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Fine tuning of intra-lattice electron transfers through site doping in tetraoxolene-bridged iron honeycomb layers<sup>†</sup>

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The precise control of intra-lattice multiple electron transfers was demonstrated in the solvated and desolvated species of the tetraoxolene-bridged Fe honeycomb layer system,  $(NPr_4)_2[Fe_2(Cl_2An)_3]$ . (solv) ( $Cl_2An^{n-} = 2,5$ -dichloro-3,6-dihydroxy-1,4-benzoquinonate;  $NPr_4^+ =$  tetrapropylammonium cation), by the site-doping of the  $Cl_2An^{n-}$  bridging unit using  $X_2An^{n-}$  units with X = Br or F.

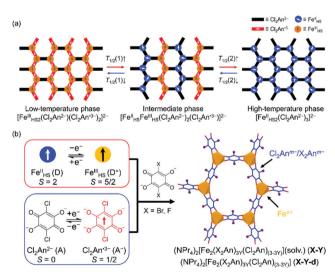
Exploring new stimuli-controllable functional systems is important in the field of materials science.<sup>1–4</sup> These systems can be obtained by achieving thermally driven electron transfers (TDETs) in a material,<sup>5</sup> which can successfully induce a state (or phase) change in materials at a transition temperature ( $T_{1/2}$ ). Moreover, modifications of the physical properties of materials such as electronic conductivity,<sup>6</sup> magnetic properties,<sup>7</sup> elasticity,<sup>8</sup> and the dielectric property can be attained.<sup>9</sup> Currently, the relevant systems in the charge-transfer metal complexes, which are composed of electron-donor (D) and -acceptor (A) moieties,<sup>9–20</sup> and organic charge transfer systems,<sup>21–24</sup> have been investigated. Nevertheless, the systematic fine tuning of the variable charge states is still challenging; however, it can potentially be used for the application of these systems in molecular devices and sensors.<sup>1–3</sup>

Recently, we have reported two kinds of tetraoxolenebridged iron honeycomb layered compounds,  $(NPr_4)_2[Fe_2 (Cl_2An)_3] \cdot 2(acetone) \cdot H_2O$  (**0**;  $Cl_2An^{n-} = 2,5$ -dichloro-3,6dihydroxy-1,4-benzoquinonate;  $NPr_4^+ =$  tetrapropylammonium cation)<sup>15</sup> and their desolvated compound  $(NPr_4)_2[Fe_2(Cl_2An)_3]$ (**0-d**).<sup>16</sup> The former exhibited a TDET at  $T_{1/2a} = 236$  K between the phases of  $[Fe^{III}_2(Cl_2An^{2-})(Cl_2An^{\cdot 3-})_2]^{2-}$  at  $T < T_{1/2a}$  and  $[Fe^{II}Fe^{II}(Cl_2An^{2-})_2(Cl_2-An^{\bullet 3-})]^{2-}$  at  $T > T_{1/2a}$ ,<sup>15</sup> while the latter underwent a stepwise TDET at  $T_{1/2b} = 317$  K and  $T_{1/2c} = 354$  K which changed phases from  $[Fe^{III}_2(Cl_2An^{2-})(Cl_2An^{\bullet 3-})_2]^{2-}$  at  $T < T_{1/2b}$  to  $[Fe^{II}_2(Cl_2An^{2-})_3]^{2-}$  at  $T > T_{1/2c}$  via the intermediate state of  $[Fe^{II}Fe^{III}(Cl_2An^{2-})_2(Cl_2An^{\bullet 3-})]^{2-}$  at  $T_{1/2b} < T < T_{1/2c}$ (Scheme 1a).<sup>16</sup> In addition, compounds **0** and **0-d** are reversible via the solvation and desolvation processes, respectively. Hence, the systematic and continuous tuning of these phases were favorable due to being unprecedented in multiple TDET systems.

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Herein, we report a precise tuning of the multiple transition temperatures  $T_{1/2a}$ ,  $T_{1/2b}$ , and  $T_{1/2c}$  based on a chemical technique of site-doping for the charge-transfer materials **0** and **0-d**. A rational incorporation of the redox metal ions and bridging



**Scheme 1** Schematic representations of the investigated tetraoxolenebridged Fe honeycomb layers. (a) The variations of charge-ordered states in **0** and **0-d**, where for **0**,  $T_{1/2}(1)$  was  $T_{1/2a}$  (*i.e.*, there was no hightemperature phase for **0**), and for **0-d**,  $T_{1/2}(1)$  and  $T_{1/2}(2)$  correspond to  $T_{1/2b}$  and  $T_{1/2c}$ , respectively. (b) Schematic route for the syntheses of sitedoped solid solution compounds in this work.

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ligands into the charge-transfer skeleton is one of the efficient strategies in achieving the fine tuning of TDET.<sup>17,18,25-28</sup> In this work, the Cl<sub>2</sub>An<sup>*m*-</sup> unit in **0** and **0-d** was replaced by X<sub>2</sub>An<sup>*m*-</sup> with X = F or Br in several ratios in the isostructural series where F<sub>2</sub>An<sup>*m*-</sup> = 2,5-difluoro-3,6-dihydroxy-1,4-benzoquinonate and Br<sub>2</sub>An<sup>*m*-</sup> = 2,5-dibromo-3,6-dihydroxy-1,4-benzoquinonate. Because the F<sub>2</sub>An<sup>2-</sup> and Br<sub>2</sub>An<sup>2-</sup> units act as stronger and weaker electron acceptors than Cl<sub>2</sub>An<sup>2-</sup>, respectively, their solid-solution series revealed inverse trends in the doping rate dependence of  $T_{1/2}$ s between the F-doped and Br-doped series, that is, proportional and inversely-proportional, respectively, despite being linearly variable in both cases.

The solvated compound series of the solid solutions  $(NPr_4)_2[Fe_2(Cl_2An)_{(3-3Y)}(X_2An)_{3Y}]\cdot 2(acetone)\cdot H_2O$  (X-Y) as X = Br with Y = 0.10 (Br-0.10), 0.18 (Br-0.18), 0.35 (Br-0.35), 0.65 (Br-0.65) and X = F with Y = 0.06 (F-0.06), 0.12 (F-0.12), 0.17 (F-0.17), 0.24 (F-0.24), 0.31 (F-0.31), 0.42 (F-0.42) were synthesized using an identical procedure for 0. Specifically,  $H_2Br_2An$  or  $H_2F_2An$  was mixed into an  $H_2Cl_2An$  solution in various ratios (Scheme 1b),<sup>15</sup> where the Y value was obtained from the halogen elemental analyses for the X-Y and their desolvated series, X-Y-d (ESI†). The Y doping ratio was approximately identical to the mixing ratio in the synthetic procedures.

All compounds of the solvated form (X-Y) were isostructural to **0** in the Y-doped region of  $0.65 \ge Y$  and  $0.42 \ge Y$  for X = Br and F, respectively. The compounds crystalized in the monoclinic  $P2_1/c$  space group with an asymmetric unit containing two crystallographically independent  $Fe^{n+}$  ions, three  $Cl_2An^{m-}/$  $X_2An^{m-}$  ligands, and two NPr<sub>4</sub><sup>+</sup> cations (Fig. S1 and Table S1, S2, ESI<sup>†</sup>). The Fe ions and Cl<sub>2</sub>An/X<sub>2</sub>An ligands formed a honeycomb anionic layer  $[Fe_2(Cl_2An/X_2An)_3]^{2-}$  in the (100) plane. The occupancy ratio (*i.e.*, the X/Cl ratios) in the Cl<sub>2</sub>An/X<sub>2</sub>An ligands was fixed using the ratio obtained using the halogen elemental analyses (ESI<sup>†</sup>). The polycrystalline samples of the X-Y series were macroscopically evaluated using the powder X-ray diffractograms (PXRD) recorded at room temperature and compared with the simulated patterns determined using the single crystal XRD (SC-XRD) analyses (Fig. 1). We investigated the highly doped compounds beyond the doping rate of 0.65 < Y and 0.42 < Y for X = Br and F, respectively. The polycrystalline samples for the highly doped compounds and X<sub>2</sub>An<sup>m-</sup>-pure compounds had different PXRD patterns from that of 0. Unfortunately, the SC-XRD analyses for the highly doped samples did not succeed because of their low quality single crystals.

The desolvated series **X-Y-d** were obtained by using the samples of **X-Y**;  $(NPr_4)_2[Fe_2(Cl_2An)_{(3-3Y)}(X_2An)_{3Y}]$  (X = Br; **Br-0.10-d** for *Y* = 0.10, **Br-0.18-d** for *Y* = 0.18, **Br-0.35-d** for *Y* = 0.35, **Br-0.65-d** for *Y* = 0.65, X = F; **F-0.06-d** for *Y* = 0.06, **F-0.12-d** for *Y* = 0.12, **F-0.17-d** for *Y* = 0.17, **F-0.24-d** for *Y* = 0.24, **F-0.31-d** for *Y* = 0.42) (Scheme 1b, Fig. S2, ESI†). The SC-XRD analysis at 103 K revealed that the **X-Y-d** compounds successfully changed to an isostructural 2D honeycomb layer structure with **0-d** (Fig. S3, S4 and Tables S3, S4, ESI†). Both the **Br-Y-d** and **F-Y-d** were crystallized in the monoclinic space group,  $P2_1/n$ , with an asymmetric unit containing one

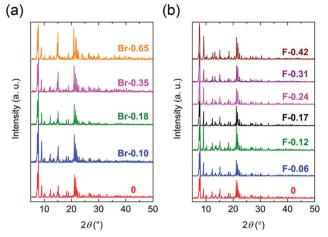
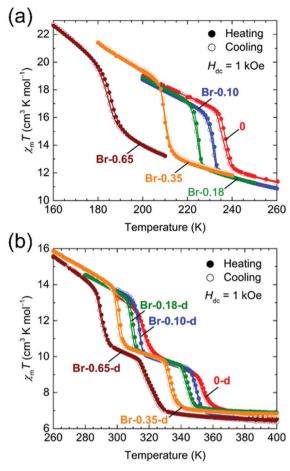


Fig. 1 PXRD patterns of  $Br\mathchar`-Y$  (a) and  $F\mathchar`-Y$  (b) measured at room temperature.

crystallographically unique  $\text{Fe}^{n+}$  (Fe  $\times$  1), one and a half  $\text{Cl}_2\text{An}^{m-}/\text{X}_2\text{An}^{m-}$  ligands, and one  $\text{NPr}_4^+$  cation (Fig. S3, ESI†)  $\text{NPr}_4^+$  was located between the honeycomb anionic layers  $[\text{Fe}_2(\text{Cl}_2\text{An}/\text{X}_2\text{An})_3]^{2-}$  formed by the Fe ions and  $\text{Cl}_2\text{An}^{m-}/\text{X}_2\text{An}^{m-}$  ligands, which were parallel to the (10–1) plane (Fig. S3, ESI†).

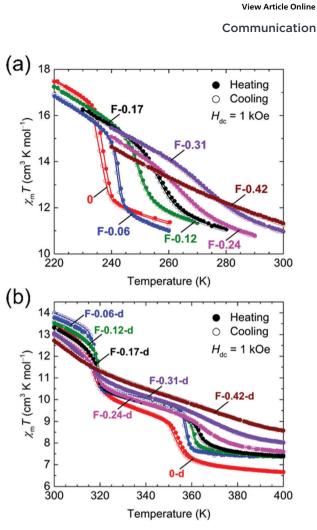
The oxidation states of the Fe<sup>*n*<sup>+</sup></sup> and X<sub>2</sub>An<sup>*m*-</sup> ligands for the series of **X-Y** and **X-Y-d** at 103 K were estimated from the local bond lengths of Fe–O and C–O in the Cl<sub>2</sub>An<sup>*m*-</sup>/X<sub>2</sub>An<sup>*m*-</sup> because their bond distances characteristically changed depending on their oxidation states.<sup>15,16,19,29</sup> At 103 K, the electronic states of all the site-doped compounds were estimated to be a low-temperature state, which was similarly observed in **0** and **0-d** as ([(Fe<sup>III</sup>)<sub>2</sub>(Cl<sub>2</sub>An<sup>2-</sup>)(Cl<sub>2</sub>An<sup>•3-</sup>)<sub>2</sub>]<sup>2-</sup>) (see the ESI<sup>†</sup>).

The temperature dependence of the magnetic susceptibilities ( $\chi = M/H$ ) of **Br-Y** (Y = 0.10, 0.18, 0.35, 0.65) and **F-Y** (Y = 0.06, 0.12, 0.17, 0.24, 0.31, 0.42) was measured by applying a 1 kOe dc field  $(H_{dc})$  in the temperature range of 300 to 1.8 K (Fig. 2a and 3a and Fig. S5 and S6, ESI<sup>†</sup>). The  $\chi_m T$  values at 300 K were in the range of 9.72-10.07 cm<sup>3</sup> K mol<sup>-1</sup> and 9.89-11.28 cm<sup>3</sup> K mol<sup>-1</sup> for the series of **Br-Y** and **F-Y**, respectively, which were varied in a narrow range (Table S5, ESI<sup>†</sup>). However, it is difficult to conclude that the variation of  $\gamma_m T$  values at 300 K was caused by the electronic state modulations in the series because of the presence of a strong antiferromagnetic spin coupling between the Fe<sup>III</sup> with S = 5/2 and  $X_2An^{\bullet 3-}$  with S = 1/2 (Fig. S5, S6 and Table S5, ESI<sup>†</sup>).<sup>15,16,29–32</sup> As the original **0** compound showed abrupt increases of  $\chi_m$  and  $\chi_m T$  at  $T_{1/2a}$  = 236 K upon cooling concomitant with the one-step transition from the intermediate phase of [Fe<sup>II</sup>Fe<sup>III</sup>(Cl<sub>2</sub>An<sup>2-</sup>)<sub>2</sub>(Cl<sub>2</sub>An<sup>•3-</sup>)]<sup>2-</sup> to a low temperature phase of  $[Fe^{III}_2(Cl_2An^{2-})(Cl_2An^{6-3-})_2]^{2-}$ , all doped X-Y compounds exhibited similar transition features in  $\chi_{\rm m}$  and  $\chi_{\rm m}T$ . However, their  $T_{1/2a}$  shifted to lower and higher temperatures with increasing doping rate Y for Br-Y and F-Y, respectively, compared with  $T_{1/2a}$  = 236 K for **0** (Fig. 2a and 3a); the  $T_{1/2a}$  values reached 184 K ( $\Delta$ 51 K from  $T_{1/2a}$  in **0**) and 270 K ( $\Delta$ 34 K) for **Br-0.65** and **F-0.31**, respectively. Specifically,  $T_{1/2a}$ s in both Br-Y and F-Y linearly varied whereas the doping rate of



**Fig. 2** Temperature dependence of  $\chi_m T$  for **Br-Y** (a) and **Br-Y-d** (b) measured at  $H_{dc} = 1 \text{ kOe} (\chi_m = M_m/H_{dc})$ , where the narrow temperature ranges of 160–260 K and 260–400 K, respectively, were used.

Y had inversely proportional and proportional trends on Y, respectively (Fig. 4 and Fig. S7, ESI<sup> $\dagger$ </sup>), where the  $T_{1/2a}$  temperatures were determined from the peak in the  $d\chi_m T/dT vs. T$  plots and their error bars were estimated from the half-value width of the peaks (Fig. S8, ESI<sup>†</sup>). The slopes of the  $T_{1/2a}$  variations were -77.3 and 110.3 K Y<sup>-1</sup> for **Br-Y** and **F-Y**, respectively. The **F-Y** series tended to have larger error bars with increasing Y than that of the Br-Y series, and  $T_{1/2a}$  was not assigned in F-0.42, which could be due to the large random domain effect of F-X. Nevertheless, these trends of  $T_{1/2a}$  in **Br-Y** and **F-Y** could be due to the site-doping of  $X_2An^{m-}$  with different electron acceptabilities, that is,  $F_2An^{2-} > Cl_2An^{2-} > Br_2An^{2-}$  (electronegativity of halogen atom: F > Cl > Br).<sup>33,34</sup> The magnetic profile in the entire temperature range from 300 to 1.8 K was similar to that for **0** excluding  $T_{1/2a}$ ; however, the magnetic correlation lengths at low temperatures were slightly different from each other. Br-Y had a tendency to have a longer correlation length in highly doped compounds, whereas it was approximately inversed for F-Y (Fig. S5 and S6, ESI<sup>†</sup>). Furthermore, the superparamagnetic behavior of single-chain magnets was observed at lower temperatures as shown in the original compound 0 (Fig. S9, ESI<sup>+</sup>), which can be attributed to the ferrimagnetic chain motif of  $[-Fe^{III}-(Cl_2An^{\bullet 3-}/X_2An^{\bullet 3-})-]$  (S<sub>Fe</sub> = 5/2, S<sub>rad</sub> = 1/2) separated by



**Fig. 3** Temperature dependence of  $\chi_m T$  for **F-Y** (a) and **F-Y-d** (b) measured at  $H_{dc} = 1 \text{ kOe} (\chi_m = M_m/H_{dc})$ , where narrow temperature ranges of 220–300 K and 300–400 K, respectively, were used.

the diamagnetic  $Cl_2An^{2-}/X_2An^{2-}$  bridges in the honeycomb layer (low-temperature phase in Scheme 1a).<sup>15</sup>

The temperature dependence of the magnetic susceptibility for the doped series of the desolvated compounds X-Y-d was measured in the temperature range of 400–1.8 K. As the  $\chi_m$ -and  $\chi_m T - T$  plots for **0-d** showed a two-step transition (abrupt increasing of  $\chi_m$  and  $\chi_m T$  upon cooling) at  $T_{1/2b}$  = 316 K and  $T_{1/2c}$  = 353 K, corresponding to the variation of the electronic state from  $[\text{Fe}_{2}^{\text{II}}(\text{Cl}_{2}\text{An}^{2-})_{3}]^{2-}$  at  $T > T_{1/2c}$  to  $[\text{Fe}^{\text{III}}_{2}(\text{Cl}_{2}\text{An}^{2-})(\text{Cl}_{2}\text{An}^{\bullet 3-})_{2}]^{2-}$  at  $T < T_{1/2b}$  via an intermediate state of  $[Fe^{II}Fe^{III}(Cl_2An^{2-})_2(Cl_2An^{\bullet 3-})]^{2-}$  at  $T_{1/2b} < T < T_{1/2c}$ (Scheme 1a),<sup>16</sup> the site-doped **X-Y-d** series showed similar twostep transitions. The error bar was larger in the highly doped species in F-X-d (Fig. 2b and 3b), owing to the characteristic domain effect. Moreover, the transition temperatures of  $T_{1/2b}$ and  $T_{1/2c}$  in **Br-Y-d** clearly shifted linearly to lower temperatures. On the other hand, the tendency in F-Y-d was slightly different from that of **F-Y**. The transition temperature of  $T_{1/2c}$ in F-Y-d tended to shift to higher temperatures in the solid solution species with lower doping rates; however,  $T_{1/2b}$  was approximately constant in the series. As mentioned above, the domain effect was characteristic in the series of F-Y-d and F-Y.

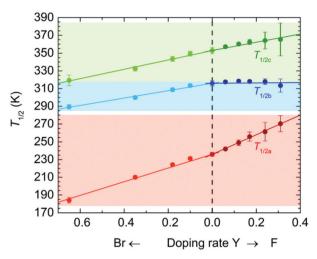


Fig. 4 Variations of the  $T_{1/2}$  value of TDETs for the site-doped rate Y in the solid solution series of X-Y and X-Y-d, whereas the basis of Y = 0 corresponded to 0 and 0-d. The left and right scales represent **Br-Y/Br-Y-d** and **F-Y/F-Y-d**, respectively. The solid lines are least-square linear fitting lines for the respective series.

Hence, it is difficult to conclude that  $T_{1/2b}$  was not affected by the site-doping of  $F_2An^{m-}$ . Although there was no clear reason, the transition temperatures in **F-Y-d** were slightly affected by the doping of the site-doping of  $F_2An^{m-}$ . The slope of the  $T_{1/2b}$  and  $T_{1/2c}$  variations *versus Y* was -41.7 and -52.7 K Y<sup>-1</sup> for **Br-Y-d** and 2.0 and 46.5 K Y<sup>-1</sup> and **F-Y-d** (Fig. S7, ESI†). The difference in tendencies can be explained by the electronacceptability of  $F_2An^{2-} > Cl_2An^{2-} > Br_2An^{2-}$ .

The site-doping of  $F_2An^{m-}$  influenced the magnetic phase of the solid-solution materials of **F-Y-d**. The **F-Y-d** species with  $Y \le 0.24$  have a maximum of  $\chi_m$  at approximately 23 K followed by a steep decrease after further cooling. This was indicative of an antiferromagnetic ground state as similarly found in **0-d** and **Br-Y-d** (Fig. S5, ESI†); however, the decrease of  $\chi_m$  did not occur in the cases with higher doping rate. Instead, the superparamagnetic behavior of single-chain magnets was observed in the low-temperature phase of **F-Y-d** with 0.24 < *Y* (Scheme 1a).

In summary, we succeeded in the precise tuning of multiple TDET at three  $T_{1/2}$ s transition temperatures through the chemical-site doping of two types of tetraoxolene-bridged Febased honeycomb layer compounds. This provided an "ondemand TDET" in a temperature range of 184 to 365 K.

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## Conflicts of interest

There are no conflicts to declare.

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