Showcasing the work on binding of ruthenium intercalators to DNA by Anna K. F. Mårtensson and Per Lincoln, Department of Chemical and Biological Engineering, Chalmers University of Technology, Sweden.

Binding of Ru(terpyridine)(pyridine)dipyridophenazine to DNA studied with polarized spectroscopy and calorimetry

Achiral Ru(tpy)(py)dppz$^{2+}$ intercalated into the minor groove of DNA has similar intermolecular interactions as opposite enantiomers of its structural isomer, the "light-switch" complex Ru(bpy)$_2$dppz$^{2+}$.

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Binding of Ru(terpyridine)(pyridine)-dipyridophenazine to DNA studied with polarized spectroscopy and calorimetry†

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Linear and circular dichroism (LD and CD) spectroscopy as well as isothermal titration calorimetry (ITC) have been used to investigate the interaction of Ru(tpy)(py)dppz^2+ (tpy = 2,2'-terpyridyl; py = pyridine; dppz = dipyrido[3,2-a:2',3'-c]phenazine) with DNA, providing detailed information about the DNA binding thermodynamics and binding geometry of the metal complex. Flow LD, CD and isotropic absorption indicate that Ru(tpy)(py)dppz^2+ bind to DNA from the minor groove with the dppz ligand intercalated between base pairs, very similar to its chiral structural isomers Δ- and Λ-Ru(bpy)_2dppz^2+ (bpy = 2,2'-bipyridine). A simple cooperative binding model with one binding geometry provide an excellent fit for calorimetric and absorption titration data. The values of the neighbor interaction thermodynamic parameters for Ru(tpy)(py)dppz^2+ suggest that complexes bound contiguously prefer to have their tpy ligands oriented towards the same strand.

Introduction

DNA-binding drugs are small molecules that recognize and interact with specific DNA sites. Many of the chemotherapeutic anticancer agents currently in use fall under this category with cisplatin being the most prevalent example. DNA intercalators that unwind DNA in order to π-stack between two base pairs have shown cytotoxicity towards cancerous cells but are often of limited therapeutic use due to their lack of specificity and frequent side effects. After the pioneering work by Barton and coworkers on the selective DNA binding of substitution-inert trisphenanthroline complexes of ruthenium, there has been an increasing interest in octahedral transition metal complexes. The discovery of the “light switch” complexes Ru(phen)_2dppz^2+ and Ru(bpy)_2dppz^2+ (phen = 1,10-phenanthroline; bpy = 2,2'-bipyridine; dppz = dipyrido[3,2-a:2'3'-c]phenazine) lead to the synthesis of many variations of dppz-ruthenium-centered tris-bidentate structures with the potential as biosensors and therapeutic agents. Interestingly, the two ruthenium complexes that have reached clinical trials have substitution-labile ligands and a proposed mode of action completely different to DNA intercalation.

Spectroscopic and biophysical methods have established that it is the dppz ligand in tri-bidentate complexes that is intercalated between the base pairs of the DNA, and this has recently been confirmed by several X-ray crystal structures. Compared to the intense research on the brightly luminescent bipyridine and phenanthroline complexes, the very low quantum yield at room temperature has led to less interest for ruthenium dppz complexes carrying the tridentate tpy ligand (tpy = 2,2':6',2''-terpyridyl), despite the fact that the absence of a stereocentre at the metal when coordinated to a tridentate ligand eliminates the need of separating Δ and Λ racemic mixtures characteristic of tris-bidentate systems. The main focus has been on tpy-based complexes with reactive oxoruthenium(iv) functionality that cleaves DNA by oxidation of guanine and the 1'-deoxyribose hydrogen. However, not until recently has the single free coordination site left on the metal atom been recognized as a potential way of fine-tuning complex–DNA interactions, either by improving the DNA cleaving ability, or as a novel light-activated drug delivery system. Although some spectrophotometric studies have aimed at the DNA binding mode of tpy-based complexes to DNA, a definite proof of the binding geometry has not yet been obtained.

Our group has previously determined that the tris-bidentate complexes affect each other, either cooperatively or anti-cooperatively, when interacting with a DNA-polymer via intercalation. A binding model that gives a satisfactory fit to the data needs two distinct binding modes, one symmetrical (perpendicular) and one unsymmetrical (polar). This model with two modes of binding is supported by the crystallographic study by
Niyazi et al. (2012) where a symmetric and a non-symmetric intercalation geometry was found for Λ-Ru(phen)$_2$dpdz$^{2+}$.

In this study, we have carried out spectrophotometric and calorimetric measurements in order to determine the binding geometry and thermodynamic characteristics of Ru(tpy)(py)-dpdz$^{2+}$ (Ru-tpy, see Scheme 1). To simplify the system by avoiding effects from DNA sequence heterogeneity, and to compare with an earlier study, we chose to primarily study the interaction of poly(dAdT)$_2$ (AT-DNA). However, since the AT-DNA was too short to orient in the flow cell for the linear dichroism study, we used calf thymus DNA (ctDNA) instead. We wanted to characterize a tpy/dmpz ruthenium complex in its simplest form, and for the purpose of comparison, a pyridine (py) ligand was attached to the single coordination site, making the complex an achiral structural isomer of the original “light-switch” complex Ru(bpy)$_2$dpdz$^{2+}$ (Ru-bpy). Once we have gained more understanding in how this mononuclear complex interacts with DNA, the substituents of this single coordination site could be varied to optimize properties that would make it and its binuclear derivatives more suitable as metallopharmaceuticals.

Results

Absorption

The absorption spectra of ΔRu-bpy and Ru-tpy in the absence and presence of AT-DNA at [base pairs]/[Ru] ratio of 5 are illustrated in Fig. 1. The broad band system centered at about 440 nm for ΔRu-bpy and 475 nm for Ru-tpy is attributed to the metal to ligand charge transfer (MLCT) transitions, which in the presence of AT-DNA show a slight hypochromicity and red shift for both complexes. The characteristic 372 nm band, assigned to the lowest π→π* transitions of the dmpz chromophore is almost identical for both Ru-tpy and ΔRu-bpy, and in presence of AT-DNA there is also a very similar pronounced hypochromicity and red-shift (see Fig. 1 and Fig. S1 in ESI†).

For Ru-tpy, the most intense band outside the overlapping DNA absorption is the band at about 310 nm, which is assigned to the long-axis polarized lowest π→π* transition of the tpy chromophore. This peak shows a red-shift of about 5 nm and in contrast to the 372 nm dmpz band, only a slight hypochromicity at DNA binding, indicating that of the two large ruthenium ligands in Ru-tpy, the dmpz ligand is the one that has the closest interaction with the nucleobases. As shown by titrating a constant concentration of Ru-tpy with AT-DNA (see Fig. S2 in ESI†), the hypochromicity of the dmpz-band at 372 nm for Ru-tpy remains virtually constant at ratios [base pairs]/[Ru] > 2, whereas the initial hypochromicity in the 310 nm band is somewhat reduced at higher ratios.

Circular dichroism

A pure enantiomer of a chiral ruthenium complex, such as Ru-bpy, will show a strong intrinsic CD signal when free in aqueous solution, but Ru-tpy, which is an achiral complex will show zero CD signal under the same conditions. However, for both chiral and achiral molecules, binding to DNA can lead to a proper induced circular dichroism signal (proper ICD) by perturbation of the chromophores of the bound molecule by the chiral arrangement of the nuclebase chromophores. Here the magnitude and the sign of the ICD will mainly be dictated by the distances and angles between the interacting electronic transition dipole moments, and similar geometries are expected to give rise to similar ICD-patterns even for opposite enantiomers. In addition, for chiral molecules, changes merely in position and intensity of the intrinsic CD bands themselves, caused by the interaction with DNA, will add an apparent ICD contribution, but this will be characterized by a mirror-image like pattern for opposite enantiomers with similar binding geometries. Changes in CD were monitored upon addition of ΔRu-bpy, ΛRu-bpy, and Ru-tpy to a [base pairs]/[Ru] ratio of 5. In order to compare the ICD spectra of the complex–DNA interactions, the spectrum of the DNA and the spectrum of the free complex were subtracted from the spectrum in the presence of DNA. Fig. 2 shows the ICD of all three complexes as well as the CD spectrum of the AT-DNA. The general shape of the ICD below 290 nm is quite similar for all three complexes.
a negative band between 275 and 290 nm and a positive band at about 260 nm, indicative of a proper ICD mechanism, while the general mirror-image relationship of the ICD curve >290 nm for Δ- and ΔRu-bppy suggests the predominance of an apparent ICD mechanism in this region. Titration of a solution with constant concentration of Ru-tpy shows an almost invariant CD spectrum >300 nm for [base pairs]/[Ru] > 2, similar to the results of the corresponding absorption titration (see Fig. S3 in ESI†).

**Linear dichroism**

Fig. 3 shows the LD spectra for Ru-tpy in ctDNA solution at [base pairs]/[Ru] ratios 8, 4 and 2. The ctDNA concentration remained constant at 270 µM nucleotides. The reduced LD (LDr), which is the LD divided by the isotropic absorbance, remained constant at 270 µM nucleotides. The reduced LD (LDr), which is the LD divided by the isotropic absorbance, is very similar for all [base pairs]/[Ru] ratios, indicating little or no change in the geometric orientation at higher saturation levels (see Fig. S4 in ESI†). Thus, the conditions were fulfilled for determining the orientation factor S using eqn (5), (8) and (11) (See Theory and methods section below) using the LD spectra of free DNA (L0) and at ratios 8 (L1) and 4 (L2). The relative values S1/S0 = 1.06 and S2/S0 = 1.16 indicate that the DNA with bound Ru-tpy becomes better oriented, as earlier found for Δ- and ΔRu-bppY.11b Finally, the pure (without DNA contribution) Ru-tpy LD spectrum, at perfect orientation and 10 mm optical path-length, was calculated as b = c2−1(S2−1L1 − S0−1L0). The weights w1 and w2 in eqn (6) were varied manually until the ddppz band at 375 nm vanished in component e1 and the sharp tpy absorption band at 310 nm vanished in component envelope spectra e2, as shown in Fig. 4. The optimal weight values were w1 = 3 ± 0.5 and w2 = −1.5 ± 0.3, the theoretical limits for parallel and perpendicular orientation of a transition dipole moment relative to the orientation axis. Thus, the results show that the tpy long-axis (z) is aligned along the ddppz long-axis (x) perpendicular to the DNA helix axis, which, since the ddppz ligand lies in the x3y plane, is consistent with intercalation of the ddppz ligand in-between the base pairs.

**Binding isotherms**

ITC profiles for the binding of Ru-tpy to AT-DNA at 20, 25, and 30 °C are shown in Fig. 5 and for comparison, the corresponding ITC-profiles for Δ- and ΔRu-bppy from Andersson et al. (2013).10 Ru-tpy shows a similar overall shape with a gradually increasing exothermic enthalpy until a negative maximum is reached at [Ru]/[base pairs] = 0.4–0.5, somewhat higher than for the other two complexes. The initial slope of the ITC profile is in contrast to the initial constant part of the sigmooidal curve expected for the simple binding as indicated above by absorption, circular dichroism and linear dichroism spectroscopies, where the spectra of Ru-tpy in the presence of DNA were found to be practically invariant with binding ratio. This suggests that intermolecular interactions between bound molecules must contribute to the binding enthalpy also for Ru-tpy, as earlier concluded for Δ- and ΔRu-bppy.10 This
observation prompted us for a more thorough analysis of the spectroscopic changes, and absorption spectra were collected for addition of Ru-tpy to a constant concentration of AT-DNA at 25 °C.

The data (shown in Fig. S5 in ESI†) were analyzed with singular value decomposition (SVD) as described in Theory and methods. The first three (normalized) singular values were $s_1 = 100$, $s_2 = 2.45$ and $s_3 = 0.23$, and the corresponding columns of $U$ and $V$ are plotted in Fig. 6 and 7, respectively.

Ru-tpy ITC and absorption data could be excellently globally fitted with the classical McGhee–von Hippel cooperative binding model (see Theory and methods), and the best fit to the data are shown with solid lines in Fig. 5 and 7. Fig. 8 shows the calculated absorption spectra for free complex, bound complex with at most one neighbor and bound complex with at least one neighbor. For comparison, the model was also used to fit the ITC data for ΔRu-bpy and ΛRu-bpy, as shown in Fig. 5. A linear fit to the calculated Δ$H^\circ$ values as a function of temperature is shown in Fig. 9, and Table 1 gives the parameters and derived thermodynamic data.
saturation in both CD and absorption spectra. Likewise, linear dichroism spectra were found to be invariant with binding ratio too, allowing determination of the orientation factor $S$ and a quantitative analysis of the angular binding geometry. In contrast to the Ru(L)$_2$dpdz complexes, where major transition moment directions have oblique angles to the plane of the dpdz ligand, in Ru-tpy the long-axis polarized tpy transition is perpendicular to the dpdz plane. With a weight $w_1 = +3 \pm 0.5$ we find that it is almost perfectly parallel oriented to the DNA helix axis, and the dpdz long-axis polarized transition, with weight $w_2 = -1.5 \pm 0.3$, perpendicularly oriented, the geometry expected for intercalation of the dpdz ligand between the base pairs of DNA. In contrast to the very minute differences in absorption, CD and LD spectra at different binding ratios, the calorimetric titration show strong effects on the heat of binding. We have previously been reported such non-classical ITC curves for Ru(L)$_2$dpdz complexes ($L = \text{bpy or phen}$) and attributed them to an additional enthalpy contribution from interaction between neighboring complexes on the DNA.\textsuperscript{10} In comparison to $\Delta$- and $\Lambda$Ru-bpy, Ru-tpy showed an ITC profile qualitatively most similar to the latter. Since a satisfactory global fit for the ITC and absorption experimental data of Ru-tpy could be obtained with the classical McGhee-von Hippel model with only one type of binding geometry, for comparison this model was also used to reanalyse our data for $\Delta$- and $\Lambda$Ru-bpy, which were originally fitted with a symmetrical and a pair of unsymmetrical intercalation geometries.\textsuperscript{10} In this model (Model 3 in ref. 10), ligands bound with only one nearest-neighbor are assumed to have one distinct binding geometry (unsymmetrical), while ligands bound either isolated or with nearest-neighbors on both sides have a second binding geometry (symmetrical). As noted in our previous study, the fit of the simpler model to the ITC-data alone is excellent, but any attempt to rationalize the observed molar fractions of the short and the long excited state lifetime fails. With the simple model, the apparent site sizes for $\Delta$- and $\Lambda$Ru-bpy were found to be 2.2 and 2.3, consistent with the neglect of the anti-cooperativity which is inherent in Model 3 for which the $n$ parameters were found to be 2.0 and 1.8, respectively.\textsuperscript{10} Interestingly, the binding site size parameter $n$ was found to be 2.0 for Ru-tpy, indicating that the distinction between symmetrical and unsymmetrical intercalation geometries might be less pronounced for this complex. Although the low emission quantum yield of Ru-tpy bound to DNA precluded time-resolved luminescence measurements, the observation that spectra (absorption, CD, LD) change very little with binding density (see Fig. 3, S2 and S4) support the conclusion that difference between intercalation geometries is small for Ru-tpy. The simple model has only one cooperativity parameter $y$, which is found to be 1 for $\Delta$Ru-bpy (i.e. non-cooperative binding) and 5.5 for $\Lambda$Ru-bpy (cooperative binding); for Ru-tpy $y = 2.8$, in-between the values of the two Ru-bpy enantiomers. The intrinsic binding constant $K$ is 10$^6$ M$^{-1}$ for Ru-tpy, very similar to that of $\Delta$Ru-bpy, while $K$ for $\Lambda$Ru-bpy is almost 50 times smaller. The latter value is about 3 times smaller than that obtained with Model 3, however,"
since simulated binding isotherms are the most sensitive to the value of the binding constant close to saturation, the best-fit value will normally be quite dependent on the binding model since the influence of the cooperativity parameters will differ.

The intrinsic binding enthalpy $\Delta H^\circ_b$ is positive for all complexes at 20 °C, i.e. the binding in the absence of neighbor interactions is endothermic. The value of $\Delta H^\circ_b$ is smaller for Ru-tpy than for either $\Delta^-$ or $\Lambda$Ru-bpy, but the temperature dependence is similar as evidenced by the negative $\Delta C_p$-values of $-680 \pm 120$ J K$^{-1}$ M$^{-1}$. The nearest-neighbor interaction enthalpy $\Delta H^\circ_{nn}$ is negative (i.e. exothermic) for all complexes, and exhibits a tendency similar to the cooperativity factor $y$, namely that the value for Ru-tpy is in-between the values for $\Delta^-$ and $\Lambda$Ru-bpy.

However, even if a simple, single binding geometry model is sufficient to model the binding data for Ru-tpy, the lack of a 2-fold axis of symmetry, in contrast to $\Delta^-$ and $\Lambda$Ru-bpy, make 2 types of neighbor interaction possible: either the tpy ligands of two consecutive complexes are oriented towards the same strand ($T_{SS}$, see Fig. 10) or they are oriented towards opposite (alternating) strands ($T_{AS}$). In comparison to Ru-bpy, neglecting the influence of the single pyridine ring of Ru-tpy, the alternating model ($T_{AS}$) would have intermolecular contacts resembling alternating $\Delta^-$-$\Delta^-$ and $\Lambda^-$-$\Lambda^-$ contacts, while the same-side model ($T_{SS}$) would everywhere have intermolecular contacts resembling $\Delta^-$-$\Lambda^-$, comparable to that found in a recent X-ray crystal structure of $\Delta^-$ and $\Lambda$-Ru(phen)$_2$dppz simultaneously intercalated to a DNA hexamer duplex.$^5$

Since the same-side model $T_{SS}$ has only one type of intermolecular contact ($\Delta^-$-$\Lambda^-$), it is logically consistent with the simple one-geometry binding model, and also consistent with the assumption that Ru-bpy interactions model those of Ru-tpy, since 7.8, the square of the Ru-tpy $y$ value of 2.8, is larger than the product 5.5 of the corresponding $\Delta^-$-$\Delta^-$ and $\Lambda^-$-$\Lambda^-$ $y$ values for Ru-bpy. However, the values 7.8 and 5.5 are similar enough in magnitude to suggest that both $T_{SS}$ and $T_{AS}$ arrangements will be significant for Ru-tpy, even if the former arrangement with tpy ligands oriented towards the same strand will predominate. Our results indicating a relatively modest cooperativity factor ($y = 2.8$) for Ru-tpy appears to be in contrast to the case of Ru(phen)$_2$dppz$^{2,5}$, for which a rather substantial $\Delta^-$-$\Lambda^-$ cooperativity can be inferred, since Cardin and coworkers report that the hexamer duplex used in their crystal structure study preferentially binds precisely one $\Delta^-$ and one $\Lambda$-Ru(phen)$_2$dppz in solution.$^5$

**Experimental**

**Materials**

All experiments were performed in aqueous solution (pH = 7.0) containing 150 mM NaCl and 1 mM cacodylate (dimethylarsinic acid sodium salt) except for the LD experiments where 10 mM NaCl was used (pH = 7.0). This was because a lower salt concentration gave a higher signal intensity without affecting the overall shape of the spectra. Stock solutions of calf thymus DNA (ctDNA) (−5 mM nucleotides) were prepared.
by dissolving highly polymerized type I sodium salt calf thymus DNA (Sigma-Aldrich) in buffer. A stock solution of poly-(dAdT)2 (AT-DNA) (~5 mM nucleotides) was prepared by dissolving the sodium salt (Sigma-Aldrich) in buffer. The solutions were filtered two times through a 0.7 µm polycarbonate filter. Stock solutions of the complexes (~1 mM) were prepared by dissolving the chloride salts in buffer. Concentrations were determined spectrophotometrically using extinction coefficients: ε258 = 6600 M⁻¹ cm⁻¹ per nucleotide for ctDNA, ε260 = 6600 M⁻¹ cm⁻¹ per nucleotide for AT-DNA, ε373 = 16 900 M⁻¹ cm⁻¹ for Ru-tpy and ε444 = 16 100 M⁻¹ cm⁻¹ for Ru(bpy)2dppz2+ (Ru-bpy). For ITC measurements the DNA solution was dialyzed against pure buffer for at least 48 hours at 8 °C. Ruthenium complex solutions of appropriate concentrations were prepared by dialution of the stock solutions in the dialysate. The dialysis membrane used had a molecular weight cut-off of 3.5–5 kDa (Spectra-Por® Float-A-Lyzer® G2, Sigma Aldrich).

Δ- and Λ-[Ru(bpy)4dppz]Cl2 used in this study were prepared as previously reported.11a Other chemicals were purchased from Sigma-Aldrich and used without purification.

**Synthesis of [Ru(tpy)(py)dppz]Cl2**

The synthetic route for preparation of [Ru(tpy)(py)dppz]Cl2 is shown in Scheme 2. The procedure for preparation of Ru(tpy)-Cl2, and [Ru(tpy)(dppz)]Cl are in accordance to the methods previously reported by Zhou et al. and Leising et al., respectively.7a,14 [Ru(tpy)(dppz)2(py)][PF6]2 was synthesized using the method previously described by Zhou et al. with some modifications. A portion of 0.0557 g of [Ru(tpy)(dppz)2]Cl and 0.0270 g of AgNO3 were refluxed in 20 mL of ethanol–water (1:1) for 3 h under N2(g). The solution was filtered after cooling and the filtrate was refluxed again under N2 (g) for another 4 h, with 0.0101 g pyridine (py) added. The solution was left in a fridge overnight for cooling. The next day the precipitate was filtered and washed with ethanol–ether (1:2). Purification of [Ru(tpy)(dppz)2(py)][PF6]2 was done using column chromatography with CH3CN and neutral Al₂O₃. The eluate containing the pure orange product was collected leaving a dark purple residue layer on top of the column.

To replace the hexafluorophosphate anion with chloride, the CH₃CN solution of [Ru(tpy)(dppz)2(py)][PF6]2 was reduced to ~1 mL evaporating with a stream of N₂ under mild heating, where after 0.5 g of [([CH₃(CH₂)₃]NCl), dissolved in 1 mL of acetone, was added in increasing portions while stirring until the solution was only weakly yellow and the precipitation complete. The product was collected by a sintered glass filter and washed first with acetone and then with diethyl ether to yield [Ru(tpy)(dppz)2(py)]Cl2 as a brown powder (36%, calculated from the starting material, Ru(tpy)Cl2). UV/vis (in water; λmax in nm, ε10³ M⁻¹ cm⁻¹ enclosed in parenthesis): 475(12.1), 372(16.9), 310(45.4), 275(69.2). 1H NMR (as PF6 salt, 400 MHz, acetone-d6): δ 10.03 (dd, J = 8.2, 1.3 Hz, 1H), 9.56–9.52 (m, 1H), 9.48 (dd, J = 5.4, 1.4 Hz, 1H), 8.94 (t, J = 8.0 Hz, 2H), 8.84–8.71 (m, 2H), 8.59–8.42 (m, 4H), 8.26–8.11 (m, 9H), 8.04–7.94 (m, 2H), 7.76 (ddd, J = 8.2, 5.5, 0.8 Hz, 1H), 7.49–7.42 (m, 3H). The 1H NMR spectrum of Ru-tpy (see Fig. S8 in ESI†) was in accordance with previous results by Zhou et al. (2009)7a and showed no significant impurities.

**Spectroscopy**

Absorption spectra were measured on a Varian Cary 4000 UV/vis spectrophotometer (path length = 1 cm). The reverse absorption titration spectra for Ru-tpy were measured with a constant concentration of 5 µM Ru-tpy in buffer. The stock solution of AT-DNA (with 5 µM of Ru-tpy to avoid dilution) added directly in the cuvette (path length = 1 cm) in aliquots up to a concentration of 80 µM nucleotides.

Linear dichroism spectra (LD) were measured on a Chirascan LD spectropolarimeter on samples oriented in an outer-rotating Couette flow cell with a 1 mm path length at a rate of 1000 rpm. The spectra of the same samples were recorded without rotation for baseline contribution and were subsequently subtracted from the LD spectra. The concentration of ctDNA used in the measurements was 266 µM nucleotides and mixed with complex solution with appropriate concentration to obtain desired [Ru]/[base pairs] ratios. Circular dichroism (CD) spectra were recorded on a Chirascan CD spectropolarimeter similarly as the absorption spectra. Five CD spectra were averaged for each sample. To be consistent with the absorption and ITC measurements, AT-DNA was used for the CD measurements (~138 µM nucleotides).

Calorimetric data was obtained using an ITC200 isothermal titration calorimeter (Microcal) controlled by Origin 7.0 software. The ruthenium complexes (~600 µM) were loaded in a syringe (40 µL) and titrated at 2 µL aliquots into 206 µL of AT-DNA in 150 mM NaCl aqueous solution (~340 µM nucleotides). By integrating the power required to maintain the reference and sample cells at the same temperature it is possible to obtain a direct measurement of the heat generated or absorbed when complex and DNA interact. The experimental raw data consists of a series of heat flow peaks, and each peak...
corresponds to one injection of complex. These heat flow spikes are integrated with respect to time, which gives the total heat exchanged per mole injectant, plotted against the ratio [Ru]/[base pairs]. The primary ITC data was corrected for the heat of ligand dilution by subtracting the average heat per injection of complex titrated into buffer. There was negligible heat arising from DNA dilution. The experiments were performed at 20 °C, 25 °C and 30 °C.

\( ^1H \) NMR spectrum of [Ru(tpy)(py)dppz](PF_6)_2 in acetone-\( d_6 \) was recorded on an Agilent 400 MHz spectrophotometer.

Theory and methods

Analysis of LD spectra

The isotropic absorption spectrum \( A_{iso} \) of a sample can be considered as a simple sum of component spectral envelopes \( e_i \), each such envelope itself being a collecting all electronic transitions with a common polarisation direction relative to the molecular coordinate system:

\[
A_{iso} = \sum_i e_i
\]

Then, the linear dichroism spectrum LD will be a weighted sum of the same component spectral envelopes:

\[
LD = \sum_i w_i e_i
\]

there for an effectively uniaxial oriented system (e.g. DNA oriented by the shear flow in a Couette cell) the weights \( w_i \) are determined by the orientation of the polarization direction characteristic for the electronic transition dipole moments of \( e_i \) relative to the macroscopic orientation axis of the system:

\[
w_i = \frac{\gamma}{2} S(3\cos^2\theta_i - 1) \quad 0 \leq S \leq 1
\]

where \( \gamma \) is the ratio of the optical path-length of the LD to the absorption measurement, \( \theta_i \) is the angle between the transition dipole moment direction of envelope \( e_i \) and the orientation axis of the molecular system (in our case the DNA helix axis), and the orientation factor \( S \) describes the degree of alignment of the molecular system orientation axis relative to the macroscopic orientation axis, with \( S = 0 \) for an unoriented and \( S = 1 \) for a perfectly aligned sample.

The reduced linear dichroism LD’ value at a certain wavelength is defined as the ratio of the linear dichroism value over the isotropic absorbance value:

\[
LD' = \frac{LD}{A_{iso}}
\]

In wavelength regions where a single component spectral envelope \( e_i \) dominates the absorption spectrum, the LD’ curve will be essentially constant and take the value \( w_i \), e.g. as observed around the 260 nm band of B-DNA. Since \( \theta \) for the

in-plane polarized \( \pi \rightarrow \pi^* \) nucleobase transitions here is close to 90°, the orientation factor \( S \) for ligand-free DNA can readily be calculated from the LD’ value.

\[
LD'_{260} = w = -\frac{\gamma}{2} S_0
\]

However, in most cases, and in particular for 3D-chromophores like ruthenium polypyridyl complexes, component spectral envelopes overlap substantially over the whole range of the spectrum, and the LD’ curve will vary strongly with wavelength. Deconvolution of the experimental \( A_{iso} \) and LD spectra into component envelopes, and thereby determining the weights \( w_i \), can in favorable cases be accomplished by the TEM-method.\(^5\)\(^6\)\(^7\) For a system with two component envelope spectra, having distinct characteristic absorption band features and distinct weights, the TEM-method in matrix notation can be formulated as the solution of the combined eqn (1) and (2):

\[
[a \ b] \begin{bmatrix} 1 & w_1 \\ 1 & w_2 \end{bmatrix}^{-1} = [e_1 \ e_2]
\]

where \( a, b, e_1 \) and \( e_2 \) are column vectors corresponding to the isotropic absorption, linear dichroism and the two component envelope spectra, respectively. The variables \( w_1 \) and \( w_2 \) in eqn (6) are varied until, as determined by visual inspection, the characteristic features of component \( e_2 \) have vanished in component \( e_1 \) and vice versa.

The orientation factor \( S \) is required for the angular orientation to be determined from the weights \( w_i \). For a dye that (a) is strongly bound to DNA; (b) has significant absorption at wavelengths >300 nm (where the DNA is transparent); and (c) has, within a certain range of [dye]/[DNA] ratios, invariant binding geometry and invariant absorption spectrum; \( S \) can readily be obtained as follows:

Let two LD spectra with different [dye]/[DNA] ratios (within the invariant range) be column vectors \( L_1 \) and \( L_2 \), and \( L_0 \) be the LD spectrum of a sample with DNA only. Then, if condition (c) is fulfilled, the three columns are linearly dependent, thus scalars \( \alpha \) and \( \beta \) can be found so that \( \alpha L_1 + \beta L_2 = L_0 \); in matrix notation:

\[
Mx = L_0
\]

where \( M = [L_1 \ L_2] \) and \( x = [\alpha; \beta] \). Solving eqn (7) by a least square projection gives \( \alpha \) and \( \beta \):

\[
(M^T M)^{-1} M^T L_0 = x
\]

When condition (a) is fulfilled, practically all added dye can be considered to be bound; and with the DNA concentration being equal in the three samples, the vectors can be written as:

\[
L_0 = \gamma S_0 d \quad L_1 = \gamma S_1 (d + c_1 b) \quad L_2 = \gamma S_2 (d + c_2 b)
\]

where (with \( \gamma = 1 \)), \( d \) is the LD spectrum at perfect orientation of DNA only, \( b \) is the LD spectrum at perfect orientation of bound dye at unit concentration, and \( c_1 \) and \( c_2 \) are
the two known total dye concentrations. Since condition (b) assures that \( d \) is zero and \( b \) is non-zero for wavelengths > 300 nm, the following two equations are obtained by inserting eqn (9) in (7):

\[
\alpha S_1 + \beta S_2 = S_0 \quad c_1 \alpha S_1 + c_2 \beta S_2 = 0
\]  

(10)

Since \( \gamma S_0 \) can be evaluated from eqn (5), \( S_1 \) and \( S_2 \) can be obtained from known quantities after solving the two equations in (10):

\[
S_1 = \frac{S_0 c_2}{\alpha (c_2 - c_1)} \quad S_2 = \frac{S_0 c_1}{\beta (c_2 - c_1)}
\]  

(11)

Analysis of binding isotherms

Although global fitting of calorimetric and excited state lifetime data for \( \Delta \) and \( \Lambda \)-Ru-bpy and their 1,10-phenanthroline analogues required a fairly complicated binding model comprising a symmetrical and a pair of unsymmetrical intercalation geometries (Model 3), fits to calorimetric data only were found to be satisfactory with a model with a single binding geometry.\(^{10} \) This model, the classical McGhee–von Hippel single ligand cooperative binding model was used also in this work, and involves three adjustable parameters: the thermodynamic binding constant \( K \), the binding site coverage parameter \( n \) and the cooperativity factor \( \gamma \).\(^{16} \) Given values of these three parameters (assumed to be constant in the small temperature range used here), and total concentrations of binding sites \([B]\)\(_{\text{tot}}\) (i.e. base pairs) and DNA-ligand \([L]\)\(_{\text{tot}}\) (i.e. ruthenium complex) for each step of the titration, the mass-balance equations were solved iteratively with a Newton–Raphson procedure, to give consistent binding densities \( \theta \) and free ligand concentrations \([L]\)\(_{\text{free}}\), as well as the conditional probabilities \( p_{ij} \), as previously described.\(^{10,17} \)

Calorimetric titration

The ITC data obtained at a certain temperature was assumed to be composed of three components:

\[
\text{ITC}(i) = (\Delta H_0)_{\text{b}}(i) + (\Delta H_{\text{nn}})_{\text{b}}(i) + \Delta H_{\text{baseline}}
\]  

(12)

where \( \Delta H_0 \) is the standard binding enthalpy change, \( \Delta H_{\text{nn}} \) is the standard nearest neighbor interaction enthalpy change and \( \Delta H_{\text{baseline}} \) is a small constant value. For titration data point \( i \), the change in concentration of bound ligand \( \Delta b(i) \) is calculated as \([B]_{\text{tot}}(i) \theta(i) - \theta(i - 1)] \) and the change in concentration of nearest neighbors \( \Delta n \) is calculated as \([B]_{\text{tot}}(i) \theta(2p_{22}(i) - \theta(i - 1)p_{22}(i - 1)] \), where \( p_{22} \) is the conditional probability that a bound ligand is followed by another bound ligand on the DNA lattice. The \( \Delta H \) values were determined by projecting the matrix with the ITC data as columns, one for each temperature, on the space spanned by the \( \Delta b, \Delta n \) and a constant vector. The simulated ITC curves were then calculated with eqn (12) and the goodness-of-fit determined as the Euclidian norm of the difference between the measured and simulated data matrices.

Spectroscopic titration

The absorption spectra were arranged in data matrix \( M \) as columns with \( w \) elements corresponding to the wavelengths recorded; the columns corresponding to the \( t \) titration steps. Singular value decomposition, using the \texttt{svd} command in the MATLAB software, factorized the data matrix into two matrices of orthonormal columns \( U \) and \( V \), and a diagonal matrix \( S \) with the singular values \( s_1 \geq s_2 \geq s_3 \cdots \geq 0 \) along the diagonal. Keeping only the \( m \) singular values that are significantly larger than zero make it possible to simplify the factorization:

\[
M = USV^T = USV^T + N
\]  

(13)

where \( S \) is the upper \( m \) by \( m \) submatrix of \( S; U \) and \( V \) are the first \( m \) columns of \( U \) and \( V \), respectively, and \( N \) is a matrix of small elements, ideally corresponding only to the noise in the measurements. Assuming thus that the titration involves \( m \) different absorbing species, with absorption spectra in \( w \) by \( m \) matrix \( A \) and concentrations in \( t \) by \( m \) matrix \( C \), the factorization can now be written as

\[
M - N = USC^TV = AC^TV = U[R]^{-1}SV^T = U[R]^{-1}SV^T
\]  

(14)

where \( R \) is a non-singular \( m \) by \( m \) matrix. Since \( A = UR \) and \( C = VS[R^{-1}]^T \), any theoretical binding model, which can calculate a matrix \( C_m \) of concentrations of \( m \) species, can be evaluated for consistency with the data by finding the \( R \) that is the least square projection of \( VS \) on the space spanned by \( C_m \):

\[
R = (C_m^T C_m)^{-1} C_m^T VS
\]  

(15)

By varying the adjustable parameters of the theoretical model, \( C_m \) is varied to minimize the norm of the residual \( VS - C_m R \), while keeping the specie spectra in \( A = UR \) physically reasonable (non-negative in the case of absorption spectra). In the present case, the concentration matrix \( C_m \) was constructed with three columns corresponding to free ligand, bound ligand with a neighboring free binding site, calculated as \([B]_{\text{tot}}(1 - p_{22}) \), and bound ligand adjacent to another bound ligand, calculated as \([B]_{\text{tot}}p_{22} \).

Conclusion

The angular orientation of the dppz and tpy ligands as determined by LD and CD spectroscopies, strongly indicate that Ru-tpy binds from the minor groove by intercalating the dppz ligand between the base pairs, as has previously been determined for the isomeric complexes \( \Delta \) and \( \Lambda \)-Ru-bpy. The strong hypochromicity in the dppz absorption band is almost identical in magnitude to that of \( \Delta \)-Ru-bpy, indicating that the alignment in the intercalation pocket is similar, too.

A simple cooperative binding model with one symmetrical binding geometry provide an excellent fit to data for calorimetric and absorption titrations of Ru-tpy into AT-DNA. The intrinsic intercalation has an equilibrium constant of \( 10^6 \text{ M}^{-1} \), close to that of \( \Delta \)-Ru-bpy, and is associated with a small endothermic enthalpy change of +3 kJ M\(^{-1}\). The cooperativity
factor is 2.8, and the associated neighbor interaction enthalpy is exothermic, $-10.2$ kJ M$^{-1}$; these values being in-between those of $\Delta^0$- and $\Lambda$-Ru-bpy, and consistent with a slight preference of a one-sided arrangement of tpy-ligands of complexes consecutively bound to DNA.

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**Notes and references**