



Cite this: *Phys. Chem. Chem. Phys.*,
2024, 26, 28819

Quantum effects in CH activation with $[\text{Cu}_2\text{O}_2]^{2+}$ complexes†

Selin Bac ^a and Shaama Mallikarjun Sharada ^{*,ab}

We investigate the mechanism of primary alkane CH bond activation with dioxo-dicopper ($[\text{Cu}_2\text{O}_2]^{2+}$) complexes, which serve as model catalysts for enzymes capable of activating CH bonds under mild conditions. As large H/D kinetic isotope effects (KIEs) are observed in enzymes and their synthetic mimics, we employ density functional theory along with variational transition-state theory with multidimensional tunneling to estimate reaction rate coefficients. By systematically varying ligand electrophilicity and substrate chain length, we examine trends in rate coefficients and kinetic isotope effects for the two proposed CH activation pathways – one-step oxo-insertion and two-step radical recombination. Although larger tunneling transmission coefficients are obtained for the radical pathway, the oxo-insertion mechanism yields higher rate coefficients on account of lower activation barriers. The question of the preferred CH activation mechanism, however, remains open: excellent agreement is observed between the predicted and known experimental KIE results for the radical pathway, while calculated Hammett slopes for the oxo-insertion pathway closely mirror experiments.

Received 24th July 2024,
Accepted 3rd November 2024

DOI: 10.1039/d4cp02929a

rsc.li/pccp

1 Introduction

Efforts to streamline the conversion of natural gas into methanol as a cleaner alternative to traditional syngas-based methods have long been pursued, with the vision of establishing a methanol economy that could revolutionize energy and chemical industries.¹ A promising avenue is the development of catalysts capable of selectively activating strong C–H bonds, mimicking the functionality of enzymes found in nature, specifically that of the enzyme particulate methane monooxygenase (pMMO).^{2–5} While significant progress has been made in elucidating the reactivity of oxygen-activated metal complexes, specifically iron and copper complexes, challenges persist in interpreting the preferred mechanisms of CH activation.^{6–10}

Two mechanisms have been proposed for CH activation reactions catalyzed by dioxo-dicopper ($[\text{Cu}_2\text{O}_2]^{2+}$) complexes: (i) one-step oxo-insertion where the activation of a C–H bond occurs in a single concerted step alongside the insertion of an oxygen atom into the metal–carbon bond, and (ii) two-step radical recombination pathway, in which the first step involves homolytic cleavage of C–H bond (hydrogen atom transfer, HAT), leading to the formation of a methyl radical and

subsequent recombination with a hydroxyl group.^{11–15} Computational studies, primarily based on density functional theory (DFT), have contributed valuable insights into these mechanisms but often face challenges in capturing spin-dependence and treating multireference character of these systems.^{16–18}

Experimental studies probing the effects of substrate and catalyst variations on reaction kinetics can help validate computational predictions of the preferred reaction pathways. For instance, Hammett studies have shown enhanced reaction rates following substitutions of aromatic substrates with electron-donating groups.^{11,19,20} However, the observed experimental barriers in these studies often fall within a narrow range ($<10 \text{ kJ mol}^{-1}$), making it challenging for density functional approximations to resolve energy differences accurately. To overcome this limitation, we previously proposed a strategy that yields a wider range of barriers *via* multiple ligand substitutions in the catalyst instead.¹⁴ This method enables a semi-quantitative comparison between the proposed pathways and experimental Hammett plots,^{11,19–21} aiding in elucidating the true CH activation mechanism.

The kinetic isotope effect (KIE) serves as another means to probe mechanisms and contrast experiments with theory.²² CH functionalization reactions are characterized by large kinetic isotope effects originating in large part from hydrogen tunneling.^{22–30} The accurate quantification of tunneling effects is crucial for discerning the preferred mechanism, as tunneling can substantially contribute to catalytic activity. Traditional one-dimensional approximations are cost-effective³¹ but often are inaccurate due to their inability to capture the shape of the

^a Mork Family Department of Chemical Engineering and Materials Science, University of Southern California, Los Angeles, CA 90089, USA.

E-mail: ssharada@usc.edu

^b Department of Chemistry, University of Southern California, Los Angeles, CA 90089, USA

† Electronic supplementary information (ESI) available. See DOI: <https://doi.org/10.1039/d4cp02929a>

effective potential for tunneling and the phenomenon of “corner-cutting,” where the tunneling path deviates from the minimum energy path (MEP).³² One needs to employ approaches such as variational transition state theory with multidimensional tunneling (VTST/MT) to quantify these effects.^{33–35}

Our objective is to contrast rate coefficients obtained using VTST/MT with conventional transition state theory (without tunneling) to examine the importance of multidimensional tunneling in both proposed pathways for CH activation with imidazole-substituted $[\text{Cu}_2\text{O}_2]^{2+}$ complexes, and probe the sensitivity of reaction kinetics to systematic variations in the catalyst and substrate. Although the radical pathway yields significantly larger tunneling transmission coefficients than the oxo-insertion mechanism, the latter yields larger rate coefficients on account of lower activation barriers. Comparing experimental Hammett slopes with computations indicates preference for the oxo-insertion type mechanism with a partially cationic substrate in the transition state. On the other hand the magnitudes of kinetic isotope effects predicted for the radical pathway are in better agreement with experiments. Owing to these conflicting outcomes, the question of which of the two mechanisms is preferred remains unanswered. That being said, this work shows that multidimensional tunneling corrections are necessary to capture kinetics at low reaction temperatures.

2 Models and methods

2.1 Models

Along similar lines to previous studies by our group, the bis-(μ -oxo) isomeric form of the dicopper-dioxo active site is employed.^{14,36,37} We point the reader to our earlier work justifying this choice over the (multireference) peroxo active site as well as the choice of the singlet spin-potential energy surface.¹⁴ The Cu(III) centers are each bound to 2 N-donor imidazole ligands as shown in Fig. 1. We expand on prior work examining the impact of ligand electrophilicity on CH_4 activation barriers to include their role in governing multidimensional tunneling and kinetic isotope effects. This is carried out by substituting the hydrogen atom at the carbon atom positioned between the two nitrogen atoms in



Fig. 1 Transition structures for the first step of the (a) two-step radical recombination (TS_{rad}) and (b) one-step oxo-insertion mechanisms (TS_{oxo}). The active site is coordinated to imidazole (labeled ‘H’ in the ‘catalyst’ column of Table 1), and the substrate is CH_4 . Color scheme: cyan, carbon; dark blue, nitrogen; red, oxygen; brown, copper; and white, hydrogen. The visualizations are created using the visual molecular dynamics (VMD) software.⁴⁰

imidazole with OCH_3 , CH_3 , CF_3 , or NO_2 . The Hammett parameters for *para*-substitution to benzoic acid, σ_p , are reported to be -0.268 , -0.17 , 0 , 0.54 , 0.78 for OCH_3 , CH_3 , H , CF_3 , and NO_2 , respectively.^{38,39} A negative (positive) Hammett parameter indicates electron-donating (withdrawing) character of the substituent. To examine the impact of alkane chain length on the barrier to primary CH activation and tunneling transmission coefficients, we examine reactants ranging from CH_4 to C_5H_{12} .

2.2 Minimum energy paths

All density functional theory simulations are carried out using the *ab initio* quantum chemistry software, Q-Chem.⁴¹ Gas phase calculations are carried out using the $\omega\text{B97X-D}/\text{def2-SVP}^{42,43}$ level of theory and the def2-ECP effective core potential for Cu that includes scalar relativistic corrections.⁴⁴ Natural bond orbital (NBO) analysis^{45,46} and energy decomposition analysis (EDA)^{47,48} are used to characterize initial and transition structures. The latter are determined using the freezing string method^{49,50} and Hessian-free optimization⁵¹ and verified using vibrational analysis. We calculate minimum energy paths (MEPs) in mass-weighted coordinates, referred to as intrinsic reaction coordinates (IRCs), employing the gradient-based predictor-corrector algorithm implemented in Q-Chem.^{52–54} IRCs are calculated to ensure that the TS indeed connects the intended reactant and product and to construct a reaction path for VTST/MT rate coefficient calculations. The maximum IRC step size is deliberately chosen to ensure a smooth path, set at 0.1 or 0.125 atomic units (a.u.). The convergence threshold for IRC (labeled “RPATH_TOL_DISPLACEMENT” in Q-Chem) is set between 0.001–0.005 a.u. The threshold is chosen in such a way that the vibrational analyses of the IRC end points yield either all real frequencies or only a small number of imaginary frequencies, all smaller than $100i\text{ cm}^{-1}$.

We perform vibrational analysis for all geometries constituting the IRC to compute free energies and couplings necessary for multidimensional (small-curvature) tunneling corrections, described below. As has been shown in prior work by our group,^{14,36} wavefunction stability analysis reveals that stationary points (except the initial state) typically suffer from spin instability. While a spin-correction scheme such as that proposed by Yamaguchi and coworkers⁵⁵ is necessary to correct the instability and obtain spin-pure solutions, we limit ourselves to using unstable solutions. This is because spin-corrected energies are noisy and can lead to artificial peaks on the MEP (Fig. S1 of ESI†). Given the large number of systems examined in this study and the absence of a systematic protocol for the elimination of points yielding anomalous energies, we elect to use spin-unstable energies, while emphasizing that the reported rate coefficients do not represent true values. Trends in barriers remain largely unaffected by this choice (Table S1 of the ESI†).

2.3 Kinetic isotope effects

The kinetic isotope effect (KIE) is a unitless ratio that quantifies the change in reaction rate coefficient when one or more atoms in the reactants are replaced by its isotope. For reactions involving the transfer of H, the KIE is typically defined as the

ratio of the rate coefficient for the reaction with the lighter isotope (protium) to the rate constant for the reaction with the heavier isotope (deuterium), expressed as:

$$\text{KIE} = \frac{k_{\text{H}}}{k_{\text{D}}} \quad (1)$$

where k_{H} and k_{D} are the rate coefficients for the reactions involving protium and deuterium, respectively. The KIE is sensitive to the masses of the isotopes, with larger values often observed when the isotopic substitution influences both zero-point energies and quantum mechanical tunneling.

In experiments designed to elucidate mechanisms, the isotopic substitution of protium with deuterium is only carried out for the specific hydrogen atom(s) in the substrate participating in the reaction. Since Q-Chem only allows all hydrogen atoms to be substituted with deuterium in IRC calculations, our current approach involves perdeuteration. While this may limit the direct comparison of predicted KIEs with experimental results, we expect that trends arising from variations in ligands and chain lengths obtained using complete isotopic substitution will be similar to those observed when only the substrate hydrogen atoms are substituted. To precisely capture tunneling corrections, we perform IRC calculations and subsequent vibrational analyses for both the all-protium and all-deuterium systems.

2.4 Reaction rate coefficients

It is not guaranteed that the potential energies of the end-points of the protiated and deuterated IRCs are identical. It is also not necessary that the energies of the IRC end points exactly match those of the reactant and product minima employed to initiate the TS search. We therefore describe the free energy barrier, $\Delta G^\ddagger$, as the difference in Gibbs free energies between the highest point on the free energy profile and the initial state free energy.

To capture trends in kinetics across catalyst and substrate variations, we choose the initial (or reactant) state free energy as the sum of free energies of isolated catalyst and substrate. This is because energy decomposition analysis (EDA) reveals that the interaction energies between catalyst and substrate fragments in the reactant state described by a pre-association complex span a wide range (-12.9 to -48.0 kJ mol $^{-1}$) when the alkane chain length is varied. Directly comparing rate coefficients by assigning the reactant as a pre-association complex will therefore not capture the overall kinetics of CH activation. For this reason, we employ what is akin to an *apparent* free energy barrier that is often reported for heterogeneous catalytic systems.

Rate coefficients obtained from conventional transition state theory, referred to simply as TST,⁵⁶ are contrasted with variational transition state theory with multidimensional tunneling (VTST/MT).³² The key distinction between the two approaches, in this work, lies in the treatment of tunneling. Conventional TST does not account for quantum mechanical tunneling, and therefore, κ is unity. Within the VTST/MT framework, we calculate two types of tunneling transmission

coefficients: (i) zero-curvature tunneling (ZCT, κ_{ZCT}), which accounts for multidimensional tunneling but ignores the curvature of the reaction path,^{57–59} and (ii) small-curvature tunneling (SCT, κ_{SCT}), which incorporates the reaction path curvature, accounting for corner-cutting effects.⁶⁰ The procedure is identical to that reported in the Pilgrim software.⁶¹ The approach employs splines to interpolate the effective potential energy along the minimum-energy path. Gaussian quadrature is utilized for integration. Our Python implementation is tested by showing that the transmission coefficients for test systems employed in Pilgrim are within 5% of the Pilgrim values.

We report rate coefficients (s $^{-1}$) using the relation:

$$k(T) = \kappa(T) \frac{k_{\text{B}}T}{h} \exp\left(-\frac{\Delta G^\ddagger(T)}{RT}\right) \quad (2)$$

where T is the temperature, k_{B} , h , and R are the Boltzmann's, Planck's, and gas constant, respectively. The equation represents first-order kinetics, or in other words we assume constant catalyst concentration that is folded into the rate expression.

3 Results

Table 1 reports kinetics data obtained for the two proposed mechanisms by varying the ligands coordinated to the $[\text{Cu}_2\text{O}_2]^{2+}$ active site and the chain length of the alkane substrate. To understand the impact of tunneling at both low and room temperatures, transmission coefficients, free energies, and kinetic isotope effects are reported at low (200 K) and ambient temperatures (300 K).

3.1 Radical recombination and oxo-insertion mechanisms

Fig. 1 shows the transition structures corresponding to the oxo-insertion ('oxo') mechanism and the first, rate-limiting step of the radical recombination ('rad') mechanism for CH_4 activation with a complex consisting of four imidazole N-donor ligands (catalyst labeled 'H' in Table 1). In the oxo-insertion pathway, methanol is formed in a single step, with the transition state (TS_{oxo}) involving simultaneous Cu–O and C–H cleavage alongside C–O and O–H bond formation. TS_{oxo} exhibits a bent C–H–O configuration (122°) and is associated with an apparent barrier ($\Delta E^\ddagger$) of 95.8 kJ mol $^{-1}$. The radical pathway encompasses a two-step process, with the rate-limiting step involving CH cleavage to form radical species, followed by a recombination step to produce methanol. The first step (TS_{rad}) displays a nearly linear C–H–O configuration (171°) with an associated barrier of 174 kJ mol $^{-1}$. The imaginary frequencies corresponding to the reaction coordinate in the TSs are 879i and 1981i cm $^{-1}$ for TS_{oxo} and TS_{rad} , respectively. These values are consistent with our earlier study where we report 1075i and 1858i cm $^{-1}$ for **oxo** and **rad** TSs, respectively.¹⁴ Natural bond orbital (NBO) analysis reveals a difference in the total natural charge on CH_4 between the transition states and initial states of +0.59 and +0.11 for the **oxo** and **rad** mechanisms, respectively, in close agreement with previous work (+0.65 and +0.18).¹⁴ On account of its charged character, the **oxo** TS is therefore

Table 1 Calculated kinetic parameters for oxo and rad mechanisms at 200 and 300 K across all ligand-substitutions and substrates examined in this work. $\Delta E^\ddagger$ is the ZPE-corrected barrier and $\Delta G^\ddagger$ is the Gibbs free energy barrier calculated using the harmonic oscillator approximation and procedure described in previous work.⁶² Barrier values are reported in kJ mol⁻¹. Kinetic isotope effects (KIE_{SCT}) are calculated using κ_{SCT} values

Mechanism	Catalyst	Substrate	$\Delta E^\ddagger$	200 K				300 K			
				$\Delta G^\ddagger$	κ_{ZCT}	κ_{SCT}	KIE _{SCT}	$\Delta G^\ddagger$	κ_{ZCT}	κ_{SCT}	KIE _{SCT}
Rad	H	C ₅ H ₁₂	148.1	144.1	4.8	5.4	5.5	138.8	1.4	1.4	2.5
		C ₄ H ₁₀	148.3	143.8	1.3×10^2	1.4×10^2	31.8	138.1	3.2	3.3	5.4
		C ₃ H ₈	151.2	146.8	5.8×10^2	6.2×10^2	26.0	141.2	2.7	2.8	3.3
		C ₂ H ₆	154.7	150.9	3.2×10^2	3.5×10^2	27.1	145.7	1.9	1.9	2.6
		CH ₄	174.0	171.9	2.1×10^4	2.4×10^4	30.4	168.4	7.4	7.8	1.4
	OCH ₃	CH ₃	171.7	168.7	5.1×10^2	5.8×10^2	45.7	164.6	2.7	2.8	2.4
		CH ₃	172.4	170.8	3.1×10^4	3.4×10^4	58.5	167.9	6.6	6.9	2.5
		CF ₃	163.9	161.9	9.2×10^2	1.2×10^3	43.5	158.7	3.4	3.6	2.0
		NO ₂	169.0	166.8	7.8×10^2	9.1×10^2	36.7	163.5	3.1	3.2	1.7
Oxo	H	C ₅ H ₁₂	44.0	37.9	3.4	4.4	8.4	30.5	1.6	1.8	3.4
		C ₄ H ₁₀	47.2	40.6	3.7	4.6	9.2	32.9	1.7	1.9	3.7
		C ₃ H ₈	52.1	46.0	4.4	5.6	8.1	38.5	1.8	2.0	3.1
		C ₂ H ₆	54.2	49.0	4.0	5.1	7.3	42.2	1.7	1.9	2.5
		CH ₄	95.6	91.3	8.2	10.7	5.2	85.6	2.2	2.4	1.7
	OCH ₃	CH ₃	98.0	92.8	9.7	12.6	8.2	86.3	2.3	2.5	2.4
		CH ₃	105.4	102.5	10.7	13.9	2.3	98.1	2.4	2.6	0.8
		CF ₃	77.2	72.9	7.2	9.1	5.0	67.4	2.2	2.4	1.9
		NO ₂	66.0	62.0	4.8	5.9	4.5	56.8	2.0	2.1	1.8



Fig. 2 Arrhenius plots for the two-step radical recombination ('rad') and one-step oxo-insertion ('oxo') pathways across a temperature range from 25 to 500 K. The vertical colored lines separate regions dominated by quantum mechanical tunneling and by thermal effects. Deviations between the dashed and solid lines highlight the significant impact of tunneling contributions (SCT), particularly at lower temperatures. The slopes of the dashed lines are -9.32 and -5.30 for radical and oxo-insertion mechanisms, respectively.

expected to be more sensitive to the electrophilicity of ligands coordinated to the metal center.

Fig. 2 presents the Arrhenius plot of temperature-dependence of the rate coefficients for the two mechanisms of the imidazole-substituted catalyst with CH₄ as the reactant. In the temperature range of 25 K to 500 K, the rate coefficients for **oxo** are higher than **rad**. Quantum mechanical tunneling dominates the **oxo** and **rad** pathways below 220 K and 330 K, respectively, based on the sharp changes in slopes in Fig. 2. The **rad** pathway possesses tunneling transmission coefficients that are larger than **oxo**, with κ_{SCT} nearly 3 orders of magnitude larger at 200 K.

3.2 Chain length effects

Fig. 3 depicts the alkane chain length dependence of the VTST/MT and TST rate coefficients at 200 K for the two mechanisms.



Fig. 3 Rate coefficients (s⁻¹) at 200 K for **rad** (blue) and **oxo** (red) mechanisms, demonstrating the contributions of multidimensional tunneling (κ_{SCT}) across various substrate chain lengths. The darker circles represent rate coefficients calculated using variational transition state theory with multidimensional tunneling (VTST/MT), while the lighter circles correspond to those obtained from conventional transition state theory (TST) that excludes tunneling contributions.

In the case of the oxo-insertion mechanism, barriers drop steeply from CH₄ to C₂H₆, followed by a gradual decline. ZCT and SCT coefficients vary less strongly with chain length, with C2–C5 values about half of those obtained for CH₄ activation. As a result, the TST and VTST/MT rate coefficients are close to each other for C2–C5 and higher than the value obtained for CH₄.

TST rate coefficients are smaller than VTST/MT in the rad mechanism because multidimensional tunneling plays a more important role when compared to the oxo mechanism. There is a sharp decline in the barrier from CH₄ to C₂H₆ followed by a more gradual decrease from C₂H₆ to C₅H₁₂, reflected in the TST rate coefficients in Fig. 3. The trend in VTST/MT rate



Fig. 4 Kinetic isotope effects at 200 K for the radical recombination (KIE_{Rad}) and oxo-insertion (KIE_{Oxo}) mechanisms for varying substrate chain length. The data points represent KIE values calculated using VTST/MT that includes κ_{SCT} (darker circles) and conventional TST (lighter circles) that excludes tunneling contributions.

coefficients, however, is not monotonic because the increase associated with a drop in barrier is partially offset by the general decrease in tunneling transmission coefficients from C1 to C5.

We explore the trends in kinetic isotope effects for both mechanisms in Fig. 4 across varying substrate chain lengths at 200 K. Owing to larger contributions from multidimensional tunneling, the KIE values for the **rad** mechanism are consistently larger than those for the **oxo** mechanism. The TST-based KIEs are much smaller than the VTST/MT KIEs. While there are no significant differences in KIEs observed across C1–C4 substrates, the KIE drops to 5.1 in C5 due to smaller tunneling contributions. For the **oxo** mechanism, KIE values increase from C1 to C4, reaching their highest value for C4, and then show a slight decrease for C5. Unlike the **rad** mechanism, there is no sharp change in KIE between C4 and C5.

3.3 Ligand effects

We report the dependence of rate coefficients on ligand electrophilicity in Fig. 5. **Oxo** barriers exhibit pronounced ligand dependence, evidenced by a large ZPE-corrected apparent activation barrier range (39.3 kJ mol^{-1}). With the exception of OCH_3 , increasing electron-withdrawing character, quantified by the Hammett parameter for *para*-substitution σ_p , leads to lower barriers and therefore higher rate coefficients. We attribute the decrease in the barrier to the stabilization of the transition state by the electron-withdrawing ligand *via* charge transfer interactions.¹⁴ The tunneling transmission coefficients across these ligand-substituted systems all lie within an order of magnitude of each other and are smaller than 15 at 200 K. As a result, there are very small differences between TST and VTST/MT rate coefficients with varying ligand electrophilicity. We note that (with the exception of OCH_3), there is a monotonic decrease in both κ_{ZCT} and κ_{SCT} with increasing substituent electron-withdrawing character, *i.e.*, $\text{CH}_3 > \text{H} > \text{CF}_3 > \text{NO}_2$.



Fig. 5 Rate coefficients (s^{-1}) at 200 K calculated using conventional TST and VTST/MT (with κ_{SCT}) for the oxo and rad mechanisms obtained by varying electrophilicity of ligands coordinated to the active site.

As reported in an earlier study, **rad** is less sensitive to the ability of a ligand to push/pull electron density, and therefore, the barriers exhibit a narrower range (10.0 kJ mol^{-1}). The TST rate coefficients, therefore, vary by less than 2 orders of magnitude across the ligand-substituted systems. VTST/MT rate coefficients are higher than those obtained from TST, owing to large contributions from $\kappa_{\text{ZCT/SCT}}$.

Fig. 6 depicts the Hammett plots for the oxo-insertion pathway using the σ_p Hammett parameter, with the σ values on the x-axis scaled by a factor of four to reflect the number of ligand substitutions.¹⁴ Due to smaller tunneling coefficients observed for the oxo-insertion mechanism, the linear fit to TST and VTST/MT data are similar. In contrast, shown in Fig. S6 of the ESI,[†] the radical pathway shows no dependence on ligands, resulting in poor linear fit for VTST/MT ($R^2 < 0.03$) and only a modest improvement for TST ($R^2 = 0.5$, $\rho = -0.56$).

Despite differences in the choice of ligands and substrate, the slope of the Hammett plot for oxo-insertion, $\rho = -2.16$,



Fig. 6 Hammett plots at 200 K for the oxo-insertion pathway with substituents using Hammett parameters for *para*-substituted benzoic acid. The summation accounts for all four substitutions made in each catalyst, with the reference rate coefficient (k_{ref}) representing the catalyst labeled 'H' ($\sigma_p = 0$). The R^2 value and slopes (ρ) of the linear fit are reported for VTST/MT. For TST, without tunneling contributions, the slope (ρ) is -2.24 , and R^2 is 0.92.

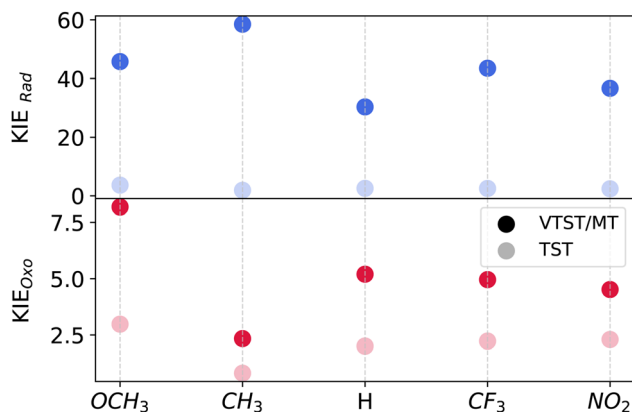


Fig. 7 A comparison of VTST/MT (with κ_{SCT}) and conventional TST-based kinetic isotope effects calculated at 200 K for oxo and rad mechanisms across different ligand substitutions.

is consistent with an electrophilic attack on the substrate reported in previous experimental studies (examining activation of weaker CH bonds) as well as prior work in our group.^{11,14,19,20} If we employ the *meta*-substituted σ_m Hammett parameter, the slope becomes -3.12 with an R^2 value of 0.96 (Fig. S5, ESI†). With σ_m as the electrophilicity descriptor, OCH₃ is electron-withdrawing ($\sigma_{m,\text{OCH}_3} = 0.12$) and no longer an outlier, and the fit is improved.

The impact of ligand electrophilicity on KIEs for oxo and rad mechanisms are shown in Fig. 7. Due to larger tunneling contributions, we observe larger VTST/MT KIEs for the **rad** pathway compared to **oxo** across all ligands. There are no discernible trends in **oxo** KIEs with variation in ligand electrophilicity.

4 Discussion

Our examination of C–H activation by $[\text{Cu}_2\text{O}_2]^{2+}$ complexes reveals that the singlet oxo-insertion pathway is energetically more favorable than the singlet radical pathway despite considerably larger tunneling coefficients and KIEs observed with the latter. We contrast these findings with experimental literature examining the activity of dioxo-dicopper complexes towards activating various CH bonds. A study of intramolecular hydroxylation of the benzyl CH bond with a $[\text{Cu}(\text{II})_2\text{O}_2]^{2+}$ center reported a substantial KIE of 35.4 at 193 K .⁶³ Similarly, a KIE of 31 ⁶⁴ was observed for the oxidation of 10-methyl-9,10-dihydroacridine (Arch₂) at 148 K . A computational study by Kim and co-workers examining the **rad** mechanism reports a VTST/MT-based KIE of 28 (277 K) for CH₄ hydroxylation with the $[\text{Cu}_2\text{O}_2]^{2+}$ active site in which each Cu(III) is bound to three N-donor amine groups. All these values closely resemble our VTST/MT-based KIE of 30.6 at 200 K for the **rad** mechanism, with the **oxo** yielding KIE = 5.2 .

Although KIEs indicate a possible **rad** pathway, barriers and Hammett plots conclude that the **oxo** pathway may be preferred. Applying spin- and zero-point corrections, the $\Delta E^\ddagger$ value for CH₄ activation with imidazole-bound $[\text{Cu}_2\text{O}_2]^{2+}$ *via* **oxo** is 25.2 kJ mol^{-1} , significantly lower than the **rad** barrier of

160.0 kJ mol^{-1} . While experimental Hammett curves are typically generated by varying the electrophilicities of substituents to an aromatic reactant, the resulting ranges in activation barriers are typically too narrow to be meaningfully interpreted using computations. Since each Cu(III) is coordinated to two N-donor ligands in this work, we can quadruple substituent effects by replacing an H- in the imidazole N-donor with an electron-donating or withdrawing group. The resulting range of zero point-corrected barriers for **oxo** is 39.4 kJ mol^{-1} . A reasonable linear fit and a finite, negative slope are observed only for the oxo-insertion pathway and not the radical pathway (Section S3 of the ESI†). Although the magnitudes of the slopes are also in agreement with experiment (ρ from -2.2 at 148 K to -1.48 at 193 K), this may be fortuitous because experimental studies are carried out for substrates containing weaker CH bonds than CH₄ and the Hammett plots are constructed using σ_p^+ values rather than σ_p .^{20,21,63}

Owing to these contrasting findings and differences in choices of substrates/catalysts in experiments and theory, it is difficult to draw conclusions regarding the preferred mechanism of CH activation. One possibility that we have not yet explored is the fact that the **rad** barrier in the triplet state is lower in energy than the singlet **rad** and singlet **oxo**, reported in an earlier study by our group.¹⁴ In other words, a two-state-reactivity (TSR) scenario, in which ligand-dependent crossover occurs from the singlet to the triplet state, remains to be explored.^{65–67} For instance, a study by Shaik and coworkers, of the **rad** mechanism-based CH activation with nonheme Fe(IV)-oxo bound to tetramethylcyclam ligands, showed that the unexpected faster kinetics with strong electron-donating ligands originates in a large probability of spin-inversion between the triplet and quintet states as well as enhanced tunneling.⁶⁸ Such a study for the dicopper system is hampered by the fact that spin-pure singlet states (unlike triplets or quintets) are difficult to determine on account of spin contamination in DFT, making it difficult to accurately identify crossing points with higher spin-states.

To explain the differences in mechanistic conclusions arising from comparing barriers, KIEs, and Hammett slopes, another testable hypothesis is that the distinction between the two mechanisms becomes blurry when the size of the substrate grows. Although rigorous testing of this hypothesis is a topic for future work, with an increase in alkane chain length in this study, we observe small changes in the ‘radical character’ of the TS. The total charge on the alkane fragment in the **rad** TS, calculated using NBO, increases by 55% from 0.11 (CH₄) to 0.17 (C₅H₁₂), indicating an increase in cationic character with increasing chain length. In contrast, the alkane fragment charge in the **oxo** TS increases only by 15% from 0.59 (CH₄) to 0.68 (C₅H₁₂). The imaginary frequency associated with the **rad** reaction coordinate decreases from 1981 cm^{-1} (CH₄) to 1733 cm^{-1} (C₅H₁₂), which in part explains the drop in tunneling transmission coefficients and KIEs with increasing chain length at 200 K (Table 1). To resolve the question of preferred mechanism therefore, we need a more extensive analysis of substrate sensitivity than what is reported in this work.

Our examination of ligand and chain length dependence is inspired by studies of CH activation with nonheme Fe(IV)-oxo complexes by Shaik and coworkers.^{23,68} In addition to TSR, they show that electron-donating ligands enhance tunneling transmission coefficients. If we assume that the substituents behave as if they are *meta*-substituted, the κ values for **oxo** decrease monotonically from CH₃ to NO₂. However, all transmission coefficients are within an order of magnitude of each other, and therefore the substituent effect on tunneling with **oxo** is not significant enough to impact rate coefficients. While a general decrease is also noted for the **rad** mechanism, with the CH₃-bound catalyst yielding κ_{SCT} that is two orders of magnitude larger than the NO₂-bound one, the decrease is not monotonic. The decrease in **oxo** and **rad** barriers with increasing chain length also leads to a decrease in tunneling transmission coefficients. Contrary to findings by Shaik and coworkers suggesting a volcano-type relationship between KIE and CH bond dissociation energy (BDE),²³ we do not find a clear correlation for either reaction pathway when the alkane chain length is varied. We note however that the range of homolytic BDEs is narrow in our work (~ 18.8 kJ mol⁻¹) compared to theirs (~ 83.7 kJ mol⁻¹).²³

Tunneling emerges as the primary factor driving the substantial kinetic isotope effects observed in C–H activation catalyzed by [Cu₂O₂]²⁺ complexes. The magnitudes of KIEs and their trends with chain length and ligand substitution exhibit at least a two-fold increase when zero- or small-curvature tunneling is incorporated (Fig. S2 of ESI†). In addition, overlapping KIEs between SCT and ZCT-based rate coefficient calculations indicate that corner-cutting does not significantly impact most systems. Therefore, even though prior computational studies of CH activation with [Cu₂O₂]²⁺ complexes estimate κ_{SCT} 's,^{22,27} the zero-curvature approximation is expected to suffice. Mandal and Shaik²³ also demonstrated, for CH activation with nonheme Fe(IV)-oxo complexes, that a simple Eckart tunneling model yields KIE values in good agreement with experiments and multidimensional models. The parity plot in Fig. S3 of the ESI† also illustrates this, with the exception of systems in which κ_{SCT} values exceed 20 000.

5 Conclusions

By examining the singlet potential energy surface of CH activation with [Cu₂O₂]²⁺ complexes, this study aims to uncover the dependence of the rate coefficients, tunneling transmission coefficients, and kinetic isotope effects of two proposed mechanisms – one-step oxo-insertion and two-step radical recombination – on catalyst electrophilicity and substrate CH bond strength. To this end, we employ DFT simulations and contrast rate coefficients and KIEs obtained using conventional TST without tunneling and VTST with multidimensional tunneling. We find that the use of multidimensional tunneling approximations is necessary to capture isotope effects in these systems, although the role of corner-cutting appears to be small. The barriers for the oxo-insertion pathway are always

lower than those for radical recombination. However, while calculated Hammett slopes are in agreement with experiment for the oxo pathway, the KIEs for the **rad** pathway align better with experimentally observed isotope effects. Future work includes exploration of two-state reactivity and expansion in substrate scope to address these conflicting mechanistic conclusions.

Author contributions

SB carried out the computational work described in this study. SMS led the design of simulations and both authors contributed to manuscript writing.

Data availability

The data supporting this article have been included as part of the ESI.†

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work is supported by the U.S. Department of Energy, Office of Science, Office of Basic Energy Sciences, under Award No. DE-SC0021417 and the USC Viterbi School of Engineering's Ershaghi Center for Energy Transition (E-CET). The authors acknowledge computational resources and support from USC's Center for Advanced Research Computing (CARC) and the National Energy Research Scientific Computing Center (NERSC). The authors thank Kaustubh Rane (USC) for images of the transition structures.

Notes and references

- 1 G. A. Olah, A. Goepfert and G. S. Prakash, *Beyond oil and gas: The methanol economy*, John Wiley & Sons, 2011.
- 2 C. E. Elwell, N. L. Gagnon, B. D. Neisen, D. Dhar, A. D. Spaeth, G. M. Yee and W. B. Tolman, *Chem. Rev.*, 2017, **117**, 2059–2107.
- 3 E. A. Lewis and W. B. Tolman, *Chem. Rev.*, 2004, **104**, 1047–1076.
- 4 L. M. Mirica, X. Ottenwaelde and T. D. P. Stack, *Chem. Rev.*, 2004, **104**, 1013–1046.
- 5 N. Kitajima and Y. Moro-oka, *Chem. Rev.*, 1994, **94**, 737–757.
- 6 P. E. Siegbahn and R. H. Crabtree, *J. Am. Chem. Soc.*, 1997, **119**, 3103–3113.
- 7 C. Krebs, D. Galonic Fujimori, C. T. Walsh and J. M. Bollinger Jr, *Acc. Chem. Res.*, 2007, **40**, 484–492.
- 8 W. Nam, *Acc. Chem. Res.*, 2015, **48**, 2415–2423.
- 9 C. J. Cramer, M. Włoch, P. Piecuch, C. Puzzarini and L. Gagliardi, *J. Phys. Chem. A*, 2006, **110**, 1991–2004.

- 10 C. J. Cramer, A. Kinal, M. Wloch, P. Piecuch and L. Gagliardi, *J. Phys. Chem. A*, 2006, **110**, 11557–11568.
- 11 S. Mahapatra, J. A. Halfen and W. B. Tolman, *J. Am. Chem. Soc.*, 1996, **118**, 11575–11586.
- 12 V. C.-C. Wang, S. Maji, P. P.-Y. Chen, H. K. Lee, S. S.-F. Yu and S. I. Chan, *Chem. Rev.*, 2017, **117**, 8574–8621.
- 13 V. W. Bowry and K. Ingold, *J. Am. Chem. Soc.*, 1991, **113**, 5699–5707.
- 14 Z. Lan and S. M. Sharada, *Phys. Chem. Chem. Phys.*, 2018, **20**, 25602–25614.
- 15 Y. Shiota, G. Juhász and K. Yoshizawa, *Inorg. Chem.*, 2013, **52**, 7907–7917.
- 16 K. Yoshizawa and Y. Shiota, *J. Am. Chem. Soc.*, 2006, **128**, 9873–9881.
- 17 Y. Shiota and K. Yoshizawa, *Inorg. Chem.*, 2009, **48**, 838–845.
- 18 S. Ye and F. Neese, *Inorg. Chem.*, 2010, **49**, 772–774.
- 19 L. M. Mirica, M. Vance, D. J. Rudd, B. Hedman, K. O. Hodgson, E. I. Solomon and T. D. P. Stack, *Science*, 2005, **308**, 1890–1892.
- 20 C. Citek, C. T. Lyons, E. C. Wasinger and T. D. P. Stack, *Nat. Chem.*, 2012, **4**, 317–322.
- 21 K. D. Karlin, M. S. Nasir, B. I. Cohen, R. W. Cruse, S. Kaderli and A. D. Zuberbuehler, *J. Am. Chem. Soc.*, 1994, **116**, 1324–1336.
- 22 Y. Kim, B. K. Mai and S. Park, *JBIC, J. Biol. Inorg. Chem.*, 2017, **22**, 321–338.
- 23 D. Mandal and S. Shaik, *J. Am. Chem. Soc.*, 2016, **138**, 2094–2097.
- 24 J. Kaizer, E. J. Klinker, N. Y. Oh, J.-U. Rohde, W. J. Song, A. Stubna, J. Kim, E. Münck, W. Nam and L. Que, *J. Am. Chem. Soc.*, 2004, **126**, 472–473.
- 25 M. P. Meyer and J. P. Klinman, *Chem. Phys.*, 2005, **319**, 283–296.
- 26 E. Derat, D. Kumar, H. Hirao and S. Shaik, *J. Am. Chem. Soc.*, 2006, **128**, 473–484.
- 27 K. Park, Y. Pak and Y. Kim, *J. Am. Chem. Soc.*, 2012, **134**, 3524–3531.
- 28 C. Citek, S. Herres-Pawlis and T. D. P. Stack, *Acc. Chem. Res.*, 2015, **48**, 2424–2433.
- 29 J. P. Klinman, *Biochim. Biophys. Acta, Bioenerg.*, 2006, **1757**, 981–987.
- 30 J. P. Layfield and S. Hammes-Schiffer, *Chem. Rev.*, 2014, **114**, 3466–3494.
- 31 R. T. Skodje and D. G. Truhlar, *J. Phys. Chem.*, 1981, **85**, 624–628.
- 32 A. Fernandez-Ramos, B. A. Ellingson, B. C. Garrett and D. G. Truhlar, *Rev. Comput. Chem.*, 2007, **23**, 125.
- 33 D. G. Truhlar and B. C. Garrett, *Annu. Rev. Phys. Chem.*, 1984, **35**, 159–189.
- 34 D. G. Truhlar and B. C. Garrett, *Acc. Chem. Res.*, 1980, **13**, 440–448.
- 35 J. L. Bao and D. G. Truhlar, *Chem. Soc. Rev.*, 2017, **46**, 7548–7596.
- 36 Z. Lan and S. M. Sharada, *Phys. Chem. Chem. Phys.*, 2020, **22**, 7155–7159.
- 37 Z. Lan and S. M. Sharada, *Phys. Chem. Chem. Phys.*, 2021, **23**, 15543–15556.
- 38 L. P. Hammett, *J. Am. Chem. Soc.*, 1937, **59**, 96–103.
- 39 C. Hansch, A. Leo and R. Taft, *Chem. Rev.*, 1991, **91**, 165–195.
- 40 W. Humphrey, A. Dalke and K. Schulten, *J. Mol. Graphics*, 1996, **14**, 33–38.
- 41 E. Epifanovsky, A. T. B. Gilbert, X. Feng, J. Lee, Y. Mao, N. Mardirossian, P. Pokhilko, A. F. White, M. P. Coons, A. L. Dempwolff, Z. Gan, D. Hait, P. R. Horn, L. D. Jacobson, I. Kaliman, J. Kussmann, A. W. Lange, K. U. Lao, D. S. Levine, J. Liu, S. C. McKenzie, A. F. Morrison, K. D. Nanda, F. Plasser, D. R. Rehn, M. L. Vidal, Z.-Q. You, Y. Zhu, B. Alam, B. J. Albrecht, A. Aldossary, E. Alguire, J. H. Andersen, V. Athavale, D. Barton, K. Begam, A. Behn, N. Bellonzi, Y. A. Bernard, E. J. Berquist, H. G. A. Burton, A. Carreras, K. Carter-Fenk, R. Chakraborty, A. D. Chien, K. D. Closser, V. Cofer-Shabica, S. Dasgupta, M. de Wergifosse, J. Deng, M. Diedenhofen, H. Do, S. Ehlert, P.-T. Fang, S. Fatehi, Q. Feng, T. Friedhoff, J. Gayvert, Q. Ge, G. Gidofalvi, M. Goldey, J. Gomes, C. E. Gonzalez-Espinoza, S. Gulania, A. O. Gunina, M. W. D. Hanson-Heine, P. H. P. Harbach, A. Hauser, M. F. Herbst, M. Hernandez Vera, M. Hodecker, Z. C. Holden, S. Houck, X. Huang, K. Hui, B. C. Huynh, M. Ivanov, A. Jasz, H. Ji, H. Jiang, B. Kaduk, S. Kahler, K. Khistyayev, J. Kim, G. Kis, P. Klunzinger, Z. Koczor-Benda, J. H. Koh, D. Kosenkov, L. Koulias, T. Kowalczyk, C. M. Krauter, K. Kue, A. Kunitsa, T. Kus, I. Ladjanszki, A. Landau, K. V. Lawler, D. Lefrancois, S. Lehtola, R. R. Li, Y.-P. Li, J. Liang, M. Liebenthal, H.-H. Lin, Y.-S. Lin, F. Liu, K.-Y. Liu, M. Loipersberger, A. Luenser, A. Manjanath, P. Manohar, E. Mansoor, S. F. Manzer, S.-P. Mao, A. V. Marenich, T. Markovich, S. Mason, S. A. Maurer, P. F. McLaughlin, M. F. S. J. Menger, J.-M. Mewes, S. A. Mewes, P. Morgante, J. W. Mullinax, K. J. Oosterbaan, G. Paran, A. C. Paul, S. K. Paul, F. Pavosevic, Z. Pei, S. Prager, E. I. Proynov, A. Rak, E. Ramos-Cordoba, B. Rana, A. E. Rask, A. Rettig, R. M. Richard, F. Rob, E. Rossomme, T. Scheele, M. Scheurer, M. Schneider, N. Sergueev, S. Mallikarjun Sharada, W. Skomorowski, D. W. Small, C. J. Stein, Y.-C. Su, E. J. Sundstrom, Z. Tao, J. Thirman, G. J. Tornai, T. Tsuchimochi, N. M. Tubman, S. P. Veccham, O. Vydrov, J. Wenzel, J. Witte, A. Yamada, K. Yao, S. Yeganeh, S. R. Yost, A. Zech, I. Y. Zhang, X. Zhang, Y. Zhang, D. Zuev, A. Aspuru-Guzik, A. T. Bell, N. A. Besley, K. B. Bravaya, B. R. Brooks, D. Casanova, J.-D. Chai, S. Coriani, C. J. Cramer, G. Cserey, A. E. DePrince, R. A. DiStasio, A. Dreuw, B. D. Dunietz, T. R. Furlani, W. A. Goddard, S. Hammes-Schiffer, T. Head-Gordon, W. J. Hehre, C.-P. Hsu, T.-C. Jagau, Y. Jung, A. Klamt, J. Kong, D. S. Lambrecht, W. Liang, N. J. Mayhall, C. W. McCurdy, J. B. Neaton, C. Ochsenfeld, J. A. Parkhill, R. Peverati, V. A. Rassolov, Y. Shao, L. V. Slipchenko, T. Stauch, R. P. Steele, J. E. Subotnik, A. J. W. Thom, A. Tkatchenko, D. G. Truhlar, T. Van Voorhis, T. A. Wesolowski, K. B. Whaley, H. L. Woodcock, P. M. Zimmerman, S. Faraji, P. M. W. Gill, M. Head-Gordon, J. M. Herbert and A. I. Krylov, *J. Chem. Phys.*, 2021, **155**, 084801.

- 42 J.-D. Chai and M. Head-Gordon, *Phys. Chem. Chem. Phys.*, 2008, **10**, 6615–6620.
- 43 J.-D. Chai and M. Head-Gordon, *J. Chem. Phys.*, 2008, **128**, 084106.
- 44 F. Weigend and R. Ahlrichs, *Phys. Chem. Chem. Phys.*, 2005, **7**, 3297–3305.
- 45 J. P. Foster and F. Weinhold, *J. Am. Chem. Soc.*, 1980, **102**, 7211–7218.
- 46 G. Frenking, K. Wichmann, N. Fröhlich, C. Loschen, M. Lein, J. Frunzke and V. M. Rayón, *Coord. Chem. Rev.*, 2003, **238**, 55–82.
- 47 P. R. Horn, Y. Mao and M. Head-Gordon, *Phys. Chem. Chem. Phys.*, 2016, **18**, 23067–23079.
- 48 Y. Mao, D. S. Levine, M. Loipersberger, P. R. Horn and M. Head-Gordon, *Phys. Chem. Chem. Phys.*, 2020, **22**, 12867–12885.
- 49 A. Behn, P. M. Zimmerman, A. T. Bell and M. Head-Gordon, *J. Chem. Phys.*, 2011, **135**, 224108.
- 50 S. M. Sharada, P. M. Zimmerman, A. T. Bell and M. Head-Gordon, *J. Chem. Theory Comput.*, 2012, **8**, 5166–5174.
- 51 S. M. Sharada, A. T. Bell and M. Head-Gordon, *J. Chem. Phys.*, 2014, **140**, 164115.
- 52 K. Fukui, *J. Phys. Chem.*, 1970, **74**, 4161–4163.
- 53 K. Ishida, K. Morokuma and A. Komornicki, *J. Chem. Phys.*, 1977, **66**, 2153–2156.
- 54 M. W. Schmidt, M. S. Gordon and M. Dupuis, *J. Am. Chem. Soc.*, 1985, **107**, 2585–2589.
- 55 K. Yamaguchi, F. Jensen, A. Dorigo and K. Houk, *Chem. Phys. Lett.*, 1988, **149**, 537–542.
- 56 D. G. Truhlar, B. C. Garrett and S. J. Klippenstein, *J. Phys. Chem.*, 1996, **100**, 12771–12800.
- 57 A. Kuppermann and D. G. Truhlar, *J. Am. Chem. Soc.*, 1971, **93**, 1840–1851.
- 58 B. Garrett, D. Truhlar, R. Grev and A. Magnuson, *J. Phys. Chem.*, 1983, **87**, 4554.
- 59 B. C. Garrett, D. G. Truhlar, R. S. Grev and A. W. Magnuson, *J. Phys. Chem.*, 1980, **84**, 1730–1748.
- 60 Y. P. Liu, G. C. Lynch, T. N. Truong, D. H. Lu, D. G. Truhlar and B. C. Garrett, *J. Am. Chem. Soc.*, 1993, **115**, 2408–2415.
- 61 D. Ferro-Costas, D. G. Truhlar and A. Fernández-Ramos, *Comput. Phys. Commun.*, 2020, **256**, 107457.
- 62 S. Bac, S. J. Quiton, K. J. Kron, J. Chae, U. Mitra and S. M. Sharada, *J. Chem. Phys.*, 2022, **156**, 184119.
- 63 S. Itoh, M. Taki, H. Nakao, P. L. Holland, W. B. Tolman, L. Que, Jr and S. Fukuzumi, *Angew. Chem., Int. Ed.*, 2000, **39**, 398–400.
- 64 C. Citek, B.-L. Lin, T. E. Phelps, E. C. Wasinger and T. D. P. Stack, *J. Am. Chem. Soc.*, 2014, **136**, 14405–14408.
- 65 D. Schröder, S. Shaik and H. Schwarz, *Acc. Chem. Res.*, 2000, **33**, 139–145.
- 66 S. Shaik, S. P. De Visser, F. Ogliaro, H. Schwarz and D. Schröder, *Curr. Opin. Chem. Biol.*, 2002, **6**, 556–567.
- 67 B. Yang, L. Gagliardi and D. G. Truhlar, *Phys. Chem. Chem. Phys.*, 2018, **20**, 4129–4136.
- 68 D. Mandal, R. Ramanan, D. Usharani, D. Janardanan, B. Wang and S. Shaik, *J. Am. Chem. Soc.*, 2015, **137**, 722–733.