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# Dual supramolecular chirogenesis based on platinum(II) metallotweezers†

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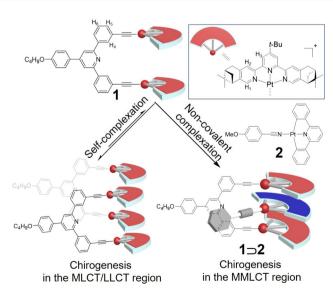
Optically active platinum(II) metallotweezers demonstrate both self-complexation and host-guest complexation capabilities, leading to two distinct supramolecular chirogenic signals in the visible region.

Chirality transfer from the molecular level to the supramolecular level is not only essential in life<sup>1</sup> but is relevant for catalytic, optoelectronic and spintronic applications in materials science.<sup>2</sup> Supramolecular chirogenesis<sup>3</sup> represents an efficient way to express chirality in artificial systems, and involves non-covalent chiral recognition between the host and guest species. Pioneering work in this field has been performed by Inoue and co-workers on the basis of achiral zinc porphyrin tweezers.<sup>3a,c</sup> Until now, supramolecular chirogenesis has extended to various artificial receptors such as macrocycles,<sup>4</sup> cages,<sup>5</sup> and helical foldamers.<sup>6</sup> Despite the progress achieved, the chirogenic signals appear primarily in the ultra-violet or high-energy visible region due to the following two reasons. One is the lack of large  $\pi$ -conjugated chromophores on host/ guest structures, and the other is the low chirality transfer efficiency because of the remoteness of the chiral center from the host-guest complexation site.<sup>7</sup> It is intriguing to shift supramolecular chirogenic signals to a low-energy absorption region, which would benefit circular polarized light detection/ emission and chiroptical switch applications.

Platinum( $\pi$ )-based metallotweezers,<sup>8</sup> with two cofacial square-planar pincers, represent an ideal candidate to attain this objective. When the cyclometalated Pt( $\pi$ ) pincers are kept

at the distance of 7 Å by a rigid spacer, the metallotweezers are capable of encapsulating a guest molecule into their cavity.<sup>9</sup> The non-covalent host–guest complexation structure is stabilized by  $\pi$ – $\pi$  stacking interactions between the Pt(II) pincers and the complementary guest (interplanar distance: ~ 3.5 Å). It can be endowed with fruitful photo-physical properties due to the spin–orbit coupling effect, leading to metal-to-ligand charge transfer (MLCT), ligand-to-ligand charge transfer (LLCT), and metal–metal-to-ligand charge transfer (MMLCT) transitions in the low-energy absorption region.<sup>10</sup> We envisage that supramolecular chirogenic signals could potentially emerge for these electronic transitions, by incorporating a stereogenic center in the receptor of the metallotweezers.

In this study, we have designed the novel Pt(n) metallotweezers **1** (Scheme 1, see Scheme S1 in the ESI<sup>†</sup> for the synthetic procedure). Unlike previous chiral Pt(n) complexes in which the



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stereogenic center was embedded in the side chains,<sup>11</sup> herein, four (1R)-pinene units are fused to the Pt(II) terpyridine pincers in 1 to strengthen the supramolecular chirogenic signals.<sup>12</sup> Interestingly, 1 is prone to associate with other molecules of 1 to form a self-complexed structure (Scheme 1). A stereospecific twist is generated because of the stacking of  $Pt(\pi)$ terpyridine  $[Pt(\pi)(N^N)]$  pincers, giving rise to the emergent chiroptical signals in the MLCT/LLCT absorption region. With the addition of compound 2 (Scheme 1) as the complementary guest, the self-complexation structure of 1 converts to the sandwich complex  $1 \supset 2$ . This consequently leads to the chirogenic signal in the MMLCT absorption region, thanks to the participation of Pt(II)-Pt(II) metal-metal interactions for the host-guest entity. Accordingly, dual supramolecular chirogenic signals form in the visible region, by taking advantage of the diverse complexation modes of the Pt(II) metallotweezers.

We first studied self-complexation behavior of the Pt(II) metallotweezers 1. In chloroform, protons H<sub>2</sub> and H<sub>3</sub> displayed downfield shifts upon varying the concentration from 0.20 mM to 20.0 mM (Fig. S12, ESI<sup>†</sup>). Concentration-dependent <sup>1</sup>H NMR measurements provided the self-association constant of  $5.34\,\times\,10^2~M^{-1}$  (±34%) for 1 (Fig. S13, ESI†). Generally, two possible self-aggregation modes exist for metallotweezers (Fig. S14, ESI<sup>+</sup>). One is the mutual stacking of pincer units in a quadruple manner.<sup>8e</sup> The other is sandwiching of the spacer unit into the cavity of the complementary tweezers.<sup>8f</sup> The latter mode is excluded in the case of metallotweezers 1, since the non-planar diphenylpyridine spacer is unable to be encapsulated into the cavity. The self-complexed structure of 1 was clarified via density functional theory (DFT) computations. Two quadruple stacking structures might form for 1, namely headto-tail and head-to-head binding modes for the neighbouring Pt(II)(N^N^N) pincers (Fig. S15, ESI<sup>+</sup>). For the optimized geometries, the head-to-head binding mode (Fig. 1a) featured a lower Gibbs free energy than that of the head-to-tail mode  $(\Delta E = 0.974 \text{ kcal mol}^{-1}, \text{ Fig. S15, ESI}^{\dagger})$ . The  $\pi$ - $\pi$  distances between the Pt(II)(N^N^N) pincers are 3.29 Å, 3.47 Å, and 3.29 Å, respectively. Apparently, the pre-organization effect of the rigid diphenylpyridine spacer, together with the strong stacking tendency of the Pt(II)(N^NN) pincers, guarantees formation of the self-complexation structure for 1. Formation of the headto-head binding structure was further demonstrated via <sup>1</sup>H-<sup>1</sup>H ROESY measurements. In particular, strong correlations exist between protons  $H_4/H_5$  and  $H_4/H_6$  (Fig. S16a, ESI<sup>†</sup>), which are absent in the <sup>1</sup>H-<sup>1</sup>H COSY spectrum (Fig. S16b, ESI<sup>†</sup>) under the same conditions.

The spectroscopic properties were further examined for **1**. In dilute chloroform (c = 0.10 mM), only 8.9% of **1** existed in the complexed form, denoting the dominance of the monomeric state. The visible light absorbance ranged between 378 and 510 nm ( $\varepsilon = 9.07 \times 10^3$  M<sup>-1</sup> cm<sup>-1</sup> at 395 nm, Fig. 1b), while the emission signal was centred at 560 nm (Fig. S17, ESI†). With reference to previous reports,<sup>9</sup> these signals were assigned to the admixture of metal-to-ligand and ligand-to-ligand charge-transfer (MLCT/LLCT) transitions of the alkynyl Pt( $\pi$ )(N^N^N) moiety. Upon switching the solvent from chloroform to

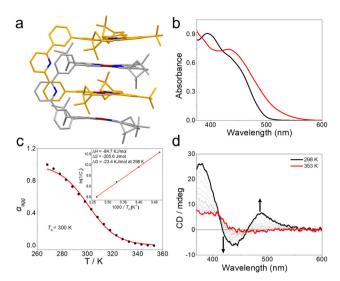


Fig. 1 (a) Optimized structure of the self-complexed dimer  $\mathbf{1}_2$ . (b) Absorption spectra of  $\mathbf{1}$  in acetonitrile (red line) and chloroform (black line) (c = 0.10 mM). (c) Degree of  $\alpha_{agg}$  for  $\mathbf{1}$  monitored at 500 nm *versus* different temperatures (c = 0.10 mM in acetonitrile). Inset: van't Hoff plot fitting for the self-complexation process of  $\mathbf{1}$ . (d) Temperature-dependent CD spectra of  $\mathbf{1}$  (c = 0.10 mM in acetonitrile). The arrows indicate the spectral change upon decreasing the temperature.

acetonitrile, the MLCT/LLCT emission signal declined for the intensity (Fig. S17, ESI<sup>†</sup>). Moreover, a low-energy shoulder band emerged for **1** ( $\varepsilon = 1.70 \times 10^3 \text{ M}^{-1} \text{ cm}^{-1}$  at 510 nm, Fig. 1b). These phenomena suggest a stronger self-complexation capability in acetonitrile. This could be ascribed to the association of the dimer **1**<sub>2</sub> into the oligomeric species, considering that  $\pi$ - $\pi$  stacking interactions are stronger in acetonitrile than those in chloroform.<sup>13</sup> The conclusion is manifested by the broadened <sup>1</sup>H NMR peaks (Fig. S18, ESI<sup>†</sup>), together with the larger hydrodynamic diameter from DLS measurements (Fig. S19, ESI<sup>†</sup>).

The intensity of the low-energy band between 510 and 600 nm declined upon increasing the temperature to 353 K, with an isosbestic point at 463 nm (Fig. S20, ESI<sup>+</sup>). The results support the reversible conversion between the monomeric state at high temperature and the complexed state at low temperature. The equal K model<sup>14</sup> was employed to fit the melting curves, acquired by plotting the absorption intensity changes at 500 nm *versus* the temperature (Fig. 1c). The  $T_{\rm m}$  values [the temperature at which the degree of aggregation  $(\alpha_{agg})$  is 0.5] increased at higher monomer concentrations (Tm: 288 K at  $2.50 \times 10^{-5}$  M versus 306 K at 2.00  $\times 10^{-4}$  M, Fig. S20, ESI<sup>†</sup>). According to a modified van't Hoff plot (Fig. 1c, inset), the enthalpy ( $\Delta H$ ) and entropy ( $\Delta S$ ) values were determined to be -84.7 kJ mol<sup>-1</sup>, and -206 J mol<sup>-1</sup> K<sup>-1</sup>, respectively. Accordingly, this provided the self-complexed binding constant of  $1.29 \times 10^4 \text{ M}^{-1}$  at 298 K, which is much higher than that in chloroform [5.34  $\times$  10<sup>2</sup> M<sup>-1</sup> (±34%)]. We rationalized that the higher self-complexion affinity in acetonitrile involved not only the dimeric stacking but the hierarchical association into oligomeric species.

Since the self-complexed structure adopts a head-to-head binding mode, it provides asymmetry by transferring chirality from the (1R)-pinenes to the Pt(II)(N^N^N) pincers. As can be seen, a weak Cotton effect below 419 nm exists for 1 at 353 K (Fig. 1d), supporting the origin of the molecular chirality from the (1R)-pinene units (Fig. S21, ESI<sup> $\dagger$ </sup>). Upon decreasing the temperature to 298 K, a bisignate CD signal appeared for 1 in the low-energy MLCT/LLCT absorption region (418-550 nm), with the positive maximum at 489 nm ( $\Delta \varepsilon = 2.04 \text{ mol}^{-1} \text{ cm}^{-1}$ ) and the negative maximum at 437 nm ( $\Delta \varepsilon = -1.89 \text{ mol}^{-1} \text{ cm}^{-1}$ , Fig. 1d). Accordingly, the self-complexation of 1 prevents carbon-carbon and carbon-platinum bond rotations, exerting a crucial impact on the supramolecular chirogenic behavior. The conclusion is further demonstrated by the weakened Cotton effect in chloroform due to its weakened self-complexation tendency (Fig. S24, ESI<sup>†</sup>).

After elucidating the self-complexation properties of 1, we turned to its host-guest complexation behavior. According to electrospray ionization mass spectrometry, an m/z value of 2501.85 was observed for  $1 \supset 2$ , corresponding to  $[1 + 2 + H]^+$ . The color of 1 in acetonitrile solution changed from yellow to orange upon adding the charge-neutral guest 2 in an equivalent ratio (Fig. S26a, ESI<sup>†</sup>). A new absorption band emerged in the low energy region, ranging from 500 to 650 nm (Fig. 2a). This is a characteristic of metal-metal-to-ligand charge-transfer (MMLCT) transitions.<sup>10</sup> Simultaneously, the MLCT/LLCT emission at 578 nm declined in its intensity, with a concomitant increase in the MMLCT emission band at 786 nm (Fig. S26b, ESI<sup>†</sup>). Depending on the molar ratio plot (Fig. S27, ESI<sup>†</sup>), the binding stoichiometry between the metallotweezer receptors 1 and guest 2 was 1:1. By fitting the collected UV/Vis absorbances at 510 nm, 525 nm and 540 nm, the "apparent" constant ( $K_d$ ) value was determined to be 3.75  $\times$  10<sup>5</sup> M<sup>-1</sup>  $(\pm 27\%)$  in acetonitrile at 298 K (Fig. 2a, inset). Since the selfcomplexation of 1 was involved in the titration process, the "real" binding constant ( $K_a$ ) value for  $1 \supset 2$  was  $6.96 \times 10^4 \text{ M}^{-1}$  $(\pm 13\%)$  (eqn (S11), ESI<sup>†</sup>).<sup>15</sup> This value was higher than that in chloroform [ $K_a$ : 1.85 × 10<sup>3</sup> M<sup>-1</sup> (±3.2%), Fig. S29, ESI<sup>†</sup>]. Although the  $K_a$  value of  $1 \supset 2$  was smaller when observed *via* <sup>1</sup>H NMR titration experiments  $[9.32 \times 10^3 \text{ M}^{-1} (\pm 28\%),$ 

Fig. 2 (a) UV/Vis absorbance changes of 1 at 298 K (c = 0.05 mM in CH<sub>3</sub>CN) upon the progressive addition of 2. Inset: intensity changes in UV/ Vis absorbance at 510 nm, 525 nm, and 540 nm. The solid lines were obtained via a Matlab-based global analysis program. (b) Optimized structure of the host-guest complex  $1 \supset 2$  on the basis of DFT calculations.

b

Fig. S30 and S31, ESI<sup>†</sup>], the value was also higher than that in chloroform [ $K_a$ : 2.54 × 10<sup>3</sup> M<sup>-1</sup> (±24%), Fig. S32 and S33, ESI<sup>†</sup>].

The energy-minimized structure of complex  $1 \supset 2$  was elucidated via DFT calculations. As expected, 2 is encapsulated into the cavity of the metallotweezers 1 to form a sandwiched complex (Fig. 2b). The inter-planar  $\pi$ -distances between 2 and the two Pt(II)(N^N^N) pincers on 1 are determined to be 3.30 Å and 3.21 Å, validating the presence of two-fold  $\pi$ - $\pi$  stacking interactions. This conclusion was further validated via <sup>1</sup>H NMR experiments. Upon addition of one equivalent of 2 to 1, the <sup>1</sup>H NMR resonances of protons H<sub>1</sub> and H<sub>2</sub> shifted upfield  $(\Delta \delta = -0.36 \text{ and } -0.58 \text{ ppm}, \text{ respectively})$ , while protons H<sub>4</sub> varied from 8.43 ppm to 8.71 ppm because of the deshielding effect (Fig. S32, ESI<sup>†</sup>). Meanwhile, the Pt-Pt distances between 1 and 2 are 3.51 and 3.22 Å, respectively. This supports the existence of Pt(II)-Pt(II) interactions in complex  $1 \supset 2$ , and is highly consistent with the emergence of the MMLCT absorption and emission signals (Fig. S26a and S26b, ESI<sup>+</sup>).

The participation of two-fold  $Pt(\pi)-Pt(\pi)$  and  $\pi-\pi$  stacking interactions contributes to the high binding affinity for complex  $1 \supset 2$ . When the control compound 3 (Fig. S21, ESI, † inset) with the mono-nuclear Pt(II)(N^N^N) unit was employed as the host instead of 1, the  $K_a$  value for the resulting complex  $3 \supset 2$ decreased to be 15.9  $M^{-1}$   $M^{-1}$  (Fig. S34 and S35, ESI<sup>+</sup>), two orders of magnitude lower than that of complex  $1 \supset 2$ . When the temperature was elevated to 353 K, the  $K_a$  value of complex  $1 \supset 2$ in acetonitrile was determined to be  $3.21 \times 10^4 \text{ M}^{-1} (\pm 5.9\%)$ (Fig. S38, ESI<sup>†</sup>), reaching one half of the value at 298 K. The high binding affinity of  $1 \supset 2$  at elevated temperature is ascribed to the weakening of the 1 self-complexation strength upon heating. This buffers the decreased host-guest complexation, and thereby the strong complexation between 1 and 2 persists.

We further investigated the supramolecular chirogenic signal for the resulting host-guest complex. Upon the gradual addition of 2 into an acetonitrile solution of 1, the positive CD signal located at 497 nm became negative ( $\Delta \varepsilon$ : from 1.79 mol<sup>-1</sup> cm<sup>-1</sup> to  $-3.18 \text{ mol}^{-1} \text{ cm}^{-1}$ , Fig. 3a and Fig. S39, ESI<sup>+</sup>). Meanwhile, the Cotton effect appeared in the MMLCT absorption region  $(\Delta \varepsilon = +1.18 \text{ mol}^{-1} \text{ cm}^{-1} \text{ at } 570 \text{ nm})$ . In stark contrast, a negligible Cotton effect was observed when employing 3 instead of 1 (Fig. S40a, ESI<sup>+</sup>), because of the weak complexation strength of complex  $3 \supset 2$ . Accordingly, metallotweezers/guest complexation

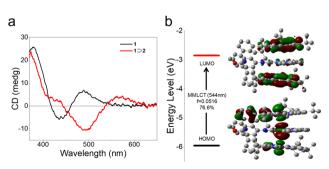


Fig. 3 (a) CD spectra of 1 and  $1 \supset 2$  at 298 K (c = 0.10 mM for each compound in CH<sub>3</sub>CN). (b) Energy-level diagram of  $1 \supset 2$  via TD-DFT computations.

а 0.8

0.6 ŝ Absorband

0.2

0.0

Equivalence of 2

600

500 Wavelength / nm

with sufficient binding affinity is a prerequisite for supramolecular chirogenesis. The Cotton shape of  $1 \supset 2$  was maintained at elevated temperatures because of the robust host-guest complexation, despite the decreased CD intensities (at 497 nm:  $\Delta \varepsilon = |1.26| \text{ cm}^{-1} \text{ M}^{-1}$  at 353 K *versus* |3.18| cm<sup>-1</sup> M<sup>-1</sup> at 298 K, Fig. S41, ESI<sup>†</sup>).

The origin of the low-energy supramolecular chiroptical signals (ranging from 444 nm to 624 nm) was clarified using time-dependent density functional theory (TD-DFT) calculations. As shown in Fig. 3b, the electron density of the LUMO is distributed over the  $Pt(\pi)(N^N^N)$  pincers of 1. Meanwhile, the electron density of the HOMO is mainly distributed on the  $5d_{z^2}$  orbitals of the Pt(II) atoms in both 1 and 2. Accordingly, the theoretical Cotton effect in the low-energy absorption region is composed of HOMO  $\rightarrow$  LUMO transitions (composition: 76.6%) at 544 nm), belonging to metal-metal-to-ligand charge transfer (MMLCT) transitions. For most of previous host-guest systems, supramolecular chirogenic signals arose from the individual or conjoint non-covalent forces of metal-ligand coordination, hydrogen bonding,  $\pi$ - $\pi$  stacking, and hydrophobic interactions. Complex  $1 \supset 2$  represents a rare type of supramolecular chirogenic system with the involvement of Pt(II)-Pt(II) metalmetal interactions.<sup>16</sup>

In summary, metallotweezer **1** with optically active  $Pt(\pi)(N^{N}N)$  pincers prefers to form a self-complexed structure *via* a head-to-head binding mode, leading to supramolecular chirogenic signals in the MLCT/LLCT absorption region. Furthermore, a metallotweezers/guest complex forms upon adding the complementary guest **2** into **1**. This is accompanied by the formation of  $Pt(\pi)$ – $Pt(\pi)$  metal-metal interactions, and thereby induces supramolecular chirogenesis in the MMLCT transition region. Therefore, a dual supramolecular chirogenesis system in the visible region has been successfully constructed by taking advantage of the diverse complexation modes of  $Pt(\pi)$  metallotweezers.

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### Conflicts of interest

There are no conflicts to declare.

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