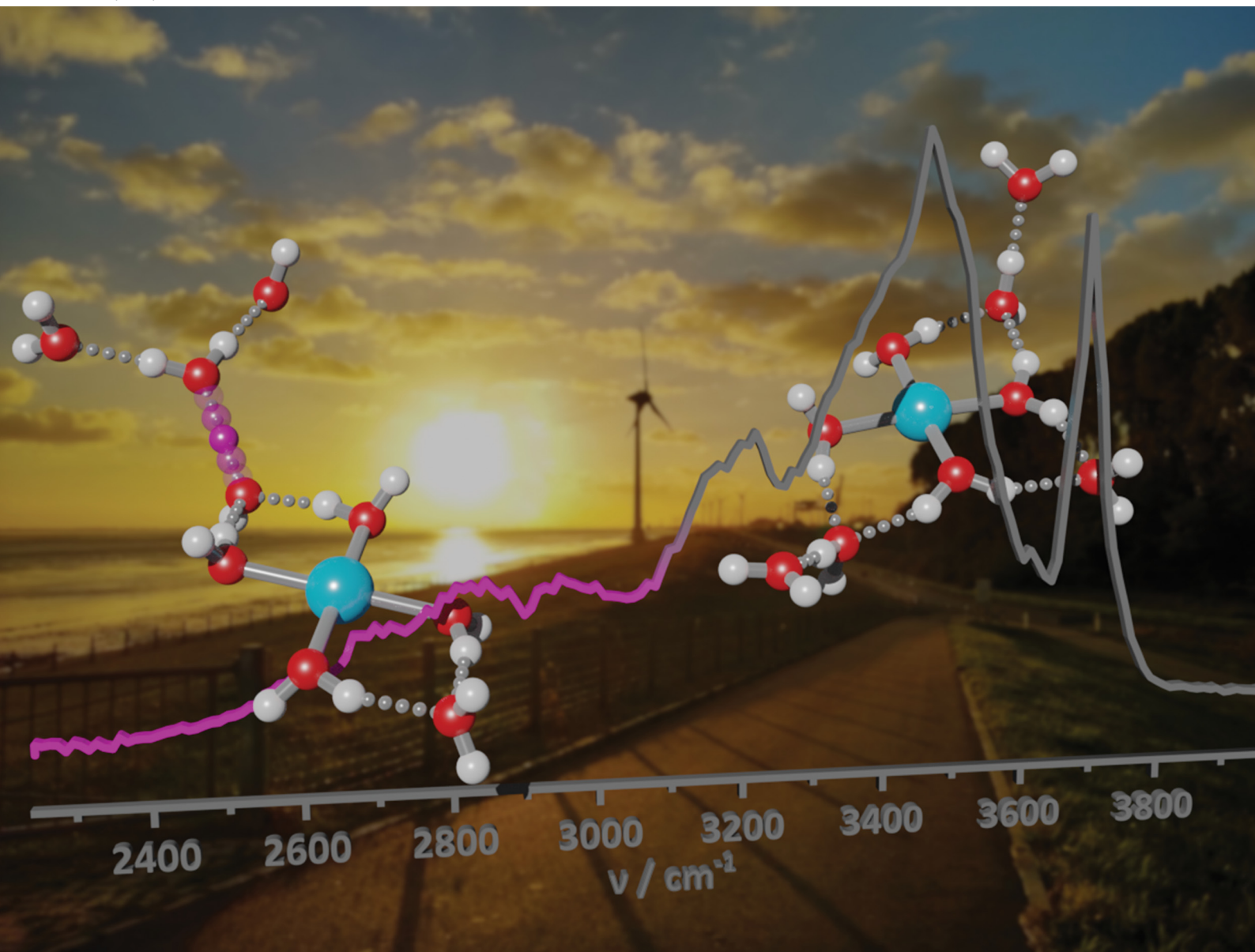


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**PAPER**

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# Size-dependent H and H<sub>2</sub> formation by infrared multiple photon dissociation spectroscopy of hydrated vanadium cations, V<sup>+</sup>(H<sub>2</sub>O)<sub>n</sub>, n = 3–51†‡

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Infrared spectra of the hydrated vanadium cation (V<sup>+</sup>(H<sub>2</sub>O)<sub>n</sub>; n = 3–51) were measured in the O–H stretching region employing infrared multiple photon dissociation (IRMPD) spectroscopy. Spectral fingerprints, along with size-dependent fragmentation channels, were observed and rationalized by comparing to spectra simulated using density functional theory. Photodissociation leading to water loss was found for cluster sizes n = 3–7, consistent with isomers featuring intact water ligands. Loss of molecular hydrogen was observed as a weak channel starting at n = 8, indicating the advent of inserted isomers, HVOH<sup>+</sup>(H<sub>2</sub>O)<sub>n–1</sub>. The majority of ions for n = 8, however, are composed of two-dimensional intact isomers, concordant with previous infrared studies on hydrated vanadium. A third channel, loss of atomic hydrogen, is observed weakly for n = 9–11, coinciding with the point at which the H and H<sub>2</sub>O calculated binding energies become energetically competitive for intact isomers. A clear and sudden spectral pattern and fragmentation channel intensity at n = 12 suggest a structural change to inserted isomers. The H<sub>2</sub> channel intensity decreases sharply and is not observed for n = 20 and 25–51. IRMPD spectra for clusters sizes n = 15–51 are qualitatively similar indicating no significant structural changes, and are thought to be composed of inserted isomers, consistent with recent electronic spectroscopy experiments.

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## Introduction

Hydrated ions are ubiquitous to many important biological and chemical processes.<sup>1</sup> Studying gas-phase hydrated ions provides a well-defined system to investigate fundamental ion–water behavior at the molecular level.<sup>2–6</sup> Important insights include solvation, hydrogen production, and corrosion mechanisms in which gas-phase studies offer the promise of bridging the gap between small clusters and bulk aqueous solution.<sup>7–14</sup>

Early experiments by Armentrout investigated the successive binding energies of water molecules to first row transition metals by collision induced dissociation (CID).<sup>15</sup> Size-dependent fragmentation behavior of a range of hydrated metal ions has been studied by blackbody infrared radiative dissociation (BIRD),<sup>16–19</sup> along with proton transfer reactions employing H<sub>2</sub>O/D<sub>2</sub>O exchange.<sup>20,21</sup> BIRD experiments on V<sup>+</sup>(H<sub>2</sub>O)<sub>n</sub> complexes

(n = 5–30) reveal strongly size-dependent intracluster redox reactivity. Vanadium was found in oxidation state +II for cluster sizes n = 9–12, resulting in V(OH)<sup>+</sup>(H<sub>2</sub>O)<sub>n</sub> ions and loss of atomic hydrogen, and the +III state for cluster sizes n = 9–23, leading to V(OH)<sub>2</sub><sup>+</sup>(H<sub>2</sub>O)<sub>n</sub> ions and loss of H<sub>2</sub>; all other cluster sizes show only water fragmentation.<sup>19</sup>

A wealth of spectroscopic studies have been performed on gas-phase solvated metal ions, including electronic spectroscopy,<sup>22–34</sup> with many more employing infrared spectroscopy, a powerful technique used to reveal structural trends of metal–ligand clusters.<sup>1,35–58</sup> The hydrated vanadium cation, however, has taken center stage and has been studied in detail by many groups using different experimental techniques. Laser ablated vanadium, when co-deposited with water molecules in an argon matrix, reacts spontaneously forming the inserted complex, HVOH.<sup>59</sup> Sasaki *et al.* investigated small V<sup>+</sup>(H<sub>2</sub>O)<sub>n</sub> complexes (n = 2–8) employing the inert-messenger technique, revealing a square-planar coordination of the vanadium center with four water molecules.<sup>58</sup> Duncan and coworkers have performed many infrared studies on solvated vanadium complexes, the first in 2003 on V<sup>+</sup>(H<sub>2</sub>O)Ar<sub>n</sub> and V<sup>+</sup>(D<sub>2</sub>O)Ar<sub>n</sub> complexes in the O–H stretching region.<sup>60</sup> The doubly-charged species, V<sup>2+</sup>(H<sub>2</sub>O)Ar<sub>n</sub> (n = 2–7), was also studied, revealing a red-shift

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in absorption frequencies compared to  $V^+(H_2O)$ , and a coordination of six to  $V^{2+}$ ; one water molecule and five argon atoms.<sup>47</sup> Later, the Duncan group also recorded the infrared spectra of  $V^+(H_2O)$  with different rare-gas messenger atoms, reporting evidence of *ortho*-*para* conversion on the water molecule.<sup>61</sup> More recently, Duncan and coworkers extended these infrared studies to larger clusters, investigating the microsolvation of  $V^+$  up to 30 water molecules.<sup>62</sup> Small clusters were probed using argon tagging, while larger clusters,  $n > 7$ , were studied *via* the photodissociation of intact water molecules. The spectra, when coupled with simulated spectra from density functional theory, present a coordination number of four, concordant with a square-planar structure. Clusters up to  $n = 8$  remain two-dimensional, however in larger clusters hydrogen bonding networks dominate, forming three-dimensional structures.

Recently, we performed electronic photodissociation experiments in the ultraviolet/visible (UV/Vis) region on  $V^+(H_2O)_n$  ions ( $n = 1-41$ ).<sup>63</sup> Clusters up to 12 water molecules exhibit intense 3d-4p transitions which red-shift until the first solvation sphere is reached at  $n = 4$ . At 9 water molecules, the intense absorption bands begin to disappear as  $V^+$  inserts into the O-H bond of a water molecule, forming  $[HVOH(H_2O)_{n-1}]^+$  whereby the oxidation state of V changes from +I to +III. Loss of water molecules, along with competing loss of atomic and molecular hydrogen, is observed for  $n \leq 12$ .

To confirm the presence of specific structural isomers inferred from the UV/Vis spectra, this study utilizes infrared multiple photon dissociation (IRMPD) spectroscopy in the O-H stretching region investigating the solvation evolution of  $V^+(H_2O)_n$  ions up to  $n = 51$ . The goal is to establish the presence of inserted isomers at particular clusters sizes and any possible infrared-driven H and  $H_2$  loss. The infrared-driven elimination of molecular hydrogen was recently observed for particular  $HALOH^+(H_2O)_{n-1}$  clusters,  $n = 9-14$ , probed in the 1400–2250  $cm^{-1}$  region.<sup>64</sup> Presence of a hydride was found to be crucial in the elimination of molecular hydrogen, and found to occur only for cluster sizes  $n = 12, 13$  and 14.

## Experimental and computational methods

The experiments are performed on a 4.7 Tesla Fourier-Transform Ion Cyclotron Resonance Mass Spectrometer (FT-ICR MS) Bruker Spectrospin CMS47X.<sup>65-68</sup> The mass spectrometer is equipped with a Bruker infinity cell<sup>69</sup> along with a laser vaporization source,<sup>70,71</sup> where a solid disk of vanadium is vaporized by a frequency doubled Litron Nano S 60-30 Nd:YAG laser (532 nm, 5 mJ per pulse, 30 Hz). The plasma containing  $V^+$  is entrained in a pulse of helium seeded with water vapor created *via* a homebuilt piezoelectric valve following the design of Proch and Trickl.<sup>72</sup> The ensuing pulse is cooled *via* supersonic expansion, whereby hydrated vanadium complexes,  $V^+(H_2O)_n$ , are created. The clusters are transferred into the ICR cell where they are stored and mass-selected within the 4.7 T magnetic field<sup>73</sup> under ultra-high vacuum conditions ( $p \approx 5 \times 10^{-10}$  mbar).

The mass-selected ions of interest are irradiated with infrared light provided by an EKPLSA NT277 Optical Parametric Oscillator (OPO) covering the range from 2240–4000  $cm^{-1}$ . Typical irradiation times are 0.1–0.2 s at 1000 Hz repetition rate. The normalized IRMPD yield is calculated from the precursor ion and fragment ion intensities<sup>74</sup> and scaled with the laser power. Ion intensities are normalized to 100% precursor ion intensity. The fragment intensities are then BIRD corrected by subtracting the fragment intensities from a control experiment without laser irradiation. IRMPD yields are calculated from the corrected ion intensities, normalized by the wavelength-specific laser power and re-normalized for display as described before.<sup>75</sup>

The laser power is measured after every mass spectrum to account for any fluctuations. The laser power drops in the range of 3500–3520  $cm^{-1}$ . Due to the complex kinetics of the IRMPD process, this may lead to small artefacts in the IRMPD yield at these wavelengths even after power correction.

Especially for larger clusters, dissociation due to BIRD from the cell walls can occur. To prevent this, the ICR cell is surrounded by a copper jacket, whereby the ions can be cooled with liquid nitrogen to a temperature of *ca.* 90 K,<sup>16,76</sup> minimizing the effects of BIRD.<sup>77-82</sup> The remaining BIRD fragmentation was taken into account by subtraction of the measured fragment ion intensities with a reference mass spectrum where the ions were trapped without irradiation from the OPO.

To complement the experimental photodissociation spectra, density functional theory (DFT) calculations were performed generating structures of  $V^+(H_2O)_n$  at the B3LYP/aug-cc-pVDZ level of theory. Simulated infrared bands are scaled by a factor of 0.96 to compensate for anharmonicity and method deficiency, and spectra are generated by implementing Gaussian functions to band positions, each with a full-width-half-maximum (FWHM) of 20  $cm^{-1}$ . All calculations were performed using the Gaussian 16 package.<sup>83</sup> Relative energies of isomers are given in  $kJ\ mol^{-1}$ , inclusive of zero-point energy. All isomers of the intact  $V^+(H_2O)_n$  structure considered in this study are present in the quintet ground state, with other multiplicities lying higher in energy. The inserted  $HVOH^+(H_2O)_n$  structures were found in triplet multiplicity, and are more stable than their singlet and quintet analogues. Wave function stabilization was performed prior to each structure optimization.

## Results and discussion

In Fig. 1, we compare the measured IRMPD spectra of  $V^+(H_2O)_n$ ,  $n = 3-51$ , Fig. 2 includes fragmentation intensities of the observed dissociation channels. Infrared spectra of Ar and  $N_2$  tagged as well as untagged  $V^+(H_2O)_n$  have been published previously in two separate studies by the groups of Duncan and Ohashi, and the original spectra from the Duncan group are included in Fig. 1.<sup>58,62</sup> Our photodissociation spectra up to  $n = 7$  are concordant with these previous studies. Since no tagging technique is used in this work, investigations of smaller systems ( $n \leq 2$ ) are not possible due to inefficient elimination





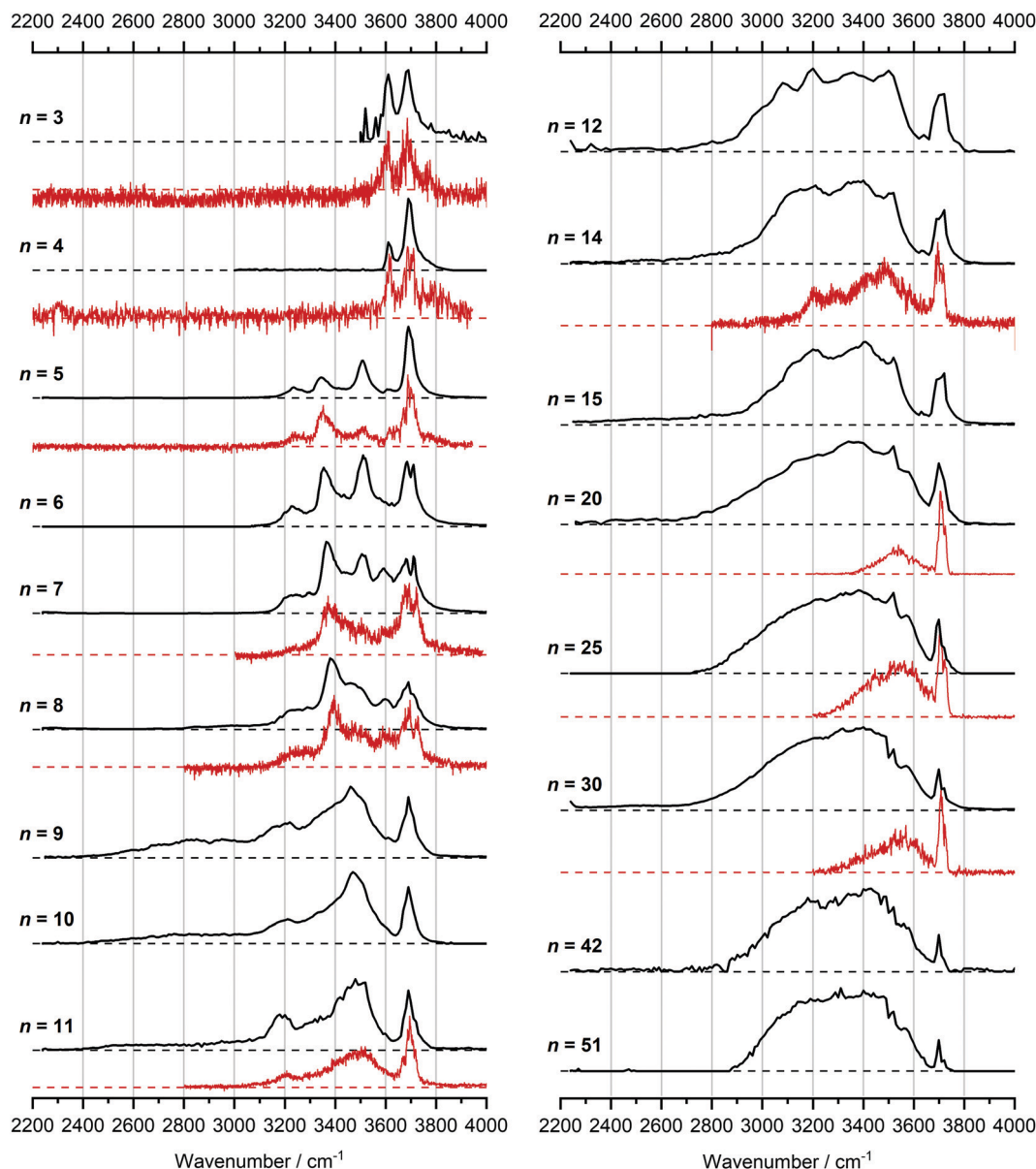


Fig. 1 Infrared multiple photon dissociation spectra of the  $V^+(\text{H}_2\text{O})_n$  complexes ( $n = 3$ –12, 14, 15, 20, 25, 30, 42 and 51) presented in black. For comparison, infrared spectra from Duncan and coworkers are also presented in red for cluster sizes  $n = 3$ –5, 7, 8, 11, 14, 20, 25 and 30.<sup>62</sup>

of  $\text{H}_2\text{O}$ . For cluster sizes  $n \geq 8$ , differences start to emerge between the IRMPD spectra measured in the present work (black spectra) and those recorded by Duncan and coworkers (red spectra).<sup>62</sup> Most notably, the broad and weak red-shifted feature at  $2500$ – $3000 \text{ cm}^{-1}$  seen in our IRMPD spectrum for  $n = 11$  is not present in the  $n = 11$  spectrum recorded by Duncan. Going larger to  $n = 30$ , the strongest feature is a broad absorption at  $2800$ – $3650 \text{ cm}^{-1}$ , with weaker, “free” O–H bands at  $3700$  and  $3720 \text{ cm}^{-1}$ . Duncan’s spectrum shows the opposite, the “free” O–H bands are the most intense, while the absorption at  $3200$ – $3650 \text{ cm}^{-1}$  is weaker and spectrally narrower. These spectral differences can be attributed to the disparities in the two experimental methods employed. In particular, a time-of-flight mass spectrometry instrument is used in the

study by Duncan, employing single-photon infrared photon dissociation (IRPD) methods. Upon increasing the number of water molecules, the heat capacity of the cluster increases, which means water evaporation from the cluster becomes more demanding energetically. Thus, under a single-photon regime, the photofragmentation rate decreases upon increasing cluster size. In contrast, clusters in the present study are trapped in the center of an ICR cell, whereby infrared multiple photon dissociation (IRMPD) spectroscopy is employed. Therefore, any increases in heat capacity can be readily overcome by introducing more photons, causing significant photodissociation.

We start our discussion with the smallest clusters,  $n = 3$ –7, for which we see structured spectra with clearly defined bands. The only observed dissociation channel is water evaporation.



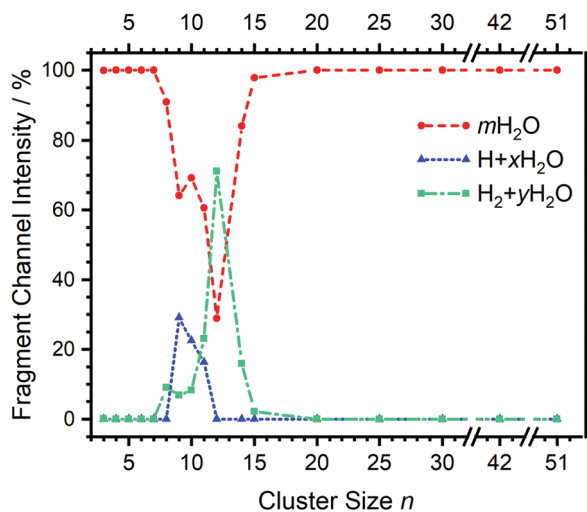


Fig. 2 Fragmentation channel intensities for all investigated clusters. The points are connected with lines to guide the eye.

In Fig. 3, we analyze the experimental as well as theoretical spectra of  $V^+(H_2O)_5$  as an example (spectra of other small clusters are analyzed in Fig. S1–S3, ESI $\ddagger$ ). The experimental spectrum is in excellent agreement with the untagged spectra measured by Duncan, as well as the spectra recorded by Ohashi.<sup>58,62</sup> As shown in Fig. 3, all three low-lying isomers of the intact  $V^+(H_2O)_5$  present a good agreement with the bands in the observed spectrum, possessing both the four- (**Va**, **Vb**) and three-coordinate (**Vc**) vanadium cation. The bands observed at 3510 and 3690  $cm^{-1}$  can be explained by isomer **Va**, consistent with Duncan's results of the  $V^+(H_2O)_5$ -Ar cluster.<sup>62</sup> The feature at 3345  $cm^{-1}$  can be ascribed to the single-acceptor water molecule of isomer **Vb**, and both features at 3345 and 3230  $cm^{-1}$  are in reasonable agreement with the two single-acceptor water molecules in **Vc**. In principle, the feature at 3230  $cm^{-1}$  can also be assigned to the bending overtone of isomer **Va**, as proposed by Ohashi and coworkers.<sup>58</sup> However, our anharmonic frequency analysis shows that, while there are indeed bending overtones in the 3100–3200  $cm^{-1}$  region, their intensity is two orders of magnitude lower compared to the intensity of the most intense O–H stretch. The four bands present in the observed IRMPD spectrum can be thus ascribed to contributions from only intact isomers.

The inserted species,  $HVOH^+(H_2O)_4$  (**iVa**–**iVc**), are considerably more stable than the intact clusters (Fig. 3d). However, they show intense red-shifted bands below 3000  $cm^{-1}$  which are not observed experimentally. Thus, the barrier to insertion must be too high to attain at these temperatures for this cluster size. This is also consistent with previous BIRD investigations on  $V^+(H_2O)_5$ , which presented only dissociation of water molecules.<sup>19</sup>

Starting with  $n = 8$ ,  $H_2$  evaporation appears as a new reaction channel, although sequential loss of water molecules is still the main channel (Fig. 2), consistent with previous BIRD studies of  $V^+(H_2O)_8$ .<sup>19</sup> The number, and composition, of ligands lost during the IRMPD process for each cluster size is given in

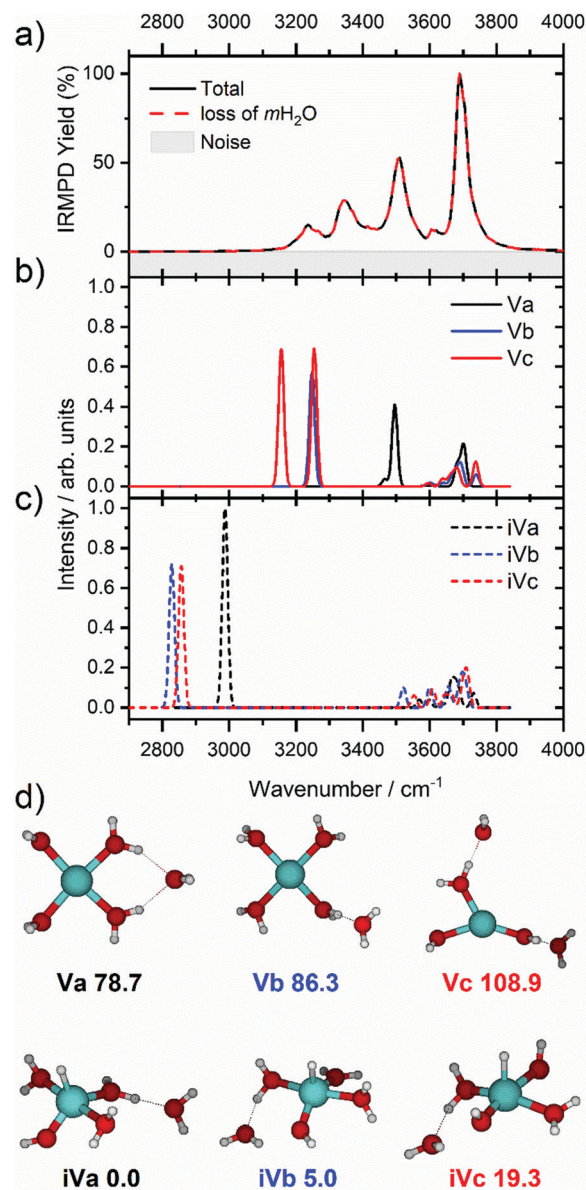


Fig. 3 (a) Infrared multiple photon dissociation spectrum of the  $V^+(H_2O)_5$  complex along with the simulated spectra of low-lying (b) intact (**Va**–**Vc**) and (c) inserted (**iVa**–**iVc**) isomers calculated at the B3LYP/aug-cc-pVDZ level of theory. (d) Structural isomers of intact (**Va**–**Vc**) and inserted isomers (**iVa**–**iVc**) with relative energies given in  $kJ\ mol^{-1}$  inclusive of zero-point energy. In each case the observed photodissociation events in a) are due to the sequential loss of intact water molecules,  $mH_2O$ .

the ESI $\ddagger$  (see Table S1). In each case,  $m$  denotes the sequential number of water molecules,  $x$  denotes the number of waters following H dissociation, and  $y$  the number of waters following  $H_2$  dissociation. Fig. 4a presents the IRMPD spectrum measured for  $V^+(H_2O)_8$ , showing a qualitatively similar IRMPD spectrum as observed for  $n = 5$ – $7$ . As discussed by Duncan,<sup>62</sup> upon increasing cluster size the structural evolution of  $V^+(H_2O)_n$  clusters retains a quasi-planar, square structural motif up to  $n = 8$ , presenting a metal coordination number of four. The two-dimensional structure is broken with inserted isomers. Different to the experimental spectrum recorded by



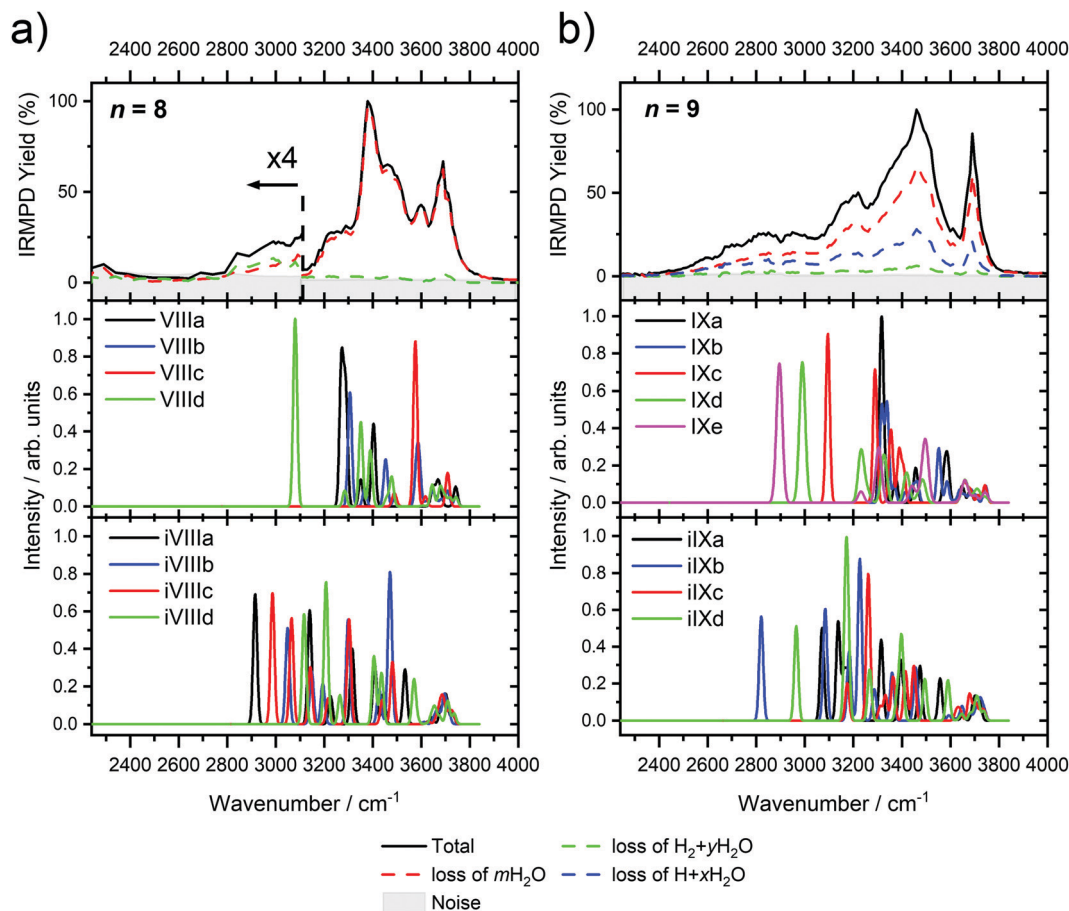


Fig. 4 Infrared multiple photon dissociation spectra of (a)  $V^+(H_2O)_8$  and (b)  $V^+(H_2O)_9$  along with the simulated photodissociation spectra of low-lying intact and inserted isomers calculated at the B3LYP/aug-cc-pVDZ level of theory. In each case, photodissociation events are due to sequential loss of either (i) loss of intact water molecules  $mH_2O$ , (ii) loss of  $H + xH_2O$ , and (iii) loss of  $H_2 + yH_2O$ , represented as dashed red, blue, and green lines, respectively.

Duncan, the spectrum in Fig. 4a shows the presence of a weak feature between 2790–3110  $cm^{-1}$ .

IR spectra provided by quantum chemical calculations offer two possible explanations for the low-intensity tail of the spectrum. The band could correspond to inserted isomers (**iVIIIa–iVIIId**, see Fig. 5a) as these spectra present red-shifted bands in this region. However, the intact isomer **VIIIId** also presents a band in this region at 3078  $cm^{-1}$ . This vibration is represented by a green arrow in Fig. 5a, showing the O–H stretch in a second solvation sphere water molecule, in a double-acceptor site bound to another third solvation sphere water molecule. The observed red-shift signifies significant weakening of the O–H bond, and could indicate the onset of H atom transfer from the second sphere water molecule, forming  $H_3O^+$  in the third sphere. Isomer **VIIIId** is also only 7  $kJ mol^{-1}$  less stable than the most stable intact isomer found, **VIIIa**.

The discrepancy between the spectrum recorded in the present study and that by Duncan could be rationalized by temperature effects; for clusters thermalized in an ICR cell, warmer clusters are present, where entropic effects dominate enthalpic effects. More entropic water binding motifs with

red-shifted bands, such as single-acceptor sites observed in **VIIIId**, are present. This effect has been observed before in our own infrared studies on hydrated zinc cations<sup>55,84</sup> and in IR studies on  $M^+(H_2O)_n$  complexes elsewhere.<sup>85,86</sup>

The main contribution to the red-shifted feature between 2790–3110  $cm^{-1}$  is the  $H_2 + H_2O$  fragmentation channel (dashed green spectrum), and is slightly more intense than the intact water loss channel (dashed red spectrum). Based on previously calculated  $H_2$  loss mechanisms of  $Al^+(H_2O)_{20}$ , pathways involving inserted isomers are more energetically favorable.<sup>87,88</sup> Briefly, a proton moves from a first to a second shell water molecule leading to  $H_3O^+$ . Through consecutive proton transfer through a “water wire”, the proton reaches a site near the hydride statistically. When the proton is pointing towards the hydride at the metal center, the hydride and the proton can recombine leading to the release of an  $H_2$  molecule. The energy released from this recombination step leads to evaporation of multiple  $H_2O$  molecules, see calculated reaction energies in Table 1. Thus, based on the weakly observed  $H_2 + H_2O$  fragment channel, there is a small contribution from inserted isomers. However, the majority of the ions present in the experiment are intact structures,  $V^+(H_2O)_8$ , as can be





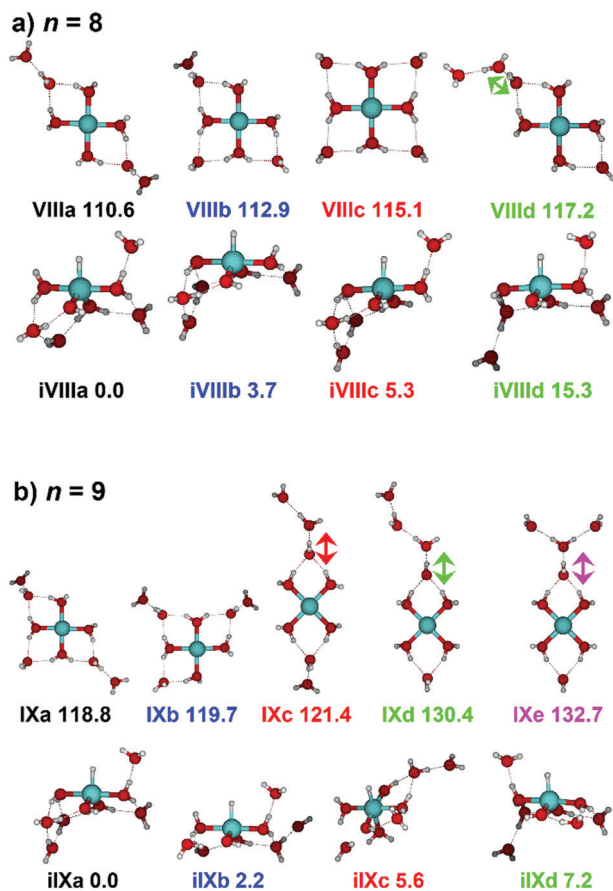


Fig. 5 Calculated intact and inserted isomers of (a)  $V^+(H_2O)_8$  and (b)  $V^+(H_2O)_9$ . All isomers were calculated at the B3LYP/aug-cc-pVDZ level of theory with relative energies given in  $\text{kJ mol}^{-1}$  inclusive of zero-point energy. Arrows indicate the O–H stretch of the most red-shifted band, in each case a water molecule in a double acceptor binding motif in the second solvation sphere.

Table 1 Calculated insertion energies ( $E_{\text{ins}}$ ) and dissociation energies (all in  $\text{kJ mol}^{-1}$ ) for intact and inserted isomers of  $V^+(H_2O)_n$  stoichiometry. Calculated at the B3LYP/aug-cc-pVDZ level of theory

$n$	$E_{\text{ins}}$	Intact ( $V^+(H_2O)_n$ )			Inserted ( $HVOH^+(H_2O)_{n-1}$ )		
		$H_2O$	H	$H_2$	$H_2O$	H	$H_2$
5	–79	50	116	–142	63	195	–63
6	–90	42	88	–156	54	178	–66
7	–103	39	63	–174	52	165	–71
8	–111	40	45	–185	48	156	–74
9	–119	35	34	–196	43	153	–77

deduced from the low-intensity red tail of the spectrum. The energetically less stable intact isomers are generated in the source chamber, while the inserted isomers,  $HVOH^+(H_2O)_7$ , are generated *via* insertion and consecutive water fragmentation of larger clusters. A possible  $2H + H_2O$  dissociation channel can be discarded as it lies by 4.48 eV ( $H_2$  dissociation energy) higher in energy and is unlikely to contribute.

The IRMPD spectrum of  $n = 9$  is presented in Fig. 4b, showing that the dissociation of water molecules is the dominant fragmentation channel, followed by  $H + H_2O$ , then the

$H_2 + H_2O$  channel. Similar fragmentation channels were observed in BIRD experiments, suggesting the same reactions take place with different heating mechanisms.<sup>19</sup> Based on the discussion above, we can attribute  $H_2$  loss to the presence of inserted isomers,  $HVOH^+(H_2O)_8$ . The H loss, on the other hand, can be assigned only to intact isomers as it is energetically too demanding for the inserted species ( $153 \text{ kJ mol}^{-1}$ , Table 1). In addition, the energy required to dissociate an H atom from an intact cluster ( $34 \text{ kJ mol}^{-1}$ ) becomes isoenergetic with  $H_2O$  evaporation for  $n = 9$  (Table 1). The  $H + H_2O$  dissociation channel, along with the water dissociation channel, is thus assigned to intact isomers,  $V^+(H_2O)_9$ . In other words, the thermochemical considerations point at the co-existence of both cluster types. The intact and inserted isomers of  $V^+(H_2O)_9$ , along with the atomic and molecular hydrogen dissociation energy pathways, are presented in Fig. 6. The lowest energy pathway starts with an intact isomer  $^5[V(H_2O)_9]^+$ , with H lost from a water molecule (*i.e.* without insertion, resulting in product  $^4[VOH(H_2O)_8]^+$ ). The insertion step is exothermic by  $119 \text{ kJ mol}^{-1}$ , which would certainly result in evaporation of multiple water molecules. Nevertheless, atomic H loss from the inserted species  $^3[HVOH(H_2O)_8]^+$  requires  $257 \text{ kJ mol}^{-1}$ , resulting in the inserted product  $^2[HV(OH)_2(H_2O)_7]^+$ . Atomic H loss from the inserted species  $^3[HVOH(H_2O)_8]^+$ , resulting in  $^4[VOH(H_2O)_8]^+$ , is also energetically demanding, calculated as  $153 \text{ kJ mol}^{-1}$ . Based on these energies, the photodissociation spectrum of  $H + H_2O$  in Fig. 4b originates from intact isomers.

Similar to  $n = 8$ , the spectrum of  $n = 9$  also shows a broad structureless absorption at  $2410\text{--}3640 \text{ cm}^{-1}$ , the “free” O–H band at  $3690 \text{ cm}^{-1}$  and a broad red-shifted feature at  $2410\text{--}3000 \text{ cm}^{-1}$  which gains intensity and an increased spectral width for  $n = 9$ . Fig. 4b also shows the simulated spectra of calculated intact isomers, **IXa–IXe**, and inserted isomers, **iIXa–iIXd**, the structures of which are shown in Fig. 5b. The broad absorption at  $2410\text{--}3640 \text{ cm}^{-1}$  is the hydrogen bonding region composed of many different O–H environments with varying degrees of red-shift, with both intact and inserted isomers presenting agreement. The “free” O–H band at  $3690 \text{ cm}^{-1}$  also shows agreement with both isomer classes.

The red-shifted feature at  $2410\text{--}3000 \text{ cm}^{-1}$  with all three reaction channels present provides particular insight into the cluster structure. As discussed above, we attribute the feature to the presence of both inserted isomers **iIXb** and **iIXd** (showing bands at  $2820$  and  $2965 \text{ cm}^{-1}$ , respectively) and the intact isomers, namely **IXc**, **IXd**, and **IXe** with bands at  $3095$ ,  $2990$ , and  $2895 \text{ cm}^{-1}$ , respectively. The O–H stretch of the latter bands are shown as arrows in Fig. 5b, in each case the O–H bond of a second-sphere water molecule in a double-acceptor motif, bound to a third-sphere water molecule. Isomer **IXe** shows the most red-shifted band, which is caused by two fourth-sphere water molecules bound to the third-sphere water molecule acting on the second-sphere water. Induced polarisation effects onto the second-sphere water molecule from the metal center, and from the third-sphere water molecule, remove electron density from the bonding orbitals in the O–H bond, which weaken the bond causing the red-shift.



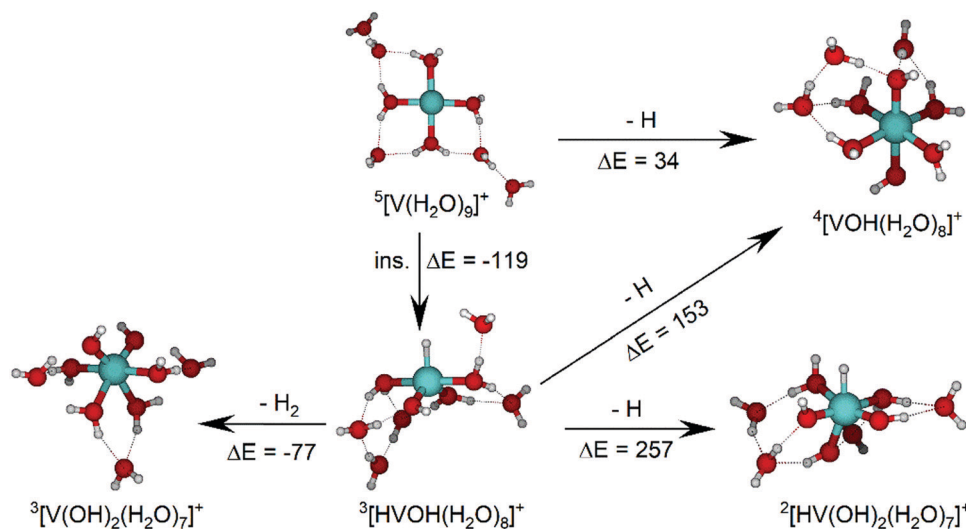


Fig. 6 Calculated intact and inserted isomers of  $V^+(H_2O)_9$ , along with insertion energy (ins.), molecular hydrogen and atomic hydrogen dissociation energy pathways. All isomers and energy pathways were calculated at the B3LYP/aug-cc-pVDZ level of theory with relative energies given in  $\text{kJ mol}^{-1}$  inclusive of zero-point energy.

We conclude that the majority of the ions for  $n = 9$  are composed of intact isomers, however  $n = 9$  represents the cluster size at which the abundance of inserted isomers starts to increase. Our conclusions are consistent with our previous investigation of  $V^+(H_2O)_n$  in the ultraviolet/visible range, with a decrease in photodissociation cross section starting at  $n \approx 9$ , decreasing further upon increasing water until  $n \geq 15$  where no photoinduced fragmentation was observed, indicating an oxidation state change from  $V(\text{I})$  to  $V(\text{III})$ .<sup>63</sup>

For  $n = 11$ , bands assigned to single-acceptor water motifs around  $3200 \text{ cm}^{-1}$  increase in intensity, as presented in Fig. 7. In terms of fragmentation channels, the  $\text{H}_2 + \text{H}_2\text{O}$  channel (green dashed spectrum) is as intense as the  $\text{H} + \text{H}_2\text{O}$  channel (blue dashed spectrum). However, the onset of  $\text{H}_2$  loss starts at lower photon energies ( $2890 \text{ cm}^{-1}$ ) than  $\text{H}$  loss ( $3120 \text{ cm}^{-1}$ ). Below  $2890 \text{ cm}^{-1}$ , only the  $\text{H}_2\text{O}$  fragment channel is detected. The dissociation patterns suggest that intact and inserted isomers still co-exist for this cluster size.

Going to  $V^+(H_2O)_{12}$ , differences in fragmentation channels are observed (Fig. 7). In the whole spectral range, the  $\text{H}_2 + \text{H}_2\text{O}$  channel is the most intense with no  $\text{H} + \text{H}_2\text{O}$  loss observed and water loss observed as a weaker channel. Additionally, the spectral band intensities change, indicating a completed structural change to the inserted clusters. The band at  $2900 \text{ cm}^{-1}$  for  $n = 9, 10$  and  $11$ , seems to decrease in intensity but it is still observable above the noise level. Since  $V^+(H_2O)_{12}$  shows the second highest BIRD rate for  $\text{H}_2 + \text{H}_2\text{O}$  loss, the signal-to-noise ratio, especially in the lower wavelength region, is lower due to BIRD dissociation, even at low temperatures.<sup>19</sup> It is possible that this band becomes lost in the noise and thus, is undetectable. Given that  $\text{H}_2 + \text{H}_2\text{O}$  is the most intense channel, the majority of ions present for  $n = 12$  are composed of inserted isomers. This would indicate that 12 water molecules are enough to form three dimensional structures, whereby the

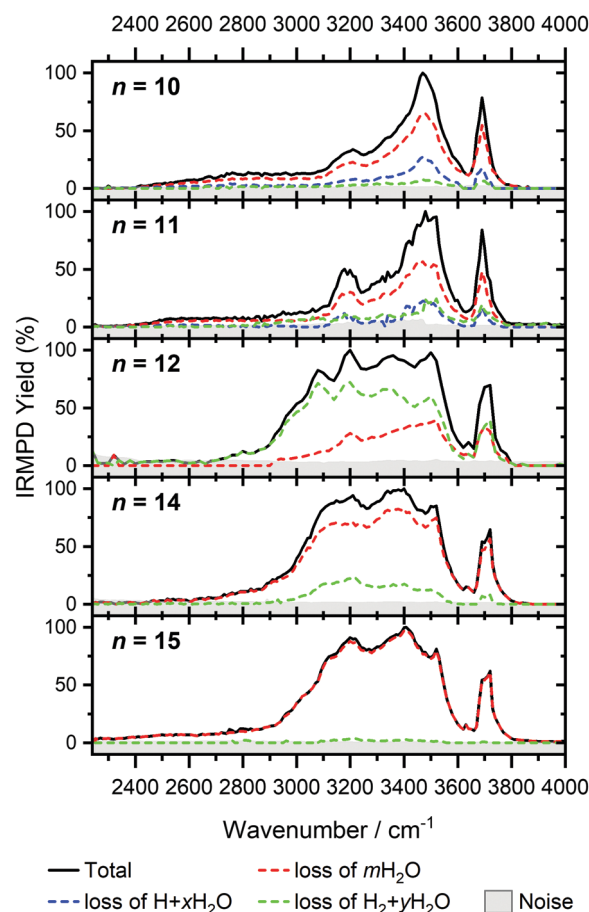


Fig. 7 Infrared multiple photon dissociation spectra of  $V^+(H_2O)_n$  complexes ( $n = 10-12, 14, 15$ ). In each case photodissociation events are due to sequential loss of either (i) loss of intact water molecules  $m\text{H}_2\text{O}$ , (ii) loss of  $\text{H} + x\text{H}_2\text{O}$ , and (iii) loss of  $\text{H}_2 + y\text{H}_2\text{O}$ , represented as dashed red, blue, and green lines, respectively.





water network can fold back and attack the metal center, cleaving the O–H bond, forming a HVOH<sup>+</sup> moiety.

The IRMPD spectra of V<sup>+</sup>(H<sub>2</sub>O)<sub>14</sub> and V<sup>+</sup>(H<sub>2</sub>O)<sub>15</sub> are qualitatively similar to the observed spectrum of  $n = 12$  (Fig. 7). One broad absorption feature from 2900–3600 cm<sup>-1</sup> along with a band within the isolated O–H stretching region around 3700 cm<sup>-1</sup> is observed. The H<sub>2</sub> + H<sub>2</sub>O loss channel decreases for  $n = 14$  when compared to  $n = 12$ , with the fragmentation intensity decreasing sharply at  $n = 15$ , consistent with BIRD studies for this cluster size.<sup>19</sup> It is likely that the majority of ions at cluster sizes  $n = 14, 15$  are inserted, however the loss of water molecules is energetically and/or entropically preferred. Thus, although the H<sub>2</sub> + H<sub>2</sub>O channel decreases upon increasing cluster size, these isomers are most likely still inserted, consistent with ultraviolet/visible spectroscopy of these cluster sizes.<sup>63</sup>

The largest clusters (up to  $n = 51$ ) investigated are shown in Fig. 1. In this cluster size range, only water loss is observed after irradiation. For  $n = 20$ –51, the IRMPD spectra are qualitatively similar to those observed for  $n = 14, 15$ ; one broad absorption feature between 2900–3600 cm<sup>-1</sup> with an additional band within the free O–H stretching region around 3700 cm<sup>-1</sup>. These spectra, however, are less structured. The onset of the broad feature varies with the cluster size. V<sup>+</sup>(H<sub>2</sub>O)<sub>20,25,30</sub> show an onset at about 2700 cm<sup>-1</sup> while for clusters greater than 30 H<sub>2</sub>O molecules the onset is blue shifted to roughly 2900 cm<sup>-1</sup>. This is in stark contrast to previous experiments,<sup>62</sup> where fragmentation in this size range was observed only above 3200 cm<sup>-1</sup>. As the  $n = 12$  cluster is mostly composed of inserted isomers, this indicates that larger clusters are also composed of inserted species, in agreement with our UV/Vis spectra.<sup>63</sup> This is not surprising, as this species provides the nucleation site to which additional water molecules bind, forming larger clusters. A counter example to this, revealed by UV/Vis spectroscopy, was found for hydrated magnesium cations, which also show the presence of inserted isomers for  $n = 12$ , (Mg<sup>+</sup>OH(H<sub>2</sub>O) <sub>$n-1$</sub> ), however with increasing cluster size ( $n > 25$ ) form intact structures, Mg<sup>+</sup>(H<sub>2</sub>O) <sub>$n$</sub> .<sup>30</sup> More revealing in the large size range of V<sup>+</sup>(H<sub>2</sub>O) <sub>$n$</sub>  are spectra in the UV/Vis range, which do not show photodissociation, meaning the HVOH<sup>+</sup>(H<sub>2</sub>O) <sub>$n$</sub>  structure is present. Also, theoretical work on hydrated aluminum showed that the insertion reaction is not affected by an increasing solvation shell.<sup>88</sup>

The signature of a mobile proton at *ca.* 2800–3500 cm<sup>-1</sup> is not visible for  $n \geq 42$ .<sup>89</sup> This could be due to the increased number of strong single- and double-acceptor O–H stretching bands in addition to a lower signal-to-noise ratio due to a smaller ion intensity and higher rate of BIRD. This also explains why the free O–H stretching bands around 3700 cm<sup>-1</sup> decrease with increasing cluster size relative to the absorption maximum.

## Conclusions

The IRMPD spectral fingerprints and size-dependent fragmentation channels of V<sup>+</sup>(H<sub>2</sub>O) <sub>$n$</sub>  yield detailed insight into the cluster structure. For  $n = 3$ –7, structured IR spectra are

observed with water loss as the only observed channel, indicating intact clusters present. H<sub>2</sub> loss is observed as a weak channel at  $n = 8$ , indicating the onset of inserted clusters, HVOH<sup>+</sup>(H<sub>2</sub>O) <sub>$n-1$</sub> . The H loss channel is observed for  $n = 9$ –11 and arises at the point where H and H<sub>2</sub>O dissociation channels become energetically competitive for the intact clusters. A noticeable change in the spectroscopic pattern and fragmentation channels between  $n = 11$  and 12 suggests that structural changes are responsible for preventing the release of H, *i.e.*, the clusters are fully converted to the inserted isomeric form. Finally, the H<sub>2</sub> loss channel disappears around  $n = 15$ , likely because of the growing H<sub>2</sub>O network; the number of possible pathways for the mobile proton increases and the probability for it to reach the hydride site required for recombination decreases. However, since only selected cluster sizes were probed for  $n > 12$ , cluster sizes  $n = 17$  and 18 may also show H<sub>2</sub> loss upon IRMPD, as observed in previous BIRD studies. In any event, significant changes in IRMPD spectra are not observed for  $n = 14$ –51. Consistent with the earlier UV/Vis results, these clusters are fully transferred to their inserted isomeric form.

## Conflicts of interest

There are no conflicts of interest to declare.

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## References

- 1 M. A. Duncan, *Int. Rev. Phys. Chem.*, 2003, **22**, 407.
- 2 M. A. Duncan, *Annu. Rev. Phys. Chem.*, 1997, **48**, 69.
- 3 A. J. Stace, *J. Phys. Chem. A*, 2002, **106**, 7993.
- 4 V. E. Bondybey and M. K. Beyer, *Int. Rev. Phys. Chem.*, 2002, **21**, 277.
- 5 M. K. Beyer, *Mass Spectrom. Rev.*, 2007, **26**, 517.
- 6 N. C. Polfer and J. Oomens, *Mass Spectrom. Rev.*, 2009, **28**, 468.
- 7 V. Artero, M. Chavarot-Kerlidou and M. Fontecave, *Angew. Chem., Int. Ed.*, 2011, **50**, 7238.
- 8 C. Tard and C. J. Pickett, *Chem. Rev.*, 2009, **109**, 2245.
- 9 S. Canaguier, V. Artero and M. Fontecave, *Dalton Trans.*, 2008, 315.
- 10 B. E. Barton, C. M. Whaley, T. B. Rauchfuss and D. L. Gray, *J. Am. Chem. Soc.*, 2009, **131**, 6942.



- 11 R. Brimblecombe, G. F. Swiegers, G. C. Dismukes and L. Spiccia, *Angew. Chem., Int. Ed.*, 2008, **47**, 7335.
- 12 K. Fuke, K. Hashimoto and S. Iwata, *Adv. Chem. Phys.*, 1999, **110**, 431.
- 13 G. Niedner-Schatteburg and V. E. Bondybey, *Chem. Rev.*, 2000, **100**, 4059.
- 14 W. A. Donald, R. D. Leib, J. T. O'Brien, A. I. S. Holm and E. R. Williams, *Proc. Natl. Acad. Sci. U. S. A.*, 2008, **105**, 18102.
- 15 N. F. Dalleska, K. Honma, L. S. Sunderlin and P. B. Armentrout, *J. Am. Chem. Soc.*, 1994, **116**, 3519.
- 16 O. P. Balaj, C. B. Berg, S. J. Reitmeier, V. E. Bondybey and M. K. Beyer, *Int. J. Mass Spectrom.*, 2009, **279**, 5.
- 17 M. Beyer, C. Berg, H. W. Görlitzer, T. Schindler, U. Achatz, G. Albert, G. Niedner-Schatteburg and V. E. Bondybey, *J. Am. Chem. Soc.*, 1996, **118**, 7386.
- 18 M. Beyer, U. Achatz, C. Berg, S. Joos, G. Niedner-Schatteburg and V. E. Bondybey, *J. Phys. Chem. A*, 1999, **103**, 671.
- 19 B. S. Fox, I. Balteanu, O. P. Balaj, H. C. Liu, M. K. Beyer and V. E. Bondybey, *Phys. Chem. Chem. Phys.*, 2002, **4**, 2224.
- 20 C. van der Linde and M. K. Beyer, *J. Phys. Chem. A*, 2012, **116**, 10676.
- 21 C. van der Linde and M. K. Beyer, *Phys. Chem. Chem. Phys.*, 2011, **13**, 6776.
- 22 B. Scharfschwerdt, C. van der Linde, O. P. Balaj, I. Herber, D. Schütze and M. K. Beyer, *Low Temp. Phys.*, 2012, **38**, 717.
- 23 D. E. Lessen, R. L. Asher and P. J. Brucat, *J. Chem. Phys.*, 1990, **93**, 6102.
- 24 M. Rosi and C. W. Bauschlicher, *J. Chem. Phys.*, 1989, **90**, 7264.
- 25 M. Rosi and C. W. Bauschlicher, *J. Chem. Phys.*, 1990, **92**, 1876.
- 26 J. M. Farrar, *Int. Rev. Phys. Chem.*, 2003, **22**, 593.
- 27 C. S. Yeh, K. F. Willey, D. L. Robbins, J. S. Pilgrim and M. A. Duncan, *Chem. Phys. Lett.*, 1992, **196**, 233.
- 28 C. T. Scurlock, S. H. Pullins, J. E. Reddic and M. A. Duncan, *J. Chem. Phys.*, 1996, **104**, 4591.
- 29 M. Ončák, T. Taxer, E. Barwa, C. van der Linde and M. K. Beyer, *J. Chem. Phys.*, 2018, **149**, 44309.
- 30 T. Taxer, M. Ončák, E. Barwa, C. van der Linde and M. K. Beyer, *Faraday Discuss.*, 2019, **217**, 584.
- 31 J. S. Daluz, A. Kocak and R. B. Metz, *J. Phys. Chem. A*, 2012, **116**, 1344.
- 32 A. Kocak, G. Austein-Miller, W. L. Pearson, G. Altinay and R. B. Metz, *J. Phys. Chem. A*, 2013, **117**, 1254.
- 33 Y. Abate and P. D. Kleiber, *J. Chem. Phys.*, 2005, **122**, 84305.
- 34 M. Sanekata, F. Misaizu and K. Fuke, *J. Chem. Phys.*, 1996, **104**, 9768.
- 35 T. Iino, K. Ohashi, K. Inoue, K. Judai, N. Nishi and H. Sekiya, *Eur. Phys. J. D*, 2007, **43**, 37.
- 36 B. Bandyopadhyay, K. N. Reishus and M. A. Duncan, *J. Phys. Chem. A*, 2013, **117**, 7794.
- 37 J. M. Lisy, *Int. Rev. Phys. Chem.*, 1997, **16**, 267.
- 38 J. P. Beck and J. M. Lisy, *J. Chem. Phys.*, 2011, **135**, 44302.
- 39 N. R. Walker, R. S. Walters, M.-K. Tsai, K. D. Jordan and M. A. Duncan, *J. Phys. Chem. A*, 2005, **109**, 7057.
- 40 R. S. Walters, N. R. Brinkmann, H. F. Schaefer and M. A. Duncan, *J. Phys. Chem. A*, 2003, **107**, 7396.
- 41 R. S. Walters, E. D. Pillai and M. A. Duncan, *J. Am. Chem. Soc.*, 2005, **127**, 16599.
- 42 T. D. Vaden, J. M. Lisy, P. D. Carnegie, E. D. Pillai and M. A. Duncan, *Phys. Chem. Chem. Phys.*, 2006, **8**, 3078.
- 43 P. D. Carnegie, A. B. McCoy and M. A. Duncan, *J. Phys. Chem. A*, 2009, **113**, 4849.
- 44 P. D. Carnegie, B. Bandyopadhyay and M. A. Duncan, *J. Chem. Phys.*, 2011, **134**, 14302.
- 45 P. D. Carnegie, B. Bandyopadhyay and M. A. Duncan, *J. Phys. Chem. A*, 2008, **112**, 6237.
- 46 P. D. Carnegie, B. Bandyopadhyay and M. A. Duncan, *J. Phys. Chem. A*, 2011, **115**, 7602.
- 47 B. Bandyopadhyay and M. A. Duncan, *Chem. Phys. Lett.*, 2012, **530**, 10.
- 48 T. Iino, K. Ohashi, Y. Mune, Y. Inokuchi, K. Judai, N. Nishi and H. Sekiya, *Chem. Phys. Lett.*, 2006, **427**, 24.
- 49 Y. Inokuchi, K. Ohshimo, F. Misaizu and N. Nishi, *J. Phys. Chem. A*, 2004, **108**, 5034.
- 50 Y. Inokuchi, K. Ohshimo, F. Misaizu and N. Nishi, *Chem. Phys. Lett.*, 2004, **390**, 140.
- 51 T. Iino, K. Ohashi, K. Inoue, K. Judai, N. Nishi and H. Sekiya, *J. Chem. Phys.*, 2007, **126**, 194302.
- 52 K. Furukawa, K. Ohashi, N. Koga, T. Imamura, K. Judai, N. Nishi and H. Sekiya, *Chem. Phys. Lett.*, 2011, **508**, 202.
- 53 M. F. Bush, R. J. Saykally and E. R. Williams, *J. Am. Chem. Soc.*, 2008, **130**, 9122.
- 54 J. T. O'Brien and E. R. Williams, *J. Phys. Chem. A*, 2008, **112**, 5893.
- 55 E. M. Cunningham, T. Taxer, J. Heller, M. Ončák, C. van der Linde and M. K. Beyer, *Phys. Chem. Chem. Phys.*, 2021, **23**, 3627.
- 56 K. R. Asmis, *Phys. Chem. Chem. Phys.*, 2012, **14**, 9270.
- 57 H. Schwarz and K. R. Asmis, *Chem. – Eur. J.*, 2019, **25**, 2112.
- 58 J. Sasaki, K. Ohashi, K. Inoue, T. Imamura, K. Judai, N. Nishi and H. Sekiya, *Chem. Phys. Lett.*, 2009, **474**, 36.
- 59 M. Zhou, J. Dong, L. Zhang and Q. Qin, *J. Am. Chem. Soc.*, 2001, **123**, 135.
- 60 N. R. Walker, R. S. Walters, E. D. Pillai and M. A. Duncan, *J. Chem. Phys.*, 2003, **119**, 10471.
- 61 T. B. Ward, E. Miliordos, P. D. Carnegie, S. S. Xantheas and M. A. Duncan, *J. Chem. Phys.*, 2017, **146**, 224305.
- 62 P. D. Carnegie, J. H. Marks, A. D. Brathwaite, T. B. Ward and M. A. Duncan, *J. Phys. Chem. A*, 2020, **124**, 1093.
- 63 J. Heller, T. F. Pascher, D. Muß, C. van der Linde, M. K. Beyer and M. Ončák, *Phys. Chem. Chem. Phys.*, 2021, **23**, 22251–22262.
- 64 J. Heller, W. K. Tang, E. M. Cunningham, E. G. Demissie, C. van der Linde, W. K. Lam, M. Ončák, C.-K. Siu and M. K. Beyer, *Angew. Chem., Int. Ed.*, 2021, **60**, 16858–16863.
- 65 A. Akhgarnusch, W. K. Tang, H. Zhang, C.-K. Siu and M. K. Beyer, *Phys. Chem. Chem. Phys.*, 2016, **18**, 23528.
- 66 M. Allemann, H. Kellerhals and K. P. Wanczek, *Int. J. Mass Spectrom. Ion Process.*, 1983, **46**, 139.
- 67 C. Berg, T. Schindler, G. Niedner-Schatteburg and V. E. Bondybey, *J. Chem. Phys.*, 1995, **102**, 4870.



- 68 A. Akhgarnusch, R. F. Höckendorf and M. K. Beyer, *J. Phys. Chem. A*, 2015, **119**, 9978.
- 69 P. Caravatti and M. Allemann, *Org. Mass Spectrom.*, 1991, **26**, 514.
- 70 V. E. Bondybey and J. H. English, *J. Chem. Phys.*, 1981, **74**, 6978.
- 71 T. G. Dietz, M. A. Duncan, D. E. Powers and R. E. Smalley, *J. Chem. Phys.*, 1981, **74**, 6511.
- 72 D. Proch and T. Trickl, *Rev. Sci. Instrum.*, 1989, **60**, 713.
- 73 A. G. Marshall, C. L. Hendrickson and G. S. Jackson, *Mass Spectrom. Rev.*, 1998, **17**, 1.
- 74 N. Polfer, B. G. Sartakov and J. Oomens, *Chem. Phys. Lett.*, 2004, **400**, 201.
- 75 J. Heller, E. M. Cunningham, C. van der Linde, M. Ončák and M. K. Beyer, *J. Phys. Chem. Lett.*, 2022, **13**, 3269.
- 76 R. L. Wong, K. Paech and E. R. Williams, *Int. J. Mass Spectrom.*, 2004, **232**, 59.
- 77 D. Thölmann, D. S. Tonner and T. B. McMahon, *J. Phys. Chem.*, 1994, **98**, 2002.
- 78 R. C. Dunbar, *Mass Spectrom. Rev.*, 2004, **23**, 127.
- 79 T. Schindler, C. Berg, G. Niedner-Schatteburg and V. E. Bondybey, *Chem. Phys. Lett.*, 1996, **250**, 301.
- 80 P. D. Schnier, W. D. Price, R. A. Jockusch and E. R. Williams, *J. Am. Chem. Soc.*, 1996, **118**, 7178.
- 81 M. Sena and J. M. Riveros, *Rapid Commun. Mass Spectrom.*, 1994, **8**, 1031.
- 82 B. S. Fox, M. K. Beyer and V. E. Bondybey, *J. Phys. Chem. A*, 2001, **105**, 6386.
- 83 M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W. Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa, M. Ishida, T. Nakajima, Y. Honda, O. Kitao, H. Nakai, T. Vreven, K. Throssell, J. A. Montgomery, Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman and D. J. Fox, *Gaussian 16 Revision A.03*, 2016.
- 84 E. M. Cunningham, T. Taxer, J. Heller, M. Ončák, C. van der Linde and M. K. Beyer, *Int. J. Mol. Sci.*, 2021, **22**, 6026.
- 85 J. Kim, S. Lee, S. J. Cho, B. J. Mhin and K. S. Kim, *J. Chem. Phys.*, 1995, **102**, 839.
- 86 K. Ohashi, J. Sasaki, G. Yamamoto, K. Judai, N. Nishi and H. Sekiya, *J. Chem. Phys.*, 2014, **141**, 214307.
- 87 B. M. Reinhard and G. Niedner-Schatteburg, *J. Phys. Chem. A*, 2002, **106**, 7988.
- 88 C.-K. Siu, Z.-F. Liu and J. S. Tse, *J. Am. Chem. Soc.*, 2002, **124**, 10846.
- 89 C. J. Johnson, L. C. Dzugan, A. B. Wolk, C. M. Leavitt, J. A. Fournier, A. B. McCoy and M. A. Johnson, *J. Phys. Chem. A*, 2014, **118**, 7590.

