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Light-driven rotary molecular motors without point chirality: a minimal design†

Jun Wang,‡ Baswanth Oruganti‡ and Bo Durbeej*

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A fundamental requirement for achieving photoinduced unidirectional rotary motion about an olefinic bond in a molecular motor is that the potential energy surface of the excited state is asymmetric with respect to clockwise and counterclockwise rotations. In most available light-driven rotary molecular motors, such asymmetry is guaranteed by the presence of a stereocenter. Here, we present non-adiabatic molecular dynamics simulations based on multi-configurational quantum chemistry to demonstrate that this chiral feature is not essential for inducing unidirectional rotary motion in molecules that incorporate a cyclohexenylidene moiety into a protonated Schiff-base framework. Rather, the simulations show that it is possible to exploit the intrinsic asymmetry of the puckered cyclohexenylidene to control the direction of photoinduced rotation.

Photoisomerization around an olefinic bond in the retinal protonated Schiff-base chromophore of the visual pigment rhodopsin is the primary event of vision.^{1,2} Mimicking this reaction *in vitro* to extract useful mechanical work involves exerting control of the direction – clockwise (CW) or counterclockwise (CCW) – of the photoinduced rotation, usually by means of molecular chirality.^{3–5} Indeed, by making the potential energy surface (PES) of the excited state asymmetric with respect to the two directions, chirality ensures that the photoinduced rotation prefers one direction over the other.⁶ Molecules that are able to accomplish a full 360° rotation through the absorption of UV or visible light are commonly known as light-driven rotary molecular motors.

Ever since their introduction in the late 1990s,³ synthetic light-driven rotary molecular motors developed from sterically overcrowded alkenes have attracted considerable interest.^{4,7–15} In these motors, the asymmetry in the excited-state PES is achieved by the presence of one or two stereocenters and one or two helical motifs that result from a substantial pretwist

(*ca.* 10–25°)⁸ of the isomerizing olefinic bond. In this work, this is referred to as a “helical” pretwist. The unidirectional rotary motion is produced by consecutive photoisomerization and thermal isomerization steps, where the direction of rotation is controlled by the configuration(s) of the stereocenter(s) (*R* or *S*).^{7,9} In two recent studies, however, it has been demonstrated that permanent point chirality is not an absolute requirement for achieving unidirectional rotary motion by an overcrowded alkene-based motor that exhibits a chiral-like feature in the form of a pseudo-asymmetric carbon,¹⁶ or incorporates an axially chiral trityl moiety.¹⁷ Interestingly, unidirectional rotary motion has also been realized or conceived for systems that contain a plane of symmetry, either by exploiting mechanically interlocked architectures¹⁸ or using polarized laser pulses^{19,20} or oscillating electric fields²¹ as the input energy source.

Although the overcrowded alkene-based motor design has shown potential for a variety of applications in nanotechnology,^{13–15} the performance of these motors is restricted in two different ways. First, under ambient conditions, the rotational frequencies are limited by the fact that the thermal isomerizations of the rotary cycles are much slower than the photoisomerizations.^{7,10} Second, the quantum yields (QYs) of the photoisomerizations are limited by the pyramidalization of one of the olefinic carbons during these processes.^{10,22,23} While much effort has been invested to improve the performance of the thermal isomerizations of overcrowded alkenes^{4,7,12,24–26} or to suggest alternative motor designs that avoid thermal steps altogether,^{6,27,28} the second limitation has only recently started to be addressed.^{6,10,11,23,29} Particularly, it has been demonstrated that a new motor design incorporating a protonated or alkylated Schiff base is able to reduce the negative influence of pyramidalization on the photoisomerization QYs.^{23,29} Moreover, for one such design by Olivucci, Léonard, Gindensperger and their co-workers,⁶ it has been found that unidirectional rotary motion around an olefinic bond can be achieved from two photoisomerization steps alone, and that a sizable helical pretwist in the associated *E* and *Z* motor isomers is not essential.⁶ However, point chirality is deemed critical for this capability.⁶

Division of Theoretical Chemistry, IFM, Linköping University, SE-581 83 Linköping, Sweden. E-mail: bodur@ifm.liu.se

† Electronic supplementary information (ESI) available: Computational details, Scheme S1, Fig. S1–S10, Tables S1–S8, description of multimedia files, and Cartesian coordinates for equilibrium geometries. See DOI: 10.1039/c6cp08484b
‡ These two authors contributed equally to this work and are co-first authors.





Scheme 1 The E_{out} , Z_{in} , E_{in} and Z_{out} isomers of motor **1** and relevant dihedral angles.

In this work, we attempt to generalize the Schiff-base motor design by demonstrating, through minimum energy path (MEP) calculations and non-adiabatic molecular dynamics (NAMD) simulations^{30–33} based on multiconfigurational quantum chemistry,³⁴ that this design can also produce unidirectional rotary motion without point chirality. To this end, we present a minimal model of a Schiff-base motor lacking a stereocenter by incorporating a puckerd cyclohexenylidene moiety into a pyrrolinium motif. This system is shown in Scheme 1 and is hereafter referred to as motor **1**. Because of the intrinsic asymmetry of the cyclohexenylidene attributable to its puckerd ring structure, both the E and Z isomers of **1** have two possible orientations of the 5' carbon relative to the approximate plane formed by the other carbons of the ring – outward (out) and inward (in). Thus, there are four different isomers of **1** that can be labeled $1-E_{\text{out}}$, $1-E_{\text{in}}$, $1-Z_{\text{out}}$ and $1-Z_{\text{in}}$, and which form two pairs of axially chiral enantiomers: $1-E_{\text{out}}/1-E_{\text{in}}$ and $1-Z_{\text{out}}/1-Z_{\text{in}}$.

The ground-state (S_0) equilibrium geometries of all isomers of **1** were obtained using the complete active space self-consistent field (CASSCF)³⁵ method in combination with the ANO-L-VDZP basis set, employing an active space of six π -electrons in six orbitals (this basis set and this active space were used for all CASSCF-based calculations of this work; see further the ESI†). The helical pretwist shown by the resulting geometries can be characterized by the $\Delta\omega$ dihedral angle, where ω is the mean of the θ and θ' dihedral angles (see Scheme 1) and $\Delta\omega$ is the signed deviation of ω from the completely untwisted ω values of 0 or $\pm 180^\circ$. Notably, with absolute $\Delta\omega$ values of around 3° only, all isomers exhibit a small helical pretwist, which is also predicted by other methods than CASSCF (see Table S3 of the ESI†).

Starting from the vertical Franck–Condon (FC) points in the optically bright first excited singlet $\pi\pi^*$ state (S_1) of the different isomers of **1**, the $E \rightarrow Z$ and $Z \rightarrow E$ photoisomerizations were first modeled by performing MEP calculations using state-averaged CASSCF (SA-CASSCF). For each geometry obtained along the respective MEP, singlepoint calculations were carried out with complete active space second-order perturbation theory



Fig. 1 Minimum energy paths (MEPs) from the S_1 Franck–Condon (FC) points of the E_{out} (a), Z_{in} (b), E_{in} (c) and Z_{out} (d) isomers of motor **1**. Shown are also the molecular geometries at the FC point and at two additional points along the paths, as well as the corresponding ω dihedral angles.

(CASPT2)³⁶ to account for dynamic electron correlation effects. The resulting MEPs are presented in Fig. 1.

As each system evolves from the FC point populated by a near-UV photon of energy ~ 3.8 eV (see Table S4 of the ESI†), the MEPs in Fig. 1 show pronounced and barrierless torsional motion around the central olefinic bond toward the S_1/S_0 conical intersection (CI) region where the system is expected to decay to the S_0 state. For the $1-E_{\text{out}}$ and $1-Z_{\text{in}}$ isomers, the torsional motion is toward increasing values of ω (here defined as CCW motion), whereas it is toward decreasing values of ω for the $1-E_{\text{in}}$ and $1-Z_{\text{out}}$ isomers (CW motion). Furthermore, S_0 geometry optimizations started from the end points of the MEPs continue the torsional motion set forth in S_1 and yield $1-Z_{\text{in}}$ and $1-E_{\text{out}}$ as the photoproducts of the $1-E_{\text{out}}$ and $1-Z_{\text{in}}$ paths, respectively, and $1-Z_{\text{out}}$ and $1-E_{\text{in}}$ as the photoproducts of the $1-E_{\text{in}}$ and $1-Z_{\text{out}}$ paths, respectively. These results suggest that **1**, despite lacking a stereocenter, produces 360° unidirectional rotary motion. Furthermore, reminiscent of how the direction of rotation of overcrowded alkenes is dictated by the R/S configuration(s) of the stereocenter(s),^{7,9} the direction of rotation of **1** is controlled by the out/in orientation of the 5' carbon, in such a way that the $[1-E_{\text{out}}, 1-Z_{\text{in}}]$ pair of isomers complete a 360° CCW rotation and the $[1-E_{\text{in}}, 1-Z_{\text{out}}]$ pair a 360° CW rotation.

Besides achieving rotary motion without a stereocenter, it is also interesting to note from Table S5 of the ESI† that **1** shows very little pyramidalization of the central olefinic carbons along the MEPs. For each MEP, the α and α' dihedral angles (see Scheme 1) remain smaller than 5° for all geometries between





Fig. 3 Influence of the sign of the initial (at $t = 0$) $\Delta\omega$ value on the direction of photoinduced torsional motion around the central olefinic bond during the first 200 fs of all 200 trajectories run for the E_{out} (a), Z_{in} (b), E_{in} (c) and Z_{out} (d) isomers of motor **1**.

directionality (or efficiency) of the rotary motion produced by the [$1-E_{out}$, $1-Z_{in}$] and [$1-E_{in}$, $1-Z_{out}$] pairs of isomers. The reason why each isomer has such a clear propensity to photoisomerize in one particular direction only is, we believe, that the S_1 PES is markedly asymmetric with respect to the CW and CCW directions (see Fig. S8 of the ESI†). Thereby, the ability of **1** to achieve unidirectional rotary motion is not a consequence of the small positive and negative helical pretwists in the different isomers of **1**, but rather a manifestation of how the out and in orientations of the cyclohexenylidene 5' carbon influence the S_1 PES.

In order for experimental realizations of **1** or its structural analogues to work, it is important that the thermal free-energy barriers separating the E_{out} and E_{in} and the Z_{out} and Z_{in} enantiomers are not too small. Otherwise, it would be difficult to synthesize or isolate an enantiopure out/in isomer, as the isomers have identical energies (see Table S3, ESI†), and similarly challenging to control the direction of photoinduced rotation through the out/in orientation of the 5' carbon. In this light, it is pleasing that neither $1-E_{out} \leftrightarrow 1-E_{in}$ nor $1-Z_{out} \leftrightarrow 1-Z_{in}$ enantiomerizations were observed during the NAMD simulations (a more rigorous assessment of the possible occurrence of such processes would require that quantum dynamics effects are accounted for, which is beyond the scope of this work).

However, to investigate this issue in further detail, the enantiomerization barriers were calculated using both density functional theory (DFT) and second-order Møller–Plesset perturbation theory (MP2) as described in the ESI†. From these calculations, room-temperature barriers of around 25 kJ mol^{-1} were predicted (see Table S7 of the ESI†), whereby efficient control of the rotary cycles of **1** would seem to require low temperatures. Alternatively, it may be possible to increase the enantiomerization barriers by introducing other substituents than hydrogens at the 5' position and/or other positions of the cyclohexenylidene moiety. To test this possibility, all four hydrogens at the 4' and 5' positions were replaced with bulkier methyl groups, with the idea that these could slow down the

enantiomerizations on steric grounds. The resulting species, labeled $2-E_{out}$, $2-Z_{in}$, $2-E_{in}$ and $2-Z_{out}$, are shown in Scheme S1 of the ESI†. Importantly, these species remain axially chiral (*i.e.*, lack a stereocenter).

By subjecting **2** to the same type of DFT and MP2 calculations performed for **1**, it was then found that the added methyl groups raise the enantiomerization barriers by about 20 kJ mol^{-1} (see Table S7, ESI†). Thus, the potential problem with facile enantiomerizations appears circumventable by introducing substituents that could raise the barriers even further. For example, complementary calculations on a tetra-brominated (at the 4' and 5' positions) analogue of motor **1**, labeled **3** in Table S7 (ESI†), predict that the barriers can actually be raised by 50 kJ mol^{-1} relative to **1**. This approach to increase the barriers by such magnitudes can be exploited in established asymmetric synthesis techniques, like dynamic kinetic resolution,^{37,38} to obtain an enantiopure analogue of **1** from a racemic mixture of **1**, using a chiral catalyst.^{37,38}

Finally, by subjecting **2** to MEP calculations and NAMD simulations, results presented in Fig. S9, S10 and Table S8 of the ESI† suggest that producing 360° unidirectional rotary motion through sequential absorption of two photons is not a behavior unique to **1**, but is a behavior that can be attributed to any structural analogue of **1** containing a puckered cyclohexenylidene moiety.

Conclusions

We have demonstrated that light-driven unidirectional rotary motion is achievable by axially chiral molecules that incorporate a cyclohexenylidene moiety into a protonated Schiff-base framework, and of which **1** is a minimal model. Despite lacking a stereocenter, the direction of photoinduced rotation of **1** is controllable by the orientation of the 5' carbon relative to the plane formed by the other carbons of the cyclohexenylidene. Furthermore, the photoisomerizations retain key favorable features, such as short S_1 lifetimes,^{6,23,29} observed in previous studies of Schiff-base motors. It is our hope that the results of this work will stimulate the further development of improved rotary motors by having brought molecules that have no point chirality into play.

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References

- 1 G. Wald, *Nature*, 1968, **219**, 800–807.
- 2 H. G. Khorana, *J. Biol. Chem.*, 1992, **267**, 1–4.
- 3 N. Koumura, R. W. J. Zijlstra, R. A. van Delden, N. Harada and B. L. Feringa, *Nature*, 1999, **401**, 152–155.



- 4 N. Koumura, E. M. Geertsema, M. B. van Gelder, A. Meetsma and B. L. Feringa, *J. Am. Chem. Soc.*, 2002, **124**, 5037–5051.
- 5 M. Guentner, M. Schildhauer, S. Thumser, P. Mayer, D. Stephenson, P. J. Mayer and H. Dube, *Nat. Commun.*, 2015, **6**, 8406.
- 6 G. Marchand, J. Eng, I. Schapiro, A. Valentini, L. M. Frutos, E. Pieri, M. Olivucci, J. Léonard and E. Gindensperger, *J. Phys. Chem. Lett.*, 2015, **6**, 599–604.
- 7 M. M. Pollard, M. Klok, D. Pijper and B. L. Feringa, *Adv. Funct. Mater.*, 2007, **17**, 718–729.
- 8 A. Kazaryan, J. C. M. Kistemaker, L. V. Schäfer, W. R. Browne, B. L. Feringa and M. Filatov, *J. Phys. Chem. A*, 2010, **114**, 5058–5067.
- 9 N. Ruangsupapichat, M. M. Pollard, S. R. Harutyunyan and B. L. Feringa, *Nat. Chem.*, 2011, **3**, 53–60.
- 10 J. Conyard, K. Addison, I. A. Heisler, A. Cnossen, W. R. Browne, B. L. Feringa and S. R. Meech, *Nat. Chem.*, 2012, **4**, 547–551.
- 11 J. Conyard, A. Cnossen, W. R. Browne, B. L. Feringa and S. R. Meech, *J. Am. Chem. Soc.*, 2014, **136**, 9692–9700.
- 12 J. Vachon, G. T. Carroll, M. M. Pollard, E. M. Mes, A. M. Brouwer and B. L. Feringa, *Photochem. Photobiol. Sci.*, 2014, **13**, 241–246.
- 13 P. T. Chiang, J. Mielke, J. Godoy, J. M. Guerrero, L. B. Alemany, C. J. Villagómez, A. Saywell, L. Grill and J. M. Tour, *ACS Nano*, 2012, **6**, 592–597.
- 14 D. J. van Dijken, J. Chen, M. C. A. Stuart, L. Hou and B. L. Feringa, *J. Am. Chem. Soc.*, 2016, **138**, 660–669.
- 15 S. J. Wezenberg, C. M. Croisetu, M. C. A. Stuart and B. L. Feringa, *Chem. Sci.*, 2016, **7**, 4341–4346.
- 16 J. C. M. Kistemaker, P. Štacko, J. Visser and B. L. Feringa, *Nat. Chem.*, 2015, **7**, 890–896.
- 17 S. C. Everhart, U. K. Jayasundara, H. Kim, R. Procúpez-Schirbu, W. A. Stanbery, C. H. Mishler, B. J. Frost, J. I. Cline and T. W. Bell, *Chem. – Eur. J.*, 2016, **22**, 11291–11302.
- 18 D. A. Leigh, J. K. Y. Wong, F. Dehez and F. Zerbetto, *Nature*, 2003, **424**, 174–179.
- 19 G. Pérez-Hernández, A. Pelzer, L. González and T. Seideman, *New J. Phys.*, 2010, **12**, 075007.
- 20 J. Zhang, A. P. Sergeeva, M. Sparta and A. N. Alexandrova, *Angew. Chem., Int. Ed.*, 2012, **51**, 8512–8515.
- 21 D. Horinek and J. Michl, *Proc. Natl. Acad. Sci. U. S. A.*, 2005, **102**, 14175–14180.
- 22 A. Kazaryan, Z. Lan, L. V. Schäfer, W. Thiel and M. Filatov, *J. Chem. Theory Comput.*, 2011, **7**, 2189–2199.
- 23 M. Filatov and M. Olivucci, *J. Org. Chem.*, 2014, **79**, 3587–3600.
- 24 G. Pérez-Hernández and L. González, *Phys. Chem. Chem. Phys.*, 2010, **12**, 12279–12289.
- 25 B. Oruganti, C. Fang and B. Durbeej, *Phys. Chem. Chem. Phys.*, 2015, **17**, 21740–21751.
- 26 B. Oruganti, J. Wang and B. Durbeej, *ChemPhysChem*, 2016, **17**, 3399–3408.
- 27 C. García-Iriepa, M. Marazzi, F. Zapata, A. Valentini, D. Sampedro and L. M. Frutos, *J. Phys. Chem. Lett.*, 2013, **4**, 1389–1396.
- 28 Y. Amatatsu, *J. Phys. Chem. A*, 2012, **116**, 10182–10193.
- 29 A. Nikiforov, J. A. Gamez, W. Thiel and M. Filatov, *J. Phys. Chem. Lett.*, 2016, **7**, 105–110.
- 30 G. Groenhof, M. Bouxin-Cademartory, B. Hess, S. P. de Visser, H. J. C. Berendsen, M. Olivucci, A. E. Mark and M. A. Robb, *J. Am. Chem. Soc.*, 2004, **126**, 4228–4233.
- 31 M. Barbatti, *WIREs Comput. Mol. Sci.*, 2011, **1**, 620–633.
- 32 I. Fdez. Galván, M. G. Delcey, T. B. Pedersen, F. Aquilante and R. Lindh, *J. Chem. Theory Comput.*, 2016, **12**, 3636–3653.
- 33 L. Liu, J. Liu and T. J. Martinez, *J. Phys. Chem. B*, 2016, **120**, 1940–1949.
- 34 F. Aquilante, J. Autschbach, R. K. Carlson, L. F. Chibotaru, M. G. Delcey, L. De Vico, I. Fdez. Galván, N. Ferré, L. M. Frutos, L. Gagliardi, M. Garavelli, A. Giussani, C. E. Hoyer, G. L. Manni, H. Lischka, D. Ma, P.-Å. Malmqvist, T. Müller, A. Nenov, M. Olivucci, T. B. Pedersen, D. Peng, F. Plasser, B. Pritchard, M. Reiher, I. Rivalta, I. Schapiro, J. Segarra-Martí, M. Stenrup, D. G. Truhlar, L. Ungur, A. Valentini, S. Vancoillie, V. Veryazov, V. P. Vysotskiy, O. Weingart, F. Zapata and R. Lindh, *J. Comput. Chem.*, 2016, **37**, 506–541.
- 35 B. O. Roos, P. R. Taylor and P. E. M. Siegbahn, *Chem. Phys.*, 1980, **48**, 157–173.
- 36 K. Andersson, P.-Å. Malmqvist and B. O. Roos, *J. Chem. Phys.*, 1992, **96**, 1218–1226.
- 37 J. L. Gustafson, D. Lim and S. J. Miller, *Science*, 2010, **328**, 1251–1255.
- 38 H. Pellissier, *Tetrahedron*, 2011, **67**, 3769–3802.

