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## Impact of dynamic covalent chemistry and precise linker length on crystallization kinetics and morphology in ethylene vitrimers†

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Vitrimers, dynamic polymer networks with topology conserving exchange reactions, have emerged as a promising platform for sustainable and reprocessable materials. While prior work has documented how dynamic bonds impact stress relaxation and viscosity, their role on crystallization has not been systematically explored. Precise ethylene vitrimers with 8, 10, or 12 methylene units between boronic ester junctions were investigated to understand the impact of bond exchange on crystallization kinetics and morphology. Compared to linear polyethylene which has been heavily investigated for decades, a long induction period for crystallization is seen in the vitrimers ultimately taking weeks in the densest networks. An increase in melting temperatures ( $T_m$ ) of 25–30 K is observed with isothermal crystallization over 30 days. Both  $C_{10}$  and  $C_{12}$  networks initially form hexagonal crystals, while the  $C_{10}$  network transforms to orthorhombic over the 30 day window as observed with wide angle X-ray scattering (WAXS) and optical microscopy (OM). After 150 days of isothermal crystallization, the three linker lengths led to double diamond ( $C_8$ ), orthorhombic ( $C_{10}$ ), and hexagonal ( $C_{12}$ ) crystals indicating the importance of precision on final morphology. Control experiments on a precise, permanent network implicate dynamic bonds as the cause of long-time rearrangements of the crystals, which is critical to understand for applications of semi-crystalline vitrimers. The dynamic bonds also allow the networks to dissolve in water and alcohol-based solvents to monomers, followed by repolymerization while preserving the mechanical properties and melting temperatures.

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## Introduction

The development of sustainable and recyclable plastics is a major challenge to handle the massive volume of products that end up in landfills.<sup>1</sup> Polymer networks, either elastomers or thermosets, are not processable with traditional chemistries but can be made recyclable *via* the incorporation of dynamic covalent bonds into the polymer. Notably, Leibler and coworkers described dynamic polyester networks which exhibit glass-like processability due to conserved ester exchange reactions, now commonly called vitrimers.<sup>2–5</sup> This concept has been widely applied to a range of dynamic bonds<sup>2,6–19</sup> including boronic esters and boroxine.<sup>20–29</sup> This ever growing toolbox of dynamic

covalent bonds which conserve the network topology is described in recent review articles.<sup>30,31</sup> A major focus thus far has been to understand how the exchange kinetics, crosslink density, and polymer backbone chemistry control the stress relaxation and reprocessability of vitrimers. The glass transition temperature ( $T_g$ ) and a hypothetical topology freezing temperature ( $T_v$ ) have been discussed to understand the interplay with bond exchange processes which governs the temperature dependent viscoelasticity. In contrast, an understanding of crystallization phenomena and melting temperatures ( $T_m$ ) in vitrimers is currently lacking. Dynamic bonds provide a mechanism for chain rearrangements within a polymer network, as well as a new timescale which can potentially facilitate crystal perfection and growth. Knowledge of how dynamic bonds impact crystallization kinetics and morphology will be critical to the development of new polymers which are easier to recycle and reprocess while still retaining desirable physical and optical properties.

The crystallinity of commodity polymers such as polyethylene (PE) and polypropylene is key to their electrical breakdown strength,<sup>32</sup> mechanical properties,<sup>33–35</sup> and thermal conductivity.<sup>36</sup> Electron<sup>37,38</sup> and ion<sup>39–41</sup> conducting functional polymers also show

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performance which is critically related to the crystal structure and amorphous fraction. In all cases, processing conditions can lead to vastly different material properties depending on the timescales of material relaxation and rate of deformation during pressing or extrusion. With the introduction of dynamic bonds in a network, an additional timescale is generated which may facilitate crystal formation by allowing smaller strands to rearrange rather than an entire polymer chain. An understanding of how bond exchange affects kinetics and final morphology is currently lacking. While some prior work has focused on semi-crystalline vitrimers based on polyethylene,<sup>18,23–25,27,42,43</sup> and poly(lactic acid),<sup>44</sup> they did not investigate any temporal evolution of crystallinity or morphology due to dynamic bond rearrangements within the matrix. Little is known about the role of bond exchange on crystallization in dynamic networks.

Another key factor in determining crystal structure and melting is the presence of precise motifs. Here the term precise refers to the exact number of carbons in the linker following a definition used in prior work on “precise” polymers.<sup>45,46</sup> For example, the melting temperatures of telechelic alkanes show a pronounced odd–even effect depending on the number of carbons between functional groups.<sup>47–49</sup> In polymers, precise polyethylenes<sup>50</sup> and polyacetals<sup>51</sup> have been made as linear polymers, including materials with periodic ionic groups which crystallize and show enhanced conductivity.<sup>52,53</sup> Precise permanent networks with alkane chain linkers have shown odd–even effects on the glass transition temperature and ionic conductivity.<sup>54</sup> All of these studies point to the potentially critical role of precision on crystallization, which has not been investigated in vitrimers. In addition to being important from an application point of view, understanding the fundamental roles of precision and dynamic bond exchange on semi-crystalline polymers could provide new insights into the design and application of such materials.

Here, we report the synthesis and systematic investigation of ethylene vitrimers with precise carbon spacing between dynamic boronic ester crosslinks. Boronic esters are a popular dynamic bond because of the relative ease by which they can be incorporated into conventional polymers. In addition, boronic esters tend to be a fast exchanging bond, even at ambient temperatures and without catalyst.<sup>29</sup> This allows one to observe the effects of dynamic bond exchange on crystallization at room temperature. This model vitrimer was chosen because linear PE has been widely studied and the crystal structure evolution is well known.<sup>55,56</sup> The initial crystallization of vitrimers is inhibited due to the network architecture, and longtime crystallization studies indicate that  $T_m$  grows continually over 30 days. An unexpected crystal–crystal phase transition is also observed *via* X-ray scattering, and is further tracked *via* polarized optical microscopy. This morphological transition, as well as the long time evolution of  $T_m$ , is not present in linear PE and a precise, permanent ethylene network shows no crystallization even after one week, implicating dynamic bonds as the origin of this phenomena. This is an extension of our previous work on the rheological investigation of ethylene vitrimers in the amorphous state.<sup>29</sup>

## Results and discussion

Precise ethylene vitrimers were synthesized by the step growth polymerization of telechelic alkane diols and boric acid following our prior work.<sup>29</sup> Networks with exactly 8, 10 or 12 carbons ( $C_8$ ,  $C_{10}$ ,  $C_{12}$ ) between boronic ester crosslinking junctions were prepared by mixing the diol and acid, followed by heating to melt the diol and then application of vacuum to drive off water condensate (Fig. 1a). The term precise only refers to the carbon spacing, not the network topology, following prior work.<sup>29,57</sup> Fourier transform infrared spectroscopy (FTIR) of the networks was performed at 80 °C to confirm the absence of the broad OH peak in the 3000  $\text{cm}^{-1}$  to 3500  $\text{cm}^{-1}$  range corresponding to unreacted species on the diols, boric acid, and water. Within the sensitivity of the instrument, no OH peak is detected indicating that the network formation reached high conversion for all samples. While the samples can take up 4–5 wt% of water under ambient conditions, they still form self-standing networks. Nevertheless, all crystallization was done in a glove box to avoid issues of moisture on kinetics. We do not claim 100% conversion of the network because the FTIR instrument would not be able to detect  $\sim 0.1\%$  impurities within the sample. Additionally, the purity of the monomers varies from 98% to 99% which is accounted for in our stoichiometric calculations. Thus free OH groups are thought to be present below the detection limit which enable exchange reactions in the network.<sup>29</sup> It is worth noting that boronic esters can also exchange *via* methathesis reactions<sup>27</sup> without free alcohol groups which can also lead to bond exchange. The emergence of a sharp peak at 1300  $\text{cm}^{-1}$  corresponds to the asymmetric B–O stretch<sup>58</sup> indicating the formation of the boronic ester crosslinks. Differential scanning calorimetry (DSC) was performed using a rapid quench from 150 to  $-80$  °C, and only a glass transition is observed in  $C_8$  and  $C_{10}$  networks on the cooling curve while  $C_{12}$  shows a weak crystallization peak. In contrast, a linear PE standard crystallizes instantly (Fig. S1b and c, ESI<sup>†</sup>).

The networks are termed vitrimers due to the well-known conserved exchange reaction of boronic esters in small molecules<sup>59</sup> and polymers.<sup>20–28</sup> In our prior work, oscillatory shear rheology was used to probe the storage ( $G'$ ) and loss ( $G''$ ) modulus of the ethylene networks in their fully amorphous state to determine the characteristic timescale for flow.

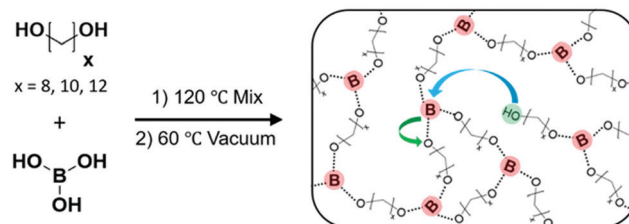


Fig. 1 Step growth polymerization of telechelic alkane diols and boric acid is carried out to make ethylene vitrimers. By using diol with an exact number of carbon atoms, networks with precise spacing between crosslinks were synthesized. Schematic has been reproduced from ref. 29 with permission.

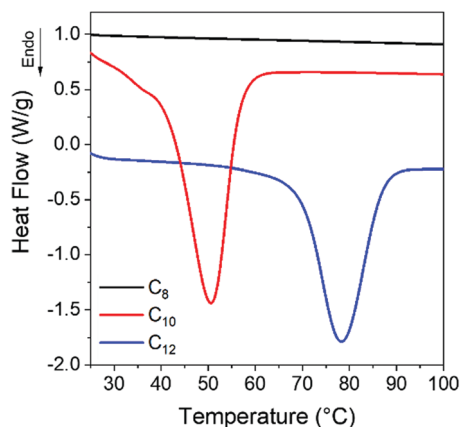


Fig. 2 After 1 day of isothermal crystallization at room temperature,  $C_{10}$  and  $C_{12}$  show clear melting peaks whereas  $C_8$  still exhibits no melting transition.

Frequency sweeps were performed in  $10^\circ\text{C}$  intervals from  $140$ – $40^\circ\text{C}$  with networks showing rubbery behavior at higher frequencies and terminal relaxation at lower frequencies (Fig. S1d, ESI $^\dagger$ ) due to dynamic bond exchange.<sup>29</sup> As expected, the modulus tracked with crosslink density/linker length. For a constant crosslink density, the modulus should increase with temperature due to entropic elasticity,<sup>60</sup> which is observed in the ethylene networks supporting the presence of topology conserving exchange reactions. We have previously observed that denser networks have longer relaxation times and slightly higher activation energies which is a function of both bond exchange kinetics and also the ability of dynamic bonds to find each other. A detailed discussion on the trends of relaxation times can be found in our previous report on amorphous ethylene dynamic networks.<sup>29</sup> The viscosity was calculated from the slope of  $G''$  vs.  $\omega$  in the low frequency limit and plotted against inverse temperature, which showed the anticipated Arrhenius behavior. This method of zero-shear viscosity determination has been compared to the complex viscosity and gives the same values and temperature dependent

trends.<sup>29</sup> All of these measurements are consistent with a network held together by associative dynamic covalent bonds.

The glass transition temperatures ( $T_g$ ) of the networks increases monotonically from  $-43^\circ\text{C}$  ( $C_8$ ) to  $-41^\circ\text{C}$  ( $C_{10}$ ) to  $-32^\circ\text{C}$  ( $C_{12}$ ), meaning there is mobility to allow the networks to crystallize at room temperature. Crystallization is observed in  $C_{12}$  networks on the cooling curve (Fig. S1b, ESI $^\dagger$ ), indicating that crystals will be present on first heating. Conversely, the  $C_8$  network shows no initial crystallinity while the  $C_{10}$  and  $C_{12}$  networks have a small melting transition at  $25$  and  $35^\circ\text{C}$ . A separate batch of samples was prepared and quenched directly to room temperature from  $150^\circ\text{C}$ . In the  $C_8$  network, no melting was observed on the initial heating ramp (Fig. S2, ESI $^\dagger$ ) while  $C_{10}$  and  $C_{12}$  gradually develop a melting peak at room temperature over  $300$  min. After 1 day of isothermal crystallization at room temperature, both  $C_{10}$  and  $C_{12}$  networks show a clear melting transition with larger enthalpy and shifted to higher temperature relative to the samples directly quenched to  $-80^\circ\text{C}$ , whereas the  $C_8$  vitrimer remains amorphous (Fig. 2). We attribute the latter result to the fact that this vitrimer has the highest crosslink density, which inhibits the ability of linkers to pack, as well as the highest viscosity which reduces translational mobility of network strands.

The melting transitions of  $C_{10}$  and  $C_{12}$  vitrimers were monitored with DSC as a function of isothermal crystallization time at room temperature over  $30$  days, and a monotonic increase in the peak melting temperature ( $T_m$ ) of the networks was observed (Fig. 3a and b). The  $C_{10}$  network shows a melting peak splitting around  $10\,080$  min, which eventually coalesces after  $43\,200$  min. In contrast,  $C_{12}$  vitrimers only show one peak for the entire series. Optical microscopy, to be discussed in detail subsequently, shows the growth of dendritic structures which appear only in the  $C_{10}$  networks and are assigned as the cause of peak splitting. A sigmoidal curve is typically seen for processes governed by a nucleation and growth mechanism,<sup>61,62</sup> which is observed in the  $C_{12}$  networks when plotting the peak melting temperatures over  $43\,200$  min (Fig. 3c). The  $C_{10}$  network has not yet plateaued even after

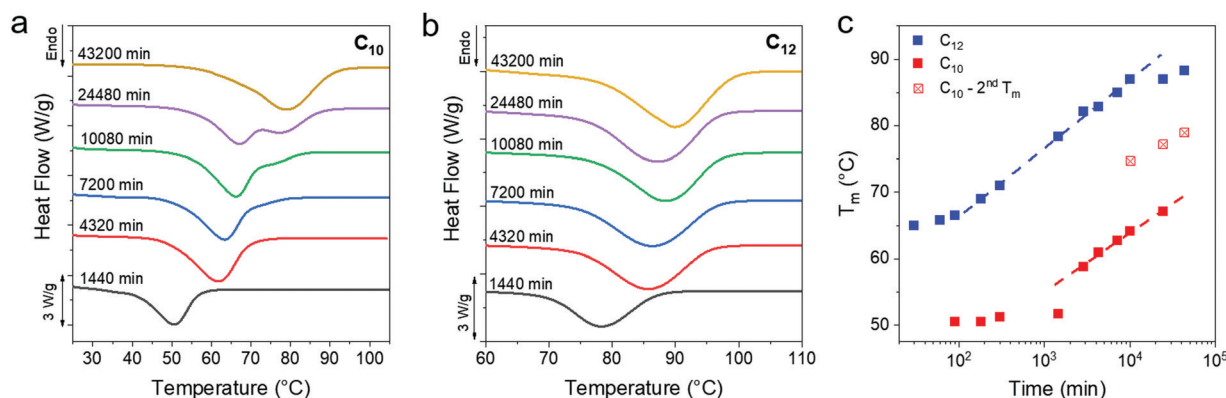


Fig. 3 (a and b) DSC of  $C_{10}$  and  $C_{12}$  after long time crystallization at room temperature. A monotonic increase in the  $T_m$  is evident from the rightward shift of the melting curve. DSC data for less than 1 day of crystallization time can be found in Fig. S2 (ESI $^\dagger$ ). (c) A plot of  $T_m$  vs. time shows similar slopes for the two networks. The 7, 17 and 30 day samples of  $C_{10}$  show two melting peaks, the lower one shown as solid squares and the higher peak as crossed squares.

$\sim 10^4$  min. In both cases, there is an initial induction period where  $T_m$  is flat, followed by an upturn at 100 and 1000 min for the  $C_{10}$  and  $C_{12}$  networks, respectively. The slow crystallization kinetics in  $C_{10}$  as compared to  $C_{12}$  can be rationalized because  $T_m$  is lower for  $C_{10}$  networks, thus they have a smaller thermodynamic driving force for crystallization. Additionally, a higher crosslink density in  $C_{10}$  vitrimers means that the system is more frustrated which would also impede crystallization and  $T_m$  evolution.

Over the observed period of 43 200 min a  $\Delta T_m$  of 25 °C for  $C_{12}$  and 30 °C for  $C_{10}$  is observed. An increase in the enthalpy of melting (and thus crystallinity) accompanies the  $T_m$  evolution, and nearly doubles for the  $C_{12}$  network (Fig. S3a, ESI†). The  $C_{10}$  network shows an  $\sim 70\%$  increase in enthalpy over the same time period. Only the relative change in crystallinity is reported and not the absolute crystallinity, as there is no appropriate reference enthalpy. Polyethylene consists of only methylene repeat units, while the present networks include a periodic placement of boronic ester crosslinks which will affect the enthalpy of fusion. Ethylene vitrimer crystals undergo a polymorphic transition, discussed later, which makes an absolute degree of crystallization determination difficult. After 150 days, the  $T_m$  reaches 82 °C and 91 °C for the  $C_{10}$  and  $C_{12}$  networks, respectively (Fig. S3b, ESI†). There are no free OH

groups within the resolution of the instrument after crystallization (Fig. S3d, ESI†). Thus, we conclude that the boronic esters are incorporated into the vitrimers crystals.

At intermediate times the  $T_m$  shows a power law dependence with crystallization time. Generally, a power law slope is attributed to the logarithmic growth of the lamellar thickness which in turn causes a logarithmic growth in  $T_m$ .<sup>63</sup> The apparent power law dependence was fit to the equation  $T_m(t_A) = T_0 D \log(t_A/t_0)$  used in prior work where  $t_A$  is annealing time and  $t_0$  is the initial time.<sup>64</sup> The  $D$  values were 10.3 and 8.8 for  $C_{12}$  and  $C_{10}$  networks using only data beyond 180 and 1440 min, respectively. These values are similar and attributed to the fact that the same dynamic bond exchange is controlling the growth of the crystal. Future work will incorporate dynamic bonds with much faster exchange reactions to understand how  $D$  and the time evolution of  $T_m$  are affected.

More analysis is required for a detailed molecular explanation, but we presently assign the slow evolution and emergence of crystals to dynamic bond facilitated reorganization of strands in the amorphous region into the crystalline phase. Different dynamic bonds with variable exchange kinetics are expected to yield different slopes for  $T_m$  evolution. Two key control experiments were performed, the first on linear PE (SRM 1475, 52 000 g mol<sup>-1</sup>) which shows instantaneous

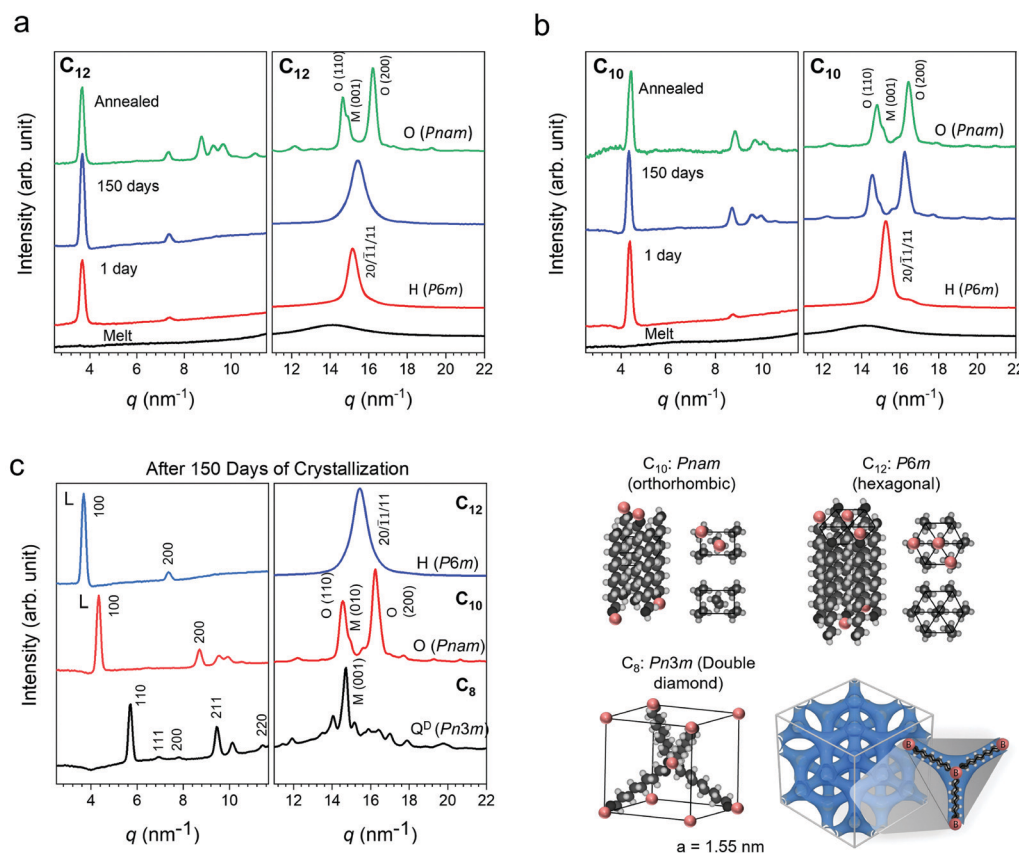


Fig. 4 (a) and (b) Room temperature WAXS patterns of  $C_{10}$  and  $C_{12}$  networks crystallized in a glove box. Samples were sandwiched and sealed between Kapton to prevent exposure to moisture. After 150 days of crystallization, samples were annealed at 50 °C ( $C_{10}$ ) and 65 °C ( $C_{12}$ ) for 2 days for the “annealed” data set. (c) A comparison of the WAXS patterns after 150 days of crystallization. (d) Schematic representation of the hexagonal ( $P6m$ ), orthorhombic ( $Pnam$ ) and double diamond ( $Pn3m$ ) crystal lattices adopted by the  $C_{12}$ ,  $C_{10}$ , and  $C_8$  polymers, respectively.

crystallization and no substantial long-term  $T_m$  evolution (Fig. S1c, ESI†). This is expected as PE exhibits fast crystallization kinetics, and  $T_m$  evolution has only been observed using flash differential scanning calorimetry with time resolution as fast as  $10^{-4}$  s.<sup>65</sup> The second control involved the synthesis of a precise, permanent ethylene network with  $C_{10}$  linkers as described in the ESI† (Fig. S4). In the absence of dynamic bonds, this network does not crystallize even after a week while the dynamic analogue shows crystallization within the first day (Fig. S5a, ESI†). A linear  $C_{10}$  polymer with the same amide linkages was also synthesized which crystallizes immediately upon cooling and indicates that the network topology, not the specific junction chemistry, are preventing crystallization (Fig. S5b, ESI†). Thus, the presence of dynamic bonds is critical to the melting temperature evolution in semi-crystalline vitrimers.

The long time evolution of  $T_m$  is atypical for polymer crystals. Increases in properties like  $T_m$ , density and lamellar thickness have been previously reported for polyethylene<sup>66,67</sup> and polyethylene terephthalate<sup>64,68–71</sup> but only with annealing or crystallization at elevated pressure. Slow reordering would require molecular motions of the chains in the amorphous regions which can reorganize to add to the crystalline interface. Such a reorganization in local structure would be possible only if mobility is available for strands to translate. Is it noteworthy that previous reports on increasing  $T_m$  emphasize the role of an external stimuli, while the present networks sit quiescently at room temperature. In addition, these reports are on linear polymers and little is known about the role of network architectures on crystallization, particularly at such high crosslink densities.<sup>72</sup>

As mentioned, along with the upward shift in  $T_m$  the melting curves for  $C_{10}$  show a transition from a single peak to a double peak and then revert to a single peak (Fig. 3a). The second peak is visible first as a small shoulder in the 5 day measurement and progressively increases in intensity in the 10 080 min and 24 480 min measurements at the expense of the first peak. This single to double peak transition is not seen in  $C_{12}$ . To probe the nature of this phenomena, wide-angle X-ray scattering (WAXS) was performed on the networks. Patterns collected over 1440 min (1 day) and 150 days of isothermal crystallization time at room temperature in an Argon glove box are shown in Fig. 4. Scattering patterns were also collected more frequently on a separate benchtop XRD at room temperature for samples crystallized under ambient conditions (Fig. S6, ESI†) and show the same general phenomena to be discussed. The presence of atmospheric moisture slightly modifies the kinetics of crystal evolution, while the patterns in Fig. 4a–c correspond to samples crystallized in a glovebox. The final crystal structures are not affected by ambient water. Prior work on PE has indicated that the semi-crystalline state at room temperature has three types of unit cells: orthorhombic, monoclinic and hexagonal.<sup>73–75</sup> The orthorhombic PE crystal phase is the densest and most commonly observed phase in PE, while the monoclinic PE crystal structure may coexist with orthorhombic.<sup>50,53,57</sup> The hexagonal phase is typically seen under high pressure crystallization.<sup>76,77</sup> The X-ray scattering patterns at 1 day show a single Bragg peak at  $15 \text{ nm}^{-1}$  for both  $C_{10}$  and  $C_{12}$  arising

Table 1 Location of low  $q$  scattering peak and analysis of peak width using the Scherrer equation

Network	$q$ ( $\text{nm}^{-1}$ )	$d$ (nm)	All-trans linker length (nm)	Correlation length, $\epsilon$ (nm)
$C_8$	5.71	1.09	1.23	31.5
$C_{10}$	4.33	1.44	1.54	28.1
$C_{12}$	3.68	1.70	1.84	27.4

from the diffraction of planes (20), ( $\bar{1}1$ ), and (11) of a PE hexagonal unit cell (lattice parameters:  $a = 4.7 \text{ \AA}$ ,  $b = 8.2 \text{ \AA}$ ;  $P6m$  space group). With increasing isothermal crystallization at room temperature, a transition to a two peak pattern is observed for  $C_{10}$  (Fig. 4a). The two Bragg peaks at  $14.8 \text{ nm}^{-1}$  and  $16.4 \text{ nm}^{-1}$  can be respectively assigned to the diffraction of planes (110) and (200) of a PE orthorhombic unit cell (lattice parameters:  $a = 7.1 \text{ \AA}$ ,  $b = 4.9 \text{ \AA}$ ,  $c = 2.5 \text{ \AA}$ ; in a  $Pnam$  space group).<sup>78</sup> The  $C_{12}$  network only shows one peak after 150 days in the glovebox but annealing at  $65 \text{ }^\circ\text{C}$  for 2 days can drive a transition to the orthorhombic phase (Fig. 4b). In both cases, the networks evolve from a metastable phase given sufficient time and temperature to orthorhombic with a coexistence of monoclinic. In prior work on ionic telechelic ethylenes, a hexagonal to orthorhombic transition was observed on cooling and the monoclinic phase coexisted with orthorhombic phase, which may explain the shoulder observed in the scattering result of the 150 days and annealed samples. In contrast, a linear PE standard immediately forms an orthorhombic crystal structure (Fig. S7, ESI†). In the low  $q$  range the lamellar (100) and (200) reflections are observed for  $C_{10}$  and  $C_{12}$  indicating formation of lamellar stacks. The X-ray scattering patterns were also collected for  $C_8$ , which shows a single dominant peak even after 150 days of isothermal crystallization (Fig. 4c). Analysis of the scattering data indicates a tetrahedral arrangement of short rods, or the  $C_8$  alkyl chain, interconnected in a cubic crystal lattice with a lattice parameter of  $a = 1.55 \text{ nm}$  and symmetry of a  $Pn3m$  space group. A double diamond cubic crystal structure has not been previously reported in polyethylene or in polymer networks. This network showed only a faint melting peak by DSC after 30 days (Fig. S3c, ESI†) due to the extremely slow kinetics. The crystal structures and proposed packing of the linkers is illustrated in Fig. 4d. The peaks for all three networks shift to higher  $q$  with annealing indicating densification of the crystalline networks, and a key finding of this work is that the crystal structure depends strongly on the choice of linker length in precise, dynamic networks.

A comparison of the scattering patterns for 150 days-crystallized networks reveals a monotonic shift of the lower  $q$  peak with increasing linker length. This is attributed to the boron-boron correlation length along the backbone based on a comparison with the calculated  $d$ -spacing ( $d = 2\pi/q$ ) which is slightly shorter than that of an all-trans configuration of the ethylene linkers (Table 1). The correlation length of the crystals was also calculated using the Scherrer equation  $\epsilon = 2\pi/\Delta q$  where  $\Delta q$  is the full width at half maximum of the peak corresponding to the spacer length. This length is  $\sim 30 \text{ nm}$  and decreases slightly as the linker length increases.

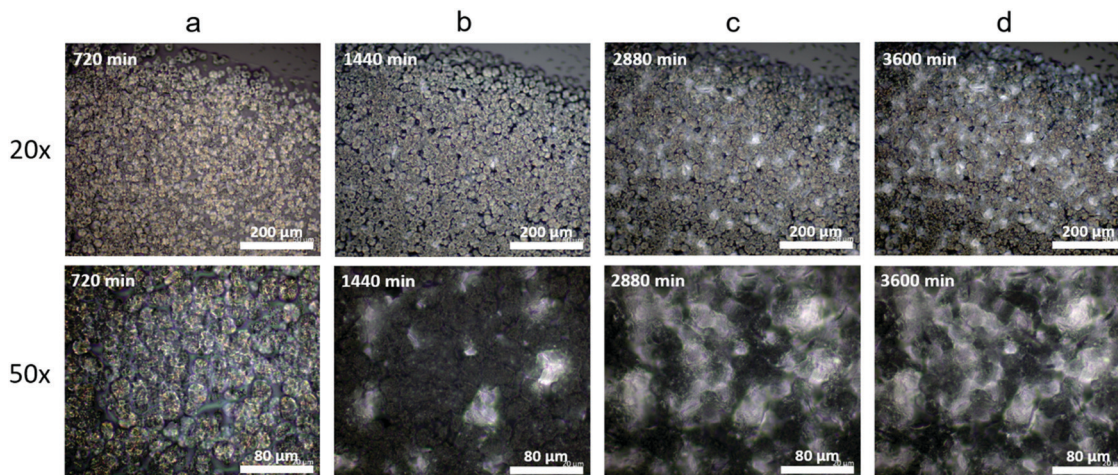


Fig. 5 Cross polarization microscopy of  $C_{10}$ , crystallization at room temperature under ambient condition. We follow the evolution in crystal structure as a function of time (from left to right). (a) Initial formation of spherulites after 720 min of crystallization time. (b) A new crystal structure is seen alongside the spherulites at 1440 min. (c) and (d) At 2400 min and 3600 min, the new crystals grow and uniformly span the image while the size of the spherulites remains constant with time.

To confirm that the network is intact after months of crystallization, FTIR spectra were collected on samples crystallized and tested inside the glovebox (Fig. S3d, ESI<sup>†</sup>). No free OH groups are observed indicating an intact network. Additionally, the scattering patterns of the starting alkane diols are shown in Fig. S9b (ESI<sup>†</sup>), which are clearly distinct from the network patterns. Thus, the alkanes are not disengaging from the network and crystallizing separately over time.

An increase in  $T_m$  can occur for multiple reasons. First, the growth of crystallites raises the melting temperature due to the well-known Gibbs-Thomson effect and surface destabilization

of a small crystal.<sup>79–82</sup> Second, a decreasing number of defects improves the crystal perfection<sup>64</sup> and decreases the free energy, leading to an increase in  $T_m$ . The long-term evolution of  $T_m$  in vitrimers is a consequence of the bond exchange reactions allowing for local rearrangements of network strands in the amorphous regions or at the crystal–amorphous interfaces which increase both the size and perfection of the crystals. To examine morphological evolution over time, cross polarization microscopy was used to observe changes in the crystal structure. The formation of uniformly distributed spherulites is observed in the  $C_{10}$  network during the initial 720 min of crystallization (Fig. 5a) which is typical of linear polymers. After 1440 min of crystallization under ambient conditions, a new crystal structure emerges (Fig. 5b). With time these crystals proliferate, while the growth of the spherulites is arrested (Fig. 5c and d). Eventually the new crystals dominate and the spherulites are no longer observable. We assign the initial crystals as a metastable hexagonal phase which transitions into the final orthorhombic crystal. The morphological development observed in optical microscopy is qualitatively consistent with the WAXS evolution further supporting a crystal–crystal transition in the absence of external stimuli and occurring over multiple days. Quantitatively relating these slow transformations to dynamic bond rearrangements is an important avenue for future work.

In select regions of the initial micrographs for  $C_{10}$  networks, dendritic structures are observed alongside the spherulites (Fig. 6a and b). When annealed at 50 °C for 2880 min, the dendrites grow and eventually span the entire crystal domain (Fig. 6c and d). Comparing the morphological changes with the appearance of the second melting peak in  $C_{10}$ , we assign the second melting peak in DSC to the percolated dendritic crystal once it grows large enough to contribute to the heat flow. Compared to  $C_{10}$ , the crystal morphology of  $C_{12}$  networks is noticeably different by optical microscopy (Fig. S8, ESI<sup>†</sup>).

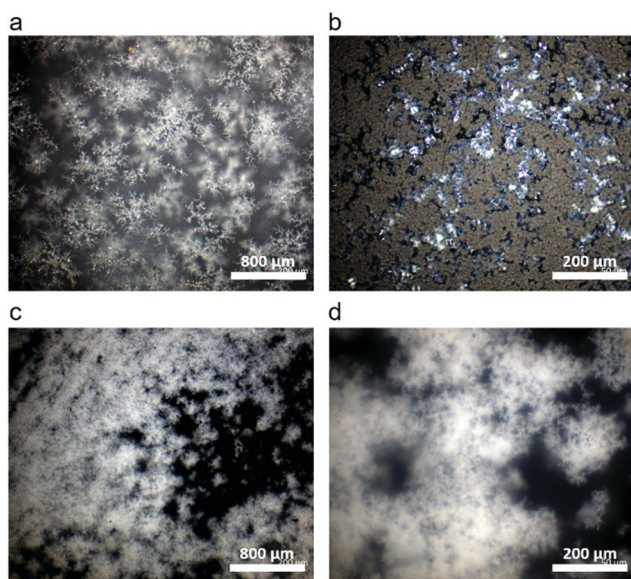


Fig. 6 (a) and (b) Dendritic structures are seen in other locations of the same  $C_{10}$  sample along with the crystals shown in Fig. 5. (c) and (d) After annealing at 50 °C for 2 days the dendrites span the entire crystal domain and the spherulites are no longer visible.

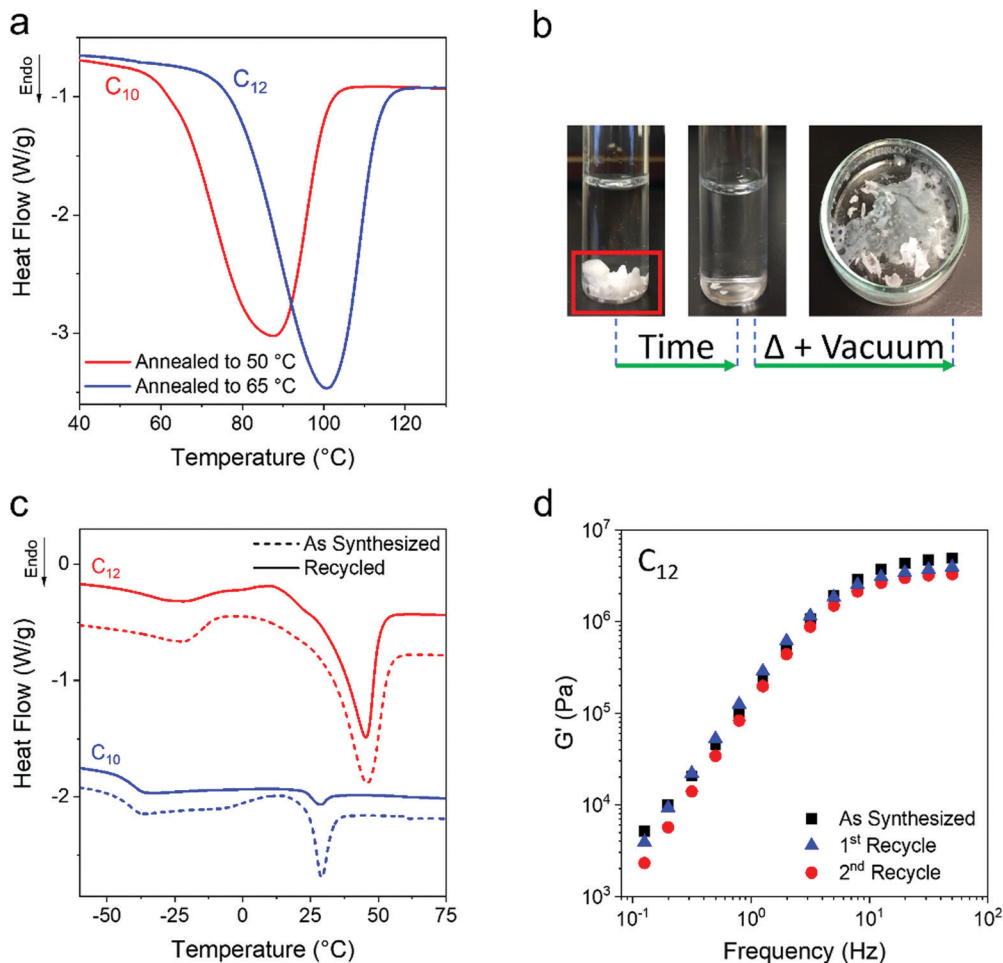


Fig. 7 (a) DSC of the annealed samples. After annealing for 2 days an increase in  $T_m$  to 89 °C and 100 °C is observed for C<sub>10</sub> and C<sub>12</sub> respectively. (b) Networks were dissolved in ethanol and recovered by solvent evaporation. (c) DSC of as synthesized vs. recycled C<sub>10</sub> and C<sub>12</sub>. The  $T_g$  and  $T_m$  of the recovered networks match the virgin networks. (d) Comparison of the rheology of recycled C<sub>12</sub> and as synthesized C<sub>12</sub> network indicates that there is minimal degradation in material properties.

Spherulites are not resolvable in C<sub>12</sub> networks indicating their size is below the resolution of our microscope. Nevertheless, WAXS confirms the presence of the same initial hexagonal crystal. Dendrites are not observed in C<sub>12</sub> at any time in agreement with the absence of a second melting peak in C<sub>12</sub> network DSC curves.

One open question for the present ethylene vitrimers is how high the  $T_m$  can be with annealing protocols. Higher temperature annealing of the C<sub>10</sub> and C<sub>12</sub> networks at 50 °C and 65 °C, respectively, was performed after the initial isothermal crystallization for 43 200 min at room temperature. Annealing for 2 days resulted in an increase of  $T_m$  to 89 °C for C<sub>10</sub> and 100 °C for C<sub>12</sub> (Fig. 7a). Future work will investigate more complex annealing protocols for manipulating the final percent crystallinity and  $T_m$ . It is worth noting that the evolved  $T_m$  of the networks is above the  $T_m$  of their corresponding monomers and the WAXS patterns of the diols do not match those of the ethylene networks (Fig. S9, ESI†). This indicates that the alkane diols are not simply being expelled from the network and crystallizing as the small molecule.

An attractive property of vitrimers is their recyclability and the ability to recover monomer in some cases.<sup>20,83–87</sup> The present

ethylene vitrimers are crosslinked by boronic esters, and water or alcohol based solvent can attack the boron and dissolve the network. As a proof of concept, ethanol was used to dissolve both a C<sub>10</sub> and C<sub>12</sub> network, followed by removal of solvent under heat and vacuum (Fig. 7b). The recovered networks were tested by DSC and oscillatory shear rheology to compare properties of freshly prepared networks ('as synthesized') to the recovered networks. As shown in Fig. 7c, the same  $T_g$  and  $T_m$  are reported for the recovered networks indicating that semi-crystalline vitrimers can be easily broken down to monomer and recovered. Fig. 7d shows that the rheological properties of the networks are mostly retained, although the modulus does drop from 4.9 to 3.2 MPa which may be due to the volatility of boric acid at elevated temperature. Future materials can incorporate bonds which are more or less sensitive to the desired solvent for recovery.

## Conclusions

Ethylene vitrimers with boronic ester bonds are shown to evolve their crystal morphology and  $T_m$  over long time periods relative

to linear PE or permanent networks. With increasing crosslink density, the initial crystallization is frustrated and there is a longer induction period before  $T_m$  increases logarithmically with isothermal annealing time. In the  $C_{10}$  networks, a crystal–crystal transformation is observed by WAXS from hexagonal to orthorhombic structure which does not occur in linear PE under ambient conditions. Depending on linker length, either double diamond ( $C_8$ ), orthorhombic ( $C_{10}$ ), or hexagonal ( $C_{12}$ ) crystals are formed after 150 days of isothermal room temperature crystallization in a glovebox indicating the importance of precision on the final structure. The long-time evolution of  $T_m$ , morphology transformation, and dendrite formation are all attributed to dynamic bond exchange as they are not present in linear PE or a precise, permanent ethylene network control sample. Annealing the vitrimers at higher temperatures following the initial crystallization leads to a further increase in  $T_m$  which reaches 100 °C for a  $C_{12}$  network. Vitri-mer recyclability is demonstrated by dissolving the networks in common laboratory solvents and then recovered by solvent evaporation. These findings on the role of dynamic bond exchange on crystallization are critical to the design, processing, and use of vitrimers as both commodity polymers, and functional materials, which are more sustainable.

## Conflicts of interest

There are no conflicts to declare.

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