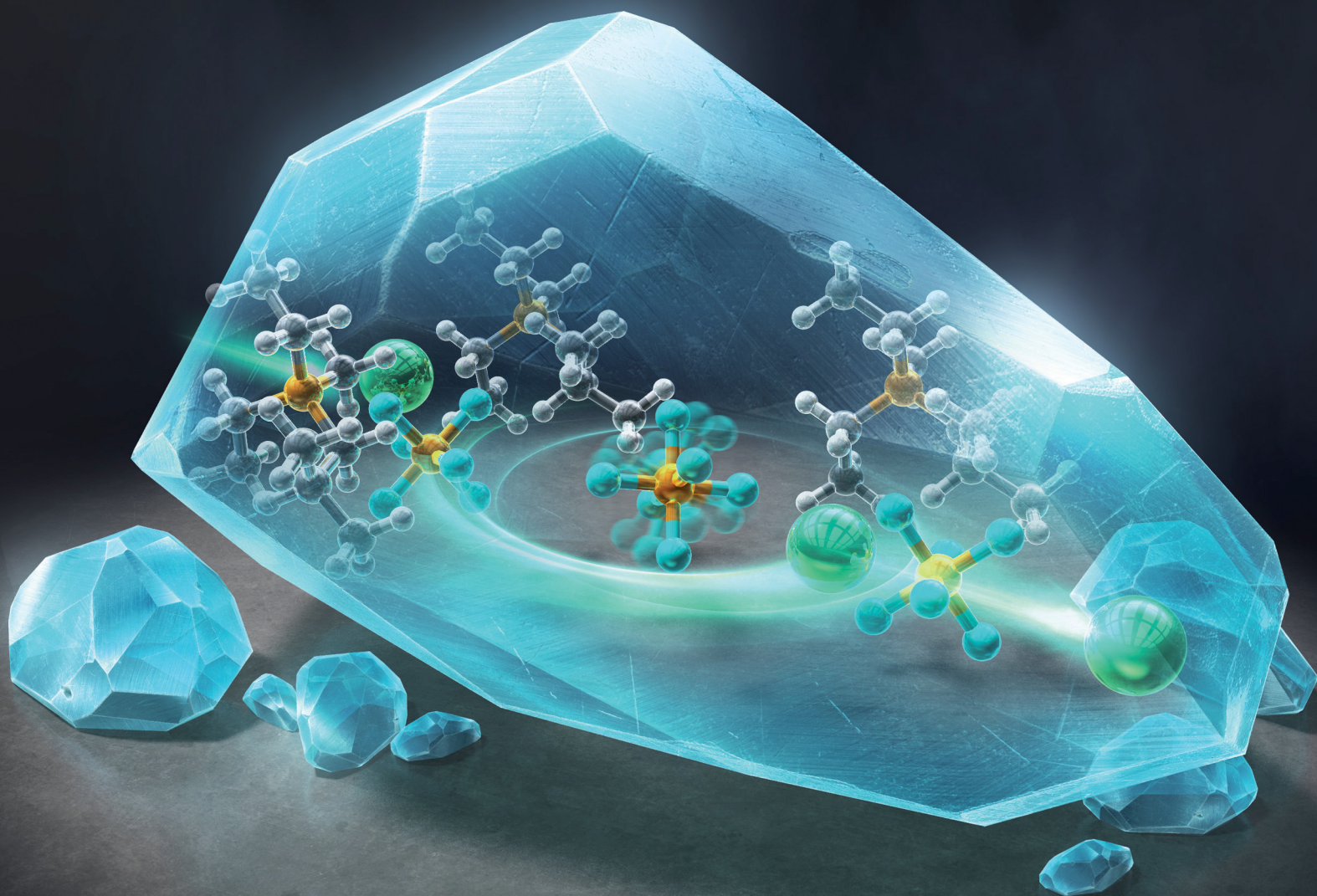


Energy Advances

Volume 2
Number 6
June 2023
Pages 741-888

rsc.li/energy-advances



ISSN 2753-1457

REVIEW ARTICLE

Masahiro Yoshizawa-Fujita *et al.*
Organic ionic plastic crystals: flexible solid electrolytes for
lithium secondary batteries



Cite this: *Energy Adv.*, 2023,
2, 748

Received 17th February 2023,
Accepted 24th April 2023

DOI: 10.1039/d3ya00078h

rsc.li/energy-advances

Organic ionic plastic crystals: flexible solid electrolytes for lithium secondary batteries

Morgan L. Thomas,^a Kan Hatakeyama-Sato,^b Shinkoh Nanbu^a and
Masahiro Yoshizawa-Fujita^{*a}

The growing global demand for energy has led to the active development of efficient energy generation and storage technologies, driving the development of electrochemical devices such as high-energy density rechargeable batteries, fuel cells and solar cells. One of the essential materials for the development of high-performance electrochemical devices is the electrolyte. Currently, flammable electrolyte solutions are used, causing problems such as leakage and ignition incidents. It would be significant if the electrolyte could be replaced with a solid electrolyte, as this would eliminate these problems. In addition, with the increasing size of electrochemical devices, there is a societal demand for safer electrochemical devices, and the development of high-performance solid electrolytes is becoming more active. Although development has mainly focused on inorganic and solid polymer electrolytes, organic ionic plastic crystals (OIPCs) are beginning to attract attention as new candidates for flexible solid electrolytes. In this review, we describe OIPCs for lithium secondary batteries. Firstly, we introduce OIPCs and OIPC/polymer composites as lithium-ion conductors and discuss the effects of ionic architecture and polymer species on their ionic conduction. Secondly, we present recent progress in the development of lithium secondary batteries with OIPC-based solid electrolytes.

1. Introduction

Growing interest in energy issues has increased demand for large batteries that can be used in electric vehicles and home power sources. Lithium-ion batteries (LIBs) have become widely used as compact power sources for mobile phones and other devices, but further improvements in safety and energy density

^a Department of Materials and Life Sciences, Sophia University, 7-1 Kioi-cho, Chiyoda-ku, Tokyo 102-8554, Japan. E-mail: masahi-f@sophia.ac.jp

^b Department of Applied Chemistry, Waseda University, 3-4-1 Okubo, Shinjuku-ku, Tokyo 169-8555, Japan

† Current address: Materials Science and Engineering, School of Materials and Chemical Technology, Tokyo Institute of Technology, 2-12-1-S8-36 Ookayama, Meguro-ku, Tokyo 152-8552, Japan.



Morgan L. Thomas

professor in 2019. His research interests are clean technologies, with a particular interest in solvent and electrolyte systems, and green chemistry.

Morgan L. Thomas undertook undergraduate studies at the University of Bath (U.K.), and he received his PhD in 2008 from the University of Nottingham (U.K.). This was followed with post-doctoral experiences at RWTH Aachen (Germany), the University of Saskatchewan, York University and McGill University (Canada), RIKEN and Yokohama National University (Japan). He joined Sophia University as a specially appointed associate



Kan Hatakeyama-Sato

Kan Hatakeyama-Sato obtained his PhD in Chemistry from Waseda University in 2018. His PhD work was conducted on redox polymers and their electrochemistry. From 2018 to 2023, he worked in the same University as an assistant professor, and began the data science study for chemistry.



are needed to meet these demands.¹ However, conventional LIBs have generally reached their performance limits and there is a need to develop new energy storage devices that can be scaled up. Against this background, all-solid-state batteries, in which the cathode, electrolyte, and anode are all made of solid materials, have attracted much attention in recent years.^{2–4}

The electrolyte is responsible for conducting the carrier ions between the positive and negative electrodes, while electrically insulating the positive electrode to prevent short circuits. Solid electrolytes are more flame retardant than the organic solvent-based electrolytes used in conventional LIBs, so all-solid-state batteries are expected to be very safe. In addition, high energy densities can be achieved by manufacturing stacked cells; in the case of conventional LIBs, several small batteries are connected in series to achieve high voltages, whereas in all-solid-state batteries, high voltages can be easily achieved by stacking the cathode, electrolyte, anode and current collector. In addition, because the solid electrolyte is not liquid, it can be used as a material in different components of a single cell, *i.e.* for the positive electrode, negative electrode and separator, providing a high degree of freedom in battery design. There is also the possibility of using high-capacity electrode active materials such as metallic lithium and sulphur,^{5–8} which have been difficult to use in conventional LIBs, and expectations are high for the realisation of next-generation batteries.

There are two main types of all-solid-state batteries: thin film and bulk. Thin film all-solid-state batteries are made by growing crystals of the cathode, electrolyte and anode on a substrate using the vapour phase method. The advantage of such thin-film batteries is that achieving a good interfacial contact between the electrode and the electrolyte is facile.^{9,10} However, the amount of active electrode material that can be introduced into a single cell is small, resulting in low battery capacity. Bulk all-solid-state batteries are made by layering fine particles of cathode, electrolyte and anode. The advantage is that a large amount of electrode active material can

be introduced, enabling the production of high capacity batteries.¹¹ However, high cell resistance is caused by the thick electrolyte layer and many particle interfaces. Therefore, the development of solid electrolytes with excellent ionic conductivity and the formation of good particle interfaces are important issues for the practical application of bulk all-solid-state batteries.

Plastic crystals (PCs) are defined as solid materials that are composed of regularly aligned three-dimensional crystal lattices, but at the molecular level, there is a local orientational/rotational disorder. Since the molecular and ionic species in the plastic crystal phase have a higher degree of freedom than in the normal crystal state, PCs are crystalline solids exhibiting plasticity and diffusivity of constituents. The features of PCs found in early studies are summarized below.

- (1) One or more solid–solid phase transitions before melting.
- (2) Small change in entropy of fusion (ΔS_f).
- (3) High plasticity.

Timmermans reported that PCs exhibit a ΔS_f value below 5 e.u. ($20.92 \text{ J K}^{-1} \text{ mol}^{-1}$)[‡] due to the initiation of rotational or translational motion of the constituents in the solid phases.¹² PCs are classified into non-ionic compounds and ionic compounds. Succinonitrile is known as a non-ionic PC, and inorganic crystals such as Li_2SO_4 and $\alpha\text{-Na}_3\text{PO}_4$ and organic crystals such as ammonium salts are known as ionic PCs. Recently, both systems are attracting much attention as soft solid electrolytes.^{13–18} In this paper, organic ionic compounds showing a PC phase will be referred to as organic ionic plastic crystals (OIPCs).

‡ The entropy of the fusion for plastic crystals is known to be less than $20 \text{ J K}^{-1} \text{ mol}^{-1}$. However, Timmermans proposed 5 e.u., which is equivalent to $20.92 \text{ J K}^{-1} \text{ mol}^{-1}$. In this review, we have adopted $20.92 \text{ J K}^{-1} \text{ mol}^{-1}$ as a more accurate value.



Shinkoh Nanbu

Shinkoh Nanbu graduated from the Department of Chemistry, Faculty of Science and Technology, Keio University with a Japan Society for the Promotion of Science (JSPS) fellowship (DC). Working at the Institute for Molecular Science Computer Center as research-associate, conducted development of a theory for quantum chemical reaction dynamics at Argonne National Laboratory in the state of Illinois, US. Worked at Kyushu University before assuming

his current position in 2009. His research themes revolve around nonadiabatic phenomena, elementary processes, isotopologues, and photochemical reactions in the condensed phase.



Masahiro Yoshizawa-Fujita

Masahiro Yoshizawa-Fujita received his PhD degree (2002) from Tokyo University of Agriculture and Technology (Japan). During his PhD studies, he received a Research Fellowship for Young Scientists at the Japan Society for the Promotion of Science (JSPS). He spent two years as a postdoctoral research fellow at Monash University (Australia). He moved to Sophia University as an Assistant Professor in 2006. He was promoted to

Associate Professor in 2011 and to Professor in 2019. His recent research activities are concerned with the design of ionic liquids, ionic plastic crystals, and solid polymer electrolytes, especially for battery technologies and biomass processing.



2. Lithium-ion conductors

2.1 Pyrrolidinium cations

The figure displays 34 chemical structures of organosilicon compounds, categorized into 24 cationic and 10 anionic species. The structures are numbered 1 through 24, with the last row containing structures labeled a through j.

- 1-10:** Cationic organosilicon compounds with various ring systems (cyclopentadienyl, cycloheptatrienyl, cyclooctatetraenyl, cyclononatrienyl, cyclodecaheptatrienyl, and a bicyclic system).
- 11-17:** Cationic organosilicon compounds with linear and branched alkyl chains, including a nitrile group in 14.
- 18-24:** Cationic organosilicon compounds with various ring systems and substituents, including a dimethylamino group in 24.
- a-j:** Anionic organosilicon compounds, including sulfonates (a, g), nitriles (h), and various silanones and silanediols (i, j).

Fig. 1 Chemical structures of OIPCs reported as lithium-ion conductors.

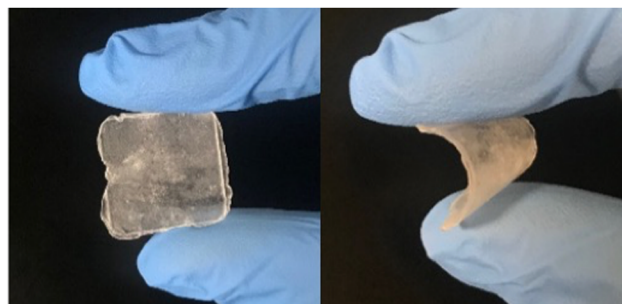


Fig. 2 Photographs of **1b** at room temperature.

© 2023 The Author(s). Published by the Royal Society of Chemistry

into a self-standing film, and the obtained film can be deformed freely.

The ionic conductivity in Phase I of **1b** increased linearly with increasing temperature, showing an Arrhenius-type temperature dependence (Fig. 3). It was suggested that the ions exhibit hopping conduction in the solid. The ionic conductivity in Phase II was as low as 10^{-8} S cm $^{-1}$ or less but increased by about an order of magnitude in T_{I-II} . It is considered that the rotational and translational motions of the constituent ions started in Phase I, leading to the high diffusivity of these ions. The ionic conductivity of **1b** was 1.23×10^{-6} S cm $^{-1}$ at 25 °C, which is about two orders of magnitude higher than that of **1a**, which has a similar anion structure. The physicochemical properties changed significantly due to the slight difference in the anion structure. The influence of the cation structure on various properties was also investigated. **2b**, **3b**, and **4b** with varying side chain lengths of pyrrolidinium cation **1** were synthesized.^{43–46} All compounds showed multiple solid–solid phase transitions, and ΔS_f was 20.92 J K $^{-1}$ mol $^{-1}$ or less, suggesting that these salts are OIPCs. The ionic conductivity of **2b**, **3b**, and **4b** in Phase I also exhibited Arrhenius type behavior. The ionic conductivity values of **2b**, **3b**, and **4b** were 1.3×10^{-8} (25 °C), 1.3×10^{-5} (25 °C), and 1.7×10^{-6} S cm $^{-1}$ (30 °C), respectively. Thus, higher ionic conductivity was observed in the order **2b** < **1b** < **4b** < **3b**, and the ionic conductivity of pyrrolidinium salts was improved with the extension of the alkyl side chain. When 5 mol% of lithium bis(fluorosulfonyl)amide (LiFSA) was added to these pyrrolidinium salts, the ionic conductivity values of **2b**/LiFSA, **1b**/LiFSA, and **3b**/LiFSA were 2.8×10^{-6} , 1.1×10^{-5} , and 5.8×10^{-4} S cm $^{-1}$ at 25 °C, respectively. In each system, the ionic conductivity was improved by about an order of magnitude as compared with the pristine salt (Fig. 3). The ionic conductivity of **4b**/LiFSA (10 mol%)

was 4.0×10^{-4} S cm $^{-1}$ at 30 °C. As with other pyrrolidinium salts, the addition of Li salt improved the ionic conductivity.

The lithium-ion transport number (t_{Li^+}) of **3b**/LiFSA (5 mol%) was 0.27.⁴³ The t_{Li^+} values of high-concentration electrolytes containing 50 mol% of LiFSA added to **3b** and **4b** have been reported. The t_{Li^+} of **3b**/LiFSA and **4b**/LiFSA were 0.39⁴⁴ and 0.46,⁴⁵ respectively, and **4b** showed higher values. It was suggested that the isopropyl group, which is a branched chain, is a functional group that promotes lithium-ion transport. Further improvement of t_{Li^+} can be expected by optimizing the side-chain structure of the cation. The t_{Li^+} of the composite of *N*-methyl-*N*-propylpyrrolidinium bis(fluorosulfonyl)amide and LiFSA, which has a chemical structure similar to that of **3b** and **4b**, is only 0.14,⁴⁷ lower than those of **3b**/LiFSA and **4b**/LiFSA. It was revealed that OIPCs, which are solids, are superior to ionic liquids in lithium-ion transport. It has been reported that a system in which the lithium salt itself becomes an OIPC is effective for further improvement of t_{Li^+} .⁴⁸ The system in which the lithium salt itself is an OIPC has a high lithium-ion concentration because half of the component ions are lithium ions. On the other hand, it is difficult to achieve a high concentration of lithium salt in OIPC/lithium-salt composites because the melting point lowers and it may become liquid at room temperature when over 10 mol% of lithium salt is added. The t_{Li^+} of the Li salt using **d** as an anion was 0.7, which is a high value that is difficult to realize not only in OIPC/Li salt composite s but also in liquid electrolyte solutions. The OIPC characteristics of certain Li salts themselves is an interesting theme.

2.2 Other nitrogen-based cations

There are many reports on nitrogen-based cations other than pyrrolidinium salts. Abu-Lebdeh *et al.* evaluated the phase transition behavior and ionic conductivity of spiroammonium salts **5a–9a**.⁴⁹ Of the five spiroammonium salts, **5a**, **6a**, and **7a** showed a plastic crystal phase. Furthermore, among the three OIPCs, **6a** showed the highest ionic conductivity, which was 2.0×10^{-6} S cm $^{-1}$ at 20 °C. The ionic conductivity of the **6a** composite with 5 mol% of lithium bis(trifluoromethylsulfonyl)amide (LiTfSA) was 1.2×10^{-6} S cm $^{-1}$ at 20 °C, and thus the ionic conductivity in the solid phase was about the same before and after the addition of the Li salt. In general, the ionic conductivity of OIPC is improved by the addition of Li salt. These spiroammonium cations are alicyclic and are considered to have a limited degree of freedom compared to pyrrolidinium cations with alkyl side-chains. The side-chain structure of pyrrolidinium cations was found to greatly affect ionic conductivity, which is consistent with the findings already mentioned.

In addition to the spiroammonium salts, **10a**, which is an alicyclic OIPC using 1,4-diazabicyclo[2.2.2]octane (DABCO), has been reported.^{50,51} It differs from the pyrrolidinium and spiroammonium salts in that a diamine is used as a starting material. When the diamine is monocationized, an ammonium salt containing a tertiary amine is obtained, and it can be expected that the Lewis basic tertiary amine moiety interacts with the Li-ion to promote the dissociation of the Li salt.

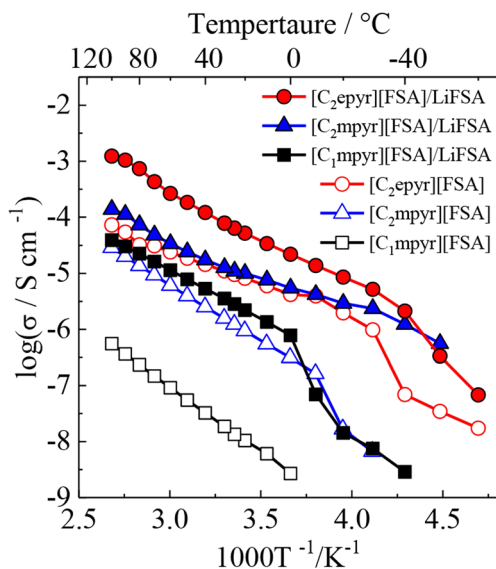


Fig. 3 Arrhenius plots of ionic conductivities for neat pyrrolidinium salts (**1b** ([C₂mpyr][FSA]), **2b** ([C₁mpyr][FSA]), **3b** ([C₂epyr][FSA])) and 5 mol% LiFSA-doped ones. Reprinted from ref. 43 with permission.



The transition temperature and melting point of **10a** were 33 and 76 °C, respectively, and the ΔS_f was 25.3 J K⁻¹ mol⁻¹. This is slightly larger than the Timmermans criterion (<20.92 J K⁻¹ mol⁻¹). It is known that ΔS_f of OIPCs having TFSA anions shows a value larger than the Timmermans criterion because of the residual entropy based on the conformations of the TFSA anion.¹³ The temperature dependence of the ionic conductivity of **10a** showed typical OIPC behavior. The ionic conductivity increased at the phase transition temperature from Phase II to I. When LiTFSA was added to **10a**, the ionic conductivity in Phase I improved by up to an order of magnitude compared to the value before the addition of LiTFSA. As a result of ⁷Li NMR measurement, a low magnetic field shift based on the interaction between Li-ion and **10** tertiary amines was observed. It is considered that the promotion of dissociation of Li salt contributes to the improvement of ionic conductivity, and it is expected that it can be utilized as one of the new molecular design guidelines of OIPCs.

Although the number of reports on aliphatic ammonium salts is small, there are many structural variations (**11–16**).^{52–54} Yunis *et al.* reported 6 types of ammonium salts by combining **11**, **12**, **13**, and **a** and **b**, and comprehensively discussed their ion conduction behavior.⁵² The results of DSC measurements suggest that 5 of the ammonium salts (all except **11a**) are OIPCs. The temperature range of Phase I of **12b** is –8 to 133 °C, and among the 6 ammonium salts, **12b** showed the highest ionic conductivity, which is 9.9×10^{-6} S cm⁻¹ at 30 °C. It is already known that the ionic conductivity in OIPCs correlates with the defect volume.⁵⁵ The correlation between the free volume and the ionic conductivity obtained by the positron annihilation method was investigated. In the TFSA system, the larger the free volume of OIPC, the higher the ionic conductivity, which is consistent with conventional knowledge. On the other hand, in the FSA system, no correlation was found between the free volume and the ionic conductivity. In the case of the FSA system, it is therefore suggested that the free volume is not a regulator of ionic conductivity, which is an interesting result. Recently, it has been reported that there is a correlation between the ionic radius ratio (ρ) and the ionic conductivity (Fig. 4),^{56,57} and the understanding of the molecular design guidelines for high ionic conductivity OIPCs is deepening. Further elucidation of the controlling factors of ionic conductivity is awaited.

Additionally, ammonium salts comprised of a combination of **14** having a cyano group introduced in the side chain and various anions (**b**, **e**, **f**, **g**, **h**),⁵⁸ and **17a**, which has organosilicon introduced into the side chain,⁵⁹ have been reported. Since the cyano group forms a protective film on the surface of metallic lithium, it is known to improve the electrochemical stability of the electrolyte and the redox cycle characteristics of lithium,^{60,61} and similar effects can be expected with such OIPCs. Four of the **14** salts (all except **14h**) showed solid–solid phase transitions in DSC measurements. The ionic conductivity in Phase I (30 °C) showed values in the following order (only **14g** was compared with Phase II values): **14d** (6.6×10^{-8} S cm⁻¹), **14b** (2.1×10^{-8} S cm⁻¹), **14g** (6.7×10^{-9} S cm⁻¹), **14f** (7.2×10^{-10} S cm⁻¹). The ionic

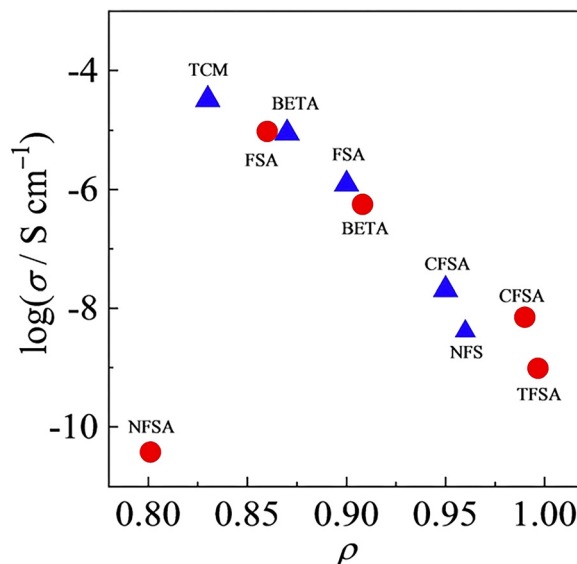


Fig. 4 Relationship between ionic conductivity value at 25 °C and ρ for **1**- (blue triangle) and **3**-based salts (red circle) with various anion species. Reprinted from ref. 57 with permission.

conductivity of the systems in which 10 mol% Li salt was added to **14b** and **14c** showed different behavior. The ionic conductivity of **14b**/LiFSA improved by more than 3 orders of magnitude compared to the value of **14b**. On the other hand, the ionic conductivity of **14c**/LiBF₄ was less than an order of magnitude greater than that of **14c**. When the Li salt concentration was increased to 50 mol%, **14b** became a liquid at room temperature. This electrolyte showed a high t_{Li^+} , 0.56 at 50 °C. Since the redox behavior of lithium was also stable, it is presumed that various properties were improved by the effect of the cyano group. As described above, the introduction of a functional group into the cationic structure leads to the addition of other properties, which is one of the advantages of organic-based OIPCs. There are many variations of nitrogen cations that can be employed to form the skeleton of the cation; pyrazolium cations (**18–23**)^{62–64} and guanidinium cations (**24**)⁶⁵ have also been reported. Volel *et al.* evaluated the mechanical properties of pyrazolium salts **18a–22a** using atomic force microscopy (AFM).⁶⁶ The Young's modulus of these pyrazolium salts was 100 MPa or less in the solid phase below the melting point and thus showed a value equivalent to that of elastomers (10 to 100 MPa). The Young's modulus changed before and after the phase transition and decreased by about 10 to 30 MPa in the solid phase on the high-temperature side. Few reports quantify the softness of OIPCs, and it is difficult to make a comprehensive comparison. In the future, it will be necessary to accumulate more data on mechanical properties.

2.3 Phosphonium cations

Although the number of variants of the cation structure is inferior to those of the nitrogen cation systems, the development of lithium-ion conductors using phosphonium salts has also been vigorously performed. Armel *et al.* synthesized



15 phosphonium salts by combining 25–28 cations with **a**, **b**, **e**, **f**, **h**, and **i** anions.⁶⁷ Among these, 12 of the phosphonium salts showed a plastic crystal phase. The cations 26 to 28 have an isobutyl group in the side chain, suggesting that the branched structure is advantageous for expression of the plastic crystal phase. However, the melting points of 25**b**, 27**b**, and 28**b**, which exhibit a transition between solid phases, are 47, 4, and 36 °C, respectively, which are lower than those of the nitrogen-based cation system. The introduction of phosphonium cations and isobutyl groups appears to lead to lower melting points for the OIPCs. The ionic conductivity of these phosphonium salts was also evaluated. Interestingly, the ionic conductivity of 26**i** at 40 °C (phase I) was comparable to that of 28**b** in the molten state ($1.75 \times 10^{-3} \text{ S cm}^{-1}$). The alkyl chain having a branched structure seems to be a functional group exhibiting a specific ionic conduction behavior and will be one of the structures to be utilized widely in the future.

Detailed analysis of the ion-conductive mechanism in a phosphonium salt was performed using solid-state NMR.³³ The solid–solid phase transition temperatures of 26**f** are 25 ($T_{\text{III–IV}}$), 70 ($T_{\text{II–III}}$), 120 °C ($T_{\text{I–II}}$), and the melting point is 150 °C. The ionic conductivity in each solid phase is $10^{-10} \text{ S cm}^{-1}$ (phase IV), 10^{-8} to $10^{-6} \text{ S cm}^{-1}$ (phase III), 10^{-5} to $10^{-4} \text{ S cm}^{-1}$ (phase II), and $10^{-3} \text{ S cm}^{-1}$ (phase I). The ionic conductivity jumped at each solid–solid phase transition temperature. Based on the analysis results by solid-state NMR, isotropic rotation of 26 alkyl groups and **e** occurs in phase IV, and lattice defects occur due to uniaxial rotation of 26 in phase III, and thus the diffusion of **e** occurs. It was clarified that in phase II, isotropic rotations of 26 occurred, and the main diffusing species was **e**, and in phase I, both **e** and 26 diffused. There are reports using magnetic resonance imaging (MRI),^{68,69} and there is no doubt that NMR is an effective tool for analyzing the mechanism of ion conduction.

3. Ion conduction in OIPCs

Understanding the ion conduction mechanisms in OIPCs is a very important research topic, as it is essential for further improving their ionic conductivity as a practical requirement for rechargeable batteries. Two promising techniques (NMR techniques and MD simulations) for elucidating the ion conduction mechanisms are presented here.

3.1 NMR techniques

The application of nuclear magnetic resonance (NMR) techniques to provide crucial insights into the thermochemical phenomena, dynamics and compositional variations of OIPCs for lithium batteries has been well-developed. Experimental progress in this field has built on a great deal of earlier work with OIPCs with simple anions, for example ammonium halides.^{70–72} The analysis of OIPCs relevant to Li battery application, and in many cases with the addition of alkali metal salts, by variable-temperature solid state NMR has been undertaken by a number of different approaches,^{73,74} including (i)

chemical shift analysis, (ii) linewidth/lineshape analysis, (iii) self-diffusion coefficient measurements, and (iv) imaging.

Regarding the determination of chemical shifts and their variation upon *e.g.* addition of alkali metal salts, various authors have employed such analyses to further develop understanding of interionic interactions in OIPCs. We note that whilst the primary focus of this review is on Li^+ -containing OIPCs, some specific finding for Na^+ are also relevant. For example, Forsyth and co-workers described the existence of a second, Na^+ -rich, phase in 1**a**/NaTFSA, based in part on the observation of splitting of the ^{19}F and ^{13}C NMR signals, measured under magic angle spinning (MAS) conditions, into multiple peaks at high Na^+ content.⁷⁵

The exploration of linewidth/lineshape for elucidation and understanding of molecular motion was explored early in the development of OIPC-based Li^+ electrolytes. For analysis of Li mobility, Forsyth *et al.* observed that for 1**a**/LiTFSA, the presence of a sharp ^7Li signal overlaying a broader peak implies the presence of both higher and lower mobility Li^+ ions (Fig. 5).²⁹ However, MacFarlane *et al.* also note that some caution is necessary, as the assumed mobility leading to line narrowing does not necessarily imply that Li itself is undergoing increased mobility or related processes; the narrowing of the peak may also be due to the influence of neighbouring nuclei/species with increasing mobility.²⁸ Nevertheless, such measurements are facile, and continue to be successfully applied for various OIPC/Li salt systems. Pringle and co-workers noted an apparently higher mobility for the 3 cation in 3**b** in the presence of LiFSA based on a narrowing of the ^1H signal,⁴⁴ and more recently applied ^7Li lineshape analysis for understanding mobilities in 1**e**/LiFSA, and its comparison with similar zwitterionic/Li salt materials.⁷⁶

Beyond qualitative consideration of linewidths, consideration of chemical shift anisotropy (CSA) can be used to derive important details about mobility. Pringle *et al.* proposed the

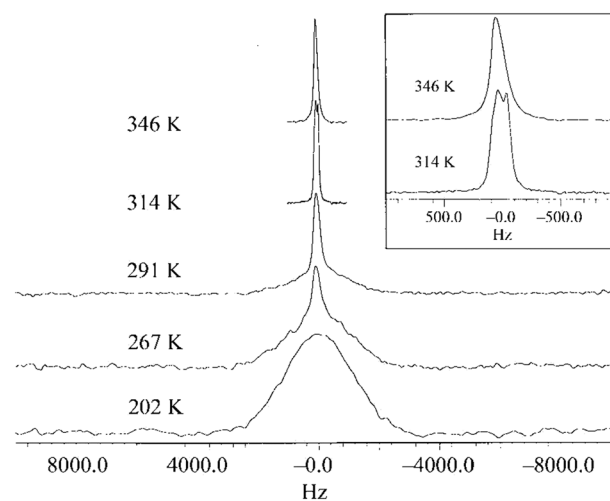


Fig. 5 ^7Li NMR spectra for 9.3 mol% LiTFSA in 1**a**, as a function of temperature. The inset is an expansion of the peaks for 314 K and 346 K. Reprinted from ref. 29 with permission.



axial rotation of the cation upon heating through consideration of CSA (specifically the shape of the peak) for ^{31}P in neat **25b**. Furthermore, the authors estimated the CSA symmetry parameter for ^{19}F at various temperatures by modelling (3-site jump process) and comparison with the experimental data, providing key insights into the anion mobility, and enabling the estimation of activation energy.⁷⁷ A further quantitative approach to linewidth analysis has also been employed, allowing the temperature-dependent correlation time (inverse of jump/hop frequency) and the corresponding activation energy to be determined from the linewidths by application of the Bloembergen, Purcell, and Pound model.⁷⁸ Quantitative analyses such as these allow for more detailed examination of the ion mobility phenomena and comparison between similar systems.

Self-diffusion coefficients, measured using pulsed field gradient (PFG) NMR, can provide important details on the relative mobility of components. For example, Pringle *et al.* reported that the self-diffusion coefficient for Li^+ was the highest of all components in **3b** with 90 mol% LiFSA, resulting in a relatively high transference number of 0.63 (in good agreement with the electrochemically determined value of 0.68).⁴⁴ Similar results for 90 mol% LiFSA in a zwitterion (*i.e.* Li cation showing the highest mobility) were also reported more recently.⁷⁶

Within the plethora of imaging strategies available to NMR users, many powerful imaging techniques are based on evaluation of the transverse relaxation time constant, T_2^* ,⁷⁹ which is a measure of the decay in transverse magnetization resulting from both spin–spin relaxation and magnetic field inhomogeneities. Specifically, Romanenko and co-workers demonstrated the utilisation of pulse sequences employing single point ramped imaging with T_1 enhancement (SPRITE) to probe anisotropy and develop understanding of the orientation of

grain boundaries in polycrystalline OIPCs such as **25b** and **26f**.⁸⁰ Indeed, their step-by-step progress in these measurements identified thermal-history derived variations in crystallite alignment,^{68,81} morphological changes upon solid–solid phase transitions (variable-temperature),⁸² and ultimately culminated in a landmark *in operando* study on a lithium cell with **28b**/LiFSA electrolyte.⁶⁹ Fig. 6 shows the variation in Li^+ ion distribution in a symmetric Li–Li cell during galvanostatic stripping and plating. The authors comment on the role of the ion conduction mechanism, grain boundaries, and presence of Li depleted domains (essentially neat OIPC) in the bulk in relation to the possible applications of OIPCs for Li batteries. Although technically challenging as compared to more routine chemical shift, lineshape and conventional diffusion measurements, further development of this (and other⁸³) imaging techniques is expected to contribute to further understanding of the importance of grain boundaries, morphology and heterogeneity in OIPC/Li salt electrolyte based devices.

3.2 Molecular dynamics and time scales

Probably the first MD (molecular dynamics) study focused on the importance of the rotational motion of anions and cations forming crystals as a dynamic behavior contributing to ionic conductivity and ion transport in OIPCs.⁸⁴ In this paper, the OIPC **11h** was investigated. A unit cell was constructed based on the crystal structure determined by X-ray structural analysis, and the molecular force field and charge distribution given by the CHARMM MD simulation package were used to perform the MD simulation with the DL_POLY simulation package; the *NPT* ensemble simulation was carried out from $T = 200$ to 600 K. The results successfully suggest that the rotational motion of both the anion and the cation starts at a slightly lower temperature for the cation than for the anion, while

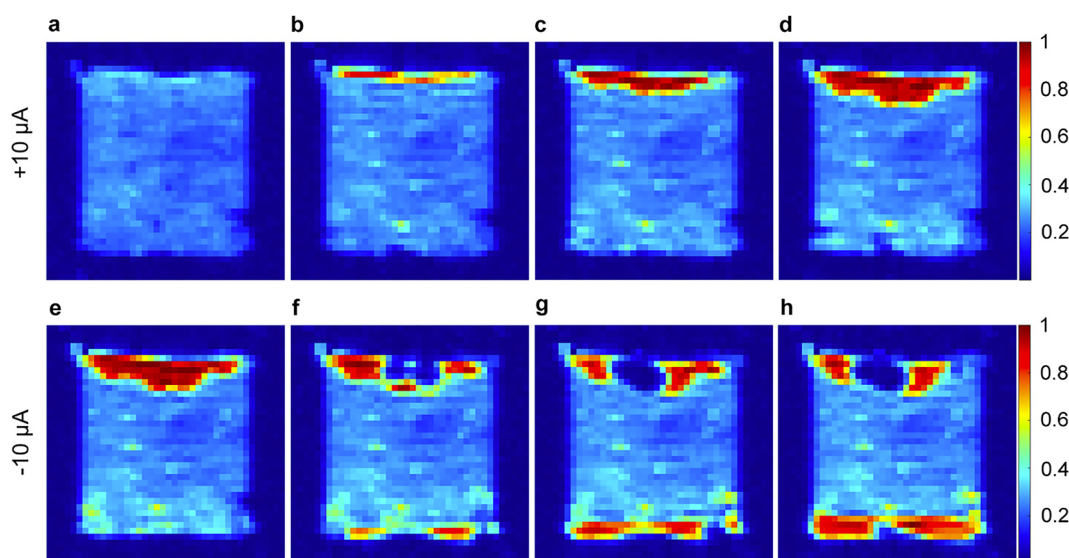


Fig. 6 ^1H SPRITE images of a symmetric cell (Li coated Cu disks) with 4 mol % LiFSA-doped **28b** electrolyte under galvanostatic conditions. (a) Initial cell then continuous application of $+10\ \mu\text{A}$ for (b) 17 h, (c) 34 h, and (d) 51 h; and reversed polarity ($-10\ \mu\text{A}$) for the same cell for (e) 6 h, (f) 40 h, (g) 57 h, and (h) 77 h. Reprinted from ref. 69 with permission.



maintaining the crystal lattice near $T = 0$ K, which is a perfect crystal, and could correspond to the temperature of phase transitions observed experimentally in similar systems. As the unit cell size expands with the phase transition, the position of the peak of the radial distribution function (RDF) is found to shift towards the long-range side. However, the theoretical result at that time could not be directly compared with the experimental results by any means and was verified by comparison with similar OIPCs. However, Jin *et al.* provided the experimental evidence for the rotational motion with fixing to the lattice site.³³ The OIPC **26f** was investigated by the analysis of several experimental data (single crystal XRD, multinuclear solid-state NMR, DSC, ionic conductivity, and SEM) and theoretical simulations based on the second moment-based solid-state NMR line-width simulation; the peak in the NMR spectrum suggests molecular rotation, which is a small sharp peak that appears above the centre of the broad peak at temperatures above 333 K. Fig. 7 shows a schematic of molecular motion in different phases (IV to I) of **26f** with increasing temperature. Four phases and melting were observed in this system, and it is attributed to the abrupt onset of uniaxial rotational motion of the cation during the transition from phase IV to phase III. This is also the dominant transport mechanism in phase III, where defect formation also results in the diffusion of a small number of anions. All the cations begin to rotate isotropically in phase II, allowing fast diffusion of a minority of cations together with the majority of anions. Perhaps, the most important point is that Jin *et al.* identified that the formation of defects in the ionic crystal was associated with this fast diffusion.

Jin *et al.* also report the effect of Li ion addition, including the investigation of defect effects, by MD simulation.⁸⁵ The effect of lithium ion doping into the structure and the cell-dynamics were theoretically analyzed with the OIPC matrix-model **11h**, at a complete crystal structure of 128 pairs arranged in a unit cell employed in the literature.⁸⁴ Here a Li addition procedure was performed by replacing the **11** cation with Li^+ to introduce impurities. The number of replacements was 5, 10, and 15. In doing so, Adebahr *et al.* obtained molecular insights into the significant increase in ionic conductivity when Li^+ was added to the plastic crystalline phase of OIPC. In particular, the RDF shows that the plastic crystalline phase is indeed the most sensitive phase to Li^+ doping. For example, in the case of a

perfect crystalline structure at low temperatures, the RDF value is completely zero at a certain radial value. But when Li^+ is added, the RDF value does not reach zero at the same low temperature, which indicates a melting process. This process is assigned to the formation of clusters between the **h** anion and Li^+ in the plastic crystalline phase, which causes a restriction of the rotational dynamics of the **h** anion on a time scale of about 10 ps. However, on longer time scales, this process creates a “free volume” in the rest of the crystal structure, which promotes the translational motion (diffusion coefficient) of all ionic species, especially **h** anions that are not trapped in clusters with Li^+ . The molecular transport model revealed in this paper is to our knowledge the first theoretical clarification of the mechanism for enhanced conduction mechanism observed in Li^+ -doped OIPCs.

The OIPC model consisting of **11h** has been theoretically elucidated so far, although Adebahr *et al.* could not compare with their theoretical results. Then, the same group performed the analogous simulations for the system consisting of **26f** reported in 2012,³³ where their experimental data were abundant.⁸⁶ This simulation revealed new dynamical heterogeneities; the van Hove self-correlation function $G_s(r, t)$ would be likely to be Gaussian distribution, if obtained from the Max-Boltzmann distribution for particles with Brownian motion in Fick's diffusion, which is a uniform diffusion. However, after approximately 50 ps in their simulations, the non-Gaussian shape appears; in other words, a non-uniform diffusion process was observed. The minority of ions with this anomalous dynamic process are identified in the plastic crystal phase. Furthermore, the introduction of vacancies in the crystal structure reveals the coexistence of mobile species in a rigid matrix, supporting the previous explanation for the non-uniform motion constriction based on NMR measurements. Finally, Chen *et al.* also evaluated the effect of Li^+ and Na^+ ion addition.⁸⁷ The mobility of Li^+ in this system appears to be higher than that of Na^+ . It becomes clear that the reason for this phenomenon is due to two different complex-coordination geometries for the first solvation shell of the alkali metal ions. The hopping motion of metal ions proposed here is a migration involving a sequence of destruction and reformation of this solvation structure; in the ion transport of Na^+ , the exchange between the tetrahedral and triangular solvation shell

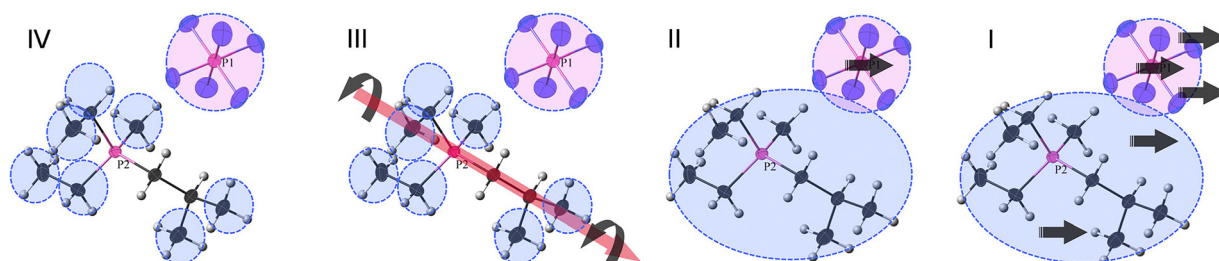


Fig. 7 Schematic of molecular motions in different phases (IV to I) of **26f** with increasing temperature. Phase IV, only the methyl and ethyl groups of the cation are rotating, but the anion is tumbling (on the NMR time scale); phase III, the whole cation can rotate around one axis, and the anion continues tumbling; phase II, the whole cation can tumble, and the anion both tumbles and diffuses; phase I, both the cation and the anion can diffuse. Reprinted from ref. 33 with permission.



coordination is dominant, while in Li^+ , only the cooperative motion between the metal ion and the triangular solvation shell is observed. As a result, Chen *et al.* proposed that the different geometries and associated potential barrier changes at the solvation structures of alkali metals are responsible for the differences in transport rates. They concluded that such an analysis could be used to screen OIPC systems with different structures and chemistries and to optimize the concentration of Li^+ or Na^+ for optimal transport. In a subsequent report, magnetic resonance imaging (MRI) analysis revealed that the domains of the crystal exhibit different orientations and anisotropic ionic conduction with high values in certain directions.⁸¹ In this paper, Forsyth *et al.* comprehensively discussed this new understanding, linking molecular-level information from theory and experiment with information at the macroscopic level.

It is very interesting to consider here a series of heterogeneous dynamics in terms of the time scales and observational methods associated with molecular motion. First, Henderson *et al.* investigated the phase behavior and ionic conductivity of OIPC crystals composed of a mixture of **13a** salt and LiTFSA.⁸⁸ In particular, they observed the temperature-dependent Raman spectra, in which a Raman peak shift attributed to the ν_s (C_4N) vibrational mode was observed, which led to distinguishing the (tg,tg) and (tt,tt) conformations of the **13** cation, where t and g indicate *trans*- and *gauche*-type. This study informs us that we can observe femtosecond molecular dynamics in OIPC. Furthermore, as mentioned above, Forsyth *et al.* have succeeded in observing sub-nanosecond to near-microsecond molecular dynamics by NMR measurement. Of course, they predicted rotational dynamics on picosecond time scales and alkali metal hopping motion on picosecond to nanosecond time scales by theoretical MD simulations. We are eagerly awaiting the observation of molecular dynamics in the picosecond order dynamics, which have not yet been reported.

Now, the addition of alkali metals leads to a large increase in ionic transport capacity, but different amounts and types of alkali metals need to be evaluated. Forsyth *et al.* performed MD simulations in mixed alkali metal salts. The addition of alkali metal ions was carried out up to 50 mol%. In both systems, the Na/Li ion concentration was found to have different effects depending on the ion species; it was suggested that alkali metal ion aggregates are formed at higher concentrations of alkali metal ions. However, the diffusion of alkali metal ions toward organic cations and anions was not affected by increasing alkali concentration. Some alkali metal ions moved most rapidly in the highest alkali metal concentration systems if the temperature was high enough to overcome the viscosity increase caused by this concentration increase. These transport features in the high concentration systems were attributed to the formation of bulky anion-alkali metal complexes. The authors stated that these complexes could enhance the dynamics of alkali metal ions in a clustered environment by increasing the opportunity to reorganize the ion coordination environment through rearrangement of the aggregates.

The ammonium-based OIPC **11b** was investigated in a recent paper; the temperature dependence of the structure

and ion dynamics were explored by a combination of single-crystal XRD, synchrotron powder XRD, and MD simulations.⁸⁹ Both MD simulations and X-ray diffraction showed consistent anisotropic changes in the crystal structure at high temperatures. The ionic force analysis in the simulations additionally revealed that there is a very strong hydrogen-bonding interaction in one direction of the unit cell and that the expansion in this direction is minimal upon heating. It is also explained that this strong interionic force restricts ionic motions such as rotation and translation compared to the **26f** system reported in others,⁹⁰ resulting in low plasticity and high structural order in the high temperature solid phase. The simplicity of the cation and anion motions, *i.e.*, the cation rotation and anion *cis-trans* transition resulting from the structural simplicity, leads to the simple phase behavior of **11b**, in contrast to the complex phase behavior of **26f**, which was derived from various levels of rotational, dislocation, and translational motions before the melting phase transition.

4. OIPC/polymer composites

As already mentioned, OIPC is a plastic solid and easily deforms when pressure is applied. When an electrochemical device using OIPC as an electrolyte is fabricated, it is assumed that if excessive pressure is applied to the device, the OIPC will be deformed, and the electrodes will come into contact with each other, resulting in a short circuit. Composites with polymers (Fig. 8) are being studied as a means of improving mechanical strength without impairing the softness of OIPC. Howlett *et al.* reported composites in which polyvinylidene fluoride (PVDF) nanofiber membranes were impregnated with **1e**.⁹¹ The yield stress and Young's modulus of the **1e**/PVDF composite were 300 and 60 MPa, respectively. Compression testing of **1e** was also performed. The compressive modulus of **1e** was 200 MPa at 30 °C. The compressive modulus of **1e** supplemented with 10 mol% LiBF_4 was 20 MPa, *i.e.*, a reduction of an order of magnitude. Furthermore, the compressive modulus of the **1e**/PVDF composite was 5 MPa, which was lower than that of the LiBF_4 -containing system. It is considered that **1e** was plasticized by mixing with PVDF. The ionic conductivity of the composite formed by impregnation of **1e** in the PVDF nanofiber membrane was about an order of magnitude higher than that of **1e** in phase I and II. In the case of the PVDF-**1e**/ LiBF_4

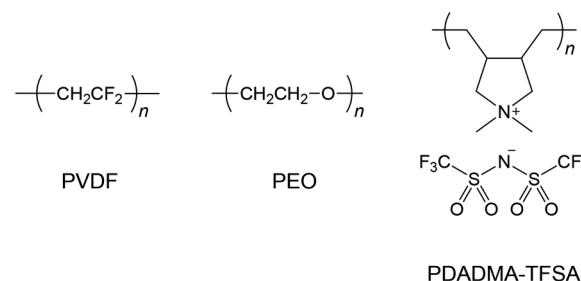


Fig. 8 Chemical structures of host polymers for the composites with OIPCs.



composite, the ionic conductivity was improved in phase II. Interestingly, mixing with PVDF, which is a non-conductive polymer, improved not only the mechanical strength but also the ionic conductivity. The ionic conductivity of the composite of **1b** impregnated in the PVDF nanofiber membrane was higher than that of **1b**.⁹² On the other hand, the ionic conductivity of the polystyrene-**1b** composite decreased with increasing polystyrene content.⁹³ When OIPC was mixed with PVDF, the ionic conductivity was improved regardless of the type of OIPC, demonstrating that it is a phenomenon peculiar to PVDF. PVDF is a ferroelectric polymer and is spontaneously polarized. It is believed that the interaction between PVDF and ions disrupted the arrangement of the constituent ions of the OIPC and increased the defects/pores on the PVDF surface, which promoted the rotational/translational motion of cations and anions. It was also suggested that the surface area of PVDF is an important factor in improving the physical properties of OIPC in the PVDF composite. Therefore, to increase the surface area of the polymer, the shape of PVDF was changed from fiber to fine particles, and a composite with OIPC was prepared.⁹⁴ When **1b** and PVDF fine particles were combined, the ionic conductivity of the **1b**/PVDF fine particle composite was higher than that of the analogous composite with PVDF fibers, demonstrating the importance of the OIPC/PVDF interface.

Formation of composites with polyethylene oxide (PEO), which is an ion conductive polymer, is also being considered. PEO is a rare polymer material in that it has both high polarity to dissociate added salts and high molecular mobility to transport the generated ions. It has been explored extensively as an electrolyte component for Li batteries.⁹⁵ However, to put these PEO/Li salt composites into practical use, it is necessary to further improve the ionic conductivity and t_{Li^+} . Li *et al.* added **1b** (0–40 wt%) to PEO/LiFSA ([Li]/[EO] = 0.05) to prepare a composite.⁹⁶ The glass transition temperature (T_g) of PEO/LiFSA decreased with increasing **1b** content, indicating that OIPC functions as a plasticizer in PEO. The tensile strength of PEO/LiFSA was about 4 MPa. Although the tensile strength of PEO/LiFSA decreased due to the addition of **1b**, it could maintain about 2 MPa. The addition of **1b** to PEO/LiFSA improved the ionic conductivity. It showed the highest value when **1b** was 30 wt% and was $1.05 \times 10^{-4} \text{ S cm}^{-1}$ at 50 °C. This is thought to be due to the decrease in T_g due to the addition of OIPC. A composite was prepared by adding **3b**⁹⁷ or **12b**⁹⁸ to PEO/LiFSA ([Li]/[EO] = 0.05). The thermal and mechanical properties of these composites were comparable to those of the **1b** composite. The ionic conductivity of the **3b** and **12b** composites was 3.02×10^{-4} and $2.14 \times 10^{-4} \text{ S cm}^{-1}$, respectively, at 50 °C. This is considered to be the result of the ionic conductivity of the added OIPC. The t_{Li^+} of the PEO/LiFSA/**1b** and **12b** composite s were 0.15 and 0.18, respectively, comparable to PEO/LiFSA. On the other hand, t_{Li^+} of PEO/LiFSA/**3b** was 0.45. It was suggested that the ion conduction mechanism in PEO changes depending on the type of OIPC added. A more detailed analysis is needed on the role of OIPC in PEO.

A composite of poly(diallyldimethylammonium) bis(trifluoromethanesulfonyl)amide (PDADMA-TFSA), which is a polycation, and

OIPC was investigated. It has been reported that a solid electrolyte obtained by adding a Li salt to PDADMA-TFSA functions as a lithium-ion conductor,⁹⁹ and PDADMA-TFSA is interesting as a matrix to which OIPC is added. Li *et al.* varied the mass percent of PDADMA-TFSA, **1b**, LiTFSA ([70/30]20, [60/40]20, [50/50]20 (LiTFSA fixed at 20%)).¹⁰⁰ As a result of DSC and XRD measurements, both composites were found to be amorphous and the **1b**-based phase transition temperature disappeared. [50/50]20, which has a high OIPC content, showed the highest ionic conductivity, with a value of $3.33 \times 10^{-4} \text{ S cm}^{-1}$ at 40 °C. It showed a higher value than the **1b**/PEO composite. On the other hand, t_{Li^+} at 40 °C was 0.17, which was comparable to that of the **1b**/PEO composite. A composite of PDADMA-TFSA with **12b** added was also investigated. Yang *et al.* altered mass percent of PDADMA-TFSA, **12b**, LiTFSA ([60/40]20, [50/50]20, [40/60]20 (LiTFSA fixed at 20%)).¹⁰¹ The **12b** composite was also amorphous, as was observed for the **1b** composite. The ionic conductivity of these **12b** composites improved with increasing OIPC content. The ionic conductivity of [40/60]20 was $2.08 \times 10^{-4} \text{ S cm}^{-1}$ at 25 °C, and t_{Li^+} was 0.18 at 40 °C. Although the two types of OIPC are compounded, the effect of PDADMA-TFSA on various characteristics is unclear and needs further investigation.

5. Lithium secondary batteries

Research on the use of OIPCs as electrolytes for electrochemical devices is becoming more active, and lithium secondary batteries, sodium secondary batteries,¹⁰² fuel cells,¹⁰³ dye-sensitized solar cells (DSSC),^{104,105} and electric double-layer capacitors,^{106,107} *etc.*, have been investigated as applicable electrochemical devices. Here, lithium secondary batteries will be described.

5.1 Preconditioning process

It has been reported that repeated cycling at a low current density (0.01 mA cm^{-2}) in symmetric cells using lithium metal for both electrodes improves OIPC/electrode interface contact and improves cell performance.^{108–110} This method, called the preconditioning process, reduced the internal resistance of the cell after the preconditioning process in symmetric cells of lithium metal using **1a**,^{108,110} **2e**,¹⁰⁹ **25a**,¹¹⁰ and **26f**.¹¹¹ Fig. 9 shows the SEM images of the **1a**/LiTFSA (1 mol%) and electrolyte/lithium metal interfaces before and after preconditioning.¹¹⁰ From the comparison of Fig. 9(a) and (b), it was observed that the particle size of **1a**/LiTFSA decreased after the preconditioning process. Furthermore, by comparing Fig. 9(c) and (d), it was observed that the electrolyte/lithium interface was obscured after the preconditioning process. It was revealed that the internal resistance was reduced by the miniaturization of OIPC particles and the good contact of the OIPC/electrode interface. It is believed that repeated dissolution and precipitation of lithium causes stress due to changes in the volume of the OIPC, resulting in a decrease in the particle size of the OIPC. It is interesting that the miniaturization of OIPC particles not only improves the contact



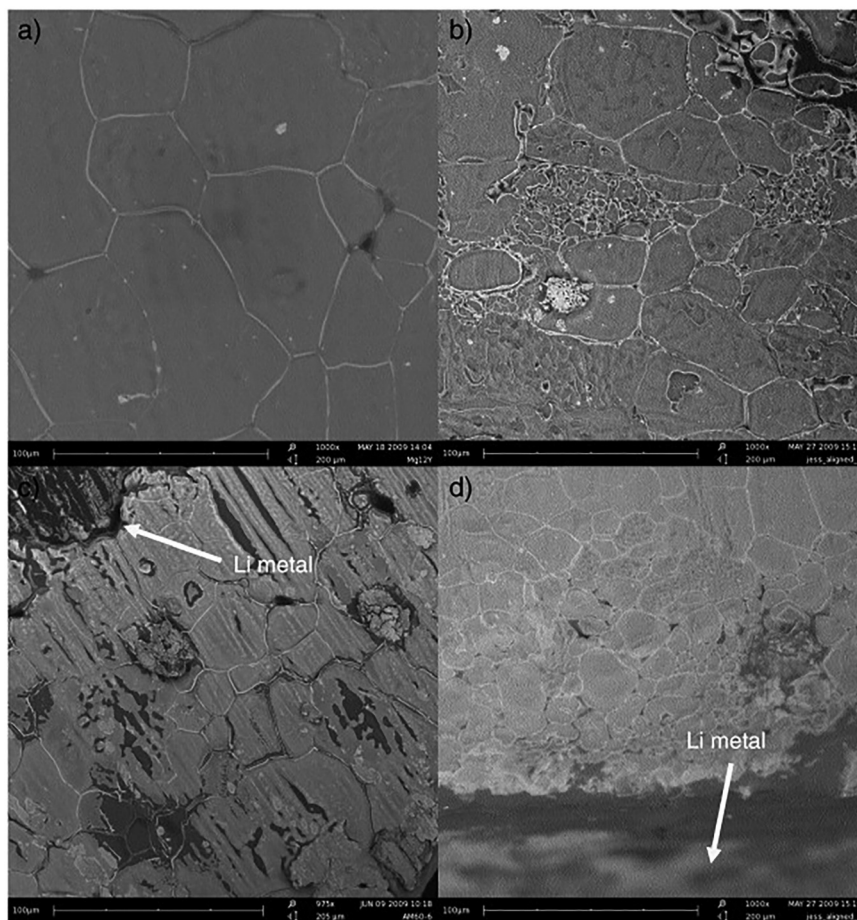


Fig. 9 SEM micrographs of a 1 mol% LiTfSA doped **1a** electrolyte. (a) Top-down view, (b) top-down view after pre-conditioning, (c) cross-section view, and (d) cross-section view after pre-conditioning. Reprinted from ref. 110 with permission.

with the electrode interface but also contributes to the improvement of ionic conductivity.

The importance of the preconditioning process was demonstrated in the Li/LiFePO₄ cell (**1a**/LiTfSA melt-impregnated in PVDF separator) using **1a**/LiTfSA (1 mol%).¹¹⁰ Repeated cycling at 0.01 mA cm⁻² at 50 °C using the above cells reduced cell resistance and resulted in good charge/discharge cycles after the preconditioning process. In charge/discharge tests (50 °C, C/5) of a Li/LiFePO₄ cell using **1a**/LiTfSA (10 mol% (0.26 mol kg⁻¹)), the initial discharge capacity was 129 mA h g⁻¹. After 50 cycles, the discharge capacity was 110 mA h g⁻¹. After 3 cycles, the coulombic efficiency improved sharply and remained around 99% thereafter. Furthermore, the effect of the preconditioning process was investigated by changing the thickness of the electrolyte (100 and 200 μm).¹¹¹ **1a**/LiTfSA (1 mol%) was used as the electrolyte. The internal resistance of the metallic lithium symmetric cell was similar before the preconditioning process, regardless of the electrolyte thickness. This indicates that the interfacial resistance is higher than the bulk resistance. After the preconditioning process, the 200 μm cell resistance determined by the Nyquist plot was twice that of 100 μm, reflecting the thickness of the electrolyte. The contribution to cell resistance was higher in bulk resistance,

indicating that preconditioning reduced interfacial resistance, demonstrating the effectiveness of the process. The effect of the separator was also examined. As a result of the charge/discharge test (80 °C, C/10) of Li/LiFePO₄ cells using **1a**/LiTfSA (10 mol%), the highest discharge capacity was shown in the order of glass filter, Separion[®], and PVDF. It is thought that this reflects the order of wettability, but it is also considered that the special effect of the PVDF surface also contributes, and it is thus clear that the selection of the separator is important.

A charge/discharge test of a Li/LiFePO₄ cell using **28b**/LiFSA (4 mol%) was conducted.³⁴ The electrolyte was melt-impregnated into a polyethylene separator. **28b**/LiFSA (4 mol%) $T_{\text{H-III}}$ is 8.2 °C, $T_{\text{I-II}}$ is 24.8 °C, and T_{m} is 34.6 °C. The ionic conductivity was 2.6×10^{-4} S cm⁻¹ at 22 °C. The discharge capacity at 30 °C (phase I) and 0.3C reached 160 mA h g⁻¹ after 5 cycles, and the coulombic efficiency improved to 99%. This is thought to be due to the formation of a stable solid electrolyte interphase (SEI) on the electrodes and the miniaturization of OIPC particles by the preconditioning process. The cell using **28b**/LiFSA showed a discharge capacity of 130 mA h g⁻¹ even at 1C. The discharge capacity at 20 °C (phase II) and C/10 reached 155 mA h g⁻¹ after 10 cycles, which is similar to cells using ionic liquid as the electrolyte,^{112,113} and it was higher than cells



using other OIPCs.^{111,114} The interfacial resistance after 10 cycles was 1000 Ω or lower than the value before the cycle, and the effect of the preconditioning process was observed more significantly than the conditions of 30 $^{\circ}\text{C}$ and 0.3C.

5.2 Charge/discharge tests

In 2004, Alarco *et al.* reported a charge/discharge test of $\text{Li}_4\text{Ti}_5\text{O}_{12}/\text{LiFePO}_4$ cells at 40 $^{\circ}\text{C}$ using **21a**/LiTFSa (20 mol%).⁶² The melting point of **21a**/LiTFSa (20 mol%) is 56 $^{\circ}\text{C}$, and the ionic conductivity at 40 $^{\circ}\text{C}$ is $6.2 \times 10^{-4} \text{ S cm}^{-1}$. To prevent cell short-circuiting, the electrolyte was melt-impregnated into Celgard[®]. The discharge capacity at C/24 was 90 mA h g^{-1} , and the discharge capacity decreased as the C-rate increased. The coulombic efficiency was almost 100% at all C-rates, suggesting that no side reactions occurred. The cell was made using **23a**, which has better ionic conductivity than **21a**.⁶⁴ The melting point of **23a**/LiTFSa (10 mol%) was 7.3 $^{\circ}\text{C}$, and the ionic conductivity at 5 $^{\circ}\text{C}$, which is a plastic crystal phase, was $6.2 \times 10^{-4} \text{ S cm}^{-1}$. The discharge capacity of the $\text{Li}_4\text{Ti}_5\text{O}_{12}/\text{LiFePO}_4$ cell using this electrolyte (melt impregnated in Celgard[®]) at 5 $^{\circ}\text{C}$ is 92 mA h g^{-1} at C/12, and **23a**/LiTFSa has better charge/discharge characteristics than **21a**/LiTFSa. On the other hand, the development of a lithium battery using **1e** was examined without a separator. **1e** has excellent mechanical strength, and it is presumed that this was taken into consideration. A charge/discharge test of a $\text{Li}/\text{LiFePO}_4$ cell using a **1e**/LiBF₄ composite was performed. The discharge capacity of the cell is about 140 and 90 mA h g^{-1} at 100 $^{\circ}\text{C}$, C/10, 80 $^{\circ}\text{C}$, and C/10, respectively, which is about 80% and 50% of the theoretical capacity (168 mA h g^{-1}). The initial capacity at each temperature was maintained after 20 cycles.

In recent years, device fabrication using OIPC/polymer composites has become mainstream. There are many merits, such as device safety and improved ionic conductivity of OIPCs, so it will be a natural flow. Howlett *et al.* prepared and evaluated $\text{Li}/\text{LiFePO}_4$ cells using a composite in which a PVDF nanofiber membrane was impregnated into **1e**.⁹¹ At 80 $^{\circ}\text{C}$ (phase I) and C/15, the cell discharge capacity was approximately 140 mA h g^{-1} . Although the C-rate is different, the discharge capacity at 80 $^{\circ}\text{C}$ increased compared to the result of **1e** alone, and it is inferred that the improvement of ionic conductivity due to the effect of PVDF nanofibers contributes. $\text{Li}/\text{LiFePO}_4$ cells were prepared using a solid electrolyte in which **1b** and PEO⁹⁶ or PDADMA-TFSA,¹⁰⁰ **3b** and PEO,⁹⁷ and **12b** and PEO⁹⁸ were combined. When a composite of PEO/LiFSA ([Li]/[EO] = 0.05) and **1b** (30 wt%) is used, the discharge capacity at 50 $^{\circ}\text{C}$ and C/5 gradually increases with the number of cycles, and 5 cycles. After that, it reached 157.3 mA h g^{-1} . After 90 cycles, it maintained a capacity of 150.3 mA h g^{-1} . The discharge capacities at 50 $^{\circ}\text{C}$, C/2, and 1C were 148.1 and 115.6 mA h g^{-1} , respectively, and the coulombic efficiency was almost 100% at both C-rates. When a composite of PEO/LiFSA ([Li]/[EO] = 0.05) and **3b** (30 wt%) was used, the discharge capacities at 50 $^{\circ}\text{C}$ were 158.4 at C/5, 151.5 at C/2, and 113.7 mA h g^{-1} 1C, respectively. When **1b** and PDADMA-TFSA ([50/50]20) were used, the discharge capacity at 40 $^{\circ}\text{C}$ and C/5 reached 152 mA h g^{-1} after several cycles. The discharge capacities at 40 $^{\circ}\text{C}$, C/2 and 1C were 147.6 and 111.2 mA h g^{-1} ,

respectively, and the coulombic efficiency was almost 100% at both C-rates. There are also reports of cells using PEO/LiFSA ([Li]/[EO] = 0.05)/**12b** (30 wt%) and **12b**/PDADMA-TFSA ([60/40]20) composites as electrolytes. The discharge capacity of these cells were in the range of 155 and 158 mA h g^{-1} at C/5, and the cell performance of **1b**/PEO, **3b**/PEO, **12b**/PEO, **1b**/PDADMA-TFSA, and **12b**/PDADMA-TFSA complexes was almost the same. It is considered that the result reflects the ionic conductivity and t_{Li^+} of these complexes, and the specific effects of OIPC and polymer structure on the cell performance could not be found. On the other hand, $\text{Li}/\text{LiFePO}_4$ cells prepared using a composite of **1b**/LiFSA (amount of substance ratio 9:1) and PVDF fine particles (40 wt%) showed different behavior (Fig. 10(a)).⁹⁴ The discharge capacity of the cell at 50 $^{\circ}\text{C}$ and C/10 was 124 mA h g^{-1} , which was lower than that of the above composite but increased to 128 mA h g^{-1} at C/2. The discharge capacities at 1, 2, and 5C were 127, 124, and 115 mA h g^{-1} , respectively. Even though the charge/discharge rate is doubled, the discharge capacities at 1 and 2C were the same within error, and even at 5C, 90% of the discharge capacity at 1C was maintained. The effect of the C-rate on the

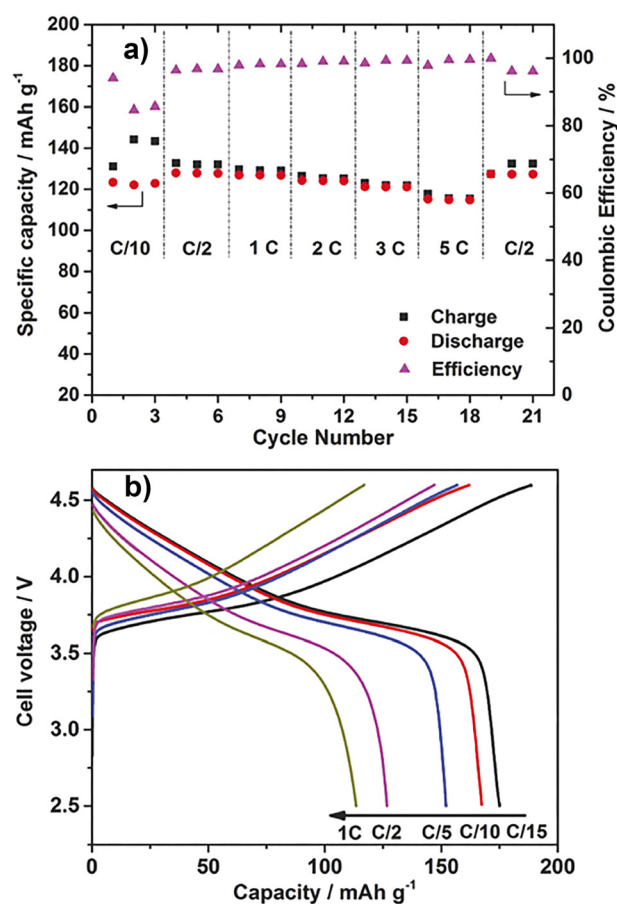


Fig. 10 Electrochemical measurements of 40 wt% $\text{Li}_{0.1}\text{Ni}_{0.9}\text{b}/\text{PVDF}$ composites. (a) Specific capacities and coulombic efficiencies of $\text{Li}/\text{LiFePO}_4$ cells at different current rates, ranging from C/10 to 5C, 50 $^{\circ}\text{C}$. (b) Charge-discharge curves of the $\text{Li}/\text{LiCo}_{1/3}\text{Ni}_{1/3}\text{Mn}_{1/3}\text{O}_2$ cell using 40 wt% $\text{Li}_{0.1}\text{Ni}_{0.9}\text{b}/\text{PVDF}$ composites (cut-off voltage 2.5–4.6 V), 50 $^{\circ}\text{C}$. (Figure partly modified from the original one.) Reprinted from ref. 94 with permission.



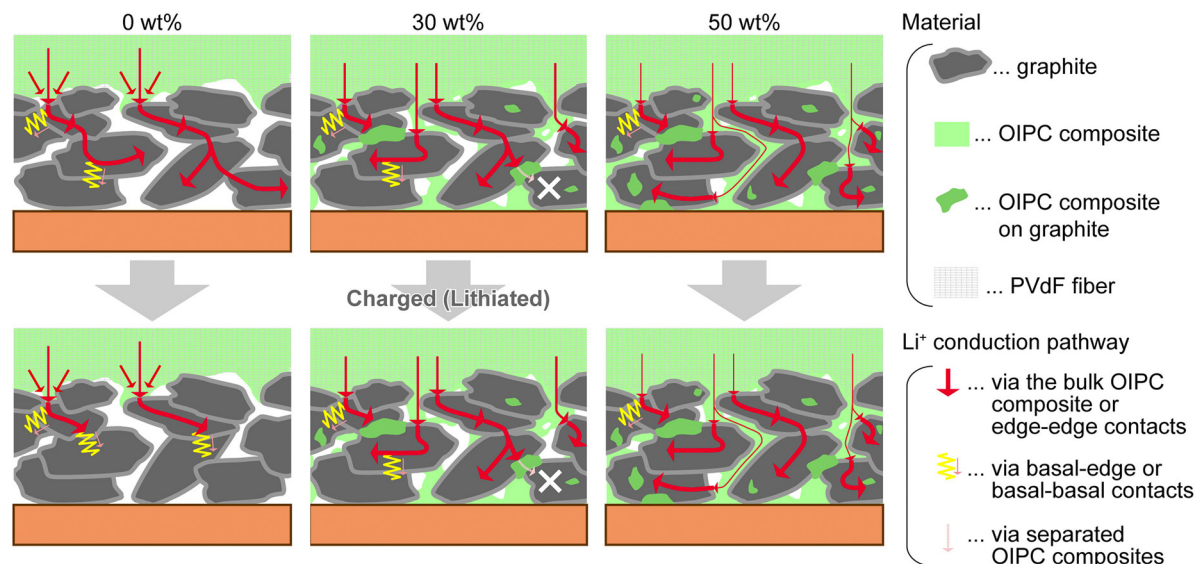


Fig. 11 Role of OIPC in the solid-state electrodes. Mixing OIPC with graphite anode promotes Li^+ conduction and electrochemical reaction in the electrode layer. Reprinted from ref. 115 with permission.

discharge capacities was small. A phenomenon peculiar to PVDF, and not found in other polymers, was observed. In addition, cycle tests were performed at room temperature and 2C. At the beginning of the cycling, the discharge capacity was about 100 mA h g^{-1} . The discharge capacity after 1200 cycles was about 80 mA h g^{-1} , and the coulombic efficiency was maintained at almost 100%. In a similar test using LP30, which is an electrolyte solution, the discharge capacity at the beginning of the cycle was about the same, but after 630 cycles, it decreased to about 70 mA h g^{-1} , showing the superiority of the OIPC/PVDF fine particle composite.

The development of a lithium secondary battery using $\text{LiNi}_{1/3}\text{Co}_{1/3}\text{Mn}_{1/3}\text{O}_2$ (NCM111) as a positive electrode active material to replace LiFePO_4 has been reported.⁹⁴ Since the OIPC potential window is wide, high voltage positive electrode materials such as NCM can be applied. A-Li/NCM111 cell was prepared using a composite of **1b**/LiFSA (amount of substance ratio 9:1) and PVDF fine particles (40 wt%), and a charge/discharge test was performed at a cut-off voltage of 2.5 to 4.6 V (charge/discharge test). Fig. 10(b). The discharge capacity of the Li/NCM111 cell at 50°C was 178 mA h g^{-1} at C/15 and 113 mA h g^{-1} at 1C. The value was higher than that of the Li/ LiFePO_4 cell. It was also reported that cells were prepared and evaluated using $\text{LiNi}_{0.6}\text{Co}_{0.2}\text{Mn}_{0.2}\text{O}_2$ (NCM622) with an increased Ni ratio for higher energy density.⁹⁷ The discharge capacity of the Li/NCM622 cell using the composite of PEO/LiFSA ([Li]/[EO] = 0.05) and **3b** (30 wt%) as the electrolyte is 60°C , C/5, C/2, and 1C. At that time, they were 155, 133, and 103 mA h g^{-1} respectively (cutoff voltage was 2.8 to 4.3 V). It was found that various positive electrode active materials are applicable.

OIPCs are beneficial for interface engineering in all-solid-state batteries.^{91,94,111,115,116} The compounds contribute to faster charge/discharge of solid-state batteries by forming continuous ion-conduction pathways throughout the electrode layer (Fig. 11).

Ueda *et al.* reported the electrode characteristics of an OIPC composite(**1b** and LiFSA) and graphite as the active material.¹¹⁵ The experimental capacity of the prepared electrode reached over 80% against the theoretical value at a rate of 1C, whereas only 50% was obtained without OIPC. The addition of the OIPC composite also increased cycle performances (>100 cycles). The improvements were explained by the ion-conducting and binding properties of the OIPC composite, accelerating the charge transport throughout the electrolyte layer, active material, and current collector.

From a practical viewpoint, OIPC may increase the energy density of solid-state batteries. The paper claimed that the volume fraction of the electrolyte components could be smaller than the usual liquid electrolyte systems; OIPCs filled the electrode voids effectively.¹¹⁵ On the other hand, suppression of side reactions at the electrode interfaces is an ongoing challenge. The initial coulombic efficiency of the OIPC electrodes tended to be lower (88–95%) than the regular liquid electrolyte system (>97%), also taking longer (>20) cycles for stable solid electrolyte interphase (SEI) formation. Molecular engineering is critical to effectively suppressing unfavorable degradation of OIPC, primarily occurring at the anode side where bond dissociation of OIPCs proceeds easily.¹¹⁷

6. Conclusions and outlook

The development of lithium secondary batteries using OIPCs as electrolytes is steadily progressing. However, the ionic conductivity of OIPC is lower than that of electrolyte solutions and inorganic solid electrolytes, so further improvement will be necessary. OIPCs exceeding $10^{-3} \text{ S cm}^{-1}$ at room temperature have been reported,¹¹⁸ and it is expected that high lithium-ion conductivity can be achieved by clarifying molecular design



guidelines in the future. An effort has already been initiated to better understand the correlation between chemical structure and motion-mode of each ion for OIPCs.¹¹⁹ The motion-mode for each ion was discussed by comparing ion symmetry and the site symmetry where the ion was positioned based on their powder crystal diffraction data. The diffraction intensities of samples with spherically rotating ions were calculated. In addition, a correlation between the ionic radius ratio of cations and anions composed of OIPCs and the ionic conductivity has been reported, providing an indicator for molecular design guidelines.^{56,57} Further understanding of the correlation between chemical structure and physicochemical properties will further improve ionic conductivity.

On the other hand, artificial intelligence (AI)-based screening of highly conductive OIPCs will be an interesting challenge. AI-assisted exploration of solid polymer electrolytes has already been reported,¹²⁰ and glass-type lithium-ion conductors with aromatic structures, which is an unfamiliar electrolyte design, showed the ionic conductivity of around 10^{-3} S cm⁻¹ at room temperature. The glassy polymer electrolyte design was contrary to the traditional concept of rubbery polymer electrolytes. The suggestions provided by machine learning models can help to discover unexpected chemical structures with a high ionic conductivity.

Furthermore, as with the synergistic effect of mixing OIPC and macromolecules, mixing OIPC with organic or inorganic materials is also an interesting technique for achieving high lithium-ion conductivity. For example, the t_{Li^+} of the composite of the covalent organic framework (COF) and OIPC was 0.58, which is a high value for a Li salt mixed system.¹²¹ There is also a report that OIPC was applied to the surface of Li metal by spin coating to reduce the contact resistance between Li-La₃Zr₂O₁₂ (LLZO), which is a solid ceramic electrolyte, and the electrode.¹²² On the other hand, the use of Li metal anodes in sulfide electrolyte (SE)-based all-solid-state lithium metal batteries (ASSLMBs) has been less successful. The main challenges are the pronounced interfacial reactions between Li metal and SE and Li dendrite formation. A solid plastic crystal electrolyte (PCE) based on succinonitrile and LiTFSa was designed as an interlayer for SE-based ASSLMB, and it was demonstrated that the PCE interlayer can prevent the interfacial reaction between SE and Li metal and the formation of Li dendrites by Wang *et al.*¹²³ These are good examples of utilizing the softness of OIPC and PC is a technology that simplifies the cell fabrication process using solid inorganic solid electrolytes. In the future, it is expected that the implementation of OIPC not only in lithium secondary batteries but also in other electrochemical devices will be achieved by elucidating the correlation between chemical structures and physicochemical properties and optimizing the combination of dissimilar materials.

Author contributions

Conceptualization: MYF; resources: MLT, KHS, SN, MYF; data curation: MLT, KHS, SN, MYF; formal analysis: MLT, KHS, SN,

MYF; supervision: MYF; funding acquisition: MLT, KHS, SN, MYF; investigation: MLT, KHS, SN, MYF; visualization: MLT, KHS, SN, MYF; methodology: MLT, KHS, SN, MYF; project administration: MYF; writing – original draft: MLT, KHS, SN, MYF; writing – review & editing: MLT, KHS, SN, MYF.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

Recent studies were supported by JSPS KAKENHI grant number 19K05604 for M. Y. F., JSPS Bilateral Program grant number JPJSBP120199977 for M. Y. F., and a Sophia University Special Grant for Academic Research. The authors also acknowledge the great efforts of staff and students working on ionic plastic crystals in our laboratories.

Notes and references

- 1 J. M. Tarascon and M. Armand, *Nature*, 2001, **414**, 359–367, DOI: [10.1038/35104644](#).
- 2 S. Chen, K. Wen, J. Fan, Y. Bando and D. Golberg, *J. Mater. Chem. A*, 2018, **6**, 11631–11663, DOI: [10.1039/C8TA03358G](#).
- 3 S. K. Sharma, G. Sharma, A. Gaur, A. Arya, F. S. Mirsafi, R. Abolhassani, H.-G. Rubahn, J.-S. Yu and Y. K. Mishra, *Energy Adv.*, 2022, **1**, 457–510, DOI: [10.1039/D2YA00043A](#).
- 4 Y. Ma, R. Shang, Y. Liu, R. Lake, M. Ozkan and C. S. Ozkan, *J. Power Sources*, 2023, **559**, 232647, DOI: [10.1016/j.jpowsour.2023.232647](#).
- 5 Y. Lu, X. Huang, Y. Ruan, Q. Wang, R. Kun, J. Yang and Z. Wen, *J. Mater. Chem. A*, 2018, **6**, 18853–18858, DOI: [10.1039/C8TA07241H](#).
- 6 Y. Liu, Q. Sun, Y. Zhao, B. Wang, P. Kaghazchi, K. R. Adair, R. Li, C. Zhang, J. Liu, L.-Y. Kuo, Y. Hu, T.-K. Sham, L. Zhang, R. Yang, S. Lu, X. Song and X. Sun, *ACS Appl. Mater. Interfaces*, 2018, **10**, 31240–31248, DOI: [10.1021/acsami.8b06366](#).
- 7 V. Lacivita, A. S. Westover, A. Kercher, N. D. Phillip, G. Yang, G. Veith, G. Ceder and N. J. Dudney, *J. Am. Chem. Soc.*, 2018, **140**, 11029–11038, DOI: [10.1021/jacs.8b05192](#).
- 8 R. Murugan, V. Thangadurai and W. Weppner, *Angew. Chem., Int. Ed.*, 2007, **46**, 7778–7781, DOI: [10.1002/anie.200701144](#).
- 9 D.-L. Xiao, J. Tong, Y. Feng, G.-H. Zhong, W.-J. Li and C.-L. Yang, *Solid State Ionics*, 2018, **324**, 202–206, DOI: [10.1016/j.ssi.2018.07.011](#).
- 10 J. Glenneberg, G. Kasiri, I. Bardenhagen, F. La Mantia, M. Busse and R. Kun, *Nano Energy*, 2019, **57**, 549–557, DOI: [10.1016/j.nanoen.2018.12.070](#).
- 11 M. Calpa, N. C. Rosero-Navarro, A. Miura and K. Tadanaga, *Electrochim. Acta*, 2019, **296**, 473–480, DOI: [10.1016/j.electacta.2018.11.035](#).



- 12 J. Timmermans, *J. Phys. Chem. Solids*, 1961, **18**, 1–8, DOI: [10.1016/0022-3697\(61\)90076-2](#).
- 13 D. R. MacFarlane and M. Forsyth, *Adv. Mater.*, 2001, **13**, 957–966, DOI: [10.1002/1521-4095\(200107\)13:12/13<957::AID-ADMA957>3.0.CO;2-#](#).
- 14 J. M. Pringle, P. C. Howlett, D. R. MacFarlane and M. Forsyth, *J. Mater. Chem.*, 2010, **20**, 2056–2062, DOI: [10.1039/B920406G](#).
- 15 J. M. Pringle, *Phys. Chem. Chem. Phys.*, 2013, **15**, 1339–1351, DOI: [10.1039/c2cp43267f](#).
- 16 H. Zhu, D. R. MacFarlane, J. M. Pringle and M. Forsyth, *Trends Chem.*, 2019, **1**, 126–140, DOI: [10.1016/j.trechm.2019.01.002](#).
- 17 P.-J. Alarco, Y. Abu-Lebdeh, A. Abouimrane and M. Armand, *Nat. Mater.*, 2004, **3**, 476–481, DOI: [10.1038/nmat1158](#).
- 18 J. Hwang, K. Matsumoto, C.-Y. Chen and R. Hagiwara, *Energy Environ. Sci.*, 2021, **14**, 5834–5863, DOI: [10.1039/D1EE02567H](#).
- 19 E. I. Cooper and C. A. Angell, *Solid State Ionics*, 1986, **18–19**, 570–576, DOI: [10.1016/0167-2738\(86\)90180-3](#).
- 20 H. Ishida, R. Ikeda and D. Nakamura, *Bull. Chem. Soc. Jpn.*, 1986, **59**, 915–924, DOI: [10.1246/bcsj.59.915](#).
- 21 S. Sato, R. Ikeda and D. Nakamura, *Bull. Chem. Soc. Jpn.*, 1986, **59**, 1981–1989, DOI: [10.1246/bcsj.59.1981](#).
- 22 T. Tanabe, D. Nakamura and R. Ikeda, *J. Chem. Soc., Faraday Trans.*, 1991, **87**, 987–990, DOI: [10.1039/FT9918700987](#).
- 23 M. Tansho, D. Nakamura and R. Ikeda, *J. Chem. Soc., Faraday Trans.*, 1991, **87**, 3255–3258, DOI: [10.1039/FT9918703255](#).
- 24 H. Ishida, N. Hayama and R. Ikeda, *Chem. Lett.*, 1992, 1333–1336, DOI: [10.1246/cl.1992.1333](#).
- 25 H. Ishida, T. Iwachido and R. Ikeda, *Ber. Bunsenges. Phys. Chem.*, 1992, **96**, 1468–1470, DOI: [10.1002/bbpc.19920961023](#).
- 26 H. Ishida and Y. Furukawa, *Z. Naturforsch., A: Phys. Sci.*, 1996, **51**, 83–86, DOI: [10.1515/zna-1996-1-212](#).
- 27 H. Ono, S. i Ishimaru, R. Ikeda and H. Ishida, *Bull. Chem. Soc. Jpn.*, 1999, **72**, 2049–2054, DOI: [10.1246/bcsj.72.2049](#).
- 28 D. R. MacFarlane, J. Huang and M. Forsyth, *Nature*, 1999, **402**, 792–794, DOI: [10.1038/45514](#).
- 29 M. Forsyth, J. Huang and D. R. MacFarlane, *J. Mater. Chem.*, 2000, **10**, 2259–2265, DOI: [10.1039/b003168m](#).
- 30 J. Golding, N. Hamid, D. R. MacFarlane, M. Forsyth, C. Forsyth, C. Collins and J. Huang, *Chem. Mater.*, 2001, **13**, 558–564, DOI: [10.1021/cm000625w](#).
- 31 A. L. Chong, M. Forsyth and D. R. MacFarlane, *Electrochim. Acta*, 2015, **159**, 219–226, DOI: [10.1016/j.electacta.2015.01.180](#).
- 32 M. Lee, Y.-H. Lee, J. H. Park and U. H. Choi, *Org. Electron.*, 2017, **48**, 241–247, DOI: [10.1016/j.orgel.2017.06.004](#).
- 33 L. Jin, K. M. Nairn, C. M. Forsyth, A. J. Seeber, D. R. MacFarlane, P. C. Howlett, M. Forsyth and J. M. Pringle, *J. Am. Chem. Soc.*, 2012, **134**, 9688–9697, DOI: [10.1021/ja301175v](#).
- 34 L. Jin, P. C. Howlett, J. M. Pringle, J. Janikowski, M. Armand, D. R. MacFarlane and M. Forsyth, *Energy Environ. Sci.*, 2014, **7**, 3352–3361, DOI: [10.1039/C4EE01085J](#).
- 35 T. Mochida, Y. Funasako, M. Ishida, S. Saruta, T. Kosone and T. Kitazawa, *Chem. – Eur. J.*, 2016, **22**, 15725–15732, DOI: [10.1002/chem.201603170](#).
- 36 T. Tominaga, T. Ueda and T. Mochida, *Phys. Chem. Chem. Phys.*, 2017, **19**, 4352–4359, DOI: [10.1039/C6CP08308K](#).
- 37 H. Kimata and T. Mochida, *Cryst. Growth Des.*, 2018, **18**, 7562–7569, DOI: [10.1021/acs.cgd.8b01390](#).
- 38 T. Mochida, M. Ishida, T. Tominaga, K. Takahashi, T. Sakurai and H. Ohta, *Phys. Chem. Chem. Phys.*, 2018, **20**, 3019–3028, DOI: [10.1039/C7CP02746J](#).
- 39 T. Hayasaki, S. Hirakawa and H. Honda, *Bull. Chem. Soc. Jpn.*, 2013, **86**, 993–1001, DOI: [10.1246/bcsj.20130040](#).
- 40 X. Chen, H. Tang, T. Putzeys, J. Snickers, M. Wübbenhorst, K. Binnemans, J. Fransaer, D. E. De Vos, Q. Li and J. Luo, *J. Mater. Chem. A*, 2016, **4**, 12241–12252, DOI: [10.1039/C6TA05472B](#).
- 41 M. Matsuki, T. Yamada, N. Yasuda, S. Dekura, H. Kitagawa and N. Kimizuka, *J. Am. Chem. Soc.*, 2018, **140**, 291–297, DOI: [10.1021/jacs.7b10249](#).
- 42 M. Yoshizawa-Fujita, E. Kishi, M. Suematsu, T. Takekawa and M. Rikukawa, *Chem. Lett.*, 2014, **43**, 1909–1911, DOI: [10.1246/cl.140833](#).
- 43 H. Yamada, Y. Miyachi, Y. Takeoka, M. Rikukawa and M. Yoshizawa-Fujita, *Electrochim. Acta*, 2019, **303**, 293–298, DOI: [10.1016/j.electacta.2019.02.076](#).
- 44 D. Al-Masri, R. Yunis, H. Zhu, L. Jin, P. Bruce, A. F. Hollenkamp and J. M. Pringle, *J. Mater. Chem. A*, 2019, **7**, 25389–25398, DOI: [10.1039/c9ta11175a](#).
- 45 D. Al-Masri, R. Yunis, A. F. Hollenkamp, C. M. Doherty and J. M. Pringle, *Phys. Chem. Chem. Phys.*, 2020, **22**, 18102–18113, DOI: [10.1039/D0CP03046E](#).
- 46 D. Al-Masri, R. Yunis, A. F. Hollenkamp and J. M. Pringle, *ChemElectroChem*, 2020, **7**, 4118–4123, DOI: [10.1002/celec.202000772](#).
- 47 H.-T. Kim, O. M. Kwon, J. Mun, S. M. Oh, T. Yim and Y. G. Kim, *Electrochim. Acta*, 2017, **240**, 267–276, DOI: [10.1016/j.electacta.2017.04.078](#).
- 48 M. Moriya, D. Kato, W. Sakamoto and T. Yogo, *Chem. Commun.*, 2011, **47**, 6311–6313, DOI: [10.1039/c1cc00070e](#).
- 49 Y. Abu-Lebdeh, E. Austin and I. J. Davidson, *Chem. Lett.*, 2009, **38**, 782–783, DOI: [10.1246/cl.2009.782](#).
- 50 M. Yoshizawa-Fujita, D. MacFarlane, P. Howlett and M. Forsyth, *Electrochem. Commun.*, 2006, **8**, 445–449, DOI: [10.1016/j.elecom.2006.01.008](#).
- 51 Y. Lauw, T. Rüther, M. D. Horne, K. S. Wallwork, B. W. Skelton, I. C. Madsen and T. Rodopoulos, *Cryst. Growth Des.*, 2012, **12**, 2803–2813, DOI: [10.1021/cg201592u](#).
- 52 R. Yunis, D. Al-Masri, A. F. Hollenkamp, C. M. Doherty, H. Zhu and J. M. Pringle, *J. Electrochem. Soc.*, 2020, **167**, DOI: [10.1149/1945-7111/ab76a2](#).
- 53 K. Nishikawa, T. Yamada, K. Fujii, H. Masu, K.-I. Tozaki and T. Endo, *Bull. Chem. Soc. Jpn.*, 2021, **94**, 2011–2018, DOI: [10.1246/bcsj.20210159](#).



- 54 Z.-B. Zhou and H. Matsumoto, *Electrochem. Commun.*, 2007, **9**, 1017–1022, DOI: [10.1016/j.elecom.2006.12.012](#).
- 55 J. Huang, A. Hill, M. Forsyth, D. MacFarlane and A. Hollenkamp, *Solid State Ionics*, 2006, **177**, 2569–2573, DOI: [10.1016/j.ssi.2006.01.041](#).
- 56 S. Yamaguchi, H. Yamada, Y. Takeoka, M. Rikukawa and M. Yoshizawa-Fujita, *New J. Chem.*, 2019, **43**, 4008–4012, DOI: [10.1039/c8nj05127e](#).
- 57 M. Yoshizawa-Fujita, H. Yamada, S. Yamaguchi, H. Zhu, M. Forsyth, Y. Takeoka and M. Rikukawa, *Batteries Supercaps*, 2020, **3**, 884–891, DOI: [10.1002/batt.202000040](#).
- 58 R. Yunis, J. M. Pringle, X. Wang, G. M. A. Girard, R. Kerr, H. Zhu, P. C. Howlett, D. R. MacFarlane and M. Forsyth, *J. Mater. Chem. A*, 2020, **8**, 14721–14735, DOI: [10.1039/D0TA03502E](#).
- 59 X.-Y. Zhao, J.-L. Wang, H. Luo, H.-R. Yao, C.-Y. Ouyang and L.-Z. Zhang, *J. Zhejiang Univ., Sci., A*, 2016, **17**, 155–162, DOI: [10.1631/jzus.A1500099](#).
- 60 M. Egashira, S. Okada, J.-I. Yamaki, D. A. Dri, F. Bonadies and B. Scrosati, *J. Power Sources*, 2004, **138**, 240–244, DOI: [10.1016/j.jpowsour.2004.06.022](#).
- 61 M. Egashira, M. Nakagawa, I. Watanabe, S. Okada and J.-I. Yamaki, *J. Power Sources*, 2005, **146**, 685–688, DOI: [10.1016/j.jpowsour.2005.03.069](#).
- 62 P.-J. Alarco, Y. Abu-Lebdeh, N. Ravet and M. Armand, *Solid State Ionics*, 2004, **172**, 53–56, DOI: [10.1016/j.ssi.2004.02.029](#).
- 63 P.-J. Alarco, Y. Abu-Lebdeh and M. Armand, *Solid State Ionics*, 2004, **175**, 717–720, DOI: [10.1016/j.ssi.2003.10.024](#).
- 64 Y. Abu-Lebdeh, A. Abouimrane, P.-J. Alarco and M. Armand, *J. Power Sources*, 2006, **154**, 255–261, DOI: [10.1016/j.jpowsour.2005.03.231](#).
- 65 K. Biernacka, D. Al-Masri, R. Yunis, H. Zhu, A. F. Hollenkamp and J. M. Pringle, *Electrochim. Acta*, 2020, **357**, 136863, DOI: [10.1016/j.electacta.2020.136863](#).
- 66 M. Volel, P.-J. Alarco, Y. Abu-Lebdeh and M. Armand, *ChemPhysChem*, 2004, **5**, 1027–1033, DOI: [10.1002/cphc.200400070](#).
- 67 V. Armel, D. Velayutham, J. Sun, P. C. Howlett, M. Forsyth, D. R. MacFarlane and J. M. Pringle, *J. Mater. Chem.*, 2011, **21**, DOI: [10.1039/c1jm10417a](#).
- 68 K. Romanenko, L. Jin, L. A. Madsen, J. M. Pringle, L. A. O'Dell and M. Forsyth, *J. Am. Chem. Soc.*, 2014, **136**, 15638–15645, DOI: [10.1021/ja508290z](#).
- 69 K. Romanenko, L. Jin, P. Howlett and M. Forsyth, *Chem. Mater.*, 2016, **28**, 2844–2851, DOI: [10.1021/acs.chemmater.6b00797](#).
- 70 M. Hattori, S.-i. Fukada, D. Nakamura and R. Ikeda, *J. Chem. Soc., Faraday Trans.*, 1990, **86**, DOI: [10.1039/ft9908603777](#).
- 71 S. Iwai, M. Hattori, D. Nakamura and R. Ikeda, *J. Chem. Soc., Faraday Trans.*, 1993, **89**, 827–831, DOI: [10.1039/ft9938900827](#).
- 72 T. Shimizu, S. Tanaka, N. Onoda-Yamamuro, S. I. Ishimaru and R. Ikeda, *J. Chem. Soc., Faraday Trans.*, 1997, **93**, 321–326, DOI: [10.1039/a605760h](#).
- 73 J. M. Chezeau and J. H. Strange, *Phys. Rep.*, 1979, **53**, 1–92, DOI: [10.1016/0370-1573\(79\)90077-2](#).
- 74 H. Zhu and L. A. O'Dell, *Chem. Commun.*, 2021, **57**, 5609–5625, DOI: [10.1039/d1cc02151f](#).
- 75 M. Forsyth, T. Chimdi, A. Seeber, D. Gunzelmann and P. C. Howlett, *J. Mater. Chem. A*, 2014, **2**, 3993–4003, DOI: [10.1039/c3ta15153k](#).
- 76 F. Makhlooghiyazad, L. A. O'Dell, L. Porcarelli, C. Forsyth, N. Quazi, M. Asadi, O. Hutt, D. Mecerreyes, M. Forsyth and J. M. Pringle, *Nat. Mater.*, 2022, **21**, 228–236, DOI: [10.1038/s41563-021-01130-z](#).
- 77 L. Jin, K. M. Nairn, C. D. Ling, H. Zhu, L. A. O'Dell, J. Li, F. Chen, A. F. Pavan, L. A. Madsen, P. C. Howlett, D. R. MacFarlane, M. Forsyth and J. M. Pringle, *J. Phys. Chem. B*, 2017, **121**, 5439–5446, DOI: [10.1021/acs.jpcc.7b02780](#).
- 78 N. Goujon, R. Kerr, C. Gervillie, Y. V. Oza, L. A. O'Dell, P. C. Howlett and M. Forsyth, *ACS Omega*, 2020, **5**, 2931–2938, DOI: [10.1021/acsomega.9b03773](#).
- 79 G. B. Chavhan, P. S. Babyn, B. Thomas, M. M. Shroff and E. M. Haacke, *Radiographics*, 2009, **29**, 1433–1449, DOI: [10.1148/rg.295095034](#).
- 80 K. Romanenko, *Annu. Rep. NMR Spectrosc.*, 2017, 67–114, DOI: [10.1016/bs.arnmr.2016.12.002](#).
- 81 M. Forsyth, F. Chen, L. A. O'Dell and K. Romanenko, *Solid State Ionics*, 2016, **288**, 160–166, DOI: [10.1016/j.ssi.2015.11.025](#).
- 82 K. Romanenko, J. M. Pringle, L. A. O'Dell and M. Forsyth, *Phys. Chem. Chem. Phys.*, 2015, **17**, 18991–19000, DOI: [10.1039/C5CP03294F](#).
- 83 M. Casalegno, F. Castiglione, G. Raos, G. B. Appetecchi, S. Passerini, A. Mele and E. Ragg, *ACS Appl. Mater. Interfaces*, 2020, **12**, 23800–23811, DOI: [10.1021/acsami.0c01890](#).
- 84 J. Adebahr, F. C. Grozema, S. W. deLeeuw, D. R. MacFarlane and M. Forsyth, *Solid State Ionics*, 2006, **177**, 2845–2850, DOI: [10.1016/j.ssi.2006.07.061](#).
- 85 L. Jin, S. de Leeuw, M. V. Koudriachova, J. M. Pringle, P. C. Howlett, F. Chen and M. Forsyth, *Phys. Chem. Chem. Phys.*, 2013, **15**, 19570–19574, DOI: [10.1039/C3CP53604A](#).
- 86 F. Chen, S. W. de Leeuw and M. Forsyth, *J. Phys. Chem. Lett.*, 2013, **4**, 4085–4089, DOI: [10.1021/jz402222j](#).
- 87 F. Chen, J. M. Pringle and M. Forsyth, *Chem. Mater.*, 2015, **27**, 2666–2672, DOI: [10.1021/acs.chemmater.5b00538](#).
- 88 W. A. Henderson, D. M. Seo, Q. Zhou, P. D. Boyle, J.-H. Shin, H. C. De Long, P. C. Trulove and S. Passerini, *Adv. Energy Mater.*, 2012, **2**, 1343–1350, DOI: [10.1002/aenm.201200130](#).
- 89 N. Sirigiri, F. Chen, C. M. Forsyth, R. Yunis, L. O'Dell, J. M. Pringle and M. Forsyth, *Mater. Today Phys.*, 2022, **22**, 100603, DOI: [10.1016/j.mtphys.2022.100603](#).
- 90 F. Chen and M. Forsyth, *Phys. Chem. Chem. Phys.*, 2016, **18**, 19336–19344, DOI: [10.1039/C6CP01411A](#).
- 91 P. C. Howlett, F. Ponzio, J. Fang, T. Lin, L. Jin, N. Iranipour and J. Efthimiadis, *Phys. Chem. Chem. Phys.*, 2013, **15**, 13784–13789, DOI: [10.1039/c3cp51986d](#).



