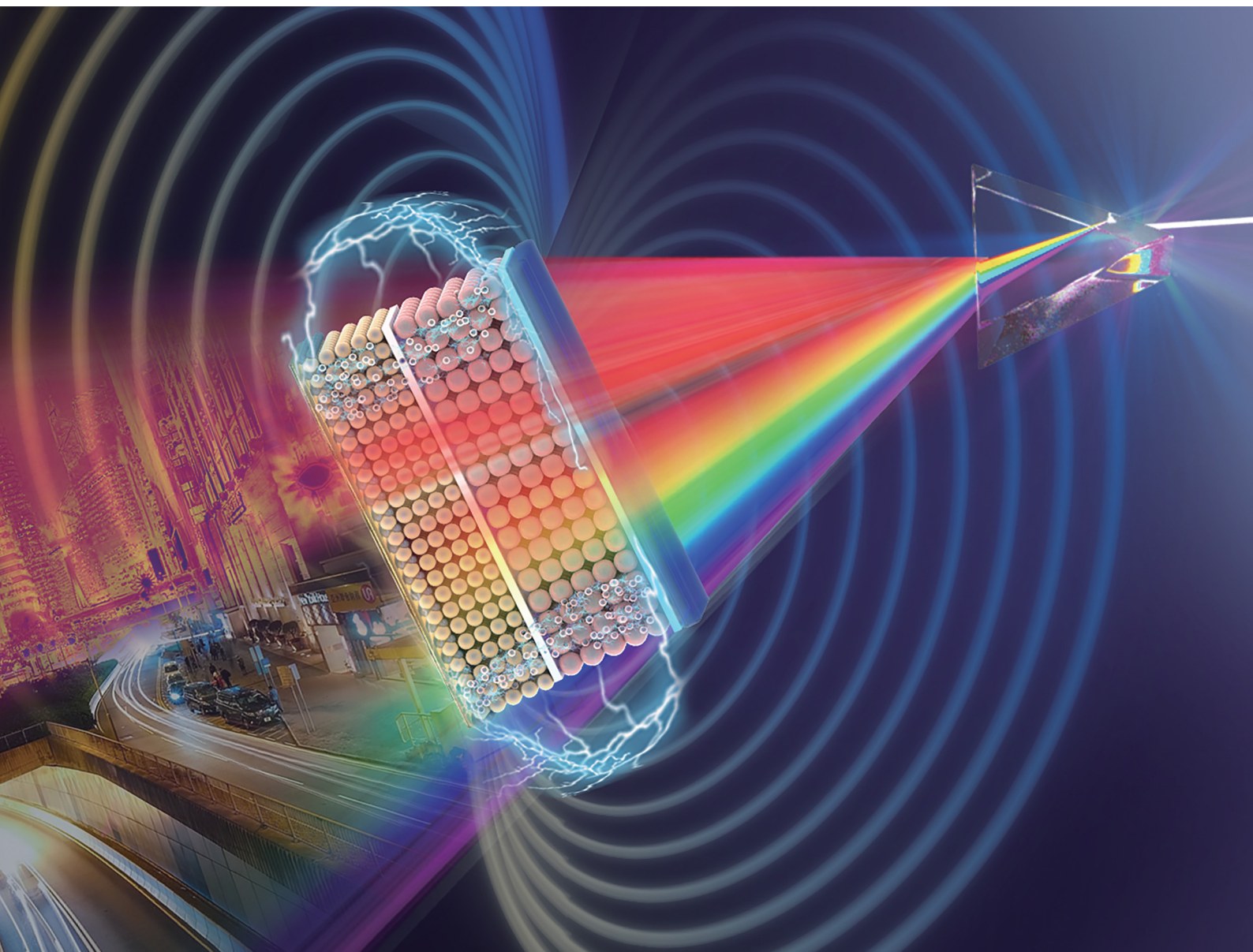


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Band-engineered dual-band visible and short-wave infrared photodetector with metal chalcogenide colloidal quantum dots†

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Dual-band photodetectors have attracted significant attention because of their potential in optical communication, biochemical detection, and environmental monitoring. Colloidal quantum dots (CQDs) are promising materials for dual-band photodetectors because of their size-tunable bandgaps with wide spectral tunability and easy solution processability. Here, we propose a dual-band visible (VIS) and short-wave infrared (SWIR) detector based on metal chalcogenide CQDs (HgTe and CdTe CQDs). An n-type ZnO layer between the CdTe and HgTe layers is introduced as a hole-blocking layer to prevent hole injection from the different sensing layers, enabling the detector to switch between the VIS and SWIR modes by changing the polarity and magnitude of the bias voltage. The device exhibits an excellent performance with high responsivities of 0.5 and 1.1 A W⁻¹ for the two bands that peak at 700 and 2100 nm, respectively. The detectivity of the device can reach 1.1 × 10¹¹ Jones at +3 V (VIS) and 4.5 × 10¹¹ Jones at -2 V (SWIR). In addition, VIS/SWIR dual-band imaging is realized, which provides more comprehensive object information than single-band detectors.

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Xin Tang

spectral/hyperspectral infrared detectors, colloidal quantum-dot focal plane array imagers, and functional photonic nanostructures.

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Introduction

Visible imaging can offer clear and intuitive details of objects and is consistent with the human visual system.^{1–3} However, the narrow spectral range in visible light limits the comprehensive acquisition of object information. Beyond the visible band, infrared can provide important optical information on, for example, the molecular composition and thermal distribution.⁴ Dual-band visible and infrared photodetectors (PDs) that can process signals of visible and infrared wavebands afford more accurate and comprehensive images of detected objects than obtained *via* single-band detection, which improves the detection and identification ability of the system.^{1,2,5,6}

The mature semiconductors such as silicon (visible, VIS, 400–700 nm), InGaAs (short-wave infrared, SWIR, 1.5–1.7 μm), and HgCdTe (mid-infrared, MIR, 3–12 μm) are the primary optoelectronic material systems that cover the broadband sensing range.^{4,7,8} However, when integrating materials with different energy gaps into the same PD pixel, the lattice mismatch and incompatible preparation processes hinder the realization of visible/infrared dual-band PDs.

Emerging materials, such as III-nitride-based semiconducting materials, two-dimensional materials, organic materials, and perovskite materials, have received a lot of attention in recent years.^{9–19} Among the III-nitride family, InN materials are attractive for optoelectronic device applications due to their

suitable bandgap and high electron mobility.¹⁹ Ali *et al.* fabricated the high-quality InN/GaN heterointerface and obtained high-performance broadband PDs.¹⁷ Two-dimensional materials with excellent transparency, extraordinary flexibility, high conductivity, and direct bandgaps have been successfully used not only as electrodes but also as active layers in optoelectronic devices.¹⁸ Zhang *et al.* employed two-dimensional materials of Ta₂NiSe₅ and GaSe to construct a VIS-near infrared (NIR) dual-band photodetector with a heterogeneous structure.⁵ Organic photovoltaic devices employing inexpensive, easily regulated polymer or small-molecule organic materials have shown promising optoelectronic applications. Lan *et al.* reported a dual-band organic PD with a P3HT:PC₇₀BM visible-light absorber/optical spacer/P3HT:PTB7-Th:PC₇₀BM NIR light absorber configuration, which can operate in the NIR and visible spectral ranges.¹ Perovskite materials exhibit superior physical properties such as large optical absorptions, tunable bandgaps, and high charge-carrier mobilities.²⁰ Huang *et al.* proposed a dual-band PD with perovskite materials of MAPbBr₃ and MAPbI₃ as the photoactive layers for potential applications in visible-light communication.¹³

However, the operating band range of these dual-band PDs is limited to the visible (380–750 nm) and NIR (0.75–1.5 μm) regions due to the bandgaps of the materials, while the SWIR (1.5–2.5 μm) spectral region exhibits unique potential for spectrometers, night vision, and pollution detection.^{4,8} It is necessary to explore other material systems to widen the band range of dual-band PDs to the SWIR region.

Colloidal quantum dots (CQDs) are promising materials for the development of high-performance optoelectronic devices in wide spectral ranges because of their unique features, such as size-tunable bandgaps and easy solution processability.^{21–28} Among these colloidal nanomaterials, metal chalcogenide CQDs, such as CdX (X = Te, Se, S) and HgTe CQDs, have been verified with wide spectral ranges that cover the visible region to the long-wave infrared (LWIR, 8–12 μm), and competitive single-band device performances with a detectivity of up to 10¹¹ Jones has been demonstrated.^{29–35} Based on quantum confinement, the size of the quantum dots tunes their energy gap, while the absolute energy levels of the valence and conduction bands are mediated by the energy levels of the bulk material and the surrounding matrix. The valence bands for the II–VI group bulk will be of the same order as the p-orbital energies of the anions.³⁶ As a result, CdTe and HgTe CQDs are chosen for better band alignment.

In this work, we proposed a visible and SWIR dual-band detector based on CdTe and HgTe CQDs. The dual-band detector is arranged in a p–n–p back-to-back photodiode configuration. Through energy band design, an interface layer with a large band-edge offset between the CdTe and HgTe layers can be introduced in the detector, which is beneficial for selective charge extraction. Thus, an n-type ZnO layer between the CdTe and HgTe layers is presented as the electron-transporting layer and hole-blocking layer to prevent hole injection from different sensing layers, which plays a critical role in the bias-switchable spectral response. Well-engineered

dual-band PDs can switch between the visible and SWIR modes upon changing the polarity and magnitude of the bias. The detectivity of the device can reach 1.1 × 10¹¹ and 4.5 × 10¹¹ Jones for the visible and SWIR modes, respectively. In addition, VIS/SWIR dual-band imaging is realized, providing more comprehensive object information than is achievable using single-band PDs.

Experimental section

Synthesis of materials

Synthesis of CdTe CQDs. A TBP–Te (tributylphosphine–tellurium) solution was prepared by dissolving 10 g Te powder in 111 ml TBP using a glove box at room temperature until the solution turned bright yellow. 1.9 g (15 mmol) CdO (cadmium oxide) was put into a three-neck flask. Then, 18.8 ml OA (oleic acid) and 20.3 ml ODE (octadecenoic acid) were added, and a magnetic stirring bar was added. The reaction system was put under a vacuum and then purged with argon gas three times. Then the reaction mixture was heated to 220 °C until the solution became clear. The temperature was set to 270 °C (thermal equilibrium for half an hour), and 9.6 ml TBP–Te solution was quickly injected. The three-neck flask was removed from the heating mantle and cooled using an air gun, followed by a water bath for cooling. The obtained solution was precipitated with toluene and absolute ethanol (toluene: absolute ethanol = 1:2) with centrifugation. The precipitated sediment was redissolved by adding hexane and then placed in the glove box for storage.

Synthesis of HgTe CQDs. In the glovebox, 27 mg (0.1 mmol) HgCl₂ (mercuric chloride) was dissolved in 4 ml OAM (oleylamine) at 100 °C. A quenching solution was prepared by adding 0.5 ml TOP (trioctylphosphine) in 4 ml TCE (tetrachloroethylene) and stored in the refrigerator. A TOPTe solution was prepared by dissolving 0.1 mmol of Te powder and 0.1 ml TOP using a nitrogen glovebox at room temperature to form a bright yellow solution. 0.1 ml of TOPTe was quickly injected into the HgCl₂ solution at 80 °C. After 20 s, the color of the solution changed to black, and the reaction lasted for 4 min. After 4 min, the quenching solution was added to cool down the reaction mixture.

Synthesis of Ag₂Te nanocrystals. 34 mg AgNO₃ (silver nitrate) was dissolved in 5 ml OAM and 0.5 ml OA and heated to 70 °C for 30 min. After 30 min, 0.5 ml TOP was injected and heated rapidly to 160 °C for 40 min. The solution changed from clear to orange. 0.1 ml TOPTe was injected, and the solution quickly turned black. After 10 min, the solution was removed from the glove box, and the reaction vessel was placed in a water bath to end the synthesis.

Device fabrication

A 100-nm-thick ITO (indium tin oxide) film was deposited on a glass substrate *via* sputtering and annealed at 300 °C for 15 min. Before drop-casting of the colloidal solution, MPTS (3-mercaptopropyltrimethoxysilane) was used to treat the substrate

for 1 min to increase the adhesion between the substrate and the CQDs. The CQDs can be stacked layer by layer *via* spin coating. Firstly, the visible CdTe CQD solution was spin-coated layer-by-layer until the thickness reached about 400 nm. Each layer was placed on a 150 °C hot plate for 1 min to remove excess water, then immersed in a saturated solution of CdCl₂ (cadmium chloride) dissolved in methanol at 60 °C for 15 s. IPA (isopropanol) was used to remove the excess CdCl₂, and the films were finally placed on a 350 °C hot plate for 20 s. 50 μl ZnO (20 mg ml⁻¹) was spin-coated on top of the CdTe film at 3000 rpm for 60 s, and it was then annealed at an 80 °C hot plate for 10 min. The SWIR HgTe CQD solution was spin-coated on top of the ZnO film with a thickness of 400 nm. Each layer of the spin-coated HgTe CQDs requires ligand exchange with an EDT (ethanedithiol)/HCl (hydrochloric acid)/IPA (1:1:50 by volume) solution for 10 s. A layer of Ag₂Te nanocrystals was spin-coated on the HgTe CQDs as the p-dopant at 3000 rpm for 60 s, followed by HgCl₂ (10 mmol L⁻¹) treatment for 10 s. Finally, a 40-nm-thick Au layer was deposited as the top contact.

Characterization

Room-temperature UV-Vis absorption spectra were measured using a V-770 UV-Vis-NIR spectrophotometer (Jasco). The absorption spectra of the CdTe CQDs and HgTe CQDs were measured in solution and on the deposited thin films, respectively. Transmittance electron microscopy (TEM) images were obtained using a Tecnai Spirit instrument with a driving voltage

of 120 kV. COMSOL Multiphysics software was used for the finite element analysis. The response spectra of the device were measured using a monochromator. The photocurrent of the device was measured using a Keithley 2602B source meter, and a Femto DLPCA-100 variable gain low-noise transimpedance amplifier in series with a DLPVA-100s low-noise preamplifier. The performance of the devices was characterized at room temperature.

Results and discussion

Device structure and operation principle

The structure and energy bands of the dual-band visible and SWIR PD are shown in Fig. 1a and b, respectively. The device consists of a visible photodiode and a SWIR photodiode in a p-n-p configuration. Previous studies have indicated that Ag₂Te nanocrystals serve as good p-dopants for the HgTe CQDs by diffusing into neighboring HgTe CQDs as Ag⁺ ions through the solid-state cation exchange method.³³ The ZnO layer has previously been used as an n-type layer in both CdTe and HgTe CQD photovoltaic devices.^{36,37} The CdTe CQDs and the HgTe CQDs are used as the visible absorption layer and the SWIR sensing layer, respectively. The energy bands of the CdTe and HgTe CQDs render them well combined into the same detector. The SWIR photodiode is above the VIS photodiode, and the two photodiodes share a mutual n-type ZnO layer. The ZnO layer acts not only as the electron-transporting layer but also as the

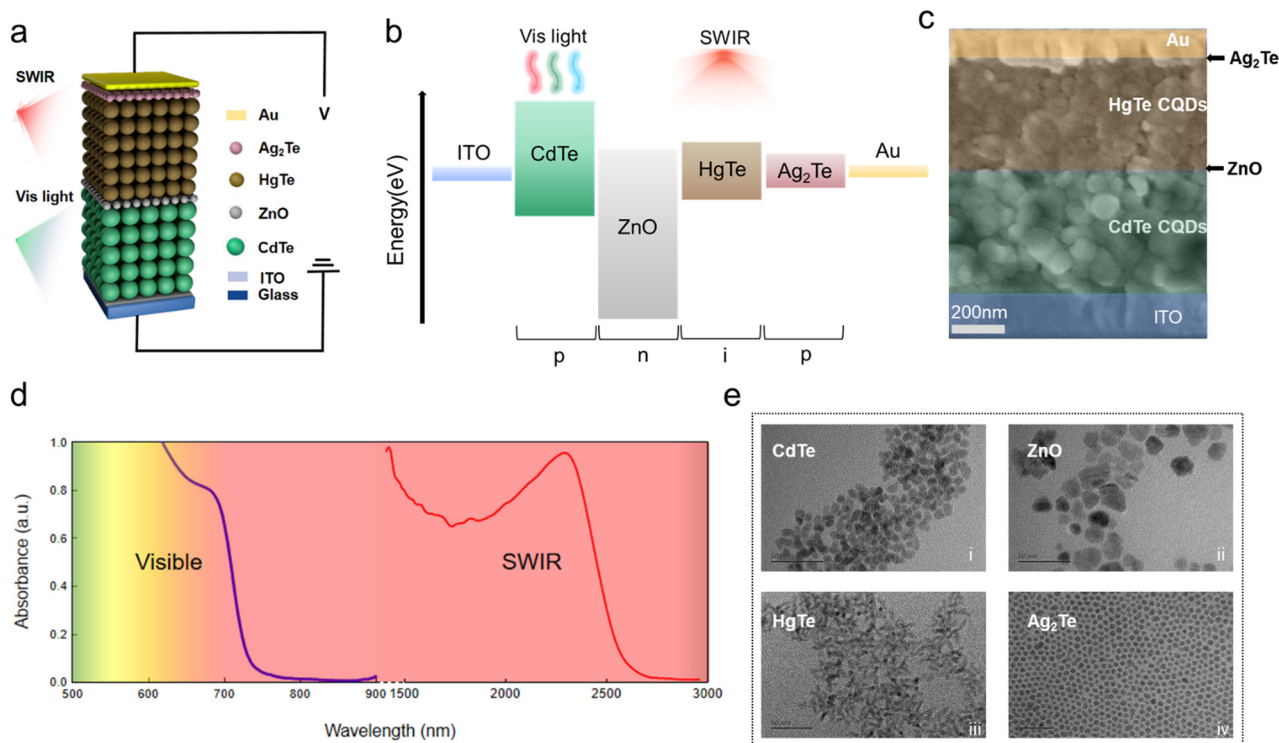


Fig. 1 Architecture of the dual-band visible and SWIR photodetector. (a) Structure diagram of the dual-band device. (b) Energy band structure diagram of the dual-band device. (c) Cross-sectional SEM image of the dual-band device. (d) Absorption spectra of CdTe CQDs and HgTe CQDs. (e) TEM images of the CdTe, ZnO, HgTe, and Ag₂Te materials.

hole-blocking layer to prevent hole injection from different visible and SWIR sensing layers, which plays a critical role in the bias-switchable spectral response. When a positive voltage is applied to the Au electrode, the VIS region of the CdTe/ZnO heterojunction is in the reverse bias, and the SWIR region is in the forward bias region. At this point, the visible channel starts to work. Conversely, when a negative voltage is applied to the Au electrode, the SWIR channel of the ZnO/HgTe/Ag₂Te heterojunction begins to operate.⁸ The cross-sectional scanning electron microscopy (SEM) image of the fabricated dual-band device is shown in Fig. 1c, and the device exhibits an apparent layered structure.

The absorption spectra of the CdTe CQDs and HgTe CQDs are shown in Fig. 1d. The cutoff wavelengths of the synthesized CdTe and HgTe CQDs are about 700 and 2300 nm, which are in the VIS and SWIR band range, respectively. Transmission electron microscopy (TEM) images of the CdTe, ZnO, HgTe, and Ag₂Te materials are shown in Fig. 1e.

Characterization of the device performance

First, we fabricated CdTe and HgTe single-band devices. The fabrication process of the CdTe/ZnO single-band devices is presented in Supplement 2 (S2) (ESI†). The ligand exchange process with pyridine of the CdTe CQDs is essential for fabricating high-performance devices. The photocurrent of the CdTe

single-band devices is increased by about ten times at zero bias after the ligand exchange process (Fig. S1, ESI†), which is attributed to the improved CdTe carrier mobility.^{38,39} The CdTe/ZnO single-band devices are capable of detecting visible light and exhibited an excellent performance under 700 nm visible light illumination with an optical power of $8.0 \times 10^{-4} \text{ W cm}^{-2}$ (Fig. 2a(i) and Fig. S2, ESI†). The detailed preparation steps of the ZnO/HgTe/Ag₂Te single-band devices are shown in S3 (ESI†).⁴⁰ With the presence of the ZnO layer, the rectifying characteristics of the devices are greatly enhanced compared with the HgTe CQDs photodiodes without ZnO, which is similar to previous studies (Fig. S3, ESI†).⁴¹ The dark current density of the ZnO/HgTe/Ag₂Te devices is significantly reduced by nearly eight times (Fig. S3, ESI†). The ZnO/HgTe/Ag₂Te single-band devices can detect SWIR with a cutoff wavelength of 2300 nm at room temperature, and under illumination with a calibrated 600 °C blackbody they demonstrate an excellent performance of $2.2 \times 10^{-3} \text{ W cm}^{-2}$ (Fig. 2a(ii) and Fig. S4, ESI†).

The dual-band visible and SWIR detector arranged in a CdTe/ZnO/HgTe/Ag₂Te back-to-back photodiode configuration were fabricated. The fabrication process is detailed in the Experimental section, and the CdTe also underwent ligand exchange. The bias-tunable spectral response curves of the dual-band devices were characterized using the measurement

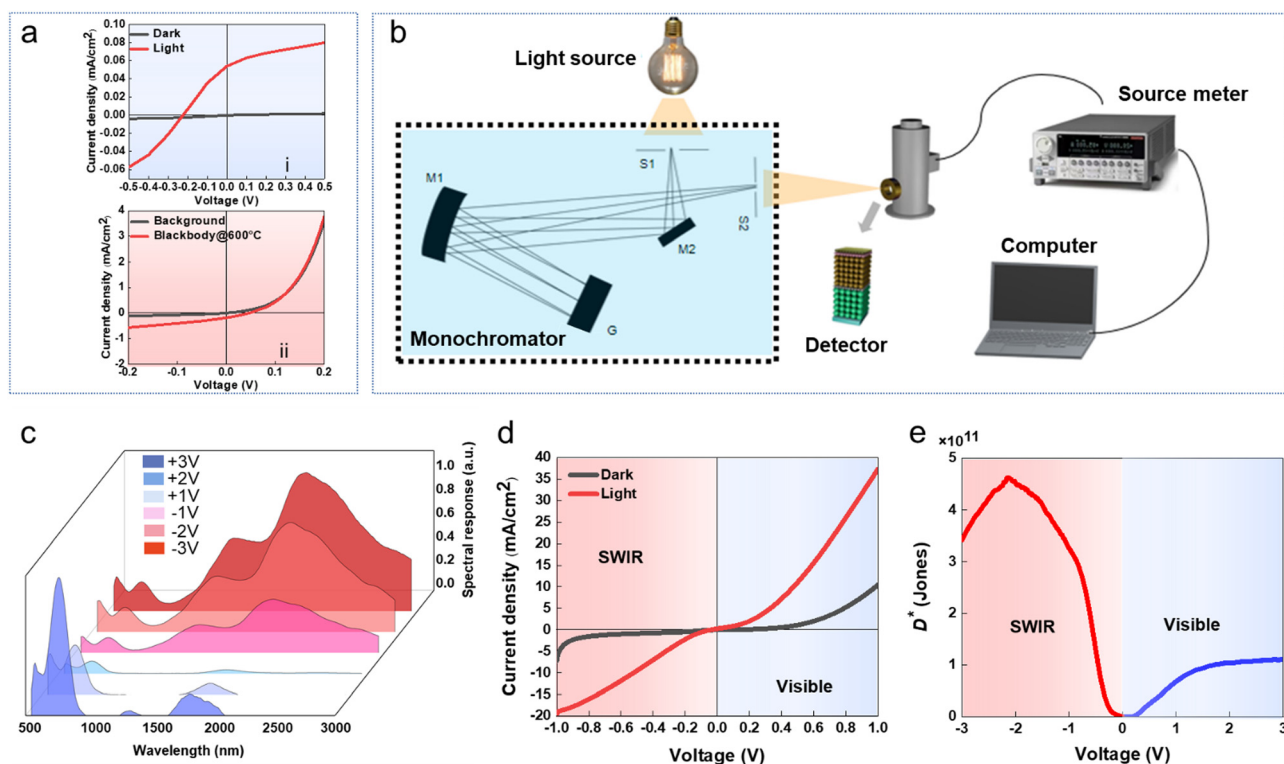


Fig. 2 Performance of dual-band visible and SWIR photodetectors. (a) Current density–voltage curves of CdTe/ZnO single-band device and ZnO/HgTe/Ag₂Te single-band device. (b) Illustration of the performance characterization system. (c) Response spectra of the dual-band device under positive and negative bias. (d) Current density–voltage curves of the dual-band device. (e) Specific detectivity of the dual-band device under positive and negative bias.

system illustrated in Fig. 2b. The detectors exhibit bias-dependent dual-band visible and SWIR spectral responses, as shown in Fig. 2c. By changing the bias voltage from -3 V to $+3$ V, the spectral response gradually switches from SWIR mode to visible mode. The response amplitude increases with the intensity of the bias voltage. Crosstalk signals exist, as shown in Fig. 2c and Fig. S5 (ESI[†]), but the magnitudes of these signals are tiny and can be ignored.

The current density–voltage curve of the dual-band detector using a tungsten lamp as the light source is shown in Fig. 2d. The black line is the dark current curve, and the dark current is suppressed for both positive and negative bias voltages due to the back-to-back diode structure. The red line presents the current density–voltage curve under illumination, and the dual-band detector exhibits a significant photocurrent at both positive and negative bias voltages. Both the responsivity \mathfrak{R} and the specific detectivity D^* are important indicators of a detector. The calculation process for the responsivity is as follows:

$$\mathfrak{R} = \frac{I_{\text{ph}}}{P} \quad (1)$$

where I_{ph} is the photocurrent, and P is the input optical power. The fabricated dual-band devices exhibit a high responsivity, reaching 0.5 and 1.1 A W^{-1} for the two bands that peak at 700 and 2100 nm, respectively (Fig. S5, ESI[†]). The detectivity is a

measure of the signal-to-noise ratio, which is given by the following equation:

$$D^* = \frac{\sqrt{A \cdot \Delta f}}{I_n} \mathfrak{R} \quad (2)$$

where A is the sensing area, Δf is the bandwidth, I_n is the noise spectral density, and \mathfrak{R} is the responsivity. The noise current of the dual-band devices is low, at 5×10^{-13} $\text{A Hz}^{-1/2}$ at 3 V and 3×10^{-13} $\text{A Hz}^{-1/2}$ at -2 V. As a result, the dual-band devices display a high detectivity of nearly 1.1×10^{11} at 3 V in the visible mode and 4.5×10^{11} Jones at -2 V in the SWIR mode, which is shown in Fig. 2e. The performance of the CdTe/HgTe CQD dual-band devices is comparable to that of dual-band devices based on organic materials, two-dimensional materials, and perovskite materials, as shown in Table 1.

Device simulation

To investigate the physics and working mechanism of the dual-band device, finite element analysis was carried out using the wave optical module and the semiconductor module. The structure diagram of the simulation model is shown in Fig. 3a. The CdTe (400 nm), ZnO (50 nm), and HgTe (400 nm) are superimposed vertically, and the periodic boundary condition is used.⁴⁸ Bulk CdTe has a bandgap of 1.45 eV.^{49,50} Besides, bulk HgTe possesses a zero gap in theory.⁵¹ The CQDs have different energy levels in their bulk forms. In our simulation,

Table 1 Comparison of the performance of dual-band PDs

| Active materials | Wavelength (nm) | Voltage (V) | Incident light intensity (W cm^{-2}) | Responsivity (A W^{-1}) | Noise current ($\text{A Hz}^{-1/2}$) | Detectivity (Jones) | Ref. |
|---|-----------------|-----------------|---|--|---|---|-----------|
| Ta ₂ NiSe ₃ /GaSe | 520/1550 | — | $1.4 \times 10^{-8}/$ 3.0×10^{-8} | 4.8/0.15 | — | 1.1×10^9 | 5 |
| SnO ₂ /CsPbBr ₃ | 320/530 | 3/3 | — | $3.47 \times 10^{-4}/2.93 \times$ 10^{-4} | — | $1.6 \times 10^{10}/1.2 \times 10^{10}$ | 42 |
| N-InAsSb/ N-GaSb | 1700/3400 | 0.3–0.8 | — | 0.3/40 | — | $1.5 \times 10^{11}/1.4 \times 10^9$ | 43 |
| InP/InGaAs/ GaAsSb | 1640/2000 | $-0.1/$ -1 | — | 0.57/0.22 | — | $2.63 \times 10^{11}/1.96 \times 10^9$ | 44 |
| MAPbI ₃ | 405/1064 | 1/0.01 | $3.31 \times 10^{-2}/$ 3×10^{-2} | 680/1.74 | Noise equivalent power: 3.2 pW $\text{Hz}^{-1/2}/$ 1.9 nW $\text{Hz}^{-1/2}$ | $1.2 \times 10^9/5.4 \times 10^7$ | 15 |
| MAPbI ₃ /IEICO | 650/820 | — | $5.6 \times 10^{-3}/$ 5.6×10^{-3} | 0.28/0.14 | 4×10^{-14} | $1.45 \times 10^{12}/7.37 \times 10^{11}$ | 3 |
| InSe/Au | 365/685 | 0 | $1.57 \times 10^{-2}/$ 1.31×10^{-1} | 0.369/0.244 | — | $5 \times 10^{12}/1 \times 10^{11}$ | 45 |
| ZnCdSe/ ZnCdMgSe | 4800/7600 | 5/1 | — | 0.011/0.007 | $\sim 2 \times 10^{-14}/\sim 1 \times 10^{-13}$ | $2 \times 10^8/2 \times 10^7$ | 47 |
| Ge/Si | 1000/1550 | 1/ -0.5 | $10^{-6}-10^{-2}/$ $10^{-5}-10^{-1}$ | 0.33/0.63 | — | $7 \times 10^{11}/2 \times 10^{10}$ | 6 |
| Cs ₃ Cu ₂ I ₅ / PdTe ₂ /Ge | 265/1550 | — | $6.8 \times 10^{-5}/$ 8.0×10^{-5} | 0.0257/0.6941 (0 V) 0.3629/0.6984 (-0.1 V) 0.5758/0.7031 (-0.2 V) 0.7442/0.7072 (-0.3 V) 0.7325/0.7125 (-0.4 V) | 1.01×10^{-12} (0 V) 4.2×10^{-10} (-0.1 V) 4.3×10^{-10} (-0.2 V) 7.8×10^{-10} (-0.3 V) | $5.08 \times 10^9/1.37 \times 10^{11}$ (0 V) $1.73 \times 10^8/3.33 \times 10^8$ (-0.1 V) $1.83 \times 10^8/2.23 \times 10^8$ (-0.2 V) $1.91 \times 10^8/1.81 \times 10^8$ (-0.3 V) | 46 |
| CdTe/ZnO/ HgTe | 700/2100 | 3/ -2 | $8.0 \times 10^{-4}/$ 2.2×10^{-3} | 0.5/0.9 (1.1 at -3 V) | $5 \times 10^{-13}/3 \times 10^{-13}$ | $1.1 \times 10^{11}/4.5 \times 10^{11}$ | This work |

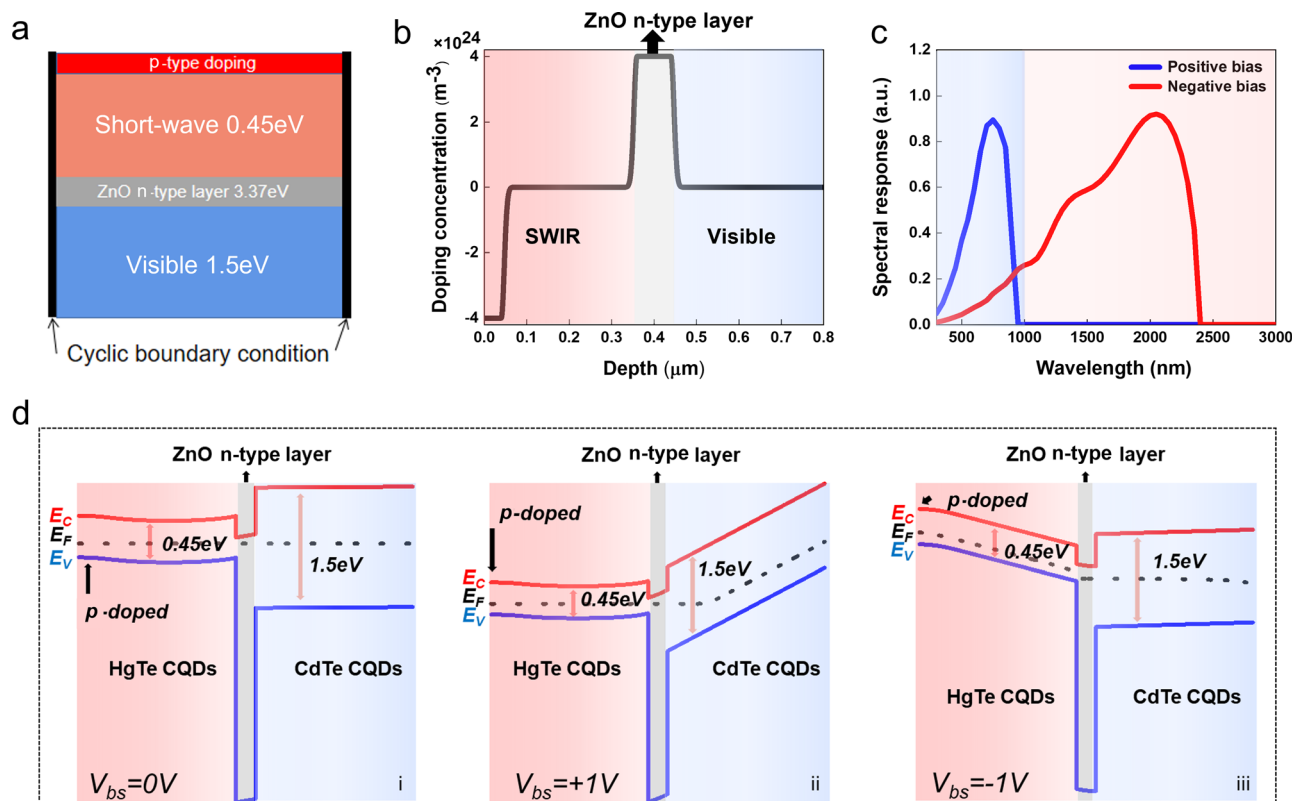


Fig. 3 Simulation of the dual-band visible and SWIR photodetector. (a) Simulation model structure of the dual-band device. (b) Distribution of the doping concentration. (c) Simulated spectral responses of the dual-band device under negative and positive bias voltages. (d) Simulated energy band diagram of the dual-band device under zero (i), positive (ii), and negative (iii) bias voltages.

the energy bandgaps of the CdTe CQDs and the HgTe CQDs are 1.5 and 0.45 eV, respectively.^{52–60} The refractive index of the CQD film is $2.3 + 0.1i$.^{61,62} Besides, the energy bandgap of ZnO is 3.3 eV in the simulation.^{52–56} Layers of p-type Ag₂Te and n-type ZnO are introduced at the top of HgTe and at the interface between HgTe and CdTe, respectively. Fig. 3b shows the doping concentration distribution of the dual-band devices. The simulated spectral response of the dual-band detector under different polarity bias voltages is shown in Fig. 3c. The device exhibits an obvious visible band response under a positive bias and an SWIR response under a negative bias, similar to the experimental results (Fig. 2c).

To understand the bias-switchable spectral response of the dual-band visible and SWIR PD, the energy band structure of the devices at zero, positive and negative bias voltages was simulated and analyzed (Fig. 3d). Under a positive bias, the visible CdTe/ZnO heterojunction is in the reverse working region. The CdTe CQD layer generates electron–hole pairs under visible light, and electrons and holes are transferred to the adjacent ZnO layer and the ITO electrode, respectively. The ZnO layer is beneficial for fast electron transport. More importantly, a large band-edge offset between the ZnO layer and the CdTe layer prevents hole injection into the HgTe layer. Only the CdTe visible-light-absorbing layer generates photogenerated carriers under a positive bias, inducing CdTe band bending. Under a negative bias, the SWIR ZnO/HgTe/Ag₂Te heterojunction begins

to operate in the reverse region. The Fermi level near the HgTe/Ag₂Te interface tends to the valence band due to the introduction of p-type Ag₂Te doping, resulting in a lower energy barrier for hole transport. The photogenerated charges are selectively extracted from the HgTe CQD layer due to the hole barrier effect of the ZnO layer. The energy bend of the HgTe layer is bent, caused by the high density of SWIR light-generated carriers. The above mechanism enables the detector to effectively access two separate visible and SWIR channels through simple bias regulation.

Dual-band imaging

The visible and SWIR dual-mode imaging of the detector is demonstrated using a single-pixel scanning imaging system with a tungsten lamp as the light source (Fig. 4a). As the projected image is scanned by the detector, the photocurrent is amplified by the amplifier and recorded by the high-speed acquisition card. The four different chemical solvents, of water, IPA, toluene, and methanol (MeOH), in the glass vials are the imaging objects, as shown in Fig. 4b(i). Under a positive bias, the dual-band detector operates in the visible mode. The gap between the current signals at the baseline of the same height is hardly noticeable (Fig. 4b(ii)). All liquids are transparent, and the liquid height is difficult to determine (Fig. 4b(ii)). The dual-

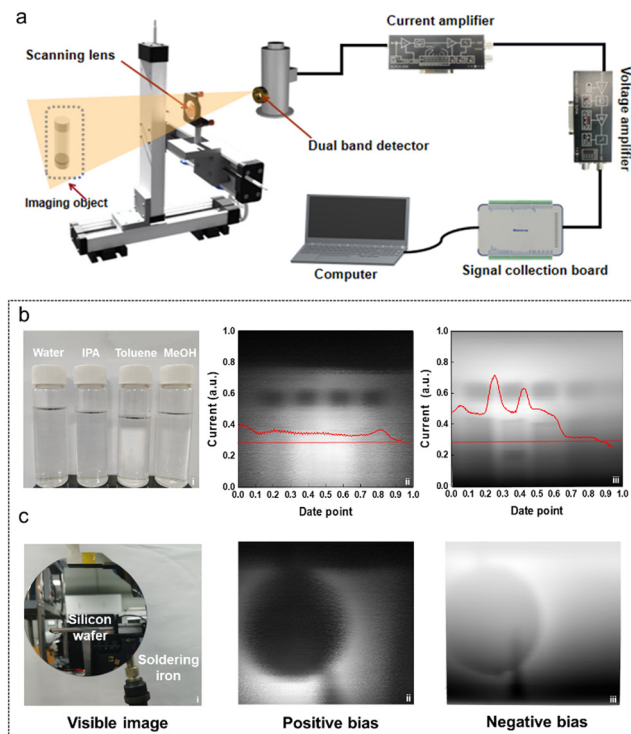


Fig. 4 Dual-band imaging. (a) Illustration of the single-pixel scanning imaging setup. (b) and (c) Visible images (i) and the images captured by the dual-band detector under positive (ii) and negative (iii) bias voltages for (b) four different chemical solvents in glass vials and (c) a silicon wafer and soldering iron arrangement.

band detector changes to the SWIR mode under a negative bias. Transparent objects in visible light can be identified in the SWIR range due to their different vibrational absorptions. As expected, there are significantly different current signals at the baseline of the same height for various chemical solvents, causing the transparent solvents in visible mode to appear differently in the SWIR mode (Fig. 4b(iii)). Besides, a soldering iron held behind the silicon wafer cannot be seen in the visible mode of the dual-band detector under a positive bias, as shown in Fig. 4c(ii). However, the soldering iron that is blocked by silicon wafer can be clearly captured by the dual-band detector when under a negative bias, indicating that the device operates in the SWIR mode (Fig. 4c(iii)).

The dual-band detector operating in the visible and SWIR modes under positive and negative bias voltages, respectively, are verified by the imaging results. When a positive voltage is applied, the visible CdTe/ZnO heterojunction is in the reverse bias and the visible channel begins to operate. The ZnO layer serves as both an electron-transporting layer and a hole-blocking layer, which not only promotes electron transport in the visible channel but also inhibits hole injection from the visible to the SWIR absorbing layer. When the device is operated under a negative bias, the SWIR ZnO/HgTe/Ag₂Te heterojunction begins to operate in the reverse region. Similarly, the large band-edge offset between the ZnO layer and the HgTe layer enables the dual-band device to respond only to SWIR

light. Thus, the detector realizes visible and SWIR dual-mode imaging through changing the bias polarity, providing more information about detected objects than is achievable from single-band images.

Conclusions

In this work, we constructed a metal chalcogenide CdTe/HgTe CQD-based dual-band detector consisting of two stacked photo-diodes, which realizes bias-controlled visible and SWIR dual-operation modes. An n-type ZnO layer is introduced between the VIS CdTe and SWIR HgTe layers as both an electron-transporting layer and a hole-blocking layer to prevent the injection of holes from one sensor layer to the other. The working mechanism of the dual-band PD is investigated through simulations. With the well-designed architecture, the dual-band device can switch between visible and SWIR modes by changing the polarity and the magnitude of the bias. The response peaks of the device are at 700 and 2100 nm in the visible and SWIR modes, respectively. The performance of the dual-band device is excellent, with a low noise current of about 10^{-13} A Hz^{-1/2}, high responsivity values of 0.5 A W⁻¹ (VIS) and 1.1 A W⁻¹ (SWIR), and a high detectivity above 10^{11} Jones both in the visible and SWIR modes at room temperature. VIS/SWIR dual-band imaging is demonstrated, providing more comprehensive object information than is achievable using single-band PDs. The obtained high-performance dual-band visible and SWIR detector offers the potential for more applications, such as autonomous driving, biomedicine, computer surveillance, and military operations.

Author contributions

The manuscript was written through the contributions of all authors. All authors have approved the final version of the manuscript.

Conflicts of interest

There are no conflicts to declare.

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References

- Z. Lan, Y. Lei, W. K. E. Chan, S. Chen, D. Luo and F. Zhu, *Sci. Adv.*, 2020, **6**, 1–10.

- 2 J. Zhao, X. Wang, Y. Xu, Y. Pan, Y. Li, J. Chen, Q. Li, X. Zhang, Z. Zhu, Z. Zhao, E. E. Elemike, D. C. Onwudiwe, B. S. Bae, S. Bin Shafie and W. Lei, *ACS Appl. Mater. Interfaces*, 2022, **14**, 25824–25833.
- 3 C. Li, J. Lu, Y. Zhao, L. Sun, G. Wang, Y. Ma, S. Zhang, J. Zhou, L. Shen and W. Huang, *Small*, 2019, **15**, 1–10.
- 4 S. Zhang, G. Mu, J. Cao, Y. Luo, Q. Hao, M. Chen, Y. Tan, P. Zhao and X. Tang, *Photonics Res.*, 2022, **10**, 1987–1995.
- 5 Y. Zhang, L. Huang, J. Li, Z. Dong, Q. Yu, T. Lei, C. Chen, L. Yang, Y. Dai, J. Zhang, W. Yu, Q. Bao and K. Zhang, *Appl. Phys. Lett.*, 2022, **120**, 261101.
- 6 E. T. Simola, A. De Iacovo, J. Frigerio, A. Ballabio, A. Fabbri, G. Isella and L. Colace, *Opt. Express*, 2019, **27**, 8529.
- 7 L. Li, H. Chen, Z. Fang, X. Meng, C. Zuo, M. Lv, Y. Tian, Y. Fang, Z. Xiao, C. Shan, Z. Xiao, Z. Jin, G. Shen, L. Shen and L. Ding, *Adv. Mater.*, 2020, **32**, 1–9.
- 8 X. Tang, M. M. Ackerman, M. Chen and P. Guyot-Sionnest, *Nat. Photonics*, 2019, **13**, 277–282.
- 9 Y. Li, Z. Shi, W. Liang, L. Wang, S. Li, F. Zhang, Z. Ma, Y. Wang, Y. Tian, D. Wu, X. Li, Y. Zhang, C. Shan and X. Fang, *Mater. Horiz.*, 2020, **7**, 530–540.
- 10 X. Xu, W. Deng, X. Zhang, L. Huang, W. Wang, R. Jia, D. Wu, X. Zhang, J. Jie and S. T. Lee, *ACS Nano*, 2019, **13**, 5910–5919.
- 11 D. Shen, Z. Guan, M. Li, S. W. Tsang, W. Zhang, M. F. Lo and C. S. Lee, *J. Mater. Chem. C*, 2021, **9**, 3814–3819.
- 12 Y. Wang, B. Siegmund, Z. Tang, Z. Ma, J. Kublitski, S. Xing, V. C. Nikolis, S. Ullbrich, Y. Li, J. Benduhn, D. Spoltore, K. Vandewal and K. Leo, *Adv. Opt. Mater.*, 2021, **9**, 1–8.
- 13 B. Huang, J. Liu, Z. Han, Y. Gu, D. Yu, X. Xu and Y. Zou, *ACS Appl. Mater. Interfaces*, 2020, **12**, 48765–48772.
- 14 C. Li, H. Wang, F. Wang, T. Li, M. Xu, H. Wang, Z. Wang, X. Zhan, W. Hu and L. Shen, *Light: Sci. Appl.*, 2020, **9**, 31.
- 15 Y. Li, Y. Zhang, T. Li, M. Li, Z. Chen, Q. Li, H. Zhao, Q. Sheng, W. Shi and J. Yao, *Nano Lett.*, 2020, **20**, 5646–5654.
- 16 H. Kim, W. Kim, Y. Pak, T. J. Yoo, H. W. Lee, B. H. Lee, S. Kwon and G. Y. Jung, *Laser Photonics Rev.*, 2020, **14**, 1–7.
- 17 A. Imran, M. Sulaman, S. Yang, A. Bukhtiar, M. Qasim, S. Elshahat, M. S. A. Khan, G. Dastgeer, B. Zou and M. Yousaf, *Surf. Interfaces*, 2022, **29**, 101772.
- 18 S. Das, D. Pandey, J. Thomas and T. Roy, *Adv. Mater.*, 2019, **31**, 1–35.
- 19 A. Imran, M. Sulaman, M. Yousaf, M. A. Anwar, M. Qasim, G. Dastgeer, K. A. A. Min-Dianey, B. Wang and X. Wang, *Adv. Mater. Interfaces*, 2022, 2200105.
- 20 C. Li, J. Lu, Y. Zhao, L. Sun, G. Wang, Y. Ma, S. Zhang, J. Zhou, L. Shen and W. Huang, *Small*, 2019, **15**, 1–10.
- 21 S. Bhowmick, R. Sen, C. P. Saini, R. Singhal, L. Walczak, M. Gupta, D. M. Phase and A. Kanjilal, *J. Phys. Chem. C*, 2021, **125**, 4846–4859.
- 22 G. Konstantatos and E. H. Sargent, *Nat. Nanotechnol.*, 2010, **5**, 391–400.
- 23 M. V. Kovalenko, *Nat. Nanotechnol.*, 2015, **10**, 994–997.
- 24 S. Keuleyan, E. Lhuillier and P. Guyot-Sionnest, *J. Am. Chem. Soc.*, 2011, **133**, 16422–16424.
- 25 P. Zhao, G. Mu, M. Chen and X. Tang, *Coatings*, 2022, **12**, 499.
- 26 Y. Luo, S. Zhang, X. Tang and M. Chen, *J. Mater. Chem. C*, 2022, **10**, 8218–8225.
- 27 A. Imran, J. Jiang, D. Eric, M. N. Zahid, M. Yousaf, M. Ahmad and S. A. Hassan, *Results Phys.*, 2018, **10**, 241–247.
- 28 A. Imran, M. Sulaman, Y. Song, D. Eric, M. N. Zahid, M. Yousaf, M. I. Saleem, M. Li and D. Li, *J. Comput. Electron.*, 2021, **20**, 310–316.
- 29 S. G. Kumar and K. S. R. K. Rao, *Energy. Environ. Sci.*, 2014, **7**, 45–102.
- 30 X. Y. Yu, B. X. Lei, D. Bin Kuang and C. Y. Su, *Chem. Sci.*, 2011, **2**, 1396–1400.
- 31 F. D. De Menezes, A. G. Brasil, W. L. Moreira, L. C. Barbosa, C. L. Cesar, R. D. C. Ferreira, P. M. A. De Farias and B. S. Santos, *Microelectron. J.*, 2005, **36**, 989–991.
- 32 X. Tang and K. W. C. Lai, *ACS Appl. Nano Mater.*, 2019, **2**, 6701–6706.
- 33 M. Chen, X. Lan, X. Tang, Y. Wang, M. H. Hudson, D. V. Talapin and P. Guyot-Sionnest, *ACS Photonics*, 2019, **6**, 2358–2365.
- 34 M. M. Ackerman, X. Tang and P. Guyot-Sionnest, *ACS Nano*, 2018, **12**, 7264–7271.
- 35 X. Tang and K. W. C. Lai, in *2016 IEEE 16th International Conference on Nanotechnology (IEEE-NANO)*, IEEE, 2016, pp. 494–497.
- 36 M. Chen and P. Guyot-Sionnest, *ACS Nano*, 2017, **11**, 4165–4173.
- 37 M. G. Panthani, J. M. Kurley, R. W. Crisp, T. C. Dietz, T. Ezzyat, J. M. Luther and D. V. Talapin, *Nano Lett.*, 2014, **14**, 670–675.
- 38 C. Livache, B. Martinez, N. Goubet, J. Ramade and E. Lhuillier, *Front. Chem.*, 2018, **6**, DOI: [10.3389/fchem.2018.00575](https://doi.org/10.3389/fchem.2018.00575).
- 39 J. Jasieniak, B. I. MacDonald, S. E. Watkins and P. Mulvaney, *Nano Lett.*, 2011, **11**, 2856–2864.
- 40 S. Keuleyan, E. Lhuillier, V. Brajuskovic and P. Guyot-Sionnest, *Nat. Photonics*, 2011, **5**, 489–493.
- 41 G. Mu, T. Rao, S. Zhang, C. Wen, M. Chen, Q. Hao and X. Tang, *ACS Appl. Mater. Interfaces*, 2022, **14**, 45553–45561.
- 42 Y. Zhang, W. Xu, X. Xu, J. Cai, W. Yang and X. Fang, *J. Phys. Chem. Lett.*, 2019, **10**, 836–841.
- 43 J. Tong, L. Y. M. Tobing, Y. Luo, D. Zhang and D. H. Zhang, *Sci. Rep.*, 2018, **8**, 1–9.
- 44 Z. Xie, Z. Deng, X. Zou and B. Chen, *IEEE Photonics Technol. Lett.*, 2020, **32**, 1003–1006.
- 45 M. Dai, H. Chen, R. Feng, W. Feng, Y. Hu, H. Yang, G. Liu, X. Chen, J. Zhang, C. Y. Xu and P. Hu, *ACS Nano*, 2018, **12**, 8739–8747.
- 46 Y. Liang, C. Xie, C. Y. Dong, X. W. Tong, W. H. Yang, C. Y. Wu and L. B. Luo, *J. Mater. Chem. C*, 2021, **9**, 14897–14907.
- 47 Y. Kaya, A. Ravikumar, G. Chen, M. C. Tamargo, A. Shen and C. Gmachl, *AIP Adv.*, 2018, **8**, 075105.
- 48 P. Zhao, G. Mu, M. Chen and X. Tang, *Coatings*, 2022, **12**, 499.
- 49 B. I. MacDonald, A. Martucci, S. Rubanov, S. E. Watkins, P. Mulvaney and J. J. Jasieniak, *ACS Nano*, 2012, **6**, 5995–6004.
- 50 G. Jain, A. Ahnood, A. Chanaewa, K. Fox and P. Mulvaney, *Phys. Lett. A*, 2019, **383**, 1199–1202.
- 51 X. Shen, J. C. Peterson and P. Guyot-Sionnest, *ACS Nano*, 2022, **16**, 7301–7308.

- 52 J. Jasieniak, B. I. MacDonald, S. E. Watkins and P. Mulvaney, *Nano Lett.*, 2011, **11**, 2856–2864.
- 53 H. Liu, Y. Tian, Y. Zhang, K. Gao, K. Lu, R. Wu, D. Qin, H. Wu, Z. Peng, L. Hou and W. Huang, *J. Mater. Chem. C*, 2015, **3**, 4227–4234.
- 54 X. Guo, Q. Tan, S. Liu, D. Qin, Y. Mo, L. Hou, A. Liu, H. Wu and Y. Ma, *Nano Energy*, 2018, **46**, 150–157.
- 55 B. Chen, J. Liu, Z. Cai, A. Xu, X. Liu, Z. Rong, D. Qin, W. Xu, L. Hou and Q. Liang, *Nanomaterials*, 2019, **9**, 626.
- 56 M. Hou, Z. Zhou, A. Xu, K. Xiao, J. Li, D. Qin, W. Xu and L. Hou, *Nanomaterials*, 2021, **11**, 2071.
- 57 A. Jagtap, B. Martinez, N. Goubet, A. Chu, C. Livache, C. Gréboval, J. Ramade, D. Amelot, P. Troussset, A. Triboulin, S. Ithurria, M. G. Silly, B. Dubertret and E. Lhuillier, *ACS Photonics*, 2018, **5**, 4569–4576.
- 58 U. N. Noubé, C. Gréboval, C. Livache, A. Chu, H. Majjad, L. E. Parra López, L. D. N. Mouafo, B. Doudin, S. Berciaud, J. Chaste, A. Ouerghi, E. Lhuillier and J. F. Dayen, *ACS Nano*, 2020, **14**, 4567–4576.
- 59 M. M. Ackerman, M. Chen and P. Guyot-Sionnest, *Appl. Phys. Lett.*, 2020, **116**, 083502.
- 60 E. Lhuillier, S. Keuleyan, P. Zolotavin and P. Guyot-Sionnest, *Adv. Mater.*, 2013, **25**, 137–141.
- 61 X. Tang, M. M. Ackerman and P. Guyot-Sionnest, *ACS Nano*, 2018, **12**, 7362–7370.
- 62 Y. Ning, S. Zhang, Y. Hu, Q. Hao and X. Tang, *Coatings*, 2020, **10**, 1–8.