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In situ Raman study of the photoinduced behavior of dye molecules on TiO₂(*hkl*) single crystal surfaces†

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In dye-sensitized solar cells (DSSCs), the TiO₂/dye interface significantly affects photovoltaic performance. However, the adsorption and photoinduced behavior of dye molecules on the TiO₂ substrate remains unclear. Herein, shell-isolated nanoparticle-enhanced Raman spectroscopy (SHINERS) was used to study the adsorption and photoinduced behavior of dye (N719) molecules on different TiO₂(*hkl*) surfaces. On TiO₂(001) and TiO₂(110) surfaces, the *in situ* SHINERS and mass spectrometry results indicate S=C bond cleavage in the anchoring groups of adsorbed N719, whereas negligible bond cleavage occurs on the TiO₂(111) surface. Furthermore, DFT calculations show the stability of the S=C anchoring group on three TiO₂(*hkl*) surfaces in the order TiO₂(001) < TiO₂(110) < TiO₂(111), which correlated well with the observed photocatalytic activities. This work reveals the photoactivity of different TiO₂(*hkl*) surface structures and can help with the rational design of DSSCs. Thus, this strategy can be applied to real-time probing of photoinduced processes on semiconductor single crystal surfaces.

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Dye-sensitized solar cells (DSSCs) have attracted much interest due to their versatility, low cost, and ease of fabrication into flexible devices.¹ However, to date, the highest certified power conversion efficiency (PCE) recorded for a DSSC under simulated sunlight illumination is only 11.9%,² while the theoretical maximum PCE is over 30%.³ Additionally, the long-term stability of DSSCs is very important for practical applications. DSSCs function *via* a charge injection mechanism between closely interacting dye/semiconductor molecules.⁴ Studying the adsorption and photoinduced behavior of N719 molecules on TiO₂ surfaces can reveal the fundamental reaction processes and help improve the conversion efficiency and lifetime of DSSCs.⁵

Previously, there have been some research studies on the adsorption configuration and electronic properties of N719 dye molecules at TiO₂ surfaces. Grätzel *et al.* concluded, based on Fourier transform infrared spectroscopy (FTIR), that N719 was anchored on the TiO₂ surface in a bridged configuration *via* two

carboxyl groups.^{4a} They further showed with density functional theory (DFT) calculations that deprotonation of the terminal carboxylic groups cause the excitation energy to increase.⁶ In addition, Tian *et al.* investigated the Ag@TiO₂/N719 interface under a series of potential controls using *in situ* Raman spectroscopy and found that the SCN ligand participates in N719 adsorption on the surface of aggregated TiO₂ nanoparticles,⁷ while Lund *et al.* employed high-performance liquid chromatography (HPLC) and ultraviolet-visible spectroscopy (UV/Vis) to determine the photoinduced degradation and regeneration rates for N719 at the TiO₂ surface. Depending on the rates achieved by the DSSC, they calculated operating times as high as 20 years or as low as 0.5 years.⁵ However, most of these studies only focused on polycrystalline TiO₂ substrate surfaces where it is difficult to study the structure–activity relationship of the substrate. For a better understanding of the reaction processes, investigating atomically flat single crystals with well-defined surface structures as well as electronic and optical properties can act as a bridge between experimental studies and theoretical simulations.

Surface-enhanced Raman spectroscopy (SERS) is a non-destructive and ultrasensitive technique that has become a powerful surface analysis tool.⁸ Unfortunately, conventional SERS techniques cannot be applied to single crystal surfaces because they lack sufficient surface plasmon resonance (SPR).⁹ In 2010, our group invented the shell-isolated nanoparticle-enhanced Raman spectroscopy (SHINERS)¹⁰ technique for

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obtaining enhanced surface Raman signals, irrespective of the surface and with single crystal surface compatibility. Shell-isolated nanoparticles (SHINs) act as surface Raman signal enhancers without detrimentally affecting the external environment or reaction processes due to their inert silica shell coating.¹¹ Since then, important catalytic reactions, such as the oxygen reduction reaction (ORR) and oxygen evolution reaction (OER) with their intermediate species, as well as the structure of interfacial water at single-crystal electrode surfaces have been observed using *in situ* SHINERS.¹²

Herein, the adsorption and photoinduced behavior of N719 molecules on three TiO₂(*hkl*) surfaces (rutile TiO₂(001), TiO₂(110), and TiO₂(111)) under 405 nm laser irradiation were investigated with *in situ* SHINERS. Combined with DFT simulations, the TiO₂(*hkl*) facet effect and the intrinsic photocatalytic mechanism were elucidated and provide a promising method for studying photoinduced surface processes in real-time.

In this system, N719 dye molecules were assembled on rutile TiO₂(*hkl*) surfaces and SHINs were deposited on top and then transferred to a Raman cell filled with acetonitrile (without N719 in solution). During *in situ* Raman measurements, the cell was illuminated with two different wavelength lasers, 638 nm and 405 nm (6600 mW cm⁻²), for detecting Raman signals and exciting the photoinduced reactions, respectively (Fig. 1a). Besides, the power density of the 405 nm laser is 66 times higher than one sun condition (100 mW cm⁻²), so the photo-induced reaction speed of N719 on the TiO₂ surface will be much higher due to more photo-generated carriers being produced under 405 nm laser illumination. Therefore, it is more likely for us to obtain the Raman signal in the current research. Fig. 1b shows the TEM image of two adjacent SHINs. The enhancement factor of SHINs on the TiO₂ surfaces was assessed using three-dimensional finite-difference time-domain (3D-FDTD) simulations (Fig. 1c). Plasmonic coupling dictates the Raman enhancement factor and is shown to be stronger between adjacent SHINs (55 nm dia. Au core @ 2 nm

thick SiO₂ shell) than between SHINs and the rutile TiO₂(*hkl*) substrate; however N719 molecules were only adsorbed on the TiO₂ single crystal surface. Therefore, all N719 Raman signals measured herein must originate from the electromagnetic field enhancement between the SHINs and the substrate which had a calculated enhancement factor of about 10³ times (Fig. 1c). Importantly, compared to a strategy with no SHINs (Fig. S1†), this SHINERS strategy provides enough enhancement for detection of surface N719.

In Fig. 2a, three Raman bands are visible on all pristine rutile TiO₂(*hkl*) surfaces at around 238 cm⁻¹, 448 cm⁻¹ and 610 cm⁻¹ assigned to the multi-phonon, E_g, and A_{1g} modes of the Ti–O–Ti bond vibrations, respectively,¹³ and their relative peak intensities ($I_{238\text{ cm}^{-1}}$, $I_{448\text{ cm}^{-1}}$ and $I_{610\text{ cm}^{-1}}$) are clearly surface facet dependent. For the N719 molecule (see Fig. S2† for the chemical structure), there are two SCN binding groups attached to the Ru atom, both of which share the highest occupied molecular orbital (HOMO) and the 2,2-bipyridyl-4,4-dicarboxylic ligand contains the lowest unoccupied molecular orbital (LUMO) of the molecule. In the UV-vis absorption spectrum of N719 dissolved in ethanol (Fig. S3†), N719 has two clear adsorption peaks. The first peak, located at around 389 nm, is associated with a mixture of ligand-to-ligand charge transfer (LLCT) and metal-to-ligand charge transfer (MLCT) processes, while the second peak at 532 nm is related to low energy MLCT processes within N719 molecules.¹⁴ The $I_{610\text{ cm}^{-1}}/I_{448\text{ cm}^{-1}}$ is the same after adding SHINs on the TiO₂(001) surface (Fig. S4†), but interestingly, after N719 is adsorbed at the surface, the $I_{610\text{ cm}^{-1}}/I_{448\text{ cm}^{-1}}$ on all three rutile TiO₂(*hkl*) surfaces changed. In Fig. 2b, $I_{610\text{ cm}^{-1}}/I_{448\text{ cm}^{-1}}$ intensities decreased on TiO₂(111) and TiO₂(001) surfaces and increased on the TiO₂(110) surface. For N719 molecules chemisorbed on rutile TiO₂(*hkl*) surfaces, the different N719 anchoring modes result in different interactions of Ti–O–Ti (Fig. 2b),¹⁵ so the ratios of $I_{610\text{ cm}^{-1}}/I_{448\text{ cm}^{-1}}$ differ between the three TiO₂(*hkl*) surfaces after N719 adsorption, manifesting as a change in Raman peak intensities.^{14a}

For N719 adsorbed at different rutile TiO₂(*hkl*) surfaces, both 1482 cm⁻¹ and 1610 cm⁻¹ peaks, belonging to the C=N and C=C stretching vibrations of the bipyridyl ring, respectively,¹⁶ are surface facet dependent (Fig. 2b). For TiO₂(110), the 1482 cm⁻¹ peak intensity ($I_{1482\text{ cm}^{-1}}$) is obviously stronger than the 1610 cm⁻¹ peak intensity ($I_{1610\text{ cm}^{-1}}$), whereas on TiO₂(111), the $I_{1482\text{ cm}^{-1}}$ is weaker than $I_{1610\text{ cm}^{-1}}$, and both peak intensities are similar on TiO₂(001). This can be attributed to the different



Fig. 1 (a) Schematic diagram of using a SHINERS nanoantenna to study the photoinduced reaction of N719 illuminated with two lasers on TiO₂. The inset shows electron transfer from the 405 nm laser excited N719 molecules to the TiO₂ single crystal surface and the 638 nm laser used to detect the changes in excited the molecules. (b) TEM image of 55 nm Au@2 nm SiO₂ SHINs. (c) 3D-FDTD simulation of the enhancement hotspot of a SHIN dimer with a nanogap of 2 nm on rutile TiO₂.



Fig. 2 (a) SHINER spectra of rutile TiO₂(001), TiO₂(110), and TiO₂(111) surfaces in acetonitrile without and (b) with N719 molecules adsorbed.



rutile single crystal surfaces having different N719 adsorption configurations, giving rise to different bond enhancements.^{11a,17}

SHINERS was then employed to study, *in situ*, the adsorption and photoinduced behavior of N719 molecules on the TiO₂(001) surface under 405 nm laser illumination. Fig. S5† shows Raman spectra of TiO₂(111) with and without 405 nm laser illumination and shows no observable interference peaks. This lack of Raman shifts under 405 nm illumination can be further verified using the formula in Fig. S6.† Furthermore, in Fig. S7,† the Raman signals for rutile TiO₂(110) without N719 adsorbed in acetonitrile were observed under extended illumination times (50 min). And peaks associated with the rutile TiO₂(110) substrate are not changed, indicating that 405 nm laser illumination does not change the TiO₂(*hkl*) substrate spectral peaks. As shown in Fig. 3a and b, as the illumination time increases, the Raman peaks of the adsorbed N719 molecules gradually weaken, while the rutile TiO₂ peaks increase significantly. Considering these spectral trends over time (Fig. 3b), we tentatively ascribe this trend to desorption of N719 molecules from the TiO₂(001) surface as opposed to conformational changes in the N719 adsorption angle. Additionally, the Raman peak of the SCN group at around 2136 cm⁻¹ blue shifts to 2146 cm⁻¹ with increasing illumination time, indicating that as the N719 surface coverage decreases the N719 molecules adsorb on the TiO₂(001) surface *via* the S-terminal of the N=C=S group, corresponding with previous studies.^{7,18}

Mass spectrometry was used to examine photoinduced reaction products. In Fig. 3c, the produced ion of *m/z* 1157 appears after 36 hours under 405 nm laser illumination which corresponds to the loss of an S atom from N719. This confirms bond cleavage of the S=C bond in the SCN group adsorbed on the TiO₂(001) surface under 405 nm laser illumination producing

a [Ru^{II}(H₂dcbpy)₂(NCS)(CN)] species.¹⁹ For comparison, the adsorption behavior of the N=C=S functional group of a Ph-N=C=S species on different rutile TiO₂(*hkl*) surfaces was simulated using DFT. This also allowed for the reaction energy change due to the S=C double bond cleavage to be directly calculated. The calculated structures of the Ph-N=C=S group and S atom adsorbed on rutile TiO₂(001) are shown in Fig. 3d. The S=C dissociation energy on TiO₂(001) was found to be lower than that on TiO₂(111) and TiO₂(110) surfaces (Fig. S8†). Thus, the TiO₂(001) peak intensities (238 cm⁻¹, 448 cm⁻¹, and 610 cm⁻¹) increase as the illumination time increases due to increasing S=C bond cleavage resulting in N719 desorption.²⁰

To further determine the facet effect, the adsorption and photoinduced behavior of N719 was also systemically studied on TiO₂(111) and TiO₂(110) with *in situ* SHINERS. In Fig. 4a, the time-dependent Raman spectra of N719 at the TiO₂(110) surface are similar to those at the TiO₂(001) surface (Fig. 3a), with N719 Raman intensities decreasing with increasing illumination times; however the trend is less pronounced. Meanwhile, the Raman peak shift of the SCN group around 2143 cm⁻¹ is also small (~6 cm⁻¹). In contrast, in Fig. 4b, the N719 Raman peak intensities gradually increase with time; meanwhile the rutile TiO₂(111) peaks around 236 cm⁻¹, 446 cm⁻¹ and 610 cm⁻¹ significantly weaken. The resulting decrease in the HOMO-LUMO gap in the adsorption complex produces a species that is more strongly Raman resonant with the 638 nm laser, resulting in the observed higher peak intensities.²¹ In addition, The SCN peak is more stable on the TiO₂(111) surface during the photoinduced reaction processes and does not shift with time. Furthermore, comparing the mass spectra between surfaces, less of the [Ru^{II}(H₂dcbpy)₂(NCS)(CN)] photoinduced degradation product is produced on TiO₂(111) than on the TiO₂(001)



Fig. 3 (a) SHINER Spectra of N719 adsorbed at TiO₂(001) in acetonitrile under 405 nm laser illumination, spectra collected every 4 minutes. (b) Plot of changes in significant N719 (blue lines and points) and TiO₂ (black line and points) spectral peak intensities with time. (c) Mass spectra of a 5 × 10⁻⁴ M N719 ethanol solution containing TiO₂(001) under 405 nm laser illumination after 0 h and 36 h, and without TiO₂(001) after 36 hours. (d) Schematic of the optimized structures of the Ph-N=C=S group and S atom adsorbed on rutile TiO₂(001).



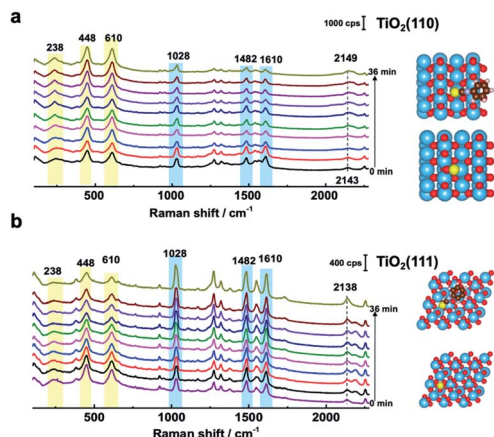


Fig. 4 (a) SHINER Spectra of N719 adsorbed on a rutile $\text{TiO}_2(110)$ surface with 405 nm laser illumination. The right shows the DFT simulated optimized structures of the Ph-N=C=S group (top) and S atom (below) adsorbed on rutile $\text{TiO}_2(110)$. (b) SHINER Spectra of N719 molecules adsorbed on a rutile $\text{TiO}_2(111)$ surface with 405 nm laser illumination. Spectra were collected every 4 minutes. The right shows the DFT simulated optimized structures of the Ph-N=C=S group (top) and S atom (below) on rutile $\text{TiO}_2(111)$.

and $\text{TiO}_2(110)$ surfaces (Fig. S9[†]), indicating that N719 adsorption is more stable on $\text{TiO}_2(111)$. Interestingly, the SCN complex has different Raman peak shifts in the order $\text{TiO}_2(111)$ (0 cm^{-1}) < $\text{TiO}_2(110)$ (6 cm^{-1}) < $\text{TiO}_2(001)$ (10 cm^{-1}), which correlates well with the photoinduced stability of the SCN group on three rutile $\text{TiO}_2(hkl)$ surfaces, and the DFT simulated thermodynamic stability of the S=C functional group. Thus, the important effect of TiO_2 structural on N719 molecules adsorbed on three different rutile $\text{TiO}_2(hkl)$ surfaces was elucidated.

In summary, using *in situ* SHINERS and DFT simulations, SCN was found to be the main adsorption group of the N719 dye molecule at different rutile $\text{TiO}_2(hkl)$ surfaces and to play a determining role in its stability under 405 nm laser illumination. Importantly, different rutile $\text{TiO}_2(hkl)$ surfaces were also shown to have obvious facet effects on N719 molecular adsorption and photoinduced behavior. Under illumination, the photoinduced process caused the N719 molecules to desorb from $\text{TiO}_2(110)$ and $\text{TiO}_2(001)$ surfaces. Desorption was shown to occur *via* bond cleavage of the S=C bond using mass spectrometry. In addition, DFT simulations showed the $\text{TiO}_2(001)$ surface to have the highest reaction activity for catalyzing S=C double bond cleavage, while the $\text{TiO}_2(111)$ surface had the lowest reaction activity and was the most stable single crystal surface for adsorbed N719 molecules. These results indicate that N719 molecules will have long-term stability adsorbed on $\text{TiO}_2(111)$, providing important guidance for the design of practical DSSCs. Our research has elucidated the mechanism of DSSC interfacial reaction and demonstrates the universality of this approach to study the photoinduced reaction on semiconductor single crystal surfaces.

Conflicts of interest

The authors declare no conflict of interest.

Acknowledgements

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Notes and references

- (a) B. O'Regan and M. Grätzel, *Nature*, 1991, **353**, 737; (b) M. L. Parisi, S. Maranghi and R. Basosi, *Renewable Sustainable Energy Rev.*, 2014, **39**, 124; (c) A. Fakharuddin, R. Jose, T. M. Brown, F. Fabregat-Santiago and J. Bisquert, *Energy Environ. Sci.*, 2014, **7**, 3952; (d) A. Hagfeldt, G. Boschloo, L. Sun, L. Kloo and H. Pettersson, *Chem. Rev.*, 2010, **110**, 6595.
- M. A. Green, Y. Hishikawa, E. D. Dunlop, D. H. Levi, J. Hohl-Ebinger and A. W. Y. Ho-Baillie, *Prog. Photovoltaics Res. Appl.*, 2018, **26**, 659.
- M. Shakeel Ahmad, A. K. Pandey and N. Abd Rahim, *Renewable Sustainable Energy Rev.*, 2017, **77**, 89.
- (a) M. K. Nazeeruddin, R. Humphry-Baker, P. Liska and M. Grätzel, *J. Phys. Chem. B*, 2003, **107**, 8981; (b) N. Takeda and B. A. Parkinson, *J. Am. Chem. Soc.*, 2003, **125**, 5559; (c) Q. Q. Zuo, Y. L. Feng, S. Chen, Z. Qiu, L. Q. Xie, Z. Y. Xiao, Z. L. Yang, B. W. Mao and Z. Q. Tian, *J. Phys. Chem. C*, 2015, **119**, 18396; (d) L. Q. Xie, D. Ding, M. Zhang, S. Chen, Z. Qiu, J. W. Yan, Z. L. Yang, M. S. Chen, B. W. Mao and Z. Q. Tian, *J. Phys. Chem. C*, 2016, **120**, 22500.
- F. Nour-Mohammadi, S. D. Nguyen, G. Boschloo, A. Hagfeldt and T. Lund, *J. Phys. Chem. B*, 2005, **109**, 22413.
- M. K. Nazeeruddin, F. De Angelis, S. Fantacci, A. Selloni, G. Viscardi, P. Liska, S. Ito, B. Takeru and M. Grätzel, *J. Am. Chem. Soc.*, 2005, **127**, 16835.
- Z. Qiu, M. Zhang, D. Y. Wu, S. Y. Ding, Q. Q. Zuo, Y. F. Huang, W. Shen, X. D. Lin, Z. Q. Tian and B. W. Mao, *Chemphyschem*, 2014, **14**, 2217.
- S. M. Nie and S. R. Emery, *Science*, 1997, **275**, 1102.
- K. A. Willets and R. P. Van Duyne, *Annu. Rev. Phys. Chem.*, 2007, **58**, 267.
- J. F. Li, Y. F. Huang, Y. Ding, Z. L. Yang, S. B. Li, X. S. Zhou, F. R. Fan, W. Zhang, Z. Y. Zhou, D. Y. Wu, B. Ren, Z. L. Wang and Z. Q. Tian, *Nature*, 2010, **464**, 392.
- (a) J. F. Li, Y. J. Zhang, A. V. Rudnev, J. R. Anema, S. B. Li, W. J. Hong, P. Rajapandiyam, J. Lipkowski, T. Wandlowski and Z. Q. Tian, *J. Am. Chem. Soc.*, 2015, **137**, 2400; (b) J. F. Li, Y. J. Zhang, S. Y. Ding, R. Panneerselvam and Z. Q. Tian, *Chem. Rev.*, 2017, **117**, 5002; (c) Y. H. Wang, M. M. Liang, Y. J. Zhang, S. Chen, P. Radjenovic, H. Zhang, Z. L. Yang, X. S. Zhou, Z. Q. Tian and J. F. Li, *Angew. Chem., Int. Ed.*, 2018, **57**, 11257.
- (a) J. C. Dong, X. G. Zhang, V. Briega-Martos, X. Jin, J. Yang, S. Chen, Z. L. Yang, D. Y. Wu, J. M. Feliu, C. T. Williams, Z. Q. Tian and J. F. Li, *Nat. Energy*, 2019, **4**, 60; (b) J. C. Dong, M. Su, V. Briega-Martos, L. Li, J. B. Le,



- P. Radjenovic, X. S. Zhou, J. M. Feliu, Z. Q. Tian and J. F. Li, *J. Am. Chem. Soc.*, 2020, **142**, 715; (c) N. Bodappa, M. Su, Y. Zhao, J. B. Le, W. M. Yang, P. Radjenovic, J. C. Dong, J. Cheng, Z. Q. Tian and J. F. Li, *J. Am. Chem. Soc.*, 2019, **141**, 12192; (d) C. Y. Li, J. B. Le, Y. H. Wang, S. Chen, Z. L. Yang, J. F. Li, J. Cheng and Z. Q. Tian, *Nat. Mater.*, 2019, **18**, 697.
- 13 S. K. Parayil, R. J. Psota and R. T. Koodali, *Int. J. Hydrogen Energy*, 2013, **38**, 10215.
- 14 (a) K. E. Lee, M. A. Gomez, T. Regier, Y. Hu and G. P. Demopoulos, *J. Phys. Chem. C*, 2011, **115**, 5692; (b) H. Rensmo, S. Lunell and H. Siegbahn, *J. Photochem. Photobiol., A*, 1998, **114**, 117; (c) A. El-Shafei, M. Hussain, A. Atiq, A. Islam and L. Han, *J. Mater. Chem.*, 2012, **22**, 24048.
- 15 J. G. Ma, C. R. Zhang, J. J. Gong, B. Yang, H. M. Zhang, W. Wang, Y. Z. Wu, Y. H. Chen and H. S. Chen, *J. Chem. Phys.*, 2014, **141**, 234705.
- 16 K. E. Lee, M. A. Gomez, S. Elouatik and G. P. Demopoulos, *Langmuir*, 2010, **26**, 9575.
- 17 (a) A. G. Thomas and K. L. Syres, *Chem. Soc. Rev.*, 2012, **41**, 4207; (b) R. B. Jaculbia, H. Imada, K. Miwa, T. Iwasa, M. Takenaka, B. Yang, E. Kazuma, N. Hayazawa, T. Taketsugu and Y. Kim, *Nat. Nanotechnol.*, 2020, **15**, 105.
- 18 J. Singh, A. Gusain, V. Saxena, A. K. Chauhan, P. Veerender, S. P. Koiry, P. Jha, A. Jain, D. K. Aswal and S. K. Gupta, *J. Phys. Chem. C*, 2013, **117**, 21096.
- 19 P. T. Nguyen, P. E. Hansen and T. Lund, *Sol. Energy*, 2013, **88**, 23.
- 20 C. Pérez León, L. Kador, B. Peng and M. Thelakktat, *J. Phys. Chem. B*, 2006, **110**, 8723.
- 21 S. Ben-Jaber, W. J. Peveler, R. Quesada-Cabrera, E. Cortes, C. Sotelo-Vazquez, N. Abdul-Karim, S. A. Maier and I. P. Parkin, *Nat. Commun.*, 2016, **7**, 12189.

