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# The unexpected mechanism of transformation from conventional room-temperature phosphorescence to TADF-type organic afterglow triggered by simple chemical modification†

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The study of transformation of photophysical behaviours in organic afterglow systems has emerged as an important topic, whereas the transformation in the reported studies only gives change of afterglow colour, duration or intensity. Herein, we report a serendipitous finding of the mechanism of afterglow transformation from conventional RTP to TADF-type organic afterglow triggered by simple chemical modification of coronene systems; usually, chemical modification can only lead to spectral shifts of luminescent systems. Coronene molecules show typical RTP behaviours when doped in organic matrices. After being substituted by difluoroboron β-diketonate moieties, the coronene-containing materials exhibit a TADF-type organic afterglow mechanism, which features a moderate $k_{RISC}$ to harvest triplet energies, enhance afterglow efficiency, and maintain long afterglow lifetimes. Interestingly, the TADF-type afterglow materials can be excited by visible lights, possess emission wavelength > 600 nm and PLQY > 40%, display excellent processability into desired patterns and aqueous dispersion, and function as high-contrast *in vivo* bioimaging agents. The present study provides a unique pathway for the manipulation of triplet excited states to fabricate high-performance organic afterglow materials.

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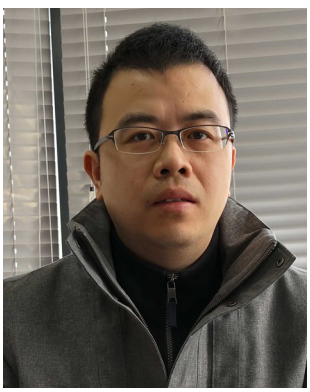
## Introduction

Control of the excited state properties is of fundamental importance for constructing novel and high-performance luminescent materials.[1–3] In room-temperature phosphorescence (RTP) and organic afterglow systems, the understanding and manipulation of triplet excited states represent a central topic since the population and decay of triplet excited states determine the photophysical properties of organic systems.[4–9] Pioneers in the fields showed that although phosphorescence emission is spin-forbidden, the introduction of n–π* transition and heavy atom effects can facilitate intersystem crossing and enhance the radiative decay of triplet excited states to allow the occurrence of organic RTP.[10–13] Crystalline and glassy environments have also been reported to inhibit the nonradiative decay of triplet excited states for the fabrication of RTP and organic afterglow materials.[14–16] A supramolecular assembly such as host–guest inclusion can be used to restrict intramolecular motions of luminescent guests to obtain efficient organic RTP



Kaka Zhang

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materials.[17–21] Recent studies exhibit the achievements of ultralong organic afterglow durations in donor–acceptor systems *via* photo-induced charge separation and slow charge recombination mechanisms.[22–25] Elegant design of molecular systems and the control of the triplet excited states have been reported to lead to very high RTP and afterglow efficiency in organic systems.[26]

Besides achieving high afterglow efficiency and long afterglow lifetimes, the study of the transformation of photophysical properties in organic RTP and afterglow systems has attracted increasing interest.[27–41] Early studies on this topic showed the preparation of oxygen-, mechanical force, pH- and ion-responsive RTP and afterglow materials with intriguing photophysical transformation behaviours.[27–35] Recently, it has been reported that colour-tunable RTP materials can be obtained by adjusting the doping concentration of luminescent molecules in organic matrices.[36,37] Emission lifetimes of RTP systems have been found to be modulated by the triplet energy levels of organic matrices.[38] The Diels–Alder reaction has been applied to control the RTP and afterglow behaviours in organic systems.[39] Supramolecular recognition has also been reported to enhance the RTP properties of crown ether systems.[40] A very recent study in our laboratory showed that the reaction between pyrylium salts and amine-containing compounds can be used as time-gated afterglow chemodosimeters.[41] Despite these advancements in RTP systems, the mechanism of transformation of afterglow induced by specific stimuli or chemical reactions has been rarely reported. The study of the mechanism of transformation of afterglow not only has significant impact on the deep understanding of triplet excited state behaviours in organic systems, but also paves the way for high-performance organic afterglow materials.

Herein, we report a serendipitous finding of the mechanism of photophysical transformation from conventional RTP to TADF-type organic afterglow triggered by simple chemical modification of coronene systems; usually, chemical modification can only lead to spectral shifts of luminescent systems. Coronene molecules show the typical RTP mechanism when doped in 4-methoxybenzophenone matrices. After being substituted by difluoroboron β-diketonate moieties, the coronene-containing materials are found to exhibit the TADF-type organic afterglow mechanism, which features a moderate $k_{RISC}$ (rate constant of reverse intersystem crossing) to harvest triplet energies, enhance afterglow efficiency, and maintain afterglow lifetimes. The TADF afterglow materials can be excited by visible lights, possess emission wavelength > 600 nm and PLQY > 40%, display excellent processability into desired patterns and aqueous dispersion, and function as high-contrast *in vivo* bioimaging agents.

## Results and discussion

By doping coronene (Cor) into 4-methoxybenzophenone (MeOBP) at 0.1 wt%, Cor-MeOBP-0.1% samples exhibited bright green organic afterglow under ambient conditions (Fig. 1A).

The Cor-MeOBP-0.1% materials displayed a fluorescence band in the range of 400 nm to 480 nm in their steady-state emission spectra, exhibited a phosphorescence band ranging from 500 to 650 nm in the delayed emission spectra (1 ms delay), and possessed a long phosphorescence lifetime up to 4.12 s (Fig. 1B, C and Fig. S1, ESI†). The afterglow of Cor-MeOBP-0.1% materials mainly originates from room temperature phosphorescence. At room temperature, the delayed emission spectra exhibit a small delayed fluorescence band ranging from 400 nm to 480 nm. Upon lowering temperature, this 400–480 nm delayed fluorescence band first decreases and then disappears at 243 K (Fig. S2, ESI†). TD-DFT calculation on the ORCA 4.2.1 program with the B3LYP/G functional and def2-TZVP(-f) basis set shows that coronene has very small oscillator strength ($f_{osc}$) in its $S_1$ state due to symmetry-forbidden $S_1$–$S_0$ transition (Fig. 1D), which would allow the occurrence of intersystem crossing to a large extent. It is noteworthy that $T_1$–$S_0$ transition shows very a small spin–orbit coupling matrix element (SOCME), which agrees well with the long phosphorescence lifetimes of Cor-MeOBP-0.1% materials. These, together with the reported studies on coronene systems,[42–44] motivate us to propose that the coronene molecules dispersed in MeOBP matrices can be excited by UV or visible light to form singlet excited states, then undergo intersystem crossing to reach triplet excited states, and subsequently emit ultralong RTP due to the small $k_P$ (rate constant of phosphorescence decay) of coronene and very small $k_{nr} + k_q$ (rate constants of nonradiative decay and oxygen quenching) resulted from the rigid environment provided by MeOBP matrices (Fig. 1E).

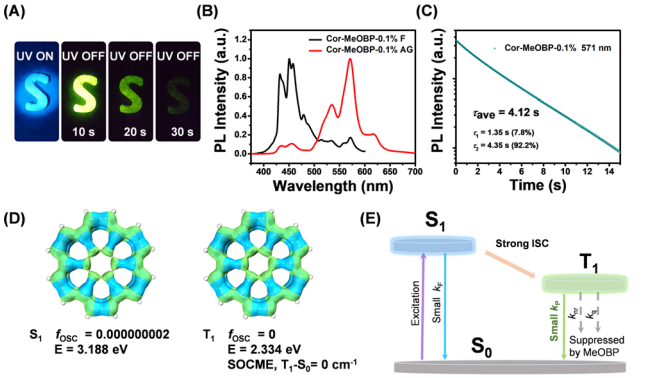


Fig. 1 (A) Photographs of Cor-MeOBP-0.1% materials under 365 nm UV light and after removal of an UV lamp. (B) The steady-state emission spectra and delayed emission (1 ms delay) spectra of Cor-MeOBP-0.1% materials. (C) Room-temperature emission decay of Cor-MeOBP-0.1% materials monitored at 571 nm. (D) Isosurface maps of the electron-hole density difference of the lowest singlet excited states and lowest triplet excited states of the coronene molecule, where blue and green isosurfaces correspond to hole and electron distributions, respectively. (E) The proposed mechanism of room-temperature phosphorescence in the Cor-MeOBP system, which features small $k_F$, small $k_P$, and very small $k_{nr} + k_q$.

Given that the long and bright afterglow of the coronene-MeOBP two-component system, we chemically modified the coronene to regulate the UV-vis and emission spectra of the materials. To this end, luminescent difluoroboron β-diketonate ($BF_2$bdk) compounds were synthesized *via* the cascade reaction developed in our laboratory,[45–47] then separated and purified by column chromatography and recrystallization successively (Fig. 2A). The structure of Cor$BF_2$ compounds was characterized by LRMS, HRMS, NMR, and FT-IR (see the ESI†). UV-vis spectra

 

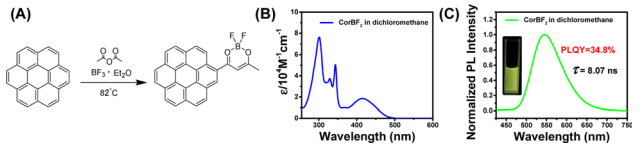


Fig. 2 (A) Cascade synthesis of luminescent $BF_2$bdk compounds. (B) UV-vis spectra of $CorBF_2$ in the dichloromethane solution. (C) Steady-state emission spectra of $CorBF_2$ in the dichloromethane solution.

of $CorBF_2$ in dichloromethane solutions showed broad absorption bands ranging from 360 nm to 500 nm with absorption maxima at 415 nm and a molar absorption coefficient ($\varepsilon$) of 18000 $M^{-1}$ $cm^{-1}$ (Fig. 2B). Upon excitation, $CorBF_2$ in the dichloromethane solution displayed greenish yellow emission with structureless bands at 547 nm in its steady-state emission spectra (Fig. 2C; PLQY, 34.8%; fluorescence lifetime, 8.07 ns), which is much red-shifted compared to the coronene solution (Table 1). TD-DFT calculation showed that the $S_1$ state of $CorBF_2$ possessed intramolecular charge transfer (ICT) characters from the coronene group to the dioxaborine ring plus localized excitation (LE) characters (Fig. S3, ESI†), which agree with the positive solvatochromicity of $CorBF_2$ solutions (Fig. S4, ESI†). The $T_1$ state of $CorBF_2$ also shows ICT (from the coronene group to the dioxaborine ring) plus LE (within the coronene group) characters. The $CorBF_2$ in the solid state and in dichloromethane solution did not show afterglow upon switching off the UV lamp under ambient conditions (Fig. S5, ESI†).

The introduction of organic matrices to control excited states of luminescent dopants has been demonstrated by us and other research groups to be significant for constructing organic room-temperature afterglow materials.[22,23,48–51] Various organic matrices have been selected to accommodate $CorBF_2$ dopants (Fig. S6, ESI†). When doping $CorBF_2$ into MeOBP at 0.01%, the $CorBF_2$-MeOBP-0.01% powders showed green emission under 365 nm UV light and performed green afterglow with duration longer than 4.5 s in a dark room after the removal of the UV lamp (Fig. 3A). The steady-state emission spectra of $CorBF_2$-MeOBP materials have been found to be almost identical to the delayed emission (1 ms delay) bands (Fig. 3B). It is an unexpected finding because, in general, the delayed emission spectra of afterglow materials show a significant spectral red-shift compared to the steady-state emission

Table 1 Photophysical data of luminescent compounds in dichloromethane at room temperature

| Entry | $\lambda_{abs}$/nm ($\varepsilon \times 10^{-4} M^{-1}$ $cm^{-1}$) | $\lambda_{em}$/nm | $\tau$/ns | $\Phi$/% |
|---|---|---|---|---|
| Coronene in dichloromethane | 304 (14.0)<br>325 (1.5)<br>341 (3.2) | 445 | 1.7 (40%)<br>20.6 (60%) | 4.8% |
| $CorBF_2$ in dichloromethane | 301 (7.6)<br>329 (3.7)<br>344 (5.0)<br>415 (1.8) | 547 | 8.07 | 34.8% |
| $R_1$ in dichloromethane | 303 (9.2)<br>341 (3.6)<br>547 (9.2) | 640 | 1.58 | 27.8% |

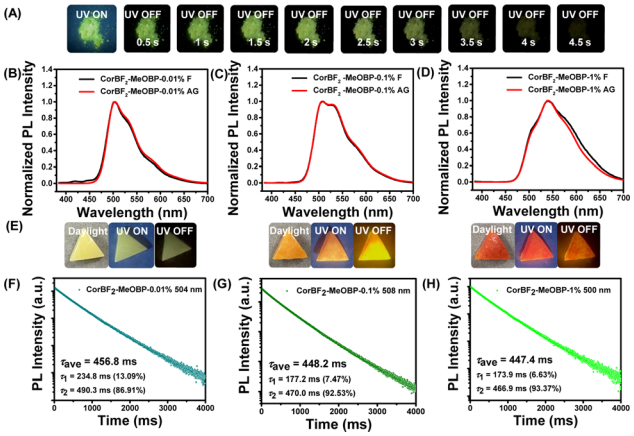


Fig. 3 (A) Photographs of the $CorBF_2$-MeOBP-0.01% powders under a UV lamp and after ceasing UV excitation. (B–D) Room-temperature steady-state and delayed emission (1 ms delay) spectra of $CorBF_2$-MeOBP materials with different doping concentrations (B, 0.01%; C, 0.1%; D, 1%). (E) Photographs of the $CorBF_2$-MeOBP melt-cast afterglow objects (left, $CorBF_2$-MeOBP-0.01%; middle, $CorBF_2$-MeOBP-0.1%; right, $CorBF_2$-MeOBP-1%) under a 365 nm UV lamp and after ceasing UV excitation, respectively. (F–H) Room-temperature emission decay of $CorBF_2$-MeOBP-0.01% materials monitored at 504 nm (F), $CorBF_2$-MeOBP-0.1% melt-cast materials monitored at 508 nm (G) and $CorBF_2$-MeOBP-1% materials monitored at 500 nm (H). These emission decay profiles were recorded using a Hitachi FL-4700 fluorescence spectrometer equipped with chopping systems.

spectra, for example, in the case of the coronene-matrix system (Fig. 1B). The steady-state emission spectra and delayed emission spectra of $CorBF_2$-MeOBP-0.01% and $CorBF_2$-MeOBP-0.001% materials, where $CoBF_2$ molecules are molecularly dispersed in MeOBP matrices, showed broad and vibronic-structured emission bands ranging from 450 nm to 650 nm (Fig. 3B and Fig. S7, ESI†), which agree with TD-DFT calculation that $CorBF_2$ possesses ICT plus LE characters (Fig. S3, ESI†).

The coincidence of steady-state and delayed emission spectra has been reported in organic systems of triplet-to-singlet excited state energy transfer from the RTP donor to the fluorescence acceptor.[52–56] It is known that the triplet excited state of benzophenone (BP) derivatives can be easily populated due to their highly efficient intersystem crossing (ISC).[11] The excited state energy transfer can only occur when the matrices are sufficiently excited. In the present study, when excited at longer wavelength such as 405 nm and 420 nm visible lights, $CorBF_2$-MeOBP materials still showed significant afterglow (Fig. S8 and S9, ESI†). MeOBP matrices possess insignificant absorption at 420 nm, so we reason that the excited state energy transfer from MeOBP to $CorBF_2$ may exist when the samples are excited at 365 nm but is not necessary for the emergence of organic afterglow and the coincidence of steady-state and delayed emission spectra in the present $CorBF_2$-MeOBP system. In the literature,[22,23] the donor–acceptor system of intermolecular charge transfer characters exhibits that the organic long persistent luminescence (OLPL) mechanism can also show identical spectra between steady-state and delayed emission spectra according to recent reported studies. Cyclic voltammetry of $CorBF_2$ and MeOBP has been performed. MeOBP matrices

(HOMO, −6.24 eV; LUMO, −2.56 eV) possess low-lying HOMOs and high-lying LUMOs when compared to $CorBF_2$ dopants (HOMO, −5.80 eV; LUMO, −2.82 eV), which suggests the absence of an exciplex between MeOBP and $CorBF_2$. UV-vis absorption spectra and excitation spectra of $CorBF_2$-MeOBP-0.01% materials also showed absorption bands with maxima similar to the UV-vis absorption spectra of $CorBF_2$ in dichloromethane solutions (Fig. S10, ESI†). On the other hand, the two-photon ionization mechanism has been reported to give rise to OLPL.[57,58] Since the two-photon ionization process involves two sequential excitation steps, the brightness and duration of OLPL are strongly dependent on excitation power. In the present study, power-dependent delayed fluorescence intensity measurements showed a quasi-linear relationship between delayed fluorescence intensity and excitation power (Fig. S11, ESI†). These results can rule out the two-photon ionization mechanism in the $CorBF_2$-MeOBP system. In addition, the afterglow decay profiles of $CorBF_2$-MeOBP materials follow exponential decay (rather than power law decay observed in some OLPL systems). Impurity has a vital contribution to organic afterglow reported in the reported studies.[59] There are recent studies showing that the existence of impurities has critical impact on some organic afterglow systems. Considering the similar maxima of excitation spectra and solid UV-vis absorption spectra of $CorBF_2$-MeOBP-0.01% materials (Fig. S10, ESI†) and the relatively high purity of $CorBF_2$ confirmed by HPLC (Fig. S12A, ESI†), the afterglow was not likely to be originated from the impurity mechanism. Furthermore, by doping higher purity of $CorBF_2$ (separated and collected fractions through HPLC) into MeOBP, we have prepared $CorBF_2$-MeOBP (HPLC) materials whose room temperature emission spectra and emission decay profile showed similar afterglow properties compared to $CorBF_2$-MeOBP materials (Table 2, Fig. 3B, F and Fig. S12, ESI†). We have also purified MeOBP matrices *via* two cycles of recrystallization. The purified MeOBP matrices showed very weak delayed emission signals and the absence of room-temperature afterglow (Fig. S13, ESI†). These results powerfully proved that the impurity mechanism can be ruled out in the present study.

Table 2 Photophysical properties of dopant-matrix afterglow materials under ambient conditions

| Sample | $\lambda_F$/nm | $\lambda_{AG}$/nm | $\tau_{AG}$ (ms) | $\Phi$ (%) |
|---|---|---|---|---|
| Cor-MeOBP-0.1% | 450 | 571 | 4113 | — |
| $CorBF_2$-MeOBP-0.01% | 504 | 504 | 456.8 | — |
| | 534 | 534 | 456.0 | |
| | 590 | 590 | 499.2 | |
| $CorBF_2$-MeOBP-0.1% | 508 | 508 | 448.2 | 42.5 |
| | 534 | 534 | 436.1 | |
| | 595 | 595 | 466.4 | |
| $CorBF_2$-MeOBP-1% | 500 | 500 | 447.4 | — |
| | 541 | 541 | 429.1 | |
| | 590 | 590 | 462.7 | |
| $R_1$-MeOBP-0.002% | 610 | 610 | 31.77 | — |
| $R_1$-MeOBP-0.01% | 619 | 619 | 26.03 | 43.7 |
| $R_1$-MeOBP-0.02% | 619 | 619 | 49.30 | — |
| $R_1$-MeOBP-0.2% | 619 | 619 | 81.10 | — |

To study the afterglow mechanism of the $CorBF_2$-MeOBP system, we performed low temperature delayed emission experiments and observed $CorBF_2$-MeOBP-0.01% powders showing orange afterglow after switching off the UV lamp at 77 K (Fig. 4A). Temperature-dependent delayed emission spectra from 77 K to 300 K have been recorded to study the triplet excited state of $CorBF_2$-MeOBP-0.01%, where the phosphorescence ($\lambda_P$ = 584 nm) is dominant in the emission spectra at 77 K, and a significant delayed emission band ranging from 450–550 nm gradually appears with the increase of the experimental temperature (Fig. 4B). Especially when the temperature has been set at 298 K, the delayed emission spectra nearly coincide with the steady-state emission spectra. In this case, the afterglow mechanism could be either triplet–triplet annihilation (TTA) or thermally activated delayed fluorescence (TADF). The TTA mechanism that is a bimolecular mechanism should have insignificant contribution to the organic afterglow due to the low doping concentration, 0.01% and even 0.001%. Temperature-dependent delayed emission lifetime measurements (monitored at 504 nm) of the $CorBF_2$-MeOBP-0.01% materials have been performed. Upon increasing the temperature, the delayed fluorescence lifetimes have been found to decrease (Fig. S14, ESI†). These observations suggest that $k_{RISC}$ increases with temperature, which supports the TADF-type afterglow mechanism in the present study. Furthermore, power-dependent delayed fluorescence intensity measurements showed a quasi-linear relationship between delayed fluorescence intensity and excitation power (Fig. S11, ESI†). These experiments together with TD-DFT calculation (*vide infra*) all supported that $CorBF_2$-MeOBP-0.01% materials undergo a TADF-type organic afterglow mechanism under ambient conditions.

By increasing doping concentrations, the $CorBF_2$-MeOBP-0.1% and $CorBF_2$-MeOBP-1% materials also exhibited the coincidence

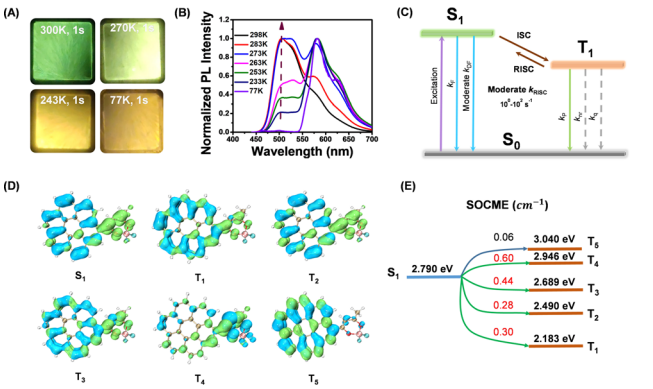


Fig. 4 (A) Photographs of the afterglow of the $CorBF_2$-MeOBP-0.01% materials at different temperatures. (B) Variable temperature delayed emission (1 ms delay) spectra of $CorBF_2$-MeOBP-0.01% materials (excited at 420 nm). (C) Schematic illustration of the photophysical processes in the TADF-type organic afterglow system which features a moderate $k_{RISC}$. (D and E) TD-DFT-calculated electron density difference of singlet and triplet excited states of $CorBF_2$. TD-DFT calculations were performed on the ORCA 4.2.1 program with the B3LYP/G functional and def2-TZVP(-f) basis set (D) and energy levels and SOCME between $S_1$ and triplet excited states (E).

 

between steady-state and delayed emission spectra (Fig. 3C and D). The lower-energy vibronic signals at around 534 nm and 590 nm were found to increase with $CorBF_2$ doping concentrations, which should be caused by the aggregation of the planar $CorBF_2$ molecules in MeOBP matrices; similar photophysical behaviours of vibronic signal change upon aggregation have also been reported in the literature.[60,61] The $CorBF_2$ aggregation gave rise to the red shifts of afterglow colours from green to orange (Fig. 3E). The lifetime of excited states of $CorBF_2$-MeOBP-0.1% materials can be fit into double exponential decay with $\tau_1$ = 53.66 ms (5.92%) and $\tau_2$ = 471.15 ms (94.08%) and the PLQY had also been measured to be up to 42.5% (Fig. S15 and S16, ESI†). The lifetimes of the $\tau_1$ part that can be attributed to prompt fluorescence measured by a microsecond flash lamp were largely overestimated, similar to the reported studies.[62,63] Since the room-temperature afterglow can be attributed to the $\tau_2$ part (delayed fluorescence) in the decay profile, the afterglow quantum yield of $CorBF_2$-MeOBP-0.1% materials can be estimated to be 40.0%. The $k_{RISC}$ value of the $CorBF_2$-MeOBP-0.1% materials can be estimated from their delayed fluorescence lifetimes (471.15 ms) to be on the order of $10^0$ to $10^1$ $s^{-1}$. The moderate $k_{RISC}$ value enables the TADF mechanism to harvest triplet energies and maintain emission lifetimes $>0.1$ s, which is the key design factor of the present study to significantly improve afterglow quantum yields (Fig. 4C). It is noteworthy that the $k_{RISC}$ values in the $CorBF_2$-MeOBP-0.01% materials are much smaller than those in TADF-type OLED systems[64–67] ($k_{RISC}$, $10^3$ to $10^6$ $s^{-1}$) that possess large $k_{RISC}$ values leading to short emission lifetimes. TADF emitters with large $k_{RISC}$ values of $10^3$–$10^6$ $s^{-1}$ are necessary to construct efficient OLED devices. We understand that it is difficult to achieve large $k_{RISC}$ values of $10^3$–$10^6$ $s^{-1}$ in organic systems with $\Delta E_{ST}$ of 0.3 eV or 0.4 eV or above.[68] Usually, organic molecular systems should simultaneously meet two requirements to realize large $k_{RISC}$: (1) small $\Delta E_{ST}$, for example, smaller than 0.2 eV; (2) large SOCME.[64–66,69] In the absence of the heavy atom effect, relatively large SOCME can be realized by incorporating triplet excited states of different symmetries from the lowest singlet excited states according to the El-Sayed rule. We reason that if organic molecular systems meet requirement (1), or meet requirement (2), or partially meet requirements (1) and (2), a moderate $k_{RISC}$ value of $10^{-1}$–$10^2$ $s^{-1}$ can be obtained. Given that $k_P$ values in ordinary organic systems in the absence of the heavy atom effect are in the range of $10^{-2}$–$10^3$ $s^{-1}$, such a moderate $k_{RISC}$ value would be suitable to initiate a TADF mechanism to harvest triplet energy. Most conventional organic LE systems have very small $k_{RISC}$ values due to their large $\Delta E_{ST}$ and small SOCME. $CorBF_2$ molecules possess a much smaller singlet–triplet splitting energy than LE systems, which enhances intersystem crossing to a large extent. Based on the rational molecular design of $CorBF_2$ and a two-component strategy, the TADF-type afterglow materials with the afterglow quantum yield of 40.0% and lifetimes of 471.15 ms have been prepared by doping $CorBF_2$ compounds into MeOBP matrices.

Polymers like poly(methyl methacrylate) (PMMA) have also been selected as organic matrices to prepare $CorBF_2$-matrix afterglow materials. The $CorBF_2$-PMMA-0.01% materials (PLQY = 48.0%) have been found to show blue-shifted steady-state emission spectra ($\lambda_F$ = 490 nm, Fig. S17A, ESI†) when compared to $CorBF_2$-MeOBP-0.01% materials ($\lambda_F$ = 504 nm); PMMA has a smaller dipole moment than MeOBP. These results agree well with the description of the dipole effect,[51] that is, the dipole–dipole interactions between dopants' $^1$ICT states and organic matrices can reduce dopants' $S_1$ levels. The room-temperature delayed emission spectra of $CorBF_2$-PMMA-0.01% materials mainly consist of room-temperature phosphorescence signals, as well as a small amount of delayed fluorescence signals (Fig. S17B, ESI†); at 77 K, the delayed emission spectra exhibit a phosphorescence band with maxima at 577 nm and shoulder at 620 nm (Fig. S17C, ESI†). In contrast, the room-temperature delayed emission spectra of $CorBF_2$-MeOBP-0.01% materials showed predominant delayed fluorescence (Fig. 3B). These observations suggest the decrease of $k_{RISC}$ in PMMA matrices, since the relative intensity of delayed fluorescence to RTP is proportional to $k_{RISC}/k_P$. These results also agree well with the description of the dipole effect,[51] that is, the dipole–dipole interactions between dopants' $^1$ICT states and organic matrices can reduce $\Delta E_{ST}$ and facilitate dopants' ISC and RISC; it has been reported that the decrease of $\Delta E_{ST}$ by 0.05 eV would enhance $k_{RISC}$ by around 10 times.[64] Therefore, the selection of MeOBP matrices is very important for fabricating TADF-type organic afterglow materials in the present study.

TD-DFT calculations (ORCA 4.2.1 program with the B3LYP/G functional and def2-TZVP basis set) of $CorBF_2$ molecules show that there are rich ISC and RISC channels with SOCME values around 0.3 $cm^{-1}$ and above (Fig. 4D and E). Compared to coronene, the values of $\Delta E_{ST}$ of $CorBF_2$ are relatively small and the SOCME values of $CorBF_2$ are much larger (Fig. S18, ESI†). $CorBF_2$ also possess excited states with different ICT and LE components. As a result, some triplet excited states with different electronic configurations from the lowest singlet excited state ($S_1$) would show increased $k_{ISC}$ and $k_{RISC}$ values to a moderate level according to the El-Sayed rule. Furthermore, since $CorBF_2$-MeOBP-0.01% materials have been found to exhibit apparent TADF-type afterglow behaviours, the $T_1$ level was estimated to be 2.12 eV from their phosphorescence maxima at 77 K, and the $S_1$ level of $CorBF_2$-MeOBP-0.01% materials was estimated to be 2.46 eV from the fluorescence maxima at room temperature collected in the steady-state emission mode (Fig. 4B). The $\Delta E_{ST}$ of $CorBF_2$-MeOBP-0.01% materials can be estimated to be moderate at 0.34 eV. In view of the relatively small $\Delta E_{ST}$ and rich ISC/RISC channels, it is understandable that the present $CorBF_2$-MeOBP system can obtain a moderate $k_{RISC}$ value on the order of $10^0$ to $10^1$ $s^{-1}$ and undergo a TADF-type organic afterglow mechanism; $k_P$ values of the $CorBF_2$-MeOBP systems should be much smaller than $k_{RISC}$ at room temperature. The relative intensity of delayed fluorescence to RTP is proportional to $k_{RISC}/k_P$ of the $CorBF_2$-matrix system. In the case of $CorBF_2$-PMMA-0.01% materials, the room-temperature delayed emission spectra suggest that $k_P$ could compete or surpass $k_{RISC}$ (Fig. S17B, ESI†). In the case of $CorBF_2$-MeOBP-0.01% materials, because MeOBP matrices can significantly enhance $k_{RISC}$, the room-temperature delayed

emission spectra showed predominant delayed fluorescence (Fig. 3B). The $k_{RISC}/k_P$ of $CorBF_2$-MeOBP-0.01% materials is also dependent on temperature; $k_{RISC}/k_P$ increases with temperature.

We also studied the impact of oxygen on afterglow properties of $CorBF_2$-MeOBP-0.01% materials, and the emission lifetime of $CorBF_2$-MeOBP-0.01% materials under degassed conditions is slightly longer than that under ambient conditions (Fig. S19, ESI†). The observation showed that the rigid matrix efficiently protected the triplet excited state of $CorBF_2$ against oxygen quenching by the encapsulation under ambient conditions. Moreover, the fluorescence intensities of steady-state emission spectra of $CorBF_2$ in the dichloromethane solution under degassed conditions have also been recorded, which showed negligible changes compared to air conditions (Fig. S20, ESI†). Such behaviours are different from conventional TADF emitters of $k_{RISC}$ on the order of $10^3$–$10^6$ $s^{-1}$. In the present system, even though the triplet excited state of $CorBF_2$ *via* intersystem crossing could be populated, it can also be quenched by the active nonradiative decay process in solution states due to the small $k_{RISC}$. These observations about the relationship between oxygen and the TADF-type afterglow materials have also been reported in our previous study.[62,70]

A further chemical modification of $CorBF_2$ has been performed to obtain $R_1$ of the donor–acceptor–donor molecular design to achieve more red-shifted emission by aldol condensation between α-methyl of $CorBF_2$ and aldehyde group (Fig. 5A). $R_1$ has been separated and purified successfully by column chromatography and recrystallization, and the structure was determined through NMR, LRMS, HRMS, and FT-IR, respectively (see the ESI†). No afterglow has been observed for $R_1$ molecules in solid or solution states under ambient conditions (Fig. S21, ESI†). Both the UV-vis spectra and steady-state emission spectra of $R_1$ in the dichloromethane solution showed significant red shifts compared to those of $CorBF_2$, which showed the success of rational molecular design (Fig. 5B and C). To study the UV-vis spectra of $R_1$ molecules, the emergence of the absorption band at 547 nm can be attributed to $^1$ICT transition, and UV-vis spectra of $R_1$ in the dichloromethane solution exhibit an intense ICT absorption at 547 nm with the molar absorption coefficient (ε) as high as 92000 $M^{-1}$ $cm^{-1}$. These agree with the TD-DFT calculations where $S_0$ to $S_1$ transition possesses ICT characters and large oscillator strength (Fig. S22, ESI†). The steady-state emission spectra of the $R_1$ dilute solution in dichloromethane displayed a broad fluorescence band ranging from 550 nm to 800 nm with the emission maxima at 640 nm (Fig. 5C). A high PLQY of 27.8% and a nanosecond lifetime of 1.58 ns of $R_1$ in the dichloromethane dilute solution have been measured (Table 1) and a significant red-shifted emission with positive solvatochromicity has also been observed (Fig. S23, ESI†).

The selection of organic matrices to accommodate luminescent dopants and to facilitate intersystem crossing of excited states of dopants have been considered to be significant to prepare organic room-temperature organic afterglow materials using a two-component strategy. Various matrices have been tested as the second component to control the triplet excited state of $R_1$ in the $R_1$-matrix system (Fig. S24, ESI†). When doping $R_1$ into MeOBP at 0.002% to 0.2%, the $R_1$-MeOBP powders have been found to exhibit bright red afterglow after the removal of the UV lamp (Fig. 6). The afterglow colour is very similar to the fluorescence colour, and the steady-state emission spectra of $R_1$-MeOBP-0.002% are almost identical to the delayed emission spectra with the emission maxima of 610 nm excited by a 365 nm UV lamp. Interestingly, significant delayed emission spectra of $R_1$-MeOBP-0.002% powders can also be obtained by 520 nm excitation (Fig. S25, ESI†) and the red afterglow can be achieved by 532 nm laser excitation (Fig. S26, ESI†). Again, these suggest that even though the process of energy transfer from MeOBP to $R_1$ exists when the samples are excited by a UV lamp, it is not necessary for the emergence of red organic afterglow in the $R_1$-MeOBP system. The formation of an exciplex *via* intermolecular charge transfer that has been reported to give rise to organic afterglow can be ruled out in the $R_1$-MeOBP system because the MeOBP matrix possesses low-lying HOMOs and high-lying LUMOs when compared to $R_1$ dopants (Fig. S27, ESI†). Besides, solid UV-vis spectra and excitation spectra of $R_1$-MeOBP-0.002% possess similar maxima to the UV-vis spectra of $R_1$ in dichloromethane solutions (Fig. S27, ESI†). The HPLC experiment has been performed to exclude the impurity effects in afterglow materials. However, the reverse-phase HPLC profile of $R_1$ has been found to be very broad with significant tailing, which complicates the purity measurement. It is found that other boron difluoride hemicurcuminoid compounds in our lab also showed a very broad peak with significant tailing under

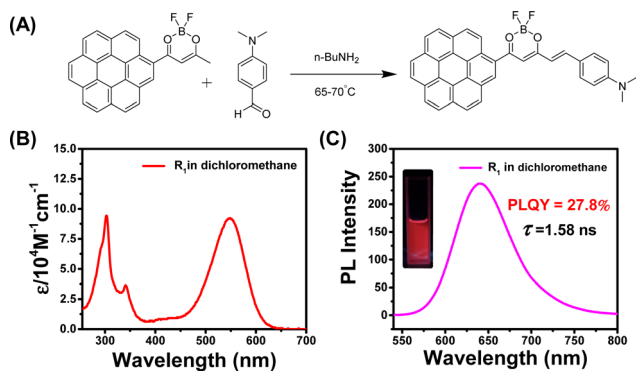


Fig. 5 (A) Synthesis of $R_1$ $BF_2$bdk luminescent compounds *via* aldol condensation. (B) UV-vis spectra of $R_1$ in the dichloromethane solution. (C) Steady-state emission spectra of $R_1$ in the dichloromethane solution.

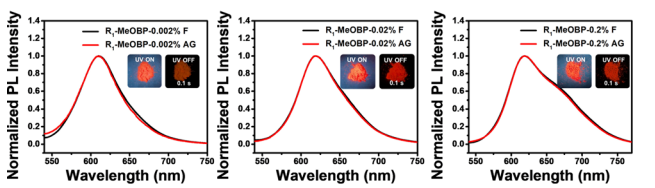


Fig. 6 The room-temperature steady-state and delayed emission (1 ms delay) spectra of $R_1$-MeOBP afterglow materials at doping concentrations of 0.002%, 0.02% and 0.2%, respectively.

 

reverse-phase HPLC measurements. We also tried normal-phase HPLC. However, because of the low polarity of $R_1$, we did not capture the elution profile of $R_1$ in the normal-phase HPLC measurement. About whether the impurity mechanism is responsible for the red afterglow, we reason that if there are impurities in $R_1$ compounds, these possible impurities and $R_1$ molecules would be well separated by MeOBP matrices after low concentration doping. The interaction between these possible impurities and $R_1$ molecules would be negligible. If these possible impurities can give rise to the significant organic afterglow when doped into MeOBP matrices, the excitation spectra would be different from the UV-vis absorption spectra. In the present system, it has been found that the UV-vis absorption and excitation spectra of $R_1$-MeOBP-0.002% and UV-vis spectra of $R_1$ in dichloromethane solutions showed similar maxima (Fig. S27, ESI†). These can rule out the impurity mechanism in the $R_1$-MeOBP system.

Low temperature experiments have been performed to study the afterglow mechanism of $R_1$ molecules, and we have observed that the $R_1$-MeOBP-0.01% powders showed red afterglow excited by a 532 nm laser at 77 K (Fig. S28, ESI†). Delayed emission spectra obtained with multiple temperature points from 77 K to 300 K have been recorded to study the triplet excited states of $R_1$ molecules (Fig. 7A). The delayed emission spectra showed that an emission band ranging from 650 to 900 nm with the emission maxima of 701 nm at 77 K originates from triplet excited states of $R_1$-MeOBP-0.01% samples (Fig. 7A). With the temperature gradually raised to 298 K, the delayed emission spectra nearly coincide with the steady-state emission spectra and show emission bands ranging from 540 nm to 650 nm with an emission maximum of 619 nm. The TTA afterglow mechanism can be ruled out due to the low concentration of $R_1$ molecules doped into MeOBP. Power-dependent delayed fluorescence intensity measurements show a quasi-linear relationship between delayed fluorescence intensity and excitation power (Fig. S29, ESI†). These observations support that the $R_1$-MeOBP-0.01% materials undergo a TADF afterglow mechanism.

The PLQY of $R_1$-MeOBP-0.01% materials was measured to be 43.7% (Fig. S30, ESI†). The excited decay profile of $R_1$-MeOBP-0.01% materials monitored at 619 nm showed tri-exponential decay that can be fit into $\tau_1$ = 2.27 ms (4.63%), $\tau_2$ = 16.73 ms (61.37%), and $\tau_3$ = 48.87 ms (31.00%) (Fig. S31, ESI†). In solid samples, there are heterogeneous microenvironments.[71] The $R_1$-MeOBP system has smaller $S_1$–$T_1$ and $T_1$–$S_0$ energy gaps than that of the $CorBF_2$-MeOBP system. Because of the energy gap law, the photophysical processes would be enhanced in systems with low $S_1$ and $T_1$ levels. Therefore, the effect of heterogeneous microenvironments on excited state decay would be enhanced in the $R_1$-MeOBP system, which possibly results in the triple-exponential decay behavior. The $\tau_1$ part and $\tau_2 + \tau_3$ part can be assigned to prompt fluorescence and delayed fluorescence, respectively. The delayed fluorescence (that is responsible for the afterglow property) quantum yield can be estimated to be 41.6% according to the PLQY of $R_1$-MeOBP-0.01% materials and the emission decay profile with a high proportion of delayed fluorescence ($\tau_2 + \tau_3$, 95.37%) (Fig. S31, ESI†). The $k_{RISC}$ value of the $R_1$-MeOBP-0.02% materials can be estimated from their delayed fluorescence lifetimes ($\tau_{average}$ = 25.41 ms) to be on the order of $10^1$ to $10^2$ $s^{-1}$, which is smaller than those of TADF-type OLED materials. This is the key design of the work to harvest the triplet excited state and maintain a relatively longer emission lifetime.

TD-DFT calculations (ORCA 4.2.1 program with the B3LYP/G functional and def2-TZVP) show that the $S_1$ state of $R_1$ mainly possesses ICT characters from coronene and hemicurcuminoid groups to the dioxaborine ring (Fig. 7B). The $T_1$ state of $R_1$ shows ICT characters from the hemicurcuminoid group to the dioxaborine ring (Fig. 7B). Because of the different excited state nature between $S_1$ and $T_1$ states, the ISC channel of $S_1$–$T_1$ possesses a relatively large SOCME value of 0.60 $cm^{-1}$ (Fig. 7B). The $\Delta E_{ST}$ can be estimated to be 0.23 eV from emission maxima; the $T_1$ level is estimated to be 1.77 eV from phosphorescence maxima at 77 K and the $S_1$ level is estimated to be 2.00 eV from fluorescence maxima at room temperature (Fig. 6 and 7A). The relatively small $\Delta E_{ST}$ together with relatively large SOCME values of 0.60 $cm^{-1}$ for $T_1$–$S_1$ RISC suggests that the $R_1$-MeOBP-0.01% materials possess a moderate $k_{RISC}$ value. Oxygen impacts on afterglow properties have also been taken into consideration, and we have observed an inactive relationship between oxygen and $R_1$-MeOBP-0.01% afterglow materials because of the triplet excited states of $R_1$ molecules well protected by the rigid matrix. The steady-state emission spectra of $R_1$ in different solvents under degassed conditions show an

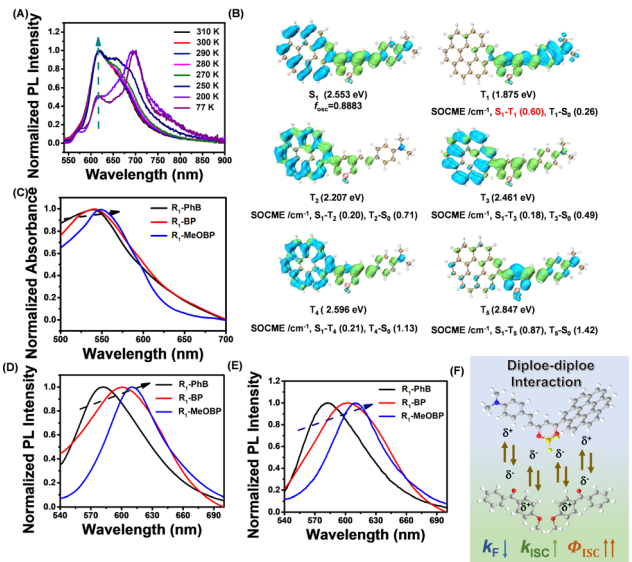


Fig. 7 (A) Variable temperature delayed emission (1 ms delay) spectra of $R_1$-MeOBP-0.01% materials (excited at 520 nm). (B) TD-DFT-calculated electron density difference of the singlet and triplet excited states of $R_1$. TD-DFT calculations were performed on the ORCA 4.2.1 program with the B3LYP/G functional and def2-TZVP(-f) basis set. (C–E) UV-vis spectra (C), steady-state emission spectra (D), and delayed emission spectra (E) of the $R_1$-matrix-0.002% samples. (F) Schematic illustration of dipole–dipole interactions between the $S_1$ states of $R_1$ and MeOBP matrices that can significantly enhance intersystem crossing of $R_1$.

insignificant change after exposure to air conditions (Fig. S32, ESI†). We reason that even though triplet excited states of $R_1$ can be populated *via* intersystem crossing, these can be easily competed by nonradiative decay due to relatively larger $k_{nr}$ and small $k_{RISC}$ in solution. These observations show that oxygen has little influence on the photoluminescence intensity of $R_1$ in a dilute solution.

The UV-vis spectra and room-temperature emission spectra of the $R_1$-matrix have been recorded to study the effects of dipole moments of organic matrices on excited states of the $R_1$ dopant. With the dipole moments of organic matrices increased, the UV-vis spectra of the $R_1$-matrix show slightly red-shifted absorption (Fig. 7C), while the steady-state emission spectra and delayed emission spectra of the $R_1$-matrix revealed red-shifted emission which indicated the reduction of the $S_1$ level of the $R_1$ dopant (Fig. 7D and E). The $S_1$ levels are estimated from their fluorescence emission maxima. Given that most of the organic systems possess $T_1$ levels that are insensitive to the medium or environment, the $T_1$ levels (1.77 eV) are estimated from the phosphorescence maxima of $R_1$-MeOBP-0.01% materials at 77 K. The $\Delta E_{ST}$ values of $R_1$-PhB (2.13 eV), $R_1$-BP (2.06 eV) and $R_1$-MeOBP (2.03 eV) can be estimated from their fluorescence emission maxima and $T_1$ level of $R_1$-MeOBP-0.01% materials. These observations indicate that dipole–dipole interaction between $R_1$ dopants and organic matrices can stabilize single excited states and effectively reduce the $\Delta E_{ST}$.[51] The emission decay profile of prompt fluorescence ($\tau_{PF}$) and delayed fluorescence ($\tau_{DF}$) of the $R_1$-matrix-0.002% samples has also been collected (Fig. S33, ESI†). When BP derivatives have been served as organic matrices, the $\tau_{PF}$ values of $R_1$-BP-0.002% (2.14 ns) and $R_1$-MeOBP-0.002% (2.20 ns) powders have been found to decrease compared to $R_1$-PhB-0.002% (2.26 ns) samples (Fig. S33, ESI†). In the literature,[72–74] $k_F$ of $^1$CT states of $BF_2$bdk compounds has been reported to decrease with the increase of dipole moments of the medium or environment. Since $\tau_{PF}$ is inversely proportional to ($k_F + k_{nr} + k_{ISC}$) of the $^1$CT states, and $k_F$ would decrease with regard to a relatively larger dipole of BP and MeOBP matrices, the decrease of $\tau_{PF}$ in $R_1$ suggests a significant enhancement of $k_{ISC}$ of $R_1$ excited states in MeOBP and BP matrices. These observations suggest that the BP and MeOBP matrices can enhance ISC and decrease the $k_F$ values of $R_1$'s $S_1$ states *via* dipole–dipole interactions, leading to a remarkable population of the triplet excited states of $R_1$, which contributes to excellent afterglow properties of $R_1$-MeOBP-0.002% materials (Fig. 7F). The MeOBP matrices in the present study not only suppress $k_{nr}$ and $k_q$ of dopants' $T_1$ by a rigid crystalline microenvironment but also enhance dopants' ISC and RISC by the dipole effect.[51] This dipole effect in enhancing both ISC and RISC has also been demonstrated by Gillett and coworkers in a very recent study.[75]

Many small-molecule RTP and afterglow materials show poor processability. In the present study, MeOBP matrices have a melting point of 60–63 °C, which allowed them to be constructed into desired objects with the aid of silicone moulds. Cat paw and bone-shaped objects of Cor-MeOBP-0.1%, $CorBF_2$-MeOBP-0.1% and $R_1$-MeOBP-0.01% materials were prepared to exhibit bright yellowish green afterglow with the duration of 30 s, yellow afterglow with the duration of 11 s and red afterglow with the duration of 3 s, respectively (Fig. 8A and Fig. S34, ESI†). Thanks to the low melting point of MeOBP, large-area films have been prepared by sandwiching hot dopants-MeOBP melts between a pair of 10 cm × 10 cm glass plates followed by cooling under ambient conditions where diverse afterglow patterns, such as fancy Christmas, sea turtle and rabbit, were obtained through predesigned masks after the removal of a UV lamp (Fig. 8B and C). The $CorBF_2$-MeOBP-0.01% and $R_1$-MeOBP-0.002% large-area afterglow films also allowed direct writing by 405 nm and 532 nm lasers to exhibit afterglow characters like "9", "1", "3" and "三", respectively (Fig. 8D and E). These suggest that the afterglow films can be employed as afterglow writing panels for data security.

Afterglow materials that can maintain afterglow properties when dispersed in an aqueous medium have shown promising biomedical applications. Generally, the majority of the reported studies showed that organic RTP or afterglow materials lose afterglow properties when being transferred in aqueous dispersions. Here, the afterglow dispersions can be obtained after the molten $CorBF_2$-MeOBP-0.1% droplets were added into an aqueous solution of the Pluronic F127 surfactant at 80 °C under

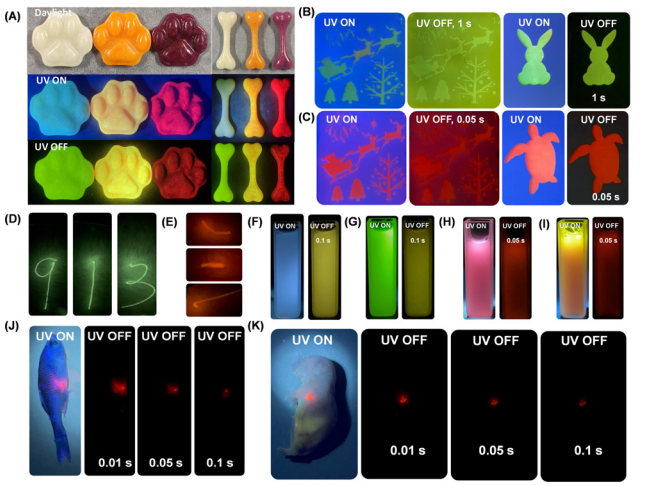


Fig. 8 (A) Photographs of Cor-MeOBP-0.1%, $CorBF_2$-MeOBP-0.1%, and $R_1$-MeOBP-0.01% afterglow objects under daylight, under 365 nm UV light, and after the removal of UV excitation, respectively. (B and C) Photographs of diverse afterglow patterns obtained through predesigned masks on a $CorBF_2$-MeBOP-0.01% film (B) and $R_1$-MeOBP-0.002% film (C) under 365 nm UV light and after the removal of a UV excitation source. (D and E) Afterglow writing of "9", "1", "3" and "三" letters on $CorBF_2$-MeBOP-0.01% (D) and $R_1$-MeOBP-0.002% films (E) (10 cm × 10 cm) by 405 nm and 532 nm laser pointers (50 mW). (F and G) Photographs of the $CorBF_2$-MeOBP-0.1% aqueous dispersion under UV light and after the removal of UV light (F) and the $CorBF_2$-MeOBP-0.1% aqueous dispersion in the presence of fluorescein sodium dyes under 365 nm UV light and after the removal of UV light (G). (H and I) Photographs of the $R_1$-MeOBP-0.02% aqueous dispersion under UV light and after the removal of UV light (H) and the $R_1$-MeOBP-0.02% aqueous dispersion in the presence of rhodamine 6G dyes under 365 nm UV light and after the removal of UV light (I). (J and K) Preliminary bioimaging studies of the aqueous dispersion of $R_1$-MeOBP-0.02% afterglow in fish (J) and mice (K).

 

sonication, then the hot dispersion was immediately frozen by liquid nitrogen and the dispersion recovered to room temperature finally (Fig. S35, ESI†). The $CorBF_2$-MeOBP-0.1% aqueous dispersion has been found to exhibit yellow afterglow after the removal of the UV lamp (Fig. 8F). Strong background fluorescence interference studies have been conducted, and the $CorBF_2$-MeOBP-0.1% aqueous dispersion can maintain afterglow properties in the presence of fluorescein sodium after the removal of the UV excitation source (Fig. 8G). The $R_1$-MeOBP-0.02% aqueous dispersion has also been prepared by following the similar procedure above, which exhibits bright red afterglow after the removal of 365 nm UV light at room temperature (Fig. 8H). The orange fluorescence of the $R_1$-MeOBP-0.02% aqueous dispersion was observed in the presence of rhodamine 6G and only red afterglow was observed after turning off the UV excitation source (Fig. 8I). These observations indicated that both $CorBF_2$-MeOBP-0.1% and $R_1$-MeOBP-0.02% aqueous afterglow dispersions have strong capability to eliminate the interference of strong background fluorescence by making use of their long-lived excited states.

The aqueous suspension of $R_1$-MeOBP-0.02% afterglow materials can be prepared using a grinding method with the assistance of the Pluronic F-127 surfactants, which also can exhibit red afterglow after the removal of an UV excitation source in a dark room. The diameter of the $R_1$-MeOBP-0.02% afterglow dispersion has been found to be 3.5 μm by fluorescence microscope studies (Fig. S36, ESI†). The aqueous afterglow suspension can be readily drawn by a plastic syringe and easily injected into living fish and mice to perform preliminary biological imaging experiments. The biological imaging in the afterglow modes can avoid the interference from background fluorescence and possible scattering from the excitation lights to display very clean backgrounds in the photographs captured (Fig. 8J and K).

## Conclusions

In summary, this study presents the unexpected mechanism of transformation from conventional RTP to efficient TADF-type organic afterglow in coronene-containing dopant-matrix systems. All the room-temperature photophysical studies, variable-temperature delayed emission experiments, the understanding of the ICT molecular design and dopant-matrix strategy to achieve $k_{RISC}$ of $10^0$–$10^2$ $s^{-1}$, TD-DFT calculations, and the discussion to rule out the mechanism of excited state energy transfer and other delayed fluorescence mechanism, as well as the supports from our previously reported TADF-type afterglow systems,[41,62,70] point to the fact that $CorBF_2$-MeOBP and $R_1$-MeOBP systems at room temperature exhibit TADF-type organic afterglow. Despite their serendipitous nature, detailed studies on the underlying photophysics reveal the strong potential of ICT molecular design to achieve intriguing photofunctional materials. ICT molecules possess much smaller singlet–triplet splitting energy than LE systems, which enhances intersystem crossing to a large extent. The incorporation of more than one electron-donating group enriches excited state characters where some $T_n$ states can mediate intersystem crossing. Furthermore, the dipole–dipole interactions between dopants' $^1$ICT states and organic matrices facilitate intersystem crossing. Meanwhile, the organic matrices suppress nonradiative decay and oxygen quenching of $^3$ICT states. This advanced ICT technology would provide a general strategy for constructing efficient TADF-type organic afterglow materials, which features a moderate $k_{RISC}$ to harvest triplet energies, enhance afterglow efficiency, and maintain long afterglow lifetimes (the advantages of TADF-type afterglow materials have been detailed in Text S1, ESI†). The novel design strategy extracted from the present study would have significant impact on the development of high-performance organic afterglow materials and their promising applications in diverse fields.

## Author contributions

Minjian Wu, Jiuyang Li, Ju Huang, Xuepu Wang, Guanming Wang, Xiuzheng Chen, Xun Li, Xuefeng Chen and Shuhui Ding performed the studies. Minjian Wu and Jiuyang Li drafted the manuscript. Hefeng Zhang supervised this work. Kaka Zhang supervised this work and drafted and revised the manuscript. All authors have given approval to the final version of the manuscript.

## Conflicts of interest

There are no conflicts to declare.

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