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Introduction

Metal oxides have recently attracted particular attention because of their unique physical and chemical properties.¹⁻³ As an important functional material, indium oxide (In_2O_3) , an n-type semiconducting binary oxide with a wide band-gap (3.5-3.7 eV), has attracted much interest for several decades due to its high electrical conductivity and good optical transparency.4,5 There are three different crystal structures of In₂O₃: cubic bixbyite-type $(c-In_2O_3)$ stable under ambient conditions, hexagonal corundum-type $(h-In_2O_3)$ and orthorhombic Rh₂O₃(II)-type structures stable only under high pressures and temperatures.6 Most previous research has been performed on the most stable c-In₂O₃ of specific morphologies for enhancing their properties using various techniques.7-9 With the recent report by Idota et al. that some tin-based oxides exhibit high capacities in lithium ion batteries (LIBs), new opportunities seem to be opening for exploring new anode materials based on metal oxide materials.10 The electrochemical reactivity of c-In₂O₃ with lithium at room temperature has been investigated because of its high theoretical capacities.¹¹⁻¹⁴ The reversible lithium storage performance of h-In₂O₃ has not been reported so far.

Large-scale synthesis of hexagonal corundum-type In₂O₃ by ball milling with enhanced lithium storage capabilities

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Hexagonal corundum-type indium oxide $(h-ln_2O_3)$ is the structure that normally exists in a high-temperature and pressure environment. This structure has been realised from ambient environment stable cubic indium oxide $(c-ln_2O_3)$ using a high-energy ball milling approach at room temperature, in which the rearrangements of lnO_6 polyhedral units take place *via* plastic deformation and large defect creation during the milling process. More interestingly, the high-temperature $h-ln_2O_3$ structure as anode materials in lithium-ion batteries exhibits lithium storage capabilities enhanced by up to 8 times compared to the $c-ln_2O_3$ phase. This study demonstrates an effective ambient environmental approach for the production of high-pressure/temperature structures, $h-ln_2O_3$, which may be extended to explore new phases and novel properties in other oxide systems.

h-In₂O₃ has been found to exhibit stable electric conductivity and high sensitivity to gases, compared to c-In₂O₃. However, the synthesis of h-In₂O₃ is very difficult because of the very high temperature and pressure synthesis conditions.¹⁵⁻¹⁸ There are two main methods for fabricating h-In₂O₃ in the literature: the static or shock-induced phase transformation method involving high pressures (>15.3 GPa) or high temperatures and high pressures (>1000 °C and >6.5 GPa)^{6,19,20} and special chemical methods under ambient pressure.^{15,18,21,22} These methods need not only an elevated temperature for the decomposition of metal salts, but also complex reaction conditions, which are disadvantageous from the manufacturing point of view.

In this letter, a large-scale and room-temperature production approach is demonstrated. The high-pressure phase $h-In_2O_3$ can be produced from the $c-In_2O_3$ phase *via* a high-energy ball milling process at room temperature. The low-environmental footprint synthesis method is based on a simple mechanical milling process with advantages of medium pressure range, low production temperature, and no need of any solvents and catalysts. More interestingly, the $h-In_2O_3$ phase exhibits enhanced reversible storage capacity up to 8 times higher and good cycling performance in lithium ion batteries compared to the initial $c-In_2O_3$, indicating the advantages of high-pressure structures as electrode materials for energy storage.

Experimental section

Synthesis of h-In₂O₃ powders

In a typical process, 9 g of c- In_2O_3 powders (99.995%, Alfa Aesar) were loaded into a stainless steel container with four hardened steel balls (diameter of 25 mm), and then filled with N_2 -15% H_2

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gas to a pressure of 300 kPa. A vertical rotating ball mill was used for the milling experiments under high-energy impact mode, which was realized by placing an external magnet under the container at a 45° position in relation to the vertical direction and the use of the mill rotation speed around 150 rpm.²³ The temperature of the container surface was about 40 °C during the milling process. The milling action was turned off every 2 days for 2 h for cooling purposes.

Characterization

The phase changes of milled samples were monitored with an X-ray powder diffraction (XRD) spectrometer (Panalytical X'Pert PRO system) using Cu K α radiation at a wavelength of 1.5418 Å. Simulation and analysis of the XRD patterns of various phases were performed using the Material Studio program. Unit-cell parameters were obtained using DICVOL91, and structure refinement was carried out using the Pawley method.²⁴ Scanning electron microscopy (SEM) analysis was performed on a Zeiss Supra 55 VP instrument. Transmission electron microscopy (TEM) analysis was performed on a JEOL 2100 (operating at 110 kV) apparatus. High resolution TEM studies were carried out with a JEOL JEM 2011 LaB₆ (operating at 200 kV) instrument. Raman spectra were recorded on a Horiba Jobin Yvon Raman microscope in a backscattering geometry at the 532 nm line of an argon ion laser, provided with a CCD detector system.

Electrochemical measurements

The working electrodes of Li-ion battery cells were fabricated by mixing 70 wt% active material (h-In₂O₃ powders), 15 wt% acetylene black (Super-P), and 15 wt% polyvinylidene fluoride (PVDF) binder in N-methyl-2-pyrrolidone (NMP) to form homogeneous slurry. A thin layer of the slurry was coated on Cu foil to form the anode electrode; the electrode was dried overnight at 120 °C in vacuum to remove the solvent. Coin cells (CR2032) were assembled using Li metal as the counter electrode in a dry glove box filled with argon. The electrolyte consisted of 1 M LiPF₆ in a non-aqueous solution of ethylene carbonate and dimethyl carbonate in a volume ratio of 1:1. Cyclic voltammetry (CV) measurements were carried out using a Solartron 1470E electrochemical workstation at a scanning rate of 0.1 mV s⁻¹. A galvanostatic cycling test of the assembled cells was performed using the voltage window of 0.01-3.0 V versus Li/Li⁺.

Results and discussion

Typical XRD patterns of In_2O_3 were collected and are plotted in Fig. 1 as a function of milling time. All diffraction peaks of the initial material can be indexed to a pure cubic bixbyite-type structure ($Ia\bar{3}$) In_2O_3 (c- In_2O_3) crystal with space group no. 206, lattice constants a = b = c = 10.11(1) Å, and unit cell volume $V_0 = 64.8(1)$ Å³, in good agreement with those previously reported in the literature.²⁵ As shown in Fig. 2a, it has two different types of In^{3+} (In_1 and In_2) occupying both octahedral and trigonal prismatic interstices within the lattice of O^{2-} ions, and O atoms at the 8b, 24d, and 48e Wyckoff positions,



Fig. 1 XRD patterns of the In₂O₃ recorded at different milling times.



Fig. 2 Pawley full-profile refinements of the diffraction patterns collected from the ln_2O_3 powders before milling (a) and after milling for 8 days (b). In the case of the high pressure phase, a good fitting with $R_{wp} = 5.2\%$ is obtained for the diffraction pattern in (b). Red, black, blue solid lines, and tick marks represent experimental, calculated, residual patterns, and the positions of calculated Bragg reflections, respectively. The insets show the respective crystal structure, and the different types of indium atoms are marked with different colors.

respectively (inset of Fig. 2a). No changes were observed in the XRD patterns of the samples after milling for 2 and 4 days except peak intensity reduction and shape broadening. The significant peak shape changes are caused by grain size reduction and structural distortion as the typical effect of ball milling

treatment. Three new diffraction peaks emerging around 32.8, 57.7, and 58.9° are indicative of a transformation to a new phase after milling for 6 days.

With further milling up to 8 days, the diffraction patterns of the sample show distinctive features indicating phase transition completion (Fig. 1). The Pawley refinement for this new phase using an $R\bar{3}c$ structural model yields a very good fit with $R_{\rm wp} = 5.2\%$ (Fig. 2b) corresponding to the high pressure and high temperature h-In₂O₃. The lattice parameters refined within this space group for the new phase are a = 5.438(2) Å, c =14.474(3) Å, and unit cell volume $V_0 = 61.78(1)$ Å³ (Z = 6). h-In2O3 has a different atomic arrangement, containing only trigonal biprism coordinated In³⁺ ions within the lattice of O²⁻ ions, with O atoms located at Wyckoff positions 12c and 18e, as shown in Fig. 2b. It can be explained that plastic deformation and increased defect concentrations created by ball milling play important roles in the occurrence of phase transformations in the In₂O₃ system. This mechanical activation initiates the lattice change of O²⁻ ions and the shift of In³⁺ ions from octahedral and trigonal prismatic to trigonal biprism sites (Fig. 2), resulting in distortion and redistribution of InO₆ octahedra units which are the building blocks of the h-In₂O₃ structure.

To further examine the different crystalline structures of the milled In_2O_3 , Raman spectroscopy measurements were performed on the samples before and after milling for 8 days. Group theory predicts the following representation for the center of Brillouin zone (q = 0) optical vibrational modes of c-In₂O₃.²⁶ The irreducible representation is given as

$$\Gamma = 4A_g + 4E_g + 14T_g + 5A_u + 5E_u + 16T_u$$

where A_g , E_g , and T_g represent Raman-active modes, A_u and E_u represent inactive modes, and T_u represents infrared-active modes.²⁷ Fig. 3 shows representative Raman spectra of the indium oxides in the frequency range of 110–900 cm⁻¹ before milling and after milling for 8 days. Five Raman active phonon modes of the initial sample are discernible and all the observed modes correspond well to the band positions of c-In₂O₃



Fig. 3 Micro-Raman spectra of In_2O_3 powders before milling and after milling for 8 days in the 100–900 cm⁻¹ frequency regions.



Fig. 4 SEM image of h-ln $_2O_3$ particles (a) and histogram of diameter distribution of the particles (b).

reported in the literature.²⁶ After milling for 8 days, different features appear in the Raman spectrum. The positions of the six Raman peaks observed at 161, 218, 304, 377, 506 and 595 cm⁻¹ are in good agreement with the previously described h- In_2O_3 structure, predicting seven Raman active modes $(2A_{g1} + 5E_g)$ by the factor group analysis.²² The Raman results actually provide additional evidence of the h- In_2O_3 phase obtained from the ambient phase c- In_2O_3 directly by a simple high-energy ball milling process.

The SEM image (Fig. 4a) shows that most h-In₂O₃ particles are spherical and have diameters in the range of 0.1–1 μ m. The histogram (Fig. 4b) shows that most particles have a diameter around 200 nm. The TEM image in Fig. 5a shows typical morphology and sizes of the starting c-In₂O₃ particles. The highresolution TEM image in Fig. 5b presents an interdistance of 0.407 nm between the parallel fringes corresponding to the *d*spacing of the (211) plane of c-In₂O₃. The TEM image presented in Fig. 5c reveals the smaller particle and grain sizes of h-In₂O₃ particles. A typical selected area diffraction (SAD) pattern is shown in Fig. 5c. All electron diffraction rings can be indexed to the h-In₂O₃ phase which is consistent with the XRD analysis



Fig. 5 TEM images of the In_2O_3 powders (a) before milling and (c) after milling for 8 days; the SAD pattern corresponding to part c is shown in the inset. HRTEM images of the In_2O_3 powders (b) before milling and (d) after milling for 8 days.

results. The HRTEM image (Fig. 5d) of an h-In₂O₃ particle shows the interdistance of the fringes of 0.283 nm, in agreement with the lattice distance of the (104) planes of the h-In₂O₃. These results further confirm the phase transition from c-In₂O₃ to h-In₂O₃ during the high-energy ball milling process.

It is well known that energy barriers often exist in chemical reactions and phase transformations. The reaction or transformation becomes possible when sufficient energy is provided to overcome these barriers. A huge energy barrier should exist for the transformation from c to $h-In_2O_3$ as it normally takes place at high temperatures and high pressures.⁵ In the case of the room-temperature milling process, a huge amount of energy is injected into the materials by repeated high-energy impacts and the energy is stored in the form of structural defects and large surface energy. After 6 days of milling, when the stored energy in the system is sufficiently high, c to $h-In_2O_3$ transformation takes place even at room temperature. Continuous milling is still needed to help with transformation kinetics until full transformation at the end of 8 days.²⁸⁻³⁰

The electrochemical performance of the h-In₂O₃ in the cell configuration of Li/h-In2O3 was evaluated. The CV measurement (Fig. 6a) was carried out to investigate the electrode reaction processes at a scan rate of 0.1 mV s^{-1} . In the first cycle, two apparent peaks are observed at 0.58 V and 1.06 V in the discharge process. The peak at 0.58 V could be assigned to the formation of LixIn alloy proceeding in a multi-step process during lithium insertion.12,13 These two peaks disappear in further cycles, indicating the irreversible reduction of h-In₂O₃ with a large irreversible capacity in the first cycle. New reduction peaks at 0.78 and 0.46 V appear after the first cycle and shift to a lower voltage during further cycles. Two oxidation peaks at 0.51 and 0.70 V could be ascribed to the de-alloying of Li/In which form in the discharge process. Another two broad oxidation peaks at 1.18 and 1.78 V are observed, which may be ascribed to the reverse alloving reaction among Li, In alloys of different stoichiometry and the further formation of metallic phase



Fig. 6 Electrochemical properties of $h-ln_2O_3$: (a) cyclic voltammogram of the $h-ln_2O_3$ electrode; (b) discharge–charge voltage curves under a current density of 30 mA g⁻¹; (c) cycling performance and coulombic efficiency of $h-ln_2O_3$ and initial c-ln_2O_3 at 30 mA g⁻¹ for 50 and 30 cycles, respectively; (d) rate performance of $h-ln_2O_3$ at rates from 30 to 300 mA g⁻¹.

indium. These two peaks fade away quickly after the first several cycles indicating the poor reversibility of this reaction. Fig. 6b shows the 1st, 2nd, 5th, and 50th cycles charge–discharge curves of the h-In₂O₃ at a constant current density of 30 mA g⁻¹ and a cutoff voltage window of 3.0 and 0.01 V. The curves show different plateaus due to the redox reactions associated with Li⁺ insertion and extraction and are consistent with the CV results. The ramp above 0.87 V in the discharge curve possibly derives from the irreversible reduction of h-In₂O₃. The discharge plateaus around 0.42 V indicate the formation of Li_xIn alloy phases and the charge plateaus around 0.69 V are also in agreement with the cyclic voltammogram profile.

The h-In₂O₃ electrode delivers initial discharge-charge capacities of 1256 and 729 mA h g^{-1} , respectively, which are higher than the corresponding capacities (1016 and 453 mA h g^{-1}) from c-In₂O₃ as shown in Fig. 6c. Subsequently, the charge and discharge capacities of the h-In2O3 electrode approach a constant level at about 390 mA h g^{-1} after 50 cycles with a coulombic efficiency of 98.8%. This reversible capacity is much higher than that of the c-In₂O₃ (49 mA h g^{-1}) electrode after 30 cycles and some other oxides such as MnO_2 (about 200 mA h g⁻¹ after 10 cycles),³¹ Co₃O₄ (about 200 mA h g⁻¹ after 20 cycles)³² and Mn₃O₄ (below 200 mA h g⁻¹ after 10 cycles).³³ In addition to the improved reversible capacity and excellent cycling behavior of the h-In₂O₃, the rate capabilities at various current rates from 30 to 300 mA g^{-1} are illustrated in Fig. 6d. A reversible capacity of 125 mA h g⁻¹ can be sustained at the highest current rate of 300 mA g^{-1} . A capacity of 382 mA h g^{-1} at 30 mA g^{-1} is retained after 25 cycles of charge and discharge at various current densities, demonstrating good reversibility and structural stability of the h-In₂O₃.

The better performance in terms of capacity and cyclability of h-In₂O₃ compared to c-In₂O₃ may be attributed to three aspects. First, the structure of c-In₂O₃ is composed of an edgeconnected network of InO₆ octahedra and trigonal prisms, as shown in Fig. 7a, in which Li ions occupy the interstitial sites in this lattice framework. Consequently, the restricted Li diffusion in three-dimensional volume leads to poor electrochemical performance of c-In₂O₃ as the LIB anode. In contrast, the h-In₂O₃ structure consists of a three-dimensional array of edges and corners sharing InO₆ trigonal prisms, giving rise to channels and cavities perpendicular to the *c*-axis (Fig. 7b) and these channels provide elegant pathways for diffusion and storage of Li-ions. Second, the h-In₂O₃ has a better electrical conductivity than c-In₂O₃ which is advantageous for high capacity and fast



Fig. 7 Polyhedral representations of (a) $c\text{-}In_2O_3$ and (b) $h\text{-}In_2O_3,$ perpendicular to the c-axis.

charging and discharging processes.³⁴ Third, the large quantities of defects and smaller grain size in the $h-In_2O_3$ particles, as shown in Fig. 5d, not only provide additional sites for the insertion of Li ions, but also facilitate the diffusion of ions.³⁵ All these properties suggest that $h-In_2O_3$ is an outstanding candidate for anode materials of LIBs.

Conclusions

High-energy ball milling is a high yield, less hazardous and straightforward route to produce the $h-In_2O_3$ from $c-In_2O_3$ at room temperature. The structural transition is supported by XRD analysis, Raman measurements, and HRTEM examinations. The structural relationships and the possible transformation mechanisms can be explained in terms of the continuous topological variation of atom coordination or continuous displacement of the atoms induced by distortion and redistribution of InO_6 polyhedra units under high-energy milling impacts. As anode materials for lithium-ion batteries the $h-In_2O_3$ exhibits an enhanced and stable capacity of 390 mA h g⁻¹ (after 50 cycles) compared to $c-In_2O_3$ (49 mA h g⁻¹, after 30 cycles). The high energy ball milling method can induce solid-state phase transformations and produce new metastable structures with enhanced or new properties.

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Notes and references

- 1 D. Portehault, S. Cassaignon, E. Baudrin and J. P. Jolivet, J. Mater. Chem., 2009, 19, 7947.
- 2 J. K. Zhang, Y. H. Han, C. L. Liu, W. B. Ren, Y. Li, Q. L. Wang, N. N. Su, Y. Q. Li, B. H. Ma, Y. Z. Ma and C. X. Gao, *J. Phys. Chem. C*, 2011, **115**, 20710.
- 3 Z. Li, P. W. Yi, Q. Sun, H. Lei, H. L. Zhao, Z. H. Zhu, S. C. Smith, M. B. Lan and G. Q. Lu, *Adv. Funct. Mater.*, 2012, 22, 2387.
- 4 J. Lao, J. Huang, D. Wang and Z. F. Ren, *Adv. Mater.*, 2004, **16**, 65.
- 5 M. Bender, N. Katsarakis, E. Gagaoudakis, E. Hourdakis, E. Douloufakis, V. Cimalla and G. Kiriakidis, *J. Appl. Phys.*, 2001, **90**, 5382.
- 6 D. Liu, W. W. Lei, B. Zou, S. D. Yu, J. Hao, K. Wang, B. B. Liu, Q. L. Cui and G. T. Zou, *J. Appl. Phys.*, 2008, **104**, 083506.
- 7 C. H. Liang, G. W. Meng, Y. Lei, F. Phillipp and L. D. Zhang, *Adv. Mater.*, 2001, **13**, 1330.
- 8 C. Li, D. H. Zhang, S. Han, X. L. Liu, T. Tang and C. W. Zhou, *Adv. Mater.*, 2003, **15**, 143.
- 9 Q. S. Liu, W. G. Lu, A. H. Ma, J. K. Tang, J. Lin and J. Y. Fang, J. Am. Chem. Soc., 2005, 127, 5276.

- 10 Y. Idota, T. Kubota, A. Matsufuji, Y. Maekawa and T. Miyasaka, *Science*, 1997, **276**, 1395.
- 11 H. Li, X. J. Huang and L. Q. Chen, *Solid State Ionics*, 1999, 123, 189.
- 12 Y. Zhou, H. Zhang, M. Xue, C. Wu, X. Wu and Z. Fu, *J. Power Sources*, 2006, **162**, 1373.
- 13 W. H. Ho, C. F. Li, H. C. Liu and S. K. Yen, *J. Power Sources*, 2008, **175**, 897.
- 14 K. R. Prasad, K. Koga and N. Miura, *Chem. Mater.*, 2004, **16**, 1845.
- 15 M. Epifani, P. Siciliano, A. Gurlo, N. Barasan and U. Weimar, *J. Am. Chem. Soc.*, 2004, **126**, 4078.
- 16 C. H. Lee, M. Kim, T. Kim, A. Kim, J. Paek, J. W. Lee, S. Y. Choi, K. Kim, J. B. Park and K. Lee, *J. Am. Chem. Soc.*, 2006, **128**, 9326.
- 17 J. Q. Xu, Y. P. Chen, Q. Y. Pan, Q. Xiang, Z. X. Cheng and X. W. Dong, *Nanotechnology*, 2007, 18, 115615.
- 18 C. Chen, D. Chen, X. Jiao and C. Wang, *Chem. Commun.*, 2006, 4632.
- 19 C. T. Prewitt, R. D. Shannon, D. B. Rogers and A. W. Sleight, *Inorg. Chem.*, 1969, **8**, 1985.
- 20 T. Atou, K. Kusaba, K. Fukuoka, M. Kiduchi and Y. Syono, J. Solid State Chem., 1990, 89, 378; M. Epifani, P. Siciliano, A. Gurlo, N. Barsan and U. Weimar, J. Am. Chem. Soc., 2004, 126, 4078.
- 21 S. S. Farvid and P. V. Radovanovic, J. Am. Chem. Soc., 2012, 134, 7015.
- 22 D. Yu, S. H. Yu, S. Zhang, J. Zuo, D. Wang and Y. Qian, *Adv. Funct. Mater.*, 2003, **13**, 497.
- 23 Y. Chen, A. T. Halstead and J. S. Williams, *Mater. Sci. Eng., A*, 1996, **206**, 24.
- 24 D. Liu, W. W. Lei, Y. W. Li, Y. M. Ma, J. Hao, X. H. Chen, Y. X. Jin, D. D. Liu, S. D. Yu, Q. L. Cui and G. T. Zou, *Inorg. Chem.*, 2009, 48, 8251.
- 25 Powder Diffraction File No. 71-2194, M. Marezio, Acta Crystallogr., 1966, 20, 723.
- 26 W. B. White and V. G. Keramidas, *Spectrochim. Acta, Part A*, 1972, **28**, 501.
- 27 R. K. Gupta, K. Ghosh, S. R. Mishra and P. K. Kahol, *Mater. Lett.*, 2008, **62**, 1033.
- 28 Y. Chen, J. S. Williams and G. M. Wang, *J. Appl. Phys.*, 1996, **79**, 3956.
- 29 Y. Chen, T. Hwang, M. Marsh and J. S. Williams, *Metall. Mater. Trans. A*, 1997, 28, 1115.
- 30 M. Puttaswamy, Y. Chen, B. Jar and J. S. Williams, J. Metastable Nanocryst. Mater., 1999, 2, 79.
- 31 X. L. Li, H. F. Song, H. Wang, Y. L. Zhang, K. Du, H. Y. Li and J. M. Huang, J. Appl. Electrochem., 2012, 42, 1065.
- 32 Z. S. Wu, W. C. Ren, L. Wen, L. B. Gao, J. P. Zhao, Z. P. Chen, G. M. Zhou, F. Li and H. M. Cheng, ACS Nano, 2010, 4, 3187.
- 33 H. L. Wang, L. F. Cui, Y. Yang, H. S. Casalongue, J. T. Robinson, Y. Y. Liang, Y. Cui and H. J. Dai, *J. Am. Chem. Soc.*, 2010, **132**, 13978.
- 34 M. Sorescu, L. Diamandescu, D. Tarabasanu-mihaila and V. S. Teodorescu, *J. Mater. Sci.*, 2004, **39**, 675.
- 35 X. H. Huang, J. P. Tu, B. Zhang, C. Q. Zhang, Y. Li, Y. F. Yuan and H. M. Wu, *J. Power Sources*, 2006, **161**, 541.