

Cite this: *Nanoscale Adv.*, 2021, 3, 2995

## Biomolecular interactions of ultrasmall metallic nanoparticles and nanoclusters†

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The use of nanoparticles (NPs) in biomedicine has made a gradual transition from proof-of-concept to clinical applications, with several NP types meeting regulatory approval or undergoing clinical trials. A new type of metallic nanostructures called ultrasmall nanoparticles (usNPs) and nanoclusters (NCs), while retaining essential properties of the larger (classical) NPs, have features common to bioactive proteins. This combination expands the potential use of usNPs and NCs to areas of diagnosis and therapy traditionally reserved for small-molecule medicine. Their distinctive physicochemical properties can lead to unique *in vivo* behaviors, including improved renal clearance and tumor distribution. Both the beneficial and potentially deleterious outcomes (cytotoxicity, inflammation) can, in principle, be controlled through a judicious choice of the nanocore shape and size, as well as the chemical ligands attached to the surface. At present, the ability to control the behavior of usNPs is limited, partly because advances are still needed in nanoengineering and chemical synthesis to manufacture and characterize ultrasmall nanostructures and partly because our understanding of their interactions in biological environments is incomplete. This review addresses the second limitation. We review experimental and computational methods currently available to understand molecular mechanisms, with particular attention to usNP–protein complexation, and highlight areas where further progress is needed. We discuss approaches that we find most promising to provide relevant molecular-level insight for designing usNPs with specific behaviors and pave the way to translational applications.

Received 3rd February 2021  
Accepted 16th April 2021

DOI: 10.1039/d1na00086a

rsc.li/nanoscale-advances

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† Electronic supplementary information (ESI) available. See DOI: 10.1039/d1na00086a



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# 1. Introduction

Metallic particles with core size below  $\sim 3$  nm are known as ultrasmall nanoparticles<sup>1,2</sup> (usNPs); those below  $\sim 2$  nm are usually called nanoclusters<sup>3,4</sup> (NCs) and may contain as few as 10–30 atoms, which gives them molecule-like properties. usNPs have garnered interest in recent years because of growing evidence suggesting a safer and more efficient performance in biomedical applications than their larger counterparts. Ultrasmall NPs have distinctive physicochemical properties and unique biological behaviors, including intrinsic catalytic activity,<sup>5</sup> efficient renal clearance,<sup>6,7</sup> exceptional tumor accumulation,<sup>8–11</sup> better cell nucleus penetration,<sup>12,13</sup> and reduced toxicity.<sup>1,14,15</sup> These attractive features, however, depend critically on the NP surface properties, including the ability to avoid exposure of the metallic core to the biological environment.<sup>16–19</sup> These are strong reasons to attract interest in the biomedical community, but their size also makes them unique in precision applications. Indeed, usNPs are protein mimics, comparable in size to an average globular protein. Their interactions with proteins, membranes, and other biological structures, known as nano–bio interactions, can, in principle, be controlled through a judicious choice of design parameters, such as core size and shape, surface chemistry, and core composition. Relative to their classical counterparts, usNPs retain features that make them comparable to drugs,<sup>20,21</sup> with the potential for target selectivity and binding specificity to regulate protein function,<sup>22–25</sup> broadening the possibilities offered by traditional small-molecule medicine. usNPs have been used in traditional and emerging areas of therapy and diagnosis, including biosensing, biolabeling, protein recognition, drug delivery, and cancer theranostics and radiotherapy. Still, significant challenges remain for regulatory approval and routine clinical applications. Some of these challenges are shared with classical NPs, such as physiological fate; others are

more pertinent to the ultrasmall regime, *e.g.*, control of NP–protein binding modes, affinities, and kinetics, all basic concerns in protein engineering as well. Unleashing the full potential of usNPs thus requires detailed knowledge of their interactions with biomolecules and their modulation by the biological microenvironment. Environmental effects include crowding, confinement, and the pH and ionic strength of the solution, all of which change as the particle migrates from the blood to the target site.

This review will focus mainly on ultrasmall gold NPs (usGNPs). However, many of the concepts described here will be valid more generally because the forces underlying usNP–bio interactions are dominated by the core morphology rather than composition and the coating layer properties. Gold as a core material offers several advantages, such as biological inertness, ease of surface modification, and exceptional electro-optical properties, which can be used for imaging,<sup>26</sup> photothermal ablation of tissues,<sup>27</sup> and other therapeutic and diagnostic applications. Progress is being made in synthesizing heterogeneous metallic nanocores,<sup>28</sup> which paves the way for multiple functionalities or anisotropic surface chemistries, adding one more design parameter to control nano–bio interactions.

We will discuss topics of general interest, highlighting areas where further progress is needed to get a clearer picture of the forces operating at the sub-nanometer length scale, especially at usNP/protein interfaces. Classical NPs will be discussed insofar as it helps illustrate contrasting or unique features of usNPs. For example, the vast literature on protein–corona formation, which has dominated the field for the last two decades,<sup>29–32</sup> will be seldom mentioned since the concept of adsorbed proteins on the surface of usNPs is not physically sound. Instead, usNP–protein complexes and hybrid usNP/protein aggregates are the relevant structures that emerge and merit consideration. Because of the fast pace of research in the field, reviews on biomedical applications and studies *in vivo* are also frequent, so these topics will be mentioned mainly to underscore the unique potentialities of usNPs that cannot be realized with classical NPs.

The focus here is on the physics and chemistry of usNP–protein interactions, with emphasis on mechanisms. Although interactions with polyelectrolytes (*e.g.*, RNA, DNA<sup>33</sup>), phospholipids (*e.g.*, membranes<sup>34</sup>), and sugars (*e.g.*, polysaccharides<sup>35</sup>) are not examined explicitly, the interaction mechanisms discussed in this review (*e.g.*, Section 3) are general. Several biophysical, spectroscopic, and computational techniques, some in development, are currently being used, often in combination, to understand and quantify usNP–protein interactions in (near) physiological conditions. This review will focus on approaches we find most promising to provide relevant molecular-level insight to help design usNPs with specific nano–bio behaviors and pave the way for translational applications. After a general introduction in Section 1, we address in Section 2 the physicochemical properties of usGNPs and the dependence of nano–bio interactions on the particle design, focusing on the effects of surface chemistry on usGNP–protein binding and complexation. In Section 3, we discuss the mechanisms that



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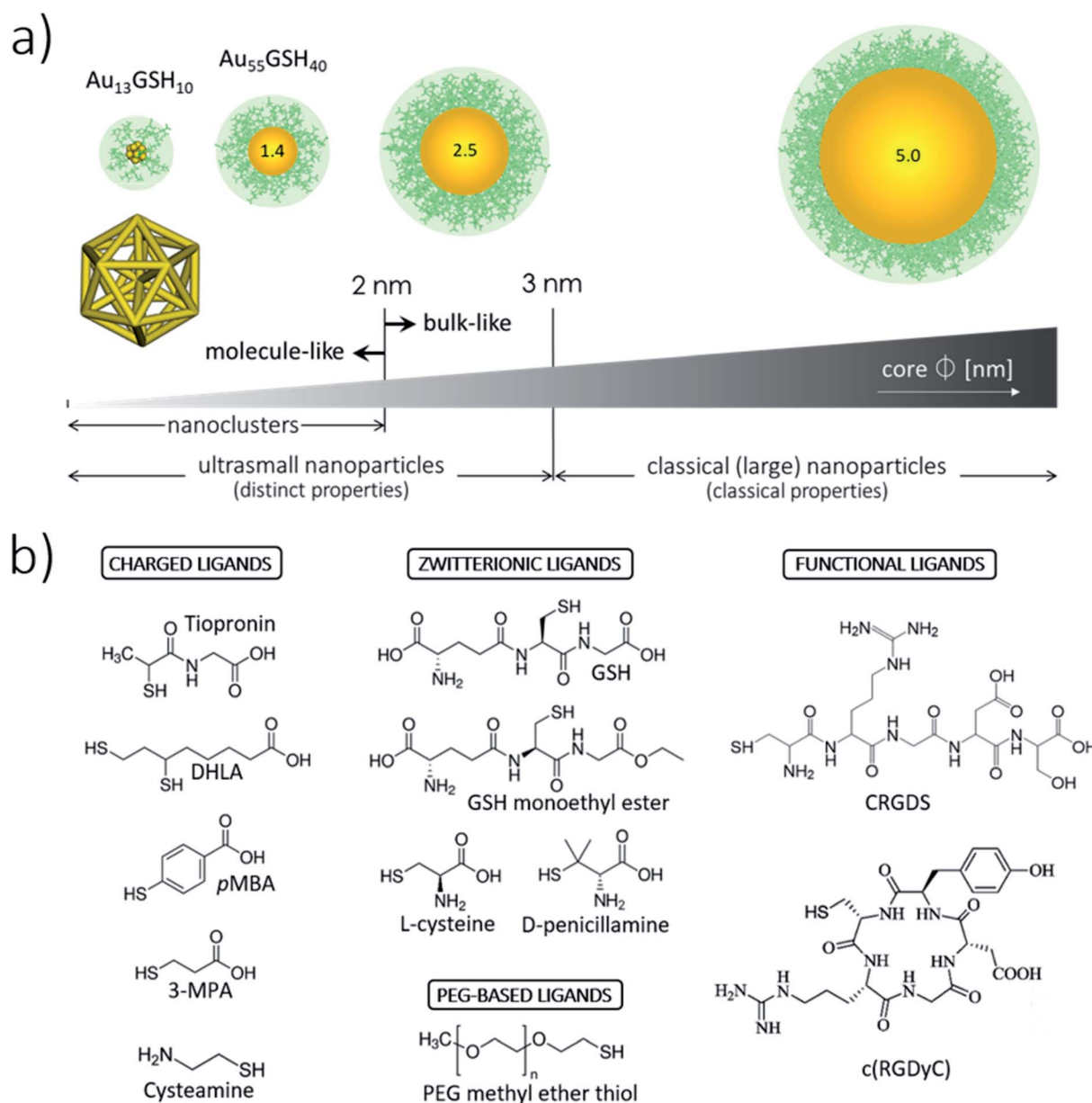
*methods to study biomolecular interactions. He recently joined the computational structural biology unit of the BCBB (NIAID/NIH) to focus on molecular mechanisms of infectious diseases, including viral replication, drug-resistant bacteria, and parasite-borne infections. His interest in ultrasmall nanoparticles stems from their potential use as antiviral and antimicrobial agents.*



govern nano-bio interactions in general. Experimental characterizations of usNP–protein association mechanisms are presented in Section 4. In Section 5, we describe advances in computational modeling and simulations and emerging approaches for reverse engineering usNPs with desired nano-bio interactions. Biomedical applications of usGNPs are presented in Section 6.

## 2. Size and surface chemistry of usGNPs and their influence on nano-bio interactions

There is no strict threshold delimiting the onset of the ultra-small size regime for metallic NPs. Assuming crystal packing, usGNPs, defined as NPs with cores <3 nm across, contain fewer



**Fig. 1** Illustration of the core and surface architectures and properties of usGNPs. (a) usGNPs have core sizes smaller than  $\sim 3$  nm. The central core is protected by an organic monolayer endowing the particles with water solubility, colloidal stability, and functionality. Shown around the cores is a monolayer of GSH molecules represented in green color. GNCs constitute a special class of usGNPs smaller than  $\sim 2$  nm. These can be assigned exact formulas and atomic structures; shown is the atomic structure of  $Au_{13}GSH_{10}$ . (b) Examples of small thiolated ligands used for the surface passivation of usGNPs. (i) Charged ligands are used in the preparation of charged particles that can interact with biomacromolecules. pMBA, *para*-mercaptobenzoic acid; 3-MPA, 3-mercaptopropionic acid; DHLA, dihydroliipoic acid. (ii) Zwitterionic and PEG-based ligands are used in the preparation of nonfouling particles. (GSH is commonly referred to as a zwitterionic ligand due to the simultaneous presence of a negative and positive charge; however, the actual net charge on the molecule is  $-1$  at pH 7). (iii) Functional ligands can endow usGNPs with special properties. Shown are RGD-based peptide ligands that recognize the integrin  $\alpha_v\beta_3$  receptor overexpressed on tumor cells.<sup>8,350</sup>



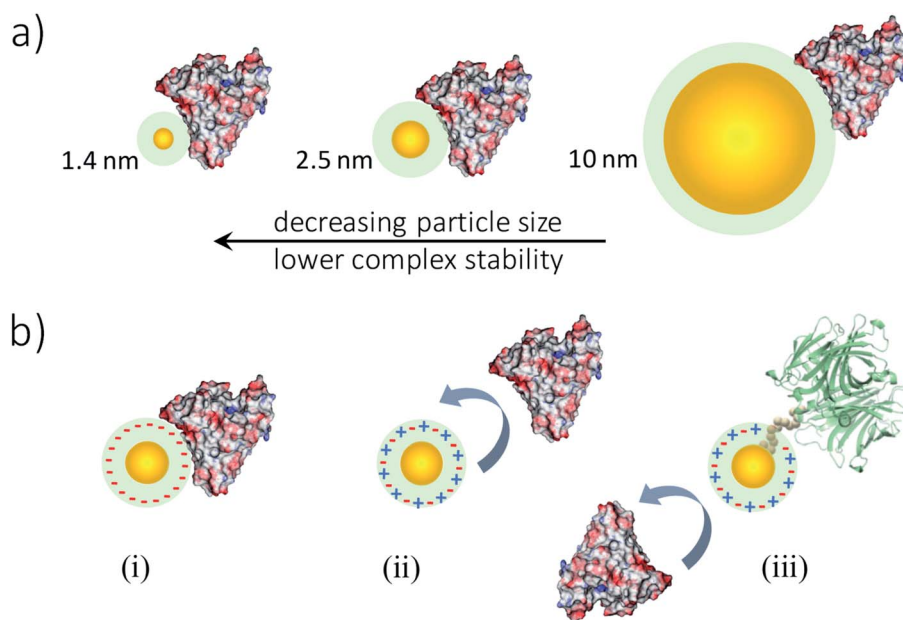
than 1000 atoms (Fig. 1a). Although encompassing a narrow size range within the broader array of nanoscale materials ( $\sim 1$ – $100$  nm),<sup>1</sup> usGNPs can be further subdivided into two distinct size regimes according to their core structures and physicochemical properties (Fig. 1a). On the one hand, usGNPs larger than  $\sim 2$  nm present fcc crystalline structures<sup>36</sup> and localized surface plasmon resonance bands resembling those in bulk crystalline gold. In contrast, gold cores smaller than  $\sim 2$  nm may display non-fcc structure (*e.g.*, icosahedral-based or amorphous-like<sup>37,38</sup>) and molecule-like electronic configurations that generate optical quantization effects, such as sharp absorbance bands and luminescence emission.<sup>39,40</sup> The special class of usGNPs under  $\sim 2$  nm, termed gold nanoclusters (GNCs), can be prepared with atomic precision and be assigned exact formulas and atomic structures.<sup>41,42</sup> Their atomically precise configurations and protein-like structural complexity<sup>43</sup> provide an ideal platform to study GNC interactions with biomolecules and reaction mechanisms. These features provide a venue to control binding specificity more accurately, *e.g.*, by fine-tuning the relative GNC/protein orientation at the binding site. Several reviews have been published on the syntheses and fundamental physical principles of GNCs.<sup>3,4</sup>

The metallic core of usGNPs is protected (passivated) by an organic monolayer endowing the particles with water solubility and colloidal stability.<sup>44</sup> In certain applications, a passivating shell formed by small molecular weight ligands is preferred to preserve the overall ultrasmall size as much as possible. For example, a 2 nm core-sized usGNP covered with the natural tripeptide glutathione (GSH) would maintain a compact hydrodynamic diameter of  $\sim 4$  nm, whereas the same core-sized

NP covered with longer and more complex ligands could reach a hydrodynamic diameter as large as  $\sim 10$  nm.<sup>45</sup> Fig. 1b illustrates a few small thiolated ligands that have been used to prepare water-soluble usGNPs. It is clear from the wide diversity of ligands that the surface chemistry of usGNPs can be tuned almost at will with the desired physicochemical properties and functionalities.

The hydrodynamic diameters of usGNPs are commensurate in scale to a small- or mid-sized globular protein. Protein binding to usGNPs, therefore, takes place through a more restricted contact area interface compared to the adsorption of proteins onto larger NPs (Fig. 2a).<sup>22</sup> Such a limited binding interface yields usGNP–protein complexes of comparatively low thermodynamic stability and short lifetimes,<sup>24,46,47</sup> and also contributes to the preservation of the native conformation and biological activity of bound proteins.<sup>48–50</sup> This is in stark contrast to the hard protein corona – the stable layer of partly denatured and permanently adsorbed proteins – that may develop on the surface of larger NPs.<sup>29,32,51</sup>

Intriguingly, small variations in nanocore size within the ultrasmall regime can influence usGNP–protein interactions substantially. Using a gel shift assay, Boselli *et al.* demonstrated that PEG-coated usGNPs of 2 nm in core diameter were effective in minimizing biomolecular interactions in human plasma, whereas 3 nm-sized particles were not.<sup>46</sup> We also observed a similarly striking effect of small size variations on particle–particle interactions leading to aggregation.<sup>52</sup> We found that GSH-coated usGNPs in biological media supplemented with  $\text{Ca}^{2+}$  ions went from being colloiddally stable to readily aggregating when crossing a size threshold around 2 nm. The forces



**Fig. 2** Influence of size and surface chemistry on usGNP–protein interactions. (a) Influence of size: particles of smaller sizes may form weaker complexes with proteins (albumin is shown as example; drawn to scale). (b) Influence of surface chemistry illustrated through a few selected examples: (i) charged usGNPs (class 1) interact with protein surface regions of opposite charge; (ii) zwitterionic usGNPs (class 2) are resistant against non-specific interactions; (iii) functional usGNPs (class 3) are designed to recognize specific proteins (in green) while avoiding non-specific interactions. See also Table 1.



underlying aggregation of usNPs have been studied by computer simulations and shown to differ from those acting on classical and colloidal particles<sup>53</sup> (cf. Section 5). These examples stress the need for a high degree of control over the uniformity of NPs within the ultrasmall size scale.

The biomolecular interactions of usGNPs of a given size are dictated by the chemical and structural features of the passivating ligands at the surface. Table 1 presents a simple classification system for usGNPs based on their surface passivation and dominant protein interaction modes. Particles in class 1 bind to proteins non-specifically (e.g., Fig. 2b, panel i) and comprehend the largest and most varied subclass owing to their diverse surface chemistries. For example, these particles can be uniformly coated with a wide array of charged ligands<sup>54</sup> (see Fig. 1b for examples), non-functional peptides,<sup>55</sup> carbohydrates,<sup>56</sup> DNA,<sup>57</sup> or combinations of different moieties, e.g., charged/hydrophobic or zwitterionic/hydrophobic,<sup>58,59</sup> to generate more complex surface architectures. Charged usGNPs, in particular, can bind with high affinity to regions of opposite charge on protein surfaces.<sup>23,45,54,60–62</sup> This feature has been exploited in several basic-science studies aimed at the design and use of synthetic usGNPs for protein surface recognition and the modulation of protein structure and function (cf. Section 6.1).<sup>23,45,54,62,63</sup> However, the inherently unspecific nature of long-range electrostatic interactions hinders the broad applicability of charged NPs to more complex environments (e.g., *in vivo*) where binding specificity is required. To overcome this limitation, the Rotello group has explored diverse strategies in monolayer design to influence the regio- and stereo-specificity of the interactions, such as the introduction of hydrophobic moieties alongside charged ligands onto the protecting monolayer or the incorporation of chiral end-groups.<sup>64–66</sup> Of note, large NP scaffolds with tunable surface chemistries have also been under development for protein surface recognition and the control of protein binding orientation.<sup>67–69</sup> However, despite successful examples in this area, realizing specific protein surface recognition with ultrasmall or larger NPs alike remains an enormous challenge. The reason is that specific recognition

requires proper chemical and surface complementarity at the NP/protein binding interface, an accepted notion in high-affinity protein–protein binding and more apparent in usNP–protein associations.<sup>70</sup> Nonetheless, by exploiting the tunability of the particle monolayer, usGNPs become well suited for sensing-type applications that require protein-binding selectivity rather than full specificity (cf. Section 6.2).<sup>66,71–74</sup> To sum up, usGNPs of class 1 offer a vast array of possibilities to control their protein binding modes, e.g., orientation, affinity, and kinetics, and their ability to modulate protein structure and function. While this can be seen as an advantage, it is also a limitation in the sense that small variations in particle size or surface chemistry (e.g., when replicating someone else's work) can translate into disproportionately large differences in biological outcome.

Class 2 comprises an important subclass of NPs for applications *in vivo*. This is due to their effective protein-repelling characteristics (Fig. 2b, panel ii), an attribute that enables them to be cleared efficiently from circulation by kidney filtration.<sup>9</sup> Class-2 NPs are typically passivated with PEG chains or zwitterionic ligands (see Fig. 1b for examples),<sup>75,76</sup> the latter being more effective as a nonfouling coating due to the strong electrostatic association of water that creates repulsive forces against protein binding<sup>77–79</sup> (cf. Section 3). usGNPs coated with GSH have been one of the most widely investigated for therapeutic applications *in vivo*, especially in the area of cancer nanomedicine.<sup>6,80</sup> However, it is important to realize that the GSH molecule has a net charge of  $-1$  at physiological conditions, *i.e.*, it is not a true zwitterion. Taken at face value, this would suggest that proteins might interact with GSH-coated particles mediated by surface patches of positive charge. At any rate, these presumed interactions are likely to be very weak, especially within the smaller nanocluster size regime. True zwitterionic molecules potentially imparting stronger resistance against protein interactions at physiological pH include, among others, betaine-containing ligands<sup>78,81</sup> and GSH monoethyl ester (Fig. 1b).<sup>82,83</sup>

**Table 1** Classification system for usGNPs based on their surface passivation and dominant protein interaction modes

NP classes	Typical surface chemistries	Dominant interaction modes	Example applications
Class 1	Passivating ligands possess charged groups or mixtures of different groups (e.g., charged/hydrophobic)	NPs interact non-specifically (but often selectively) with proteins	Protein recognition for the modulation of protein structure and function; biosensing; antimicrobial
Class 2	Passivating ligands include PEG and zwitterionic ligands	NPs are resistant against protein interactions	<i>In vivo</i> disease diagnosis and treatment, especially where renal clearance is required
Class 3	Functional ligands can be integrated into class-2 NPs. Alternatively, NPs can be uniformly coated with functional ligands (e.g., peptides) imparting both surface passivation and protein-recognition properties	Functional NPs bind to specific proteins while showing resistance against non-specific interactions	Recognition of cell surface receptors <i>in vivo</i> (e.g., for cancer theranostics); bio-imaging probes <i>in vitro</i>
Class 4	NPs are protected by proteins acting both as template and reducing agents	NP interactions are mostly dictated by the protein template	<i>In vivo</i> disease diagnosis and treatment; biosensing



Class 3 comprises NPs that incorporate functional ligands (*e.g.*, peptides) for the specific recognition of proteins (Fig. 2b, panel iii), such as cell surface receptors in cancer cells. usGNPs of class 3 can be prepared, on the one hand, by conjugating the desired functionality directly to class-2 NPs *via* either chemical crosslinking or ligand exchange.<sup>84</sup> Besides this more traditional conjugation approach, usGNPs (GNCs in particular) containing bioactive peptides can be prepared *via* a one-step biomineralization process, wherein a given peptide acts as both reducing agent and template.<sup>85–87</sup> NCs obtained in this way are uniformly coated with the desired peptide, serving the purpose of both surface passivation and molecular recognition.

Class 4 includes the special subclass of protein-protected GNCs, in which proteins are used both as stabilizing and reducing agents during synthesis.<sup>88–90</sup> These nanomaterials are particularly attractive for biomedical applications *in vivo* owing to their expected biocompatibility. Interestingly, fluorescent GNCs can be prepared using immunoglobulin G as a template, which confers the particles with intrinsic recognition properties against specific antigens.<sup>91</sup>

### 3. Molecular mechanisms of usNP–protein interactions

Because of their critical size, usNPs have features common to colloidal particles and proteins. As protein-mimics,<sup>22,92,93</sup> the forces induced by the aqueous solution control most aspects of their nano-bio interactions, including association and dissociation mechanisms (*cf.* Section 4). In general, the interaction energy between a NP, classical or ultrasmall, and a protein (or any other structures) in a biological environment can be conceptually divided into the following contributions<sup>94–112</sup> (*cf.* Fig. 3 and ESI†),

$$V = V_{\text{elec}} + V_{\text{vdW}} + V_{\text{s}} + V_{\text{c}} + V_{\text{d}} + V_{\text{x}} \quad (1)$$

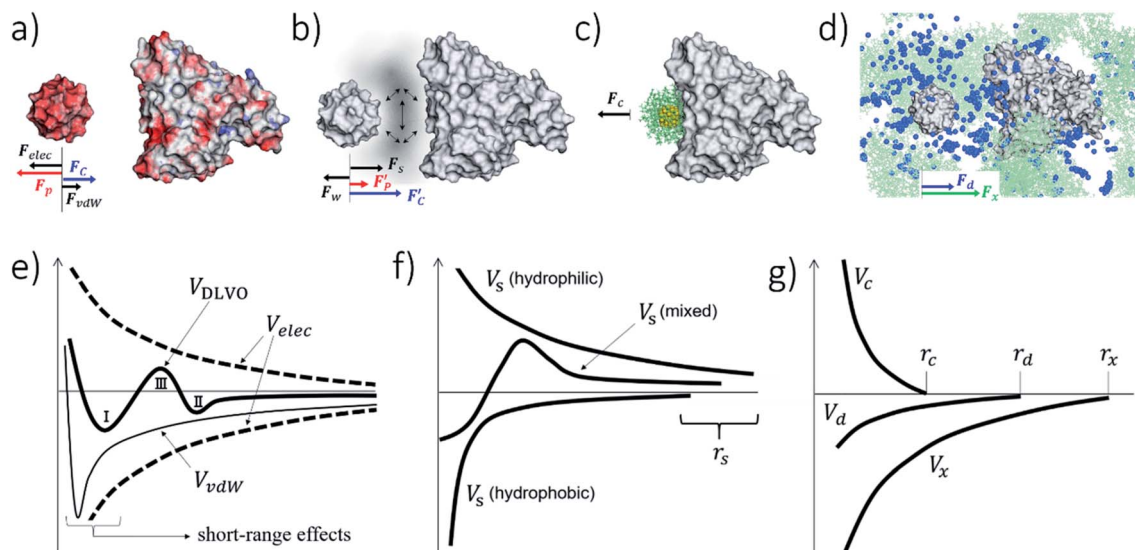
where  $V_{\text{elec}}$  and  $V_{\text{vdW}}$  are the electrostatic and van der Waals terms,  $V_{\text{s}}$  is the contribution from solvent-induced forces,  $V_{\text{c}}$  are the entropic effects of the coating layers, and  $V_{\text{d}}$  and  $V_{\text{x}}$  are the depletion and crowding effects of the solution components. In turn, each of these terms may be subdivided into different physical contributions, sometimes acting in opposite directions, *e.g.*, the counter-ionic pressure and direct Coulombic forces<sup>94</sup> in  $V_{\text{elec}}$  (*cf.* ESI†). Depending on the NP size, shape, and surface physicochemical composition (*e.g.*, types of functional groups, coating density, and length and flexibility of the layer molecules), each term would have a different weight. A central question for understanding NP behavior in the ultrasmall length scale is the right balance between these contributions and how they differ for classical NPs or colloidal systems. Differences are also expected within the ultrasmall regime itself, since the terms in eqn (1) will have different relative weights depending on the type of biomolecular structure interacting with the particle, *e.g.*, another usNP (*cf.* Section 5.1), a protein, or a membrane. Answering these questions requires examining the physical origin of each contribution in greater detail, which we discuss below (*cf.* ESI for quantitative details†).

#### 3.1. Electrostatics and van der Waals forces

In the partition of eqn (1), the  $V_{\text{elec}}$  and  $V_{\text{vdW}}$  terms arise from interactions in a structureless (continuum) water medium in which (point) ions adopt a spatial distribution determined, for example, by the Poisson–Boltzmann (PB) equation.<sup>94,98</sup> These assumptions are common in the treatment of protein–protein interactions and form the basis for the theoretical modeling of colloidal systems. However, even with this simplified view of the aqueous solution, an atomistic representation of the usNP and protein is needed, especially at close contact, because surface topography and specific local interactions at the interfaces (*e.g.*, H-bonds) play a crucial role. These interactions are highly system-dependent and cannot be captured in an analytical form. Thus, in general,  $V_{\text{elec}}$  and  $V_{\text{vdW}}$  have to be described in terms of atom-pair interactions, especially at interfaces (*cf.* Section 5). At longer separations, however, both terms can be described macroscopically using standard approximations in colloidal science, in which case the atomic details at the interface can be ignored. This approach helps make a qualitative connection with classical NPs and colloidal particles, which may be useful for interpreting experimental data qualitatively. In this case, the electrostatic term arises from two separate effects<sup>94</sup> (Fig. 3a and e): on the one hand, the water-screened Coulomb interactions between the usNP and the protein and between each of them and the ions; and, on the other hand, the electrostatics-driven pressure of the ions on the usNP and the protein due to the spatial variations of the ions concentration as the usNP approaches the protein (*cf.* ESI for discussion†). The screened Coulomb interactions between the ions are implicit in their spatial distributions (PB equation); additional interionic forces, *e.g.*, dispersion, would lead to ionic redistribution and a change in this balance of forces. Such ion-correlation effects<sup>94,113</sup> can be significant for multivalent ions, leading to a weakening or even a reversal in the direction of the net interparticle forces. The van der Waals term<sup>114</sup> (Fig. 3a and e) is determined mainly by the density and polarizability of the core atoms relative to water, properties embedded in the Hamaker constant, and by the particles' shape and size,<sup>94</sup> resulting in a relatively long-ranged attractive dispersion decay.

Under these approximations, the first two terms of eqn (1) ( $V_{\text{elec}}$  and  $V_{\text{vdW}}$ ) comprise the well-known DLVO theory; the other terms are, by definition, non-DLVO contributions, although this name is typically given to  $V_{\text{s}}$  (discussed below). The DLVO theory suffices to explain many experimental observations in colloidal systems and is often invoked in the context of classical NPs. The balance between electrostatic (which may be attractive or repulsive) and dispersion (always attractive) may produce a primary or a secondary minimum in the potential, typically associated with coagulation or flocculation of colloidal particles (Fig. 3e).<sup>95–97,99</sup> Computer simulations of usNP–usNP interactions have indeed shown the presence of both minima under certain conditions<sup>52,115,116</sup> and may also exist in some usNP–protein interactions.<sup>93</sup> However, the competition of forces that gives rise to these minima in the ultrasmall regime may be different from that in the DLVO theory, as atomistic simulations have shown,<sup>53</sup> because solvent-mediated effects ( $V_{\text{s}}$  in eqn (1); *cf.*





**Fig. 3** Components of the usNP–protein interaction energy in an aqueous solution as a function of separation (*cf.* eqn (1)). These effects are general and applicable to interactions with membranes and other NPs, although their relative importance is system-dependent (*cf.* ESI†). (a) Electrostatic ( $F_{\text{elec}}$ ) and van der Waals ( $F_{\text{vdW}}$ ) forces can be described within a continuum formalism, leading, *e.g.*, to the DLVO theory (e).  $F_{\text{elec}}$  is composed of a coulombic force ( $F_c$ ) and electrostatic-driven (entropic, often called osmotic pressure) force ( $F_p$ ). The combination of dispersion (attractive component of  $F_{\text{vdW}}$ ; thin solid line in e) and electrostatics (attractive or repulsive; dashed lines) lead, in general, to a primary (I) and a secondary (II) energy minimum, and an energy barrier (III) in the potential (thick solid line; *cf.* ESI†). These features are well characterized in colloidal systems. The continuum approach needs modifications as the size of the particles or the surface-to-surface separation decrease. At close usNP/protein contact, short-range effects (e), such as surface topography and specific interactions (*e.g.*, H-bonds), play a crucial role and need to be described in atomic detail. The major correction to the continuum solvent approximation comes from solvent-induced forces ( $F_s$ ), originating in the liquid restructuring controlled by the water H-bond network (b). Solvent-induced forces can be dominant in the ultrasmall regime and modulate the thermodynamics and kinetics of the association and dissociation processes (*cf.* Sections 4 and 5). In pure water,  $F_s$  is monotonically repulsive for hydrophilic surfaces and attractive for hydrophobic surfaces (f); usNP and protein surfaces have a mixed composition of polar/nonpolar groups that both attract and repel water. In solution, water–water and water–ion interactions (represented by the vertical arrows in b) determine the interfacial liquid structure, affecting the forces that water and ions exert, independently on the usNP and protein (diagonal arrows). Thus,  $F_s$  has a water component ( $F_w$  in b) and corrections to both the coulombic ( $F'_c$ ) and osmotic ( $F'_p$ ) forces. This combination leads to a “desolvation” barrier (f), seen both in small molecules and usNPs (*cf.* ESI†). The onset ( $r_s$ ) of  $F_s$  depends on the solution makeup and the surface characteristics. Solvent-induced effects are difficult to capture in analytical form and must be studied by computer simulations (*cf.* Section 5). Entropic effects are of two types: (i) those generated by the coating layers upon binding (c), which induce a repulsive force ( $F_c$ ) due to the restriction of movement of the layer molecules; and (ii) those generated by the solution components (d), either depletants, which induce an attractive force ( $F_d$ ), or crowders, which typically induce attraction but may induce repulsion ( $F_x$ ; *cf.* Section 3 and ESI†). The onset of these forces ( $r_c$ ,  $r_d$ , and  $r_x$  in g) depends on the layers’ thickness and the shape, size, and concentration of depletant and crowding agents. Panels (a)–(d) show a 1.4 nm GSH-coated usGNP interacting with an albumin protein, drawn to scale; in (d), depletants are the components of a cell-culture solution represented by atom-based vdW spheres (blue); crowders are albumin proteins at physiological concentration shown in green.<sup>55</sup> The relative magnitude and direction of forces depicted in (a)–(d) are for illustration purpose only as they are system-dependent.

Section 3.2) become more prominent for usNPs (*cf.* Section 5). Discrepancies between DLVO theory and experiments at small length-scales are not surprising and have led to corrections and alternative formalisms,<sup>117,118</sup> although still within the continuum solvent approximation.

### 3.2. Solvent-mediated forces

Although  $V_{\text{elec}}$  and  $V_{\text{vdW}}$  incorporate some of the effects mediated by the aqueous solution, the term solvent-mediated (or-induced) forces refer to explicit corrections ( $V_s$ ) to the continuum approximation that stem from the structure and dynamics of the interfacial liquid (Fig. 3b and f), which can be quite different from the bulk. The fundamental importance of these forces has long been recognized,<sup>94,119–123</sup> and their role in NP interactions discussed.<sup>106,107,112</sup> These forces depend critically on the chemical nature of the interfaces and are thus expected to be very different

among usNP classes (*cf.* Table 1). Surface-specific effects become increasingly important at decreasing length-scale and may be dominant in the ultrasmall regime. Aqueous interfaces are the most difficult to characterize owing to the behavior of water itself, dominated by its hydrogen-bond network.

The behavior of water at protein interfaces has long been a topic of basic research. There are still open questions, especially in crowded conditions, *e.g.*, how different is intracellular water from bulk water. The structure and dynamics of water are known to be modified up to several nanometers from certain surfaces, *e.g.*, on colloidal particles or DNA arrays,<sup>94,120,124</sup> but the extent of changes at protein surfaces has been contentious. One reason is the conflicting results from the different techniques used to probe local water behavior in complex media.<sup>125–129</sup> Data have supported a range of views, from non-bulk behavior affecting large portions of intracellular water to only a few hydration shells;<sup>123</sup> the answer likely depends on the local level of



crowding or confinement and solution composition. Answering these questions is critically important, especially for charged usNPs, since their interactions would be very different depending on the properties of the liquid in which they are immersed.

Water structure and dynamics affect the interactions between a usNP and a protein directly (hydration forces, dielectric response) and indirectly through their effects on the solution components, ions in particular. The ionic atmospheres that develop in the vicinity of complex surfaces can depart significantly from idealized distributions used in the derivations of  $V_{elec}$  and  $V_{vdW}$  above. The layers of water and ions that separate a usNP from a protein can be more or less stable depending on their structures, affecting binding affinity and kinetics and, conceivably, the reaction path towards association (*e.g.*, first encounters and secondary binding modes; *cf.* Section 4.1).

There is no simple analytical expression for the contributions of solvent-induced forces that can be used in practice as successfully as the DLVO theory, especially for usNPs. Each system is unique in its local topography and physicochemical properties, leading to specific liquid behaviors. Nonetheless, a few general comments can be made on the case of water-induced forces. These forces originate in the rearrangement of the hydrogen-bond (HB) network as a usNP approaches a protein. The strength with which water holds on their surfaces can make a difference in the thermodynamic and kinetics of usNP–protein binding. These forces are generally strong and repulsive between hydrophilic surfaces and weaker and attractive between hydrophobic ones (Fig. 3f).<sup>94</sup> These trends are seen in a variety of systems and size scales (*cf.* ESI†). However, usNP and protein surfaces are typically composed of a mix of polar and nonpolar moieties, leading to a complex interplay between hydrophilic repulsion and hydrophobic attraction. As a usNP approaches a protein, there is always a dehydration barrier to overcome because both the NP and protein surfaces contain hydrophilic groups that bind to water. The height of the barrier reflects the water's resistance to being removed from the space between the usNP and the protein. Overcoming this barrier requires a disruption of the water–water, water–usNP, and water–protein HB network. Once this resistance is overcome, a swift or mild attraction may occur as the removed water molecules find new favorable interactions, *e.g.*, with bulk water (Fig. 3f; *cf.* ESI for details†).

Ion-induced forces are more elusive and difficult to characterize but behave along similar lines. Their spatial distributions are determined indirectly by the structure and dynamics of interfacial water and directly by the surface themselves. Like water, ions must first be removed for the usNP and protein to come in close contact, and this process ultimately determines the binding affinity and kinetic mechanism.<sup>47,93</sup> Probing these effects in nano-bio interactions require special techniques, as discussed in Sections 4 and 5.

### 3.3. Surface-layer forces

The coating layer may constitute a large portion of a usNP and be a non-negligible component of a classical NP (*e.g.*, about 70% in volume for a 3 nm GSH-coated GNP and 10% for a 30 nm NP

with the same coating). Thus, the effects of the layer's configurational (or conformational) entropy upon binding need to be considered (Fig. 3c and g). Depending on the usNP class (Table 1), these coating-specific contributions can have different weights in the total energy and likely contribute more to usNP–protein binding than the changes in configuration entropy during protein–protein complexation. The thermodynamic penalty associated with the concomitant restriction of movement of the layer molecules can contribute to usNP–protein repulsion (Fig. 3c and g). Such entropic effects are well characterized in colloidal science and play a significant role in the (steric) stabilization against aggregation and deposition, although the molecular origin is different, namely, mixing and elastic compression of grafted or adsorbed polymers (*cf.* ESI†).<sup>95,97,99</sup>

It is unclear how these forces differ among the usNP classes. They are likely more prominent in layers composed of long, flexible chains (*e.g.*, PEG- or some peptide-based) and relatively low coverage density, in such a way that the layer molecules can effectively experience a measurable restriction of movement and elastic adaptation during binding. The effect is probably more limited for dense layers composed of relatively short molecules (*e.g.*, Fig. 1b) with few degrees of freedom. The challenges in estimating the coating layer's entropic contributions are similar to those involved in estimating changes in the configurational entropy during protein–protein binding (*e.g.*, the conformational and vibrational contributions of side chains and flexible loops).<sup>130</sup>

### 3.4. Depletion forces

The entropic effects associated with the solution components also affect the nano-bio interactions. Depletants are inert co-solvent molecules much smaller than the NPs, so that they do not restrict the NPs movement (Fig. 3d).<sup>109,110</sup> Thus, globular proteins or non-adsorbing polymers can act as depletants for classical NPs, but only smaller species would qualify as depletants for usNPs (see below). Depletion forces stem from changes in the volume accessible to the depletants as a NP approaches a protein or another particle.<sup>111,131</sup> The size and concentration of the depletants determine the range and strength of the force. At large usNP–protein separations, the volume excluded by the usNP and protein is maximal; as their separation decreases below a certain threshold (onset of depletion), the volume of the excluded region between them begins to drop, and the configurational entropy of the depletants rises accordingly. The associated thermodynamic reward drives usNP–protein attraction (Fig. 3g). A possible enthalpic component to depletion forces has also been discussed.<sup>132</sup> An equivalent, mechanistic interpretation in terms of osmotic pressure of the depletants can also be given (*cf.* ESI†).

Depletion forces have found use in practice to drive or control various physicochemical processes,<sup>133–135</sup> including phase separation and self-assembly and the properties of soft materials, *e.g.*, nanocomposite and polymer–colloid mixtures. Besides these practical applications, depletion forces can be used to modulate the interaction between NPs *in vitro*, *e.g.*, by





systematically varying the size ratio between NPs and depletants. The role of depletion forces in cell biology has begun to be recognized as a possible mechanism in the assembly of subcellular structures, such as the cytoskeleton and chromosomes.<sup>109</sup>

Depletion forces have been assumed to be important in nano-bio interactions of classical NPs, but their significance in usNPs is less clear. In living organisms, the type or concentrations of the solution constituents cannot be controlled. Besides water itself, only ions, amino acids, short peptides, and other small species are present, with concentrations of  $\sim 250$  mM or less. We recently used simulations to estimate the strength of depletion attraction between usNPs in a cell-culture solution.<sup>53</sup> We found the interaction energy (defined as the minimum of the interparticle potential of mean force at close contact) to be small ( $< kT$ ) for cores with diameters of up to  $\sim 2.5$  nm. However, the interaction energies became larger than  $kT$  for the larger pairs when the solution was crowded with macromolecules at physiological concentrations (Fig. 3d; see Section 3.5). Because depletion forces increase with the number of NPs, they could stabilize small usNP aggregates or usNP/protein complexes in physiological media.

### 3.5. Crowding forces

Unlike depletants, crowders are components of the solution comparable in size to the particles, or larger, and restrict the particles' mobility (Fig. 3d); confinement is an extreme case. Crowding effects have both entropic and enthalpic contributions (see below), although it is the former that is most commonly acknowledged and more intuitive to understand.<sup>100–102</sup>

Like depletion, the entropic effects of crowding stem from changes in volume exclusion. Unlike depletion, however, crowding forces can induce attraction or repulsion, depending on whether the usNP–protein association increases or decreases the system's total entropy.<sup>102</sup> Attraction is more likely for globular crowders and proteins and occurs when more space becomes available to the crowders as the usNP and protein approach one another, increasing the crowders' configurational entropy (Fig. 3g). The size and shape of the interacting species, and the size, shape, and concentration of the crowders, determine the range and strength of the effective force. This dependence affects the binding modes and affinities of the resulting complexes.<sup>136,137</sup>

Crowding forces can be exploited in practice to control the structure of supramolecular materials,<sup>138</sup> including the assembly of nanomaterials, and improve their performance. Their biological role has been established<sup>100,139,140</sup> and are believed to control subcellular processes and organization.

As NPs migrate through the body, they encounter environments that vary broadly in composition and degree of crowding. For example, blood serum contains a total protein concentration of  $60\text{--}80$  mg mL<sup>-1</sup>,<sup>141</sup> whereas the protein content in the intracellular space ranges from  $50$  to  $250$  mg mL<sup>-1</sup>.<sup>142</sup> Addressing the influence of crowding (in concert with depletion) on usNP–protein complexation, both at the experimental and theoretical levels, is therefore vital to draw a more detailed

understanding of the behavior of usNPs in naturally crowded biological milieus (*cf.* Section 4.2).

We recently used simulations to examine the strength of these entropic forces on usNPs association in an albumin solution at concentrations comparable to that in blood serum.<sup>53</sup> Like depletion, the effects were modest, yielding an interaction energy at close contact close to thermal energy for cores of up to  $\sim 2.5$  nm in diameter. However, when the solution was mixed with a cell culture (Fig. 3d), the simulations showed that the energy became larger than  $kT$  for the larger pairs. Crowding and depletion are difficult to separate from their combined effects, but the higher concentration of depletants in the (reduced) available volume of the crowded solution may be responsible for the difference. Because both effects are cumulative, albeit not pair additive, they may stabilize small aggregates and complexes containing just a few usNPs. However, our recent experimental studies *in vitro* show that crowding does have a measurable effect on usNP–protein complexation.<sup>93</sup> It may thus be that the non-entropic effects of crowding (see below) rather than the entropic effects just discussed are more significant for usNP–protein association in near-physiological solutions.

Another effect of crowding is the reduced translational and rotational diffusion of usNPs and proteins, which affect all associated processes.<sup>102</sup> Thus, if a reaction is diffusion-controlled, the reaction rate will decrease; if reaction-controlled, it will increase because the entropic attraction lowers the energy barrier. The result of these opposing effects is difficult to predict, especially if the path to the final complex involves multiple intermediates, as seems to be the case of usNP–protein associations;<sup>47,93</sup> in this case, each reaction and metastable binding mode may be affected separately.

Unlike the entropic effects, enthalpic contributions of crowding and confinement are more difficult to conceptualize. One such effect stems from the fact that crowders and confining structures remove large amounts of water, which is the dielectric medium that modulates the electrostatic interactions. These electrostatic effects of water exclusion are always present and traditionally described with the PB equation, *e.g.*, by assigning a low dielectric permittivity to proteins and a high permittivity to water, or through a position-dependent permittivity,  $\epsilon(\mathbf{r})$ , if a more detailed representation of interfacial water is desired.<sup>143,144</sup> Other enthalpic effects of crowding are closely related to the solvent-induced forces ( $V_s$ ) discussed in Section 3.2, due to the proximity of the interacting molecules to the crowders or confining surfaces; these effects can be quite intriguing and counter-intuitive,<sup>103,123,145</sup> potentially changing the usNP–protein interaction from attraction to repulsion.

Systematic studies of the effects of confinement, *i.e.*, the restriction in the movement of usNPs and proteins by macroscopic surfaces (*e.g.*, inverse micelles, parallel membranes, porous media, or crevices in assemblies, such as the cytoskeleton) might be more challenging than in crowding conditions but could provide valuable information on the nano-bio behaviors of usNPs *in vivo*. The intracellular environment contains many surfaces and crevices, and a large proportion of intracellular water is interfacial.<sup>123</sup> Because usNP interactions



within such regions can differ from those in the bulk, an account of these effects is needed to understand and quantify usNP behavior in the living cell or their interactions with cells. This problem has been recognized in the study of protein function. Failing to reproduce realistic intracellular environments in the test tube may conceal features of the usNP–protein interactions critical for *in vivo* applications.

Studies relevant for *in vivo* applications also require accounting for other effects not discussed here, *e.g.*, hydrodynamic forces. During circulation, NPs are subject to strong shear forces under flow or in contact with macromolecular complexes and cells. These forces can induce degradation or detachment of the surface molecules,<sup>146</sup> or affect the behavior of water, ions, and cosolutes in the local environment, modifying all the solvent-mediated mechanisms of usNPs–protein interactions discussed in this section.

## 4. Experimental characterization of usNP–protein interactions

To gain proper insight into the molecular mechanisms involved in usNP–protein interactions (*cf.* Section 3), a detailed experimental characterization of the thermodynamics and kinetics of the binding reaction is required.<sup>147–149</sup> Such studies may help dissect the different physical contributions to general nano-bio interactions (*cf.* eqn (1)). Examining the different effects driving usNP–protein association faces conceptual and technical challenges similar to those encountered in protein–protein association. Electrostatic effects can be probed by studying the interactions at varying ionic strengths, whereas the contributions of H-bond, hydrophobic, and van der Waals forces on complex stability may be gauged from calorimetric data collected at different temperatures.<sup>150–152</sup> The effects of crowding and depletion can be evaluated by investigating the strength of the interactions in solutions containing varying concentrations of inert species of suitable sizes, *e.g.*, dextran/PEG and glycerol,<sup>101</sup> respectively. Ultimately, a judicious combination of experiments and atomistic simulations is likely the most effective approach to untangle the different physical contributions to the interactions (*cf.* Section 5 for a systematic simulations study of the terms of eqn (1) in the case of usNP–usNP interactions<sup>53</sup>).

Because usGNPs are similar to globular proteins in terms of both size and surface chemistry, it is not surprising that many well-established characterization techniques and methods used to study protein–protein complexation can be adapted to investigate usGNP–protein interactions.<sup>45,61,93,153–158</sup> However, usNPs are not proteins, and one should be aware of the differences and the potential impact of the implementation of such techniques, especially in quantification. We additionally recall that usGNP–protein complexes are short-lived and thus need to be studied *in situ*. Experimental methods frequently used to study protein adsorption onto conventionally NPs, such as centrifugal separation followed by dynamic light scattering (DLS) and  $\zeta$ -potential analysis, perturb the binding equilibrium and are not appropriate.<sup>32,159,160</sup>

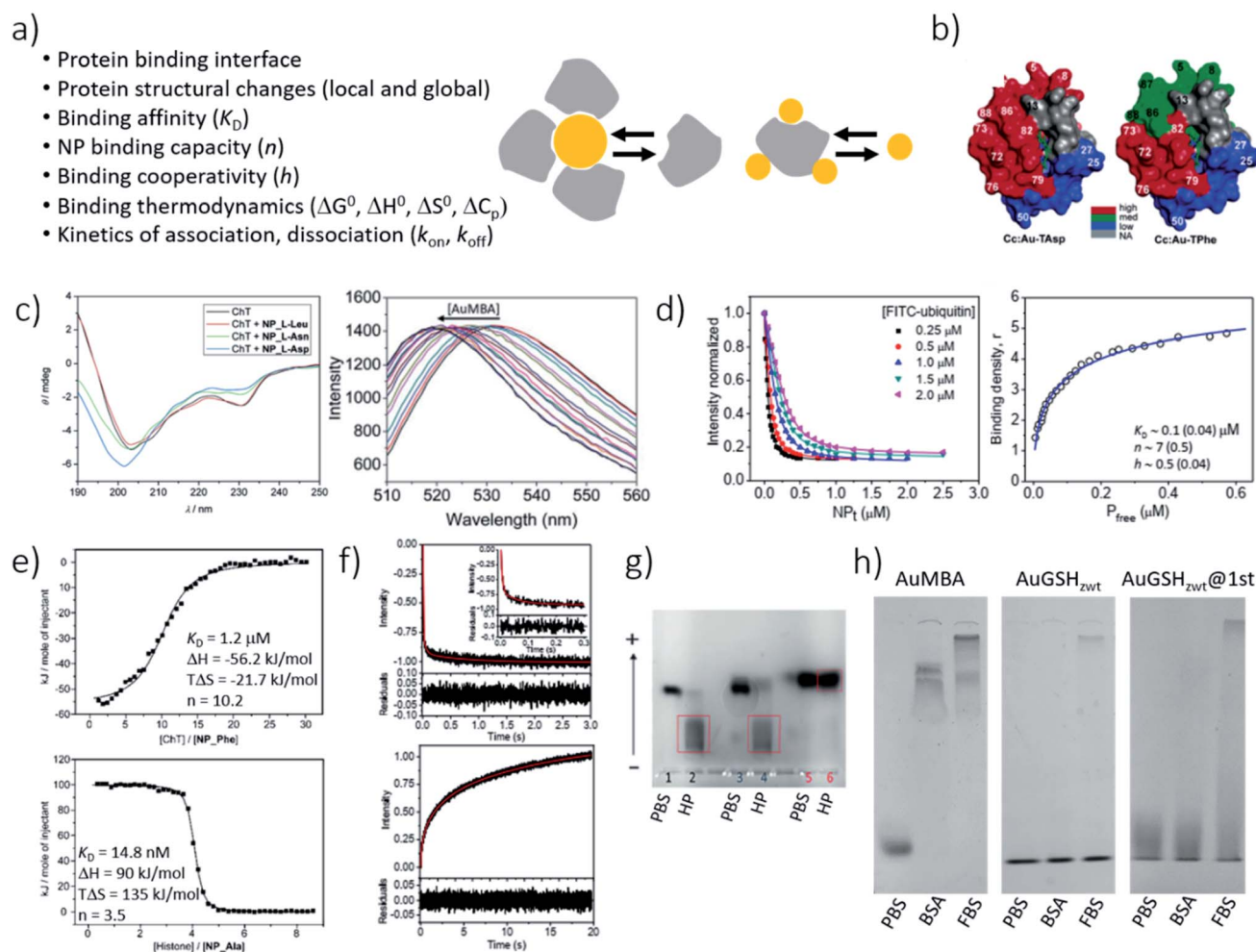
It is important to mention that preceding such detailed investigations, usGNPs must be well characterized with respect to their molar concentrations in solution,<sup>161</sup> physicochemical properties,<sup>2</sup> and colloidal stability.<sup>52</sup> usGNPs uniformly coated with a single ligand type must be characterized at least in relation to  $\zeta$ -potential, average size, and size uniformity, whereas usGNPs engineered with more complex surface architectures (*e.g.*, containing mixtures of different ligands) demand additional characterization.<sup>162,163</sup> At this point, we note that conventional transmission electron microscopy (TEM) and dynamic light scattering (DLS), as commonly employed, do not have adequate sensitivity to characterize the size and uniformity of usGNPs (especially within the NC size regime), whereas annular dark-field scanning TEM and analytical ultracentrifugation (AUC) constitute more powerful choices.<sup>52</sup> usGNPs displaying overly broad size distributions are not suitable for rigorous quantitative interaction studies and require prior size purification/separation.<sup>164,165</sup> The colloidal stability of usGNPs in the appropriate buffer or biological medium must also be carefully evaluated. A number of biophysical techniques (*e.g.*, fluorescence spectroscopy) cannot yield reliable quantitative data unless NPs and proteins remain individually dispersed. Again, AUC is the gold standard for evaluating usGNP stability in solution due to its high hydrodynamic resolution and sensitivity for trace aggregates.<sup>166,167</sup> The reader is referred to many excellent publications covering the issue of NP characterization in greater detail.<sup>2,162,163,168–171</sup>

### 4.1. Biomolecular interactions of usGNPs with isolated proteins in dilute solutions

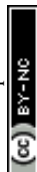
Here, we briefly examine the kinds of structural, thermodynamics and kinetics information that are needed for a detailed description of usGNP–protein complex formation. We focus our discussions on the interactions of class-1 usGNPs with isolated proteins in dilute buffered solutions, where depletion and crowding effects can be ignored. Ideally, a systematic characterization of the binding reaction would entail determination of the following biophysical interaction parameters (Fig. 4a): (i) protein interaction site(s), ideally with knowledge about the amino acid residues involved; (ii) protein conformational changes, both at the global and local levels; (iii) maximum binding capacity of usGNPs; (iv) apparent equilibrium dissociation constant ( $K_D$ ); (v) binding cooperativity; (vi) binding thermodynamics (free-energy, enthalpy, entropy, and heat capacity changes); and (vii) binding kinetics (rate constants for complex association and dissociation). Besides, investigating the dependence of the aforementioned interaction parameters on solution conditions (*e.g.*, temperature, pH, ionic strength) provides additional insight into the molecular mechanism of the binding reaction.

**Protein interaction sites.** usGNPs may interact with proteins selectively at defined protein surface domains so as to achieve the most thermodynamically stable final state. For example, we showed that anionic usGNPs bind to thrombin primarily through a couple of highly positively charged surface domains known as exosites 1 and 2.<sup>23,63</sup> This was verified by fluorescence





**Fig. 4** Experimental characterization of usGNP–protein interactions. (a) Structural, thermodynamics and kinetics information needed for a detailed description of protein binding to usGNPs of class 1 in dilute solutions (see 'b' through 'f' for illustrative examples). Depending on the physicochemical characteristics of both the usGNPs and the proteins, several proteins (gray) can bind to a single NP (yellow), or vice versa. The occurrence of mutual binding leads to aggregation (cf. Fig. 5c). (b) Protein interaction sites. Solvent protection for cytochrome c in complex with Asp- or Phe-decorated usGNPs determined by H/D exchange mass spectrometry.<sup>64</sup> Surface regions depicted in red represent the preferred binding sites for the usGNPs. Facial specificity was observed in the binding of the Phe-decorated particles. (c) Structural changes. Left panel: CD spectra of chymotrypsin (ChT) in the absence or presence of usGNPs decorated with Leu, Asn or Asp.<sup>45</sup> Right panel: normalized FITC emission spectra of active site-labeled thrombin titrated with AuMBA. A clear blue-shift was observed upon NP binding.<sup>23</sup> (d) Determination of  $K_D$ ,  $n$  and  $h$ . Left panel: Set of fluorescence quenching titrations of FITC-labeled ubiquitin with AuMBA. Right panel: thermodynamic binding isotherm calculated from the set of titration curves by a model-independent method of analysis. Fitting of the Hill equation (eqn (4)) yielded the binding parameters annotated in the figure.<sup>47</sup> (e) Binding thermodynamics. Integrated heats of reaction obtained by ITC for the binding of ChT and histone proteins to Phe- and Ala-amino acid terminated usGNPs, respectively.<sup>152</sup> Fitting according to an independent binding model yielded the thermodynamic parameters annotated in the figure. Binding of ChT to Phe-NPs was enthalpically driven with an unfavorable entropic contribution, whereas the opposite was true for the binding of histone to Ala-NPs. (f) Binding kinetics. Stopped-flow characterization of AuMBA-ubiquitin association and dissociation kinetics.<sup>47</sup> Example forward (upper panel) and reverse (lower panel) time traces; red lines represent fits to tri- and four-exponential functions, respectively. The multi-phasic association and dissociation reactions suggested a complex, multi-step binding mechanism (cf. Section 4.1). (g and h) Assessment of interactions in protein-crowded media. (g) Native gel electrophoresis characterization of carboxyl-terminated PEG-coated gold NPs in human plasma (HP).<sup>46</sup> Numbered lanes refer to 5 nm (1 and 2), 3 nm (3 and 4) and 2 nm (5 and 6) core-sized NPs. Only the 2 nm particles exhibited no signs of interactions when dispersed in HP. (h) Native gel electrophoresis characterization of usGNPs in 40 mg mL<sup>-1</sup> albumin (BSA) and undiluted fetal bovine serum (FBS).<sup>83</sup> A significant band shift appeared for AuMBA in both BSA and FBS. No changes in mobility were observed for NPs coated with zwitterionic GSH monoethyl ester (AuGSH<sub>zwt</sub>), in agreement with their strong resistance against protein binding. The smearing of AuGSH<sub>zwt</sub> particles monofunctionalized with Strep-tag (AuGSH<sub>zwt</sub>@1st) was observed in FBS, suggesting the occurrence of weak nonspecific interactions in FBS (but not in BSA). Adapted with permission from ref. 64 ((b); Copyright 2007 American Chemical Society), ref. 45 ((c), left panel; Copyright 2005 American Chemical Society), ref. 23 ((c), right panel; Published by the Royal Society of Chemistry), ref. 47 ((d and f); Published by the Royal Society of Chemistry), ref. 152 ((e); Copyright 2007 American Chemical Society), ref. 46 ((g); Copyright 2017 Wiley-VCH Verlag GmbH & Co. KGaA), and ref. 83 ((h); Copyright 2020 American Chemical Society).



quenching titrations using exosite-directed aptamers in competition experiments with the usGNPs. In an earlier study, Rotello and co-workers used native gel electrophoresis to demonstrate electrostatically selective binding of charged usGNPs to the highly basic cytochrome *c* and acidic cytochrome *c* peroxidase proteins.<sup>54</sup> Using amide hydrogen/deuterium (H/D) exchange mass spectrometry, these authors additionally showed facial specificity in the binding of a Phe-decorated anionic particle to cytochrome *c* (Fig. 4b).<sup>64</sup> Nuclear magnetic resonance (NMR) is another powerful technique with the potential to unveil high resolution, residue-specific information on protein binding domains.<sup>172–178</sup> For example, NMR-based chemical shift perturbation was employed to identify the preferred binding poses of ubiquitin onto usGNPs surface-coated with multimodal anionic ligands.<sup>179</sup> Mass spectrometry-based proteolysis studies have been also carried out to elucidate the preferred binding orientations of proteins adsorbed onto large NPs;<sup>180–184</sup> however, similar studies have yet to be implemented in the characterization of usGNP–protein complexes.

**Protein structural changes.** Proteins can change conformation upon complexation with usGNPs, and understanding these changes plays an essential part in most biointeraction studies. Circular dichroism (CD) spectroscopy in the far-UV spectral region<sup>185</sup> and Fourier transform infrared spectroscopy<sup>186,187</sup> are two techniques widely used to monitor global changes in protein secondary structure. For example, CD spectroscopy was used previously to study the impact of different amino acid-terminated usGNPs on chymotrypsin conformation.<sup>45</sup> It was observed that NPs bearing polar amino acids destabilized the protein in a time-dependent manner, whereas NPs decorated with hydrophobic amino acids had little influence on protein structure (Fig. 4c, left panel). Fluorescence spectroscopy, on the other hand, can detect NP-induced changes in local protein structure occurring around a fluorophore reporter (intrinsic tryptophan residues or extrinsic probes). The basis of the method rests on the high sensitivity of fluorescence emission to local structural perturbations affecting the local dielectric environment.<sup>188,189</sup> For example, using fluorescein as a reporter of environmental changes at the active site of thrombin, we demonstrated the occurrence of allosterically transmitted structural changes to the active site as a result of usGNP interactions to thrombin's exosites (Fig. 4c, right panel).<sup>23,63</sup> Techniques such as NMR spectroscopy and proteolytic cleavage followed by mass spectrometry can be used to reveal NP-induced changes to protein structure at the amino acid residue level; however, these methods have generally been limited to investigating protein adsorption onto large NPs.<sup>183,184,190,191</sup> One should finally consider the possibility that usGNP complexation could modulate the conformational entropy/internal dynamics of proteins without major accompanied structural changes. Techniques suitable for probing dynamics include time-resolved fluorescence anisotropy,<sup>192,193</sup> NMR spectroscopy<sup>194</sup> and hydrogen/deuterium exchange mass spectrometry.<sup>195</sup> However, the characterization of protein dynamics has remained largely unexplored in usGNP–protein interaction studies.

**Binding affinity, stoichiometry, and cooperativity.** In the following discussion, the common situation is assumed in which multiple proteins (defined as the ligands) bind around the surface of a single usGNP (defined as the receptor). We then discuss determination of the microscopic binding affinity ( $K_D$ ) of the interactions, the NP binding capacity or stoichiometry ( $n$ ), and the degree of binding cooperativity ( $h$ ) (reflecting whether protein binding gets progressively enhanced or suppressed with each binding event).

Typically,  $K_D$ ,  $n$  and  $h$  can be experimentally determined by fitting of the Hill equation to a titration curve,<sup>196,197</sup> which is, in turn, obtained by monitoring NP–protein complex formation through changes of some physicochemical signal. For example, steady-state fluorescence spectroscopy,<sup>155,198</sup> capillary electrophoresis,<sup>199,200</sup> isothermal titration calorimetry<sup>152,201</sup> and analytical ultracentrifugation<sup>202</sup> are some of the techniques that can be used to monitor complex formation through changes in fluorescence emission, electrophoretic mobility shift, evolved heat, and sedimentation coefficient, respectively. However, implicit in the application of these (and other) methods are assumptions on the relationship between the observed signal change (*e.g.*, fluorescence, sedimentation coefficient) and the degree of binding.<sup>155,203,204</sup> If underlying assumptions are not met, then fitting of binding models to a single titration curve provides only fitting parameters and not real thermodynamic interaction parameters. To better illustrate these ideas, we discuss below the application of fluorescence quenching titration and analytical ultracentrifugation (AUC) in the quantitative characterization of usGNP–protein interactions.

Characterization of NP–protein interactions by fluorescence quenching titration can be performed in either the direct<sup>61,157</sup> or reverse<sup>31,199</sup> titration modes. The latter applies when the fluorescence signal originates from the proteins (ligands), in which case their concentration is kept fixed in solution while titrating the NPs (receptors). Upon complex formation, the protein fluorescence is quenched by energy transfer to the metallic core of the particles.<sup>205</sup> Fluorescence quench data obtained by reverse titration are often analyzed with the following Hill-type equation expressing the observed fluorescence signal,  $F$ , as a function of the free NP concentration,  $[NP]^h$ :

$$F = F_0 - (F_0 - F_s) \frac{[NP]^h}{[NP]^h + (K_D)^h} \quad (2)$$

where  $F_0$  is the initial fluorescence signal and  $F_s$  is the signal at saturation. However, the use of eqn (2) in the present case is wrong.<sup>155</sup> Eqn (2) would be actually suitable to analyze data obtained in the direct titration scheme in which the NPs were the ligands and not the receptors. In order to correctly apply the Hill equation toward the analysis of data obtained under reverse titration, it is first necessary to calculate a thermodynamic binding isotherm ( $r$  vs.  $[P]$ , that is, fractional receptor saturation vs. concentration of unbound protein) according to eqn (3) below, and only then fit the Hill equation (eqn (4)) to the so-obtained isotherm to find the apparent binding parameters  $K_D$ ,  $n$  and  $h$ .<sup>155,204</sup>



$$r = \frac{F_0 - F}{F_0 - F_s} \frac{[P]_t}{[NP]_t} \quad (3a)$$

$$[P] = \left(1 - \frac{F_0 - F}{F_0 - F_s}\right) [P]_t \quad (3b)$$

$$r = n \frac{[P]^h}{[P]^h + (K_D)^h} \quad (4)$$

Even so, implementation of this procedure would require that the measured fluorescence signal be proportional to the molar concentration of bound ligand. Unfortunately, it is typically not known *a priori* whether this proportionality relation holds in any given system. For example, proteins might assume different binding configurations (in terms of orientation, conformation, or both) as a function of binding density, which in turn could affect their net emission signals in the bound state. Fluorescence quench data can be more rigorously analyzed with the Lohman–Bujalowski's and Schwarz's model-independent method of analysis, which enables the conversion of a set of titration curves into a true thermodynamic binding isotherm ( $r$  vs.  $[P]$ ).<sup>203,204,206–208</sup> Once the isotherm is generated, it can be modeled according to the Hill equation (eqn (4)) to find the relevant interaction parameters (Fig. 4d).

A different situation is encountered in sedimentation velocity AUC, where a change in the sedimentation coefficient of the NP can be measured as a function of protein concentration and the resulting NP occupancy by protein ligands. A first difficulty arises from the fact that colloidal expressions for sedimentation coefficients as a function of total mass traditionally rely on assumptions of spherically shaped particles.<sup>209</sup> This is well-known to be a poor hydrodynamic model for proteins, which are typically highly aspherical and have usually frictional ratios varying anywhere between 1.2 and 2.0. Consequently, sphericity will be a poor assumption for the characterization of usGNP/protein complexes, where the protein contour significantly dominates the shape of the complex. In fact, usGNP with different occupancy will most likely differ from spherical shape to different extent.

As in the case for fluorescence binding signals, the measured change in the sedimentation coefficient ( $\Delta s$ ) as a function of protein occupancy ( $\Delta n_p$ ) is not linear. Owing to the greater density contrast of usGNPs, complexes will usually sediment slower than the bare usGNPs, which is well documented in the literature. Opposite from protein–protein interactions, the added protein ligand will (at first) make a relatively greater contribution to the hydrodynamic friction than to the total buoyant mass of the NP–protein complex. However, with increasing protein occupancy, a point will be reached where added frictional force from further increase in protein occupancy will match the further increase in buoyant mass, and therefore no further change in sedimentation velocity will be observed. Further increase in protein occupancy will then lead to behavior typical for protein–protein interactions with an increase in sedimentation velocity, *i.e.* reversing the sign of the binding signal  $\Delta s/\Delta n_p$ .<sup>83</sup> At the inflection point – in the absence

of binding signal – it is clear that fitting different binding models will lead to poorly defined binding constants and stoichiometries. Whether or not such a “blind spot” is experimentally observed will depend on the mass and density of the usGNP and the mass and shape of the bound protein. For spherical particles it can be shown theoretically to occur at

$$M_{P,\text{bound}} = M_{NP} \times \frac{1}{2} \frac{\rho_P}{\rho_{NP}} \left[ \frac{\rho_{NP} - \rho_S}{\rho_P - \rho_S} - 3 \right] \quad (5)$$

where  $M_{P,\text{bound}}$  is the total protein ligand mass,  $M_{NP}$  is the mass of the unbound NP,  $\rho_{NP}$ ,  $\rho_P$ , and  $\rho_S$  are the density of NP, protein, and solvent, respectively. Assuming densities of typical proteins and bulk Au in water the method becomes insensitive when protein mass is 1.6-fold the NP mass; it will occur at lower protein occupancy for coated usGNPs with lower average density, and at higher protein occupancy for non-spherical protein complexes. These problems may be avoided in future developments using compositional information from exploiting spectral properties in multi-signal analyses, as frequently applied to multi-protein complexes.<sup>210,211</sup>

Finally, an additional difficulty in AUC analysis arises for proteins that do not form stable complexes with usGNPs but are in a dynamic equilibrium of bound and free states with lifetime shorter than 1000 s. In this case, the sedimentation behavior of free and complexed particles is coupled, and the measured sedimentation velocity does not correspond to that of a complex with ‘average occupancy’. Usually this leads to a ‘reaction boundary’ migrating with the time-average velocity of free and bound usGNP states, which adopts a binding isotherm as a function of total protein of a different form, as described by effective particle theory.<sup>211,212</sup> Unfortunately, the case where complexes sediment slower than one of the unoccupied components has not yet been theoretically covered.

The above discussion illustrates how proper quantification of multi-site NP–protein interactions may not be possible based solely on the analysis of a single equilibrium titration curve, or if simplifying assumptions from protein–protein interactions or colloid characterization are uncritically adopted. This limitation could be mitigated to some extent by adopting more sophisticated data acquisition and analysis schemes, as well as pursuing the combination of different techniques. In any event, there is great value to be derived from even a semi-quantitative characterization of the interactions, *e.g.*, in comparative analyses to assess relative differences in binding affinities.

**Binding thermodynamics.** Isothermal titration calorimetry (ITC) enables a detailed characterization of the binding thermodynamics and a determination of driving forces in NP–protein interactions.<sup>152,179,213–216</sup> The basic principle of the technique relies on the *in situ* measurement of the evolved heat changes (either heat absorbed or released) that accompany binding events during a step-wise titration series. The interaction enthalpy ( $\Delta H^0$ ), the apparent  $K_D$  and the maximum binding stoichiometry ( $n$ ) are obtained by fitting of the ITC titration data to an appropriate binding model (typically, to an independent binding site model). The obtained  $\Delta H^0$  reflects the sum of individual energetic contributions from the molecular events taking place during complexation. In particular, non-covalent



bond formation (van der Waals interactions, electrostatics, hydrogen bond formation) at the binding interface is an exothermic process contributing favorably to complex formation, whereas desolvation and ion removal are unfavorable endothermic processes (*cf.* Section 3). For example, a net  $\Delta H^0 < 0$  implies that the enthalpy gained from direct intermolecular interactions is greater than the enthalpic penalty for desolvation and ion removal.  $\Delta G^0$  and  $\Delta S^0$  are obtained from  $K_D$  and  $\Delta H^0$  using the relations  $\Delta G^0 = RT \ln K_D$  and  $\Delta G^0 = \Delta H^0 - T\Delta S^0$ , respectively. Similarly to  $\Delta H^0$ ,  $\Delta S^0$  reflects the sum of contributions from different molecular events. In particular, local conformational restriction and loss of global rotational freedom upon complexation are unfavorable entropic processes, whereas desolvation and ion removal contribute favorably to the overall system's entropy. Previous ITC studies have revealed that the driving force for usGNP–protein complex formation can be either of a primary enthalpic or entropic nature, thus giving insight into the relative thermodynamic contributions of direct intermolecular interactions *vs.* desolvation/ion removal toward binding (Fig. 4e).<sup>62,152,153</sup>

At this point, a word of caution is needed regarding the accurateness of the experimentally determined thermodynamic parameters. First, we note that  $\Delta H^0$  can be usually determined accurately from ITC since it can be directly determined from the data as the difference between the plateaus at the beginning and end of the titrations; as a result,  $\Delta H^0$  is not strongly influenced by the choice of binding model.  $\Delta G^0$  and  $\Delta S^0$ , on the other hand, are determined indirectly from  $K_D$ . However, the value of  $K_D$  may depend on the choice of binding model. For example, use of the commonly used independent binding model may produce incorrect estimates of  $K_D$  in the presence of multi-site binding and cooperative effects.<sup>215</sup> In addition, the analysis of single titration curves in a multi-site reaction model in the NP literature implicitly assumes that the cumulative heat of the reaction is proportional to the degree of binding,<sup>207,208</sup> but which might not be true in a given system.

As an alternative to the use of ITC, the enthalpy change ( $\Delta H^0$ ) of a biomolecular binding reaction can be determined from independent measurements of  $K_D$  (*e.g.*, by fluorescence titration) as a function of temperature and assuming the van't Hoff approximation:  $\ln(1/K_D) = -\Delta H^0/RT + \Delta S^0/R$ . However, application of this approach to NP–protein interactions (or to any other multi-site interaction system for that matter)<sup>151</sup> is problematic given the uncertainties in the determination of  $K_D$  as discussed above.

**Binding kinetics.** While a detailed thermodynamic dissection of usGNP–protein complex formation is undoubtedly important and necessary, equilibrium properties alone convey limited understanding about microscopic mechanisms of interactions. In contrast, a detailed knowledge of the kinetics of binding can afford unique insights into interaction mechanisms.<sup>217,218</sup> For example, given a simple one-step binding model, the association and dissociation reactions can be phenomenologically described by the rate constants  $k_{on}$  and  $k_{off}$ , which, in turn, are connected to the binding affinity by  $K_D = k_{off}/k_{on}$ . Briefly,  $k_{on}$  is related to the size of the energy barrier for

association, whereas  $k_{off}$  probes the stability of the binding interface, hence reporting on complex residence time.

Techniques most suitable for the real-time characterization of interaction kinetics include surface plasmon resonance (SPR) biosensing, quartz crystal microbalance biosensing, both of which rely on formation of complexes with surface-immobilized binding partners, and fluorescence stopped-flow spectroscopy, which probes binding kinetics in solution. However, relatively few studies to date have attempted to use these techniques to understand the kinetics of protein binding to and unbinding from nanomaterials.<sup>24,93,219–225</sup> Next, we briefly describe our recent work on the use of SPR biosensing and fluorescence stopped-flow spectroscopy for the characterization of usGNP–protein binding kinetics. With regards to SPR, it was implemented to investigate the electrostatically-driven interactions of the model proteins CrataBL and thrombin with pMBA-coated and GSH-coated usGNPs (AuMBA and AuGSH, respectively).<sup>24,93</sup> It was determined that protein binding to AuGSH was  $\sim 1000$ -fold weaker in comparison to AuMBA ( $\sim 30 \mu\text{M}$  *vs.*  $30 \text{ nM}$ ). This was found to be mainly due to an  $\sim 100$ -fold smaller apparent  $k_{on}$  of the model proteins toward AuGSH. Noting the larger number of ions and the more complex nature of the ionic atmosphere surrounding AuGSH, the difference in  $k_{on}$  was explained in terms of a larger energy penalty for desolvation and ion removal from AuGSH (*cf.* Section 3). Protein interactions with AuGSH were therefore identified as reaction-limited, whereas interactions with AuMBA were characterized as mainly diffusion-limited. The bound complexes were short-lived, displaying residence times in the range of 2–20 s. Borrowing from the terminology used to describe protein adsorption onto larger NPs, one could therefore refer to the proteins as forming a “soft corona” around AuMBA and AuGSH. However, this should not be mistaken for “soft interactions” (ultraweak interactions), which would actually involve  $K_D$ s in the mM range.

At this point, we note that SPR biosensing suffers from a number of shortcomings such as mass-transport limitations, various surface-related artifacts, and a limited time resolution of the order of 1 second.<sup>226,227</sup> Stopped-flow spectroscopy is a solution-based technique that circumvents most of the limitations of SPR. Of particular interest, the millisecond time resolution afforded by the technique enables identification of discrete intermediate steps along the association pathway in a multi-step reaction mechanism. Stopped-flow spectroscopy was therefore applied to investigate the kinetics of AuMBA binding to fluorescently-labeled ubiquitin, which were used as model systems (Fig. 4f).<sup>47</sup> The results supported a complex, multi-step reaction model; briefly: (i) complexation proceeded through the formation of a weakly-bound and mostly solvated first-encounter complex (apparent  $K_D \sim 8.7 \mu\text{M}$  at  $[\text{NaCl}] = 5 \text{ mM}$ ). This was followed by discrete tightening steps of partial desolvation/ion removal and conformational rearrangement, yielding the final bound state (final apparent  $K_D \sim 0.1 \mu\text{M}$ ); (ii) the bound complex was weakly stabilized, displaying an average lifetime in the range of seconds; (iii) the interactions were characterized by multiple  $k_{off}$  rate constants, which could imply the existence of either multiple intermediate states or multiple



final states (or still, a combination of both); (iv) weakening of the interactions at higher ionic strengths was due (in great part) to a destabilization of the encounter complex, whereas the average lifetime of the bound complex was not affected. Overall, these results supported the general notion that NPs in the ultrasmall size regime can be viewed, in some regards, as protein mimics.

#### 4.2. Biomolecular interactions of usGNPs in protein-crowded media

usGNPs of classes 2 and 3 may be used in applications where a high degree of resistance against protein binding is required. For example, usGNPs used as cancer nanomedicines must resist nonspecific interactions in the presence of a high concentration of proteins and other biomolecules (e.g., the plasma concentration of albumin alone is 500–700  $\mu\text{M}$ ). Here, resistance against nonspecific interactions is defined as an apparent binding affinity in the mM range ( $K_D > 1 \text{ mM}$ ) between a usGNP and a protein. Quantifying such ultraweak interactions is important to fully understand the behavior of usGNPs in the crowded biological milieu. However, there are well-known technical difficulties of characterizing ultraweak interactions in highly concentrated environments.

Native gel electrophoresis is one technique that is commonly employed to assess the behavior of usNPs in protein-rich biological media, such as human plasma (HP) or fetal bovine serum (FBS).<sup>46,83,228,229</sup> A typical experiment consists in running usGNPs diluted in buffer *vs.* HP or FBS side-by-side in a gel, then observing any differences in the electrophoretic mobility of the particles as a result of nonspecific protein binding (Fig. 4g and h). As a general rule, “strong” protein interactions produce obvious band shifts, “weak” interactions cause only the smearing of bands, and “ultraweak” interactions generate minor or even no discernable changes in particle mobility (Fig. 4h).<sup>83</sup> However, there are important limitations to take into consideration in these types of measurements, such as that samples are not strictly maintained *in situ* due to the separation of bands, ‘cage effects’ of polymers on the NP interaction,<sup>230</sup> and the qualitative nature of the results interpretation.

In contrast to gel electrophoresis, migration in sedimentation velocity AUC can be observed without separation of usGNPs from the bath of surrounding serum proteins, allowing dynamic dissociation/rebinding processes to proceed. If as little as 10% protein occupancy of usGNPs is detected, this will allow the detection and quantitation of transient interactions with  $K_D$  in the low mM range.<sup>83</sup> One difficulty for the precise analysis is posed by long-range hydrodynamic interactions altering the sedimentation process of concentrated solutions. Recently, mean-field corrections have been introduced that bring this solution regime in focus for quantitative analysis of weak interactions,<sup>231</sup> which we envision will be applicable to usGNP studies.

Compared to class-2 usGNPs, the characterization of class-3 particles in protein-crowded media presents additional challenges. Class-3 usGNPs should be tested not only for their

ability to resist nonspecific interactions but also for their capacity to recognize the desired protein receptor under crowding conditions. A lack of receptor recognition could signify either a loss of activity of the functional ligand (e.g., due to detachment from the surface or degradation) or the screening of the interactions by nonspecific protein binding. As a test model, we recently prepared usGNPs passivated with GSH monoethyl ester and functionalized with either single or multiple Strep-tag II peptides.<sup>82,83,232</sup> We then verified the ability of the strep-tagged particles to resist nonspecific interactions and to bind Streptactin in the presence of serum. Our results revealed a lower binding efficiency of the multiply functionalized usGNPs over time (apparently due to increased levels of nonspecific serum interactions), suggesting the need for optimizing ligand density onto the surface of usGNPs for improved performance.<sup>83</sup>

Overall, much remains to be understood about the behavior of class-2 and class-3 usGNPs in the highly concentrated biological milieu. For example, it is typically not known what the order of magnitude is of the nonspecific ultraweak interactions; what the lifetimes of the particle–protein complexes are once they form; the extent to which depletion and crowding effects modulate the interactions (*cf.* Section 3); whether surface-attached functional ligands behave similarly as when they are free in solution; the extent to which avidity effects may play a role in the ultrasmall particle size regime; and so forth. Tackling such complex questions will require novel ways to addressing usGNP–protein interactions experimentally.

## 5. Computational modeling and simulations of usNPs

Deeper insight into usNP–bio interactions can be gained with a judicious combination of *in vitro* experiments and computer simulations. Simulations combine two elements, sampling and forcefield, and this section covers the latter. Sampling techniques are chosen depending on the problem at hand.<sup>233–235</sup> Stochastic,<sup>236</sup> nonequilibrium,<sup>237</sup> and enhanced<sup>238</sup> sampling techniques may be used to reduce the statistical errors in a calculation or better explore specific regions of the configurational space, whereas plain equilibrium molecular dynamics (MD) is sometimes the most reliable approach to elucidate mechanisms, as it enables direct observation of the sequence of events, including transition states and the structure and dynamic behavior of water and other solvent components. However, the forcefield (FF) embeds the physics of the system and is ultimately the basis for a correct molecular interpretation of experimental data. Some fundamental themes on the simulations of usNP–protein systems are discussed below, with emphasis on results most relevant to the molecular interpretation of the experiments discussed in previous sections.

### 5.1. All-atom simulations

As protein mimics, simulations of usNP require an atomistic representation. There has been little conceptual innovation in all-atom FFs, mainly *ad hoc* modifications and parametrizations



for specific applications, so the discussion will be limited to recent studies that provide big-picture insight. Of more interest from a methodological perspective are recent developments in multiscale approaches (Section 5.2).

Classical NPs require dealing with two main processes: protein adsorption, which is the primary mechanism of soft and hard corona formation, and the structural changes that usually follow. Different approaches, *e.g.*, QM (quantum mechanical) and hybrid QM/MM (quantum mechanics/molecular mechanics)<sup>239,240</sup> for chemisorption or MM<sup>241–244</sup> for physisorption and protein restructuring, are used to model binding to surfaces of varying compositions, including oxides and mineral, alloys and metals, and, more relevant to biomedicine, polymeric or self-assembled monolayers of organic molecules (reviewed in ref. 245).

By contrast, usNPs require dealing with problems common to protein simulations. NCs may also require optical or electronic characterization of the nanocore, typically by quantum chemical calculations common in the study of small molecules, *e.g.*, density functional theory or *ab initio* QM methods. The computational demands of an all-atom representation limit the study of usNPs to relatively small, diluted systems. As summarized below, a great deal can nonetheless be learned from model systems, which have implications for more general nano-bio behaviors, including usNP–membrane, usNP–polyelectrolyte (*e.g.*, DNA and RNA), and usNP–protein binding.

Accurate calculations of potentials of mean force (PMF) in aqueous solutions have shown<sup>115,116,246</sup> that, under certain conditions, usNPs coated with either anionic or cationic molecules (class 1) can attract each other, whereas usNPs with charges of opposite signs can repel each other. usNPs covered with net-neutral hydrophilic layers (class 2) repel each other due to the solvent-induced forces discussed in Section 3. As the coating molecules increase in length and flexibility (*e.g.*, PEG-based; class 2), usNP pairs tend to show only a stable secondary minimum and a strong repulsion as the layers begin to overlap. In these simulations, the enthalpic contributions to interparticle repulsion have been identified, an interplay of ionic, surface charges, and hydrophobic interactions; it is also possible that the entropic repulsion of the coating layers ( $V_c$  in eqn (1)) play a role as they start to overlap.

The observation that like-charged particles can attract each other is well known in colloidal science, leading to flocculation (secondary minimum) or coagulation (primary minimum), depending on the conditions of the solution, as described by the DLVO and related theories (*cf.* Section 3.1). It is also known that such theories cannot be applied to usNPs because of the dominant effects of solvent-induced forces and the specific interactions at close contact (*cf.* Section 3.2). Also, the only attractive force opposing the direct electrostatic repulsion would be (effective) dispersion, which is relatively small for coated NPs with cores < 10 nm and negligible for usNPs;<sup>53</sup> calculations<sup>115,116,246</sup> have also shown that usNP–usNP attraction begins at a surface-to-surface separation of  $\sim 2$ –4 nm, which is quite long range. The increasing concentration of counterions in the interparticle space undoubtedly plays a role in like-charge attraction, and the stabilizing effects of ions forming bridges

between layer molecules help stabilize the close-contact configurations.<sup>52,115,116</sup> Also, the dependence on NP size, *e.g.*, smaller particles tending to form less stable dimers than larger ones, can be partially explained<sup>52</sup> in terms of the difference in contact surface areas. Despite these partial explanations, the exact interplay of forces underlying aggregation of usNPs in biological fluids remains obscure.

To better understand the origin of these forces and evaluate the relative importance of each term in eqn (1), including the partial contributions within each of these terms (*e.g.*, coulombic forces *vs.* pressure in  $V_{elec}$ ), we conducted MD simulations of two specific pairs of usGNPs.<sup>53</sup> The nature of the simulations (plain, equilibrium, brute-force) allowed us to track the behavior of individual ions and water molecules as the usGNPs associate. Each NP pair was composed of identical particles with core diameters of either 1.4 nm or 2.5 nm, covered with  $\gamma$ -glutathione, immersed in a cell-culture solution at physiological conditions, so the particles were negatively charged. These particles are ideal for the intended study because the dependence of interparticle forces on the core size was observed experimentally as a reversal from repulsion to attraction at a core diameter of  $\sim 2$  nm.<sup>52</sup> The analysis showed the dominant role of solvent-induced forces and revealed interaction mechanisms that contrast sharply with the accepted view of classical NPs and colloids. It was shown that Coulomb interactions alone are insufficient to induce attraction, and the van der Waals forces of the solution components (water and ions) are a major factor to drive aggregation. Although the mobile counterions (diffuse layer) induced attraction, the negative co-ions make a substantial contribution. Transient ion-bridging of layer molecules was observed, but no single ion type was responsible for stabilizing the close contact configurations, as is the case for smaller systems (*e.g.*, amino acids<sup>247</sup>). Direct interparticle dispersion was, as expected, small and almost entirely compensated by water (*i.e.*, the excess polarizability of the NP in the solvent is negligible or, equivalently, the Hamaker constant is nearly zero). Moreover, the contribution of ion dispersion to the interparticle forces was modest, albeit not negligible, and the strongest effects were due to the short-range repulsion from the bound ions (Stern layer, in colloids terminology). The latter creates an imbalance of pressure between each particle's inner and outer hemispheres, which induces interparticle attraction. This imbalance, which vanishes at long interparticle separations (isotropic ion pressure), was not due to any measurable changes in the ion concentrations at the particle surfaces but to a change in either the rate or strength of ion "bombardment" on the surfaces.

The mechanisms just described contrast with those expected from a continuum approximation (*e.g.*, DLVO; *cf.* Section 3), whereby osmotic pressure of ions induces repulsion and Coulombic forces attraction, with the former typically overcoming the latter, leading to "electrostatic" repulsion opposed to (and, in some cases, overcome by) the attractive interparticle dispersion. Moreover, it was shown<sup>53</sup> that just a few ions (1–3) entering or exiting the interparticle space are enough to affect the interparticle forces when the NPs are very close to one another. This sensitivity is expected because the water mobility





in the interparticle space is reduced, lowering the water's local dielectric permittivity. Water was shown to form long-range H-bonded wires connecting and even enveloping the usNPs as they draw near each other. These water chains could shift the  $pK_a$  of ionizable groups on the inner hemispheres, modifying the layers' charge distribution.<sup>248</sup> This charge modulation can affect the interparticle forces<sup>94</sup> and can be studied by constant-pH MD.<sup>249</sup>

All these observations, obtained from atomistic dynamics simulations in simple systems, may help develop a parsimonious theory of usNPs in biological media that could be used in practice as readily and intuitively as the DLVO theory is in colloidal systems. Such a theory would entail finding simple analytical approximations for each of the terms in eqn (1).

## 5.2. Coarse-grained, multiscale, and adaptive multiscale methods

Experimental studies of nano-bio interactions typically involve many particles and proteins, as exemplified by protein-corona formation on classical NPs. Modeling the behavior of usNPs in blood serum or the cytoplasm, or in test-tube setups designed to reproduce such microenvironments, requires dealing with states of high local concentration or crowding. Simulations of these systems need simplifications, typically through a coarse-grained (CG) model.

Several CG models for biomolecular simulations have been proposed since the early developments on united-atom<sup>250</sup> and united-residue<sup>251,252</sup> protein FFs. CG approaches involve some structural simplifications where parts of the system are grouped into beads treated as rigid units. The beads interact with one another through bonded and non-bonded energy terms. Carefully developed parameters aim to reproduce atomistic PMFs, so such CG models are knowledge- or statistical-based empirical models.<sup>245,253–260</sup> They range from highly coarse ones, *e.g.*, proteins treated as spherical particles (used to study protein adsorption and the kinetics of corona formation<sup>261</sup>), to models that retain the overall morphology of the proteins (used to study protein/usNP aggregations<sup>262</sup> or the effects of protein shape on crowding forces<sup>53</sup>), to models in which amino acids are treated as single units interconnected along the backbone chain (to study protein conformational changes upon NP adsorption<sup>263</sup>) or hydrophobic and hydrophilic groups in lipids treated as multi-bead chains (to study passive or active NP transmembrane transport<sup>264</sup>). Further subdivisions into key groups are introduced depending on the application, *e.g.*, to study specific binding of organic polymers or small ligands or even solvent molecules.

All these models have provided molecular explanations to experimental observations, showing that certain phenomena at the nanometer or even subnanometer length scales are insensitive to structural or chemical details.<sup>265</sup> However, it is difficult to know *a priori* which systems or phenomena fall into this category. Ultimately, different parts of a system may need to be treated at different resolutions to capture its essential physics, leading to multiscale modeling.

One characteristic of multiscale simulations is that the CG model(s), once chosen, is kept fixed throughout the simulation.

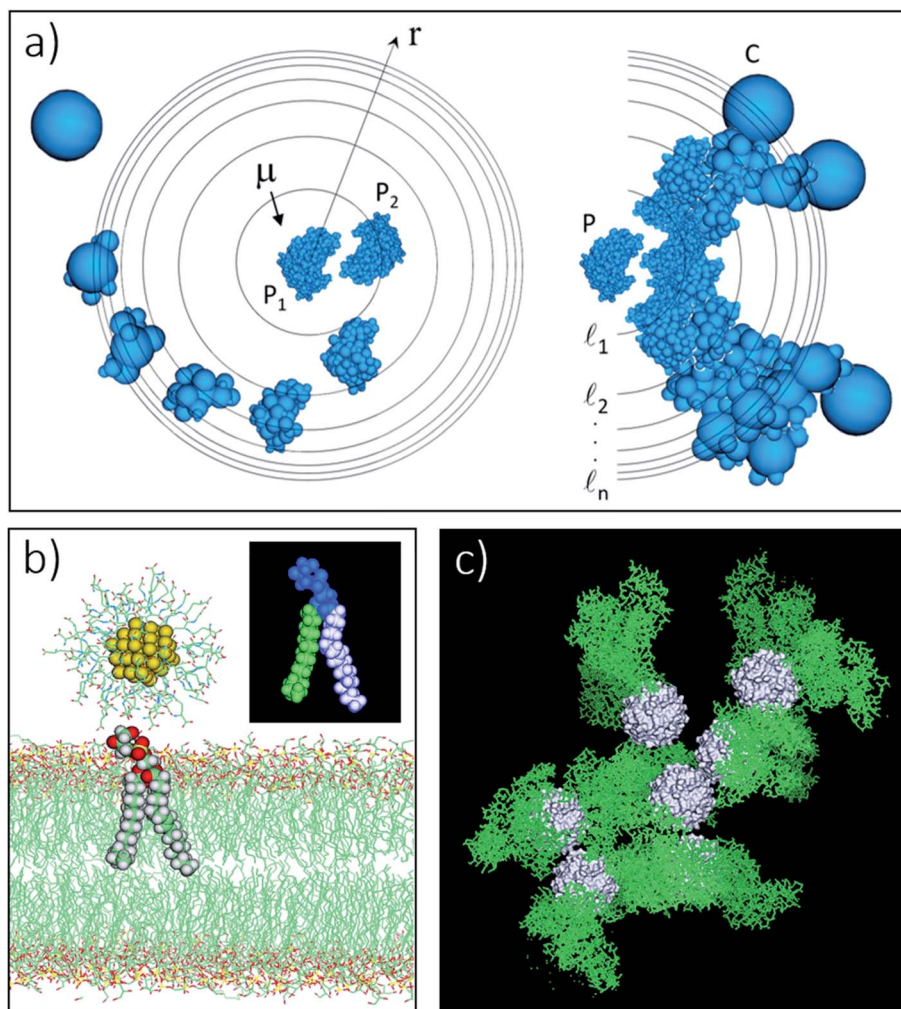
This approach may be inefficient if parts of the system are treated at high resolution when not needed or, worse, physically incomplete if a low resolution is used when an atomic representation is required. The solution to this practical problem is a self-adaptive multiscale algorithm to enable transitioning between molecular resolutions in the course of a simulation, speeding up computation while preserving the short- and long-range structural requirements for a complete physical description.<sup>266,267</sup> Of particular interest are self-adaptive methods that provide a smooth, reversible transition between the atomic, mesoscopic, and macroscopic scales, as discussed below.

The need for an adaptive multiscale approach is better exemplified by charged usNPs (class 1; *cf.* Table 1). Electrostatic forces can still operate far from their surfaces, especially in anisotropic or sub-optimal hydration conditions, *e.g.*, during complexation or aggregation or in proximity to membranes. Indeed, nearby structures can exclude large amounts of water, which affects the dielectric screening of the interactions.<sup>143</sup> These long-range effects are mostly insensitive to atomic details and can thus be modeled at low resolution. However, the forces that determine the nano-bio interactions depend more critically on the details at interfaces. Modest changes in surface chemistry can have profound effects on usNP-protein interactions.<sup>47,63,93</sup> As discussed in Sections 3 and 4, particles of nearly the same sizes and charges can interact with the same protein differently depending on the ions and water behaviors at the usNP/protein interface during binding. In a system composed of hundreds or even thousands of proteins and usNPs, both low and high resolutions need to be accounted for within the same simulation according to demands dictated solely by the system configuration at each point in the simulation.

We have developed a multiscale method based on a self-adaptive coarsening algorithm that enables on-the-fly adjustment of solvent effects to changing resolutions in the system (Fig. 5a, left panel).<sup>34,267</sup> The method prioritizes structural details at the molecular surfaces (*e.g.*, a NP, protein, or lipid) while treating the molecular cores at lower resolutions. Thus, the key elements that determine the nano-bio interactions are maintained, whereas a significant computational speedup can be achieved in crowded systems containing many molecules and NPs, as expected in real biological media. The method preserves physically relevant features of the molecules across scales,<sup>262</sup> including hydrodynamic radii, charge distributions, and moments of inertia, and a connection can thus be made between the microscopic (atomic) and macroscopic (thermodynamic, hydrodynamic) representations.<sup>149,268</sup>

The computational efficiency can be optimized by introducing intramolecular gradients of resolution (Fig. 5a, right panel). This is done by partitioning the molecule in non-overlapping substructures and treating each as an independent molecule subject to multiscale. This partition can be entirely arbitrary or follow a structure-based criterion, *e.g.*, by treating the nanocore and the layer of a usNP or the head group and hydrocarbon chains in a lipid molecule separately (Fig. 5b). Such an approach allows the simulations of highly crowded systems, *e.g.*, hybrid usNP/protein aggregates (Fig. 5c) or states of local condensation. The study of aggregates,<sup>262,269</sup> in





**Fig. 5** Adaptive multiscale method for the simulations of nano-bio interactions. (a) Left: two interacting molecules see each other at different resolutions depending on the surface-to-surface separation ( $r$ ), with the highest resolution, typically atomistic, at close separation (in practice, when the minimally containing spheres begin to overlap; region  $\mu$ ). The example illustrates how protein  $P_1$  sees  $P_2$  as their relative configuration changes and crosses different thresholds  $\lambda_i$ ; the situation is symmetrical. In a multiprotein system, a protein sees each of the other proteins at a different resolution depending on their relative spatial configurations and is seen by each of the other proteins at a different resolution. Thus, not only the components of the system are represented at different scales (the standard definition of “multiscale”), but each element of the system is represented at multiple scales, each of which changes as the system evolves in time or configuration (adaptive multiscale). The algorithm captures desirable properties, including a continuous, reversible, and self-guided adaptation of resolutions, all of which are incorporated through the forcefield as a continuous function of the coordinates. Thus, energy is conserved, the forces are continuous, and detailed balance is preserved throughout the simulation. Right: the algorithm can be adapted to accommodate intramolecular gradients of resolutions for a more efficient treatment of large systems; this can be accomplished in practice by dividing  $C$  into parts, either arbitrarily or with a pre-established criterion. (b) A GSH-coated usNP interacting with a POPG lipid bilayer; each lipid molecule may be considered a single molecule subjected to multiscale<sup>34</sup> or as a composite of, e.g., the head group and the hydrocarbon chains (inset). (c) A hybrid usNP/albumin cluster.<sup>262</sup> The interactions between the highly charged usNPs and the proteins require both long-range interactions and atomic details at all the interfaces to be accounted for. An adaptive multiscale approach provides a computationally efficient, physically robust description of such crowded systems. Internal dynamics could be introduced by switching on and off degrees of freedoms as the atoms enter or exit the region  $\mu$ ; this must be done through a thermostat to keep the energy equally partitioned among the active degrees of freedom.<sup>266</sup>

particular, may include their morphologies and kinetics of assembly and dissolution and thus require a statistical analysis of an ensemble of aggregates. Each aggregate may contain many usNPs and proteins, so the computational challenge is great. This bottleneck can be overcome while preserving the system's physical description (short- and long-range interactions) by adaptive multiscale with intramolecular variations in resolution.

The method just described permits the analytical calculation of forces during dynamics and satisfies detailed balance in MC simulations. The algorithm treats molecules as rigid bodies, limiting its current use to the early stages of complexation, *i.e.*, before interfacial relaxation events. However, this is still a critical step in the hierarchy of relaxation times because it helps identify preferential first-encounter modes<sup>34</sup> from which subsequent dynamic processes can follow (*e.g.*,



endocytosis<sup>264,270</sup> or protein adaptation upon NP adsorption). Simulating such events would require switching to a suitable CG model and a new dynamic simulation started. This discontinuity could be avoided if internal degrees of freedom can be switched on and off on-the-fly within the same simulation, enabling internal dynamics of molecules in the context of multiscaling. Such ideas have been proposed and applied to liquids.<sup>266</sup> In this method, the system is divided into regions, with molecules in each region treated at different resolutions. When a molecule moves from one region to another, it must adapt to the new scale, and this is accomplished in practice through an interpolating algorithm. The method accounts for the number of degrees of freedom associated with the different resolutions as the particles hop between regions, so MD simulations can be performed at constant temperature using a thermostat. The conceptual connection between both methods lies in how the system morphs between scales: in one case, through the interpolation of forces,<sup>266</sup> in the other, through interpolation of the potentials.<sup>267</sup>

### 5.3. Machine learning: prediction of nano-bio behavior and inverse NP design

The *in vitro* and *in silico* approaches discussed so far in this review can be combined with a data-driven method in a three-pronged strategy. This third approach falls under the general name of Machine Learning (ML).<sup>271,272</sup> It uses available data, learns from them, models a pattern, and produces new data not immediately apparent from the input set. The new information can then be used to design *in vitro* experiments, novel materials with desirable properties, or improve the FFs for more reliable simulations.

Despite their long history, ML approaches applied to the nanoscale have only recently begun to be exploited and still face many challenges. ML methods encompass a vast array of techniques. A recent survey<sup>271</sup> has identified over forty approaches that can be used in nanoscale engineering and inverse design. Among these are traditional, well-established methods, such as artificial neural networks (ANNs), including convolution, recurrent, and deep nets; regression methods; support vector machines; and others still in development.

ML methods used for signal processing (*e.g.*, vision or speech recognition, medical images) benefit from vast amounts of available data, even artificially generated, for training and testing. By contrast, systematic data acquisition in physical or chemical systems is difficult, especially in the nanoscale, for which fine-tune manufacturing and characterization are technically challenging.

The scarcity of training data can be overcome with CG simulations. Multiscaling, in particular, can be used to explore features of the usNP design that are difficult to probe with available technology (*e.g.*, high aspect ratios or coatings with specific densities or compositions). The process was illustrated in a recent systematic study of the interactions of usNPs with phospholipid bilayers.<sup>34</sup> The efficiency of the multiscaling method discussed in Section 5.2 allowed us to investigate the dependence of binding modes, energies, and early-stage

membrane penetration on usNP shape, size, and layer composition. The data thus generated were used to train an ANN for reverse engineering usNPs with preferential nano-membrane behaviors.<sup>70</sup> The use of ANNs for pattern recognition is well established,<sup>273</sup> versatile enough to study a variety of physicochemical problems, and can incorporate increasing levels of complexity into the NP design through network arrays (see below). ANNs have recently been used to design multi-layer nanophotonic devices,<sup>274</sup> graphene-based metamaterials,<sup>275</sup> and metasurfaces<sup>276</sup> with desired properties. Different types of ANN architectures have been used depending on the application. We used a single, shallow, multi-output net, the connectivity of which was suggested by the known physics of the system. Incorporating as much physicochemical information as possible in the net design may improve training, reduce redundancy, and enhance efficiency and predictive power.

Safe and efficient use of synthetic usNPs in diagnostics and therapeutics requires multiple independent conditions to be satisfied. For example, aggregation of NPs must be limited in different biological environments; the structural integrity of the coating layer must be preserved during circulation; NPs must have low affinity for plasma proteins and escape detection by the immune system; NPs should accumulate efficiently in target cells, be selective for proteins or membranes, and bind with specificity to target sites to control function or morphology. The ability to integrate data from each of these areas, possibly coming from different laboratories and techniques, may help design usNP for *in vivo* applications. This data integration could be accomplished through interconnected arrays of ANNs, each optimized independently for a particular property. The array can then be used to find design parameters that best satisfy the many independent behaviors imposed on the NPs. The inherent complexity of the biological environment is thus reduced to several smaller, independent problems, each simpler to handle with its own specially designed ANNs. The concept was illustrated by the design of several usNPs displaying preferential interactions with membranes.<sup>70</sup> This kind of approach can be used, for example, in the design of usNPs as antimicrobial agents (*cf.* Section 6.3), as the mechanisms of action and the structural requirement for the particle to reach and disrupt the inner bacterial membrane is multifaceted.<sup>277,278</sup> Similar ML-based ideas are currently being used with encouraging results in the design of antimicrobial peptides.<sup>279,280</sup>

## 6. Biomedical applications of metallic usNPs and NCs

Classical NPs have advanced through decades of research and development, and are now used successfully in several biomedical applications, with over a hundred types having already met regulatory approval or currently undergoing clinical trials.<sup>281</sup> usGNPs are a relatively new concept, but have found remarkable use in biological and biomedical applications, including but not limited to protein recognition, biosensing, biolabeling, cellular imaging, drug delivery, and disease diagnosis and therapy. Next, we briefly highlight emerging



biomedical applications of usGNPs, limiting the discussion to those cases in which non-covalent usGNP–protein interactions are involved in some respect. Due to space limitations, we only provide illustrative examples of applications rather than an exhaustive review of published papers. The reader is referred to additional reports where the myriad biomedical applications of usGNPs are reviewed more broadly.<sup>1,14,87,282–285</sup>

### 6.1. Modulation of protein function

usGNPs can be generally viewed as protein mimics. Indeed, both experiments and simulations have shown that usGNP–protein interactions may retain features of a typical biomolecular complexation event, such as a “well-defined” binding site,<sup>22,23,64,286</sup> binding reversibility,<sup>24</sup> reversibility of conformational changes upon dissociation,<sup>63</sup> and similar thermodynamic<sup>152</sup> and kinetic signatures.<sup>47</sup> usGNPs thus constitute attractive model systems to mimic the structure and action of biomolecules. For example, the Rotello group has systematically explored usGNPs as artificial scaffolds for protein surface recognition and the modulation of protein function, looking specifically into the effects of the NP surface chemistry (including charge and hydrophobicity) on the structure, stability and catalytic properties of enzymes.<sup>45,287</sup> One interesting observation was the much higher rates of hydrolysis of cationic substrates over anionic ones when chymotrypsin was complexed to anionic usGNPs.<sup>287</sup> To explain the results, it was proposed that the NP monolayer regulated both the capture of substrate by the active site and the release of product *via* electrostatic interactions. In another study, Jiang and co-workers looked at the effects of the NP surface chemistry on the structure and enzymatic activity of chymotrypsin.<sup>25</sup> They reported that DHLA-coated GNCs caused a significant change in the secondary structure of chymotrypsin and a corresponding ~80% decrease in the activity, whereas GSH-coated GNCs produced only a slight change in structure and no significant inhibition. This difference in behavior was attributed to the occurrence of hydrophobic interactions between the chymotrypsin surface and the carbon chains of DHLA.

Recently, we demonstrated that usGNPs could perform as allosteric modifiers of enzyme activity (Fig. 6a).<sup>23,63</sup> Allosterism is a powerful concept: it entails changes in substrate binding affinity or catalytic activity at the active site in response to effector binding to allosteric (distal) sites.<sup>288,289</sup> Allosteric regulation is unique in that it can access a wide range of functional responses, including enzyme activation, partial or full inhibition. Using  $\alpha$ -thrombin as a model enzyme, we showed that anionic AuMBA and AuECYN (passivated with the peptidic Ac-ECYN biomimetic coat) bound to the cationic exosites 1 and 2 on the surface of thrombin, with AuECYN demonstrating a remarkable binding selectivity to exosite 2. NP complexation allosterically transmitted long-range structural changes to the active site, leading to a partial, submaximal inhibition of the activity. It was further established that long-term (24 h) interactions with AuMBA disrupted the optimum active-site geometry and irreversibly inactivated the enzyme, whereas AuECYN sustained a constant level of inhibition over time. In sum, these

results revealed how the unique binding modes of different NP types could differentially modulate the structure and function of an enzyme's active site, while all at the same time preserving the overall folded state of the protein.

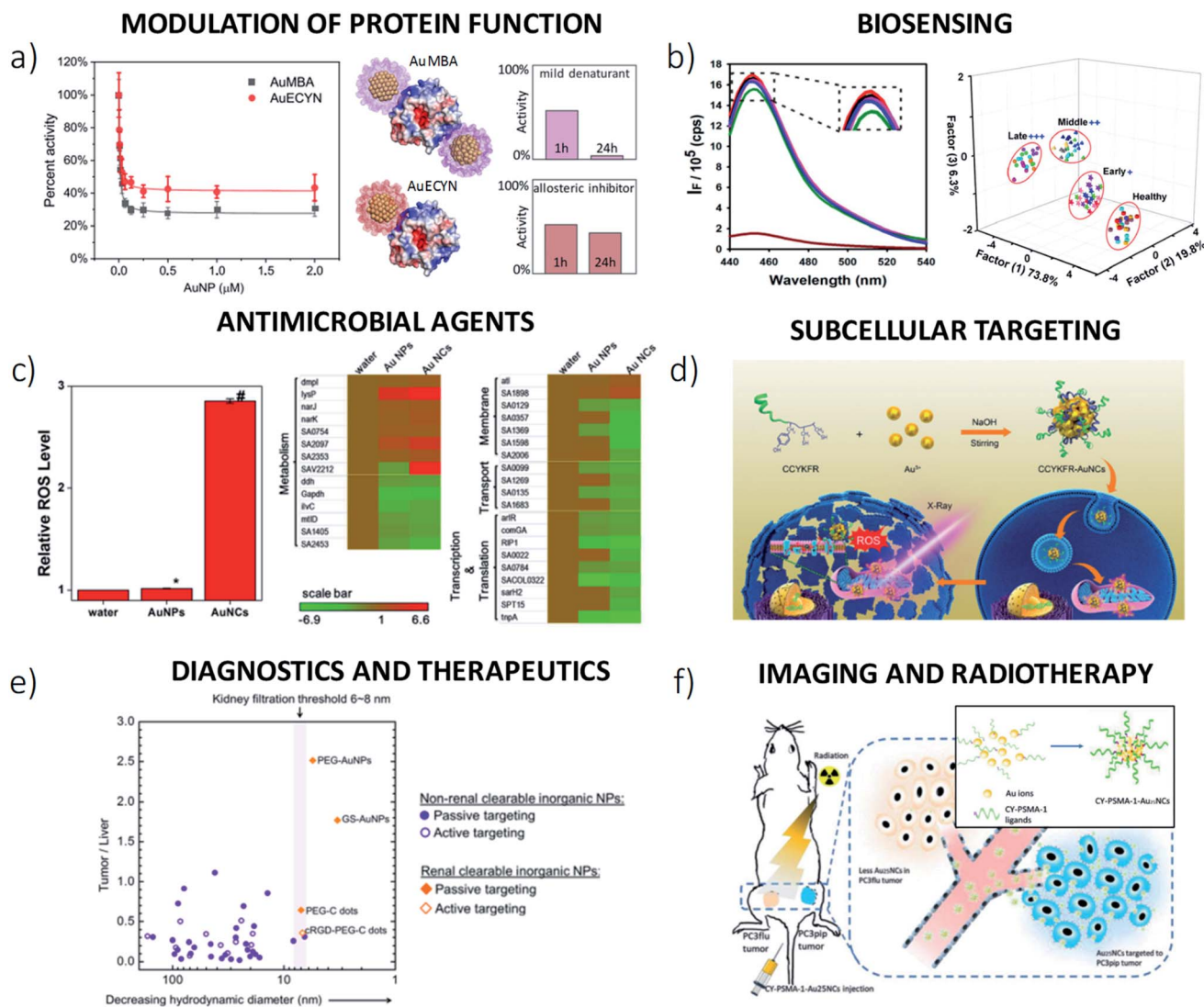
Besides usGNPs, classical NPs also offer an incredible sum of possibilities for the control of protein structure, stability, and function. For example, NP-induced enzymatic enhancement has been impressively demonstrated by means of the covalent attachment of either enzyme or substrate onto colloidal inorganic NPs. More information on this topic can be found in the works of Medintz and Algar, among others.<sup>290,291</sup>

### 6.2. Biosensing applications

usGNPs have been successfully developed for both the specific and selective (array-based) sensing of proteins. The former pertains to the use of usGNPs functionalized with recognition elements (such as peptides, proteins, or aptamers) for binding to specific proteins in biofluids, while keeping interactions with nonspecific proteins to a minimum. For example, Wen *et al.* demonstrated the use of peptide-templated fluorescent gold nanoclusters (GNCs) for sensing of the post-translational modification enzymes histone deacetylase 1 and protein kinase A (Fig. 6b, left panel).<sup>292</sup> Accordingly, enzymatic modification of the peptides damaged the protective coating around the metallic core and quenched the particle fluorescence, hence enabling the real-time and label-free detection of the tested enzymes. Peptide-templated usGNPs have also been developed into protease-responsive nanosensors for disease monitoring *in vivo*, as recently reported by Stevens, Bathia and co-workers.<sup>5</sup> Specifically, usGNPs were templated with a peptide substrate containing recognition motifs for thrombin or matrix metalloproteinase 9, as well as a biotin handle for conjugating to a neutravidin carrier protein. Disassembly of the larger neutravidin-particle complex by peptide cleavage in response to dysregulated protease activity led to the release of the usGNPs in the urine, thus enabling a direct calorimetric readout of disease state.

The array-based sensing approach, in turn, relies on selective and cross-reactive analyte recognition by multiple sensor elements to generate unique fingerprints of the sample under study.<sup>293–296</sup> The response pattern of sensor arrays is controlled in a synergistic manner by sensor interactions with all sample constituents, including proteins, small molecules, and ions. This sensing approach (reminiscent of olfaction/taste) complements the more traditional detection systems, in which peptides, antibodies or aptamers are employed to identify biomolecules within complex mixtures with high specificity. Array-based sensing is a powerful strategy for disease detection in situations where traditional immunosensing has limitations, such as in the simultaneous detection of multiple proteins for proper diagnosis, or when specific biomarkers are not available. Fluorescent GNCs constitute attractive building blocks for the generation of protein sensing arrays since they can singly act as selective recognition elements and signal transducers (the latter owing to the sensitivity of the fluorescence emission of the GNCs to protein binding). For example, Yuan *et al.* synthesized





**Fig. 6** Applications of usGNPs. (a) Allosteric regulation of  $\alpha$ -thrombin enzymatic activity.<sup>23</sup> Complexation to exosites 1 and 2 caused partial, sub-maximal inhibition of activity (left panel). Over time, AuMBA inactivated the enzyme, whereas AuECYN sustained a constant level of inhibition (right panels). (b) Biosensing applications. Left panel: sensing of a post-translational modification enzyme, histone deacetylase 1 (HDAC).<sup>292</sup> Shown are fluorescence spectral responses of peptide-templated fluorescent GNCs to different proteins: blank (red), HDAC (brown), HDAC plus its specific inhibitor trichostatin A (green), other proteins (black, pink, and blue). Fluorescence quenching observed with HDAC only. Right panel: array sensing for disease diagnosis using six amino acid-decorated fluorescent GNCs.<sup>72</sup> Shown is a canonical score plot discriminating serum samples from breast cancer patients at different stages of the disease (early, middle, and late) and healthy subjects. (c) Antimicrobial activity of GNCs.<sup>302</sup> GNCs (core size < 2 nm) induced intracellular ROS production in *Staphylococcus aureus* leading to cell death, whereas standard 6 nm-sized NPs did not (left panel). GNCs greatly affected gene expression related to cell metabolism, substrate transport, membrane integrity, and transcriptomic processes (middle and right panels); red and green designate gene up- and down-regulation, respectively. (d) Targeting of subcellular compartments in cancer cells.<sup>316</sup> Schematics of a GNC templated with a short peptide enabling selective targeting of mitochondria following endocytosis. The GNCs acted as radiosensitizers, leading to cell death upon X-ray irradiation. (e) Tumor-to-liver ratios of conventional NPs and renal-clearable usNPs in subcutaneous tumor mouse model following intravenous injection.<sup>9</sup> (f) Active *in vivo* tumor targeting of GNCs.<sup>11</sup> Schematics of GNCs templated with a highly selective PSMA ligand for prostate cancer. The GNCs showed enhanced targeting and improved radiation therapy toward PSMA-expressing tumor cells (blue) relative to PSMA-negative tumor cells (light orange). Adapted with permission from ref. 23 ((a); Copyright 2007 American Chemical Society), ref. 292 ((b), left panel; Copyright 2013 American Chemical Society), ref. 72 ((b), right panel; Copyright 2017 Elsevier), ref. 302 ((c), Copyright 2017 American Chemical Society), ref. 316 ((d), published by the Royal Society of Chemistry), ref. 9 ((e), Copyright 2015 American Chemical Society), and ref. 11 ((f), Copyright 2019 Wiley-VCH Verlag GmbH & Co. KGaA).

eight dual-ligand GNCs and demonstrated the ability of such a sensing system to discriminate among 8 different proteins in buffer.<sup>297</sup> In another study, Luo and co-workers prepared six different amino acid-decorated GNCs to use as selective sensors

for protein discrimination and disease diagnosis. Significantly, they showed that serum samples from healthy subjects and from breast cancer patients at different stages of the disease could be well distinguished by the sensor array (Fig. 6b, right



panel).<sup>72</sup> Array-based sensing can also be applied to discriminate cell types (e.g., healthy vs. cancerous) or to assess cellular response to drug administration.<sup>298,299</sup> For example, Tao *et al.* implemented seven dual-ligand GNCs in a sensor array to distinguish healthy, cancerous and metastatic human breast cancer cells.<sup>300</sup>

The integration of specificity and selectivity within the same sensor platform could provide superior discriminating power and higher stability against background variations.<sup>295</sup> For example, class specificity – that is, the ability to interact preferentially to certain classes of analytes – could be incorporated into usGNPs by rational design of their surface chemistries.

### 6.3. Detection and killing of pathogenic bacteria

usGNPs can be exploited for the efficient detection and killing of pathogenic bacteria.<sup>301</sup> For example, Xie and co-workers prepared anionic usGNPs with effective wide-spectrum antimicrobial activity.<sup>302</sup> The particles killed both Gram-negative and Gram-positive bacteria by way of inducing a metabolic imbalance in the cells and generating intracellular reactive oxygen species (ROS). Interestingly, a related antimicrobial activity was absent with larger ( $d = 6$  nm) NPs (Fig. 6c). These same authors later investigated the effects of surface chemistry on the antimicrobial properties of GNCs of identical core sizes.<sup>303</sup> They found (somewhat surprisingly) that GNCs displaying more negatively charged surfaces exhibited higher bacterial killing efficiencies. This behavior was ascribed to the generation of higher levels of ROS with the more anionic particles. In another study, Rotello and co-workers demonstrated the importance of surface hydrophobicity on the antimicrobial activity of cationic usGNPs against multi-drug-resistant bacteria.<sup>304</sup>

As illustrated above, usGNPs can kill pathogenic bacteria without having to rely on the recognition of surface receptors. Nonetheless, biospecificity can be conferred to usGNPs by engineering recognition elements on the particle monolayer. For instance, Mukherji *et al.* synthesized GNCs decorated with quorum sensing signal molecules for the selective detection of bacterial strains expressing quorum sensing receptors.<sup>305</sup> Gao and co-workers prepared fluorescent GNCs surface-coated with a peptide having high binding specificity to *Staphylococcus aureus*.<sup>306</sup> Besides enabling bacterial detection, the peptide-templated GNCs displayed effective antimicrobial activity by virtue of membrane disruption and increased ROS production. Chan *et al.* showed that mannose-capped GNCs could bind selectively to an uropathogenic *Escherichia coli* strain containing mannose-binding receptors.<sup>307</sup> Presumably, usGNPs decorated with different glycan moieties could be employed to detect and kill bacteria expressing distinct glycan receptors.

Finally, Jiang and co-workers reported the preparation of GNCs terminated with quaternary ammonium cations to serve as potent antibiotics against multi-drug-resistant Gram-positive bacteria *in vivo*.<sup>308</sup> The antimicrobial activity of the particles derived from a combination of membrane disruption, metabolic disturbance and ROS generation. The GNCs did not induce drug resistance and, remarkably, showed no detectable toxicity *in vivo*, thus raising the prospect for further clinical use.

### 6.4. *In vitro* cell biology applications

Luminescent GNCs have been intensively investigated as novel nanoscale probes for cell biology applications, offering many potential advantages in relation to conventional organic dyes and quantum dots, such as ultrasmall sizes, facile synthesis, large Stoke shifts, color tunability, long emission lifetimes (up to the  $\mu$ s range), environmentally sensitive emissions, photo and chemical stability, and biocompatibility.<sup>309</sup> Possible cell biology applications of usGNPs and GNCs include bioimaging,<sup>309</sup> gene delivery,<sup>310</sup> monitoring intracellular pH and temperature changes,<sup>311–313</sup> among many others. Of particular interest, GNCs coated with targeting moieties can be used to label specific subcellular compartments for imaging or even the manipulation of cell state and behavior. In an earlier study, Gao and co-workers developed a peptide-templated fluorescent GNC containing a TAT nuclear localization sequence for the targeted staining of cell nuclei.<sup>314</sup> The same authors later prepared GNCs functionalized with a short peptide sequence specific for  $\alpha_{11b}\beta_3$  integrin receptors.<sup>315</sup> The number of receptors per cell was then estimated by quantifying the gold content in single cells using laser ablation ICP-MS. In another application, Liu and co-workers prepared peptide-templated fluorescent GNCs and demonstrated their targeting/accumulation on the mitochondria of cancer cells following endocytosis.<sup>316</sup> The internalized GNCs acted as radiosensitizers, inducing significant mitochondrial oxidative stress and DNA damage upon X-ray irradiation, leading to cell death (Fig. 6d).

At this point, we recall that usGNPs functionalized with antibodies or Fab fragments have a long history of use as immunolabeling reagents for electron microscopy.<sup>317</sup> Given the new developments in the area of luminescent metallic NCs, it is conceivable that peptide- or protein-templated GNCs could turn into valuable tools for correlative light-electron microscopy studies, insofar as these novel nanomaterials may singly integrate the properties of specific antigen binding, fluorescence emission (thus supporting light microscopy imaging), and high electron scattering power (thus supporting high-resolution electron microscopy imaging<sup>318</sup>). Incidentally, such a powerful combined approach could furnish a better understanding of the mechanisms of cell entry and intracellular trafficking of usGNPs.<sup>319–322</sup>

### 6.5. *In vivo* disease diagnosis and therapy

The last decade has seen a surge in the development of usGNP-based platforms for *in vivo* disease diagnosis and treatment. Due to space limitations, here we consider only applications of usGNPs for cancer therapy.<sup>14,323</sup>

Studies so far have confirmed that usGNPs are biocompatible,<sup>15</sup> at least partly due to the intrinsic short-lived nature of their interactions with proteins.<sup>24</sup> usGNPs can be rapidly cleared from blood circulation through the kidneys (clearances over 50% at 24 h p.i. have been routinely reported) and escape liver uptake.<sup>9,324</sup> Tumor uptake levels of non-targeted usGNPs are generally around  $\sim 1$ –15% ID per g of tissue, which is a typical range for larger NPs as well.<sup>9</sup> The combined result is that



usGNPs may display significantly higher tumor-to-liver ratios relative to their classical counterparts (Fig. 6e).<sup>9</sup>

In a groundbreaking study, Zheng and co-workers reported that the *in vivo* behavior of atomically precise GNCs was size-dependent.<sup>35</sup> They found that smaller particles were more efficiently trapped by the glomerular glycocalyx (an effect similar to the separation principle in size-exclusion chromatography) and displayed longer blood retention times. As a result, a decrease of just a few atoms (Au<sub>25</sub> to Au<sub>10</sub>) in NC size enhanced the passive tumor uptake of GNCs from about 2 to 8% ID g<sup>-1</sup>. These results underscore the complexity and, sometimes, unpredictability of nano-bio interactions *in vivo*.

So far, usNPs have been developed as self-therapeutic agents,<sup>325</sup> drug delivery vehicles,<sup>326</sup> radiosensitizers for cancer radiotherapy,<sup>6,327,328</sup> among other applications in cancer nanomedicine. However, despite promising results so far, commonly encountered limitations such as restricted tumor uptake and off-target localization may hamper the clinical applications of usGNPs. Researchers have therefore been developing targeted usGNPs in an attempt to improve their tumor accumulation profiles. For example, Liang *et al.* observed a higher accumulation of usGNPs modified with a cyclic-RGD peptide in  $\alpha_v\beta_3$ -positive tumors relative to their non-targeted counterparts (6.4 ± 1.3% *vs.* 2.0 ± 1.8% ID/g tissue).<sup>8</sup> The enhanced tumor localization of the targeted usGNPs produced a stronger radiosensitizing effect leading to a reduced tumor mass over time. In another study, Basilion and co-workers prepared GNCs covered with a high-affinity targeting ligand for the prostate-specific membrane antigen (PSMA) receptor, which is over-expressed in most prostate cancers. This active targeting strategy significantly improved both tumor accumulation in mice (8.9% ID/g tissue) and the outcome of radiotherapy relative to a passive targeting approach (Fig. 6f).<sup>11</sup>

Further advancements in the field might benefit from a more quantitative knowledge of the interactions of both targeted and non-targeted usGNPs with proteins. For example, it is typically not known whether targeted ligands on usGNPs are able to maintain the same affinity of binding toward their receptors in comparison to free ligands in solution. In addition, although usGNPs may be deemed highly inert against nonspecific protein binding, the extent of this resistance is generally not known. We recently implemented a simple pharmacokinetic model simulation to understand how the tumor uptake of usGNPs could be modulated by ultraweak ( $K_D > 1$  mM) nonspecific interactions between usGNPs and serum proteins.<sup>329</sup> The results revealed that under “high” rates of NP plasma clearance and “slow” rates of vascular extravasation, ultraweak interactions – but not active targeting – would increase the tumor accumulation of the particles (by essentially slowing-down particle clearance from circulation). On the other hand, under “high” rates of plasma clearance and “high” rates of vascular extravasation and intravasation, high-affinity active targeting – but not ultraweak interactions – would enhance tumor uptake levels. Despite its underlying simplicity and many approximations, the model has provided a basic framework to quantitatively understand the blood and tumor pharmacokinetics of usGNPs under the influence of ultraweak protein interactions.

## 7. Outlook and concluding remarks

Compared to their larger counterparts, usGNPs display unique physicochemical properties and interaction modalities with biomolecular structures, providing a new paradigm in nano-biomedical research with potential impact in translational medicine. On the one hand, usGNPs represent powerful tools for fundamental research, helping to generate new knowledge on biosystems, all the way from single proteins to whole-organism physiology. At the same time, usGNPs have been intensively exploited in a remarkable number of biological and biomedical applications, both *in vitro* and *in vivo*, including biosensing, bioimaging, drug delivery, and cancer theranostics.

We have emphasized how unleashing the full potential of usGNPs in both fundamental and applied research will require a more in-depth understanding of their biomolecular interactions, much in the same way that a detailed knowledge of protein–protein complexation is a prerequisite to understanding higher-order biological processes. Notwithstanding the current gaps in our knowledge, it is clear that much is already known about the properties and behavior of usGNPs in complex biosystems and how such properties and behavior can be reasonably manipulated by nanoengineering. All in all, usGNPs appear ready to impact real-world clinical applications, as has already been the case for classical NPs.

Future advances in usNP-based biomedicine will likely depend on systematic studies through an integrated approach of experiments *in vitro* for strict control of environmental effects, computer simulations with advanced multiscale forcefields, and data-driven techniques for reverse-engineering NPs and NCs with desirable properties. However, progress is needed in all these fronts, including better techniques to manufacture and characterize ultrascale nanostructures (nanoengineering, chemical synthesis) and a clearer picture of the interfacial forces at the sub-nanometer length scale, particularly in usNP/liquid boundaries (theory and experiments). We have addressed some of the topics in this review.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

AAS was supported by the São Paulo Research Foundation (FAPESP, grant #19/04372-6); PS was supported by the NIH Intramural Research Program of NIBIB; and SAH was supported by the NIH Intramural Research Program through the CIT and NIAID and utilized the high-performance computational capabilities of the Biowulf HPC cluster at the NIH.

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