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Cite this: RSC Adv., 2019, 9, 1976

Solid-state AIEnh-circularly polarised luminescence of chiral perylene diimide fluorophores†

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Solid-state organic fluorescent materials are important for the development of electroluminescent sensing devices. Herein, we report that *N*,*N'*-bis((*R*)-1-phenylethyl)perylene-3,4,9,10-tetracarboxylic diimide [(*R*,*R*)-**BPP**] and its antipode [(*S*,*S*)-**BPP**], which contain extended π -electrons through planar perylenes, emit solid-state aggregation-induced-enhanced (AIEnh) circularly polarised luminescence (CPL) in inorganic (KBr) pellets and organic-polymer-film (PMMA- and *myo*-IPU-film) states; this CPL is difficult to observe in solution. These chiral perylene fluorophores emit AIEnh-CPL with high dissymmetry factors (g_{CPL}) (up to 2.4 × 10⁻³) and high quantum yields (Φ_{F} , up to 0.43) in the three solid matrices.

Received 28th November 2018 Accepted 23rd December 2018

DOI: 10.1039/c8ra09785b

rsc.li/rsc-advances

Introduction

Because of their potential applications to electroluminescence (EL) and optical-sensing devices, solid-state organic fluorescent materials have attracted much attention.¹ In particular, optically active organic and inorganic luminophores that emit circularly polarised luminescence (CPL) with high dissymmetry factors (g_{CPL}) and high quantum yields $(\Phi_{\rm F})$ have received significant attention in the field of circularly polarised organic light-emitting diodes (CP-OLEDs).² Recently, aggregationinduced-enhanced (AIEnh) CPL has also attracted much interest. However, AIEnh-CPL is almost always emitted from highly concentrated solutions.3 Therefore, systems that emit solid-state AIEnh-CPL are highly desired. With the aim of providing a common design approach to novel chiral solid-state AIEnh organic fluorophores, in this study we prepared solidstate AIEnh circularly polarised luminophores by doping chiral N,N'-bis((R)-1-phenylethyl)perylene-3,4,9,10tetracarboxylic diimide [(R,R)-BPP] and its antipode [(S,S)-BPP] into various types of solid matrix. The two BPPs are chiral bisubstituted perylene fluorophores composed of extended π electron planar perylenes. We then examined the AIEnh-CPL emissions from these imbedded BPP systems as AIEnhfluorophore models.4 KBr was used as the solid inorganic matrix, while poly(methyl methacrylate) (PMMA) and polyurethane (*myo*-IPU) were chosen as solid organic polymer matrices.⁵ *Myo*-IPU is derived from naturally occurring *myo*inositol (a *meso*-hexose) and characteristically has a higher glass transition temperature ($T_{\rm g} \sim 150$ °C) than conventional PMMA ($T_{\rm g} \sim 105$ °C).

Herein, we report that (R,R)- and (S,S)-**BPP** emit solid-state AIEnh-CPL with high quantum efficiencies when imbedded in solid inorganic and organic polymer matrices (Fig. 1).

Experimental

General

The KBr and PMMA matrices for solid pellets and films were purchased from JASCO (Tokyo, Japan) and the FUJIFILM Wako Pure Chemical Corp. (Osaka, Japan), respectively. Polyurethane (*myo*-IPU) was prepared by a previously reported method.⁵ Chloroform (CHCl₃) used for the determination of solutionstate optical properties was purchased from the FUJIFILM Wako Pure Chemical Corp. (Osaka, Japan).

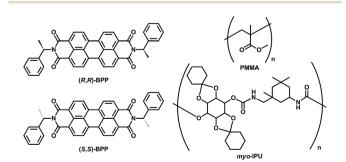


Fig. 1 The chiral perylene fluorophores ((R,R)-BPP and (S,S)-BPP) and organic polymer matrices (PMMA and myo-IPU) used in this study.

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[†] Electronic supplementary information (ESI) available. See DOI: 10.1039/c8ra09785b

Syntheses of N,N'-bis((R)-1-phenylethyl)perylene-3,4,9,10-tetracarboxylic diimide [(R,R)-BPP] and (S,S)-BPP

(R,R)-BPP and (S,S)-BPP were prepared with reference to a previously reported method.⁴ A mixture of 3,4,9,10-perylenetetracarboxylic dianhydride (430.0 mg, 1.1 mmol), (R)-1phenylethylamine (0.3 mL, 2.4 mmol), and zinc acetate dihydrate (240.1 mg, 1.1 mmol) in quinolone (20 mL) was heated with stirring at 180 °C under nitrogen for 4 h. After cooling to room temperature, the reaction mixture was treated with 2 N HCl and stirred at room temperature overnight. The resulting dark-red precipitate was collected by filtration, washed thoroughly with water and methanol (MeOH), and dried in vacuo. The crude product was purified by silica-gel column chromatography using CHCl₃/EtOAc (20: 1, v/v) as the eluent, to give 523.2 mg (79% yield) of (R,R)-BPP as a deep-red solid. The Senantiomer, namely (S,S)-BPP, was synthesised from (S)-1phenylethylamine by the same procedure to provide a deep-red solid (76% yield): ¹H NMR for (R,R)-BPP (CDCl₃ 500 MHz) δ 2.06 (d, J = 7.5 Hz, 3H), 6.60 (q, J = 7.5 Hz, 2H), 7.29 (t, J = 7.5 Hz, 2H)2H), 7.38 (t, J = 7.5 Hz, 4H), 7.57 (d, J = 7.5 Hz, 4H), 8.58 (d, J = 8.0 Hz, 4H), 8.66 (d, J = 8.0 Hz, 4H); ¹³C NMR (CDCl₃ 125 MHz) δ 16.3, 50.5, 122.9, 123.5, 126.1, 127.2, 127.5, 128.2, 129.2, 131.3, 134.3, 140.5, 163.3.

Unpolarised photoluminescence (PL) quantum yields

Samples were dispersed in KBr pellets in the following way. Microcrystalline powdered samples (10 μ g) were finely ground with 120 mg of KBr, and the mixed powder was compressed at 80 MPa cm⁻² under vacuum for 25 min to provide a 10 mmdiameter, 0.7 mm-thick transparent disk. Samples dispersed in PMMA or *myo*-IPU films were prepared using PMMA or *myo*-IPU (1.0 \times 10⁻² M) in a spin coater at 3000 rpm (Opticoat MS-A100, Mikasa, Tokyo, Japan).

Absolute unpolarised photoluminescence (PL) quantum yields ($\Phi_{\rm F}$) of the KBr-pellet-dispersed states, the PMMA- or *myo*-IPU-film-dispersed states, and CHCl₃ solutions were determined using a Hamamatsu Photonics C9920-02 absolute quantum yield spectrometer (Hamamatsu, Japan) in air at room temperature, at excitation wavelengths of 480 nm (KBr) and 450 nm (polymer film), respectively, and at 450 nm for the CHCl₃ solution. A 10 mm path length was used for solution-phase spectroscopy.

Solid-state PL and CPL spectroscopy

Solid-state PL and CPL spectra of the KBr-pelleted samples were acquired by comprehensive chiroptical spectrophotometry (CCS, JASCO, Tokyo, Japan).⁶ CPL spectra were recorded with a slit width of 33 mm and a 10 nm spectral bandwidth for the excitation and emission monochromators. The PL and CPL spectra of the PMMA or *myo*-IPU films and the CHCl₃ solutions were acquired using a JASCO CPL-300 spectrofluoropolarimeter (Tokyo, Japan) at room temperature, at a scattering angle of 0° upon excitation with unpolarised, monochromated incident light with a 10 nm bandwidth. An excitation wavelength of 480 nm was used for the KBr-pelleted samples, while 450 nm

was used for the PMMA or *myo*-IPU films, and for the $CHCl_3$ solutions. A 10 mm path length was used for solution-phase spectroscopy.

Circular dichroism (CD) and UV-Vis absorption spectroscopy

Circular dichroism (CD) and UV-Vis absorption spectra of the KBr-pelleted samples were acquired by CCS at room temperature.⁶ The CD and UV-Vis absorption spectra of PMMA- or *myo*-IPU-film-dispersed states and CHCl₃ solutions were recorded with a JASCO J-820 spectropolarimeter at room temperature. The KBr-pelleted and PMMA (or *myo*-IPU) films were prepared as described for the samples used for solid-state PL and CPL spectroscopy. A 1 mm path length was used for solution-phase spectroscopy.

Results and discussion

Prior to elucidating the solid-state chiroptical properties of (R,R)-**BPP** and (S,S)-**BPP**, we first investigated their PL, CPL, and CD spectral characteristics in CHCl₃ solution, at room temperature. To determine whether or not the obtained results are due to artefacts introduced by (R,R)-**BPP** and other factors, the CPL spectrum of (S,S)-**BPP** was also acquired. The CPL and PL spectra of (R,R)-**BPP** and (S,S)-**BPP** are shown in Fig. 2 ((R,R)-**BPP** in blue; (S,S)-**BPP** in red). When dissolved in CHCl₃ (1.0 × 10⁻⁴ M), (R,R)-**BPP** exhibited PL at 549, 576, 623 nm with an absolute photoluminescence quantum yield (Φ_F) of 0.76. Unfortunately, meaningful CPL spectra of chiral **BPP** could not be acquired under these conditions.

The CD and UV-Vis absorption spectra of (R,R)-**BPP** and (S,S)-**BPP** in CHCl₃ are shown in Fig. 3; as expected, the CD spectra are almost mirror images.

Absorption bands at 459, 490, and 527 nm that originate from π - π * transitions in the perylene units are commonly observed, irrespective of the chirality of the sample. As same as CPL properties in the solution-dissolved state, these corresponding CD bands could not be mainly observed.

For the chiral **BPPs** to emit CPL, we turned to solid-state AIEnh-CPL spectroscopy. The solid-state PL and CPL spectra

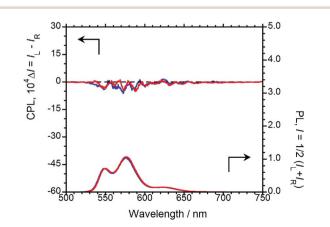


Fig. 2 CPL (upper) and PL (lower) spectra of (*R*,*R*)-**BPP** (blue) and (*S*,*S*)-**BPP** (red) in CHCl₃ (1.0 \times 10⁻⁴ M).

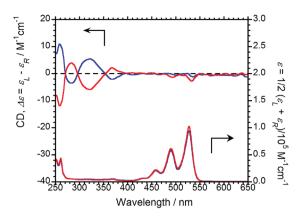


Fig. 3 CD (upper) and UV-Vis absorption (lower) spectra of (R,R)-BPP (blue) and (S,S)-BPP (red) in CHCl₃ (1.0 × 10⁻⁴ M).

of (R,R)-**BPP** were acquired in the inorganic KBr matrix, the results of which are shown in Fig. 4. In contrast to the results in CHCl₃ solution, solid-state CPL was observed from the KBr-pelleted sample; (R,R)-**BPP** in the KBr pellet exhibited a λ_{CPL} of 657 nm with a corresponding Φ_F of 0.09. When the spectrum of the KBr-pellet-dispersed state is compared with that of the CHCl₃ solution, quite different λ_{PL} values are observed (576 nm for CHCl₃ and 632 nm for KBr). The λ_{PL} of the KBr pellet suggests that CPL from this state is AIEnh CPL derived from the aggregation of (R,R)-**BPP** in the solid state. Unfortunately, the value of Φ_F in the KBr-dispersed state was smaller than that in CHCl₃ solution, which is ascribable to fluorescence quenching under high-density conditions in the solid compared to the low-concentration conditions in solution. In addition, the external heavy-atom (Br) effect may also be partially responsible.

A negative (-) CPL spectrum was observed for (R,R)-**BPP** in the solid state. To determine whether or not these results are due to artefacts introduced by (R,R)-**BPP** and the KBr matrix, the CPL spectrum of (S,S)-**BPP** was acquired (Fig. 4). (R,R)-**BPP** and (S,S)-**BPP** in KBr matrices exhibit almost mirror-image CPL spectra. The magnitude of the circular polarisation in the excited state is defined as $g_{CPL} = \Delta I/I = (I_L - I_R)/(I_L + I_R)$, where I_L and I_R are the output signals for left and right circularly Paper

polarised light under unpolarised photoexcitation conditions. The CPL efficiency, $|g_{\rm CPL}|$, of **BPP** in the KBr pellet is of the order of 10^{-3} , while the $|g_{\rm CPL}|$ of **BPP** is $\sim 2.0 \times 10^{-3}$ in the KBr dispersed state.

To examine molecular chirality in the ground state, the CD and UV-Vis absorption spectra of (R,R)-**BPP** and (S,S)-**BPP** in their KBr-pellet-dispersed states were acquired, the results of which are shown in Fig. 5.

Several UV-Vis bands between 400 and 650 nm, which are characteristic of π - π * transitions in the pervlene groups of **BPP**, can be observed. In this state, the Cotton CD bands of (R,R)-BPP and (S,S)-BPP, are almost mirror images and are due to $\pi - \pi^*$ transitions, as evidenced by their UV-Vis spectra. The CD spectra of (R,R)-BPP and (S,S)-BPP are almost unaffected by KBrpellet dispersion. Interestingly, the CD spectra of the KBrdispersed states are quite different to those of the CHCl₃ solutions are similar to those of the previously reported nanoscale aggregates in the solution-dispersed state.4g We conclude that the CD spectra of the KBr-dispersed states also result from intermolecular interactions between the pervlene units of multiple aggregated BPP molecules within the KBr pellets. To quantitatively evaluate the CD amplitude in the ground state, we evaluated the anisotropy factor. The magnitude of the circular polarisation in the ground state is defined as $g_{\rm CD} = ({\rm Abs}_{\rm L} - {\rm Abs}_{\rm R})/[({\rm Abs}_{\rm L} + {\rm Abs}_{\rm R})/2]$, where ${\rm Abs}_{\rm L}$ and ${\rm Abs}_{\rm R}$ are the absorbances of the left and right circularly polarised light, respectively. The value of $|g_{CD}|$ at the first Cotton CD band of **BPP** in KBr was determined to be \sim 5.4 \times 10⁻⁴ at 561 nm. The observed CD results in KBr support the CPL observations mentioned above. Interestingly, a comparison of these anisotropy factors, *i.e.*, the $|g_{CPL}|$ and $|g_{CD}|$ values, reveals that the $|g_{\text{CPL}}|$ is larger than the corresponding $|g_{\text{CD}}|$ value, indicating that, although the perylene units are stacked and twisted in the ground state,^{4f} the twists increase in the photoexcited state, as confirmed by the intense excimer CPL bands.

In order to emit stronger CPL, we next examined CPL from organic polymer-film states. As an organic polymer matrix, poly(methyl methacrylate) (PMMA) and polyurethane (*myo*-IPU) derived from naturally occurring *myo*-inositol were used. The CPL and PL spectra of (R,R)-**BPP** and (S,S)-**BPP** dispersed in

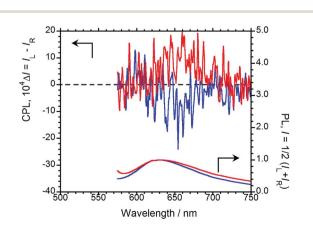


Fig. 4 CPL (upper) and PL (lower) spectra of (*R*,*R*)-**BPP** (blue) and (*S*,*S*)-**BPP** (red) in KBr-pellet-dispersed states.

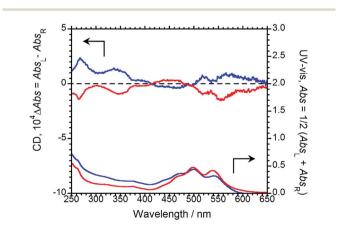


Fig. 5 CD (upper) and UV-Vis (lower) absorption spectra of (R,R)-**BPP** (blue) and (S,S)-**BPP** (red) in KBr-pellet-dispersed states.

PMMA and *myo*-IPU films are shown in Fig. 6(a) and (b), respectively.

(*R*,*R*)-**BPP** in PMMA emitted at 654 nm (λ_{CPL}) with a Φ_{F} of 0.40, while in *myo*-IPU, it emitted at 637 nm (λ_{CPL}) with a $\Phi_{\rm F}$ of 0.43. As expected, (R,R)-BPP exhibits solid-state AIEnh-CPL in both polymers, with several regressional phonon sidebands originating from the lowest π - π * transitions between perylene groups observed. The CPL spectra of (R,R)-BPP in PMMA and myo-IPU films mostly exhibited negative (-) CPL bands (Fig. 6(a) and (b), respectively). As expected, (S,S)-BPP exhibited positive (+) CPL spectra in these polymers (Fig. 6(a) and (b)). Although the CPL spectral shapes of (R,R)-BPP and (S,S)-BPP in the KBr-pellet and PMMA- and myo-IPU-film states are nearly mirror images, subtle differences in the shapes of the CPL signals are noted. These differences may be due to the effect of humidity in the pellets and films during the preparation of the pellet and film samples. The $|g_{CPL}|$ values of **BPP** in the PMMA and *myo*-IPU states are $\sim 2.4 \times 10^{-3}$ and $\sim 1.3 \times 10^{-3}$, respectively, which are similar to that observed in the KBr state (\sim 2.0 $\times 10^{-3}$).

As expected, despite the CPL wavelengths of (R,R)-**BPP** in the KBr and polymer states being similar, the values of $\Phi_{\rm F}$ in these polymers are significantly higher than that observed for the KBr state (0.09). We assume that these approximately five-times-higher values are due to the suppression of thermal deactivation modes and effective packing arrangements because the molecular geometry of the chiral **BPP** is effectively and suitably fixed by the surrounding solid polymer matrix without quenching.

The CD and UV-Vis absorption spectra of (R,R)-**BPP** and (S,S)-**BPP** in the PMMA and *myo*-IPU films are shown in Fig. 7(a) and (b), respectively; the CD spectra in these polymers are similar in

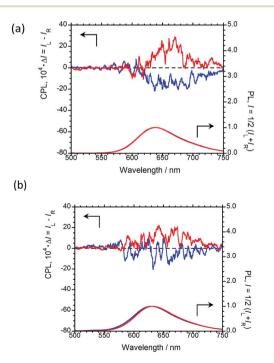


Fig. 6 CPL (upper) and PL (lower) spectra of (*R*,*R*)-**BPP** (blue) and (*S*,*S*)-**BPP** (red) in (a) PMMA-film- and (b) *myo*-IPU-film-dispersed states.

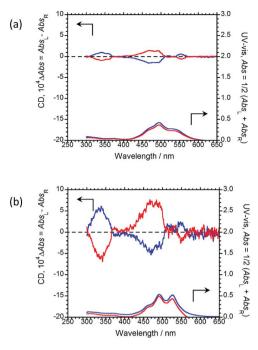


Fig. 7 CD (upper) and UV-Vis absorption (lower) spectra of (R,R)-BPP (blue) and (S,S)-BPP (red) in (a) PMMA-film- and (b) *myo*-IPU-film-dispersed states.

their long-wavelength tails and are similar to those of the KBrdispersed states. These results suggest that the CD spectra of the polymer-dispersed states are essentially due to intermolecular interactions between multiple molecules rather than intramolecular interactions within individual BPP molecules. The characteristic UV peaks originating from π - π * transitions of the pervlene groups are commonly observed between 400 and 600 nm. The CD spectra of (R,R)-BPP and (S,S)-BPP in both matrices are almost mirror images, as expected. The signs of the first Cotton CD bands of (R,R)-BPP, at 553 nm in PMMA and at 551 nm in myo-IPU, are both positive (+), while the CD signs of the corresponding bands for (S,S)-BPP in PMMS and myo-IPU are both negative (–). The $|g_{CD}|$ value at the first Cotton CD bands are ~3.1 \times 10⁻⁴ and ~2.9 \times 10⁻⁴ for the PMMA- and *myo*-IPU-dispersed states, respectively. BPP in these matrices exhibit $|g_{CD}|$ values of similar magnitude.

As was observed in the KBr matrix, the $|g_{CPL}|$ values are higher than the corresponding $|g_{CD}|$ values in the polymer films. This shows that intermolecular twists between perylene units in the ground state increase upon photo-excitation, as evidenced by the intense excimer CPL bands in spectra of the polymer-matrix imbedded samples.

Conclusions

Chiral (*R*,*R*)- and (*S*,*S*)-*N*,*N*'-bis(1-phenylethyl)perylene-3,4,9,10tetracarboxylic diimide (**BPP**) dispersed in solid KBr and PMMA (or *myo*-IPU) matrices exhibited solid-state AIEnh-CPL with high $\Phi_{\rm F}$ values at room temperature. The value of $\Phi_{\rm F}$ from the PMMA ($\Phi_{\rm F}$ 0.40) and *myo*-IPU matrices ($\Phi_{\rm F}$ 0.43) are significantly higher than that observed from the KBr matrix ($\Phi_{\rm F}$

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0.09). The present knowledge provides guidance for the design of novel solid chiral organic fluorophoric systems that emit solid-state AIEnh-CPL.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This study was supported by a Grant-in-Aid for Scientific Research (no. 18K05094) from MEXT/Japan Society for the Promotion of Science, the MEXT-Supported Program for the Strategic Research Foundation at Private Universities (2014–2018), and a research grant from the International Polyurethane Technology Foundation.

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