


 Cite this: *RSC Adv.*, 2023, **13**, 33797

Dopant engineering for ZnO electron transport layer towards efficient perovskite solar cells

Nurul Aliyah Zainal Abidin,^{*a} Faiz Arith,^{†a} N. Syamimi Noorasid,^a Hafez Sarkawi,^b A. Nizamuddin Mustafa,^{†bc} N. E. Safie,^b A. S. Mohd Shah,^d M. A. Azam,^{†ef} Puvaneswaran Chelvanathan^g and Nowshad Amin^h

The conventional electron transport layer (ETL) TiO₂ has been widely used in perovskite solar cells (PSCs), which have produced exceptional power conversion efficiencies (PCE), allowing the technology to be highly regarded and propitious. Nevertheless, the recent high demand for energy harvesters in wearable electronics, aerospace, and building integration has led to the need for flexible solar cells. However, the conventional TiO₂ ETL layer is less preferred, where a crystallization process at a temperature as high as 450 °C is required, which degrades the plastic substrate. Zinc oxide nanorods (ZnO NRs) as a simple and low-cost fabrication material may fulfil the need as an ETL, but they still suffer from low PCE due to atomic defect vacancy. To delve into the issue, several dopants have been reviewed as an additive to passivate or substitute the Zn²⁺ vacancies, thus enhancing the charge transport mechanism. This work thereby unravels and provides a clear insight into dopant engineering in ZnO NRs ETL for PSC.

Received 18th July 2023
 Accepted 30th October 2023

DOI: 10.1039/d3ra04823c

rsc.li/rsc-advances

1. Introduction

Solar energy is an essential primary energy source that continues to gain worldwide recognition for providing energy to human life. In comparison to fossil fuels, solar energy is regarded as the key renewable energy source that assists in reducing greenhouse gas emissions and curbing climate change. Consequently, solar energy technology has received considerable attention, and the solar industry has experienced enormous expansion and sustainable development globally. According to the International Energy Agency (IEA), solar energy continues to be the most prominent among all renewable energy sources with the highest growth of net renewable power capacity over the past 5 years, as illustrated in Fig. 1. Solar

energy has dominated more than half of all other technologies and is expected to grow significantly in the future.¹

Solar cells and other photovoltaic (PV) devices generate power directly from sunlight, which applies the photovoltaic effect. Energy in the form of photons from the sun is absorbed and excites the electrons and produces a current that generates electricity. In general, solar cell technology is categorized into

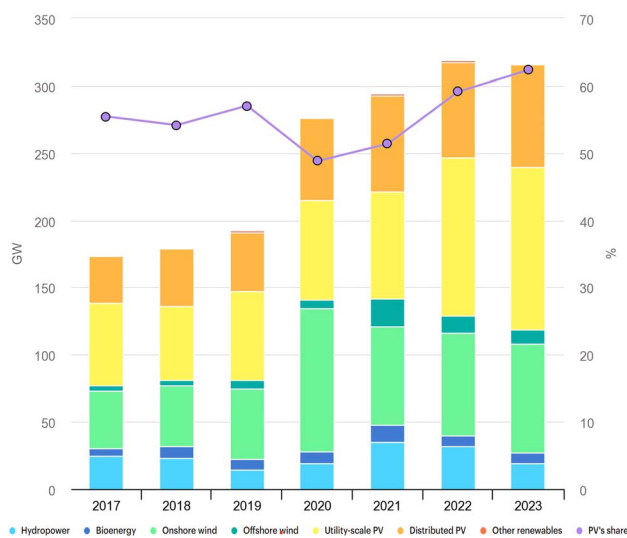


Fig. 1 The growth in the net renewable power capacity by different technologies. Data derived from the International Renewable Energy Agency report.¹

^aFaculty of Electronic and Computer Engineering, Universiti Teknikal Malaysia Melaka, Hang Tuah Jaya, Durian Tunggal, 76100, Melaka, Malaysia. E-mail: faiz.arith@utem.edu.my

^bFaculty of Electrical and Electronic Engineering Technology, Universiti Teknikal Malaysia Melaka, Hang Tuah Jaya, Durian Tunggal, 76100, Melaka, Malaysia

^cDepartment of Materials, Faculty of Engineering, Imperial College London, London SW7 2AZ, UK

^dDepartment of Electrical Engineering, College of Engineering, Universiti Malaysia Pahang, Lebuhraya Tun Razak, Gambang, Kuantan, Pahang, 26300, Malaysia

^eFaculty of Manufacturing Engineering, Universiti Teknikal Malaysia Melaka, 76100 Durian Tunggal, Melaka, Malaysia

^fCenter for Promotion of Educational Innovation, Shibaura Institute of Technology, 3-7-5 Toyosu, Koto-ku, Tokyo, 135-8548, Japan

^gSolar Energy Research Institute, The National University of Malaysia, Bangi, 43600, Selangor, Malaysia

^hDepartment of Electrical and Electronic Engineering, University of Science Engineering and Technology (USTC), Foy's Lake, Chattogram 4202, Bangladesh



three generations, with single and multi-crystalline silicon-based solar cells from the first generation being the most prevalent and widely available on the market. In response to the excessive material consumption and expense of silicon solar cells, second-generation solar cells have been developed, where the maximum layer thickness has been reduced to a nanometer scale, thereby reducing the amount of materials used. Meanwhile, numerous researchers have experimented with the concept of light control utilising dye-sensitive solar cells (DSSCs), perovskite, organic solar cells, quantum dots (QDs), nanostructuring, and nanopatterning, which is considered the third-generation solar cell. This technology, referred to as emerging photovoltaic technology, has been studied extensively for 25 years but has not yet been widely marketed owing to poor cell performance and instability limitations. Nevertheless, emerging photovoltaic technologies offer more competitive fabrication costs, allowing for mass marketing and are also superior in generating energy at low wavelengths, therefore being suitable for indoor use.

Recent developments have demonstrated the prowess of PSCs among other solar cells, which have shown a significant rise in the performance in less than a decade, as displayed in Fig. 2.² This clearly indicates the great potential of PSCs to develop and revolutionize the renewable energy technology. Nevertheless, PSC-related research has largely focused on the perovskite absorber layer and the hole transport layer (HTL), and has designated the mesoporous TiO₂ layer as the ETL. Although TiO₂ ETL has proven its superiority in terms of efficiency, stability and cost-effectiveness, there is still room for improvement in continuing to improve the overall PSC performance. The ETL plays an important role in facilitating the efficient extraction and transport of electrons from the perovskite absorber to the external circuit. In spite of that, the

addition of dopant elements or dopant engineering in ETL offers a viable route in promoting a more efficient charge transport mechanism by means of passivating or replacing the bulk defects in the ETL layer. This allows a way to tailor the layer properties and increase the full potential of that ETL layer, thereby improving the device performance.

2. Emerging perovskite photovoltaic

The PSCs layers of materials are printed, coated, or vacuum-deposited onto a support layer beneath the perovskite cell, known as the substrate in order to construct the cells. For the most part, they are simple to construct and able to achieve efficiency levels comparable to crystalline silicon. For the record, the PSC efficiency has increased from merely 3% in 2009 to over 25% by 2020, demonstrating a dramatic rise as a newly introduced solar cell. PSCs are based on the solid-state dye-sensitized solar cell technology introduced by Gratzel. In 1991, O'Regan and Gratzel invented a low-cost photo-electrochemical solar cell, in which nanocrystalline TiO₂ films were sensitized and reacted with the molecular dyes.³ In 2009, a methylammonium lead halide perovskite was introduced as an absorbent layer in a liquid-based DSSC, which successfully reacted with a nanocrystalline TiO₂ layer, resulting in a PCE of only 3–4%.⁴ Two years later, the structure was named as a PSC, as shown in Fig. 3, and has shown a dramatic improvement in performance, increasing in PCE value by two-folds.⁴ The front contact is formed from an FTO substrate, stacked up with a layered material structure comprising an ETL, a perovskite layer and a HTL. The electron and HTL provide a charge carrier conduction layer, while the perovskite material absorbs photons, depending on the wavelength range correlated with the bandgap of the material.



Fig. 2 The evolution of silicon, perovskite, and perovskite–silicon tandem device efficiency over the past decade, as derived from the efficiency table published by the National Renewable Energy Laboratory. This figure has been reproduced from ref. 2 with permission from the Royal Society of Chemistry, copyright 2021.



Fig. 3 A basic schematic diagram of PSC: (a) hole transport layer, (b) perovskite layer, (c) electron transport layer, (d) FTO/ITO coated glass, and (e) front contact.



Usually, conventional PSCs are in liquid-based form, using Spiro Ometad as the HTL layer. However, these PSCs are still hampered by stability concerns, especially the immediate dissolution of the perovskite material in the liquid-based electrolyte. To address this issue, some researchers proposed the idea of substituting the liquid-based electrolyte with a solid-state perovskite conductor. As a result, a PCE value of 10% was reached in 2012, and continued to rise as high as 18% in less than 2 years, manifesting the great potential of the structure. The stability issue has also been resolved, where cells can function efficiently for long periods of time.⁴ Recently, using low-cost organometal halides combined with silver nanoparticles (AgNPs) as HTL in PSCs, Elseman *et al.* produced good PCE values exceeding 20%. The success achieved in this short period of time shows the high potential of the use of inorganic HTL materials in reducing fabrication costs and alleviating PSC reliability, even improving the PCE value.⁵

In general, solar cells are built as PIN diodes, where one part will be the HTL (p-type) and the other part will be the ETL (n-type). In the basic structure of a PSC as shown in Fig. 3, a perovskite layer (i) is sandwiched between a HTL (p) and an ETL (n), followed by other layers consisting of the FTO/ITO coated glass, back contact and front contact. The HTL and ETL both must have high conductivity, so that the energy level of both layers are properly aligned with the photoactive perovskite layer. The HTL in p-i-n architectures plays an important role in: (i) extracting and transporting holes from the active layer to the electrode; (ii) blocking the transfer of electrons to the anode; (iii) optimizing device stability by reducing the risk of degradation and corrosion; and (iv) increasing the open-circuit voltage (V_{oc}) of the well-matched HOMO energy level.^{6,7} Meanwhile, the ETL in n-i-p architectures also plays a crucial role in the fabrication of high-performance PSC by: (i) transporting photogenerated electrons from the PSC layer to the cathode or anode, and (ii) modifying the contact between the

perovskite layer and the electrode to reduce charge recombination.^{8,9} The development of a high performance HTL for PSC was reported using MEA (monoethanolamine) as the solvent for CuI (copper(i) iodide), demonstrating the effectiveness of combining the solvent and annealing temperature to obtain desirable CuI structures. The utilization of MEA as a solvent enables the simple, inexpensive, and rapid preparation of HTL films, ultimately reducing the cost of solar cell production.¹⁰ Besides that, the simulations were conducted using SCAPS modelling software to analyse the impact of various parameters on the PSC performance. The authors focused on the thickness, doping, operating temperature, and interface defect of different HTLs and absorber layers. The simulations revealed that an inorganic material and Cu₂O in the HTL yielded the highest PCE value of 17.71%, surpassing other HTL materials. Furthermore, using a CH₃NH₃SnI₃ absorber layer, a PCE value as high as 27.72% was achieved, enabling its potential as a replacement for CH₃NH₃PbI₃ due to its higher efficiency and being Pb-free and cost-effective.¹¹

Apart from HTL and ETL layers, the transport mechanism (diffusion) of the charge carrier in PSC materials is one of the important components of a highly efficient PSC. The net movement of the particles (holes and electrons) induced by the density gradient is called diffusion, where the diffusion of these particles involves the movement of either holes or electrons to locations with lower densities, as described in Fig. 4(a). The diffusion will exist as long as there is a concentration gradient in the particles, and it can be described by Fick's law of diffusion:¹²

$$J_e = qD_e \frac{dn}{dx} \quad (1)$$

J_e – electron current density, q – elementary charge, D_e – diffusion coefficient, $\frac{dn}{dx}$ – density gradient in the direction of x .



Fig. 4 (a) The diffusion of particles moving into the low-density area forms a concentration gradient. (b) A shorter diffusion length caused by the high recombination rate of particles. (c) A longer diffusion length caused by the low recombination rate of particles.



Similarly, for holes:

$$J_h = -qD_h \frac{dp}{dx} \quad (2)$$

J_h – hole current density, q – elementary charge, D_h – diffusion coefficient, $\frac{dp}{dx}$ – density gradient in the direction of x .

The holes and electrons in the solar cell are excited by light absorption. When the excited holes and electrons collide, recombination occurs. The mechanism of loss (such as recombination) determines the lifetime of the charge carrier, and the lifetime of the charge carrier determines the average distance that excited charges can travel (diffusion length). A higher recombination rate creates a shorter diffusion length (Fig. 4(b)), while a lower recombination rate creates a longer diffusion length (Fig. 4(c)). The diffusion length and lifetime are related with the following simple equation, $L = \sqrt{D\tau}$, where D is the diffusion coefficient and τ is the lifetime.¹³ It is worth remarking here that a longer diffusion length is an indication that the material will last longer, increasing the solar cell current output (J_{sc}) and open circuit voltage output (V_{oc}), improving the PCE and becoming an important property to look for in choosing solar cell materials.^{14–16} As reported in previous studies, the balanced diffusion length of electrons and holes in perovskite absorbers with more than 100 nm has the ability to prevent the rate of recombination of charges in the perovskite absorber.^{14,15} In 2015, Dong *et al.* found that the diffusion length of electron holes in $\text{CH}_3\text{NH}_3\text{PbI}_3$ was greater than 175 microns using a solution-growth method. This diffusion length led to greater carrier mobility with a longer lifetime, thus enhancing the cell efficiency.¹⁶ However, in other cases, on the surface before ETL, a compact hole-blocking layer is utilised.^{17,18} On the surface before HTL, an electron-blocking layer is utilised.^{19,20} This aims to block excess holes or electrons on both sides, respectively. To this extent, electron and hole ambiguities in the device architecture that cause high recombination rates can be avoided. In principle, the PCE of PSC can be significantly improved by providing appropriate paths for both electrons and holes with longer diffusion length.

Noorasid *et al.* investigated the application of titanium dioxide (TiO_2) as ETL and CuI as HTL using SCAPS simulation, extensively examining and optimizing each layer of the structure, considering factors such as type, thickness, doping concentration, and metal-work functions. The results showed that the acceptor concentration and defect density in the absorber layer had a significant impact on the PSC performance compared to other layers. Increasing the acceptor concentration improved the charge mobility of the absorber layer, but excessive concentration led to increased recombination rates, negatively affecting the PSC performance.²¹

2.1. PSC working principle

The working principle of a general PSC is described in Fig. 6, where each layer (Fig. 5) of the structure is carefully chosen to achieve optimal efficiency. This allows for the formation of different material energy levels, in addition to being thermodynamically compatible for a more robust device.



Fig. 5 A generic structure of PSC for achieving the optimal efficiency based on the energy levels of materials from each layer.

Basically, PSCs work by initiating photon absorption in the perovskite layer, which originates from photons propagating through the electrode layer and the transparent glass (ITO/FTO coated glass). Decent absorption features can be achieved with a larger incident photon energy than that of the perovskite material, as this stimulates the excitation in the active layer of the perovskite. The important elements to consider in order to obtain optimal absorption depend on the HOMO (highest occupied molecular orbital) and LUMO (lowest unoccupied molecular orbital). When an electron moves from the HOMO (valence band edge of the PSC) to the LUMO (conduction band edge of the PSC), this sets the perovskite in an oxidized state and is then neutralized by electrons exiting the HOMO of the neighbouring HTL.

Due to the different work functions between the metal electrode and the transparent electrode, it creates an internal potential which leads the exciton to separate into holes and electrons.²² Both electrons and holes are then shifted from their current positions to the ETL and HTL, causing the holes to migrate to the metallic electrode, while the electrons migrate to the transparent electrode.²² Thus, the motion of the electron moving through the network between these two electrodes causes an electric current to be generated. It is important that the energy levels in each layer of the PSC are designed in such a way to avoid the recombination of excited charge carriers for

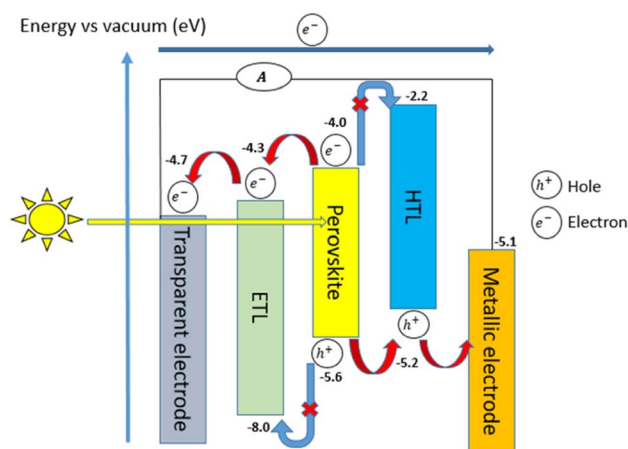


Fig. 6 A schematic of the operational principle of PSC.



the PSC to operate effectively under any conditions. Proper stratification to prevent charge carriers from mixing as they move in different directions through the cell can be achieved by introducing a LUMO energy level ETL that is below the perovskite energy level LUMO. Due to this, the electrons can move slightly further, allowing them to travel in the desired direction. In the meantime, the HOMO in HTL must exceed the HOMO in perovskite as this contributes to an efficient path for the hole to travel. In particular, each layer of cells is identical in this regard, and each layer requires a different HOMO or LUMO energy level for the charge carrier transport chain to function normally, as described in Fig. 5.²²

3. PSC ETL characteristics

The ETL is an important layer in the PSC in generating high-efficiency performance, which plays a vital role as a photo-generated electron transporter. Nevertheless, the device behaviour and performance of PSC photovoltaics still experience challenges and limiting factors, such as carrier mobility, energy band alignment, defect morphology, trap states and related interface properties.²³ Electrons are injected from the absorber layer, carried *via* ETL, and collected by the electrode in the ETL. Occasionally, an ETL is replaced with a composite photo electrode in order to improve the efficiency of the charge collection by reducing the amount of series resistance.^{24–26} In addition to this, a one-dimensional NRs ETL enables a solar cell to have a rapid charge transport. It has been discovered that managing the gaps that exist between the NRs and nanosheets can result in excellent pore filling, which is expedient for increasing both the fill factor (FF) and thereby the PCE of the material.^{27–34} Moreover, the concentrations of the fluid and the seed nanoparticle films

play a role in determining the gap between the NRs. An ideal current–voltage behaviour is expected to be achieved as a result of the effective collection of charges. This is a method for growing the external quantum efficiency, also known as EQE, where it possibly provides the utmost efficiency while avoiding thermalization.

An ETL is undoubtedly necessary to satisfy the band alignment with the perovskite layer to be effective in charge conduction. This means that both its LUMO and HOMO have higher values than the perovskite active layer, as this is the most essential quality that an ETL must possess. In order for photons to simply pass through and be absorbed by the perovskite absorber, it is necessary for this material to have a high transmittance in the UV-Vis region. In addition, the excitation generated by light absorption across the perovskite layer is split prior to the ETL or HTL collection, and the excitons are said to dissociate either at the ETL/perovskite or perovskite/HTL interface sites, as reported by Savenije *et al.*³⁵

3.1. Metal oxide as ETL materials in PSC

3.2. Titanium dioxide (TiO₂)

Generally, mesoporous TiO₂ has been widely used as the ETL in PSC due to the general features mentioned earlier, yet the energy band is the main factor of the matter. Theoretically, TiO₂ possesses a conduction band minimum of 4.2 eV and a valence band maximum of 7.63 eV (Fig. 7), which are ideal values in ensuring not only efficient electron transport from the perovskite layer, but also efficient hole-blocking capabilities at the perovskite interface.³⁶ A large number of methods have been published to improve the performance of solar cells using TiO₂ as the hole blocking layers, including spin-

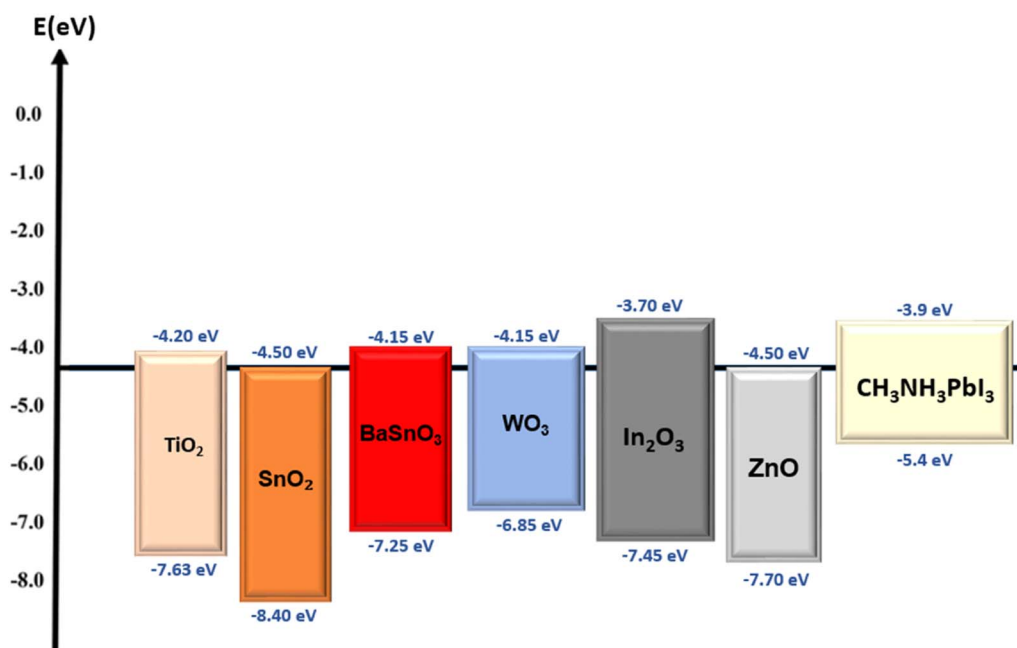


Fig. 7 A summary of the energy levels of metal oxide materials as ETL and perovskite.



coating,³⁷ spray pyrolysis,^{38,39} sol-gel,⁴⁰ chemical bath deposition,⁴¹ and atomic layer deposition (ALD).⁴² The ability to enhance the PCE of the solar cells is one of the main advantages of these methods, along with reducing the higher recombination rate and decreased photocurrent, which negatively affects the solar cells performance. In addition, the first and unique mesoscopic solar cells based on a heterojunction between perovskite and TiO₂ as the active layer are being developed,⁴³ in which the PSC nanoparticles act as both light absorber and hole conductor, eliminating the need for an additional hole transport material. This method introduced simple solution-processed mesoscopic heterojunction solar cells-based TiO₂, paving the way for low-cost and high-efficiency solar cell technology.

Besides the hole blocking layer, TiO₂ is also widely used as ETL in PSC, although there are problems with the high recombination rate⁴⁴ due to the limited electron mobility and surface defects in mesoporous TiO₂ nanoparticles.⁴⁵ In addition, the mesoporous TiO₂ layer requires a high-temperature annealing process above 450 °C for complete crystallization, and the difficulty of achieving a uniform thickness with spin-coating technique, particularly for large-scale device production.^{46,47} On the other hand, PSC-based TiO₂ as ETL also becomes unstable when exposed to UV light (sun or UV lamp) or when subjected to a forward bias for an extended period of time,⁴⁸ as the charges generated by the perovskite layers are trapped at defects and impurities within the TiO₂ layer. These TiO₂ issues are challenges that researchers are actively working on to overcome the performance and stability of these devices. One way to resolve this problem is to introduce doping into TiO₂ layers (Fig. 8), which has been shown to enhance the efficiency, V_{oc} (open circuit voltage), J_{sc} (short-circuit current density), and fill factor, resulting in an increase in the absorption, electron mobility and reducing recombination rates.⁴⁹



Fig. 8 J–V characteristics of PSC based on (a) bare TiO₂, (b) 1% Mg-doped TiO₂, (c) 2% Mg-doped TiO₂, and (d) 3% Mg-doped TiO₂. Reproduced with permission from ref. 49. Copyright 2021 Elsevier.

3.3. Tin oxide (SnO₂)

Tin oxide, SnO₂ is one of the most promising candidates of ETL for PSCs due to its superior optical and electrical properties,⁵⁰ favourable band alignment with perovskite, and flexibility. The use of SnO₂ as an ETL was initially demonstrated in dye-sensitized solar cells (DSSC), but it garnered little attention due to its relatively poor PCE when compared with TiO₂-based DSSC.⁵¹ The preliminary investigations on the effects of SnO₂ on PSCs were carried out almost simultaneously by several research groups, including Dai *et al.*⁵² and W. Ke *et al.*,⁵³ who reported efficiencies of 10.18% and 17.21%, respectively. On the other hand, Hagfeldt *et al.*⁵⁴ developed a combination of spin-coating and chemical bath post-treatment for PSC-based SnO₂ in 2016, which achieved an efficiency of 20.7%, showing that the simple solution-processed tin oxide offers advantages in terms of scalability and cost effectiveness. Furthermore, in 2017, You *et al.* achieved the highest efficiency ever recorded in the planar PSC with an efficiency of 21.6% by using a two-step method to precisely control the PbI₂ content because the excessive amount of PbI₂ leads to significant hysteresis and poor stability. This method, which provides the novel design with specific optimization techniques, can lead to a power conversion of more than 21%.⁵⁵

However, there are still issues related to the limited electron mobility of SnO₂ as an ETL in PSC, which relates to the ability of electrons to move through the material and participate in efficient charge transport processes. To address this issue, researchers have explored various strategies to enhance the electron mobility of SnO-based ETL. These include surface engineering techniques⁵⁶ and doping with other materials,⁵⁷ as shown in Fig. 9(a) and (b), respectively, where the roughness of the PSC images deposited on nickel doped-SnO₂ is smoother than on bare SnO₂, while the surface morphology by SEM for



Fig. 9 Material images of (a) AFM for bare SnO₂, (b) AFM for Ni-doped SnO₂, (c) SEM for bare SnO₂, and (d) SEM for Ni-doped SnO₂. Reproduced with permission from ref. 57. Copyright 2022 American Chemical Society.



nickel-doped SnO₂ is greater than undoped SnO₂, as illustrated in Fig. 9(c) and (d), respectively. These results are achieved through a low-temperature fabrication process with different nickel doping concentrations on SnO₂-ETL, affecting the morphology and structure of the features. This result proves that the low-temperature process and doping implied in the fabrication of Ni-doped SnO₂ ETL enables compatibility with flexible substrates, and simplifies the fabrication process of PSCs.

3.4. Barium stannate (BaSnO₃)

Another reliable metal oxide material as an ETL for PSCs is barium stannate, BaSnO₃ (BSO), in which the oxygen vacancies in BSO are naturally formed as an n-type semiconductor with a wide band gap of 3.1 eV (Fig. 7), and high electron mobility compared to other particles with suitable energy levels.⁵⁸ Due to its unique properties and versatility, BaSnO₃ is widely used in various applications, such as photocatalytic activity, sensors, gas sensing, dielectric and infrared luminescence.^{59–63} In addition, La-doped BSO single crystals and epitaxial thin films exhibit mobility values of about 320 cm² V⁻¹ s⁻¹ and 70 cm² V⁻¹ s⁻¹, respectively, which are unprecedentedly high.⁶⁴ Due to these advantages, BSO has demonstrated utility in dye-sensitized solar cells,^{65,66} where it has shown promising results leading to improved device performance and increased energy conversion efficiency. In 2015, Zhu *et al.* compared the performance of the synthesis of mesoporous BSO (mp-BSO) with the conventional mesoporous TiO₂ (mp-TiO₂) as ETLs in PSCs. As a result, it has been discovered that mp-BSO offers the highest PCE with 12.3%, and has better contact with MAPbI₃ by developing the crystals with greater cohesion. The mp-BSO device was found to have comparable performance to the mp-TiO₂, but still had certain issues. For example, under high bias conditions, the mp-BSO possesses a faster charge recombination rate.⁶⁷

To solve this problem, Zhu *et al.* fabricated novel mesoporous BaSnO₃ ETL using solvent engineering technology in 2017 by doping with lanthanum (Fig. 10) and achieved the highest PCE of 15.1%, along with the highest produced J_{sc} of 19.1 mA cm⁻².⁶⁸ To date, various morphologies of BaSnO₃ (such as NRs, nanoparticles and mesoporous) could be achieved by controlling the synthesis condition, as reported in a study from

2022,⁶⁹ which investigated the morphology of brookite TiO₂ and BaSnO₃ as ETL. This finding reinforces the notion that the morphology design of materials can improve the electron properties of the layers, leading to higher efficiency and stability of the photovoltaic devices.

3.5. Tungsten oxide (WO₃)

Tungsten oxide (WO₃) is an n-type, wide-bandgap and low-oxygen semiconductor material that has received considerable attention in recent years, with an electrical band gap of 2.6 eV to 3.6 eV (Fig. 7).^{70–73} Due to its high specific surface area and high surface permeability, nanostructured WO₃ is an excellent material for a wide range of applications. This material is a strong contender for use in optical and electrical applications^{74,75} due to its high melting temperature, photo-electrochromic properties, robustness and mechanical properties. In 2014, researchers made the first attempt to use WO₃ and TiO₂ as the core-shell electron transporting material in PSC, in which the novel structure of the WO₃-TiO₂ core-shell enhanced the electron extraction and transport capabilities.³²

Thereafter, in 2015, K. Wang *et al.* investigated the use of a low-temperature, solution-processed amorphous WO₃ thin film as ETL in PSCs in comparison to the conventional TiO₂ ETL, resulting in higher electrical conductivity and comparable photoelectric conversion efficiency with an average of 8.99% efficiency.⁷⁶ However, it also shows a lower open-circuit voltage and identifies inherent charge recombination as a limitation. Subsequently, Alexandre *et al.*⁷⁷ in 2017 proposed the printable WO₃ ETL as an alternative to TiO₂ in PSC, and found that up to 9.5% PCE can be achieved with spin-coated or inkjet-printed WO₃ compared to 15% for TiO₂ reference. The advantage of WO₃ in this experiment compared to TiO₂ is that it allows a low-temperature process to be performed under ambient conditions, although they are more sensitive to ambient humidity when illuminated, resulting in a rapid drop in photocurrent.⁷⁷ Furthermore, W-doped TiO₂ has been proposed as an ETL for PSC,⁷⁸ which uses a low-temperature solution-based method and shows an excellent performance of optimized PCE with 18.85% (Fig. 11), improving the carrier velocity and charge collection, JSC, and fill factor (FF) of the PSC. This research study suggests a promising approach for fabricating high-performance PSCs with tungsten-doped TiO₂ materials.

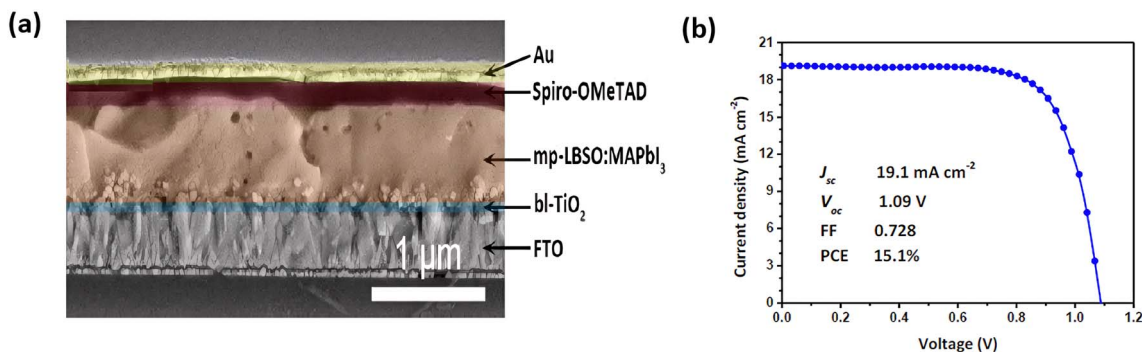


Fig. 10 (a) A cross-sectional view of SEM images from La-doped BSO in PSC. (b) J - V curves of La-doped BSO in PSC under AM 1.5 illumination. Reproduced with permission from ref. 68. Copyright 2017 Royal Society of Chemistry.



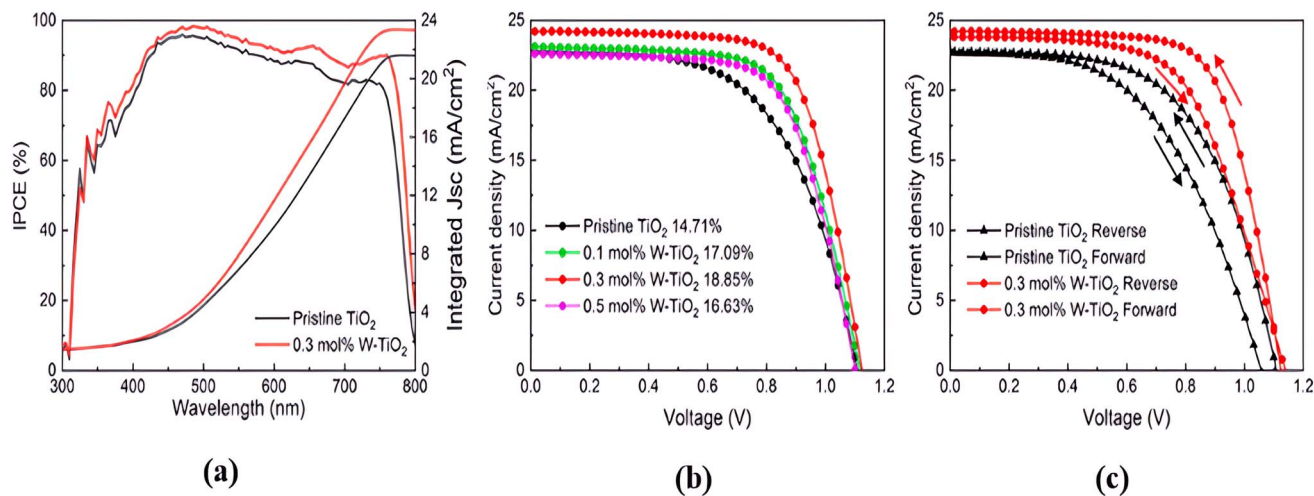


Fig. 11 (a) IPCE curves of pristine TiO₂ and 0.3% mol W-doped TiO₂. (b) IV curve of PSC with different W concentrations. (c) Reverse and forward direction of the IV curve. Reproduced with permission from ref. 78. Copyright 2021 Elsevier.

3.6. Indium oxide (In₂O₃)

Indium oxide (In₂O₃) is another promising n-type semiconductor material, with a wide band gap of 3.75 eV (Fig. 7), possessing a high mobility of 20 cm² V⁻¹ s⁻¹, and excellent thermal stability with good transparency properties.^{79–81} These properties allow In₂O₃ to be an attractive and suitable candidate for future optoelectronic applications. In 2017, Yoon *et al.* fabricated In₂O₃ ETL-based PSC and organic solar cells by solution-process, with In₂O₃ showing better photo-stability than TiO₂-PSC owing to the larger band gap. PSC showed retentions of 74% and 46% compared to the initial values measured under 1 sun illumination for 4 h.⁸²

Other than photovoltaic application, In₂O₃ is also used in gas sensing,⁸³ as reported by M. B. *et al.* in 2023, who developed a methane gas sensor-based mesoporous belt-like In₂O₃ (Fig. 12) from single-step electrospinning fabrication. The

researchers studied how the annealing temperature affected the performance of the sensor by comparing the morphology, surface defects and structural behavior of the belt-like In₂O₃ to understand the trends in gas sensing. Among the annealing temperatures, the sensor annealed at 550 °C showed the best sensing capabilities with a response of 1.1 to 90 ppm of methane at a lower operating temperature of 100 °C, and had the best response and recovery times of 36 and 44 seconds, respectively. The unique mesoporous structure of the In₂O₃ materials contributed to their enhanced capabilities by providing numerous active sites for gas interaction. The findings provide insights into optimizing the sensor's design and fabrication parameters to enhance its sensitivity, selectivity and overall performance in detecting methane gas with implications for multiple areas, including industrial safety and power generation.

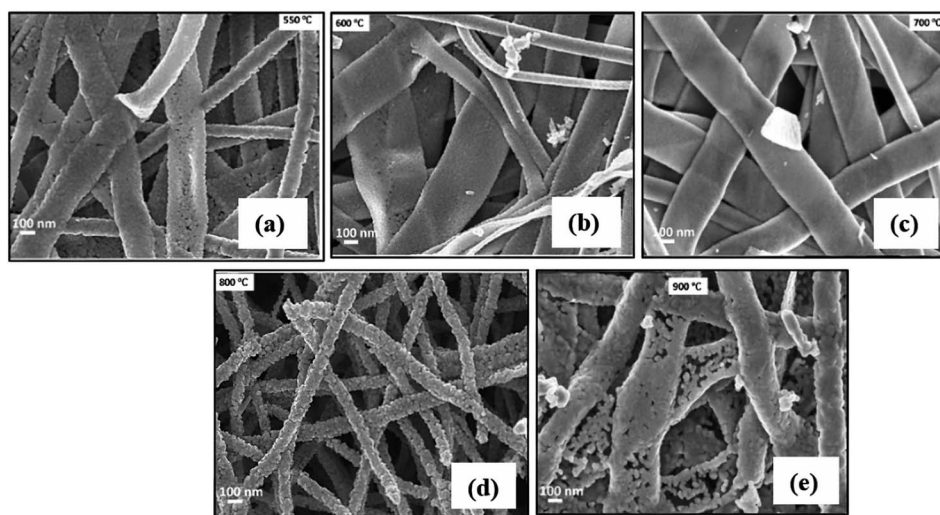


Fig. 12 FE-SEM images of the In₂O₃ belt-like products annealed at (a) 550 °C, (b) 600 °C, (c) 700 °C, (d) 800 °C, and (e) 900 °C. Reproduced with permission from ref. 83. Copyright 2022 Elsevier.



3.7. Zinc oxide (ZnO)

Solar cell applications have extensively explored the utilization of ZnO materials owing to their exceptional stability, high electrical conductivity,⁸⁴ elevated electron affinity, and remarkable electron mobility. Within the realm of binary compound semiconductors, ZnO stands as a prominent member classified under the II–VI category. It is noteworthy that ZnO materials do not exhibit significant absorption of visible light, instead manifesting as wide-bandgap semiconductors processing a bandgap range of 3.1 to 3.3 eV (Fig. 7). However, by combining ZnO with materials featuring smaller energy gaps, such as dye sensitizers and semiconductors with narrower band gap, it becomes feasible to

shift the light absorption capabilities of ZnO into the visible range, thus enhancing its suitability for solar cell applications. As previously reported, the high electron mobility of ZnO ranges from 205 to 300 cm² V⁻¹ s⁻¹ for bulk,^{85,86} and even reaches an impressive value of 1000 cm² V⁻¹ s⁻¹ for ZnO NRs. This exceptional electron mobility renders ZnO a highly attractive material for utilization in solar cell applications, particularly when compared to the electron mobility values commonly used in TiO₂, which typically range at around 0.14 cm² V⁻¹ s⁻¹. The significantly higher electron mobility values associated with ZnO underline its favourable characteristics for efficient charge transport in solar cells.

Table 1 provides a comprehensive overview of the basic properties of ZnO, a chemical compound characterized by a molar mass of 81.408 g mol⁻¹, and presents as a white solid with odourless properties. ZnO exhibits a density of approximately 5.606 g cm⁻³ and possesses a relatively high melting point of 1975 °C, at which point it undergoes decomposition, breaking down into its constituent elements, zinc (Zn) and oxygen (O). Conversely, the boiling point of ZnO is 2360 °C, indicating the temperature at which it transitions from a liquid to a gaseous state. ZnO showcases limited solubility in water, as evidenced by its solubility of approximately 0.16 mg per 100 ml of water at a temperature of 30 °C. This signifies that only a small quantity of ZnO dissolves in water under these particular conditions. One notable property of ZnO is its band gap, which denotes the energy difference existing between the valence band and the conduction band

Table 1 Basic properties of ZnO

Properties	
Molecular formula	ZnO
Molar mass (g mol ⁻¹)	81.408
Appearance	White solid
Odour	Odourless
Density (g cm ⁻³)	5.606
Melting point (°C)	1975 (decomposes)
Boiling point (°C)	2360
Solubility in water	0.16 mg/100 ML (30 °C)
Band gap (eV)	3.3 (direct)



Fig. 13 Various ZnO morphologies from SEM images of the ZnO nanostructures: (a) and (b) tetrapod structures, (c) variable diameter structures, (d) nanosheets, (e) nanoshells, (f) multipods, (g)–(i) nanorods. Reproduced with permission from ref. 94. Copyright 2010 Elsevier.



within the electronic structure of the material. In the case of ZnO, the band gap is approximately 3.3 eV, and it possesses a “direct” nature. This signifies that during transitions between the valence and conduction bands, both the momentum and energy of an electron can be conserved.⁸⁷

3.7.1. ZnO nanostructures. Nanorods, nanoflowers, nanowires, nanospheres, nanotubes, nanocrystals, and 3D nanostructures encompass the diverse range of ZnO nanomorphologies, as depicted in Fig. 13. The exceptional properties exhibited by ZnO have facilitated its extensive utilization across various fields, including solar cells. Notably, the incorporation of ZnO with different nanostructures leads to performances in PSCs, mainly due to the significant influence of the material structure on the PCE of PSCs. The presence of grain boundaries in the architecture of ZnO nanoparticles leads to numerous electron-capture events, and thereby hinders electron transfer.^{88–90}

In particular, Zheng *et al.*⁹¹ demonstrated that ZnO nanoparticles prepared by the spin coating method led to diminished PCE. This decline was attributed to the formation of pinhole surfaces, which generated defective interfaces and resulted in the loss of charge carriers. In contrast, one-dimensional nanostructured ZnO, such as NRs, exhibit significantly enhanced performance when compared to ZnO nanoparticles. This superiority primarily stems from the presence of one-dimensional single-crystal structures in ZnO NRs, facilitating direct electron pathways for electronic transport within PSC.⁹² Moreover, Mustafa *et al.* discussed the concept of piezoelectric nanogenerators-based ZnO NRs in 2022, and have demonstrated their excellent performance due to their unique characteristics, such as a wide bandgap semiconductor energy and the ability to form ordered and uniform structures.⁹³ Thus, these NRs well-defined pathways enable the observation of improved results in terms of electron transport and overall device performance.

4. Fabrication techniques in ZnO growth process

4.1. Hydrothermal

The hydrothermal fabrication method (Fig. 14) has garnered significant attention from researchers due to its inherent advantages: simplicity of equipment, cost-effectiveness, gentle preparation conditions, and an environmentally friendly approach.⁹⁵ S. Mohan *et al.*⁹⁶ conducted a study to explore the impact of reaction parameters on the structure of ZnO nanoparticles and their optical properties. They placed emphasis on the hydrothermal reaction conditions, varying the duration from 1 to 5 hours and the temperature from 100 °C to 150 °C. To analyse the influence of pH on nanomaterial development,^{97–101} a pH level of 12 was maintained. The researchers observed that the hydrothermal synthesis resulted in diverse nanostructures under different reaction parameters. Specifically, nanoflowers, NRs, and nanospheres were formed at temperatures of 100–120 °C for 2 hours, 120 °C for 1, 3, and 5 hours, and 150 °C for 2 hours, respectively.¹⁰² Additionally, the study concluded that nanoflowers and NRs exhibited superior photoluminescence, ultraviolet (UV) absorption, and an optimal bandgap.^{103,104} NRs, in particular, demonstrated enhanced heat stability compared to nanoflowers. Hence, it can be inferred that the hydrothermal approach enables the modification of nanomaterial structures in terms of size and shape. The duration of the reaction, temperature, and concentration of the reacting liquids are the primary factors determining the morphology of ZnO nanoparticles.

4.2. Sol-gel spin coating

The sol-gel process, using spin-coating techniques as illustrated in Fig. 15, is highly attractive and widely used due to its straightforward nature, flexibility, and ease of doping. Moreover, it presents the opportunity to produce large-scale coatings at an affordable cost.¹⁰⁵ In their study, Mursal *et al.*¹⁰⁵ used sol-



Fig. 14 A schematic diagram of the ZnO NRs hydrothermal process.



Fig. 15 Sol-gel spin coating in the ZnO growth process.



gel spin coating techniques to fabricate ZnO thin films. They investigated the influence of the annealing temperature on the optical properties, such as transparency, absorbance, and bandgap. The results revealed a notable correlation between the annealing temperature and the properties of the ZnO thin films. Notably, the optical bandgap of the ZnO films decreased as the annealing temperature increased. Additionally, the optical absorption coefficient of the ZnO films was found to be affected by the annealing temperature.

In another study, M. A. Shafi *et al.*¹⁰⁶ produced ZnO thin films for CZTS-based solar cells by varying the number cycles of spin coats to obtain different thicknesses of ZnO thin films on FTO substrates. A ZnO layer was formed after 8 and 10 cycles, and exhibited excellent surface density, crystallinity, roughness, and absence of pinholes with the observed bandgap ranging between 3.24 and 3.29 eV. Consequently, it can be inferred that both annealing temperature and rotation cycle significantly impact changes in the bandgap and structure of the solar cells.

4.3. Sputtering

Magnetron sputtering (Fig. 16) is a highly suitable technology for the physical vapour deposition process, and enables the deposition of uniform ZnO layers over large areas at room temperature. It has been reported that magnetron-sputtered ZnO films significantly reduce the presence of hydroxyl groups on the film surface, positioning them as promising ETL for PSCs, as mentioned previously.^{107,108} H. Niu *et al.*¹⁰⁹ used magnetron sputtering to produce high-quality ZnO films as efficient ETLs for PSC in 2021, and formed a ZnO/SnO₂ bilayer structure by spin-coating a thin layer of colloidal SnO₂ onto the ZnO underlayer. The bilayer ZnO/SnO₂ ETL showed an impressive efficiency of 15.82%, surpassing 13.04% PCE achieved by the ZnO-based device. This study provides evidence that bilayer magnetron sputtering significantly improves PCE by facilitating good energy alignment in the cascade and efficient electron extraction at the ETL/perovskite interface. Thus, the bilayer ZnO/SnO₂ ETL achieved a significantly higher efficiency of 15.82% compared to the PCE of the ZnO-based device of 13.04%, underscoring the potential of magnetron-sputtered bilayers for substantial PCE improvements.



Fig. 16 A schematic of the ZnO RF magnetron sputtering system.

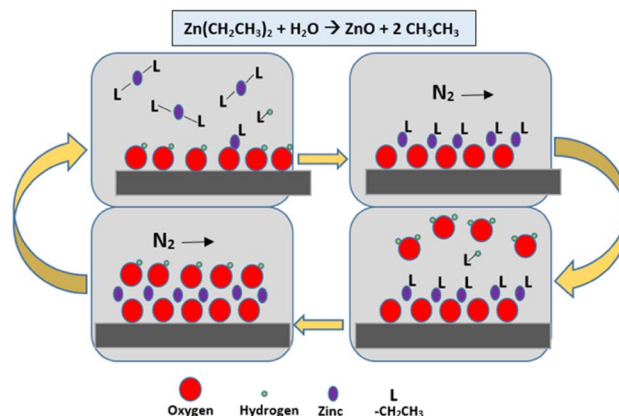


Fig. 17 A schematic illustration of ZnO atomic layer deposition synthesis.

4.4. Atomic layer deposition

Atomic layer deposition, often known as ALD (Fig. 17) is a vapour phase thin film deposition process that is performed at low substrate temperatures, allowing for the uniform deposition of thin films. This scalable method is particularly attractive for industrial applications due to its suitability for large substrates. R. Pietruszka *et al.*¹¹⁰ explored the concept of utilizing ALD and hydrothermal processes to fabricate solar cell components (Zn/Si heterojunction solar cell), presenting greener and cost-effective alternatives that eliminate the need for extreme heat or potentially harmful chemicals. The textured structures achieved a maximum efficiency of 14%, while the planar structures reached 10%. In 2017, Shibin *et al.*¹¹¹ used ALD to deposit Al₂O₃ monolayers on hydrothermally synthesized ZnO NRs to passivate interface defects. The PSC based on annealed ZnO-NRs passivated with ALD-Al₂O₃ demonstrated a maximum PCE of 16.08% and an average PCE of 15.08%.¹¹¹ Therefore, the ALD of monolayers presents a novel technique for defect passivation in PSC, and it is worth noting that the different fabrication methods employed for ZnO NRs arrays can affect the performance of the PSC.

5. Influences of dopants on ZnO as ETL in PSC

The introduction of impurities, also known as dopants, into semiconductor materials is defined as doping engineering. This allows the properties of the semiconductor to be modified to passivate defects in the bulk and at the interface, and improve bandgap alignment for enhanced charge carrier mechanism. Doping engineering involves incorporating specific elements of atoms into the semiconductor crystal lattice, thereby influencing the behavior of charge carriers (electrons or holes) in the semiconductor material and subsequently the performance of the solar cell device. Several authors have developed more sophisticated methods to enhance the electronic characteristics of ZnO-based ETL in PSCs through doping. Doping is typically employed to improve free charges and conductivity in solar





Fig. 18 A schematic image of the ZnO^{2+} cation substitution or passivation in the doping engineering process with lanthanum as an additive dopant.

cells, where it can drastically alter the structure of ZnO.¹¹² As shown in Fig. 18, doping of ZnO can be achieved by exchanging either the Zn^{2+} cation or the O^{2-} anion, where metals typically serve as cationic dopants and non-metals act as anionic dopants. Substituting Zn^{2+} with another cation can transform the structure of the conduction band (CB), while replacing the O^{2-} anion with a different anion alters the valence band (VB) energy.¹¹³ In addition, numerous studies have demonstrated the effectiveness of ZnO doping for PSC applications, utilizing various dopant materials including metals (alkali metals, metalloids, transition metals), non-metals (halogens), and lanthanides (rare earth elements) in the ETL of solar cells.

On the other hand, doping of ZnO as an ETL in PSCs can be understood in the context of Fick's law of diffusion, a fundamental principle that describes the movement of substances within materials, as explained in eqn (1) (Section 2) of this paper. Fick's law provides valuable insights into the movement of charge carriers, especially electrons, within the ZnO ETL of PSCs. According to eqn (1), Fick's law establishes a mathematical relationship between the flux (J), representing the rate of diffusion, and the concentration gradient (dn/dx) of the diffusing substance. In the context of PSCs, this gradient arises between electron-rich regions (n-type ZnO) and regions with fewer electrons (p-type perovskite or other layers), driving

electron diffusion through the ETL. When dopants are introduced into the ZnO ETL, they significantly affect the diffusion properties, particularly the diffusion coefficient D_e , which measures the electron mobility. Doping increases the electron mobility and accelerates electron diffusion, resulting in higher D_e values and faster electron transport. Furthermore, dopants reduce the trapping and recombination centers within the ZnO ETL, enhancing the electron diffusion efficiency by preventing temporary electron capture in traps. Dopants can also align energy levels within the ETL, reducing energy barriers and allowing for more efficient electron diffusion. The overall goal of doping engineering is to optimize the electron diffusion efficiency within the ETL. This optimization leads to improved charge extraction and collection, minimizes losses due to recombination, and ultimately increases the overall efficiency of PSCs. In summary, Fick's law provides a fundamental framework for understanding how dopants in the ZnO ETL influence the charge carrier diffusion, a critical factor for improving the PSC performance.

5.1. Aluminium-doped ZnO (ZnO:Al)

Both aluminium (Al) and zinc (Zn) are abundant, cost-effective, and non-toxic elements with ample availability. Doping with Al offers multiple benefits, including enhancing the stability of the ZnO/perovskite interface¹¹⁴ by modifying the fundamental properties of ZnO,¹¹⁵ but also significantly improves the carrier concentration and electron mobility of the material performance, as discussed in Section 5 regarding diffusion theory.¹¹⁶ In 2014, Mahmood *et al.* revealed that the PCE of PSC with AZO films deposited *via* electro spraying rises from 10.8% to 12.0%.¹¹⁷ The authors were able to achieve a high efficiency of 12% in PSC by precisely controlling the surface morphology of ZnO and Al-doped ZnO thin films through electrostatic spraying. Thereafter, A. Baltakesmez *et al.* presented the use of an inverted hybrid perovskite solar cell with an Al-doped ZnO substrate and anode electrode. The researchers found that AZO substrates prepared using RF magnetron sputtering with optimal growth parameters have similar electrical and optical

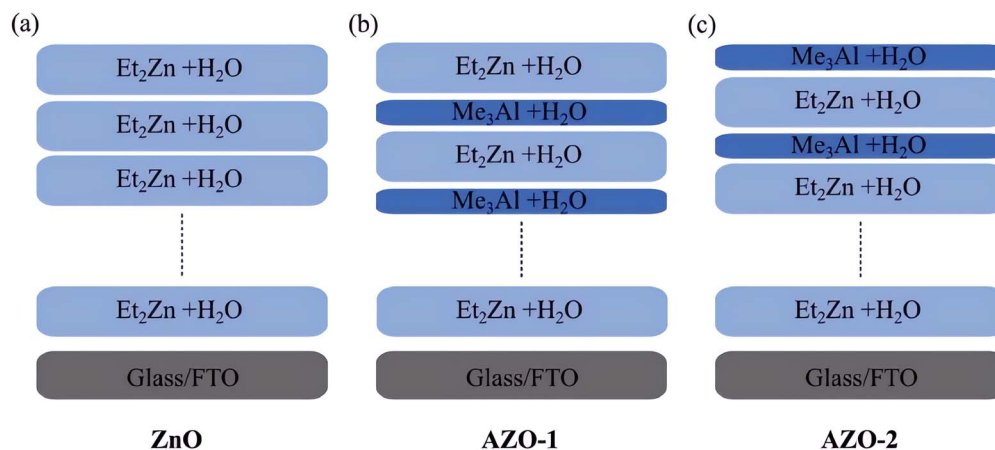


Fig. 19 A schematic representation of the growth process of ZnO and AZO films by ALD. This figure has been reproduced from ref. 119 with permission from Wiley, copyright 2022.



properties to commercial indium tin oxide (ITO), with the PCE at approximately 10%, which is close to the PCE of around 13% for the cell based on the ITO substrate under one-sun illumination.¹¹⁸ Later in 2022, Joanna *et al.* reported the potential of atomic layer engineering techniques (Fig. 19) to improve the performance and stability of PSC through the precise control of Al-doped ZnO thin films. The authors demonstrate that by tailoring the properties of AZO films, such as optimizing their conductivity, transparency, and interface characteristics, the efficiency and long-term stability of PSC can be significantly improved.¹¹⁹ In the context of gas sensing for H₂S detection, the gas sensing performance of the Al-doped ZnO thin films exceeded that of undoped ZnO films. This increase in performance is primarily due to the introduction of p-type conductivity and the consequent expansion of the films' surface area.¹²⁰

5.2. Silver-doped ZnO (ZnO:Ag)

Silver (Ag) dopants are of particular interest because of their characteristic surface plasmon resonance (SPR) effect, high surface area, and stable colloidal solution in water or organic solvents.¹²¹ In 2013, Ozlem *et al.* reported the influence of the silver content on the structural and optical properties of the resulting ZnO nanoparticles by fabricating silver-doped ZnO (Ag:ZnO) using precipitation method.¹²² It was revealed that the excessive amount of Ag led to a decrease in photocatalytic activity, and choosing the optimal amount of silver dopant is important to achieve the maximum photocatalytic performance. Another study¹²³ investigated the impact of reduced graphene oxide (rGO) sheets on the enhancement of PCE in PSC using nickel (Ni)-doped ZnO and silver (Ag)-doped ZnO as ETL. The results indicated that Ag doping at 5 wt% played a beneficial role in separating charge carriers and reducing the recombination rate by introducing defect states in ZnO-based PSC.

Besides that, according to the findings of G. Bagha *et al.*, shallow trap generation by Ag doping contributes to charge extraction. As a result, a rGO/Ag-doped ZnO bilayer ETL was created as a new approach to fabricate stable planar PSCs with low charge recombination rate and good performance, as discussed in Section 5 regarding diffusion theory. After the addition of rGO sheets and Ag doping, V_{oc} increases and causes the carrier recombination to decrease, thus leading to an increase in PCE by 30% from 7.40% to 11.03%, as shown in Fig. 20.¹²⁴ Therefore, Ag is not a dopant in its own right; in order to obtain a greater PCE, the Ag-doped substance must be modified.

5.3. Iodine-doped ZnO (ZnO:I)

Iodine is a chemical element with the symbol I and belongs to the halogen group of the periodic table, which is a lustrous, purplish-black solid at room temperature. It is highly reactive, classified as a non-metal known for its distinctive violet colour,

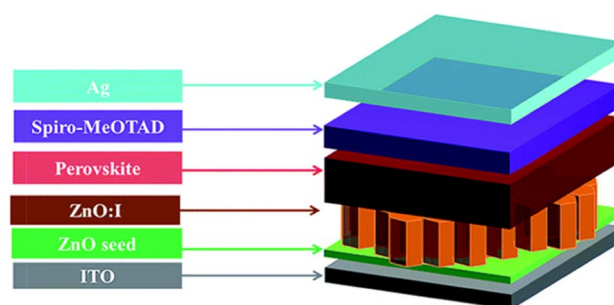


Fig. 21 A schematic illustration of the PSC configuration, in which a dense and homogeneous perovskite cap layer completely covers the ZnO:I nanopyllars. Reproduced with permission from ref. 125. Copyright 2017 Royal Society of Chemistry.

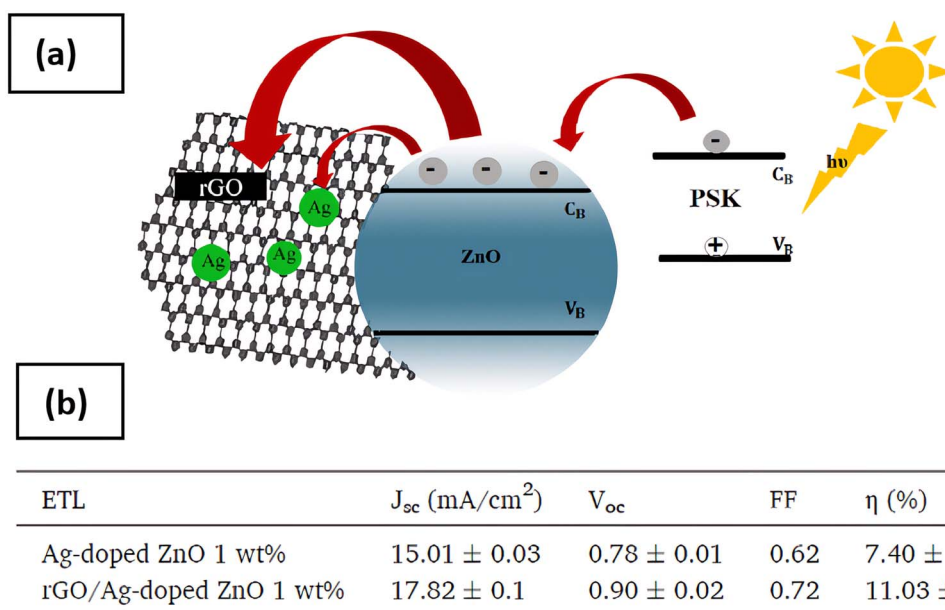


Fig. 20 (a) The functionality of rGO/Ag-doped (1 wt%) bilayer ETL of PSC ZnO. (b) J - V characteristics of Ag-doped and rGO/Ag-doped ZnO (1 wt%). Reproduced with permission from ref. 124. Copyright 2021 Elsevier.



and forms a vapour when in contact. In 2017, Zheng *et al.*¹²⁵ fabricated one dimensional (1D) ZnO:I nanopillars ETL (Fig. 21) for PSC using a one-step spin coating process, and the PCE was boosted up to 18.24%, surpassing the previous highest efficiency achieved with 1D ETL-based PSC. It has been proven that the exceptional cell performance is due to the following properties of the low aspect ratio (LAR) ZnO:I nanopillars ETL with a refined flat surface. This research introduces a novel and efficient 1D ETL approach by incorporating alien element doping and controlling the crystal orientation growth, leading to the maximization of the PSC performance. However, there are only a few studies that have reported on iodine as a dopant in PSC because most of them are used in dye-sensitized solar cells that greatly improved the PCE.^{126–128} As dye-sensitized solar cells DSSCs have a longer research and development history compared to PSCs, iodine has been extensively studied and demonstrated as an effective dopant in DSSCs, showing improved charge transport and stability, as discussed in Section 5 regarding diffusion theory. Therefore, many studies focused on utilizing iodine in DSSCs rather than in PSC, where its efficacy was already established. In addition, the utilization of iodine-doped ZnO as ETL in PSC research has ceased since 2017. This is primarily due to a combination of factors including iodine migration and degradation, as well as challenges associated with iodine loss. Iodine poses a unique challenge to the stability of PSC, especially in the presence of metallic electrodes. When these PSCs are exposed to such electrodes, the electric field generated between them can initiate a process in which iodine ions are displaced. This can lead to the migration of I^- ions, which ultimately leads to degradation of PSC performance.¹²⁹ Besides that, iodine loss resulting from a range of factors such as moisture, oxygen, and

light exposure, poses a substantial challenge in the pursuit of enhancing both the PCE and long-term durability of PSCs.^{130,131} Nevertheless, in the field of sensor technology, there has been a recent breakthrough involving the use of I-doped ZnO for detecting ethyl paraben (EP),¹³² where EP is commonly employed as a preservative in a wide range of personal care products, and the ability to detect it at low concentrations is of utmost importance for ensuring product quality and safety. This study identified an optimal level of iodine doping in ZnO that significantly increases the performance of the sensor and positions it as a promising material for EP detection. I-doped ZnO has also proven to be highly effective in detecting high-quality latent fingerprints on a variety of substrates, demonstrating its significant potential in forensic investigations.¹³³

5.4. PbX_2 -doped ZnO ($ZnO:PbX_2$)

PbX_2 compounds are ionic solids that exhibit solubility in water and polar solvents. These compounds possess high melting and boiling points, and demonstrate stability at elevated temperatures.¹³⁴ In materials science, PbX_2 compounds find diverse applications, particularly as precursors for the synthesis of perovskite materials utilized in solar cells.¹³⁵ In the context of PSC, an innovative methodology involving the incorporation of PbX_2 ($X = I^-$, Cl^- , and CH_3COO^-) compounds as dopants within ZnO layers has been reported by Wu *et al.*, where the ZnO films were prepared using a sol-gel method. PbX_2 ($X = I$, Cl , and CH_3COO) was then incorporated into the ZnO films by adding the PbX_2 precursor to the ZnO precursor solution under 200 °C annealing temperature. From the study, the efficiency of the all-inorganic perovskite solar cells is boosted from 17.64% to 20.04% with a fill factor of 84.14% and a highest reported open-circuit voltage (V_{oc}) of 1.28 V among highly efficient all-



Fig. 22 (a) Energy level alignment of ETLs and the perovskite layer. (b) Visualization of the impact of the enhanced oxygen vacancy in PbX_2 :ZnO. This figure has been reproduced from ref. 136 with permission from Wiley, copyright 2022.



inorganic PSCs. The PbX_2 PSCs also exhibit enhanced stability, retaining above 97% of the initial PCE after 1000 hours and demonstrating improved photostability. The integration of PbX_2 leads to increased density of oxygen vacancies (Fig. 22(a)), consequently leading to a higher electron concentration through n-type doping, with the Fermi levels becoming closer to the conduction band minimum (CBM), as shown in Fig. 22(b).¹³⁶ A higher concentration of oxygen vacancies on the ZnO surface leads to an increase in the Fermi level through n-type doping, which in turn increases the wettability and facilitates the subsequent deposition of an advanced perovskite film on the surface. Besides that, it has been reported that using an ETL with a shallower energy level has increased the VOC, while establishing favorable conduction band alignments with the perovskite layer.^{137,138} Consequently, charge recombination was effectively suppressed and charge extraction became more efficient, demonstrating the significant improvement in both VOC and FF, as discussed in Section 5 regarding diffusion theory. Thus, this significant improvement not only represents a new milestone in the field of purely inorganic PSCs, but also represents a groundbreaking advancement in photovoltaic

technology. It expands the applicability of PSCs to a wider range of critical applications that require high-voltage outputs, thereby enhancing their versatility and utility in the field of renewable energy.

6. Recent advances and future perspective

6.1. Lanthanum-doped ZnO (ZnO:La)

Lanthanum (La), a cationic rare earth metal oxide in powder form, exhibits a wide range of useful industrial qualities, with the particle size of the powder being in the nanoscale of 1–100 nm. In electrochemistry, lanthanum oxides are used to produce organic catalysts. It is used as an electrode, and the burning rate of propellants can also be increased by lanthanum oxide.^{139,140} Furthermore, previous research suggested that adding a small amount of La to doped ZnO can significantly improve its photocatalytic properties by: (i) enhancing charge separation; (ii) increasing surface area and active sites; (iii) shifting in band-edge positions; and (iv) reducing charge carrier recombination,^{141–143} as discussed in Section 5 regarding

Table 2 Lanthanum as a dopant in solar cell applications

Solar cell	ETL	Morphology	Method	PCE (%)	Year	Ref.
PSC	La/BaSnO ₃	Bare surface nanocrystal	Laser manufactured	23.74	2022	145
Planar PSC	La/SSO	Smooth crystal grain	Low-temperature solution	18.7	2021	146
Cs ₂ TiBr ₆ -based PSC	La/BaSnO ₃	—	Numerical analysis by SCAPS 1D	29.13	2021	147
Dye-sensitized solar cell (DSSC)	La/ZnO, Dy/La	Nanorods	Spin coating, hydrothermal	6.22, 6.52	2021	144
Inverted organic solar cell (IOSC)	La/ZnO	Nanospheres	Sol-gel	4.34	2020	148
Planar PSC	Al/La/SnO ₂ (co-doped)	—	Low-temperature solution	19.13	2020	149
Planar structure hybrid PSC	La/SnO ₂	Hexagonal disc	Low-temperature solution	17.08	2019	150
Dye-sensitized solar cell (DSSC)	La/SrSnO ₃	Nanocoral	Wet-chemical route	4.7	2018	151
PSC	La/TiO ₂	Hexagonal disc	Spray-pyrolysis	17.2	2017	152
Inverted organic solar cell (IOSC)	La/TiO ₂	Small hole grain	Sol-gel	2.45	2017	153
PSC	La/BaSnO ₃	Nanobead	Hydrothermal	15.1	2017	68

Table 3 La-doped ZnO in other applications

Application	Method	Morphology	Advantages	Year	Ref.
UV photodetectors	Spray technique	Nanorods	5% of La-ZnO layers excel in UV photodetection with a sensitivity of 5.07	2023	154
Rhodamine B degradation	Electrospinning-calcination	Nanofibers	La-doped ZnO catalyst more efficient than undoped ZnO in Rh B degradation Great potential in Rh B dye wastewater	2022	155
Dye-photodegradation	Electrospinning-calcination	Nanofibers	La-doped ZnO showed higher photocatalytic activity compare to undoped ZnO Successfully used in photodegradation of CR dye	2019	156
Piezoelectric nano-generator	Wet chemical route	Nanorods	The output of piezoelectric was enhanced by 4–5 times with 5% of La-doped ZnO	2019	157
Antibacterial study	Co-precipitation	Nano-sphere	High stability in sensor analysis with more than 4500 cycles Lanthanum doping concentration controls ZnO nanoparticle size, band gap, and photoluminescence Suitable for antibacterial activities as treatment for kidney stone and typhoid fever	2017	158
Metasystox degradation	Microwave assisted	Nanorods	La-doped ZnO completes metasystox mineralization faster than pure ZnO La-doped ZnO metasystox safe for plant growth	2014	159



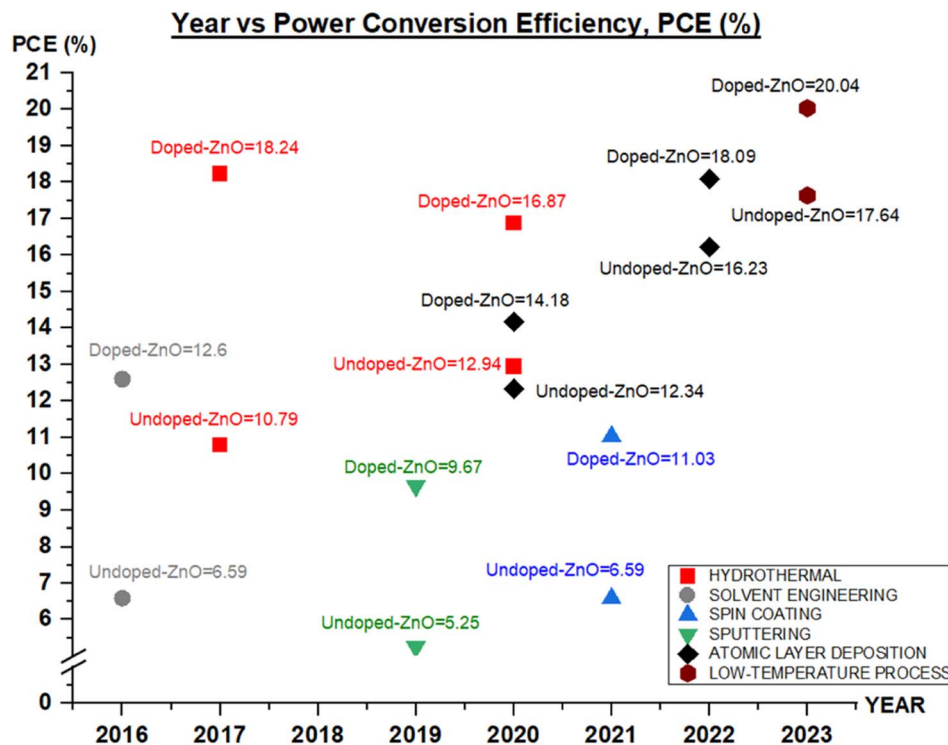


Fig. 23 Current status of PCE ETL with doped-ZnO and undoped-ZnO in PSC with different fabrication techniques in the last decade.

diffusion theory. In a study by C. Chamanzadeh *et al.*,¹⁴⁴ La-doped ZnO were synthesized using the hydrothermal method. Varying the concentration of La dopants resulted in a remarkable improvement of the PCE with a 63.6% increase compared to conventional DSSC utilizing TiO₂ nanoparticles. On the other hand, the PCE increment of La-doped ZnO is almost 97% when compared to the pure ZnO. Synthesis of optical property measurements showed enhanced absorption, with the changes in the bandgap energy almost similar to the pure ZnO with 3.218 eV, showing higher absorbance coefficient and resulting in higher PCE. The photovoltaic performance data for DSSCs, such as improved PCE, current-voltage (*I*-*V*) curves, fill factor, short-circuit current (*J*_{sc}), and open-circuit voltage (*V*_{oc}), demonstrated significant enhancements achieved through the use of La-doped ZnO NRs compared to conventional

photoanodes. These comprehensive findings collectively underscore the promise of utilizing La-doped ZnO nanorods in enhancing the photovoltaic performance of DSSCs. Given the promising results observed in this study, particularly in the enhancement of solar cell performance through the use of ZnO nanorods doped with La, it has become intriguing to consider the application of these findings in the field of PSCs since the application of La-doped ZnO as an ETL in PSCs remains relatively unexplored. However, in the context of other solar cell types, the incorporation of La (frequently in conjunction with other elements) has exhibited notable enhancements in PCE, as outlined in Table 2. The utilization of La-doped ZnO extends beyond the realm of solar cell applications. It has found diverse applications, ranging from piezoelectric nano-generators to dye photodegradation, drug delivery systems, wastewater

Table 4 A summary of ZnO doped in PSCs

Solar cell	ETL	Nanostructure	Dopant	Ionic radius (Å)	Method	Year	PCE (%)	Ref.
PSC	ZnO	Amorphous	Aluminium (Al)	0.54	Sol-gel	2023	7	160
CsPb(I _{1-x} Br _x) ₃ , PSC	ZnO	Pinhole-free	PbX ₂ , X = I, Cl, CH ₃ COO	1.19	Low-temperature process	2022	20.04	136
PSC	ZnO	—	Aluminium (Al)	0.54	Atomic layer deposition	2022	18.09	119
Planar PSC	ZnO	Nanospheres	Graphene oxide silver-doped bilayer, *rGO/Ag-doped	1.26	Modified hummers, combustion, spin-coating	2021	11.03	124
PSC	ZnO	Nanorods	Nitrogen (N), graphene (G), N/G-doped	<i>N</i> = 1	Hydrothermal	2020	16.87	161
Flexible PSC	ZnO	Nanowires	Aluminium (Al)	0.54	Atomic layer deposition	2020	14.18	162
Planar PSC	ZnO	Amorphous	Gallium (Ga)	0.62	Graded sputtering	2019	9.67	163
PSC	ZnO	Nanopillars	Iodine (I)	2.2	Hydrothermal	2017	18.24	125
PSC	ZnO	Hexagonal disc	Aluminium (Al)	0.54	Solvent engineering	2016	12.6	115



Table 5 A summary of other doping elements in other solar cells

Solar cell	ETL	Nanostructure	Dopant	Ionic radius (Å)	Method	Year	PCE (%)	Ref.
Carbon-based PSC	TiO ₂	Pin-hole free	Lithium (Li)	0.76	One-step	2023	8.09	164
PSC	SnO ₂	Nanocrystal	Chlorine (Cl)	1.81	Chemical bath deposition	2023	25	165
PSC	TiO ₂	Porous	Erbium (Er), neodymium (Nd)	0.881, 0.995	Solid-state grinding	2023	12.50	166
PSC	TiO ₂	Mesoporous	Tin (Sn)	0.71	Sol-gel	2022	20.55	167
PSC	TiO ₂	Nanograins	Nitrogen-based group (N)	1	Sputtering, sol-gel	2022	5.33	168
Inorganic halide PSC	TiO ₂	Nanorods	Cesium (Cs)	1	Hydrothermal	2022	12.72	169
Planar PSC	SnO ₂	Nanoflowers	Nickel (Ni)	0.69	Low temperature solution	2022	19.7	57
Planar PSC	SnO ₂	Right bipyramids	Copper (Cu)	0.73	Low-temperature, spin coating	2022	21.35	170
Heterojunction solar cell	ZnO	—	Aluminium (Al)	0.54	Sol-gel	2022	8.984	171
Pbs quantum dots solar cell	ZnO	Nanograins	Magnesium (Mg)	0.72	Hydrothermal, hot injection	2021	9.46	172
Inverted polymer solar cell	ZnO	—	Indium (In)	0.8	Solution-processed, sol-gel	2021	5.99	173
Inverted polymer solar cell	ZnO	Smooth grain	Indium (In), gallium (Ga), IG/ZO	In = 0.8, Ga = 0.62	Sol-gel spin coating	2021	8.72	174
Flexible PSC	ZnO	Nanorods	Cobalt (Co)	0.745	Hydrothermal	2020	7	175
Polymer solar cell	ZnO	Nanospheres	Gold (Au)	0.85	Seed-growth	2020	8.801	176
DSSC	ZnO	Branched network	Manganese (Mn)	0.46	Hydrothermal	2020	4.20	177
Quantum dot-sensitized solar cell (QDSSC)	TiO ₂	Nanospheres	Lanthanum (La)	1.061	Co-precipitation	2020	3.28	178
DSSC	ZnO	Nanobead	Copper (Cu)	0.73	Hydrothermal	2020	1.50	179
Polymer solar cell	ZnO	—	Cadmium (Cd)	0.97	Solution-processed	2020	14.7	180
Inverted polymer solar cell	ZnO	—	Gallium (Ga)	0.62	Low-temperature sol-gel	2019	7.7	181
Organic solar cell	ZnO	—	Lithium (Li), cesium (Ce), Li/Ce	Li = 0.76, Cs = 1.67	Sol-gel	2019	9.074	182
PSC	ZnO	Nano-ginger	Aluminium (Al)	0.54	Low-temperature	2019	7	183
Organic solar cell	ZnO	Nanospheres	Aluminium (Al)	0.54	Sol-gel	2019	3.24	184
Inverted polymer solar cell	ZnO	—	Aluminium (Al)	0.54	Sol-gel	2016	3.25	185

purification, and antibacterial activity, as comprehensively summarized in Table 3.

6.2. Doping elements of ZnO-ETL in PSC and other doping elements in other solar cells

Doping typically leads to increased conductivity in ZnO, and this effect depends both on the electronegativity of the dopant element and the difference in radius between the dopant and zinc ion. Additionally, to address the instability of ZnO at high temperatures, doping with other elements has been employed to enhance its properties. In particular, ZnO elemental doping has shown promising results in improving the efficiency of PSCs through various fabrication methods, as depicted in Fig. 23 and detailed in Table 4. Doping is widely recognized as a crucial technique for enhancing solar cell performance, leading to its extensive application across different types of solar cells, as summarized in Table 5.

7. Discussion

In this review, we thoroughly discuss various materials of the ETL in PSCs and other emerging solar cell technologies. ZnO NRs have demonstrated superior features for the extraction and

conduction of electrons, primarily due to their one-dimensional single-crystal structures, which facilitate efficient electronic transport in PSCs. Various fabrication techniques have been

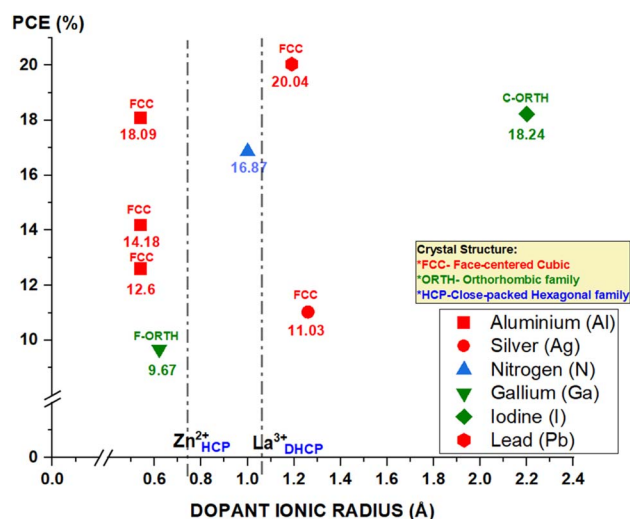


Fig. 24 Summary of impacted factors on PCE for various types of dopants in ZnO ETL-PSC.



Table 6 Impacted factors on PCE for various types of dopants in ZnO ETL-PSC

Dopants	Crystal structure	Valence state	Ionic radius (Å)	Electronegativity	PCE (%)
Zinc (Zn)	HCP	+2	0.74	1.65	—
Lanthanum (La)	DHCP	+2	1.061	1.10	—
Aluminium (Al)	FCC	+3	0.54	1.61	18.09 (ref. 119) 14.18 (ref. 162) 12.60 (ref. 115)
Silver (Ag)	FCC	+1	1.26	1.93	11.03 (ref. 124)
Nitrogen (N)	NA	−3	1	3.04	16.87 (ref. 161)
Gallium (Ga)	F-ORTH	+3	0.62	1.81	9.67 (ref. 163)
Iodine (I)	C-ORTH	−1	2.2	2.66	18.24 (ref. 125)
Lead (Pb)	FCC	+4	1.19	2.33	20.04 (ref. 136)

explored for the production of highly oriented ZnO NRs, where the hydrothermal synthesis method has appeared as the best method due to its simplicity, cost-effectiveness, mild conditions and environmental friendliness. Moreover, the ZnO NRs have exhibited remarkable thermal stability compared to other nanostructures during the hydrothermal process.

To enhance the functionality of ZnO as an ETL, doping engineering offers a viable approach to tuning the properties of the semiconductors. Apparently, doping involves the process of implanting certain elements into the crystal lattice, which offers the capability to modulate the energy band gap and improve the photocatalytic properties of ZnO, thereby increasing the PCE of solar cells. The physical and chemical properties of doped ZnO strongly depend on the choice of doping elements. A wide range of dopants on ZnO ETL PSCs, characterized by ionic radii closely resembling that of ZnO (0.74 Å), have been extensively studied as listed in Table 4, and their promising performance in terms of efficiency is shown in Fig. 24. This figure of merit highlights the importance of carefully selecting dopants for the ZnO ETL in PSCs, where the ionic radii of dopant ions and host Zn ions should be remarkably similar, allowing for nearly identical substitution within the ZnO vacancy. This process effectively involves the impurity (dopants) occupying the voids or interstices among the host Zn atoms, particularly within the ZnO vacancy sites. This criterion aims to maintain the structural integrity and stability of the ETL, which can significantly impact the overall PCE of the PSCs. Nevertheless, the use of lanthanum with 1.061 Å ionic radius as a dopant in ZnO for PSC has not yet been studied and its effectiveness has not yet been explored. Therefore, this paper proposes a novel approach for incorporating lanthanum into ZnO as an ETL in PSC. Due to the ionic interaction between Zn^{2+} and O^{2-} , pure ZnO does not have a free carrier charge. However, if Zn^{2+} ions are replaced by La^{3+} ions in the ZnO matrix, this leads to an increase in the number of free electrons in the valence band. This could instigate the electrons in the valence band to require more energy to jump to the free states at the conduction band, as the lowest state of the conduction band is already occupied by electrons. This decreases the amount of energy required to excite the electron into the lowest energy empty band in the semiconductor, and therefore increases the number of electrons that have enough energy to cross this gap. Besides that, when Zn^{2+} replaces La^{3+}

in the crystal lattice, the average ionic radius of the material may decrease because zinc has a smaller ionic radius than lanthanum.

Other than the ionic radius, other factors are crucial in assessing the compatibility of La-doped ZnO within the context of its application as a dopant, particularly in solar cell technologies. These impacted factors, including the crystal structure, valence band alignment and electronegativity, are compared with other dopants, as shown in Fig. 24 and Table 6. It is noteworthy that among the various dopants used, La bears a striking similarity to Zn due to their shared crystallographic lattice, where both ZnO and La belong to the hexagonal crystal family. This shared hexagonal symmetry promotes their compatibility as dopants within the host lattice. A matching crystal structure minimizes the lattice defects and strain, ensuring better crystalline quality and electrical properties. Furthermore, La can seamlessly integrate into the ZnO valence band with a +2 valence state. This similarity in valence states aligns their energy levels, facilitating charge transport and enhancing compatibility. This allows La to contribute charge carriers to the valence band, which is crucial for optimizing the electronic properties. Additionally, the +2 valence state of La makes it relatively easy to dissolve and incorporate into the ZnO lattice. The presence of La ions with a +2 charge helps maintain charge balance within the host material. This ease of dissolution promotes homogeneous distribution, preventing the formation of defects that could impair the structural integrity. Despite La having a lower electronegativity than oxygen, it does exhibit some degree of electronegativity compatibility with Zn. This modest electronegativity similarity between La and Zn can facilitate stable dopant–host interactions. As a result, it helps to minimize the charge imbalances and reduce electrostatic forces, which, in turn, contributes to the overall compatibility of La-doped ZnO.

8. Conclusion

The improved quality of ETL will lead to an effective PCE, and allow the cells to become genuinely revolutionary in PSCs technologies. In this paper, several types of metal oxide as electron transport materials were discussed, including TiO_2 , SnO_3 , BaSnO_3 , WO_3 , In_2O_3 and ZnO. Among these materials,



ZnO with NRs morphology has demonstrated favourable performance as an ETL due to its direct electron pathway in solar cells. As part of this comprehensive review, we have discussed various fabrication techniques that have been prominently utilized over the past decade. These techniques span a spectrum of methods, including hydrothermal synthesis, solvent engineering, spin coating, sputtering, atomic layer deposition, and low-temperature processes. Hydrothermal synthesis has emerged as the preferred method compared to other syntheses for producing ZnO NRs owing to its cost-effectiveness, ease of implementation, low energy consumption, and minimal environmental impact. The prevailing issue of the Zn²⁺ defect vacancy is believed to be improved by adding doping elements into the ZnO ETL layer. To address this, we have explored the addition of various dopants into the ZnO ETL layer. Our findings demonstrate that the incorporation of these dopants not only mitigates Zn²⁺ defect vacancies, but also yields a multitude of benefits, including enhanced electron diffusion, improved stability of PSCs, optimized charge extraction and collection, reduction in losses due to recombination, and, ultimately, a substantial increase in the overall efficiency of PSCs. With careful selection and optimization of dopants, the resulting solar cells have the potential to increase the charge carrier conduction, reduce recombination and improve the energy level alignment with the perovskite layer. Lanthanum is a promising dopant derived from rare earth materials, and has shown significant improvement in the photocatalytic properties of solar cells. The compatibility of La with Zn in the crystal structure, valence state, and electronegativity significantly contributes to its effectiveness as a dopant, thereby establishing it as an indispensable component for improving the photocatalytic properties of solar cells and further developing the field of photovoltaics. Nevertheless, the rare earth material lanthanum is still not widely studied and explored as a dopant in ZnO ETL in PSC. Therefore, this review presents a new paradigm with a doping engineering approach as an alternative means for the fabrication of highly efficient ETL in PSC. While this assumption holds true for most practical scenarios, further investigations can be undertaken to explore and validate this approach.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was supported by Universiti Teknikal Malaysia Melaka and Ministry of Education Malaysia under Grant FRGS/1/2020/TK0/UTEM/03/12.

References

1 International Renewable Energy Agency, Net renewable capacity additions by technology, 2017–2023, <https://www.iea.org/data-and-statistics/charts/net-renewable->

- [capacity-additions-by-technology-2017-2023](#), accessed 12 April 2023.
- N. Shrivastav, J. Madan, R. Pandey and A. E. Shalan, *RSC Adv.*, 2021, **11**, 37366–37374.
 - B. O'Regan and M. Grätzel, *Nature*, 1991, **353**, 737–740.
 - N. G. Park, *Mater. Today*, 2015, **18**, 65–72.
 - A. M. Elseman, W. Sharmoukh, S. Sajid, P. Cui, J. Ji, S. Dou, D. Wei, H. Huang, W. Xi, L. Chu, Y. Li, B. Jiang and M. Li, *Adv. Sci.*, 2018, **5**, 1800568.
 - T. H. Schloemer, J. A. Christians, J. M. Luther and A. Sellinger, *Chem. Sci.*, 2019, **10**, 1904–1935.
 - J. Urieta-Mora, I. García-Benito, A. Molina-Ontoria and N. Martín, *Chem. Soc. Rev.*, 2018, **47**, 8541–8571.
 - W. Zhang, Y. C. Wang, X. Li, C. Song, L. Wan, K. Usman and J. Fang, *Adv. Sci.*, 2018, **5**, 1800159.
 - Y. Wang, Y. Yue, X. Yang and L. Han, *Adv. Energy Mater.*, 2018, **8**, 1–25.
 - O. V. Aliyaselvam, F. Arith, A. N. Mustafa, P. Chelvanathan, M. A. Azam and N. Amin, *J. Mater. Sci.: Mater. Electron.*, 2023, **34**, 1274.
 - N. S. N. M. Alias, F. Arith, A. N. M. Mustafa, M. M. Ismail, S. A. M. Chachuli and A. S. M. Shah, *Appl. Opt.*, 2022, **61**, 4535.
 - M. A. Green, *Solar Cells: Operating Principles, Technology, and System Applications*, Prentice-Hall, 1982.
 - K. Devan and A. Bachchan, in *Physics of Nuclear Reactors*, ed. P. Mohanakrishnan, O. P. Singh and K. Umasankari, Academic Press, 2021, pp. 193–262.
 - S. D. Stranks, G. E. Eperon, G. Grancini, C. Menelaou, M. J. P. Alcocer, T. Leijtens, L. M. Herz, A. Petrozza and H. J. Snaith, *Science*, 2013, **342**, 341–344.
 - G. Xing, N. Mathews, S. Sun, S. S. Lim, Y. M. Lam, M. Grätzel, S. Mhaisalkar and T. C. Sum, *Science*, 2013, **342**, 344–347.
 - Q. Dong, Y. Fang, Y. Shao, P. Mulligan, J. Qiu, L. Cao and J. Huang, *Science Express*, 2015, **347**, 967–970.
 - H. S. Kim, C. R. Lee, J. H. Im, K. B. Lee, T. Moehl, A. Marchioro, S. J. Moon, R. Humphry-Baker, J. H. Yum, J. E. Moser, M. Grätzel and N. G. Park, *Sci. Rep.*, 2012, **2**, 1–7.
 - P. Zhao, J. Xu, H. Wang, L. Wang, W. Kong, W. Ren, L. Bian and A. Chang, *J. Appl. Phys.*, 2014, **116**, 194901.
 - K. C. Wang, P. S. Shen, M. H. Li, S. Chen, M. W. Lin, P. Chen and T. F. Guo, *ACS Appl. Mater. Interfaces*, 2014, **6**, 11851–11858.
 - A. S. Subbiah, A. Halder, S. Ghosh, N. Mahuli, G. Hodes and S. K. Sarkar, *J. Phys. Chem. Lett.*, 2014, **5**, 1748–1753.
 - N. S. Noorasid, F. Arith, A. N. Mustafa, P. Chelvanathan, M. I. Hossain, M. A. Azam and N. Amin, *Appl. Phys. A: Mater. Sci. Process.*, 2023, **129**, 1–16.
 - D. Eli, D. S. Dogo, S. U. Michael, F. O. Uloko and A. O. Salawu, *East European Journal of Physics*, 2021, **4**, 4–17.
 - G. Yang, H. Tao, P. Qin, W. Ke and G. Fang, *J. Mater. Chem. A*, 2016, **4**, 3970–3990.
 - K. Wang, C. Liu, P. Du, J. Zheng and X. Gong, *Energy Environ. Sci.*, 2015, **8**, 1245–1255.



- 25 C. Liu, K. Wang, P. Du, C. Yi, T. Meng and X. Gong, *Adv. Energy Mater.*, 2015, **5**, 1402024.
- 26 Y. Liu, Z. Hong, Q. Chen, W. Chang, H. Zhou, T.-B. Song, E. Young, Y. (Michael) Yang, J. You, G. Li and Y. Yang, *Nano Lett.*, 2015, **15**, 662–668.
- 27 T. Leijtens, B. Lauber, G. E. Eperon, S. D. Stranks and H. J. Snaith, *J. Phys. Chem. Lett.*, 2014, **5**, 1096–1102.
- 28 Q. Jiang, X. Sheng, Y. Li, X. Feng and T. Xu, *Chem. Commun.*, 2014, **50**, 14720–14723.
- 29 J. Qiu, Y. Qiu, K. Yan, M. Zhong, C. Mu, H. Yan and S. Yang, *Nanoscale*, 2013, **5**, 3245–3248.
- 30 S. Dharani, H. K. Mulmudi, N. Yantara, P. T. Thu Trang, N. G. Park, M. Graetzel, S. Mhaisalkar, N. Mathews and P. P. Boix, *Nanoscale*, 2014, **6**, 1675–1679.
- 31 D. Son, J. Im, H. Kim and N. Park, *J. Phys. Chem. Lett.*, 2014, **118**, 16567–16573.
- 32 K. Mahmood, B. S. Swain, A. R. Kirmani and A. Amassian, *J. Mater. Chem. A*, 2015, **3**, 9051–9057.
- 33 K. Manseki, T. Ikeya, A. Tamura, T. Ban, T. Sugiura and T. Yoshida, *RSC Adv.*, 2014, **4**, 9652–9655.
- 34 X. Zhang, Z. Bao, X. Tao, H. Sun, W. Chen and X. Zhou, *RSC Adv.*, 2014, **4**, 64001–64005.
- 35 T. J. Savenije, C. S. Ponseca, L. Kunneman, M. Abdellah, K. Zheng, Y. Tian, Q. Zhu, S. E. Canton, I. G. Scheblykin, T. Pullerits, A. Yartsev and V. Sundstro, *J. Phys. Chem. Lett.*, 2014, **5**, 2189–2194.
- 36 A. You Seok Seo, C. Lee, K. H. Lee and K. B. Yoon, *Angew. Chem.*, 2005, **44**, 910–913.
- 37 A. Hayakawa, O. Yoshikawa, T. Fujieda, K. Uehara, S. Yoshikawa, A. Hayakawa, O. Yoshikawa and T. Fujieda, *Appl. Phys. Lett.*, 2007, **163517**, 1–4.
- 38 P. J. Cameron and L. M. Peter, *J. Phys. Chem. B*, 2003, **107**, 14394–14400.
- 39 S. Ito, P. Liska, P. Comte, P. Pe, U. Bach, L. Schmidt-mende, S. M. Zakeeruddin, A. Kay, M. K. Nazeeruddin and M. Graetzel, *Chem. Commun.*, 2005, 4351–4353.
- 40 S. Lee, I. Cho, J. H. Lee, D. H. Kim, D. W. Kim, J. Y. Kim, H. Shin, J. Lee, O. H. S. Jung, N. Park, K. Kim, J. Ko and K. S. Hong, *Chem. Commun.*, 2010, 1958–1965.
- 41 A. Yella, L. Heiniger, P. Gao and M. K. Nazeeruddin, *Nano Lett.*, 2014, **14**, 2591–2596.
- 42 G. N. Parsons, D. H. Kim, M. Woodroof, K. Lee and G. N. Parsons, *ChemSusChem*, 2013, **6**, 1014–1020.
- 43 L. Etgar, P. Gao, Z. Xue, Q. Peng, A. K. Chandrair and B. Liu, *J. Am. Chem. Soc.*, 2012, **132**, 17396–17399.
- 44 A. Abrusci, S. D. Stranks, P. Docampo, H. Yip, A. K. Jen and H. J. Snaith, *Nano Lett.*, 2013, **13**(7), 3124–3128.
- 45 W. Hu, S. Yang and S. Yang, *Trends Chem.*, 2020, **2**, 148–162.
- 46 A. Hernández-Granados, A. N. Corpus-Mendoza, P. M. Moreno-Romero, C. A. Rodríguez-Castañeda, J. E. Pascoe-Sussoni, O. A. Castelo-González, E. C. Menchaca-Campos, J. Escorcia-García and H. Hu, *Opt. Mater.*, 2019, **88**, 695–703.
- 47 D. G. Lee, M. Kim, B. J. Kim, D. H. Kim, S. M. Lee, M. Choi, S. Lee and H. S. Jung, *Appl. Surf. Sci.*, 2019, **477**, 131–136.
- 48 H. J. Jung, D. Kim, S. Kim, J. Park and V. P. Dravid, *Adv. Energy Mater.*, 2018, **1802769**, 1–10.
- 49 Z. Arshad, A. Hussain, S. Shakir, A. Afzal, M. A. Mujtaba, M. Elahi, M. Soudagar, H. Fayaz, C. Ahamed Saleel and S. Farukh, *Case Studies in Thermal Engineering*, 2021, **26**, 101101.
- 50 E. J. Yeom, S. S. Shin, W. S. Yang, S. J. Lee, W. Yin, D. Kim, J. H. Noh, T. K. Ahn and S. Il Seok, *J. Mater. Chem. A*, 2017, **5**, 79–86.
- 51 Y.-F. Wang, X.-F. Li, D.-J. Li, Y.-W. Sun and X.-X. Zhang, *J. Power Sources*, 2015, **280**, 476–482.
- 52 J. Zhu, F. Liu, M. Lv, S. Chen, L. Hu, J. Tang, J. Yao and S. Dai, *RSC Adv.*, 2015, **5**, 28424–28429.
- 53 W. Ke, G. Fang, Q. Liu, L. Xiong, P. Qin, H. Tao, J. Wang, H. Lei, B. Li, J. Wan, G. Yang and Y. Yan, *J. Am. Chem. Soc.*, 2015, **137**, 6730–6733.
- 54 E. H. Anaraki, A. Kermanpur, L. Steier, K. Domanski, T. Matsui, W. Tress, M. Saliba, A. Abate, M. Grätzel, A. Hagfeldt and J.-P. Correa-Baena, *Energy Environ. Sci.*, 2016, **9**, 3128–3134.
- 55 Q. Jiang, Z. Chu, P. Wang, X. Yang, H. Liu, Y. Wang and Z. Yin, *Adv. Mater.*, 2017, **29**, 1–7.
- 56 Y. Ai, W. Liu, C. Shou, J. Yan, N. Li, Z. Yang and W. Song, *Sol. Energy*, 2019, **194**, 541–547.
- 57 H. V. Quy and C. W. Bark, *ACS Omega*, 2022, **7**, 22256–22262.
- 58 C. Park, U. Kim, C. J. Ju, J. S. Park, Y. M. Kim, K. Char, C. Park, U. Kim, C. J. Ju, J. S. Park and Y. M. Kim, *Appl. Phys. Lett.*, 2014, **105**(20), 203503.
- 59 Y. Yuan, J. Lv, X. Jiang, Z. Li, T. Yu, Z. Zou, J. Ye, Y. Yuan, J. Lv and X. Jiang, *Appl. Phys. Lett.*, 2010, **094107**, 27–30.
- 60 U. Lampe, J. Gerblinger and H. Meixner, *Sens. Actuators, B*, 1995, **25**, 657–660.
- 61 M. D. Aguas, L. Morris and I. P. Parkin, *J. Mater. Sci.*, 2002, **37**, 375–379.
- 62 É. Bévilion, A. Chesnaud, Y. Wang, G. Dezanneau and G. Geneste, *J. Phys.: Condens. Matter*, 2008, **20**, 145217.
- 63 H. Mizoguchi, P. M. Woodward, C.-H. Park and D. A. Keszler, *J. Am. Chem. Soc.*, 2004, **126**, 9796–9800.
- 64 H. J. Kim, U. Kim, T. H. Kim, H. S. Mun, B. G. Jeon, K. T. Hong, W. J. Lee, C. Ju, K. H. Kim and K. Char, *Appl. Phys. Express*, 2012, **5**, 61102.
- 65 D. W. Kim, S. Shin, S. Lee, S. Cho and H. Kim, *ChemSusChem*, 2013, **6**, 449–454.
- 66 S. S. Shin, J. S. Kim, J. H. Suk, K. D. Lee, D. W. Kim, J. H. Park, I. S. Cho, K. S. Hong and J. Y. Kim, *ChemSusChem*, 2013, **7**, 1027–1035.
- 67 L. Zhu, Z. Shao, J. Ye, X. Zhang, X. Pan and S. Dai, *Chem. Commun.*, 2016, **52**, 970–973.
- 68 L. Zhu, J. Ye, X. Zhang, H. Zheng, G. Liu, X. Pan and S. Dai, *J. Mater. Chem. A*, 2017, **5**, 3675–3682.
- 69 S. Bhandari, A. Roy, T. K. Mallick and S. Sundaram, *Chem. Eng. J.*, 2022, **446**, 137378.
- 70 P. Yu, H. Yang, X. Chen, Z. Yi, W. Yao, J. Chen and Y. Yi, *Renewable Energy*, 2020, **158**, 227–235.
- 71 F. Zhao, X. Chen, Z. Yi, F. Qin, Y. Tang and W. Yao, *Sol. Energy*, 2020, **204**, 635–643.
- 72 H. Naja, A. Bahari and S. Ghasemi, *J. Electroanal. Chem.*, 2016, **774**, 14–21.
- 73 H. Naja and A. Bahari, *Opt. Mater.*, 2016, **58**, 210–218.



- 74 S. Yue, H. Pan, Z. Ning, J. Yin, Z. Wang and G. Zhang, *Nanotechnology*, 2011, **22**(11), 115703.
- 75 L. Gaik, J. Shieh, W. Hao, I. M. Hung and M. Hsiung, *J. Alloys Compd.*, 2005, **396**, 251–254.
- 76 K. Wang, Y. Shi, Q. Dong, Y. Li, S. Wang, X. Yu and M. Wu, *J. Phys. Chem. Lett.*, 2015, **6**, 755–759.
- 77 A. Gheno, T. Thi, T. Pham, C. Di Bin, J. Bouclé, B. Ratier and S. Vedraïne, *Sol. Energy Mater. Sol. Cells*, 2017, **161**, 347–354.
- 78 H. Wang, C. Zhao, L. Yin, X. Li, X. Tu, E. Gee and Y. Liu, *Appl. Surf. Sci.*, 2021, **563**, 150298.
- 79 L. N. Lau, N. B. Ibrahim and H. Baqiah, *Appl. Surf. Sci.*, 2015, **345**, 355–359.
- 80 G. Korotcenkov, V. Brinzari, M. Ivanov, A. Cerneavski, J. Rodriguez and A. Cirera, *Thin Solid Films*, 2005, **479**, 38–51.
- 81 P. Qin, G. Fang, N. Sun, X. Fan, Q. Zheng, F. Cheng, J. Wan and X. Zhao, *Thin Solid Films*, 2012, **520**, 3118–3124.
- 82 S. Yoon, S. J. Kim, H. S. Kim, J.-S. Park, I. K. Han, J. W. Jung and M. Park, *Nanoscale*, 2017, **9**, 16305–16312.
- 83 M. B. Kgomo, K. Shingange, M. I. Nemufulwi, H. C. Swart and G. H. Mhlongo, *Mater. Res. Bull.*, 2023, **158**, 112076.
- 84 N. S. N. M. Alias, F. Arith, N. S. Noorasid, H. Sarkawi, A. N. M. Mustafa, M. M. Ismail and M. K. Nor, *TELKOMNIKA*, 2022, **20**, 1117–1124.
- 85 D. C. Look, D. C. Reynolds, J. R. Sizelove, R. L. Jones, C. W. Litton, G. Cantwell and W. C. Harsch, *Solid State Commun.*, 1998, **105**, 399–401.
- 86 J. D. Albrecht, P. P. Ruden, S. Limpijumnong, W. R. L. Lambrecht and K. F. Brennan, *J. Appl. Phys.*, 1999, **86**, 6864–6867.
- 87 J. Wojnarowicz, T. Chudoba and W. Lojkowski, *Nanomaterials*, 2020, **10**(6), 1086.
- 88 D. Liu and T. L. Kelly, *Nat. Photonics*, 2013, **8**, 133–138.
- 89 X. Yang, R. Wang, C. Fan, G. Li, Z. Xiong and G. E. Jabbour, *Org. Electron.*, 2014, **15**, 2387–2394.
- 90 L. Zuo, Z. Gu, T. Ye, W. Fu, G. Wu, H. Li and H. Chen, *J. Am. Chem. Soc.*, 2015, **137**, 2674–2679.
- 91 E. Zheng, Y. Wang, J. Song, X. Wang, W. Tian, G. Chen and T. Miyasaka, *J. Energy Chem.*, 2018, **27**, 1461–1467.
- 92 T. Pauporté, *Chapter 1 - Synthesis of ZnO Nanostructures for Solar Cells—A Focus on Dye-Sensitized and Perovskite Solar Cells*, Elsevier Inc., 2018.
- 93 M. A. Mustafa, F. Arith, N. S. Noorasid, M. S. I. M. Zin, K. S. Leong, F. A. Ali, A. N. M. Mustafa and M. M. Ismail, *Micromachines*, 2022, **13**(12), 2200.
- 94 A. M. C. Ng, X. Y. Chen and A. B. Djuri, *Prog. Quantum Electron.*, 2010, **34**, 191–259.
- 95 K. S. Babu and V. Narayanan, *Chem. Sci. Trans.*, 2013, **2**, 33–36.
- 96 S. Mohan, M. Vellakkat, A. Aravind and U. Reka, *Nano Express*, 2020, **1**, 30028.
- 97 H. O. Chu, Q. Wang, Y. Shi and S. Song, *Trans. Nonferrous Met. Soc. China*, 2020, **30**, 191–199.
- 98 P. M. Aneesh, K. A. Vanaja and M. K. Jayaraj, *Nanophotonic Mater.*, 2015, **6639**, 1–9.
- 99 Y. Zhou, L. Xu, Z. Wu, P. Li and J. He, *Optik*, 2017, **130**, 673–680.
- 100 H. Wang, W. Song, M. Zhang, Q. Zhen, M. Guo, Y. Zhang and X. Du, *J. Chromatogr. A*, 2016, **1468**, 33–41.
- 101 R. Raji and K. G. Gopchandran, *J. Sci.: Adv. Mater. Devices*, 2017, **2**, 51–58.
- 102 S. P. Shrestha, R. Ghimire, J. J. Nakarmi and Y. Kim, *Chem. Soc. Rev.*, 2010, **31**, 112–115.
- 103 N. S. Noorasid, F. Arith, S. N. Alias, A. N. Mustafa, H. Roslan, S. H. Johari, H. R. A. Rahim and M. M. Ismail, in *Intelligent Manufacturing and Mechatronics: Proceedings of SympoSIMM 2020*, Springer, 2021, pp. 895–905.
- 104 N. A. Z. Abidin, F. Arith, A. A. S. Amri, A. N. Mustafa, H. Sarkawi, M. S. M. Suan, A. S. M. Shah, S. A. Junos and F. Salehuddin, *J. Adv. Res. Appl. Sci. Eng. Technol.*, 2023, **31**, 339–351.
- 105 Z. Jalil, *J. Phys.: Conf. Ser.*, 2018, **1116**, 32020.
- 106 M. Aamir, A. Bouich, K. Fradi, J. Mari, L. Khan and B. Mari, *Optik*, 2022, **258**, 168854.
- 107 W. Zhao, H. Li, D. Li, Z. Liu, D. Wang and S. F. Liu, *J. Power Sources*, 2019, **427**, 223–230.
- 108 W. Zhao, K. Wang, H. Li, Z. Yang, Z. Liu and J. Sun, *Sol. Energy Mater. Sol. Cells*, 2018, **178**, 200–207.
- 109 H. Niu, C. Fang, X. Wei, H. Wang, L. Wan, Y. Li, X. Mao, J. Xu and R. Zhou, *Dalton Trans.*, 2021, **50**, 6477–6487.
- 110 R. Pietruszka, B. S. Witkowski, E. Zielony, K. Gwozdz, E. Placzek-popko and M. Godlewski, *Sol. Energy*, 2017, **155**, 1282–1288.
- 111 S. Li, P. Zhang, Y. Wang, H. Sarvari, D. Liu, J. Wu and Y. Yang, *Nano Res.*, 2017, **10**, 1092–1103.
- 112 J. Luo, Y. Wang and Q. Zhang, *Sol. Energy*, 2018, **163**, 289–306.
- 113 D. P. Norton, Y. W. Heo, M. P. Ivill, K. Ip, S. J. Pearton, M. F. Chisholm and T. Steiner, *Mater. Today*, 2004, **7**, 34–40.
- 114 N. S. N. M. Alias, F. Arith, A. N. Mustafa, M. M. Ismail, N. F. Azmi and M. S. Saidon, *J. Eng. Technol. Sci.*, 2022, **54**(4), 220409.
- 115 X. Zhao, H. Shen, Y. Zhang, X. Li, X. Zhao, M. Tai, J. Li, J. Li, X. Li and H. Lin, *ACS Appl. Mater. Interfaces*, 2016, **8**, 7826–7833.
- 116 K. Mahmood and S. Bin Park, *Electron. Mater. Lett.*, 2013, **9**, 161–170.
- 117 K. Mahmood, B. S. Swain and H. S. Jung, *Nanoscale*, 2014, **6**, 9127–9138.
- 118 A. Baltakesmez, M. Biber and S. Tüzemen, *J. Radiat. Res. Appl. Sci.*, 2018, **11**, 124–129.
- 119 J. Kruszyńska, J. Ostapko, V. Ozkaya, B. Surucu, O. Szawcow, K. Nikiforow, M. Hołdyński, M. M. Tavakoli, P. Yadav, M. Kot, G. P. Kołodziej, M. Wlazło, S. Satapathi, S. Akin and D. Prochowicz, *Adv. Mater. Interfaces*, 2022, **2200575**, 1–8.
- 120 P. S. Kolhe, A. B. Shinde, S. G. Kulkarni, N. Maiti, P. M. Koinkar and K. M. Sonawane, *J. Alloys Compd.*, 2018, **748**, 6–11.
- 121 V. K. Sharma, R. A. Yngard and Y. Lin, *Adv. Colloid Interface Sci.*, 2009, **145**, 83–96.
- 122 C. D. Ozlem Altuntas Yildirim and H. E. Unalan, *J. Am. Ceram. Soc.*, 2013, **773**, 766–773.



- 123 G. Bagha, H. Naffakh-moosavy and M. R. Mersagh, *J. Alloys Compd.*, 2021, **870**, 159658.
- 124 G. Bagha, M. Rezaee and H. Naffakh-moosavy, *Ceram. Int.*, 2021, **47**, 16111–16123.
- 125 Y. Zheng, E. Zhao, F. Meng, X. Lai, X. Dong, J. Wu and X. Tao, *J. Mater. Chem. A*, 2017, **5**, 12416–12425.
- 126 J. Zhao, Y. Zheng, X. Lu, J. Chen and X. Tao, *ChemPhysChem*, 2013, 1977–1984.
- 127 Y. Zheng, H. Ding, X. Tao and J. Chen, *Mater. Res. Bull.*, 2014, **55**, 182–189.
- 128 Y. Zheng, J. Zhao, S. Bi, X. Tao and M. Huang, *Electrochim. Acta*, 2015, **157**, 258–265.
- 129 T. A. Chowdhury, M. A. Bin Zafar, M. S.-U. Islam, M. Shahinuzzaman, M. A. Islam and M. U. Khandaker, *RSC Adv.*, 2023, **13**, 1787–1810.
- 130 S. Chen, X. Xiao, H. Gu and J. Huang, *Sci. Adv.*, 2023, **7**, eabe8130.
- 131 X. Wang, Q. Zhao, Z. Li, D. Liu, C. Chen, B. Zhang, X. Sun, X. Du, L. Hao, C. Gao, Y. Li, S. Mao, Z. Shao, X. Wang, G. Cui and S. Pang, *Chem. Eng. J.*, 2023, **451**, 138559.
- 132 L. Bouaziz, K. Si-Ahmed, M. Özacar, M. Trari and Y. Bessekhoud, *Microchem. J.*, 2022, **183**, 108132.
- 133 L. Bouaziz, S. Boulahlib, M. Özacar, K. Si-Ahmed and Y. Bessekhoud, *Mater. Today Commun.*, 2022, **31**, 103450.
- 134 D. Zakiryanov and N. Tkachev, *Comput. Theor. Chem.*, 2019, **1156**, 20–24.
- 135 D. B. Mitzi, *Chem. Rev.*, 2019, **119**, 3033–3035.
- 136 X. Wu, J. Zhang, M. Qin, K. Liu, Z. Lv, Z. Qin, X. Guo, Y. Li, J. Xu, G. Li, H. Yan and X. Lu, *EcoMat*, 2022, **4**, 1–11.
- 137 Y. Lin, B. Chen, F. Zhao, X. Zheng, Y. Deng, Y. Shao, Y. Fang, Y. Bai, C. Wang and J. Huang, *Adv. Mater.*, 2017, **29**, 1700607.
- 138 X. Jiang, F. Wang, Q. Wei, H. Li, Y. Shang, W. Zhou, C. Wang, P. Cheng, Q. Chen, L. Chen and Z. Ning, *Nat. Commun.*, 2020, **11**, 1245.
- 139 N. K. Mogha, S. Gosain and D. T. Masram, *Arabian J. Chem.*, 2020, **13**, 1367–1376.
- 140 F. Subhan, S. Aslam, Z. Yan, A. Ahmad, U. J. Etim, L. Zhen, M. Ikram and M. Yaseen, *Chem. Eng. J.*, 2020, **384**, 123271.
- 141 S. Anandan, A. Vinu, K. L. P. S. Lovely, N. Gokulakrishnan, P. Srinivasu, T. Mori, V. Murugesan, V. Sivamurugan and K. Ariga, *J. Mol. Catal.*, 2007, **266**, 149–157.
- 142 E. Cerrato, C. Gionco, I. Berruti, F. Sordello, P. Calza and M. C. Paganini, *J. Solid State Chem.*, 2018, **264**, 42–47.
- 143 V. H. Thi and B. Lee, *Mater. Res. Bull.*, 2017, **96**, 171–182.
- 144 Z. Chamanzadeh, V. Ansari and M. Zahedifar, *Opt. Mater.*, 2021, **112**, 110735.
- 145 X. Yang, W. Zhao, M. Li, L. Ye, P. Guo, Y. Xu, H. Guo, H. Yu, Q. Ye, H. Wang, D. Harvey, D. Shchukin, M. Feng, T. C. Sum and H. Wang, *Adv. Funct. Mater.*, 2022, **32**, 2112388.
- 146 H. Guo, H. Zhang, J. Yang, W. Gong, H. Chen, H. Wang, X. Liu, F. Hao, X. Niu and Y. Zhao, *ACS Appl. Energy Mater.*, 2020, **3**, 6889–6896.
- 147 K. Shivesh, I. Alam, A. K. Kushwaha, M. Kumar and S. V. Singh, *Int. J. Energy Res.*, 2022, **46**, 6045–6064.
- 148 M. Zafar, B. Kim and D. Kim, *Mater. Chem. Phys.*, 2020, **240**, 122076.
- 149 P. Sakthivel, S. Foo, M. Thambidurai, P. C. Harikesh and N. Mathews, *J. Power Sources*, 2020, **471**, 228443.
- 150 Z. Xu, S. Hwa, L. Gao, Z. Guo, Y. Kamata, S. Hayase and T. Ma, *Org. Electron.*, 2019, **73**, 62–68.
- 151 A. A. Kumar, A. Kumar and J. K. Quamara, *Solid State Commun.*, 2018, **269**, 6–10.
- 152 H. Li, B. Zheng, Y. Xue, S. Liu, C. Gao and X. Liu, *Sol. Energy Mater. Sol. Cells*, 2017, **168**, 85–90.
- 153 J. Sun, S. Yu, Q. Zheng and S. Cheng, *J. Mater. Sci.: Mater. Electron.*, 2017, **28**, 2272–2278.
- 154 S. Roguai and A. Djelloul, *Inorg. Chem. Commun.*, 2023, **157**, 111372.
- 155 S. Jian, Z. Tian, J. Hu, K. Zhang, L. Zhang, G. Duan, W. Yang and S. Jiang, *Adv. Powder Mater.*, 2022, **1**, 100004.
- 156 P. Pascariu, M. Homocianu, C. Cojocaru, P. Samoila, A. Airinei and M. Sucheaa, *Appl. Surf. Sci.*, 2019, **476**, 16–27.
- 157 R. Pandey, G. Khandelwal, I. A. Palani, V. Singh and S.-J. Kim, *Nanoscale*, 2019, **11**, 14032–14041.
- 158 A. Manikandan, E. Manikandan, B. Meenatchi, S. Vadivel and S. K. Jaganathan, *J. Alloys Compd.*, 2017, **723**, 1155–1161.
- 159 P. V Korake, R. S. Dhabbe, A. N. Kadam, Y. B. Gaikwad and K. M. Garadkar, *J. Photochem. Photobiol., B*, 2014, **130**, 11–19.
- 160 L. Yang, T. Xu, Z. Bai and S. Qin, *J. Phys. Chem. C*, 2023, **127**, 7492–7500.
- 161 P. S. Chandrasekhar, A. Dubey and Q. Qiao, *Sol. Energy*, 2020, **197**, 78–83.
- 162 H. Yang, H. Kwon, S. Ma, K. Kim, S. Yun, G. Jang, J. Park, H. Lee, S. Goh and J. Moon, *ACS Appl. Mater. Interfaces*, 2020, **12**, 13824–13835.
- 163 H. Seok, A. Ali, J. H. Seo, H. H. Lee, N. Jung and H. Kim, *Sci. Technol. Adv. Mater.*, 2019, **20**, 389–400.
- 164 F. Zhao, Y. Guo, P. Yang, J. Tao, J. Jiang and J. Chu, *J. Alloys Compd.*, 2023, **930**, 167377.
- 165 Y. Wang, M. Feng, H. Chen, M. Ren, H. Wang, Y. Miao, Y. Chen and Y. Zhao, *Adv. Mater.*, 2023, 2305849.
- 166 V. S. Katta, R. K. Dileep, E. Ramasamy, G. Veerappan and S. S. K. Raavi, *Sol. Energy*, 2023, **262**, 111801.
- 167 S. Chen, C. Ho, Y. Chang, K. Lee and M. Wu, *Chem. Eng. J.*, 2022, **445**, 136761.
- 168 Y. Wang, X. Zhang, Q. Wang, H. Wang, W. Jiang and W. Ding, *Mater. Lett.*, 2022, **327**, 133055.
- 169 Z. Xu, Z. Zhang and X. Zhou, *J. Alloys Compd.*, 2022, **922**, 166186.
- 170 X. Zhou, W. Zhang, X. Wang, P. Lin and S. Zhou, *Appl. Surf. Sci.*, 2022, **584**, 152651.
- 171 S. Park, R. Kang and S. Cho, *Curr. Appl. Phys.*, 2020, **20**, 172–177.
- 172 M. Xing, Y. Wei, D. Wang, Q. Shen and R. Wang, *Curr. Appl. Phys.*, 2021, **21**, 14–19.
- 173 J. Luo, Z. Feng, X. Liu and Z. Li, *IEEE J. Photovoltaics*, 2021, **11**, 374–378.
- 174 J. Y. Kim, S. Biswas, Y. Lee, H. W. Lee, J. M. Jeon and H. Kim, *Adv. Sci.*, 2021, **5**, 2000673.
- 175 F. Bouhjar, L. Derbali and B. Mari, *Nano Res.*, 2020, **13**, 2546–2555.



Review

- 176 Y. Sun, G. Ren, S. Han, X. Zhang, C. Liu, Z. Li, D. Fu and W. Guo, *Sol. Energy*, 2020, **202**, 80–85.
- 177 E. Akman, *J. Mol. Liq.*, 2020, **317**, 114223.
- 178 R. Shwetharani, T. Sushmitha, G. U. Preethi and R. G. Balakrishna, *New J. Chem.*, 2020, **44**, 20441–20448.
- 179 Z. Ge, C. Wang, T. Chen, Z. Chen, T. Wang, L. Guo, G. Qi and J. Liu, *J. Phys. Chem. Solids*, 2021, **150**, 109833.
- 180 J. Li, M. Ren, J. Qing, Y. Wang, Z. Liang, J. Tong, C. Yang and Y. Xia, *Mol. Cryst. Liq. Cryst.*, 2020, **692**, 74–82.
- 181 R. Diana, B. Panunzi, P. Morvillo, E. Bobeico and U. Caruso, *Int. J. Electrochem. Sci.*, 2019, **14**, 10281–10288.
- 182 M. Wang, Y. Sun, J. Guo, Z. Li, C. Liu and W. Guo, *Org. Electron.*, 2019, **74**, 258–264.
- 183 M. Spalla, E. Planes, L. Perrin, M. Matheron, S. Berson and L. Flandin, *ACS Appl. Mater. Interfaces*, 2019, **2**, 7183–7195.
- 184 B. Y. Kadem, M. J. M. Ali and A. F. Abdulameer, *12th International Conference on Developments in eSystems Engineering*, 2019, pp. 741–746.
- 185 X. Yu, Y. Shi, X. Yu, J. Zhang, Y. Ge, L. Chen and H. Pan, *Optoelectronics Letters*, 2016, **12**, 106–109.

