

# Structure, ligands and substrate coordination of the oxygen-evolving complex of photosystem II in the $S_2$ state: a combined EPR and DFT study†

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The  $S_2$  state of the oxygen-evolving complex of photosystem II, which consists of a  $Mn_4O_5Ca$  cofactor, is EPR-active, typically displaying a multiline signal, which arises from a ground spin state of total spin  $S_T = 1/2$ . The precise appearance of the signal varies amongst different photosynthetic species, preparation and solvent conditions/compositions. Over the past five years, using the model species *Thermosynechococcus elongatus*, we have examined modifications that induce changes in the multiline signal, i.e.  $Ca^{2+}/Sr^{2+}$ -substitution and the binding of ammonia, to ascertain how structural perturbations of the cluster are reflected in its magnetic/electronic properties. This refined analysis, which now includes high-field (W-band) data, demonstrates that the electronic structure of the  $S_2$  state is essentially invariant to these modifications. This assessment is based on spectroscopies that examine the metal centres themselves (EPR,  $^{55}Mn$ -ENDOR) and their first coordination sphere ligands ( $^{14}N/^{15}N$ - and  $^{17}O$ -ESEEM, -HYSCORE and -EDNMR). In addition, extended quantum mechanical models from broken-symmetry DFT now reproduce all EPR,  $^{55}Mn$  and  $^{14}N$  experimental magnetic observables, with the inclusion of second coordination sphere ligands being crucial for accurately describing the interaction of  $NH_3$  with the Mn tetramer. These results support a mechanism of multiline heterogeneity reported for species differences and the effect of methanol [*Biochim. Biophys. Acta, Bioenerg.*, 2011, **1807**, 829], involving small changes in the magnetic connectivity of the solvent accessible outer  $Mn_{A4}$  to the cuboidal unit  $Mn_3O_3Ca$ , resulting in predictable changes of the measured effective  $^{55}Mn$  hyperfine tensors.  $Sr^{2+}$  and  $NH_3$  replacement both affect the observed  $^{17}O$ -EDNMR signal envelope supporting the assignment of O5 as the exchangeable  $\mu$ -oxo bridge and it acting as the first site of substrate inclusion.

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## 1 Introduction

In oxygenic photosynthesis, light-driven water oxidation is catalysed by the oxygen-evolving complex (OEC) of the transmembrane protein complex photosystem II (PSII). The OEC

consists of a  $\mu$ -oxo-bridged tetramanganese-calcium cofactor ( $Mn_4O_5Ca$ ), embedded in its protein matrix. This matrix includes the redox-active tyrosine residue  $Y_Z$  (D1-Tyr161), which couples electron transfer between the  $Mn_4O_5Ca$  cluster and  $P680^{*+}$ , the photooxidant of the PSII reaction centre. The four-electron water oxidation reaction requires four consecutive light-induced charge separation events, driving the catalytic cycle of the OEC.<sup>1,2</sup> This cycle involves five redox intermediates, the  $S_n$  states, where  $n = 0-4$  indicates the number of stored oxidizing equivalents. All S-state transitions represent oxidations of the  $Mn_4O_5Ca$  cluster by  $Y_Z^*$  with the exception of the regeneration of  $S_0$  from  $S_4$ , which proceeds spontaneously under the release of molecular triplet oxygen and the rebinding of at least one substrate water molecule. The rate-limiting step, oxidation of  $S_3$  by  $Y_Z^*$ , has prevented the transient, fast-decaying  $S_4$  state from being characterized yet. For a general introduction into water oxidation by the OEC, see ref. 3–8.

The structure of the  $Mn_4O_5Ca$  inorganic core<sup>9</sup> resembles a 'distorted chair' where the base is formed by a  $\mu$ -oxo-bridged

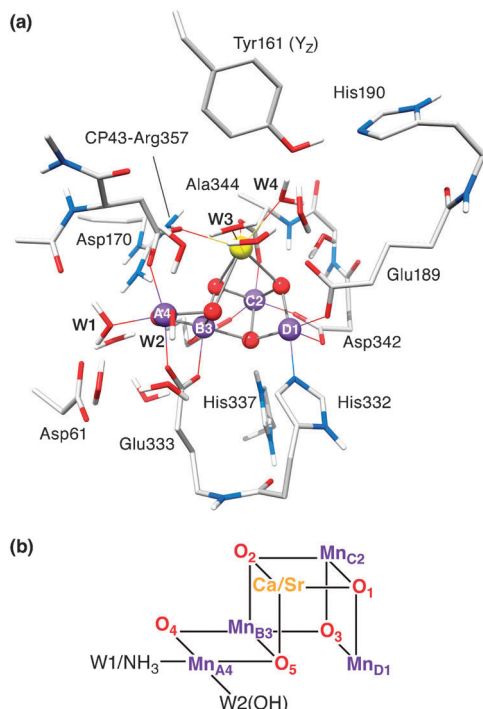
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† Electronic supplementary information (ESI) available: PSII sample preparation; data processing: baseline correction and light-minus-dark subtraction; spectral simulations; theoretical background; DFT calculations:  $NH_3$  binding modes and spatial coordinates of the  $S_2$  state models; Q-band  $^{55}Mn$ -ENDOR: comparison of instrumental settings; electronic structures: exchange couplings and spin states of the BS-DFT models; the  $Mn_{D1}$ -His332-imino-N interaction: Q-band three-pulse ESEEM, Q-band HYSCORE, W-band EDNMR and simulations; X-band  $^{14}N$ - and  $^{15}N$ -three pulse ESEEM experiments on the  $NH_3$ -modified  $S_2$  state; exchangeable  $^{17}O$  species: W-band EDNMR and X-band EPR. See DOI: 10.1039/c3cp55017f





**Fig. 1** (a) DFT cluster model of the native OEC and its immediate surroundings in the  $S_2$  ( $S_T = 1/2$ ) state. Mn ions are depicted in purple, O in red, Ca in yellow, C in light grey, N in blue and H in white; most H atoms are omitted for clarity. (b)  $Mn_4O_5Ca/Sr-H_2O/NH_3$  cores showing the atom and ligand labels used in this work.

cuboidal  $Mn_3O_4Ca$  unit (Fig. 1), while the fourth, ‘outer’ Mn,  $Mn_{A4}$  is attached to this core structure *via* two  $\mu$ -oxo bridges, O4 and O5. Spectroscopic evidence suggests that all five oxygen bridges are deprotonated in the  $S_2$  state.<sup>10–15</sup> In addition to the  $\mu$ -oxo-bridged network, the  $Mn_4O_5Ca$  scaffold is held together by six carboxylate ligands, four of which (D1-Asp342, D1-Asp170, D1-Glu333 and CP47-Glu354) form  $\mu$ -carboxylato bridges between Mn sites. The cluster has only one immediate nitrogen ligand, the imidazole imino-N of D1-His332.

An understanding of the mechanism of water oxidation catalysis is intimately linked to an understanding of the electronic states of the paramagnetic tetranuclear Mn complex. Electron paramagnetic resonance (EPR) spectroscopy represents a powerful methodology in this endeavour.<sup>16–19</sup> The  $S_2$  state, being readily generated from the dark-stable  $S_1$  state, is the most extensively characterized S state by EPR. It is intrinsically heterogeneous,<sup>20–22</sup> existing in two ground state configurations. The lowest energy form exhibits a ground state of total spin  $S_T = 1/2$ . This spin configuration gives rise to the well-known  $S_2$  state EPR multiline signal centred at  $g \approx 2.0$ , with a hyperfine pattern of at least 24 peaks, with a line spacing around 87 G.<sup>23</sup> Depending on the conditions used, the  $S_2$  state can also exhibit other broad EPR signals, centred at  $g \geq 4.1$ , which have been assigned to (an)  $S_T \geq 5/2$  spin state(s).<sup>20,24–28</sup> The  $g \geq 4.1$  signal can also be induced by near-infrared illumination of the  $S_2$  multiline state at temperatures  $\leq 160$  K.<sup>24,25,29</sup> Pulse electron nuclear double resonance (ENDOR) experiments,<sup>30</sup> especially at Q-band

frequencies,<sup>31–34</sup> have enabled a more detailed examination of the electronic structure by uniquely determining four  $^{55}Mn$  hyperfine interactions that give rise to the multiline pattern of the corresponding  $S_T = 1/2$  EPR signal. This ENDOR analysis has strongly suggested that the OEC contains one  $Mn^{III}$  ion and three  $Mn^{IV}$  ions when poised in the  $S_2$  state.<sup>30–34</sup>

OEC models developed from broken-symmetry (BS) density functional theory (DFT) have been crucial for the interpretation of EPR and related magnetic resonance results.<sup>12,14,18,33,35–44</sup> These calculations now allow different S-state models to be assessed based on the spin multiplicities and relative energies of their lowest magnetic levels, and, by means of the spin projection formalism (see ref. 18 and 36), the metal and ligand hyperfine couplings. This strategy enables the assignment of individual Mn oxidation states and coordination geometries and represents a method to develop unified models of the OEC that combine constraints from X-ray diffraction, EXAFS and magnetic resonance.

Site perturbation of the OEC provides a means to further characterize the global geometric and electronic structure of the Mn tetramer and obtain information about localized changes associated with the introduced modification. The two best described OEC perturbations in the literature are (i) the replacement of the  $Ca^{2+}$  with a  $Sr^{2+}$  ion<sup>33,45–47</sup> and (ii) the binding of  $NH_3$  to the cluster.<sup>48–55</sup> Both of these perturbations *do not* inhibit functional turnover of the catalyst, but do modify the kinetics of O–O bond formation, substrate water uptake and its subsequent deprotonation. A brief description of these two modified OEC forms is given below:

(i)  $Sr^{2+}$  can be introduced into the OEC through chemical removal of the  $Ca^{2+}$  ion<sup>45,46,56</sup> or alternatively by biosynthetic incorporation in cyanobacterial cultures grown in the presence of  $SrCl_2$  instead of  $CaCl_2$ .<sup>47</sup> Without  $Ca^{2+}$  (or  $Sr^{2+}$ ) the OEC is inhibited,<sup>56–63</sup> not progressing further than a modified form of the  $S_2$  state, *i.e.* the  $S_2/Y_Z^*$  state.<sup>60–62</sup>  $Sr^{2+}$  is unique as it is the only ion which can replace the  $Ca^{2+}$  ion while retaining catalytic activity, albeit at a lower enzymatic rate.<sup>45,47,56</sup> Presumably this is because  $Sr^{2+}$  has a similar size and Lewis acidity as  $Ca^{2+}$ .<sup>64</sup> While slowing the turnover rate of the catalyst,  $Sr^{2+}$  substitution at the same time enhances the exchange rate with bulk water of at least one of the bound substrates,<sup>65,66</sup> as observed by time-resolved membrane inlet mass spectrometry (MIMS).<sup>67,68</sup> This behaviour suggests the  $Ca^{2+}$  ion may play an important role in substrate water binding and possibly proton release (for reviews, see ref. 3 and 69).

(ii) Ammonia binding to the Mn cluster (in the presence of high  $Cl^-$  concentrations)<sup>48,50,51,54</sup> only occurs upon formation of the  $S_2$  state. It is subsequently released at some point during the S-state cycle ( $S_3 \rightarrow S_0 \rightarrow S_1$ ), such that it is not bound upon return to  $S_1$ .<sup>51</sup> As with  $Sr^{2+}$  replacement,  $NH_3$  binding does not inhibit catalytic function. In the higher plant electron spin echo envelope modulation (ESEEM) study of Britt *et al.*,<sup>70</sup>  $NH_3$  was shown to bind as a direct ligand of the Mn tetramer. The precise binding site and coordination mode of the  $NH_3$  molecule was the subject of a recent study on cyanobacterial PSII from our laboratory.<sup>44</sup> By employing electron electron double resonance (ELDOR)-detected NMR (EDNMR), it was concluded that  $NH_3$



replaces the water ligand of  $\text{Mn}_{\text{A4}}$  *trans* to the O5 bridge (W1, Fig. 1b). As the binding of  $\text{NH}_3$  was also shown not to affect substrate exchange rates, these results suggest W1 does not represent a substrate water. One or more additional  $\text{NH}_3$  binding sites, which are inhibitory, are known but are less well characterized.<sup>49,50,71</sup>

Here, we present an extension of our earlier multifrequency EPR studies<sup>33,34,44</sup> of different ‘archetypal’ multiline forms, namely the native ( $\text{Mn}_4\text{O}_5\text{Ca}$ ),  $\text{Sr}^{2+}$  substituted ( $\text{Mn}_4\text{O}_5\text{Sr}$ ) and  $\text{NH}_3$ -treated ( $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$ )  $\text{S}_2$  states, providing a comprehensive analysis of all present X-, Q- and W-band data of  $^{55}\text{Mn}$ ,  $^{14}\text{N}/^{15}\text{N}$  and  $^{17}\text{O}$  signals and for the first time including additional data on the doubly modified ( $\text{Mn}_4\text{O}_5\text{Sr-NH}_3$ )  $\text{S}_2$  state. Improved  $^{55}\text{Mn}$  ENDOR experimental conditions provide more reliable spectral line shapes than before, confirming our previous general  $\text{S}_2$  state model. Extended BS-DFT models are shown to reproduce all EPR,  $^{55}\text{Mn}$  and  $^{14}\text{N}$  magnetic spectroscopic observables for the native and the modified systems, a feature not achieved previously. The experimental results and calculations for  $^{14}\text{N}/^{15}\text{N}$  ligands of the various  $\text{S}_2$  state forms serve to prove that the basic electronic structure is not perturbed by these modifications, a result crucial for the interpretation of concomitant perturbations of the  $^{17}\text{O}$  EDNMR signal envelope. This combined experimental and theoretical approach supports our qualitative model for multiline heterogeneity, demonstrating that the magnetic connectivity between the two subunits and also within the trimeric moiety governs the structure of the multiline signal.<sup>34</sup> This basic structural template also explains the apparent orientations of the  $^{55}\text{Mn}$  hyperfine tensors, as inferred from spectral simulations and single crystal measurements,<sup>32</sup> and potentially provides a framework to further examine substrate binding. The different OEC forms represent a starting point to examine the energetics of higher S-state transitions, as they differ with regard to substrate binding and the kinetics of O–O bond formation and  $\text{O}_2$  release.

## 2 Materials and methods

### 2.1 PSII sample preparation

$\text{Ca}^{2+}$ - and  $\text{Sr}^{2+}$ -containing PSII core complex preparations from *T. elongatus*<sup>72</sup> were isolated as reported before<sup>47,73,74</sup> with the same modifications for the X-band samples as described in ref. 44. Universal  $^{15}\text{N}$ -labelling of the PSII proteins was achieved by growing the cyanobacteria in a modified BG11 or DTN medium that contained  $^{15}\text{NH}_4\text{Cl}$  and  $^{15}\text{NO}_3^-$  salts as the sole nitrogen source.<sup>75</sup> PSII preparations were stored at  $-80^\circ\text{C}$  until use. Dark-adapted samples were placed in X-, Q- or W-band quartz tubes with inner diameters of 3.0, 1.6 and 0.6 mm, respectively, and kept at 77 K (liquid  $\text{N}_2$ ) until use. A sample concentration of 3.0–4.0 mg chlorophyll per ml was used throughout this study. All work was conducted in the dark or under dim green light.  $\text{NH}_3$  modification was conducted as described in ref. 44, see also Section S1 of the ESI.† Isotopically enriched  $\text{H}_2^{17}\text{O}$  buffer exchange was achieved as described in ref. 14.

### 2.2 X-, Q-, and W-band EPR and ENDOR measurements

X-band ( $\approx 9$  GHz) continuous-wave (CW) EPR spectra were recorded at liquid He temperatures on a Bruker ELEXSYS E500 spectrometer, equipped with an ESR 900 liquid helium flow cryostat and an ITC503 helium flow temperature controller (Oxford Instruments Ltd). X-band pulse EPR measurements were carried out at 4.3 K using a Bruker ELEXSYS E580 spectrometer, equipped with a CF935 cryostat and an ITC503 temperature controller. Q-band ( $\approx 34$  GHz) pulse EPR measurements were performed around 5 K using an ELEXSYS E580 spectrometer, equipped with a homebuilt  $\text{TE}_{011}$  microwave cavity,<sup>76</sup> a CF935 liquid helium cryostat, an ITC-503 temperature controller and a radiofrequency (RF) amplifier ENI 5100L. W-band ( $\approx 94$  GHz) EPR experiments were performed at 4.8 K using a Bruker ELEXSYS E680 EPR spectrometer. All W-band experiments were carried out using a homebuilt ENDOR microwave cavity, which contained a solenoid of Teflon coated silver wire integrated into a commercial ENDOR probe head (Bruker). In order to ensure broadband microwave excitation and minimize distortions, the loaded quality factor  $Q_L$  was lowered to 700 to obtain a microwave frequency bandwidth of 130 MHz.

Electron spin echo-detected (ESE) field-swept EPR spectra were measured using the pulse sequence  $t_p-\tau-2t_p-\tau$ -echo,<sup>77</sup> three-pulse ESEEM spectra by use of  $t_p-\tau-t_p-T-t_p-\tau$ -echo<sup>78</sup> and hyperfine sublevel correlation (HYSCORE) spectra by employing  $t_p-\tau-t_p-T_1-2t_p-T_2-t_p-\tau$ -echo.<sup>79</sup> The lengths of the  $\pi/2$  microwave pulses were generally set to  $t_p = 16$  ns (X-band), 12 ns (Q-band) and 24 ns (W-band), respectively. For ESE-detected EPR experiments, inter-pulse distances were  $\tau = 260$  ns (Q-band) and 300 ns (W-band). For the three-pulse ESEEM measurements, multiple  $\tau$  values in the ranges  $\tau = 136$ –248 ns (X-band) and 200–356 ns (Q-band) and an optimum  $\tau = 260$  ns for the HYSCORE experiments were chosen to account for blind-spotting artefacts. Q-band  $^{55}\text{Mn}$ -ENDOR spectra were acquired employing the Davies-type pulse sequence  $t_{\text{inv}}-t_{\text{RF}}-T-t_p-\tau-2t_p-\tau$ -echo<sup>80</sup> using a length  $t_{\text{inv}} = 24$  ns for the  $\pi$  inversion microwave pulse and a radio frequency  $\pi$  pulse length  $t_{\text{RF}} = 3.5$   $\mu\text{s}$ . The length of the  $\pi/2$  microwave pulse in the detection sequence was generally set to  $t_p = 12$  ns and the inter-pulse delays to  $T = 2$   $\mu\text{s}$  and  $\tau = 268$  ns. A shot repetition time of 1 ms was used for all measurements. EDNMR measurements were collected using the pulse sequence  $t_{\text{HTA}}-T-t_p-\tau-2t_p-\tau$ -echo.<sup>81</sup> The high turning angle (HTA) microwave pulse was applied at microwave frequencies  $\nu_{\text{mw}}$ . The Hahn echo detection pulse sequence  $t_p-\tau-2t_p-\tau$ -echo, at a microwave frequency  $\nu_{\text{mw}}^{(0)}$  matched to the cavity resonance, was set at a sufficient time  $T$  after the HTA pulse to ensure near-complete decay of the electron spin coherencies. The  $\pi/2$  pulse length used for detection was  $t_p = 200$  ns ( $^{14}\text{N}$ ,  $^{17}\text{O}$ ) or 80 ns ( $^{15}\text{N}$ ) and an inter-pulse separation of  $\tau = 500$  ns was used. The echo was integrated  $\approx 600$  ns around its maximum. The spectra were acquired *via* continuously sweeping the HTA frequency  $\nu_{\text{mw}}$  at a fixed magnetic field in steps of 78.1 kHz ( $^{14}\text{N}$ ), 128.9 kHz ( $^{15}\text{N}$ ) or 162.1 kHz ( $^{17}\text{O}$ ). A HTA microwave pulse of length  $t_{\text{HTA}} = 14$   $\mu\text{s}$  ( $^{14}\text{N}$ ,  $^{17}\text{O}$ ) and 8  $\mu\text{s}$  ( $^{15}\text{N}$ ) and an amplitude  $\omega_1 = 12$ –16  $\times 10^6$   $\text{rad s}^{-1}$  was used.



### 2.3 Spectral simulations

Spectra were fit assuming an effective spin  $S_T = 1/2$  ground state (Section S4.2, ESI†). The basis set that describes the  $^{55}\text{Mn}$  tetramer-single electron spin manifold (eqn (1)) and the  $^{14}\text{N}$ ,  $^{15}\text{N}$  and  $^{17}\text{O}$  single nucleus-single electron spin manifolds (eqn (2)) can be built from the product of the eigenstates of the interacting spins:

$$|\frac{1}{2}MI_1m_1I_2m_2I_3m_3I_4m_4\rangle \quad (1)$$

$$|\frac{1}{2}MIm\rangle \quad (2)$$

Here,  $M$  refers to the electronic magnetic sublevel,  $\pm\frac{1}{2}$ ;  $I$  takes the values  $5/2$  for  $^{55}\text{Mn}$ ,  $1$  for  $^{14}\text{N}$ ,  $\frac{1}{2}$  for  $^{15}\text{N}$  and  $5/2$  for  $^{17}\text{O}$ ;  $m_i$  takes the values  $-I_i, 1 - I_i, \dots, I_i - 1, I_i$ . The spin manifolds can be described by the following spin Hamiltonian:

$$\hat{H} = \beta_e \vec{B}_0 \cdot \hat{G} \cdot \vec{S} + \sum_i \left( -g_{n,i} \beta_n \vec{B}_0 \cdot \vec{I}_i + \vec{S} \cdot \hat{A}_i \cdot \vec{I}_i + \vec{I}_i \cdot \hat{Q}_i \cdot \vec{I}_i \right) \quad (3)$$

It contains (i) the Zeeman term for the total electronic spin, (ii) the hyperfine and (iii), except for the EPR spectra, nuclear Zeeman terms for either the metal  $^{55}\text{Mn}$  nuclei, or the ligand  $^{14}\text{N}$ ,  $^{15}\text{N}$  or  $^{17}\text{O}$  nuclei and (iv), for the  $^{14}\text{N}$  nuclei ( $I = 1$ ), the nuclear quadrupole interaction (NQI) term; the NQI splitting is not resolved in the  $^{55}\text{Mn}$  and  $^{17}\text{O}$  spectra. Spectral simulations were performed numerically using MATLAB® (R2010a, The MathWorks, Natick, MA, USA), a vector-based linear algebra package, and the EasySpin toolbox.<sup>82</sup> For further information on data processing, details of the simulations and theory, see Sections S2, S3 and S4 (ESI†), respectively.

### 2.4 DFT calculations

All calculations were performed with ORCA.<sup>83</sup> The DFT models of the OEC systems consist of 238 or 239 atoms (with  $\text{H}_2\text{O}$  or  $\text{NH}_3$  at the W1 position, respectively) and were constructed as described in Pantazis *et al.*<sup>40</sup> Alternative ammonia binding modes, including terminal or bridging amido and imido substitution, can be rejected on energetic grounds alone (see Fig. S2, ESI†). Geometry optimizations of the cluster models used the BP86 density functional<sup>84,85</sup> with the zeroth-order regular approximation (ZORA)<sup>86–88</sup> and specially adapted segmented all-electron relativistically contracted basis sets<sup>89</sup> (ZORA-TZVP for Mn, O and N; ZORA-SVP for C and H). The conductor-like screening model (COSMO)<sup>90</sup> was used with a dielectric constant  $\epsilon = 8.0$ . D3 dispersion corrections<sup>91</sup> were applied in all calculations. Exchange coupling constants, assuming an isotropic Heisenberg Hamiltonian of the form

$$\hat{H} = -2J_{ij} \sum_{i < j} \hat{S}_i \cdot \hat{S}_j, \quad (4)$$

and all EPR properties were computed with the TPSSH hybrid meta-GGA functional<sup>92,93</sup> from BS-DFT calculations.<sup>35,94–98</sup> The resolution of identity (RI)<sup>99</sup> approximation was used in the calculation of Coulomb integrals and the chain-of-spheres approximation (COSX)<sup>100</sup> was used for Hartree–Fock exchange, employing completely decontracted def2-TZVP/J auxiliary basis sets.<sup>101</sup> Tight SCF

convergence criteria and increased integration grids (Grid6 and GridX6) were applied throughout.

For the calculation of the hyperfine tensors, triple-zeta ZORA-recontracted basis sets were used on all atoms, while locally dense radial grids were used for Mn, N and O atoms (integration accuracy of 11 for Mn and 9 for N and O as per ORCA nomenclature). Picture change effects were applied for the calculation of EPR parameters and the complete mean-field approach was used for the spin–orbit coupling operator. The results were transformed into on-site or spin-projected values as detailed previously.<sup>36</sup> To compare computed  $^{55}\text{Mn}$  hyperfine coupling constants using the methods described above with experimental results, a scaling factor of 1.78 was calculated from a set of twelve  $\text{Mn}^{\text{III}}\text{Mn}^{\text{IV}}$  mixed-valence dimers.<sup>102</sup>

## 3 Results and discussion

### 3.1 DFT models of different OEC forms in the $S_2$ state

Geometric parameters of optimized DFT cluster models of the  $S_2$  state of the OEC in the  $S_T = 1/2$  configuration<sup>40</sup> are shown in Fig. 2 (for coordinates, see Section S5, ESI†). Four variants were considered in this study: (i) the native cofactor system ( $\text{Mn}_4\text{O}_5\text{Ca}$ , also see Fig. 1), (ii) the  $\text{Sr}^{2+}$ -substituted system obtained by replacing  $\text{Ca}^{2+}$  with  $\text{Sr}^{2+}$  ( $\text{Mn}_4\text{O}_5\text{Sr}$ ), (iii) the  $\text{NH}_3$ -modified system obtained by replacing the  $\text{H}_2\text{O}$  in the W1 position with  $\text{NH}_3$  ( $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$ ), and (iv) the combined  $\text{Sr}^{2+}$ -substituted and  $\text{NH}_3$ -modified system ( $\text{Mn}_4\text{O}_5\text{Sr-NH}_3$ ). In all models, W2 was considered to be an  $\text{OH}^-$  ligand, as determined previously.<sup>12</sup> Mulliken spin population analysis of all four variants confirms that the only  $\text{Mn}^{\text{III}}$  ion of the tetramanganese complex is  $\text{Mn}_{\text{D1}}$ . The three  $\text{Mn}^{\text{IV}}$  ions ( $\text{Mn}_{\text{A4}}$ ,  $\text{Mn}_{\text{B3}}$  and  $\text{Mn}_{\text{C2}}$ ) represent coordinatively saturated, 6-coordinate octahedral sites, whereas the  $\text{Mn}_{\text{D1}}^{\text{III}}$  is 5-coordinate square-pyramidal, with a Jahn–Teller elongation along the axis of the  $\text{Mn}_{\text{D1}}\text{-Asp342}$  carboxylate ligand, opposite to its open coordination site.

In accordance with previous DFT and QM/MM structures,<sup>12,37,38,40–42,103–106</sup> the optimized Mn–Mn and Mn–Ca distances of the  $\text{Mn}_4\text{O}_5\text{Ca}$  model are consistent with those determined from EXAFS spectroscopy.<sup>107–110</sup> Only minor changes are observed between the  $\text{Mn}_4\text{O}_5\text{Ca}$  and the  $\text{Mn}_4\text{O}_5\text{Sr}$  models (Fig. 2). As a result of the larger radius of  $\text{Sr}^{2+}$ , the O–Sr bond lengths increase by 0.04 Å, while the Mn–Sr distances also increase by 0.04 Å except for  $\text{Mn}_{\text{D1}}\text{-Sr}$ , which is 0.03 Å shorter than the  $\text{Mn}_{\text{D1}}\text{-Ca}$  distance. On average, this is in line with observations from EXAFS spectroscopy<sup>108</sup> and with the recent 2.1 Å resolution crystallographic model of  $\text{Sr}^{2+}$ -substituted PSII.<sup>111</sup> The Mn–Mn distances are almost entirely unaffected, with the exception of  $\text{Mn}_{\text{D1}}\text{-Mn}_{\text{C2}}$ , which is shortened by 0.02 Å in the  $\text{Mn}_4\text{O}_5\text{Sr}$  model.

Upon  $\text{NH}_3$  substitution of W1 ( $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$ ), only the  $\text{Mn}_{\text{D1}}\text{-Mn}_{\text{B3}}$  distance and the  $\text{Ca}^{2+}$  distance from the terminal Mn ions change notably, albeit by less than 0.05 Å (Fig. 2). Only one structural element is more significantly perturbed, *i.e.* the position of O5, the  $\mu$ -oxo bridge *trans* to the binding position of  $\text{NH}_3$ . The  $\text{Mn}_{\text{A4}}\text{-O5}$  distance increases by 0.05 Å with concomitant





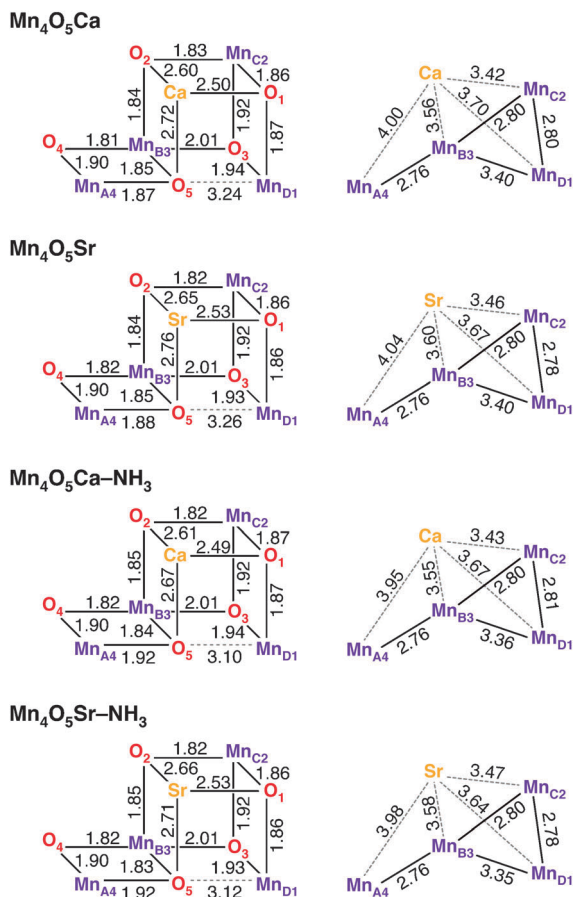


Fig. 2 Bond lengths and metal-metal distances in Å of the four geometry-optimized DFT models.

decrease of the Mn<sub>D1</sub>-O<sub>5</sub> distance by 0.14 Å. Other ligands of Mn<sub>A4</sub>, such as the second water-derived ligand, W2, remain unaffected. A similar modification to the connectivity of the Mn<sub>4</sub>O<sub>5</sub> core was seen for the smaller OEC models reported in our previous study.<sup>44</sup> We note that EXAFS data for the NH<sub>3</sub>-modified OEC have only been reported for samples purified from spinach, not from the cyanobacterial model systems *T. elongatus/vulcanus*, and suggested an elongation of one of the short Mn-Mn distances of 0.02 Å.<sup>113</sup> This type of perturbation is not observed in our optimized Mn<sub>4</sub>O<sub>5</sub>Ca-NH<sub>3</sub> model. The structural model including both Sr<sup>2+</sup> and NH<sub>3</sub> (Mn<sub>4</sub>O<sub>5</sub>Sr-NH<sub>3</sub>) is found to replicate both effects seen in the singly modified structures.

### 3.2 Multifrequency EPR and <sup>55</sup>Mn-ENDOR of the S<sub>2</sub> states of the OEC variants

Multifrequency EPR/<sup>55</sup>Mn-ENDOR experiments spanning the microwave range from  $\approx 9$  to  $\approx 90$  GHz were employed to experimentally characterize the electronic structures of the different S<sub>2</sub> state forms described above. Fig. 3A (black solid traces) depicts X-band CW EPR spectra of the S<sub>2</sub> state of native PSII (Mn<sub>4</sub>O<sub>5</sub>Ca),<sup>33</sup> Sr<sup>2+</sup>-substituted PSII (Mn<sub>4</sub>O<sub>5</sub>Sr),<sup>33</sup> NH<sub>3</sub>-modified (annealed) native PSII (Mn<sub>4</sub>O<sub>5</sub>Ca-NH<sub>3</sub>)<sup>44</sup> and NH<sub>3</sub>-modified Sr<sup>2+</sup>-substituted PSII (Mn<sub>4</sub>O<sub>5</sub>Sr-NH<sub>3</sub>). Shown are light-minus-dark spectra, generated by taking the difference between the

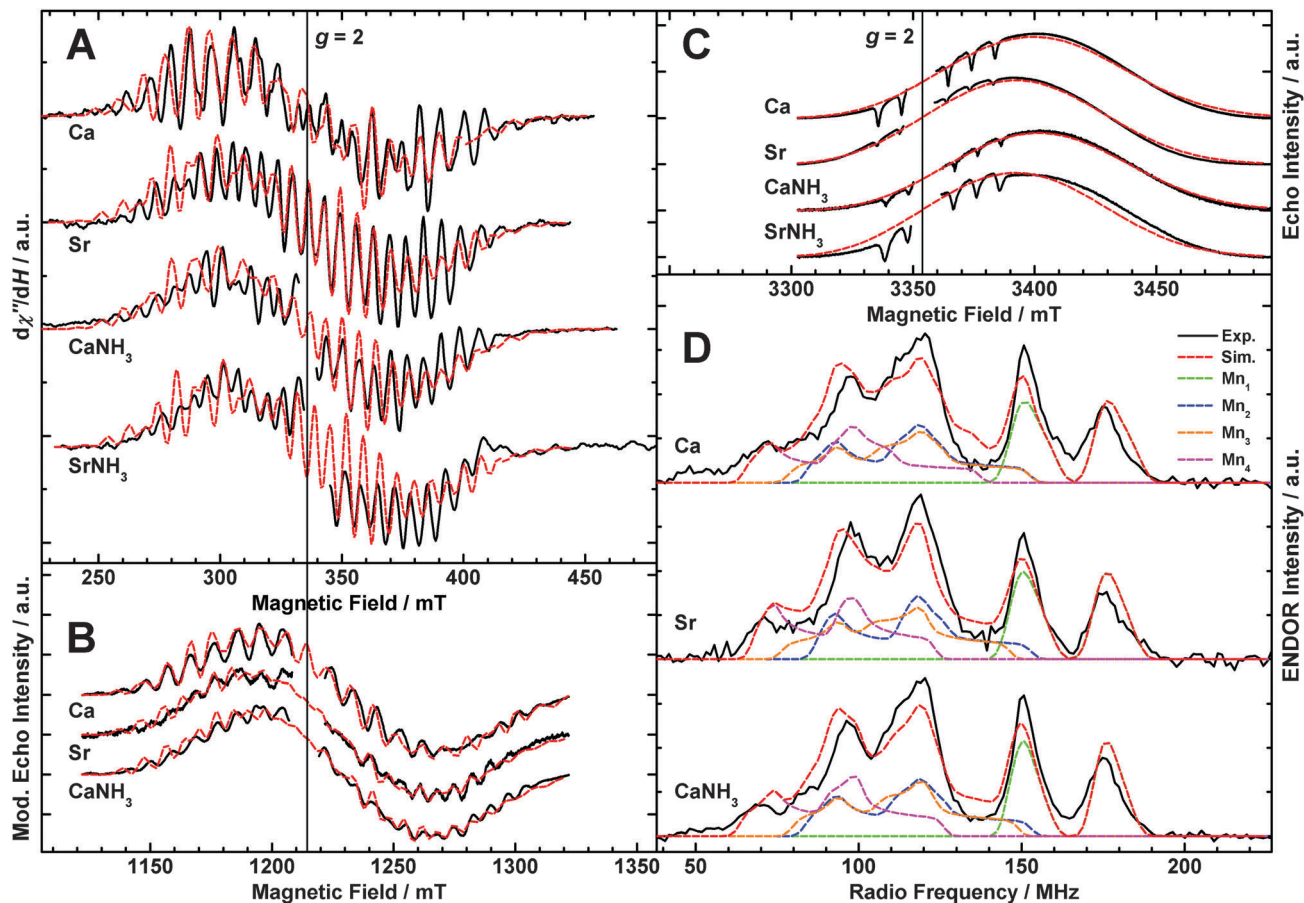
illuminated spectrum (S<sub>2</sub>) and the dark-state spectrum (S<sub>1</sub>) in order to remove background signals, such as from the cytochromes *b*<sub>559</sub> and *c*<sub>550</sub>. The modified multiline (Mn<sub>4</sub>O<sub>5</sub>Sr, Mn<sub>4</sub>O<sub>5</sub>Ca-NH<sub>3</sub>) displays 26 lines of altered intensity as compared to the native multiline signal with 24 lines. The Mn<sub>4</sub>O<sub>5</sub>Sr-NH<sub>3</sub> S<sub>2</sub> state yields essentially the same modified multiline signal; the simulation superimposing this data trace uses the parameters that fit the Mn<sub>4</sub>O<sub>5</sub>Sr dataset.

Fig. 3B shows the corresponding Q-band ESE-detected S<sub>2</sub> state multiline EPR signals of the Mn<sub>4</sub>O<sub>5</sub>Ca, Mn<sub>4</sub>O<sub>5</sub>Sr and Mn<sub>4</sub>O<sub>5</sub>Ca-NH<sub>3</sub> OEC forms. Pseudo-modulated (CW-like) spectra are shown in order to more clearly visualize differences in the hyperfine structures between the three forms. Compared to earlier published data by Cox *et al.*,<sup>33</sup> the spectra are essentially free of contaminating hexaquo-Mn<sup>2+</sup> signals. Furthermore, there is a small difference of the centre positions of the multiline spectral envelopes, presumably due to inaccuracy in the microwave frequency calibration of this earlier study.

Fig. 3C shows the W-band ESE-detected multiline signals of the Mn<sub>4</sub>O<sub>5</sub>Ca, Mn<sub>4</sub>O<sub>5</sub>Sr, Mn<sub>4</sub>O<sub>5</sub>Ca-NH<sub>3</sub> and Mn<sub>4</sub>O<sub>5</sub>Ca-NH<sub>3</sub> S<sub>2</sub> state forms. These spectra do contain small hexaquo-Mn<sup>2+</sup> contaminations. This contamination manifests itself as six inverted hyperfine lines centred at  $g \approx 2$ , as the high-spin Mn<sup>2+</sup> signal ( $S = 5/2$ ) is over-rotated when using optimal instrumental settings to visualize the  $S = 1/2$  multiline signal. No <sup>55</sup>Mn hyperfine structure is observable in the W-band multiline EPR spectra. Thus, these signals provide no additional information on the hyperfine matrices. The utility of these high-frequency data instead is to constrain the *G* tensor. While the spectra of the Mn<sub>4</sub>O<sub>5</sub>Ca and the Mn<sub>4</sub>O<sub>5</sub>Ca-NH<sub>3</sub> cluster show similar signals, centred at  $g = 1.975$ , the Mn<sub>4</sub>O<sub>5</sub>Sr spectrum is shifted to higher field and centred at  $g = 1.980$ , similar to the shift observed in W-band spectra from the Bittl laboratory.<sup>114</sup> The NH<sub>3</sub>-modified Sr<sup>2+</sup>-substituted S<sub>2</sub> state signal is centred at  $g = 1.979$ . The almost identical high-field shift, illustrated by the superimposed Mn<sub>4</sub>O<sub>5</sub>Sr simulation, indicates that the *G*-tensor shift is dependent only on the presence of Sr<sup>2+</sup> but not NH<sub>3</sub>. This is in contrast to the result at X-band, which showed that the hyperfine structure is approximately the same for both modifications.

Fig. 3D shows the Q-band Davies <sup>55</sup>Mn-ENDOR light-minus-dark spectra, measured at the central magnetic field position of the corresponding EPR spectra (Fig. 3B). The line shape of the <sup>55</sup>Mn-ENDOR signal shows only a small field dependence over the 1190 to 1260 mT range (not shown), consistent with its assignment to the tetranuclear Mn cluster.<sup>31</sup> In contrast to the X-band CW EPR spectra described above (Fig. 3A), the S<sub>2</sub> states of the native, Sr<sup>2+</sup>-substituted, NH<sub>3</sub>-annealed and doubly modified (not shown) OEC give rise to highly similar <sup>55</sup>Mn-ENDOR spectra. Five peaks are observed for all three sample types appearing at approximately the same frequency positions and of similar intensities. Small differences in the region of the largest peak ( $\approx 115$  MHz) may represent residual Mn<sup>2+</sup> contaminations. Importantly, no large difference is seen with regard to the total spectral breadth of the signal ( $\approx 55$  to  $\approx 195$  MHz). For the NH<sub>3</sub>-modified S<sub>2</sub> state, the results are nominally consistent with the earlier X-band ENDOR data of Peloquin *et al.*<sup>30</sup>





**Fig. 3** EPR and  $^{55}\text{Mn}$ -ENDOR spectra of the  $S_2$  states of the native (Ca),  $\text{Sr}^{2+}$ -substituted (Sr),  $\text{NH}_3$ -modified (CaNH $_3$ ) and  $\text{NH}_3$ -modified  $\text{Sr}^{2+}$ -substituted (SrNH $_3$ ) OEC in PSII isolated from *T. elongatus*. In panels A, B and D, the black solid traces depict the light-minus-dark subtractions of the experimental spectra. If present, the  $g \approx 2$  radical signal of  $\text{Y}_D^*$  (D2-Tyr160) was removed from the EPR spectra for clarity of presentation. Least-squares fittings to the EPR and  $^{55}\text{Mn}$ -ENDOR datasets using a model based on the spin Hamiltonian formalism (see Section 2.3 and Sections S3 and S4, ESI $^\dagger$ ) are represented by superimposing red dashed lines. In panel D, coloured dashed lines represent a decomposition of the simulation showing contributions from the individual  $^{55}\text{Mn}$  nuclei. Simulations superimposed on the SrNH $_3$  spectra are those fitted to the Sr dataset. The optimized parameter sets are listed in Table 1. (A) X-band CW EPR. In the Ca and Sr samples,  $\text{Y}_D$  had been replaced by a phenylalanine, removing the  $\text{Y}_D^*$  signal from the spectra,<sup>112</sup> which were taken from Cox *et al.*<sup>53</sup> The CaNH $_3$  spectrum was originally published in ref. 44. Experimental parameters: microwave frequencies: 9.4097 GHz (Ca), 9.4213 GHz (Sr), 9.4075 GHz (CaNH $_3$ ), 9.4970 GHz (SrNH $_3$ ); microwave power: 20 mW; modulation amplitude: 25 G; time constant: 82 ms; temperature: 8.6 K. (B) Q-band ESE-detected EPR. The experimental data are presented as pseudo-modulated, derivative-shaped spectra. Experimental parameters: microwave frequencies: 34.0368 GHz (Ca), 34.0430 GHz (Sr), 34.0162 GHz (CaNH $_3$ ); shot repetition time: 1 ms; microwave pulse length ( $\pi$ ): 24 ns;  $\tau$ : 260 ns; temperature: 5.2 K. (C) W-band ESE-detected EPR. Contaminating  $\text{Mn}^{2+}$ , present in the samples in varying concentrations, is evident as over-rotated hyperfine features of negative signal intensity. Experimental parameters: microwave frequencies: 93.9894 GHz (Ca), 93.9781 GHz (Sr), 94.0669 GHz (CaNH $_3$ ), 94.0615 GHz (SrNH $_3$ ); shot repetition time: 1 ms (Ca, Sr), 5 ms (CaNH $_3$ , SrNH $_3$ ); microwave pulse length ( $\pi$ ): 48 ns;  $\tau$ : 300 ns; temperature: 4.8 K. (D) Q-band Davies ENDOR. Experimental parameters: microwave frequencies: 33.9678 GHz (Ca), 33.9950 GHz (Sr), 34.0053 GHz (CaNH $_3$ ); magnetic field: 1220 mT; shot repetition time: 1 ms; microwave pulse length ( $\pi$ ): 32 ns;  $\tau$ : 268 ns; RF pulse length ( $\pi_{\text{RF}}$ ): 3.5  $\mu\text{s}$ ; temperature: 4.8 K.

The Q-band ENDOR spectra presented here do slightly differ from those presented in ref. 33, 34, 44 and 115 with regard to line intensities, discussed in detail in the Section S6 (ESI $^\dagger$ ).

Spectral simulations of the complete EPR and  $^{55}\text{Mn}$ -ENDOR datasets using the spin Hamiltonian formalism are also shown in Fig. 3 (red dashed lines); the fitted effective  $G$  and hyperfine tensors  $A$  of the Mn clusters in the low-spin  $S_2$  state are listed in Table 1. The effective  $G$  tensors of all three spin systems are in the same range as inferred from EPR measurements on PSII single crystals at W-band, *i.e.*  $G = [1.997, 1.970, 1.965]$ <sup>116</sup> and  $G = [1.988, 1.981, 1.965]$ .<sup>117</sup> As found previously,<sup>30,31,33,34,44</sup> the inclusion of four hyperfine tensors of approximately the

same magnitude and near-axial symmetry is required to simultaneously fit the X-, Q- and W-band EPR and Q-band  $^{55}\text{Mn}$ -ENDOR line widths and shapes. The  $z$  component represents the principal component for the fitted  $G$  and all four hyperfine tensors. Comparison of the fitted parameters demonstrates that the three samples basically exhibit the same electronic structure. The sets of the four isotropic values  $A_{i,\text{iso}}$  deviate only by  $\leq 4\%$  between the three different systems and the signs and magnitudes of the anisotropies  $A_{i,\text{aniso}}$  are broadly similar, suggesting that there are no significant differences in the electronic exchange coupling schemes of the  $\text{Mn}_4\text{O}_5\text{Ca}/\text{Sr}(\text{NH}_3)$  clusters.



**Table 1** Principal values of the effective  $G$  and  $^{55}\text{Mn}$  hyperfine tensors  $A_i$  for the simulations of the  $S_2$  state spectra of the  $\text{Mn}_4\text{O}_5\text{Ca}$ ,  $\text{Mn}_4\text{O}_5\text{Sr}$  and  $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$  clusters in PSII from *T. elongatus*<sup>a</sup>

		$G$	$A_i/\text{MHz}$			
			$A_1$	$A_2$	$A_3$	$A_4$
Native	x	1.989	350	214	214	173
	y	1.978	329	195	184	157
	$\perp^b$	1.983	339	204	199	165
	$z^b$	1.956	321	282	282	251
	iso <sup>c</sup>	1.974	333	230	227	194
	aniso <sup>d</sup>	0.028	19	-78	-83	-87
$\text{Sr}^{2+}$ -substituted	x	1.992	328	213	215	161
	y	1.981	347	201	180	175
	$\perp^b$	1.986	338	207	197	168
	$z^b$	1.963	320	283	270	224
	iso <sup>c</sup>	1.978	332	232	221	187
	aniso <sup>d</sup>	0.024	17	-76	-73	-56
$\text{NH}_3$ -modified	x	1.989	326	214	215	154
	y	1.978	345	195	187	175
	$\perp^b$	1.984	336	204	201	164
	$z^b$	1.956	322	283	275	229
	iso <sup>c</sup>	1.975	331	231	225	186
	aniso <sup>d</sup>	0.027	13	-79	-74	-65

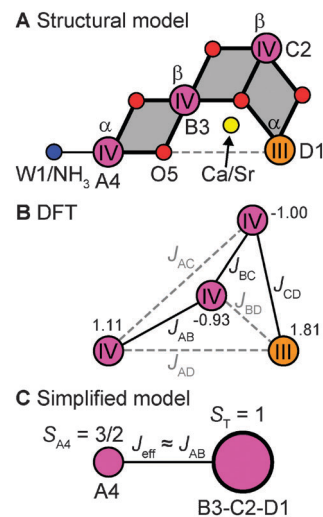
<sup>a</sup> All  $G$  and  $A$  tensors are collinear. <sup>b</sup> The equatorial and axial  $G$  and  $A_i$  values are defined as  $G_{\perp} = (G_x + G_y)/2$ ,  $G_{\parallel} = G_z$ , and  $A_{i,\perp} = (A_{i,x} + A_{i,y})/2$ ,  $A_{i,\parallel} = A_{i,z}$ . <sup>c</sup> The isotropic  $G$  and  $A_i$  ( $i = 1-4$ ) values are the averages of the individual values:  $G_{\text{iso}} = (G_x + G_y + G_z)/3$  and  $A_{i,\text{iso}} = (A_{i,x} + A_{i,y} + A_{i,z})/3$ . <sup>d</sup> The anisotropy in the  $G$  and  $A_i$  values is expressed as the difference  $A_{i,\text{aniso}} = A_{i,\perp} - A_{i,\parallel}$ .

### 3.3 Calculated magnetic properties for the native and modified $S_2$ states of the OEC

The electronic structure of the coupled OEC spin system is defined by the set of six pairwise Mn–Mn exchange interaction terms  $J_{ij}$ , which can be calculated using BS-DFT. For all four computational models describing the set of native and chemically perturbed  $S_2$  state clusters, the calculations reveal that the  $\alpha\beta\beta\alpha$  spin configuration (Fig. 4A and B) is the lowest in energy. Sets of  $J_{ij}$  coupling constants are given in Table S1 of the ESI.† Diagonalization of the Heisenberg Hamiltonian to obtain the complete spin ladder confirms that all four models exhibit an effective total spin  $S_T = 1/2$  ground state, as observed experimentally, and an  $S_T = 3/2$  first excited state. The estimated energy differences between the two lowest states of the spin ladder are on the order of 24–26  $\text{cm}^{-1}$  for the  $\text{Mn}_4\text{O}_5\text{Ca/Sr}$   $S_2$  state structures, lowering by 7  $\text{cm}^{-1}$  upon exchange of W1 for  $\text{NH}_3$  (Table S1, ESI†). These values are in the range inferred from experiments.‡<sup>33,34,118,119</sup>

For all four  $S_2$  state OEC forms, the  $J$ -coupling topology consists of three main coupling pathways (Table S1, ESI† and Fig. 4B): (i) an antiferromagnetic coupling pathway between  $\text{Mn}_{\text{D1}}$  and  $\text{Mn}_{\text{C2}}$  ( $J_{\text{CD}}$ ); (ii) a ferromagnetic coupling pathway between  $\text{Mn}_{\text{C2}}$  and  $\text{Mn}_{\text{B3}}$  ( $J_{\text{BC}}$ ); (iii) and an antiferromagnetic

‡ Differences in the experimental numbers between the  $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$  ( $\geq 30 \text{ cm}^{-1}$ ) and the  $\text{Mn}_4\text{O}_5\text{Ca/Sr}$  ( $\approx 25 \text{ cm}^{-1}$ )  $S_2$  states are likely due to different sample types, *i.e.* higher plant *vs.* cyanobacterial PSII with and without EtOH, respectively.



**Fig. 4** The magnetism of the low-spin  $S_2$  state. (A) Connectivity of the inorganic centres showing the three  $\mu$ -oxo-bridged dimanganese units, defining the main electronic exchange pathways of the cluster. (B) Spin coupling scheme showing the main (black solid lines) and weak (grey dashed lines) exchange couplings  $J_{ij}$  within the OEC. The numbers next to the Mn ions represent their spin projection factors  $\rho_{i,\text{iso}}$  for the native  $S_2$  state BS-DFT model. (C) Simplified spin coupling model in which the trimer unit  $\text{Mn}_{\text{B3}}\text{Mn}_{\text{C2}}\text{Mn}_{\text{D1}}$  interacts with the outer  $\text{Mn}_{\text{A4}}$  via an effective coupling  $J_{\text{eff}}$  that may be approximated as  $J_{\text{AB}}$ .

coupling pathway between  $\text{Mn}_{\text{B3}}$  and  $\text{Mn}_{\text{A4}}$  ( $J_{\text{AB}}$ ). The ferromagnetic exchange pathway  $J_{\text{BC}} = 19\text{--}28 \text{ cm}^{-1}$  is the largest in absolute magnitude, while the antiferromagnetic pathways  $J_{\text{CD}} = -16$  to  $-18 \text{ cm}^{-1}$  and  $J_{\text{AB}} = -12$  to  $-16 \text{ cm}^{-1}$  are slightly weaker. The remaining exchange coupling constants  $J_{\text{AC}}$ ,  $J_{\text{AD}}$  and  $J_{\text{BD}}$  are small, as can be expected from geometric considerations (see Table S1, ESI†).  $J_{\text{CD}}$  and  $J_{\text{BC}}$  represent the two largest exchange interactions within the cuboidal trimer unit ( $\text{Mn}_{\text{B3}}\text{Mn}_{\text{C2}}\text{Mn}_{\text{D1}}$ ) of the cluster, whereas  $J_{\text{AB}}$  can be considered to a good approximation as being representative of an effective exchange interaction between this cuboidal unit and the outer  $\text{Mn}_{\text{A4}}$ , as shown in Fig. 4C.

Systematic differences are observed for the exchange pathways upon the two chemical perturbations, replacement of  $\text{Ca}^{2+}$  by  $\text{Sr}^{2+}$  and  $\text{NH}_3$  exchange at W1 (Table S1, ESI†). When comparing the structure pairs that differ in terms of the presence of  $\text{Ca}^{2+}$  or  $\text{Sr}^{2+}$ , *i.e.*  $\text{Mn}_4\text{O}_5\text{Ca}$  *vs.*  $\text{Mn}_4\text{O}_5\text{Sr}$  and  $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$  *vs.*  $\text{Mn}_4\text{O}_5\text{Sr-NH}_3$ , it is seen that only the major coupling pathways  $J_{\text{CD}}$  and  $J_{\text{BC}}$  are modified, decreasing by 2  $\text{cm}^{-1}$  and 5  $\text{cm}^{-1}$ , respectively.  $J_{\text{AB}}$  remains unchanged. By contrast, for the corresponding structure pairs where  $\text{NH}_3$  is exchanged for W1, the  $J_{\text{CD}}$  pathway is unchanged, while  $J_{\text{BC}}$  and  $J_{\text{AB}}$  increase by 4  $\text{cm}^{-1}$ . It is noted that the perturbation of the O5 position upon  $\text{NH}_3$  substitution, as shown in Fig. 2, results also in an enhancement of  $J_{\text{BD}}$  by 3  $\text{cm}^{-1}$ . In both cases, the changes in the magnetic interactions can be understood within the geometric changes discussed in Section 3.1 (see Fig. 2):  $\text{Sr}^{2+}$  substitution mostly affects the structure of the cuboidal unit, thus perturbing principally the exchange pathways within the Mn-trimer unit, whereas  $\text{NH}_3$  binding perturbs mostly the connectivity between the trimeric moiety and the outer  $\text{Mn}_{\text{A4}}$  (Fig. 4C).





**Table 2** Isotropic and anisotropic spin projection factors  $\rho_{i,\text{iso}}^a$  and  $\rho_{i,\text{aniso}}^b$ , and calculated and experimental isotropic and anisotropic on-site hyperfine values  $a_{i,\text{iso}}^a$  and  $a_{i,\text{aniso}}^b$  in MHz for the Mn ions of the BS-DFT models of the  $\text{Mn}_4\text{O}_5\text{Ca}$ ,  $\text{Mn}_4\text{O}_5\text{Sr}$  and  $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$  clusters in the low-spin  $S_2$  state configuration<sup>c</sup>

Mn ion	Structure	Spin projections		BS-DFT		Experiment	
		$\rho_{i,\text{iso}}$	$\rho_{i,\text{aniso}}$	$a_{i,\text{iso}}$	$a_{i,\text{aniso}}$	$a_{i,\text{iso}}$	$a_{i,\text{aniso}}$
$\text{Mn}_{\text{A4}}$ ( $\text{Mn}^{\text{IV}}$ )	$\text{Mn}_4\text{O}_5\text{Ca}$	1.11	0.23	-247	-33	-206	25
	$\text{Mn}_4\text{O}_5\text{Sr}$	1.13	0.19	-247	-32	-204	31
	$\text{Mn}_4\text{O}_5\text{Ca-NH}_3$	0.94	0.36	-242	-30	-245	-9
$\text{Mn}_{\text{B3}}$ ( $\text{Mn}^{\text{IV}}$ )	$\text{Mn}_4\text{O}_5\text{Ca}$	-0.93	0.29	194	-22	207	26
	$\text{Mn}_4\text{O}_5\text{Sr}$	-0.92	0.24	194	-21	202	8
	$\text{Mn}_4\text{O}_5\text{Ca-NH}_3$	-0.86	0.44	193	-19	221	-33
$\text{Mn}_{\text{C2}}$ ( $\text{Mn}^{\text{IV}}$ )	$\text{Mn}_4\text{O}_5\text{Ca}$	-1.00	0.32	212	-17	226	9
	$\text{Mn}_4\text{O}_5\text{Sr}$	-0.99	0.27	212	-18	224	11
	$\text{Mn}_4\text{O}_5\text{Ca-NH}_3$	-1.01	0.46	213	-16	226	-26
$\text{Mn}_{\text{D1}}$ ( $\text{Mn}^{\text{III}}$ )	$\text{Mn}_4\text{O}_5\text{Ca}$	1.81	0.38	-128	-142	-186	-46
	$\text{Mn}_4\text{O}_5\text{Sr}$	1.78	0.32	-138	-144	-188	-41
	$\text{Mn}_4\text{O}_5\text{Ca-NH}_3$	1.92	0.54	-127	-142	-176	-52

<sup>a</sup> The isotropic  $\rho_{i,\text{iso}}$  and  $a_{i,\text{iso}}$  values are the averages of the individual tensor components  $\rho_{i,\text{iso}} = (2\rho_{i,\perp} + \rho_{i,\parallel})/3$  and  $a_{i,\text{iso}} = (a_{i,x} + a_{i,y} + a_{i,z})/3$ .

<sup>b</sup> The anisotropies of the  $\rho_i$  and  $a_i$  tensors are expressed as the differences  $\rho_{i,\text{aniso}} = \rho_{i,\perp} - \rho_{i,\parallel}$  and  $a_{i,\text{aniso}} = a_{i,\perp} - a_{i,\parallel}$ , *i.e.* between the perpendicular and parallel tensor components. <sup>c</sup> The intrinsic fine structure values of  $\text{Mn}^{\text{IV}}$  ions were assumed to be  $d_{\text{A4}} = d_{\text{B3}} = d_{\text{C2}} = 0 \text{ cm}^{-1}$ . For the  $\text{Mn}_{\text{D1}}^{\text{III}}$  ion, a value of  $d_{\text{D1}} = -1.43 \text{ cm}^{-1}$  was fitted, with  $e_{\text{D1}}/d_{\text{D1}} = 0$ .

Calculated  $^{55}\text{Mn}$  on-site (intrinsic/not spin-projected) hyperfine tensors as in the full (non-effective) spin Hamiltonian based on the BS-DFT models are listed in Table 2 for the four  $S_2$  state variants. The calculated isotropic hyperfine values  $a_{i,\text{iso}}$  for the three  $\text{Mn}^{\text{IV}}$  ions fall within the range seen in  $\text{Mn}^{\text{IV}}$  model compounds experimentally, *i.e.*  $|a_{\text{iso}}| = 187\text{--}253 \text{ MHz}$  (see ref. 30 and 33). The anisotropy of the calculated hyperfine tensors for these three sites is also small, of the order seen in octahedral  $\text{Mn}^{\text{IV}}$  model complexes, *i.e.*  $|a_{\text{aniso}}| < 30 \text{ MHz}$ .<sup>33</sup> For the  $\text{Mn}_{\text{D1}}^{\text{III}}$  ion, the calculated isotropic hyperfine value ( $\approx 130 \text{ MHz}$ ) is smaller than that for  $\text{Mn}^{\text{IV}}$ , as expected, and lies just outside the range seen in  $\text{Mn}^{\text{III}}$  compounds, *i.e.*  $|a_{\text{iso}}| = 165\text{--}225 \text{ MHz}$ .<sup>30,33</sup> As typical for  $\text{Mn}^{\text{III}}$ , it exhibits a significant hyperfine anisotropy, more pronounced than for the  $\text{Mn}^{\text{IV}}$  ions. However, it is noted that the calculated values for the  $\text{Mn}_{\text{D1}}^{\text{III}}$  site are unexpectedly large. Nevertheless, the computed parameters correlate with the inferred site geometry of  $\text{Mn}_{\text{D1}}$ , namely that of a square-pyramidal 5-coordinate  $\text{Mn}^{\text{III}}$  ion. Such a coordination environment generally yields a small isotropic  $^{55}\text{Mn}$  on-site hyperfine coupling and a negative anisotropy (see Table 2), consistent with an effective local  $^5\text{B}_1$  electronic ground state for the  $\text{Mn}_{\text{D1}}^{\text{III}}$  ion.<sup>30,33,120</sup>

The effective hyperfine couplings measured by EPR spectroscopy for oligonuclear metal complexes reflect the on-site hyperfine couplings of the individual metal ion nuclei scaled by the contribution of the electronic spin of each metal ion to the effective spin state:  $A_i = \rho_i a_i$ . The set of scaling factors  $\rho_i$ , termed spin projection coefficients, are primarily determined by the set of pairwise exchange couplings as detailed in ref. 18,

30, 31, 33, 34, 36, 37 and 63. However, additional terms must be included to correctly estimate such spin projections for the OEC, specifically the relevant on-site fine structure parameters  $d_i$  for the individual  $\text{Mn}_i$  ions,<sup>18,30,33,34,63</sup> yielding what are more accurately described as spin projection tensors. As the coordination geometries of the three  $\text{Mn}^{\text{IV}}$  ions of the  $S_2$  state are all octahedral, their local electronic structure should be of approximate spherical symmetry, their orbitals of  $t_{2g}$  origin ( $d_{xy}$ ,  $d_{xz}$  and  $d_{yz}$ ) being half-filled (local high-spin  $d^3$  configuration). As such, the  $\text{Mn}^{\text{IV}}$  ions are expected to only display small fine structure parameters  $d_i$  ( $< 0.3 \text{ cm}^{-1}$ )<sup>121</sup> and hence do not need to be explicitly considered. Thus, the set of parameters which define the spin projection tensors in the  $S_2$  state are the six pairwise exchange interaction terms and the fine structure parameter of the  $\text{Mn}^{\text{III}}$  ion,  $d_{\text{D1}}$ .

Using these spin projection tensors, the fitted projected  $^{55}\text{Mn}$  hyperfine tensors were scaled back to on-site hyperfine tensors to allow comparison to the BS-DFT values discussed above (Table 2). The only plausible assignment for all three forms of the OEC is that  $A_1$ ,  $A_2$ ,  $A_3$ , and  $A_4$  correspond to  $a_{\text{D1}}$ ,  $a_{\text{A4}}$ ,  $a_{\text{C2}}$  and  $a_{\text{B3}}$ , respectively. In our previous work,<sup>33</sup> using BS-DFT structural models predating the latest crystal structure,<sup>9</sup> values of  $-1.2$  to  $-1.3 \text{ cm}^{-1}$  were estimated for a supposedly axially symmetric  $d_{\text{D1}}$  in the native and  $\text{Sr}^{2+}$ -substituted  $S_2$  states. Using the same approach,  $d_{\text{D1}}$  was re-estimated. It was possible to obtain on-site hyperfine anisotropies in the ranges characteristic for  $\text{Mn}^{\text{III}}$  and  $\text{Mn}^{\text{IV}}$  ions employing a single value of  $-1.43 \text{ cm}^{-1}$  for the three OEC systems, well within the range typically seen for  $\text{Mn}^{\text{III}}$  model complexes. As discussed above with regard to the hyperfine tensor anisotropy of  $\text{Mn}_{\text{D1}}^{\text{III}}$ , a negative  $d$  value requires an effective local  $^5\text{B}_1$  state for the  $\text{Mn}_{\text{D1}}^{\text{III}}$  ion and is thus consistent with the square-pyramidal 5-coordinate ligand geometry of  $\text{Mn}_{\text{D1}}^{\text{III}}$ , as present in all computational models. The experimental on-site hyperfine tensor values (Table 2) generally agree well with the BS-DFT estimates and  $\text{Mn}^{\text{III}}$  and  $\text{Mn}^{\text{IV}}$  model compounds, with smaller isotropic values  $a_{i,\text{iso}}$  for the  $\text{Mn}^{\text{III}}$  ion than for the  $\text{Mn}^{\text{IV}}$  ions. In the native and  $\text{Sr}^{2+}$ -substituted models,  $\text{Mn}_{\text{A4}}$  displays a smaller  $a_{\text{A4,iso}}$  than computed and in the  $\text{NH}_3$ -modified system, where it exhibits the largest coupling of the  $\text{Mn}^{\text{IV}}$  ions. Compared to the calculations, the anisotropic components for the  $\text{Mn}^{\text{IV}}$  ions show a larger variance within  $|a_{i,\text{aniso}}| \lesssim 30 \text{ MHz}$ . For  $\text{Mn}_{\text{D1}}^{\text{III}}$ ,  $a_{\text{D1,aniso}} < -40 \text{ MHz}$  is less negative than calculated. Overall, the experimental results confirm that the computed spin coupling schemes serve as a valid description of the native and modified  $S_2$  states.

### 3.4 The $\text{Mn}_{\text{D1}}$ -His332-imino-N interaction

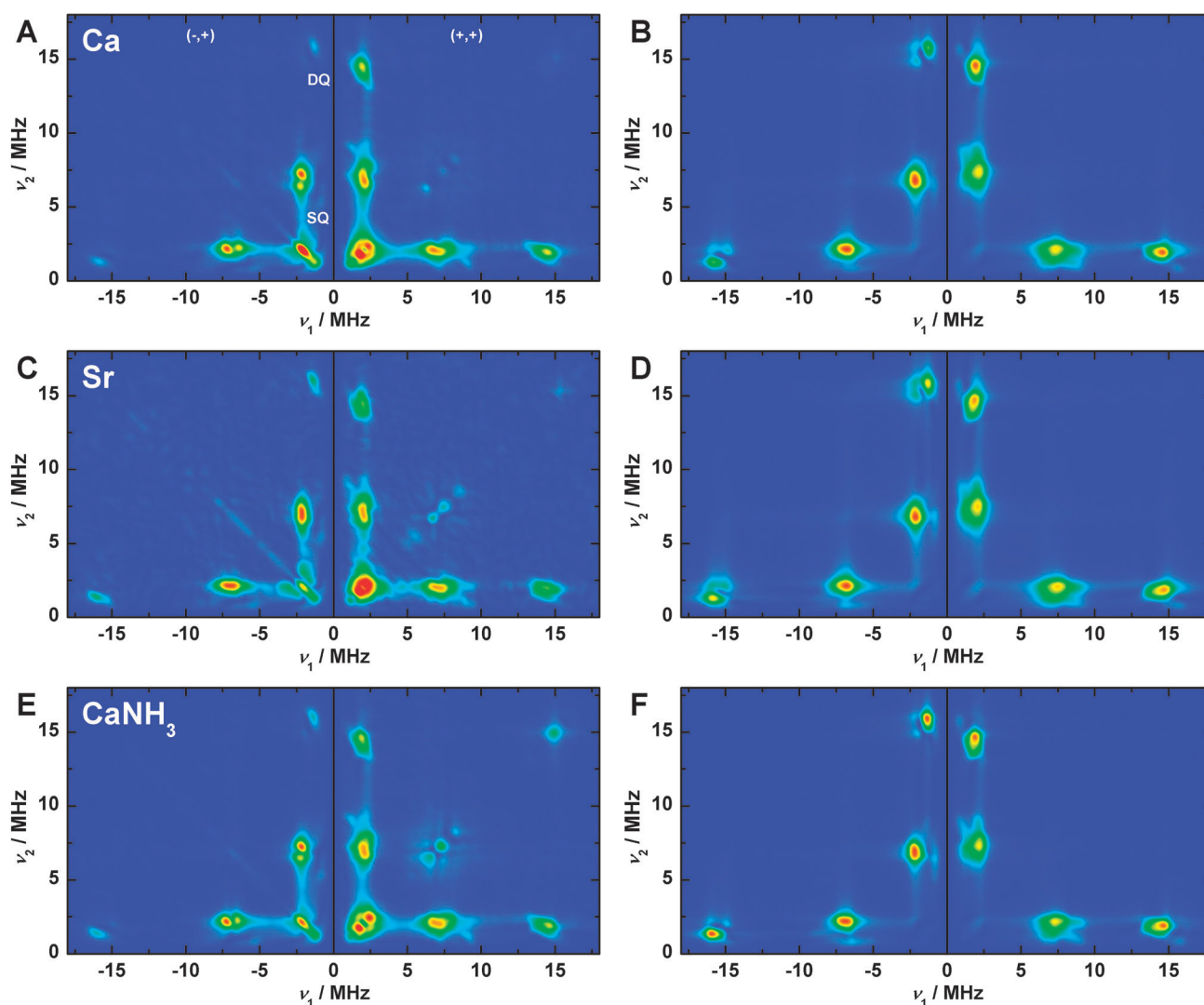
Three-pulse ESEEM measurements were performed to characterize the imino-N signal of His332 associated with the OEC variants in the  $S_2$  state. Fig. S4 and S5 (ESI<sup>†</sup>), respectively, depict  $\tau$ - and magnetic-field-dependent ( $g \approx 2.10\text{--}1.90$ ) light-minus-dark- subtracted spectra and simulations of the native  $\text{Mn}_4\text{O}_5\text{Ca}$  (A, B),  $\text{Sr}^{2+}$ -substituted  $\text{Mn}_4\text{O}_5\text{Sr}$  (C, D) and  $\text{NH}_3$ -modified  $\text{Mn}_4\text{O}_5\text{Ca-NH}_3$  (E, F)  $S_2$  state samples. As noted in Pérez Navarro *et al.*,<sup>44</sup> the  $^{14}\text{N}$  nitrogen signal observed for the native  $S_2$  state





from *T. elongatus* is very similar to that measured in PSII from both higher plants (spinach)<sup>122</sup> and the mesophilic cyanobacteria *Synechocystis* sp. PCC 6803,<sup>123</sup> assigned to the imino-N of His332 *via* mutagenesis.<sup>123,124</sup> The signals are essentially the same in the native, Sr<sup>2+</sup>-substituted and NH<sub>3</sub>-modified OEC clusters with regard to both their  $\tau$  and magnetic-field dependence. The His332 imino-<sup>14</sup>N signal at Q-band nearly fulfils the cancellation condition, where  $A_{\text{iso}}$  is twice the <sup>14</sup>N nuclear Larmor frequency ( $\nu_n = 3.75$  MHz at 1.22 T). The spectra are characterized by three features: the lines centred at frequencies below 2.5 MHz ( $\nu_\alpha = \nu_n - |A_{\text{iso}}|/2$ ), single-quantum transitions around 7.5 MHz ( $\nu_\beta = \nu_n + |A_{\text{iso}}|/2$ ) and smaller double-quantum resonances around 15 MHz ( $\nu_{2\beta} = 2\nu_n + |A_{\text{iso}}|$ ). The line structuring is defined both by the <sup>14</sup>N hyperfine anisotropy and the NQL.

HYSCORE spectroscopy (a two-dimensional ESEEM technique) was performed on the three S<sub>2</sub> state OECs at different magnetic-field positions ( $g \approx 2.07$ – $1.93$ ) of the corresponding Q-band EPR envelopes to further constrain the <sup>14</sup>N hyperfine and quadrupolar interaction matrices. Panels A, C and E in Fig. 5 show the Fourier-transformed spectra and simulations at the centre field position; low- and high-field spectra and simulations are presented in Fig. S6 and S7 in the ESI.† As seen for the three-pulse ESEEM spectra, their appearance is highly similar for all three variants of the OEC in the S<sub>2</sub> state. In two dimensions, the three features that comprise the Q-band ESEEM spectra appear as cross peaks at corresponding frequencies both in the (–,+) and the (+,+) quadrants. As the <sup>14</sup>N hyperfine coupling matches the cancellation condition, the cross peaks are shifted away from the diagonal, instead appearing near the



**Fig. 5** (–,+) and (+,+) quadrants of the Fourier-transformed Q-band HYSCORE experimental spectra (A, C, E) and spin Hamiltonian-based simulations (B, D, F) of the S<sub>2</sub> state Mn<sub>4</sub>O<sub>5</sub>Ca (A, B), S<sub>2</sub> state Mn<sub>4</sub>O<sub>5</sub>Sr (C, D) and annealed S<sub>2</sub> state Mn<sub>4</sub>O<sub>5</sub>Ca–NH<sub>3</sub> (E, F) clusters in PSII samples isolated from *T. elongatus* at central magnetic field. SQ and DQ point out the regions of single- and double-quantum transitions, respectively. The optimized parameter sets for the simulations, as described in Section 2.3 and Sections S3, S4 and S8.4 (ESI†), are listed in Table 3 and, in detail, in Table S2 (ESI†). Experimental parameters: microwave frequencies: 34.0370 GHz (Ca), 34.0433 GHz (Sr), 34.0151 GHz (NH); magnetic fields: 1220 mT (Ca), 1222 mT (NH<sub>3</sub>, Sr); shot repetition time: 1 ms; microwave pulse length ( $\pi/2$ ): 12 ns;  $\tau$ : 260 ns;  $\Delta T$ : 100 ns; temperature: 5.2 K.



frequency axes. Overall, virtually no orientation dependence is seen comparing the spectra at the three different magnetic fields (Section S8.2, ESI†), consistent with the electron–nuclear interaction being dominated by the isotropic component of the hyperfine coupling as compared to the anisotropic part and the traceless NQI, as in ref. 122, 123 and 125. Thus, the orientation of the His332 imino-<sup>14</sup>N hyperfine tensor relative to the *G* tensor cannot be determined from this dataset.

Fitted spin Hamiltonian parameters derived from the simultaneous simulation of both the ESEEM and HYSOCORE datasets are given in Table 3 together with BS-DFT estimates. To directly compare DFT values with experiment, the calculated site hyperfine tensor for the His332 was multiplied by the axial Mn<sub>D1</sub> spin projection tensor described in Section 3.3. All DFT models yield virtually the same hyperfine and quadrupole values. The calculated *A*<sub>iso</sub> underestimates experimental results by <20%, but the dipolar component *A*<sub>dip</sub> and the rhombicity *A*<sub>η</sub> nominally agree with experiment. It is noted that the on-site <sup>14</sup>N hyperfine tensor *a* is expected to be axial with its unique component *a*<sub>1</sub> aligned along the Mn<sub>D1</sub>–N bond, as seen in our calculations (Table S2, ESI†). As such, the axial <sup>14</sup>N hyperfine component *a*<sub>1</sub> is essentially rotated 90° relative to that of the Mn<sub>D1</sub> spin projection tensor, which is expected to coincide with the Jahn–Teller axis of Mn<sub>D1</sub><sup>III</sup>. This 90° rotation explains why the <sup>14</sup>N hyperfine tensor *A* is rhombic in the projected (experimental) reference frame. For a more detailed description of the simulations, see Section S8.4 (ESI†). Importantly, the near-invariance of the imino-N spin Hamiltonian parameters for the three S<sub>2</sub> state forms requires the His332 ligation, the electronic structure of the Mn<sub>D1</sub><sup>III</sup> ion and by extension the whole tetranuclear Mn cluster, to not be significantly perturbed by Ca<sup>2+</sup>/Sr<sup>2+</sup> and NH<sub>3</sub>/W1 exchange, in line with the EPR/<sup>55</sup>Mn ENDOR results described in Section 3.2.

While the orientation of the hyperfine tensor relative to the *G* tensor cannot be determined using ESEEM/HYSOCORE at Q-band frequencies, it can be measured at W-band, *e.g.* using

EDNMR. In our earlier study,<sup>14</sup> it was found that the hyperfine tensor is orientated such that its principal, *i.e.* the smallest component *A*<sub>1</sub> is aligned such that it is mid-way between *G*<sub>xz</sub> and *G*<sub>z</sub>. Importantly though, it is noted that the set of spin Hamiltonian parameters deduced from Q-band ESEEM/HYSOCORE (Table 3) does not reproduce the W-band data sets (Section S8.3, ESI†). This is not due to the inclusion/exclusion of the NQI term, which, for the W-band EDNMR data, mainly contributes to the spectral line width. To reproduce the field dependence of the <sup>14</sup>N- and <sup>15</sup>N-EDNMR signals (Fig. S8 and S9, ESI†), the values determined from Q-band ESEEM/HYSOCORE needed to be scaled: *A*<sub>iso</sub> was decreased by 10%, whereas *A*<sub>dip</sub> was increased by a factor of two (Table S2, ESI†). The same results were observed for all three S<sub>2</sub> state forms, which basically exhibit the same <sup>14</sup>N-EDNMR spectra. A possible reason for this difference comes from the observation that the ground spin state, an effective spin *S*<sub>T</sub> = 1/2 state, is not very well separated energetically from higher spin states in the regime of the W-band excitation energy (≈3 cm<sup>-1</sup>), consistent with DFT estimates for the ground-to-first excited state energy splitting Δ*E* (Section S7, ESI†). Excited-state mixing due to a small Δ*E* has the consequence of altering spin Hamiltonian observables such as effective <sup>55</sup>Mn and <sup>14</sup>N hyperfine tensors. Alternatively, the rhombicity of the effective *G* tensor as inferred from the EPR/<sup>55</sup>Mn-ENDOR simulations may be artificial, a consequence of using collinear *G* and <sup>55</sup>Mn hyperfine tensors. This latter suggestion would also explain why the *G* tensors inferred from W-band measurements on PSII single crystals<sup>116,117</sup> differ from those inferred from our multifrequency measurements on frozen solution PSII samples.

The lack of agreement between the two <sup>14</sup>N datasets brings into question whether the W-band <sup>14</sup>N/<sup>15</sup>N-EDNMR signals can be used to assign the exchangeable μ-oxo bridge <sup>17</sup>O signal based on the relative orientations of the <sup>14</sup>N and <sup>17</sup>O hyperfine tensors, as suggested by Rapatskiy *et al.*<sup>14</sup> Thus, further experimental results, particularly from single crystals of PSII, are needed to test this proposal (see Section 3.6).

**Table 3** Fitted and calculated effective/projected <sup>14</sup>N hyperfine and NQI tensors in MHz for the electron–nuclear couplings of the His332 imino-N and of NH<sub>3</sub> with the various cluster forms studied in the S<sub>2</sub> state in PSII from *T. elongatus*

S <sub>2</sub> state	<sup>14</sup> N	Method	<i>A</i> <sub>iso</sub> <sup>a</sup>	<i>A</i> <sub>dip</sub> <sup>b</sup>	<i>A</i> <sub>η</sub> <sup>c</sup>	<i>e</i> <sup>2</sup> <i>Qq/h</i>	η <sup>c</sup>
Native	His332 <sup>d</sup>	Exp.	7.1	0.75	0.81	1.97	0.75
		DFT	5.8	0.59	0.74	1.65	0.91
Sr <sup>2+</sup> -substituted	His332 <sup>d</sup>	Exp.	7.3	0.69	0.83	1.98	0.79
		DFT	5.8	0.57	0.61	1.65	0.91
NH <sub>3</sub> -modified	His332 <sup>d</sup>	Exp.	7.2	0.75	0.89	1.96	0.80
		DFT	6.1	0.71	0.99	1.68	0.88
	NH <sub>3</sub>	Exp.	2.36	0.33	0.22	1.52	0.47
		DFT	2.68	−0.65	0.02	0.94	0.87
NH <sub>3</sub> -modified, Sr <sup>2+</sup> -substituted	NH <sub>3</sub>	Exp.	2.37	0.28	0.23	1.58	0.45
		DFT	2.68	−0.59	0.03	0.93	0.87

<sup>a</sup> *A*<sub>iso</sub> is defined as the average of the principal components of the hyperfine tensor: *A*<sub>iso</sub> = (*A*<sub>1</sub> + *A*<sub>2</sub> + *A*<sub>3</sub>)/3. <sup>b</sup> *A*<sub>dip</sub> is defined in terms of *T*<sub>1</sub>, *T*<sub>2</sub>, and *T*<sub>3</sub> as *A*<sub>dip</sub> = (*T*<sub>1</sub> + *T*<sub>2</sub>)/2 = −*T*<sub>3</sub>/2. <sup>c</sup> The rhombicity is defined by *A*<sub>η</sub> or η = (*T*<sub>1</sub> − *T*<sub>2</sub>)/*T*<sub>3</sub>, respectively. *T*<sub>1</sub>, *T*<sub>2</sub>, and *T*<sub>3</sub> represent the three principal components of the hyperfine tensors minus *A*<sub>iso</sub> and of the NQI tensors and are labelled such that |*T*<sub>1</sub>| ≤ |*T*<sub>2</sub>| ≤ |*T*<sub>3</sub>|. <sup>d</sup> The Euler rotation angles [α, β, γ] of the NQI relative to the *A* tensors are [20, 12, 0]°, [18, 9, 0]° and [16, 16, 0]° for the Mn<sub>4</sub>O<sub>5</sub>Ca, Mn<sub>4</sub>O<sub>5</sub>Sr and Mn<sub>4</sub>O<sub>5</sub>Ca–NH<sub>3</sub> clusters, respectively.



### 3.5 NH<sub>3</sub> binding to the Ca<sup>2+</sup>- and the Sr<sup>2+</sup>-containing OEC

In the NH<sub>3</sub>-modified S<sub>2</sub> state, a second nitrogen nucleus is bound to the Mn cluster as a terminal ligand, as described in Pérez Navarro *et al.*<sup>44</sup> Its binding can be observed using X-band (three-pulse) ESEEM, as shown in Fig. S11 (ESI<sup>†</sup>) for <sup>14</sup>NH<sub>3</sub>/<sup>15</sup>NH<sub>3</sub>-bound, Ca<sup>2+</sup>- and Sr<sup>2+</sup>-containing PSII. The <sup>14</sup>NH<sub>3</sub> resonances comprise three characteristic single-quantum lines at 0.5, 0.95 and 1.45 MHz split by the NQI and smaller double-quantum transitions centred at 4.9 MHz, highly similar to the higher plant data.<sup>70</sup> Due to the lack of the NQI, the <sup>15</sup>NH<sub>3</sub> signal is clearly less complicated, consisting only of one single-quantum hyperfine peak centred at 0.3 MHz. As seen for the His332 imino-<sup>14</sup>N signal at Q-band, the NH<sub>3</sub> interaction at X-band fulfils the cancellation condition, leading to a narrow  $\nu_\alpha$  line while the  $\nu_\beta$  line is broadened beyond detection.<sup>70</sup> Most importantly, the spectra of the <sup>14</sup>NH<sub>3</sub>-modified Ca<sup>2+</sup>- and Sr<sup>2+</sup>-containing <sup>14</sup>N-PSII samples are essentially identical. Thus, NH<sub>3</sub> binding to the Sr<sup>2+</sup>-substituted S<sub>2</sub> state cluster is the same as in the native S<sub>2</sub> state.

In our first report on NH<sub>3</sub> binding to the OEC, only the <sup>14</sup>NH<sub>3</sub> interaction was considered.<sup>44</sup> Here, we simultaneously fit the spectra of both the <sup>14</sup>NH<sub>3</sub>-modified <sup>15</sup>N-PSII and the <sup>15</sup>NH<sub>3</sub>-modified <sup>14</sup>N-PSII in the S<sub>2</sub> state (Table 3, Fig. S11 and Table S3, ESI<sup>†</sup>). This resulted in an optimized hyperfine tensor  $A = [2.76 \ 2.62 \ 1.69]$  MHz for <sup>14</sup>NH<sub>3</sub> (and for <sup>15</sup>N scaled by the ratio of the nuclear  $g$  values of <sup>15</sup>N and <sup>14</sup>N) and <sup>14</sup>N NQI parameters  $e^2Qq/h = 1.52$  MHz and  $\eta = 0.47$ . Highly similar hyperfine and NQI values reproduce the NH<sub>3</sub> signal observed in the annealed S<sub>2</sub> state of Sr<sup>2+</sup>-containing <sup>14</sup>N-PSII. The isotropic <sup>14</sup>N hyperfine coupling  $A_{\text{iso}} = 2.36$  MHz and  $A_{\text{dip}} = 0.33$  MHz,  $e^2Qq/h$  and the asymmetry parameter  $\eta$  are the same as reported before for PSII from *T. elongatus*<sup>44</sup> and similar to the values from the analysis of higher plant X-band ESEEM spectra<sup>70</sup> (Table 3).

An axial projected <sup>14</sup>NH<sub>3</sub> hyperfine tensor is obtained from BS-DFT calculations, as seen in the experiment. This is because (i) the on-site <sup>14</sup>NH<sub>3</sub> hyperfine tensor is axial, and (ii) its axial and equatorial components are essentially coincident with those of the Mn<sub>A4</sub> spin projection tensor (Table S3, ESI<sup>†</sup>), unlike the case for the His332 imino-<sup>14</sup>N  $a$  tensor (see Section 3.4). The BS-DFT calculations also reproduce the comparably large and rhombic NQI parameters (Table 3), although the sign of the hyperfine anisotropy is inverted compared to experiment. For more details, see Section S9 in the ESI<sup>†</sup>.

### 3.6 Interactions with exchangeable <sup>17</sup>O species

As we have recently shown,<sup>14,19</sup> EDNMR spectroscopy at W-band, due to its comparatively high sensitivity, is the preferred method to measure the interactions of exchangeable <sup>17</sup>O nuclei with fast-relaxing electronic species such as the S<sub>2</sub> state of the OEC. Fig. 6 shows these spectra and simulations (see Section 2.3 and Sections S3 and S4, ESI<sup>†</sup> for details) of the single-quantum region for the native, the Sr<sup>2+</sup>-substituted, the NH<sub>3</sub>-annealed and the Sr<sup>2+</sup>- and NH<sub>3</sub>-modified S<sub>2</sub> state variants after H<sub>2</sub><sup>17</sup>O buffer exchange in the S<sub>1</sub> state (see Fig. S12 (ESI<sup>†</sup>) for the double-quantum region). The spectrum of the native system exhibits the single- and double-quantum resonances of the

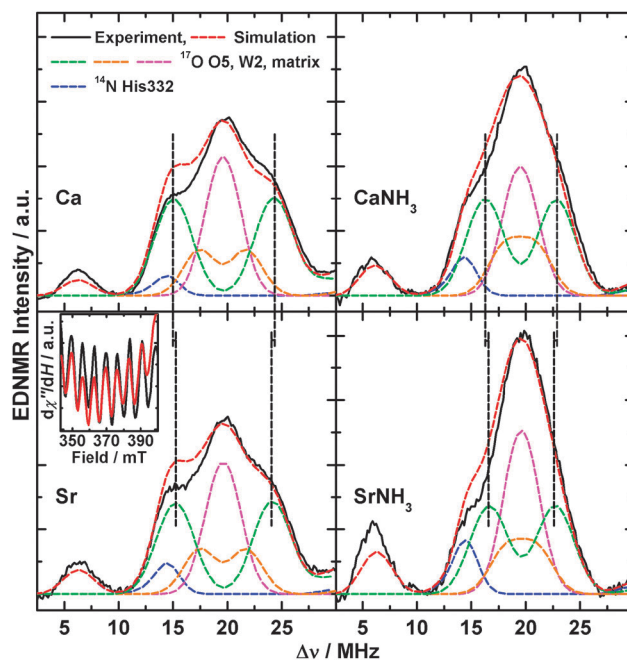


Fig. 6 <sup>17</sup>O-EDNMR spectra of the native (Ca) the Sr<sup>2+</sup>-substituted (Sr), the NH<sub>3</sub>-annealed (CaNH<sub>3</sub>) and the Sr<sup>2+</sup>-substituted NH<sub>3</sub>-annealed (SrNH<sub>3</sub>) S<sub>2</sub> states in PSII samples isolated from *T. elongatus*. Black solid traces depict the single-quantum region of background-corrected experimental spectra; superimposing red dashed traces represent simulations based on the spin Hamiltonian formalism as outlined in Section 2.3 and Sections S3 and S4 (ESI<sup>†</sup>). Coloured dashed lines represent a decomposition of the simulation showing contributions from the individual <sup>14</sup>N and <sup>17</sup>O nuclei. Black dashed lines highlight the decrease of the strong <sup>17</sup>O interaction upon Ca<sup>2+</sup>/Sr<sup>2+</sup> exchange. The optimized parameter sets are listed in Table S4 (ESI<sup>†</sup>). Double-quantum regions are shown in Fig. S12 (ESI<sup>†</sup>). Inset: a section of the X-band CW EPR spectra of the Sr<sup>2+</sup>-substituted S<sub>2</sub> state in PSII samples in the absence (black) and presence (red) of H<sub>2</sub><sup>17</sup>O showing no line broadening upon <sup>17</sup>O exchange. Experimental parameters: see Fig. S12 and for the inset Fig. S14 (ESI<sup>†</sup>).

imino-<sup>14</sup>N of His332 (blue) and of three different classes of <sup>17</sup>O species,<sup>14</sup> *i.e.* (i) a strongly coupled, bridging species (green), (ii) an intermediately coupled terminal O-ligand (orange), and (iii) a weakly coupled terminal class (cyan). These were assigned to the  $\mu$ -oxo bridge O5, the hydroxide ion W2<sup>12</sup> and the H<sub>2</sub>O matrix (comprising ligand W1 of Mn<sub>A4</sub> and two H<sub>2</sub>O ligands at the Ca<sup>2+</sup> ion), respectively. NH<sub>3</sub> binding causes a narrowing of the <sup>17</sup>O single- and double-quantum envelopes, reproduced by a decrease of the hyperfine couplings of O5 and W2 and concomitant reduction of the matrix line intensity, which was interpreted by NH<sub>3</sub> binding to Mn<sub>A4</sub> in exchange for W1.<sup>44</sup> Comparing these two spectral forms to those of the corresponding Sr<sup>2+</sup>-substituted W1- and NH<sub>3</sub>-containing clusters (Fig. S13, ESI<sup>†</sup>), we see a systematic narrowing of the single-quantum envelope by  $\approx 0.5$  MHz and a corresponding narrowing of the double-quantum envelope. This can be reproduced by spectral simulations in which the hyperfine couplings of the  $\mu$ -oxo bridge are reduced accordingly (W1: 9.2 MHz *vs.* 9.7 MHz, NH<sub>3</sub>: 6.5 *vs.* 7.0 MHz), while the other <sup>17</sup>O interactions remain unaltered (for a complete set of hyperfine parameters, see Table S4, ESI<sup>†</sup>). Although weaker than the NH<sub>3</sub> effect, the





narrowing was found to be reproducible in all Sr<sup>2+</sup>-substituted PSII samples. It clearly shows that Ca<sup>2+</sup>/Sr<sup>2+</sup> exchange perturbs the  $\mu$ -oxo bridge, in addition to a simultaneous modification by NH<sub>3</sub> binding.

The inset in Fig. 6 depicts a section of the X-Band CW EPR spectra of the Sr<sup>2+</sup>-substituted S<sub>2</sub> state, which exhibits an intrinsically smaller average line width ( $\approx 3.6$  mT peak-to-peak) than the native form ( $\approx 4$  mT), in the presence and absence of <sup>17</sup>O (see Fig. S14 (ESI<sup>†</sup>) for the entire spectra). No EPR line broadening is observed upon <sup>17</sup>O exchange. This demonstrates that the largest <sup>17</sup>O coupling represents only one exchangeable oxo bridge. In the case of two hyperfine interactions of  $\approx 10$  MHz, the effective line broadening due to the combined <sup>17</sup>O couplings would be larger than 120 MHz or 4.3 mT, exceeding the actual line width.

## 4 Discussion

### 4.1 A common electronic structure of the S<sub>2</sub> state variants

Our DFT results show that the Sr<sup>2+</sup>-substituted, the NH<sub>3</sub>-annealed and the Sr<sup>2+</sup>- and NH<sub>3</sub>-modified low-spin S<sub>2</sub> states basically represent the same structure on both a geometric and electronic level. This result is not immediately obvious from their X-band EPR signals. Indeed historically, the Sr<sup>2+</sup>-substituted and NH<sub>3</sub>-modified forms were explained in terms of a change of the valence state distribution within the Mn tetramer and thus of the coordination environment of the Mn<sup>III</sup> ion.<sup>126</sup> The comprehensive approach pursued in this study conclusively rules out such a mechanism for electronic structure perturbation. Instead, as proposed by our group,<sup>33,34,63</sup> multiline heterogeneity reflects rather subtle changes of the Mn-tetramer structure. The similarity of the perturbed multiline forms suggest a common mechanism for electronic perturbation, which probably also explains S<sub>1</sub> state heterogeneity. This is discussed below, with reference to solvent access, substrate binding and exchange.

**4.1.1 The mechanism of structural perturbation.** In Su *et al.*,<sup>34</sup> a qualitative model for multiline heterogeneity was proposed. In this model, the electronic structure of the S<sub>2</sub> state was considered in terms of a simplified model of two spin fragments: (i) a cuboidal trimer unit, made up of Mn<sub>D1</sub>, Mn<sub>C2</sub> and Mn<sub>B3</sub>, and (ii) a 'monomeric' Mn unit, consisting of the outer Mn<sub>A4</sub>. Therein, variation of the electronic structure in S<sub>2</sub> was attributed to changes in the connectivity of the outer Mn<sub>A4</sub> to the cuboid, changing the properties of the electronic ground state by altering the mixing-in of excited spin state character. The physical rationale for this observation was that the outer Mn<sub>A4</sub> represents the solvent accessible end of the cluster. It has two water-derived ligands and solvent channels that begin or terminate at this site.<sup>9,127</sup> As such, it is this site and its connection to the rest of the cluster that is most likely to vary amongst different sample conditions and possibly different PSII species. In terms of this 'monomer-trimer' model with regard to the two modifications discussed here, Ca<sup>2+</sup>/Sr<sup>2+</sup> forms part of the linkage between the cuboidal and outer fragments, as mediated by the  $\mu$ -oxo bridge O5, whereas NH<sub>3</sub> binds to the

outer Mn<sub>A4</sub>, also perturbing O5, the bridge *trans* to its binding position W1.<sup>44</sup>

The magnetic observable that is altered upon Ca<sup>2+</sup>/Sr<sup>2+</sup> replacement and/or W1/NH<sub>3</sub> exchange, leading to the perturbed multiline forms is the <sup>55</sup>Mn hyperfine anisotropy. Small perturbations of the four hyperfine tensors result in a change in the hyperfine peak superposition, altering the apparent structure of the X-band EPR signal (Fig. 3A). Importantly these changes are subtle, as demonstrated by the invariance of the <sup>55</sup>Mn-ENDOR spectra (Fig. 3D). The <sup>55</sup>Mn hyperfine anisotropy is not a site property, but instead an indirect measure of the fine structure splitting of the Mn<sup>III</sup> ion<sup>30,33,34,63</sup> or, in the 'monomer-trimer' model, the zero-field splitting of the whole trimer unit. Within this model, its contribution is modulated by the electronic connectivity between the two fragments, predominantly the exchange pathway *J*<sub>AB</sub>, the coupling that mostly defines the energy splitting  $\Delta E$  between the ground state and the first excited state (Fig. 4C). Our BS-DFT results support this basic mechanism for electronic structure perturbation and, for the first time, describe the changes on the molecular level that impart this variation, and which differ for the two modifications. Upon replacement of Ca<sup>2+</sup> by Sr<sup>2+</sup>, the slight distortion of the cuboidal moiety leads to a perturbation of the intra-cuboidal exchange network and possibly the Mn<sub>D1</sub><sup>III</sup> site fine structure splitting. It is noted that this, besides changing the <sup>55</sup>Mn hyperfine anisotropy, also manifests itself in terms of the *G* tensor, also contributing to the altered multiline appearance and the *g* shift of the W-band EPR signal. Exchange of W1 by NH<sub>3</sub> affects the connectivity of the outer Mn<sub>A4</sub> to the cuboidal unit, as modulated by the  $\mu$ -oxo bridge O5, perturbing the *J*<sub>AB</sub> exchange pathway, thus changing the <sup>55</sup>Mn hyperfine anisotropy. In the case of the S<sub>2</sub> state variant that contains both these modifications, their effects on the electronic structure are additive. It is noted that it is the properties of the cuboidal unit that define the *G* tensor as opposed to the outer Mn<sub>A4</sub>, which presumably has an isotropic on-site *g* value. This is expected, as it is the Mn<sub>D1</sub><sup>III</sup> ion, which is part of the trimer fragment in the *S* = 1/2 configuration, that should form the dominant contribution to the anisotropy of the *G* tensor in all four systems.

### 4.1.2 The Mn<sub>D1</sub>-His332-imino-N and Mn<sub>A4</sub>-NH<sub>3</sub> interactions as local probes for the electronic structure

**4.1.2a The Mn<sub>D1</sub>-His332-imino-N interaction.** As described by Stich *et al.*,<sup>123</sup> the magnitude of the Mn<sub>D1</sub>-His332 imino-N hyperfine interaction, as compared against mixed-valence Mn<sup>III</sup>Mn<sup>IV</sup> model compounds and protein cofactors with imidazole ligands to Mn<sup>III</sup> (*A*<sub>iso</sub>  $\leq 13$  MHz) and Mn<sup>IV</sup> ions (*A*<sub>iso</sub> = 1.5–3.3 MHz),<sup>128–134</sup> favours assigning Mn<sub>D1</sub> as the only Mn<sup>III</sup> ion of the S<sub>2</sub> state, consistent with the EPR/<sup>55</sup>Mn-ENDOR/DFT results already reported in the literature and detailed above (Sections 3.1–3.3). The large hyperfine couplings seen for ligands coordinating to Mn<sup>III</sup> in S<sub>2</sub> (and model systems) comes from the fact that the Mn<sup>III</sup> ion carries the largest spin projection coefficient, *i.e.* in Mn dimers  $\rho_{\text{iso}}(\text{Mn}^{\text{III}}) = 2$  and  $\rho_{\text{iso}}(\text{Mn}^{\text{IV}}) = -1$ . Interestingly, the hyperfine and quadrupole couplings of imidazole ligands of Mn<sup>III</sup> ions differ depending on whether they represent axial (*A*<sub>iso</sub> = 9–13 MHz,  $e^2Qq/h = 2.1\text{--}3.0$  MHz)<sup>128–132</sup> or equatorial ligands (*A*<sub>iso</sub> = 1.5–6.6 MHz,  $e^2Qq/h = 1.5\text{--}2.5$  MHz).<sup>131–134</sup> The values seen for the His332

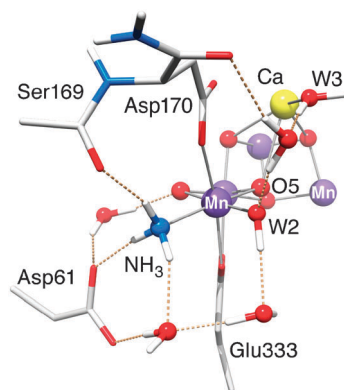


imino-N ( $A_{\text{iso}} = 7.1$  MHz,  $e^2Qq/h = 1.97$  MHz) fall closer to the equatorial range supporting its assignment as an equatorial ligand consistent with DFT structural models. In such models,<sup>12,37,40,103,106</sup> the local Jahn–Teller axis of the  $\text{Mn}_{\text{D1}}^{\text{III}}$  ion is aligned along the open coordination site, thus considered a pseudo-Jahn–Teller axis, perpendicular to the  $\text{Mn}_{\text{D1}}\text{--N}$  bond. It is supposed that the reason why the  $^{14}\text{N}$  couplings measured for the His332 do not exactly fall within the range seen in model complexes is that all model complexes measured thus far represent 6-coordinate  $\text{Mn}^{\text{III}}$  ions whereas the  $\text{Mn}_{\text{D1}}^{\text{III}}$  ion in the  $\text{S}_2$  state is 5-coordinate.

**4.1.2b The  $\text{Mn}_{\text{A4}}\text{--NH}_3$  interaction.** As recently shown in Pérez Navarro *et al.*,<sup>44</sup> the binding site of  $\text{NH}_3$  is likely the W1 site. The small effective isotropic  $^{14}\text{N}$  hyperfine coupling ( $A_{\text{iso}} = 2.36$  MHz) and the axially of the hyperfine tensor are both consistent with a terminal ligand to a  $\text{Mn}^{\text{IV}}$  ( $d^3$ ) ion.<sup>44,70</sup> The similar  $A_{\text{iso}}$  in the  $\text{Mn}_4\text{O}_5\text{Sr--NH}_3$  cluster confirms that the oxidation state of the  $\text{Mn}_{\text{A4}}$  ion is not altered by  $\text{Sr}^{2+}$  substitution. Moreover, the binding mode and perturbation mechanism of  $\text{NH}_3$  is the same in the  $\text{Ca}^{2+}$ - and  $\text{Sr}^{2+}$ -containing Mn clusters. The non-axiality of the electric field gradient ( $\eta = 0.47$ ) is characteristic for this ligand. A large asymmetry parameter is uncommon for a terminal ligand of  $\text{Mn}^{\text{IV}}$  (although our value is already  $\approx 20\%$  smaller than that reported earlier<sup>70</sup>). The latest crystal structure<sup>9</sup> suggests that such an asymmetric distortion could be present for the W1 site due to the charged residue D1-Asp61, in H-bonding distance to W1/ $\text{NH}_3$ , as seen for other protein systems.<sup>135</sup> Indeed, upon inclusion of the Asp61 residue, which was not included in our previous, smaller BS-DFT model,<sup>44</sup> the asymmetric quadrupole tensor is now reproduced, and the hyperfine coupling constant shows better agreement with experiment (Fig. 7 and Table 3). In contrast, such an asymmetric distortion is not seen for the W2 ligand as a similar charged amino-acid residue partner is not present to provide a H-bond.

## 4.2 The exchangeable $\mu$ -oxo bridge

Both modifications investigated,  $\text{Ca}^{2+}/\text{Sr}^{2+}$  substitution and  $\text{NH}_3/\text{W1}$  replacement, perturb the  $^{17}\text{O}$ -EDNMR signals of



**Fig. 7** Hydrogen bonding network in the vicinity of the  $\text{Mn}_{\text{A4}}$ -bound  $\text{NH}_3$ , as observed in the optimized DFT structure of the  $\text{Mn}_4\text{O}_5\text{Ca--NH}_3$  model. Mn ions are depicted in purple, O in red, Ca in yellow, C in light grey, N in blue and H in white. Only relevant residues, water molecules and protons are shown.

exchangeable oxygen species of the OEC, specifically the exchangeable  $\mu$ -oxo bridge. It is this bridge that likely represents one of the substrate water sites of the Mn tetramer. As the electronic structure of the OEC is essentially invariant for all four OEC forms, the change in hyperfine coupling for this  $\mu$ -oxo bridge must represent a site modification, near or at the oxygen nucleus.  $\text{NH}_3$  binding primarily affects the connectivity of the outer  $\text{Mn}_{\text{A4}}$  to the cuboidal trimer, whereas  $\text{Sr}^{2+}$  substitution instead perturbs the exchange network within the cluster. Thus, it can be surmised that the exchangeable  $\mu$ -oxo bridge must both coordinate to the outer  $\text{Mn}_{\text{A4}}$  and be associated with the  $\text{Ca}^{2+}/\text{Sr}^{2+}$  ion itself as a structural element of the cuboidal trimer. Only the bridge O5 fulfils both these criteria. As a ligand to the  $\text{Ca}^{2+}/\text{Sr}^{2+}$  ion, O5 is affected by the exchange of these ions of the same charge but different sizes. Similarly, as argued in Pérez Navarro *et al.*,<sup>44</sup>  $\text{NH}_3/\text{W1}$  exchange perturbs O5 by binding *trans* to this bridge position, distorting the  $\text{Mn}_{\text{A4}}\text{--O5}$  bond length.

It is noted that these results exclude the possibility that  $\text{NH}_3$  displaces the exchangeable  $\mu$ -oxo bridge as a bridging  $\text{--NH}_2\text{--}$  species, an alternative rationale for the narrowing of the  $^{17}\text{O}$  signal envelope in line with earlier suggestions.<sup>70</sup>  $\text{Ca}^{2+}/\text{Sr}^{2+}$  and  $\text{W1}/\text{NH}_3$  exchange are *additive* in terms of their effect on the width of the  $^{17}\text{O}$ -EDNMR envelope, modelled here as defined by the  $\mu$ -oxo bridge hyperfine coupling  $A_{\text{iso}}$ . This result mirrors the structural modifications observed for the doubly modified  $\text{Mn}_4\text{O}_5\text{Sr--NH}_3$  OEC model; *i.e.*, the model contains *additive* structural modifications reflecting both singularly modified  $\text{Mn}_4\text{O}_5\text{Sr}$  and  $\text{Mn}_4\text{O}_5\text{Ca--NH}_3$  structures. If instead the  $\text{NH}_3$  did indeed replace the bridge, the width of the  $^{17}\text{O}$ -EDNMR envelope would be now defined by the W2 hyperfine coupling, and as such should be invariant to  $\text{Ca}^{2+}/\text{Sr}^{2+}$  substitution. It is also noted that  $\text{NH}_3$  replacement of the exchangeable bridge O5 cannot quantitatively explain the virtually unaltered  $^{14}\text{NH}_3$  signal upon exchange of the O5-binding  $\text{Ca}^{2+}$ , and the  $^{17}\text{O}$  hyperfine changes. Assuming an unaltered spin density on the bridge position, as follows from the similar spin projection factors for the four Mn ions, the measured  $^{14}\text{N}$  hyperfine coupling for the bound ammonia at this position is far too small. On the other hand, the  $^{17}\text{O}$  coupling of 6.5–7 MHz seen for the  $\text{NH}_3/\text{W1}$ -exchanged system is in the range of those observed in Mn model complexes with a N-ligand *trans* to the oxo bridge.<sup>14</sup>

## 5 Conclusions

Time-dependent mass spectrometry experiments indicate that the early binding substrate ( $\text{W}_s$ ) is associated with all intermediate states of the OEC.<sup>66,136</sup> Furthermore, the relatively slow exchange and the S-state dependence of this bound substrate with bulk water suggests that it represents a ligand of (a) Mn ion(s). As  $\text{Ca}^{2+}/\text{Sr}^{2+}$  substitution also perturbs its exchange rate,  $\text{W}_s$  is also supposed to coordinate to the  $\text{Ca}^{2+}$  ion.<sup>65,66</sup> Of the exchangeable oxygen species identified here by  $^{17}\text{O}$ -EDNMR, only O5 is a ligand to both Mn and  $\text{Ca}^{2+}$ . Similarly, only the O5 spectral signature is perturbed by  $\text{Ca}^{2+}/\text{Sr}^{2+}$  exchange. Thus, O5 is the most likely candidate for  $\text{W}_s$ . This assignment limits the



possible reaction pathways for photosynthetic water splitting, and lays a foundation for studies of higher oxidized S states, which will serve to identify the second, fast exchanging substrate and eventually elucidate the mechanism of O–O bond formation. Currently two pathways are envisaged: O–O bond formation could proceed as a coupling between O5 and either (i) Mn<sub>A4</sub>-bound W2 or Ca<sup>2+</sup>-bound W3, or (ii) a further oxygen not present yet in the S<sub>2</sub> state.

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