

Cite this: *Sustainable Energy Fuels*,  
2023, 7, 3239Received 2nd February 2023  
Accepted 5th June 2023

DOI: 10.1039/d3se00130j

rsc.li/sustainable-energy

# Core-shell Fe<sub>3</sub>O<sub>4</sub>@CoFe<sub>2</sub>O<sub>4</sub> nanoparticles as high-performance anode catalysts for enhanced oxygen evolution reaction†

Lisa Royer,<sup>ab</sup> Iryna Makarchuk,<sup>b</sup> Simon Hettler,<sup>cd</sup> Raul Arenal,<sup>cde</sup> Tristan Asset,<sup>a</sup> Benjamin Rotonelli,<sup>a</sup> Antoine Bonnefont,<sup>f</sup> Elena Savinova<sup>a</sup> and Benoit P. Pichon<sup>ib\*bg</sup>

Water electrolysis is a promising and environmentally friendly means for renewable energy storage. Recent progress in the development of anion exchange membranes (AEMs) has provided new perspectives for high-performance anode catalysts based on transition metal oxides (TMOs) for the sluggish anodic oxygen evolution reaction (OER). Here, we report on core-shell nanoparticles (Fe<sub>3</sub>O<sub>4</sub>@CoFe<sub>2</sub>O<sub>4</sub>) which allow combining an electrocatalytic shell (CoFe<sub>2</sub>O<sub>4</sub>) with a conductive core (Fe<sub>3</sub>O<sub>4</sub>). Such an original approach significantly minimizes critical Co content in the catalyst and avoids addition of unstable conductive carbon black. The core-shell nanoparticles outperform Co<sub>(1-x)</sub>-Fe<sub>(2+x)</sub>O<sub>4</sub> nanoparticles and show an exceptional OER activity per Co unit mass (2800 A g<sub>Co</sub>at<sup>-1</sup> at 1.65 V vs. RHE) suggesting synergistic interaction between the core and the shell. Along with the core-shell structure, the size of the Fe<sub>3</sub>O<sub>4</sub> core is a critical parameter, with a large conductive Fe<sub>3</sub>O<sub>4</sub> core being beneficial for OER enhancement.

Water electrolysis provides an environmentally friendly way to store renewable energies through H<sub>2</sub> production. While proton exchange membrane water electrolysis is well suited to renewable energy storage, it relies on the utilization of scarce and expensive iridium (Ir) to accelerate the sluggish kinetics of the

anodic oxygen evolution reaction (OER). Thanks to the recent progress in the development of anion-exchange membranes (AEMs), transition metal oxides (TMOs) appear as promising substitutes for noble metals.<sup>1-4</sup> Various approaches have been proposed to enhance their OER activity, notably varying the composition, particle size, and defect concentration, or, recently, forming core-shell structures.<sup>5-9</sup>

Owing to their composition-dependent and widely tunable properties, TMOs with spinel structures appear very attractive for various energy-related applications. While Co<sub>3</sub>O<sub>4</sub> and CoOOH thin films deposited on an Au(111) substrate have demonstrated significant surface-weighted OER activity,<sup>10</sup> core-shell nanoparticles (NPs) combining a gold core and a TMO shell reached 10 000 A g<sub>TM</sub><sup>-1</sup> at 1.65 V vs. RHE. According to the authors, such a high OER activity of TMOs may result from an electronic effect of the gold core that favors the oxidation of the TM up to the 4+ state.<sup>11,12</sup> The use of core-shell nanoparticles might be an efficient way to reduce the amount of expensive and rare electrocatalysts by circumscribing them to a thin shell,<sup>13,14</sup> but the exclusive use of cost-effective and abundant materials is the next challenge for the worldwide dissemination of water electrolysis.

TMO NPs are of significant interest since the catalytic activity can be greatly enhanced by adapting their size and shape, *i.e.* the active surface area. The valence of the metallic cations and their chemical environment are also critical to enhancing the OER activity.<sup>3,15,16</sup> Beyond the fine control of these parameters, recent advances in nanoparticle synthesis in liquid media avoid aggregation, thus also favoring high surface area, *i.e.* enhanced mass-weighted catalytic activity. Consequently, the challenge lies in controlling the formation and the chemical structure of NPs in order to systematically study their electrochemical properties.

Although most TMOs are insulating, magnetite (Fe<sub>3</sub>O<sub>4</sub>) is conductive due to electron hopping between Fe(II) and Fe(III).<sup>17</sup> Therefore, it may be considered a promising support material to avoid carbon black, which is often added to circumvent the lack of conductivity of TMO catalysts, but is unstable under anode

<sup>a</sup>ICPEES, UMR 7515 CNRS-ECPM-Université de Strasbourg, 25, rue Becquerel, F 67087, Strasbourg Cedex 2, France

<sup>b</sup>Université de Strasbourg, CNRS, Institut de Physique et Chimie des Matériaux de Strasbourg, UMR 7504, F-67000, Strasbourg, France. E-mail: benoit.pichon@unistra.fr

<sup>c</sup>Instituto de Nanociencia y Materiales de Aragon (INMA), CSIC-Universidad de Zaragoza, Calle Pedro Cerbuna 12, 50009, Zaragoza, Spain

<sup>d</sup>Laboratorio de Microscopías Avanzadas (LMA), Universidad de Zaragoza, Calle Mariano Esquillor, 50018, Zaragoza, Spain

<sup>e</sup>ARAID Foundation, 50018, Zaragoza, Spain

<sup>f</sup>Institut de Chimie, UMR 7177, CNRS-Université de Strasbourg, 4 rue Blaise Pascal, CS 90032, 67081, Strasbourg Cedex, France

<sup>g</sup>Institut Universitaire de France, 5 rue Descartes, 75015, Paris, France

† Electronic supplementary information (ESI) available: Nanoparticle synthesis, TEM images, XRD patterns, cyclic voltammetry curves, impedance measurements, EELS and STEM HAADF micrographs. See DOI: <https://doi.org/10.1039/d3se00130j>



operation conditions. Furthermore, iron cations in TMO phases may result in synergistic effects with the surrounding atoms, thus enhancing the OER.<sup>18,19</sup> Hence, nanostructures combining  $\text{Fe}_3\text{O}_4$  as a conductive core and a TMO as a catalyst shell would be promising high performance noble metal-free OER catalysts.

Here we report an original approach to design a new type of carbon-free, noble-metal-free nanostructured material for the OER. The core-shell structure of  $\text{Fe}_3\text{O}_4@ \text{CoFe}_2\text{O}_4$  NPs allows combining a conductive magnetite core and a catalytically active cobalt ferrite shell. Co being classified as a critical raw material by the European Union, such a design of NPs with a core-shell structure efficiently contributes to reducing the amount of Co used for the OER. Such a cost-effective material resulted in an excellent OER activity per unit mass of Co. The fine control of the core size and the shell thickness led us to investigate the influence of the core-shell structure on the OER activity which is markedly enhanced by a larger  $\text{Fe}_3\text{O}_4$  core.

Core-shell NPs were synthesized by the thermal decomposition of metal complexes in high-temperature-boiling solvents (around 300 °C) as we recently reported<sup>20</sup> (see the ESI† for details). Two different approaches were applied: (i) the seed-mediated growth of a  $\text{CoFe}_2\text{O}_4$  shell onto the surface of pristine  $\text{Fe}_{3-\delta}\text{O}_4$  nanoparticles which led to samples CS-1, CS-3, and CS-4. (ii) The diffusion of cobalt cations into vacancies at the surface of partially oxidized pristine  $\text{Fe}_{3-\delta}\text{O}_4$  nanoparticles which led to sample CS-2.<sup>21</sup> The core size was varied by adjusting the experimental conditions upon synthesis of  $\text{Fe}_{3-\delta}\text{O}_4$  nanoparticles as we reported earlier.<sup>22,23</sup> The shell thickness was modified by adjusting the amount of Co and Fe metal complexes with a Co/Fe molar ratio of 0.5. An increase in the amount of Co and Fe complexes by 2.5 from CS-3 to CS-4 resulted in a two-fold shell volume increase which agrees with the larger size of pristine  $\text{Fe}_{3-\delta}\text{O}_4$  nanoparticles used for CS-4. The transmission electron microscopy (TEM) images showed faceted-spherical shapes and narrow size distributions both for pristine and for core-shell NPs (Fig. 1). Granulometry measurements showed stable colloidal suspensions of non-aggregated NPs thanks to the efficient coating of oleic acid which was added to the reaction medium (Fig. S3†).

Energy-dispersive X-ray (EDX) spectroscopy confirmed the presence of Co between 12 and 19 at% with respect to Fe in core-shell NPs (Table S1†). For CS-1, CS-2, and CS-3, these values are much higher than those corresponding to the size variation between pristine and core-shell NPs (Table S1†) which is indicative of the diffusion of cobalt in the  $\text{Fe}_{3-\delta}\text{O}_4$  core as we have reported earlier.<sup>24</sup> X-ray diffraction (XRD) patterns recorded for core-shell NPs display peaks that are all indexed to the spinel crystal structure (Fig. S4†). Although  $\text{Fe}_3\text{O}_4$  and  $\text{CoFe}_2\text{O}_4$  phases cannot be distinguished because of similar cell parameters (8.396 Å and 8.3919 Å, respectively),<sup>20</sup> cell parameters of core-shell NPs are relatively close to those of  $\text{Fe}_3\text{O}_4$  (Table S2†). This is ascribed to the low extent of the surface oxidation of pristine nanoparticles which are protected by the  $\text{CoFe}_2\text{O}_4$  shell when exposed to air. The high crystallinity was confirmed by lattice fringes all-across core-shell NPs as observed in scanning transmission electron microscopy (STEM) micrographs

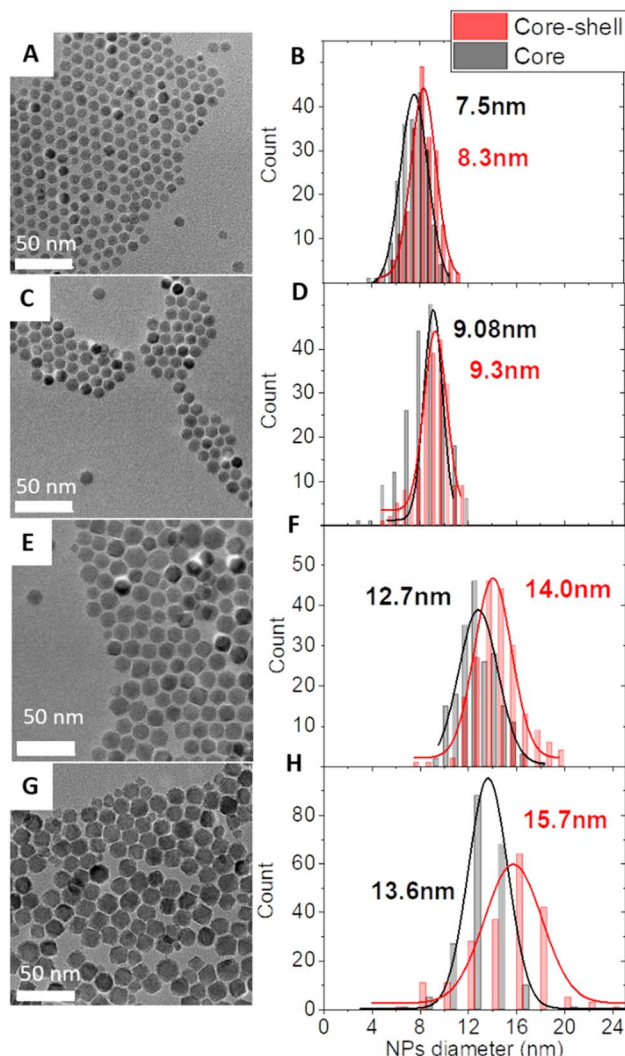


Fig. 1 TEM micrographs of core-shell NPs: CS-1 (A), CS-2 (C), CS-3 (E), and CS-4 (G) and the corresponding size distributions (B, D, F and H) of the core-shell NPs and their pristine core NPs.

recorded in the high angle annular dark field (HAADF) mode (Fig. S6†).

The local composition of these NPs was further investigated by electron energy-loss spectroscopy (EELS) analysis in STEM (Fig. 2 and S6†). Spatially resolved EELS maps show the homogeneous distribution of O in NPs. In contrast, the Co content is the highest at the edge while the Fe content is at its lowest. Cross-section profiles show that the Co fraction reaches a maximum value of 15–20% with respect to Fe and O on the edge while it quickly goes down to 5% (CS-2 and CS-4) and 2% (CS-3) in the center of the NP, in agreement with a much thinner shell than the core size. Considering stoichiometric  $\text{CoFe}_2\text{O}_4$  (14% at. Co) and the resolution of the measurement (5 Å), Co is certainly distributed as a graded concentration ( $\text{Co}_{1+x}\text{Fe}_{2-x}\text{O}_4$ ) which decreases from the NP surface to the center. This result is supported by the variation of the Co content over a longer distance from the surface to the center than the size variation usually ascribed to the shell thickness calculated from TEM





**Fig. 2** Scanning transmission electron microscopy (STEM) high angle annular dark field (HAADF) images of CS-2 (A–C) and CS-3 (D–F). Spatially resolved electron energy loss spectroscopy (EELS) analyses (B and E) showing the spatial distributions of oxygen O (red), iron Fe (green) and cobalt Co (blue). Quantitative cross-section profiles (C and F) recorded along the arrows depicted in (B and E) respectively. The left y-scale refers to percentages of Fe, Co and O atoms. (G) Fe atomic fraction (*i.e.* the ratio of Fe/(Co + Fe)) calculated from the depth profiling analysis performed for different incident energies on CS-1 and CS-3, the photon energy of the X-ray source was varied from 850 to 1350 eV in order to analyse Co 2p and Fe 2p peaks. Schematic illustration of the nanoparticle depth (H) probed as a function of the inelastic mean free path ( $\lambda$ ).

micrographs. A slight increase in the O content can also be observed at the edges of nanoparticles which can be correlated to some cationic vacancies and surface defects as usually observed for such nanoparticles.

Complementary information on the average near-surface composition of NPs was obtained by performing X-ray photoelectron spectroscopy (XPS) measurements. The variation of the photon energy from 850 to 1350 eV allowed changing the inelastic mean-free path of the photoelectron emission ( $\lambda$ ) from 0.6 to 1.3 nm thus allowing different NP depths to be probed. Considering that 68% of the emitted photoelectrons arise from a thickness  $\lambda$  and 99% emerge from  $3\lambda$ ,<sup>25,26</sup> the Fe atomic fraction was calculated as the concentration of Fe normalized to the sum of Fe and Co from the area of Co 2p and Fe 2p peaks which were recorded for CS-1 and CS-3 (Fig. 2G). The Fe atomic fraction at small photon energies (hence short  $\lambda$ ) is characteristic of the shell, but then gradually increases, since the contribution of the  $\text{Fe}_3\text{O}_4$  core to the XP spectra increases with the photon energy. It is interesting to note that the Fe atomic fraction (Fe/(Fe + Co)) in the particle shell ( $\sim 0.5$ ) is higher than stoichiometry (0.66 for  $\text{CoFe}_2\text{O}_4$ ), which may be ascribed to the decomposition of Co stearate which happens at a higher temperature than that of Fe stearate.<sup>27</sup> According to the La Mer theory, iron oxide starts growing before cobalt oxide at the surface of pristine NPs. This results in a gradient in the Co concentration in the shell. Nevertheless, both decomposition temperatures (of the iron and cobalt stearate) are sufficiently close to avoid phase segregation since no CoO phase was observed in the XRD patterns<sup>24</sup> (Fig. S4†).

The electrochemical properties of core-shell NPs were studied in order to determine their OER activity and compare

it to the activity of Co ferrite ( $\text{Co}_{(1-x)}\text{Fe}_{(2+x)}\text{O}_4$ ) NPs (for characterization see the ESI†). Cyclic voltammograms (CVs) were recorded after deposition of a thin layer of NPs on a glassy carbon (GC) electrode. To accurately determine the intrinsic electrocatalytic activity and avoid artefacts arising from charge and mass transport effects associated with thick NP layers (see *e.g.* ref. 28) the NP loading was set between  $\sim 1$  and  $4 \mu\text{g cm}^{-2}$  (for more details regarding the influence of the NP loading on the apparent activity the reader is referred to ref. 29). In order to compare our results with the literature, current-potential curves were normalized in three different ways: to the mass of metal oxide, to the mass of cobalt (Co being considered a ‘strategic’ element), and to the NP surface area (estimated from the TEM mean diameter), see Fig. 3A. Considering the low NP loading on the electrode, the mass- and surface-weighted OER activities were measured at a relatively high potential of 1.65 V *vs.* RHE (Fig. 3D–F). While all core-shell NPs exhibit excellent OER activities (much exceeding the  $28 \text{ A g}_{\text{oxide}}^{-1}$  activity of the  $\text{Fe}_3\text{O}_4$  NPs of 10 nm), the largest CS-4 and CS-3 NPs display Co-mass weighted-activities (2300 and 2800  $\text{A g}_{\text{cobalt}}^{-1}$ , respectively) which are twice those of CS-1 and CS-2 (about 1300  $\text{A g}_{\text{cobalt}}^{-1}$ ). This trend is even clearer when considering the surface-weighted activity which has increased 3–4-fold. It is noteworthy that both the mass- and the surface-weighted OER activity of the best (CS-3) core-shell NPs is significantly higher than that of the  $\text{Co}_{(1-x)}\text{Fe}_{(2+x)}\text{O}_4$  ferrite NPs synthesized through a similar protocol (see the ESI†). Hence, the OER activity of NPs unambiguously arises not only from the cobalt ferrite shell but also from the combination of  $\text{Fe}_3\text{O}_4$  and  $\text{CoFe}_2\text{O}_4$  into a core-shell structure. The difference in activity between CS-1





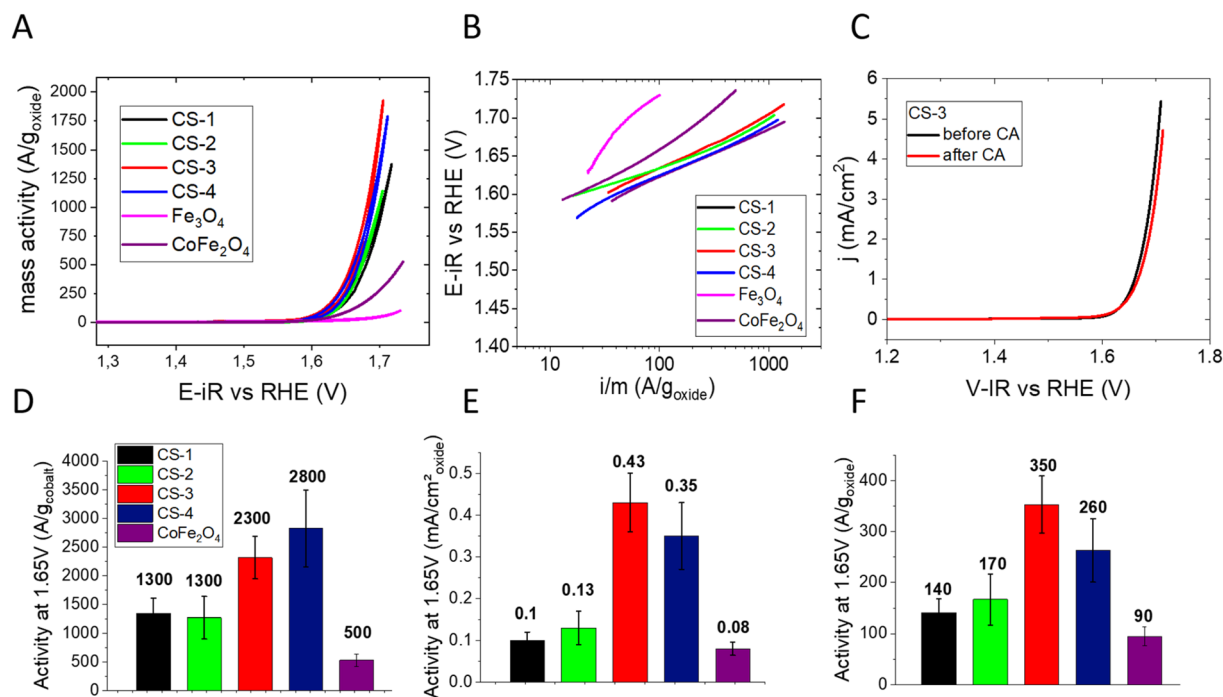


Fig. 3 OER activity in 0.1 M NaOH: (A) cyclic voltammograms performed on CS-1, CS-2, CS-3, CS-4 and Co<sub>(1-x)</sub>Fe<sub>(2+x)</sub>O<sub>4</sub> nanoparticles deposited on a glassy carbon electrode (oxide loading between ~1 and 4  $\mu\text{g cm}^{-2}$ , see the ESI<sup>†</sup>), scan rate 10  $\text{mV s}^{-1}$ . The electrode potential is ohmic drop-corrected, and the current is divided by the oxide mass. (B) Tafel slopes of these NPs (in  $\text{mV decade}^{-1}$ ). (C) Positive scans of CVs acquired for CS-3 NPs before (black) and after chronoamperometry measurements at 1.65 V vs. RHE for 3 hours. NP loading: 1.3  $\mu\text{g cm}^{-2}$  and scan rate: 10  $\text{mV s}^{-1}$ . Panels (D)–(F) compare OER activity of core–shell and Co<sub>(1-x)</sub>Fe<sub>(2+x)</sub>O<sub>4</sub> NPs at 1.65 V vs. RHE (extracted from CVs) per gram of cobalt (D), per  $\text{cm}^2$  of oxide (E) and per gram of oxide (F).

and CS-2 vs. CS-3 and CS-4 likely originates from the core size (7.5–9.1 vs. 12.7–13.6 nm). A larger core is likely to be essential for efficient OER since it results in a higher fraction of conductive magnetite.<sup>30,31</sup> This hypothesis is supported by the low activity of Fe<sub>3</sub>O<sub>4</sub>@CoO, consisting of a 4.5 nm core.<sup>32</sup> Indeed, our recent *operando* soft X-ray absorption spectroscopy data suggest that the synergy results from cooperative redox interaction between Fe(II) from the core and Co(II) from the shell.<sup>33</sup> Additionally, we also showed that the OER activity of such Fe<sub>3</sub>O<sub>4</sub>@CoFe<sub>2</sub>O<sub>4</sub> NPs is significantly enhanced by a thicker CoFe<sub>2</sub>O<sub>4</sub> shell although the resistivity of the NP layer increases.<sup>29</sup> It seems that the insulating properties of Co ferrite are compensated for by a larger number of active sites on the NP surface. It is worth noting that the Co-weighted activity of the most active Fe<sub>3</sub>O<sub>4</sub>@CoFe<sub>2</sub>O<sub>4</sub> core–shell NPs (CS-3) greatly outperforms that of TMO catalysts, either with uniform or core–shell structures, which were reported in the literature (Table S4<sup>†</sup>). Indeed only CoO<sub>x</sub>(OH)<sub>y</sub> NPs<sup>34</sup> and core–shell NPs with a gold core<sup>6,8</sup> (Au@CoFeO<sub>x</sub> and Au@NiCo<sub>2</sub>S<sub>4</sub>) display higher activity although they are unstable when mixed with carbon black under the OER operation.

With Tafel slopes being often used for mechanistic interpretations, one should note that they are often affected by ohmic losses in catalytic layers.<sup>29</sup> Indeed, Tafel slopes increase with the electrode potential (Fig. 3B), notably for the Co<sub>(1-x)</sub>Fe<sub>(2+x)</sub>O<sub>4</sub> NPs. The Tafel slopes were extracted from the anodic scans of CVs at two electrode potentials: 1.62 (where

the OER just started) and 1.67 V vs. RHE (where the OER fully began). At the lower potential, all core–shell NPs display Tafel slopes in the range of 50 to 65  $\text{mV decade}^{-1}$  which is slightly less than that for Co<sub>(1-x)</sub>Fe<sub>(2+x)</sub>O<sub>4</sub> and much lower than those of the OER-inactive Fe<sub>3</sub>O<sub>4</sub> NPs (Table S3<sup>†</sup>). The increase in the Tafel slope at higher potentials may be related to the ohmic resistance across the catalytic layer.<sup>29</sup> At the higher potential, the Tafel slope of the Co<sub>(1-x)</sub>Fe<sub>(2+x)</sub>O<sub>4</sub> particles exceeds the one for core–shell NPs, which is in agreement with the lower conductivity of the cobalt ferrite NPs. Finally, the medium-term stability of core–shell CS-3 NPs was studied using chronoamperometry at 1.65 V vs. RHE. Fig. 3C shows that anodic CV scan after 3 hours of polarization is within the reproducibility limit ( $\pm 10\%$ ).

To sum up, Fe<sub>3</sub>O<sub>4</sub>@CoFe<sub>2</sub>O<sub>4</sub> core–shell NPs with a narrow size distribution and uniform shape were successfully synthesized through two alternative approaches (diffusion and crystal growth). They present excellent OER activities per unit mass of Co (from 1300 to 2800  $\text{A g}_{\text{cobalt}}^{-1}$ ) which greatly outperform the OER activities of transition metal oxide nanoparticles reported in the literature. Along with the core–shell structure, the size of the Fe<sub>3</sub>O<sub>4</sub> conductive core and the thickness of the Co-containing shell are critical parameters for efficient OER activity. Thus, synthesizing TMO based core–shell nanostructures with fine control of the chemical structure is a viable approach for high performance alkaline OER electrocatalysts.



## Conflicts of interest

The authors declare no competing financial interests.

## Acknowledgements

This project was financially supported by the Foundation for Frontier Research in Chemistry. We thank the Helmholtz-Zentrum Berlin für Materialien und Energie for the allocation of synchrotron radiation beamtime. The authors acknowledge Juan Velasco Velez for his help on the synchrotron beamline as well as Vinavadini Ramnarain, Barbara Freis and Théo Lucante for the TEM images. STEM and EELS measurements have been performed at the Laboratorio de Microscopías Avanzadas (LMA) at the Universidad de Zaragoza (Spain). I. M. acknowledges the funding from IdEx Université de Strasbourg. R. A. and S. H. acknowledge the funding from the Spanish MICINN (project grant PID2019-104739GB-I00/AEI/10.13039/501100011033), from the Government of Aragon (project DGA E13-20R) and from the European Union H2020 program “ESTEEM3” (823717).

## References

- 1 C. Feng, M. B. Faheem, J. Fu, Y. Xiao, C. Li and Y. Li, *ACS Catal.*, 2020, **10**, 4019–4047.
- 2 M.-I. Jamesh and X. Sun, *J. Power Sources*, 2018, **400**, 31–68.
- 3 J. S. Kim, B. Kim, H. Kim and K. Kang, *Adv. Energy Mater.*, 2018, **8**, 1702774.
- 4 N.-T. Suen, S.-F. Hung, Q. Quan, N. Zhang, Y.-J. Xu and H. M. Chen, *Chem. Soc. Rev.*, 2017, **46**, 337–365.
- 5 S.-C. Chou, K.-C. Tso, Y.-C. Hsieh, B.-Y. Sun, J.-F. Lee and P.-W. Wu, *Materials*, 2020, **13**(12), 2703.
- 6 A. L. Strickler, M. Escudero-Escribano and T. F. Jaramillo, *Nano Lett.*, 2017, **17**, 6040–6046.
- 7 Z. Luo, S. Martí-Sánchez, R. Nafria, G. Joshua, M. de la Mata, P. Guardia, C. Flox, C. Martínez-Boubeta, K. Simeonidis, J. Llorca, J. R. Morante, J. Arbiol, M. Ibáñez and A. Cabot, *ACS Appl. Mater. Interfaces*, 2016, **8**, 29461–29469.
- 8 Y. Lv, S. Duan, Y. Zhu, P. Yin and R. Wang, *Nanomaterials*, 2020, **10**, 611.
- 9 Z. Zhuang, W. Sheng and Y. Yan, *Adv. Mater.*, 2014, **26**, 3950–3955.
- 10 F. Reikowski, F. Maroun, I. Pacheco, T. Wiegmann, P. Allongue, J. Stettner and O. M. Magnussen, *ACS Catal.*, 2019, **9**, 3811–3821.
- 11 Y. Gorlin, C.-J. Chung, J. D. Benck, D. Nordlund, L. Seitz, T.-C. Weng, D. Sokaras, B. M. Clemens and T. F. Jaramillo, *J. Am. Chem. Soc.*, 2014, **136**, 4920–4926.
- 12 B. S. Yeo and A. T. Bell, *J. Am. Chem. Soc.*, 2011, **133**, 5587–5593.
- 13 L. Gloag, T. M. Benedetti, S. Cheong, R. F. Webster, C. E. Marjo, J. J. Gooding and R. D. Tilley, *Nanoscale*, 2018, **10**, 15173–15177.
- 14 H. N. Nong, L. Gan, E. Willinger, D. Teschner and P. Strasser, *Chem. Sci.*, 2014, **5**, 2955–2963.
- 15 D. Guo, H. Kang, P. Wei, Y. Yang, Z. Hao, Q. Zhang and L. Liu, *CrystEngComm*, 2020, **22**, 4317–4323.
- 16 F. Song, L. Bai, A. Moysiadou, S. Lee, C. Hu, L. Liardet and X. Hu, *J. Am. Chem. Soc.*, 2018, **140**, 7748–7759.
- 17 A. I. Dmitriev, S. I. Alekseev and S. A. Kostyuchenko, *Inorg. Mater.*, 2019, **55**, 576–581.
- 18 S. Anantharaj, S. Kundu and S. Noda, *Nano Energy*, 2021, **80**, 105514.
- 19 M. Li, Y. Gu, Y. Chang, X. Gu, J. Tian, X. Wu and L. Feng, *Chem. Eng. J.*, 2021, **425**, 130686.
- 20 K. Sartori, A. Musat, F. Choueikani, J.-M. Grenèche, S. Hettler, P. Bencok, S. Begin-Colin, P. Steadman, R. Arenal and B. P. Pichon, *ACS Appl. Mater. Interfaces*, 2021, **13**, 16784–16800.
- 21 X. Liu, B. P. Pichon, C. Ulhaq, C. Lefèvre, J.-M. Grenèche, D. Bégin and S. Bégin-Colin, *Chem. Mater.*, 2015, **27**, 4073–4081.
- 22 W. Baaziz, B. P. Pichon, S. Fleutot, Y. Liu, C. Lefevre, J.-M. Grenèche, M. Toumi, T. Mhiri and S. Begin-Colin, *J. Phys. Chem. C*, 2014, **118**, 3795–3810.
- 23 K. Sartori, D. Gailly, C. Bouillet, J.-M. Grenèche, P. Dueñas-Ramirez, S. Begin-Colin, F. Choueikani and B. P. Pichon, *Cryst. Growth Des.*, 2020, **20**, 1572–1582.
- 24 K. Sartori, G. Cotin, C. Bouillet, V. Halté, S. Bégin-Colin, F. Choueikani and B. P. Pichon, *Nanoscale*, 2019, **11**, 12946–12958.
- 25 B. R. Strohmeier, *Surf. Interface Anal.*, 1990, **15**, 51–56.
- 26 T. A. Carlson and G. E. McGuire, *J. Electron Spectrosc. Relat. Phenom.*, 1972, **1**, 161–168.
- 27 W. Baaziz, B. P. Pichon, Y. Liu, J.-M. Grenèche, C. Ulhaq-Bouillet, E. Terrier, N. Bergéard, V. Halté, C. Boeglin, F. Choueikani, M. Toumi, T. Mhiri and S. Begin-Colin, *Chem. Mater.*, 2014, **26**, 5063–5073.
- 28 D. Y. Chung, S. Park, P. P. Lopes, V. R. Stamenkovic, Y.-E. Sung, N. M. Markovic and D. Strmcnik, *ACS Catal.*, 2020, **10**, 4990–4996.
- 29 L. Royer, J. Guehl, M. Zilbermann, T. Dintzer, C. Leuvrey, B. P. Pichon, E. Savinova and A. Bonnefont, *Electrochim. Acta*, 2023, **446**, 141981.
- 30 M. Pauly, J.-F. Dayen, D. Golubev, J.-B. Beaufrand, B. P. Pichon, B. Doudin and S. Bégin-Colin, *Small*, 2012, **8**, 108–115.
- 31 G. Rydzek, D. Toulemon, A. Garofalo, C. Leuvrey, J.-F. Dayen, D. Felder-Flesch, P. Schaaf, L. Jierry, S. Begin-Colin, B. P. Pichon and F. Boulmedais, *Small*, 2015, **11**, 4638–4642.
- 32 L. Zhou, B. Deng, Z. Jiang and Z.-J. Jiang, *Chem. Commun.*, 2019, **55**, 525–528.
- 33 L. Royer, A. Bonnefont, T. Asset, B. Rotonelli, J.-J. Velasco-Vélez, S. Holdcroft, S. Hettler, R. Arenal, B. Pichon and E. Savinova, *ACS Catal.*, 2023, **13**, 280–286.
- 34 F. T. Haase, A. Bergmann, T. E. Jones, J. Timoshenko, A. Herzog, H. S. Jeon, C. Rettenmaier and B. R. Cuenya, *Nat. Energy*, 2022, **1**–9.

