_4$</td>
<td>3.18</td>
<td>0.77</td>
<td>{38.7, 86.2, −135.0}</td>
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<td>−22.6</td>
<td>0.67</td>
</tr>
<tr>
<td>Na$_2$SO$_4$</td>
<td>3.20</td>
<td>0.60</td>
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<tr>
<td>$\alpha$-Al$_2$O$_3$</td>
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<td>{0, 0, 0}</td>
<td>538.1</td>
<td>7.90</td>
<td>0.0</td>
</tr>
</tbody>
</table>

This table lists the 11B NMR parameters of bis(catecholato)diboron (C$_2$H$_5$B$_2$O$_4$) and Na$_2$SO$_4$, and the 27Al NMR parameters of $\alpha$-Al$_2$O$_3$ as determined experimentally and by GIPAW-DFT calculations. The quadrupolar coupling constant and asymmetry parameter are denoted $\delta_{\text{iso}}$ and $\delta_{\text{aniso}}$, respectively. $\eta$ is the experimental isotropic chemical shift, whereas $\delta_{\text{iso}}$ and $\eta_0$ are the calculated absolute isotropic chemical shielding. The reported values of the anisotropic chemical shift $\delta_{\text{aniso}}$ and the asymmetry parameter $\eta$ conform to the Haarberg convention. The Euler angles $\Omega_{BS}$ give the relative orientation of the principal axis system of the interaction tensor and a crystal fixed frame.
primitive unit cell prior to the DFT calculations with a $6 \times 6 \times 6$ k-point grid.

4 Results

4.1 The $^{11}\text{B}^{11}\text{B}$ pair in bis(catecholato)diboron ($I = 3/2$)

We recently demonstrated that CT 2Q sideband NMR spectroscopy allowed for accurately determining the internuclear distance in the pair of directly bonded boron atoms in bis(catecholato)diboron by iterative fitting of numerically exact $^{11}\text{B}^{11}\text{B}$ two-spin simulations to experimental results, i.e., using approach (A) in Section 2.2.1. Here we analyze the source of these discrepancies further by employing numerically exact multi-spin simulations (comprising up to 4 spins) to account for the contributions from weaker intermolecular $^{11}\text{B}^{11}\text{B}$ dipolar couplings between neighboring molecules, whose inclusion will be shown to provide excellent agreement with the experimental sideband patterns. Indeed, Holland et al. reported similar effects of intensified low-order sideband amplitudes in the context of $^1\text{H}$ NMR, and attributed them to the presence of distributions of dipolar-coupling strengths. We further demonstrate that the relatively effortless summation of sideband NMR spectra from a larger set of spin-pairs (i.e., strategy (B) in Section 2.2.1) reproduce the results from the rigorous multi-spin analysis. Noteworthy, the estimated internuclear distance of the directly bonded B–B pair remains largely unaffected by the approach for numerical analysis.

Fig. 2a shows the structure of a single bis(catecholato)diboron molecule together with the principal axis systems of the $^{11}\text{B}$ chemical shift and quadrupolar tensors as obtained by GIPAW-DFT calculations. Fig. 2b displays the nearest neighboring molecules in the crystal structure, where the four boron sites in closest spatial contact are labeled by 1–4. From the XRD structure,55 the internuclear distances are determined to $r_{12}^{\text{XRD}} = 168\,\text{pm}$, $r_{23}^{\text{XRD}} = 382\,\text{pm}$ and $r_{13}^{\text{XRD}} = r_{24}^{\text{XRD}} = 475\,\text{pm}$. The corresponding $^{11}\text{B}^{11}\text{B}$ dipolar coupling constants are given by $b^{11}_{12}/2\pi = -2619\,\text{Hz}$, $b^{11}_{23}/2\pi = -222\,\text{Hz}$ and $b^{11}_{13}/2\pi = -116\,\text{Hz}$. Note that the intramolecular constant of the directly bonded (1,2) boron sites is $\approx 12$ times larger than the second largest intermolecular counterpart. In ref. 46 we therefore analyzed the experimental CT 2Q spinning sideband amplitudes by solely employing spin-pair simulations.

4.1.1 Spin-pair analysis. Fig. 3b shows the integrated odd-order CT 2Q experimental $^{11}\text{B}$ sideband amplitudes (dashed sticks) obtained from bis(catecholato)diboron for $\tau_{\text{exc}} = 480\,\mu\text{s}$ at an external field of 11.75 $\text{T}$ and spinning frequency of 25 $\text{kHz}$. The solid sticks depict the corresponding best-fit simulation for a $^{11}\text{B}^{11}\text{B}$ spin pair with $b_{12}/2\pi = -2580$ $\text{Hz}$ (equivalent to $r_{12}^{\text{NMR}} = 168.6\,\text{pm}$). While the calculated intensities of all high-order sidebands agree very well with the experimental ones, the experimental ±1 sidebands are more intense relative to those of the two-spin simulations. We performed a series of simulations by varying the dipolar coupling constant in steps of 10 $\text{Hz}$ between $-3080$ and $-2080$ $\text{Hz}$. In each case, we calculated the sideband pattern and located the minimum of sum-squared deviation (SSD) between the simulated $(a_{\text{sim}}^{(k)})$ and integrated experimental $(a_{\text{exp}}^{(k)})$ sideband amplitudes, only accounting for the sideband orders $k = \pm 3, \pm 5, \pm 7$, and $\pm 9$:

$$\text{SSD} = \sum_{k} \left( a_{\text{exp}}^{(k)} - A a_{\text{sim}}^{(k)} \right)^2$$

where $A$ represents a free parameter that scales the simulated sideband amplitudes.
coupling constant in the range of \((-3.5, -2.1)\) MHz had a very weak effect on the sideband amplitudes. This generally also applies to the orientations of the quadrupolar interaction, as well as to the magnitude and orientation of the CSA tensor. After accounting for the potential contributions from these interactions, our analysis amounted to a reliable estimate of the \(^{11}\text{B}^{-^{11}\text{B}}\) dipolar coupling constant as \(-2580 (\pm 140, -120)\) Hz, i.e., corresponding to a B-B distance of (168.6 \pm 3) pm, with relative uncertainties of about \(\pm 5\%\) and \(\pm 2\%\) in the dipolar coupling constant and the internuclear distance, respectively.

The observed dipolar coupling constant is about 1.5\% smaller than that calculated from the XRD structure, and the corresponding NMR determined distance is only about 0.5\% larger than that obtained from XRD. This small discrepancy can be attributed to the thermal motion of the boron positions: for the case of glycine, Ishii et al. calculated by molecular dynamics simulations that direct bond distances between non-proton nuclei measured by NMR at 273 K are typically 1.9–2.9\% longer than the distances estimated by diffraction techniques. Furthermore, the anisotropy of the \(J\)-coupling \(\langle J \rangle\) has the same symmetry under sample rotation and rf pulses as the direct dipolar interaction. Hence, in principle, only the effective spin–spin interaction \(D_{\text{eff}} = b_{12}/2\pi - \langle J \rangle /3\) can be determined with the CT 2Q sideband experiment. However, DFT calculations (not shown here), indicate that \(J\) is in the order of 13–17 Hz for the \(^{11}\text{B}^{-^{11}\text{B}}\) pair in bis-catecholato-diboron. Consequently, its potential bearings on the estimation of the dipolar coupling is negligible compared to the effects from thermal motions.

### 4.1.2 Multi-spin analysis

Owing to the high natural abundance of 80.1\% for \(^{11}\text{B}\) and the close molecular packing in the crystal structure of bis(catecholato)diboron (see Fig. 2b), long-range \(^{11}\text{B}^{-^{11}\text{B}}\) dipolar interactions need to be considered for explaining the intense \(\pm 1\) CT 2Q spinning sidebands. Table S1 of the ESI\(^\dagger\) summarizes the various probabilities of finding a certain multiple \(^{11}\text{B}\) spin system for these 4 boron sites, revealing that isolated \(^{11}\text{B}^{-^{11}\text{B}}\) spin pairs are rarely encountered in the structure. While the remaining 19.9\% boron isotopes are spin-3 \(^{10}\text{B}\) nuclei, the heteronuclear \(^{11}\text{B}^{-^{10}\text{B}}\) dipolar interactions are effectively decoupled during the 2Q-\(J\)-recoupling sequence on the \(^{11}\text{B}\).

Fig. 3c shows the calculated CT 2Q sideband amplitude patterns obtained from the various \(^{11}\text{B}\) multiple-spin systems by employing numerically exact simulations; each is scaled by employing numerically exact simulations; each is scaled by the corresponding probability of finding that particular constellation of coupled \(^{11}\text{B}\) sites (see Table S1 in the ESI\(^\dagger\)). Fig. 3d contrasts their sum (solid sticks) with the experimental amplitudes (dashed sticks): an excellent agreement is observed. We stress that the present summation procedure is only required due to the presence of the \(^{10}\text{B}\) isotope in the sample; it must not be confused with the analogous spin-pair summation protocol described in Section 2.2.3 and utilized in Section 4.1.3, as well as for the 100\% abundant \(^{27}\text{Al}\) and \(^{23}\text{Na}\) nuclides below.

### 4.1.3 Spin-pair summation

We further applied approach (B) of Section 2.2.2 that derives the NMR spectrum by summation over all NMR responses from the individual spin-pairs...
encountered within a given cut-off distance $R$. Fig. 2c–f show the boron sites in the crystal structure of bis(catecholato)diboron. Each reveals the different $^{11}$B–$^{11}$B site pairs with a particular internuclear distance and dipolar coupling. Site pairs that contribute to the multiplicity $M_{jk}$ of the respective coupling are depicted as solid lines, whilst equivalent pairs not contributing to that particular multiplicity are drawn with dashed lines. Table S2 in the ESI† summarizes all distinct $^{11}$B–$^{11}$B site-pair and their corresponding multiplicities up to $R = 621$ pm, as well as the orientation of each dipolar vector with respect to the crystal frame.

Fig. 3e shows the RMS deviation between the experimental 2Q$_{CT}$-sideband amplitudes and the sum of two-spin simulations according to eqn (1), plotted against $R$. The RMS calculation accounted for all odd sideband orders $k = \pm 1, \pm 3, \pm 5, \pm 7, \pm 9$, with each RMS value normalized with respect to its minimum (RMS$_{min}$) found at $R = 382$ pm; Fig. 3f contrasts the respective sideband pattern obtained by summing all two-spin simulations (solid sticks) with the corresponding experimental result (dashed sticks). The very good agreement observed between the simulation and experiment is only marginally worse than that by using numerically exact multi-spin simulations in Fig. 3d. This validates the approach leading to eqn (1), i.e. that the spin dynamics of a large number of mutually dipolar-coupled spins may in these applications be reasonably well approximated by a sum over pair-wise responses.

### 4.2 The $^{23}$Na ($I = 3/2$) network in Na$_2$SO$_4$

We next consider application of CT 2Q sideband NMR spectroscopy to the dipolar-coupled network of $^{23}$Na spins ($I = 3/2$) in Na$_2$SO$_4$, which involves one crystallographically unique Na site. Fig. 4a illustrates the single-crystal XRD-derived structure, also indicating the directions of the principal axes of the $^{23}$Na chemical shift and quadrupolar tensors, as obtained by GIPAW-DFT calculations. Fig. 4b–e convey the spatial relationship between the various equivalent and inequivalent $^{23}$Na–$^{23}$Na pairs. Table S2 in the ESI† summarizes all inequivalent $^{23}$Na–$^{23}$Na dipolar interactions, the corresponding multiplicities, and vector directions over a cut-off radius of $R = 590$ pm.

Fig. 5 plots the experimental $^{23}$Na 2Q$_{CT}$-filtered (2QF$_{CT}$) efficiencies against the excitation interval $\tau_{exc}$. The results were obtained at a spinning frequency of 20 kHz, an external field of 11.75 T and two distinct rf carrier frequency positions (relative to the center of mass of the $^{23}$Na powder lineshape); this produced resonance offsets of 0 Hz and 1080 Hz, and the associated maximum 2QF$_{CT}$ efficiencies of 15.9% and 23.4%, respectively. However, regardless of the precise value of the resonance offset, an optimum signal intensity was observed at $\tau_{exc} = \tau_{exc,max} = 0.8$ ms.

The left panel of Fig. 6 shows RMS deviations between experimental 2Q$_{CT}$-sideband amplitudes and those calculated by summing over simulated two-spin responses [eqn (1)], plotted against the cut-off distance $R$, and with 2Q$_{CT}$ excitation periods increasing from top to bottom; each curve is normalized with respect to its respective minimum. The RMS deviations were determined both by including (squares) and omitting
Structural analyses must rather aim at validating if the atomic coordinates (and its associated set of dipolar interactions) of a proposed structural model may reproduce the experimental results. At a selected 2\textsuperscript{Q}CT excitation interval, the experiment provides the number of interacting spins within the macroscopically large spin ensemble. For a finite-sized cluster in a structure, \( N_{\text{pairs}} \) corresponds to its number of spins. The value of \( N_{\text{pairs}} \) may then be estimated by locating the \( R \)-value that minimizes the RMS deviation between the experimental 2\textsuperscript{Q}CT sideband manifold and the sum over all simulated spin-pair responses therein. Table 2 lists the cluster-size \( N_{\text{pairs}} \) of the Na\textsubscript{2}SO\textsubscript{4} structure assessed at each excitation period; \( N_{\text{pairs}} \) increases from 12 to 22 when \( \tau_{\text{exc}} \) increases from 1.6 ms to 2.4 ms.

We note that Duer\textsuperscript{37} used a related approach to analyze the \(^{23}\text{Na} - ^{23}\text{Na} \) dipolar coupling-network in Na\textsubscript{2}SO\textsubscript{4}, by modeling the experimental triple-quantum MAS (3\text{QMAS})\textsuperscript{33} NMR spectrum via numerical simulations that also accounted for multiple homonuclear couplings. This experiment produces a sideband manifold stemming from the rotor-modulated evolution of (single-spin) \( 3Q \) under the, by slow-MAS incompletely averaged, homonuclear \(^{23}\text{Na} - ^{23}\text{Na} \) couplings, which in the absence of explicit homonuclear dipolar-recoupling were truncated to solely comprise the mutually commuting \( I_{ij} I_{kz} \) operators of the dipolar Hamiltonians in the analysis.\textsuperscript{37} By including the homonuclear couplings to the five nearest neighbors, Duer successfully reproduced the isotropic projection of the 3\text{QMAS} experiment by exploiting the known dipolar coupling-constants and orientations in the crystal structure of Na\textsubscript{2}SO\textsubscript{4}. This approach differs to ours, both concerning the origin of the sideband-formation, as well as in the numerical analysis, where we include the full form of the (high-field) dipolar Hamiltonian, yet assuming a set of mutually non-interacting spin-pairs.

**4.3 The \( ^{27}\text{Al} \) \( (I = 5/2) \) network in \( \alpha\text{-Al}_2\text{O}_3 \)**

We further evaluated the CT 2\text{Q} sideband NMR experiment and its accompanying “spin-counting” analysis in the context of the \(^{27}\text{Al} \) spins \((I = 5/2)\) in \( \alpha\text{-Al}_2\text{O}_3 \), which involves one crystallographically unique Al site coordinated by oxygen atoms. Fig. 7a reveals a very densely populated network of Al, stemming from the presence of face-shared AlO\textsubscript{6} octahedra. The GIPAW-DFT-derived principal axis systems of the \(^{27}\text{Al} \) chemical shift and quadrupolar tensors are shown in Fig. 7a and b.

<table>
<thead>
<tr>
<th>Sample</th>
<th>( \tau_{\text{exc}} ) [ms]</th>
<th>( R ) [pm]</th>
<th>( N_{\text{pairs}} )</th>
<th>( N_{\text{spins}} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \text{C}_2\text{H}_4\text{B}_2\text{O}_4 )</td>
<td>0.48</td>
<td>382</td>
<td>8</td>
<td>12</td>
</tr>
<tr>
<td>Na\textsubscript{2}SO\textsubscript{4}</td>
<td>1.6</td>
<td>321</td>
<td>8</td>
<td>12</td>
</tr>
<tr>
<td></td>
<td>2.4</td>
<td>420</td>
<td>31</td>
<td>22</td>
</tr>
<tr>
<td>( \alpha\text{-Al}_2\text{O}_3 )</td>
<td>0.4</td>
<td>266</td>
<td>2</td>
<td>4</td>
</tr>
<tr>
<td></td>
<td>0.8</td>
<td>279</td>
<td>14</td>
<td>16</td>
</tr>
<tr>
<td></td>
<td>1.0</td>
<td>350</td>
<td>50</td>
<td>34</td>
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</tbody>
</table>

Fig. 6 Left column: normalized RMS deviations between experimental \(^{23}\text{Na} \) CT 2\text{Q} spinning sideband amplitudes and summed spin-pair simulations in Na\textsubscript{2}SO\textsubscript{4} plotted against the maximum Na–Na internuclear distance included in the summation. A and B correspond to the results obtained in the absence of presence of the CSA in the simulations, respectively. Right column: dashed sticks depict the experimental CT 2\text{Q} sideband amplitudes. Solid sticks represent the sum of numerical simulations obtained by including all Na–Na distances up to the \( R \)-value marked in the left column.
A selection of inequivalent Al–Al pairs is shown in Fig. 7c–f, while all coupling constants, multiplicities and dipolar vector orientations out to $R = 591$ pm are listed in Table S2 in the ESI.†

Fig. 8 plots experimental $^{27}$Al 2QF CT efficiencies against increasing $t_{\text{exc}}$-values, as obtained from $\alpha$-$\text{Al}_2\text{O}_3$ at a magnetic field of 11.75 T and 10 kHz MAS frequency. A maximum 2QCT efficiency of 7% was obtained around $t_{\text{exc}} = 0.4$ ms, which is considerably lower than the $^{23}$Na 2QF$_\text{CT}$ efficiency in Na$_2$SO$_4$. While these signal losses stem partially from a more rapid $^{27}$Al relaxation during rf application compared to that for $^{23}$Na in Na$_2$SO$_4$, they primarily originate from a significantly higher CT-magnetization leakage out to the satellite transitions for a spin-$5/2$ compared to the spin-$3/2$ case, as discussed in ref. 15 and 28. Furthermore, for a fixed dipolar-coupling constant, the 2QC buildup rate is higher for a $I = 5/2$ pair compared to that of two spins-$3/2$ (vide infra);$^{28}$ together with the overall stronger $^{27}$Al–$^{27}$Al dipolar coupling constants in the $\alpha$-$\text{Al}_2\text{O}_3$ structure (see Table S2, ESI†), this accounts for the nearly doubled 2QC generation rate; compare the results of Fig. 5 and 8.

Fig. 9 displays the experimental and simulated 2Q$_\text{CT}$-sideband spectra for increasing 2Q$_\text{CT}$ excitation periods. At the shortest value $t_{\text{exc}} = t_{\text{max}}$ = 0.4 ms where the 2QF$_\text{CT}$ efficiency is optimized (Fig. 9a and b), the experimental 2Q$_\text{CT}$-sideband spectrum reveals only the ±1 sideband orders. The best-fit spin-pair sum was obtained from the set of couplings over $R = 266$ pm, corresponding to only two Al–Al pairs and a cluster of four $^{27}$Al sites (i.e., $N_{\text{spins}} = 4$); see Table 2. On the other hand, at $t_{\text{exc}} = 0.8$ ms (Fig. 9c), the sideband manifold comprises both ±1 and ±3 sideband orders. Now only a fair agreement is observed for the “best-fit” NMR sideband pattern, which resulted by summing the simulated responses from all pairs over $R = 279$ pm ($N_{\text{pairs}} = 14$), and translating into a cluster comprising 16 unique $^{27}$Al sites. At the longest 2Q$_\text{CT}$ excitation period $t_{\text{exc}} = 1.0$ ms (Fig. 9f), the experimental 2Q$_\text{CT}$-sideband pattern includes all
sideband orders up to $\pm 5$, although the amplitudes of the highest order are very weak. The experimental result is well-reproduced by a sum over 50 spin-pairs within $R = 350$ pm, corresponding to a cluster comprising 34 spins (Table 2).

5 Discussion

Here we discuss the prospects of applying each of strategies (A) and (B) of Section 2.2 for analyzing dipolar-coupled spin networks and estimating internuclear distances.

5.1 Estimating internuclear distances

The results for the $^{11}$B spin-3/2 pair in bis(catecholato) diboron demonstrates that CT 2Q spectroscopy is an excellent method to accurately determine internuclear distances within (isolated) pairs of half-integer quadrupolar nuclei. The sensitivity of the method increases with the number of 2QCT sideband orders that are considered in the analysis. Given that the number of spinning sidebands with significant intensity depends on the product $b_3 \tau_{\text{exc}}$, the excitation period may be progressively increased to ensure that a sufficiently large sideband manifold is available for the estimation of weak dipolar coupling constants, whereas the upper limit of the $\tau_{\text{exc}}$ is only limited by signal-damping effects from relaxation. For a truly isolated spin-pair, the dipolar coupling constant is then readily extracted through effortless two-spin simulations: the calculated 2QCT sideband manifold that most faithfully reproduces the experimental spectrum is located, while accounting for all sideband orders in the iterative fitting.

Despite that truly isolated half-integer spin-pairs are very rare, spin-pair approximations still provide reliable results provided that the targeted dipolar coupling constant is much larger (\(\gtrsim 10\) times) than those of the other couplings in the system: for example, the directly bonded B–B pair in the bis(catecholato)diboron molecule represents a very strong $^{11}$B–$^{11}$B dipolar coupling that is $\simeq 12$ times larger than its second-largest counterpart. The longer-range intermolecular interactions manifest themselves primarily in the $\pm 1$ sideband amplitudes that become overemphasized relative to those of the higher-orders; such effects account for most of the deviations between experimental results and those obtained by spin-pair simulations. Yet, an accurate estimate of the large dipolar coupling is still feasible by simply omitting the lowest-order sidebands in the RMS calculation and only evaluating all higher order ($\pm 3, \pm 5, \ldots$) sideband amplitudes. This effectively amounts to ignoring all weaker dipolar interactions; see Fig. 3 and discussions thereof.

5.2 Probing spin-cluster sizes

We next consider the more common scenario of macroscopically large quadrupolar-spin systems encountered in inorganic materials, which implies the absence of well-defined “clusters”. Here, analysis of the 2QCT experiment via strategy (B) of Section 2.2 provides information about the number of unique spin sites that contributes to the experimental sideband pattern, the number of which increases with the 2QCT preparation period. As may be verified both from the CT 2Q NMR results of the Na$_2$SO$_4$ and $\alpha$-Al$_2$O$_3$ structures in Fig. 6 and 9, there are no 2QCT sideband intensities beyond the lowest ($\pm 1$) order at $\tau_{\text{exc}} \approx \tau_{\text{exc}}^{\text{max}}$, i.e., when the selected excitation period optimizes the total 2QCT signal amplitude. For such cases, the number of contributing spins is necessarily small and also most difficult to determine, because the discrimination between simulations with distinct cut-off radii is poor.

For increasing $\tau_{\text{exc}}$ the best-fit $R$-value grows and thereby also its associated number of contributing spins. The number of spinning sidebands of significant amplitude is proportional to the product of the 2QCT excitation interval and the effective dipolar interaction of the multi-spin system, i.e., a fictive dipolar coupling constant that approximates the net effect from all spin-pairs in the system;\textsuperscript{34–36,72,94–96} it is dictated by the magnitudes and numbers of all contributing dipolar interactions over a certain $R$-value. Hence, for progressively increased excitation intervals, inspection of the sideband intensity-distributions at the corresponding ratios of $\tau_{\text{exc}}/\tau_{\text{exc}}^{\text{max}}$ gives direct qualitative information about the magnitude of the effective dipolar coupling constant of the probed spin network. For increasing $\tau_{\text{exc}}/\tau_{\text{exc}}^{\text{max}}$ ratio, Table 2 reveals that the observed number of interacting spins grows significantly faster in the $\alpha$-Al$_2$O$_3$ structure relative to its Na$_2$SO$_4$ counterpart, in full accordance with the known higher packing density of Al in $\alpha$-Al$_2$O$_3$ compared to that of Na in Na$_2$SO$_4$.

Thus far, the prevailing approach to estimate spin-cluster sizes—“multiple-quantum spin counting”\textsuperscript{57–64}—relies on high-order multiple-quantum coherences (MQC) excitation in stationary or rotating powders comprising multi-spin-1/2 systems. The highest generated MQC order provides a (directly determined) lower bound of $M$ interacting spins present in the “real” cluster. However, only one such attempt at using MQC excitation to count the number of half-integer spins is hitherto reported, then restricted to low-resolution applications in stationary powders.\textsuperscript{64} Our proposed strategy of comparing experimental 2QCT-sideband intensities with approximative spin-pair summations constitutes the first general means for estimating the size of a cluster of half-integers spins under MAS conditions. The spin-pair summation strategy is not necessarily restricted to quadrupolar nuclei,\textsuperscript{56} we are currently exploring its prospects for spin-1/2 applications.

The main limitation of the present spin-pair summation strategy is its time-consuming numerical analysis.\textsuperscript{56} Yet, despite its feature of providing the spin-cluster sizes indirectly through an (approximate) numerical approach, it exhibits a decisive advantage relative to MQ spin counting in that it avoids the cumbersome high-order MQC excitation. Also noteworthy is that techniques for producing high-order multiple-quantum coherences solely involving the central transitions of half-integer spins remain to be developed. The current spin-1/2 MQC excitation strategy is also prone to underestimating the size of large spin-systems due to the impact of NMR relaxation and pulse imperfections, particularly in rotating solids where the progressive high-order MQC generation is slow. Furthermore, also MQ spin-counting data-analyses often resorts to indirect
estimates of the “real” spin-cluster size by fitting the experimental MQC amplitude envelope to a Gaussian decay.\(^{57,58}\) This procedure may be associated with non-negligible uncertainties, as discussed further in ref. 58 and 63.

6 Conclusions

In this contribution, we consolidated the application of CT 2Q sideband NMR spectroscopy for estimating internuclear distances within (nearly) isolated pairs of half-integer spin quadrupolar nuclei, and demonstrated that this experimental approach may be combined with numerical analyses to probe multiple-spin effects and estimate the sizes of local clusters of mutually dipolar-coupled quadrupolar nuclei in network structures.

All these aspects were first explored for the case of bis-(catecholato)diboron, where each molecule involves one pair of directly bonded boron atoms, while the internuclear B–B distances are sufficiently long to ensure a nearly isolated \(^{11}B–^{11}B\) spin pair in each molecule. Since the largest coupling is an order of magnitude stronger than its second largest counterpart, we demonstrated the feasibility of determining its magnitude by fitting results from two-spin simulations employing variable dipolar couplings to the experimental results. Considering only the higher odd-order (±3, ±5, ...) CT sidebands in the numerical fitting, the B–B internuclear distance was determined as \((168.6 ± 3)\) pm, i.e., with an uncertainty of ±2%. However, the strong amplitudes of the lowest order (±1) sidebands cannot be reproduced by single spin-pair simulations; they were shown to stem from the longer range internuclear \(^{11}B–^{11}B\) dipolar interactions. We presented two approaches to account for their contributions, both of which successfully reproduced the entire experimental manifold of 2Q\(_{\text{CT}}\) sidebands: (i) performing computationally expensive multi-spin simulations, or (ii) apply the spin-pair summation approach in ref. 56, in which the complete calculated NMR spectrum is obtained by summation of individual simulations for spin-pairs found within a certain cut-off distance.

Besides justifying the approximative numerical spin-pair simulation protocol by rigorous multi-spin \(^{11}B–^{11}B\) simulations for the case of bis-(catecholato)diboron, we furthermore successfully demonstrated it by \(^{23}\)Na and \(^{27}\)Al NMR on the macroscopic networks of Na\(_2\)SO\(_4\) and Al\(_2\)O\(_3\), respectively. For progressively lengthened 2Q\(_{\text{CT}}\) excitation intervals, each respective \(^{23}\)Na and \(^{27}\)Al cluster size may be extracted by locating its corresponding cut-off distance, thereby allowing the monitoring of a growing number of interacting nuclear spins. Altogether, this analysis yields similar information as the MQ spin-counting experiment,\(^{57-64}\) but without the cumbersome need to excite high-order MQC, which for MAS applications is hitherto only demonstrated for spin-1/2 systems.\(^{59-63}\) However, we emphasize that the MQ spin-counting technique provides direct estimates of spin-cluster sizes, whereas the 2Q\(_{\text{CT}}\) sideband experiment requires a priori knowledge about the atomic coordinates of the structure for determining the number of interacting spins for the given 2Q\(_{\text{CT}}\) excitation interval.

The approach of analyzing experimental CT 2Q sideband data by comparison with sums of numerically simulated spin-pair responses may in its present form only be used for discriminating between various proposed structural models through the adherence of their corresponding numerically simulated sideband patterns to the experimental results. Nevertheless, this protocol makes an important dent towards elucidating the dependence of the quadrupolar-spin dynamics on a potentially wide range of structural parameters, such as the positions and local geometries of the quadrupolar spin-sites. As already demonstrated in the context of 2Q–1Q correlation spectroscopy, a spin-pair summation analysis allows for determining the absolute orientations of the electric field gradient (EFG) tensor at each quadrupolar spin-site,\(^{56}\) hence, the encoding of distance-information in the sidebands from the closely related CT 2Q sideband protocol may be linked to the EFG tensor orientations of the recoupled spins.

Our method prepares the grounds for “ab initio distance-measurements” by numerical fitting of experimental CT 2Q spinning sideband data from disordered structures, where direct information about the coordinates of the quadrupolar sites is intractable by other means; for many structures, this technique may constitute the sole option for obtaining (semi-) quantitative information about internuclear distances and spin-cluster sizes. The strategy of analyzing results from 2Q sideband experiments with a spin-pair summation numerical protocol is generally applicable: we are currently exploring it for spins-1/2 applications to oxide-based inorganic network structures.

Acknowledgements

The authors would like to thank Igor Moudrakovski, Victor Terskikh and Chris Ratcliffe for helpful discussions. Access to CASTEP was provided by the National Ultrahigh-Field NMR Facility for Solids (Ottawa, Canada), a national research facility funded by the Canada Foundation for Innovation, the Ontario Innovation Trust, Recherche Québec, the National Research Council Canada, and Bruker BioSpin and managed by the University of Ottawa (http://www.nmr900.ca). This work was supported by the National Research Council Canada (NRC), the Swedish Research Council (contract VR-NT 2010-4943) and the Faculty of Sciences at Stockholm University.

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