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## Phase stability and charge compensation in disordered rock salt compounds based on nickel and titanium†

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Disordered rock salt compounds (DRX) are increasingly investigated as an alternative to established layered oxides, due to the significant compositional flexibility allowing the use of cheap elements, as well as the demonstrated high specific discharge capacities. In this work we investigate DRXs based on Ni as the main redox active element and Ti as a charge compensator and structure stabilizer, according to composition  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  ( $0.50 \leq y \leq 0.67$ ). We prepare such materials by solid state synthesis and carefully investigate their stability as single phases by combining theoretical calculations based on density functional theory and *in situ* X-ray diffraction, elucidating the role of entropy in the synthesis process. By a combination of SEM, ICP-AES and X-ray and neutron diffraction we reveal that the above-mentioned formula indicating a solid solution is only appropriate at high temperatures. When increasing the Li content and lowering the temperature, a secondary  $\text{Li}_2\text{TiO}_3$ -like phase emerges. The electrochemical performances showed a low capacity in the first cycle, which appears to be kinetically limited as it strongly increases at 50 °C. A strong voltage hysteresis is nonetheless present regardless of the temperature and is the main bottleneck limiting energy efficiency in these compounds. By combined *operando* X-ray diffraction and X-ray absorption near edge structure (XANES) spectroscopy, we demonstrate such hysteresis can be observed both structurally in the unit cell parameters as well as electronically at the Ni K-edge, hence in the Ni oxidation state. Our study confirms the peculiar, coupled behaviour of Ni and O redox in this class of materials, which needs to be overcome if new improved compositions are to be designed. We therefore provide a clear structural analysis of the synthesis and electrochemical behaviour of the  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  family of compounds, identifying the main bottlenecks towards their improvement.

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## Introduction

The successful commercialization of Li-ion batteries as electrochemical energy storage solutions with high energy density made a wide range of applications possible, including portable electronic devices, power tools and electric vehicles.<sup>1</sup> In the past decade, many efforts have been made to develop high energy density cathode active materials (CAMs), which are nevertheless still considered as one of the main bottlenecks for further improvements.<sup>2</sup> Commercial lithium transition metal oxides are compounds with formula  $\text{Li(TM)}\text{O}_2$  (with TM being a mixture of Mn/Al, Co and Ni) that crystallize in a layered crystal structure.<sup>3,4</sup> Despite their success, layered CAMs have a practical gravimetric capacity in the range of 160–230  $\text{mA h g}^{-1}$ , that is insufficient for some applications (for example for achieving extended driving ranges at affordable price in electric vehicles), along with safety issues that increase with increasing achievable reversible capacity and that have not been overcome yet.<sup>2</sup> To address these drawbacks, the extensive exploration of new materials brought attention to Li-rich cation disordered rock salt (DRX) materials,<sup>5,6</sup> due to the larger capacities achievable (even over 300  $\text{mA h g}^{-1}$ )<sup>2,7</sup> often relying on the oxygen redox contribution beyond the one from the transition metals alone.<sup>8–10</sup> In terms of stoichiometry, DRX materials can also be written as  $\text{Li}_{1+x}(\text{TM})_{1-x}\text{O}_2$ , yet now Li and TM occupy statistically the same crystallographic site. Concerning Li diffusion, the hopping mechanism is the same between layered and rock salt structures: lithium ions jump through a tetrahedral site between two octahedral sites (o–t–o diffusion pathway).<sup>11</sup> The  $\text{Li}^+$  in this tetrahedral site represents an activated state for lithium diffusion and the energy of this state will affect the lithium mobility in the crystal. The energy migration barrier is strongly dependent on the electrostatic repulsion between the  $\text{Li}^+$  and the element face-sharing the tetrahedron, and consequently increases when there are one or several TMs close to the activated  $\text{Li}^+$ . In layered oxides, the tetrahedron is always face-sharing with one TM from the layer above/below, yet Li diffusion is still possible within the Li layers. On the other hand, when no face-sharing TM are present, Li migration proceeds through 0-TM channels leading to an enhanced hopping rate, especially if compared to the 1-TM or 2-TM channels with a high migration barrier, thus inactive for Li diffusion in the structure. Therefore, the Li excess ( $x > 0$ ) in the composition is key to the formation of a percolating network of 0-TM channels (it was shown that it should be  $x > 0.1$  to allow percolation).<sup>8</sup> Moreover, in contrast to layered oxides, where a large amount of lithium removal affects the structure and consequently lithium transport due to strong electrostatic interactions between  $\text{Li}^+$ ,  $\text{O}^{2-}$  and  $\text{TM}^+$  along the  $c$  unit cell parameter,<sup>12,13</sup> in DRX compounds the more homogeneous cation distribution generates less structural changes and potentially less deterioration during cycling.<sup>14</sup> Finally, one of the common degradation mechanisms of layered materials is their reduction at the surface, making them transform into a disordered rock salt phase; such a mechanism is hence expected to play a minor role in DRX materials.<sup>4,15</sup> Despite the above-

mentioned benefits of DRX materials such as the oxygen activity along cycling, volume stability and high specific capacity, there are still drawbacks that need to be overcome. These are in particular the slow kinetics often reported (but not always),<sup>16</sup> the low average voltage and large voltage hysteresis, as well as the voltage decay associated with the anionic reactions.<sup>17</sup>

In terms of design criteria, the DRX formula can further be expanded as  $\text{Li}_{1+x}(\text{M}'\text{M}'')_{1-x}\text{O}_2$ , where  $\text{M}'$  represents a redox active species and  $\text{M}''$  is a redox inactive element ideally with a  $d^0$  electronic configuration, which has been reported to be often essential to promote the disorder and stabilize the crystal structure.<sup>18</sup> Most DRX compositions reported use Ni or Mn as redox-active species, while common  $d^0$  elements are Ti, Nb, Mo and few others. So far, Mn-based DRX proved superior in terms of observed capacities,<sup>19</sup> yet Ni is also promising due to its  $2e^-$  redox capability, high potential expected for  $\text{Ni}^{2+/4+}$  vs.  $\text{Li}^+/\text{Li}$ , and the large know-how developed in the field of Ni-rich layered CAMs, especially concerning their synthesis at scale. In this work, we investigate DRX compounds with a target composition  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  ( $0.50 \leq y \leq 0.67$ ) where the Ni redox activity  $\text{Ni}^{2+} \rightarrow \text{Ni}^{3+/4+}$  is exploited and  $\text{Ti}^{4+}$  is used as charge compensator. Ti being lighter and cheaper than most other  $d^0$  elements,<sup>20</sup> it may be expected to partially offset the higher cost of Ni with respect to Mn. Fig. 1 illustrates a Gibbs triangle in the chemical space comprising  $\text{NiO}$ – $\text{Li}_2\text{NiO}_3$ – $\text{Li}_2\text{TiO}_3$ . Here, Ni in the reported compounds is the only element which formally changes oxidation state from 2+ ( $\text{NiO}$ ) to 4+ ( $\text{Li}_2\text{NiO}_3$ ). Along the  $\text{NiO}$ – $\text{Li}_2\text{NiO}_3$  tie line, we find the important cathode material  $\text{LiNiO}_2$  (LNO), which is a prototypical Ni-rich layered oxide.<sup>4</sup> Under sufficiently oxidizing conditions, one can over-lithiate LNO to reach  $\text{Li}_2\text{NiO}_3$ . While this compound should be redox-inactive based on the Ni valence state, it has been shown to possess significant capacity, mostly attributed to O redox activity, which however fades quickly.<sup>21</sup> Therefore, in our aim to design a DRX CAM exploiting mostly Ni redox, we target the lower portion of the Gibbs triangle where Ni is divalent. From  $\text{LiNiO}_2$ , it is possible to substitute Ti for Ni according to  $\text{LiNi}_{1-z}\text{Ti}_z\text{O}_2$ . This substitution (doping) has been already heavily investigated and will not be discussed further, as overall it decreases the available capacity significantly.<sup>22,23</sup> On the other hand, we target the region at the bottom of Fig. 1 between  $\text{LiNi}_{0.5}\text{Ti}_{0.5}\text{O}_2$  and  $\text{Li}_2\text{TiO}_3$ . We start for  $y = 0.5$  at  $\text{LiNi}_{0.5}\text{Ti}_{0.5}\text{O}_2$ , which is a stoichiometric compound ( $\text{Li}/\text{TM} = 1$ ) and has been reported to possess a disordered rock salt structure.<sup>24–27</sup> Since this compound is not Li-rich, it possesses few 0-TM channels and therefore no percolation pathways for Li, hence it displays very low specific capacity.<sup>27</sup> On the other hand, in the direction followed by the materials we investigate (marked with a red arrow in Fig. 1) the compounds become Li-rich with increasing  $y$ , while Ni maintains his 2+ oxidation state throughout the series. The compounds have been labelled as LNT050, LNT055, LNT060 as function of the  $y$  value (where the number represents the amount of Ti and half the amount of Li), whereas LTO67 corresponds to  $\text{Li}_{1.33}\text{Ti}_{0.67}\text{O}_2$  (or  $\text{Li}_2\text{TiO}_3$ ). The same compositional region was studied very recently by Li *et al.*,<sup>28</sup> whom however indicated it as a two phase region with  $x\text{Li}_2\text{TiO}_3 \cdot (1-x)\text{LiTi}_{0.5}\text{Ni}_{0.5}\text{O}_2$  composition. Their findings highlight



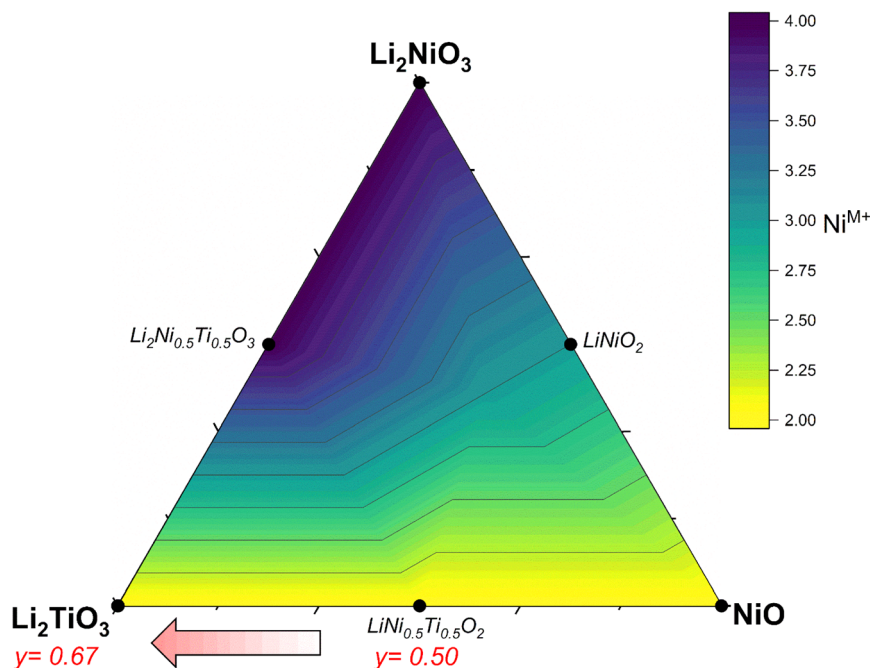


Fig. 1 Gibbs triangle (pseudo phase diagram) in the system  $\text{Li}_2\text{TiO}_3$ – $\text{Li}_2\text{NiO}_3$ – $\text{NiO}$ . The red arrow depicts the compositions of the materials  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  under investigation. The color scale indicates the expected nominal Ni oxidation state.

the activation of the anionic redox mechanism through a ligand-to-metal charge transfer process. However, it is at present unclear whether these materials are indeed a biphasic mixture, or rather single phase. Hence, we undertook our study to verify the single vs. 2-phase behavior of such samples, and to further highlight the structural, electrochemical and charge transfer behavior of these materials.

## Experimental section

### Materials synthesis

All the materials were synthesized by solid-state reaction preceded by a mixing of the precursors in a Spex ball mill for 1.50 h using zirconia jars. The used precursors are nickel(II) hydroxide (61% Ni, Alfa Aesar), titanium(IV) oxide (nanopowder, Sigma-Aldrich  $\geq 99.5\%$ ) and lithium carbonate (1% excess, Thermo Scientific 99.998%). The mixture was pelletized and calcined in air at three different annealing temperatures (700 °C, 800 °C and 900 °C) for 10 h in an alumina crucible with a heating rate of 7 °C  $\text{min}^{-1}$  and cooling rate of 5 °C  $\text{min}^{-1}$ .

### Computational details

The optimization of the structural models for  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  ( $y = 0.5; 0.6; 0.67$ ) were performed by using the Vienna *Ab initio* Simulation Package (VASP)<sup>29–31</sup> within the projector augmented-wave approach<sup>32</sup> employing the Perdew–Burke–Ernzerhof (PBE)<sup>33</sup> with the  $r^2\text{SCAN}$  functional.<sup>34</sup> Correlation effects of the 3d-Ni orbitals (corresponding parameter  $U_{\text{Ni}} = 2.1$  eV)<sup>35</sup> as well as long-range dispersion with Grimme's D4 correction<sup>36</sup> were taken into account. For the calculations of the total energy a plane-wave energy cut-off of 600 eV and  $4 \times 4 \times 4$  gamma centered  $k$ -point grid were used.

The initial structural models are based on cubic  $\text{Li}_{0.98}\text{Ni}_{1.02}\text{O}_2$  (S.G.  $Fm\bar{3}m$ ,  $a = 4.1708(5)$  Å<sup>3</sup>) and layered  $\text{Li}_{0.98}\text{Ni}_{1.02}\text{O}_2$  ( $R\bar{3}m$ ,  $a = 2.87758(3)$  Å,  $c = 14.1980(1)$  Å<sup>3</sup>). For each composition ( $y = 0.5; 0.6; 0.67$ ) the occupancies of the metal cations Li, Ni and Ti were adjusted accordingly. In the cubic system ( $y = 0.5$ :  $\text{Li}_4\text{Ni}_2\text{Ti}_2\text{O}_8$ ;  $y = 0.6$ :  $\text{Li}_4\text{NiTi}_3\text{O}_8$ ) the  $2 \times 2 \times 2$  supercell was chosen as starting point ( $P1$ ,  $a = b = c = 5.857$  Å,  $\alpha = \beta = \gamma = 60^\circ$ ). With  $y = 0.67$  ( $\text{Li}_8\text{Ti}_4\text{O}_{12}$ ), a  $2 \times 2 \times 3$  supercell was employed ( $P1$ ,  $a = b = 5.862$  Å,  $c = 8.794$  Å,  $\alpha = \beta = \gamma = 60^\circ$ ). For all the possible configurations generated by Supercell the total energy was calculated using VASP. The final structure with the most stable configuration was then used for comparison with the hexagonal one. For the layered structure type ( $y = 0.5$ :  $\text{Li}_4\text{Ni}_2\text{Ti}_2\text{O}_8$ ;  $y = 0.6$ :  $\text{Li}_4\text{NiTi}_3\text{O}_8$ ) the starting point was a  $2 \times 2 \times 1$  supercell ( $P1$ ,  $a = b = 10.042$  Å,  $c = 5.021$  Å,  $\alpha = \beta = \gamma = 33.29^\circ$ ), whereas for  $y = 0.67$ :  $\text{Li}_2\text{TiO}_3$  a  $2 \times 2 \times 2$  supercell was chosen ( $P1$ ,  $a = b = 10.042$  Å,  $c = 5.021$  Å,  $\alpha = \beta = \gamma = 33.29^\circ$ ). Again, the total energy of all possible configurations given by supercell was calculated and the most stable one was selected for comparison with the cubic phase. For  $\text{Li}_2\text{TiO}_3$  ( $y = 0.67$ ), which is layered but with a honeycomb ordering resulting in a monoclinic unit cell, only one cell was calculated, since there is no partial occupancy of Li or Ti (optimized cell:  $P1$ ,  $a = 5.071$  Å,  $b = 8.796$  Å,  $c = 9.706$  Å,  $\alpha = \gamma = 90^\circ$ ,  $\beta = 100.16^\circ$ <sup>37</sup>).

### Material characterization methods

Powder X-ray diffraction (XRD) patterns were collected using a Stoe StadiP diffractometer with a Mythen2K detector in Debye–Scherrer geometry in  $2\theta$  angular range of 5–50° equipped with a monochromatic  $\text{Mo-K}\alpha_1$  radiation (wavelength  $\lambda = 0.70932$  Å). The acquisition was performed on powder packed in



a 0.5 mm diameter borosilicate glass capillary. Neutron powder diffraction (NPD) was carried out on the D2B neutron diffractometer at the Institut Laue-Langevin (ILL). Powder samples were put into cylindrical vanadium cans and measured in transmission geometry at  $\lambda = 1.59485(3)$  Å (<https://doi.org/10.5291/ILL-DATA.5-23-77>), where the  $\lambda$  value was opportunely refined using FullProf software.<sup>38</sup> For both XRD and NPD, structural refinements of the rock salt crystalline models were performed against the collected data and the error bars reported as  $3\sigma$  accordingly to the value obtained from the FullProf software. *In situ* temperature resolved XRD was performed by using the Stoe ST2K furnace mounted on the Stoe StadiP diffractometer. Patterns were recorded in the  $2\theta$  range of 5–35° at a temperature between 30 °C and 900 °C. The temperature is hold at each temperature every 10° for 10 minutes, during which the acquisition is carried out. A sapphire capillary with two open ends was filled with the powder and the XRD data were collected in air as function of temperature. Quartz wool was placed on each side of the powder sample to prevent his movement during the experiment. The furnace undergoes a rocking movement during acquisition to improve powder averaging.

The morphological features of the samples sputtered with platinum were observed using a LEO-1530 electron microscope in secondary electron detector mode and a magnification of 75 000×. Particle size distribution was estimated using the ImageJ software.

### ICP-AES

The metal elements were determined by ICP-AES (Inductively Coupled Plasma – Atomic Emission Spectroscopy with a spectrometer iCap 6500 from Thermo Instruments). About 2–3 mg of the samples was dissolved in 1 mL of hydrochloric acid, hydrofluoric acid and nitric acid at 180 °C for 12 hours. All samples were dissolved three times. The digestions were diluted, and analysis of the elements was accomplished with different calibration solutions and an internal standard. 4 wavelengths of each element were used for calculation/analysis.

### Electrochemical measurements

Cathodes were prepared by slurry casting on Al foil. The slurry was obtained by dispersing 80 wt% active material, 10 wt% Super C65 carbon black (Timcal) and 10 wt% Solef polyvinylidene fluoride (PVDF) binder in *N*-methyl-2-pyrrolidone. The slurry was dried overnight and the as-obtained tapes were cut into round electrodes using a steel punch. For the electrochemical testing, coin half-cells consisting of cathode, porous glass fiber (Whatman, GF/D grade), Li metal anode (Albemarle Germany GmbH) and 90  $\mu$ L LP30 electrolyte (1 M LiPF<sub>6</sub> in 1 : 1 EC/DMC) were assembled inside an Ar-filled glovebox. The loading of the electrodes was  $\approx 3$  mg cm<sup>-2</sup>. The electrochemical behavior of every compound was tested in the voltage range between 1.5 and 4.8 V *versus* Li<sup>+</sup>/Li at C/10 (1C corresponding to a current density of 200 mA g<sup>-1</sup>). The tests were performed at room temperature (25 °C controlled in a binder chamber) and at 50 °C under the same experimental conditions. *Operando* XRD and X-ray absorption

near-edge structure (XANES) spectra at the Ni K-edge were collected semi-simultaneously on the same coin cell in transmission mode at the beamline NOTOS of the ALBA Synchrotron in Barcelona. The half-cell was charged and discharged for 1 cycle at C/15 (1C corresponding to a current density of 149 mA g<sup>-1</sup>) while XRD patterns and XANES patterns were alternately acquired (0.5 s for XRD in the 7 to 50° angular range, 11 min for XANES in the energy range from 8200 eV to 8700 eV).

## Results

We begin our study with a computational investigation to evaluate the likelihood of obtaining rock salt structures; in other word to quantify at a given composition the phase stability of a rock salt structure *versus* a layered one, which is also likely for the materials in this phase diagram. Two different structural models (NaCl-type and layered-type) for Li<sub>2y</sub>Ti<sub>y</sub>Ni<sub>2-3y</sub>O<sub>2</sub> were investigated by DFT on supercells generated by the software Supercell.<sup>39</sup> All details on the supercells used can be found in the Experimental section. For Li<sub>2</sub>TiO<sub>3</sub> ( $y = 0.67$ ), which is layered but with a honeycomb ordering resulting in a monoclinic unit cell, the energy of the ordered monoclinic cell was also calculated. The results of our calculations are summarized in Fig. 2a.

The computed relative energies of each phase were calculated at 0 K to understand whether the layered structure (LS) or rock salt structure (RS) was more stable as a function of the different Li content in the ground state. The results displayed in Fig. 2a show that the LS exhibits lower energy values than the RS, which for this reason would be energetically less stable. Therefore, one may expect the synthesis of our target compounds as a single RS phase to be challenging, and that phase separation may be preferable. On the other hand, the difference in energy between rock salt and ground state layered phase is small (0.26 eV, 0.38 eV and 0.47 eV for  $y = 0.5, 0.6$  and 0.67, respectively), growing slightly with increasing  $y$ . One should then notice that in the reported calculations there is no entropic component. In fact, a different case could be expected when the configurational entropy contribution for the two different phases is included. In particular, the configurational entropy corresponds to all the distinct ways in which the lattice sites can be occupied at a specific composition.<sup>40</sup> In the disordered RS phase, the cations are randomly distributed in the same site, while in the layered phase, they are ordered. Therefore, there is no configurational entropy for a “perfect” layered structure, while one may expect a large configurational entropy for a rock salt phase. In light of these considerations, since the energetic difference at 0 K between the phases is small, we expect that by sufficiently increasing the temperature, the Gibbs free energy of the rock salt phase will drop to more negative values as compared to the layered one (Fig. 2b). When  $y > 0.5$ , as shown in Fig. 2c, this transition temperature might rise due to the higher initial enthalpic difference to overcome, but it would occur nonetheless. Working under the above-mentioned assumptions, we predict to observe a pure rock salt phase above a certain synthesis temperature, and that the exact value would increase with increasing  $y$ .





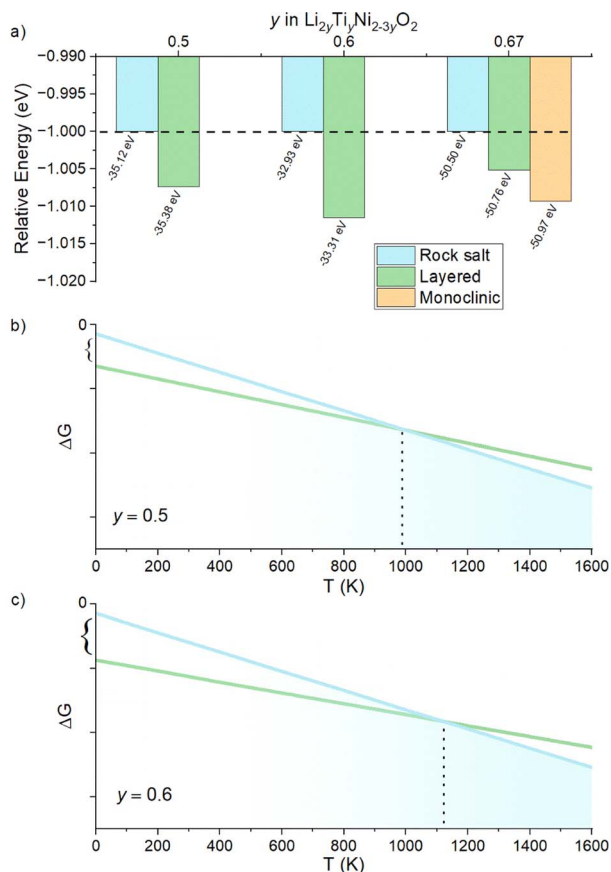


Fig. 2 (a) Relative energies of the  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  system given per formula unit, as computed with VASP. The layered phase shows the lowest energy and thus highest stability compared to the NaCl type. (b and c) Schematic representation of the Gibbs free energy vs. temperature plotted for different Li contents ( $y = 0.5$  in (b),  $y = 0.6$  in (c)). In both cases a transition  $T$  might be reached where the RS overcome to the LS in stability.

The samples were then synthesized by solid state synthesis at 700, 800 and 900 °C as described in the Experimental section, and further investigated by X-ray and neutron diffraction, as gathered in Fig. 3.

The samples synthesized at 900 °C and 800 °C (Fig. S1 and S2†) exhibit characteristic diffraction peaks of a rock salt-type structure with  $Fm\bar{3}m$  space group. The end member  $\text{Li}_2\text{TiO}_3$  crystallizes into a monoclinic structure with  $C2/c$  space group<sup>41</sup> regardless of the temperature used. For this end-member, one may notice the low 002/104 peak intensity ratio: this can be used to indicate a high degree of Li/Ti mixing.<sup>42</sup> However, because of the significant impact of other structural defects such as stacking faults, the exact Li/Ti distribution is not investigated further. By decreasing the temperature to 700 °C (Fig. 3), the samples  $y = 0.5, 0.55$  and  $0.6$  still appear as rock salt, yet a secondary phase starts to emerge associated to a  $\text{Li}_2\text{TiO}_3$ -like phase, observed for all the synthesized samples at this temperature but with a clearly increasing trend for increasing Li content. Namely, the second phase grows upon Li and Ti addition as evidenced from the evolution of the 002<sub>LTO</sub> peak in the XRD pattern in Fig. 3. All diffraction data were analysed by

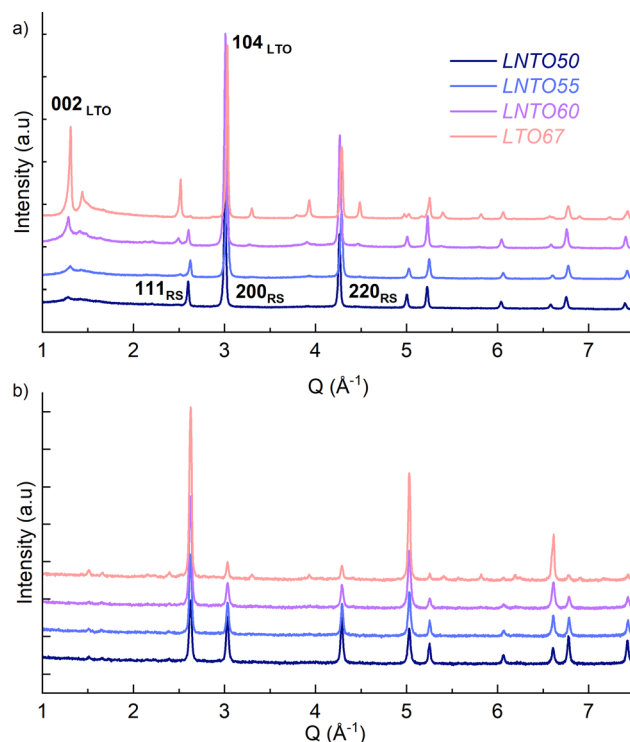


Fig. 3 X-ray diffraction (a) and neutron diffraction (b) patterns of the samples annealed at 700 °C. The  $2\theta$  angle is converted to the scattering vector  $Q$  to compare the two experiments despite the different wavelength (XRD:  $\lambda = 0.70932$  Å, NPD:  $\lambda = 1.59485(3)$  Å).

Rietveld refinement assuming a single rock salt phase at 800 °C and 900 °C, while using two-phase fitting for the samples annealed at 700 °C. As shown in Fig. 4a, with this method we are able to estimate the phase fractions and unit cell parameters of the different phases observed. Here, the results clearly show how by decreasing the temperature and increasing the  $y$  content, the monoclinic phase amount consequently grows, conforming to our DFT-based predictions. As can be seen in Fig. 3, the superstructure peaks related to the  $C2/c$  phase are barely present in the NPD data, even for pure  $\text{Li}_2\text{TiO}_3$ , while they are clear in the XRD data around  $1.3$  Å<sup>-1</sup>. This may seem surprising, but it stems from the structure factors for NPD being all essentially zero below  $Q = 2.5$  Å<sup>-1</sup>. For this reason, and due to the higher precision on the wavelength obtainable with X-rays, the phase fractions reported in Fig. 4a and the unit cell volumes reported in Fig. 4b are obtained from refinement of the XRD data alone.

In Fig. 4b we gather the unit cell volumes obtained *via* Rietveld refinement. The behaviour of the samples annealed at 900 °C and 800 °C revealed a decreasing trend as function of increasing  $y$  that can be ascribed to a single solid solution behavior. However, as expected from the average ionic radius of the cations (barely changing from 0.704 to 0.708 across the materials' series), only a very slight trend is observed. Instead, for the samples synthesized at 700 °C, fitted with a 2-phase model (refining only the structural parameters of the rock salt phase and maintaining constant those found from the



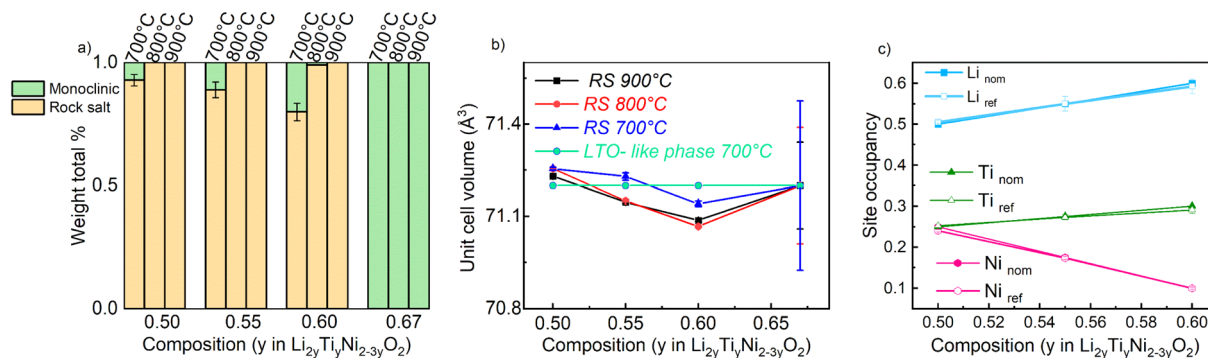


Fig. 4 Evolution of the phase content from Rietveld refinement in the  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  system. A pure rock salt phase is detected for the samples annealed at 900 °C and 800 °C, while a  $\text{Li}_2\text{TiO}_3$ -like phase appears and grows with increasing  $y$  at 700 °C. (b) Unit cell volume of the cubic phases. The value obtained for the  $\text{Li}_2\text{TiO}_3$ -like phase is divided by 6 to account for the larger  $Z$  in this structure. (c) Evolution of the site occupancy factors for the samples at 700 °C. Value in (a) and (b) are from refinement of XRD data, while those in (c) from combined refinement of XRD and NPD data.

refinement of the as-prepared  $\text{Li}_2\text{TiO}_3$  phase), a fully linear trend was not detected.

The fact that we observe an increasing amount of  $\text{Li}_2\text{TiO}_3$ -like phase with increasing  $y$  raises the question: what is the composition of the actual rock salt phase we observe? To answer this question, we firstly verified the nominal composition of the synthesised materials using ICP-AES measurements performed for all the samples synthesized at 700 °C. The different cationic ratio (Li/Ti, Li/Ni, Ni/Ti) are reported in Table S1.† The resulting data reveal that, while the Li/Ti ratio was in good agreement with the theoretical value, the Li/Ni and Ni/Ti ratio slightly differed from the expected ones, pointing towards a slight Ni-deficiency in the samples. We then used our X-ray and neutron diffraction data to carry out a combined refinement against a unique (rock salt) structural model. Since three atoms occupy the 4a Wyckoff site, we chose the following constraints for the refinement: (i) the cations fully occupy the site ( $\text{Li} + \text{Ni} + \text{Ti} = 1$ ) and (ii) given that the Li/Ti ratio was found as expected by ICP, the Li content fixed to two times the Ti content ( $\text{Li} - 2\text{Ti} = 0$ ). Having three unknown atoms contents on a site and two constraints, we proceed to the refinement. It should be mentioned that in this case, the added sensitivity of neutrons to light elements (Li, O) and to Debye–Waller factors makes a combined refinement a better choice to determine the cations distribution in the rock salt-type phase. The site occupancies are reported in Fig. 4c, as obtained by combined Rietveld refinement for the samples at 700 °C. Tables 1 and 2 gather the structural results from the combined refinement. The combined refinement of XRD and NPD for an exemplary sample, LNT055 made at 700 °C, is also displayed in Fig. 5. The structural parameters of the compounds annealed at 900 °C and 800 °C are instead shown in Table S2.†

From this analysis, as can be clearly seen in Fig. 4c, we can state that the refined values in general do not exhibit a discrepancy from the nominal ones, *i.e.* the synthesized rock salt-type phases, despite a  $\text{Li}_2\text{TiO}_3$ -like impurity, have composition very close to the desired one. Nonetheless, it is worth mentioning that we cannot determine the composition of the  $\text{Li}_2\text{TiO}_3$ -like side phase, hence STEM/EDX might be the ideal

Table 1 Exemplary structural parameters obtained from Rietveld refinement of combined XRD and NPD of the sample  $\text{Li}_{1.1}\text{Ni}_{0.35}\text{Ti}_{0.55}\text{O}_2$  (LNT055) synthesized at 700 °C

$\text{Li}_{1.1}\text{Ni}_{0.35}\text{Ti}_{0.55}\text{O}_2$						
S.G.: $Fm\bar{3}m$ , $Z = 4$			$\chi^2: 1.42$			
$a, b, c = 4.14382(24) \text{ \AA}$			XRD		NPD	
Vol = $71.154(6) \text{ \AA}^3$			$R_{\text{Bragg}}: 2.16\%$		$R_{\text{Bragg}}: 6.80\%$	
			$R_{\text{wp}}: 4.81\%$		$R_{\text{wp}}: 16.3\%$	
Atomic positions						
Atoms	Wyckoff site	$x$	$y$	$z$	$B_{\text{iso}}$	Occupancy
O	4b	0.5	0.5	0.5	1.15(12)	1.00000
Ti	4a	0.0	0.0	0.0	0.83(9)	0.269(6)
Ni	4a	0.0	0.0	0.0	0.83(9)	0.169(3)
Li	4a	0.0	0.0	0.0	0.83(9)	0.547(18)
Constraints: $\text{Li} + \text{Ni} + \text{Ti} = 1$ ; $\text{Li} - 2\text{Ti} = 0$						

method to further verify in the future the elemental composition of both phases.

To further elucidate the reason for the presence of a secondary  $\text{Li}_2\text{TiO}_3$ -like phase at low temperature, *in situ* XRD during synthesis of the  $y = 0.50$  sample was performed, as reported in Fig. 6. The experiment was carried out in the temperature range between 25 °C and 900 °C (where the cubic phase can be synthesized phase-pure). First, the mixture of the precursors was pre-annealed at 350 °C to dehydrate  $\text{Ni}(\text{OH})_2$  into  $\text{NiO}^{43}$  and avoid the occurrence of an exothermic reaction from the release of water.  $\text{NiO}$  presents a cubic structure with space group  $Fm\bar{3}m^{44}$  and the presence of the Bragg peaks associated to this phase are already visible at the beginning of the experiment. At this initial stage, the additional detected reflections were indexed to the precursors  $\text{Li}_2\text{CO}_3$  and  $\text{TiO}_2$ . Upon annealing, these reflections start vanishing at 350 °C and 400 °C, respectively. More interestingly, when the temperature is raised to 600 °C, a clear new peak is detectable at low angles (around  $2\theta = 8.30^\circ$ ), which can be indexed to an intermediate  $\text{Li}_2\text{TiO}_3$  phase that quickly forms during synthesis. This is in



Table 2 Structural parameters from combined XRD and NPD for the samples LNT050, LNT055 and LNT060 synthesized at 700 °C

Sample	Method	Unit cell parameters (Å)	Unit cell volume (Å <sup>3</sup> )	$B_{\text{iso}}$ (Å <sup>2</sup> )	Site occupancy
LNT050 700 °C	XRD	4.1457(2)	71.254(6)	—	—
	NPD	4.1458(1)	71.256(3)	—	—
	Combined	4.1451(2)	71.220(6)	Li, Ni, Ti: 1.06(6); O: 1.20(6)	Li: 0.505(3) Ni: 0.240(3) Ti: 0.252(3)
LNT055 700 °C	XRD	4.1451(4)	71.22(1)	—	—
	NPD	4.1442(2)	71.172(62)	—	—
	Combined	4.1438(2)	71.154(6)	Li, Ni, Ti: 0.83(9); O: 1.2(1)	Li: 0.55(2) Ni: 0.169(3) Ti: 0.269(6)
LNT060 700 °C	XRD	4.1430(3)	71.115(9)	—	—
	NPD	4.1432(2)	71.123(6)	—	—
	Combined	4.1423(3)	71.074(9)	Li, Ni, Ti: 1.05(9); O: 1.4(1)	Li: 0.60(2) Ni: 0.097(3) Ti: 0.293(9)

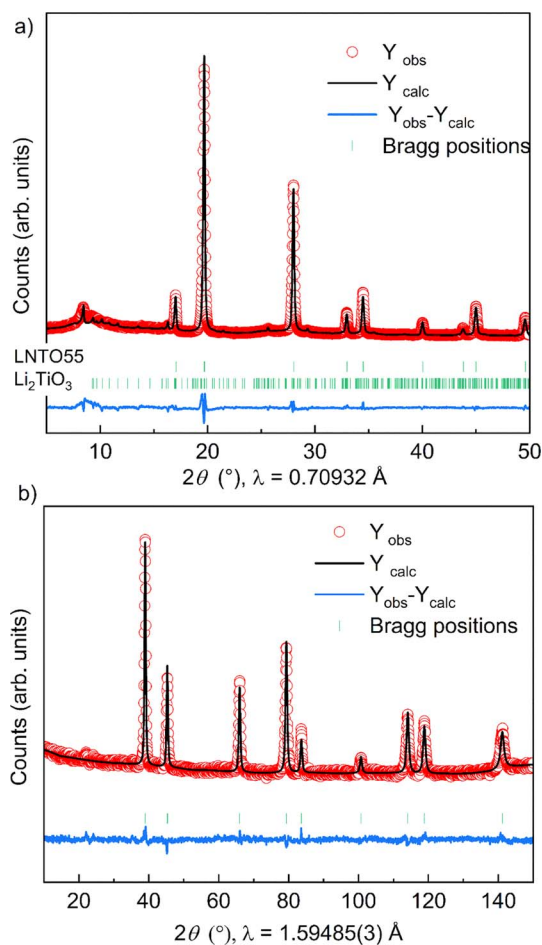


Fig. 5 Combined Rietveld refinement of XRD (a) and NPD (b) patterns for the exemplary sample LNT055 synthesized at 700 °C. Due to the very weak peak intensities, the secondary  $\text{Li}_2\text{TiO}_3$ -like phase is only included in the XRD refinement.

agreement with our DFT calculations, showing that at low  $T$  a layered-type phase would be preferred. Furthermore, it clearly shows the importance of intermediate compounds during solid

state synthesis, which template and ultimately affect which final phases can be accessed and at which temperature.<sup>45</sup> Again in agreement with our expectations, as the temperature increases beyond 700 °C and entropic effects start to dominate, the  $\text{Li}_2\text{TiO}_3$ -like phase disappears and a pure rock salt phase forms. Finally, the peaks 200 and 220 of the cubic crystal structure become much more pronounced as the compound further crystallizes, exhibiting a stronger intensity that persist to the final temperature at 900 °C and further during the cooling step. By decreasing the temperature, the reflections of the rock salt phase remain strong and shift to higher angles indicating a contraction of the lattice parameters.

Finally, to complete our structural investigation, scanning electron microscopy (SEM) data of the samples annealed at 700 °C were collected to analyse their morphology. The acquired images show that the primary particle sizes of the as prepared compounds are highly agglomerated into secondary particles. Roughly, the sample with  $y = 0.50$  have average primary particles of less than 80 nm, that increases with Li and Ti content as shown in Fig. S3.† However, such primary particles are agglomerated to clusters of several hundreds of nm.

### Electrochemical behavior and charge compensation

Galvanostatic charge/discharge tests were carried out for the pristine materials annealed at 700 °C to determine their electrochemical behaviour. Fig. 7 gathers the first cycle voltage profiles of LNT050, LNT055 and LNT060 within the voltage range 1.5–4.8 V at  $C/10$  and room temperature. LTO was not tested due to the expected limited capacity due to the low electrical conductivity.<sup>46</sup> In Fig. 7a, the voltage profile *versus* Li/formula unit is provided. All the curves revealed large hysteresis and rather low specific discharge capacity. For LNT050 ( $\text{LiNi}_{0.50}\text{Ti}_{0.50}\text{O}_2$ ) the capacity was expected to be limited<sup>25</sup> based on previous reports.<sup>27</sup> However, one could expect an improvement with increasing  $y$ , assuming a contribution from O redox is also possible. In the LNT055 sample, approximately 0.6 Li are extracted per formula unit ( $122.1 \text{ mA h g}^{-1}$ ), higher if compared to LNT050 and LNT060 materials ( $86 \text{ mA h g}^{-1}$  and



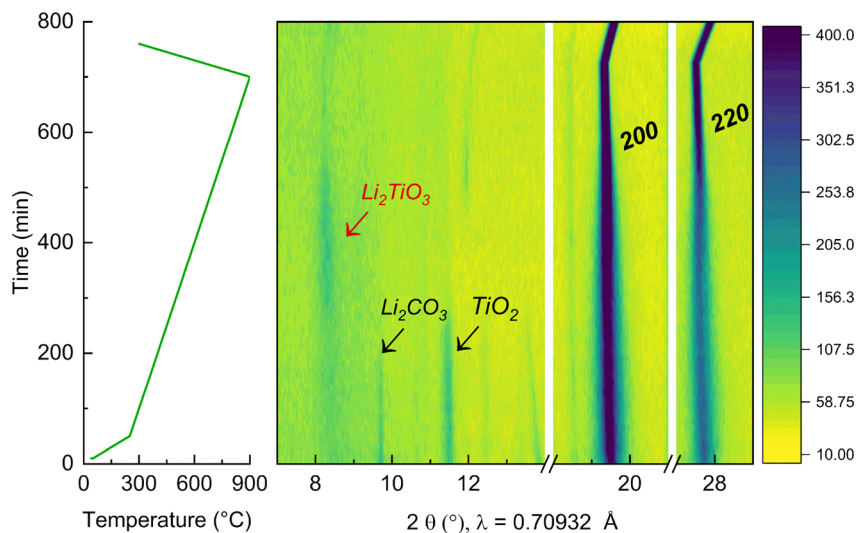


Fig. 6 *In situ* XRD during the synthesis of the sample LNT050 starting from NiO,  $\text{Li}_2\text{CO}_3$  and  $\text{TiO}_2$  as precursors. On the left the temperature evolution over time is shown. The central panels contain contour plots showing the Bragg peaks evolution during the synthesis from 25 °C to 900 °C. The color scale bar appears on the right.

102.2  $\text{mA h g}^{-1}$  respectively). This proves the Li excess is beneficial to Li percolation, yet the obtained capacities are lower than the theoretical capacity expected by Ni redox contribution, hence it is unclear if O is also contributing to the charge compensation. The overall performance of the Li-rich electrodes shows limited performance as compared to the results shown by Li *et al.*,<sup>28</sup> most likely because of the different synthesis process (in their case a sol-gel process yields smaller and less agglomerated crystallites). Nonetheless, the shape of

our curves (with their large voltage hysteresis) is closely related to the ones that Li *et al.* reported, suggesting the O redox contribution is present. To overcome possible kinetic limitations, we further examined the electrochemical behavior of the electrodes at 50 °C, leaving the other parameters unchanged. The reversible capacity tested in this condition is greatly improved as shown in Fig. 7b. Again, the sample LNT055 delivers the highest discharge capacity with about 223.7  $\text{mA h g}^{-1}$ , followed by LNT060 (213.4  $\text{mA h g}^{-1}$ ) and

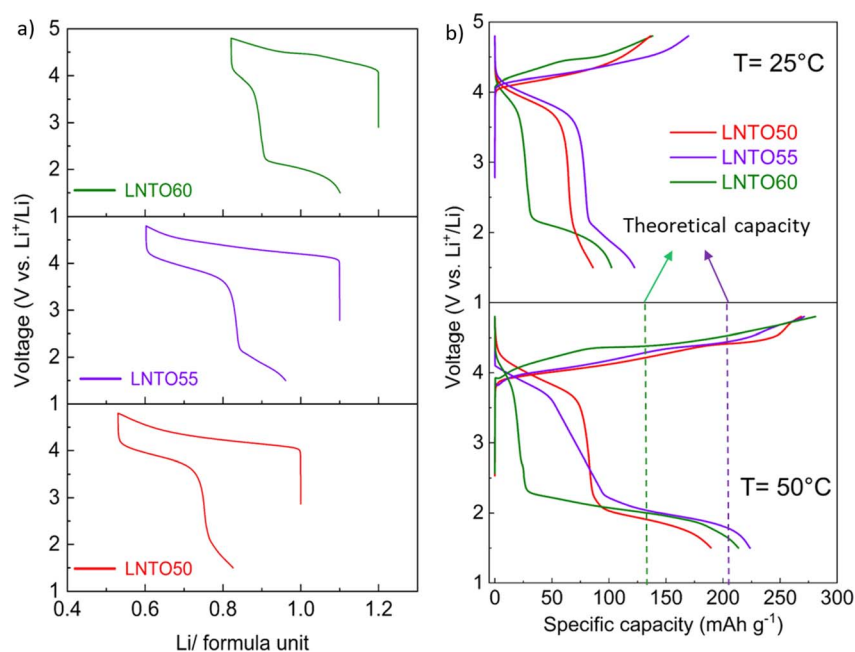


Fig. 7 Galvanostatic cycling of LNT050, LNT055 and LNT060 samples synthesized at 700 °C tested in Li half cells. (a) Voltage profiles versus lithium amount in the voltage range 1.5–4.8 V at 25 °C and  $C/10$  current rate. (b) Comparison of voltage profiles vs. specific capacity measured at 25 °C and 50 °C (same rate and voltage window as in (a)). The theoretical capacity based on the redox activity of Ni alone is indicated with dashed lines for LNT055 (violet) and LNT060 (green).





finally LNT050 ( $189.3 \text{ mA h g}^{-1}$ ) although the voltage hysteresis still appears very large. From these results we can conclude that (i) clear kinetic limitations are present in the studied materials, and (ii) a capacity larger than the theoretical one based on Ni redox can be observed for both LNT055 and LNT060 compounds ( $g_{\text{theo}} = 216.9 \text{ mA h g}^{-1}$  and  $132.7 \text{ mA h g}^{-1}$ , respectively), implying a clear contribution of oxygen to the charge compensation process.

To better understand the charge compensation mechanism at the structural and electronic level, a semi-simultaneous *operando* XRD and XANES study was conducted during the first electrochemical cycle. Fig. 8a displays the XRD data, giving us insight on the material's structural evolution, while Fig. 8b exhibits the XANES spectra at the Ni K edge, informing us on the Ni redox behavior. We chose to study the LNT055 sample, which showed the best electrochemical performances. The *operando* XRD results (Fig. 8a) didn't show noticeable phase changes in the crystal structure during both charge and discharge, indicating that a single rock salt phase is well preserved along the first cycle as shown by the evolution of the 200 and 220 reflections. To further investigate the structural parameters, we performed a Le Bail fit of the acquired patterns, whose results are displayed in Fig. 8c and d. When charging to

4.36 V, the unit cell parameters revealed a shift towards higher angular values, consistent with a unit cell shrinkage due to oxidation of Ni ( $\text{Ni}^{3+}$  being smaller than  $\text{Ni}^{2+}$ ). Notably, from around 4.36 V to 4.8 V, the unit cell parameter variation becomes less pronounced, suggesting the structural evolution is also less pronounced at higher voltage.<sup>47</sup> During discharge to 2.0 V the cell parameters shift back to lower angles, consistent with a unit cell expansion, attributed to the elongated Ni–O bond distance, then followed by a steep variation at the end of discharge at 1.5 V. It is interesting to notice here the change in structural parameters, *i.e.* after one cycle the unit cell parameter is larger than in the pristine state, pointing to a somewhat irreversible behavior. The lattice evolution as a function of the composition is shown in Fig. 8d, where a strong structural hysteresis is observed and a bimodal linear trend is observed in different voltage regions, presumably coinciding with the regions where nickel and oxygen are redox active. To corroborate the structural observation, with XANES at the Ni K-edge we observed directly the Ni oxidation state. These data are displayed in Fig. 8e and reveal a similar bimodal linear behavior as in the XRD data, namely there is a linear trend with two different slopes in different regions of the charge/discharge (linear fits are provided in Fig. S4 and S5†). The evaluation of

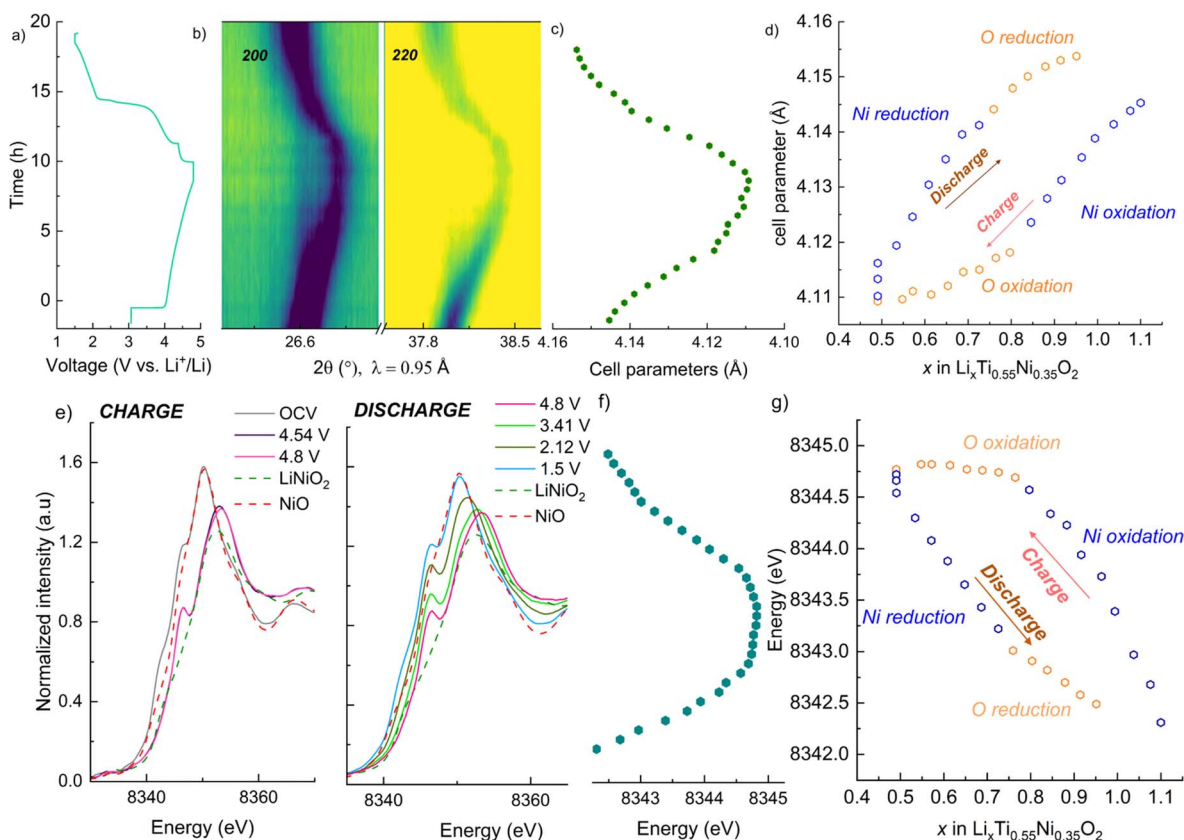


Fig. 8 Semi-simultaneous *operando* synchrotron XRD and XANES of the sample LNT055 during the initial cycle, with charge to 4.8 V and discharge to 1.5 V. (a) Voltage vs. time during the *operando* XRD/XANES experiment. (b) Contour plot showing the evolution of the 200 and 220 rock salt reflections. (c) consequent cell parameters evolution in charge and discharge, obtained from Le Bail fit. (d) Unit cell parameters vs. Li content. (e) Ni K-edge spectra collected during the initial cycle during charge (left) and discharge (right). (f) Observed change in the K-edge energy, obtained tracking the white line position. (g) Evolution of the K-edge energy vs. Li content. In (d) and (g), the choices of the points belonging to the O and Ni regions is given by the highest  $R^2$  observed in the linear fit (Fig. S4 and S5†).



Ni oxidation state was performed by determining the position of the white line for each spectrum. In Fig. 8e, the standard spectra of  $\text{LiNiO}_2$  and  $\text{NiO}$  are also added for a qualitative comparison. During charge, a noticeable edge shift to higher energy was observed from the pristine state of 3.0 V to a charged state at 4.8 V, indicating Ni oxidation from 2+ to 3+ and even beyond.<sup>48</sup> On the other hand, during discharge the edge undergoes a shift back towards lower energy values when Ni is reduced to 2+ valence state, demonstrating a good reversibility of the Ni redox process. Fig. 8d summarizes the observed change in cell parameters in charge and discharge states showing a similarity with the Ni K-edge evolution (Fig. 8g). As already mentioned in the introduction, the findings of Li *et al.*<sup>28</sup> helped us to rationalize the peculiar behavior detected pointing out an explanation by the unusual Ni and O redox mechanism observed. A kinetically activated  $\text{Ni}^{3+/4+}$  intermediate state is formed in charge and it triggers the activation of the oxygen oxidation, observed at the top of charge, reflected by the barely noticeable variation in cell parameters and K-edge energy. In discharge, contrary to expectations, the  $\text{Ni}^{3+/4+}$  is reduced before the oxygen (as evidenced by the severe edge shift) due to his high instability which leads to an inversion of cationic and anionic redox during discharge.

## Discussion

In our investigation we have shown for the first time how in these series of materials the composition, the temperature and consequently the configurational entropy played a strong role in reaching a pure disordered rock salt structure, to the disadvantage of the layered  $\text{Li}_2\text{TiO}_3$ -like phase encountered at lower synthesis temperature. Therefore whether one writes these compounds as a solid solution or as a biphasic mixture also strongly depends on the synthesis route. Unfortunately, high temperature does not only lead to single phase rock salt materials, but also to crystallites growth. Here we show that already at 700 °C the samples display kinetic limitations, therefore it is unlikely that using high synthesis temperature could be of any use to improve electrochemical performances.

Our study also further demonstrated that the variation of the compositions from a stoichiometric lithium amount to a lithium rich one is necessary to create Li percolation pathways and enhance the electrochemical performances. Nonetheless, even when the synthesis temperature is kept low to avoid crystallites growth, Li-rich samples are made and the cycling temperature is raised to 50 °C, one can obtain high specific capacity but cannot avoid the huge voltage hysteresis. This is terribly detrimental to the energy density of the cathode material as the average voltage is more than 1 V lower in discharge than in charge, determining an extremely poor energetic efficiency. Our simultaneous *operando* XRD and XANES results clearly confirmed the results of Li *et al.*, showing that this behavior is intrinsic to the coupled Ni–O charge compensation in these materials and likely hard to avoid. Nonetheless, strategies to change the kinetics of Ni–O charge transfer, by doping or other means, should be the primary objective to reduce the voltage hysteresis in this class of compounds.

More in general, Ni redox in DRX revealed high polarization that may be ascribed to the unfavourable Li diffusion kinetics, but also to significant voltage and structural hysteresis.<sup>8,49</sup> The detected voltage hysteresis is associated with anionic redox, which in this case mostly takes place asymmetrically between charge and discharge.<sup>17</sup> Interestingly, Ji *et al.* observed similar features during cycling when  $\text{Nb}^{5+}$  is used as  $d^0$  element. In fact, the first-charge voltage showed a clear plateau at  $\sim 4.3$  V attributed to the anionic redox contribution where accordingly to their XANES results, Ni oxidation barely is responsible to the observed capacity, suggesting sluggish oxygen oxidation as a charge compensation mechanism.<sup>50</sup> It's worth mentioning that although Li enrichment in the structure and the formation of the Li–O–Li configuration are considered to be the necessary condition to activate anionic redox, Fukuma *et al.*<sup>51</sup> pointed out that the structural disorder and the formation of a linear Ni–O–Ni configuration can also contribute to the anionic redox.

On the other hand, other studies revealed that when  $\text{Ru}^{5+}$  is exploited as charge compensation element, significantly lower hysteresis is observed, and the anionic redox chemistry of  $\text{O}^{2-}/\text{O}^-$  redox is identified at high potential after Ni oxidation, leading to a resulting capacity which rely in both cathodic and anionic redox processes separately but reversibly.<sup>52</sup> Therefore the choice of charge-compensating element, and its electronic structure ( $\text{Ru}^{5+}$  being not a  $d^0$  element), as well as the possibly induced short range-order in the materials, are all key elements to be considered for the future design of DRXs that are based on Ni redox but are not plagued by significant voltage hysteresis.

## Conclusions

In conclusion, we investigated a family of Li–Ni–Ti–O rock salt compounds with  $\text{Li}_{2y}\text{Ti}_y\text{Ni}_{2-3y}\text{O}_2$  as target composition. We evidenced a two-phase system (disordered rock salt and a small amount of  $\text{Li}_2\text{TiO}_3$ -like phase) rather than a solid solution unless the temperature is higher than 700 °C. Increasing the Li and Ti content increases the amount of  $\text{Li}_2\text{TiO}_3$ -like secondary phase. Our findings helped rationalize how the effect of the temperature and composition can play a key role in achieving a pure cubic crystal structure, thanks to DFT computations combined with monitoring the structural evolution during synthesis by *in situ* XRD. The results showed the reflection of the  $\text{Li}_2\text{TiO}_3$ -like phase forming first, as expected from a more energetically favoured phase, and its gradual disappearance by further increasing the temperature promoting the formation of the rock salt structure due to its larger configurational entropy. The compounds, whom we have verified to have the expected stoichiometry, have been tested electrochemically and exhibit large voltage hysteresis and low specific capacity at room temperature, which can however be enhanced by increasing the temperature, pointing to clear kinetic limitations of the reaction in this material. The simultaneous *operando* XRD and XANES study conducted on the  $\text{Li}_{1.1}\text{Ni}_{0.35}\text{Ti}_{0.55}\text{O}_2$  sample highlights how the rock salt phase is maintained stable along the first cycle. We also tracked the behaviour of the cell parameters and Ni K-edge position, which confirmed that the O redox mechanism plays a role in the charge compensation and evidenced the



clear presence of a structural hysteresis as well, on top of the poor Li diffusion kinetics. At the moment, such hysteresis given by the peculiarity of the electronic distribution in the oxidized Ni–O bond, and the ensuing voltage one, are the main bottleneck to the exploitation of these materials. To change this behavior, a vast chemical space of high-energy cathodes can be opened up by cation-doping, exploiting a mix of redox TM and by substituting  $Ti^{4+}$  with other high-valent non-redox-active ions (which could be  $d^{10}$  or  $d^0$ , or even  $d^3$ ).<sup>53</sup> On the other hand, fluorine substitution can be also performed by partially replacing  $O^{2-}$  ions with  $F^-$  ions allowing to decrease the total anionic charge and expand the compositional space, as already extensively demonstrated.<sup>16</sup> This strategy can also improve the electrochemical performance by maximizing the TM redox amount.<sup>54</sup>

## Conflicts of interest

There are no conflicts to declare.

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