



Cite this: *CrystEngComm*, 2021, 23, 2491

Received 2nd February 2021,
Accepted 24th February 2021

DOI: 10.1039/d1ce00168j

rsc.li/crystengcomm

Recent advances in probing host–guest interactions with solid state nuclear magnetic resonance

Ashlea R. Hughes ^a and Frédéric Blanc ^{*ab}

An overview of the recent role of solid state nuclear magnetic resonance (NMR) spectroscopy in the field of supramolecular chemistry to probe host-guest interactions is provided. Over the last few years, solid state NMR methodologies have delivered unique insights into the atomic level structure and dynamics of guest molecules adsorbed in solid materials that are not available by other approaches. This *Highlight* discusses the range of NMR interactions that enable access to this information and provides a number of illustrating examples that highlight their applications in a wide range of chemical systems and porous materials covering metal organic frameworks, porous molecular solids and zeolites.

1. Introduction to host-guest interactions and supramolecular chemistry

Host-guest interactions are prevalent in nature and play important roles in many critical functions. These non-covalent interactions form supramolecular complexes between multiple molecules to occur by hydrogen bonding,

hydrophobic forces, van der Waals forces, π - π interactions and electrostatic effects and are responsible for the three-dimensional structure of biomacromolecules, molecular recognition processes, self-assembly, and many further examples which would not be able to perform their function without these crucial interactions.¹ These examples listed above have inspired chemists to exploit host-guest interactions for drug discovery, molecular separation and catalysis leading to the growing field of supramolecular chemistry. To the best of our knowledge, the first example of a synthetic compound capable of complexing alkali metal cations was the crown ether dibenzo-18-crown-6 (illustrated in Fig. 1), discovered by Charles J. Pedersen in 1967.² The

^a Department of Chemistry, University of Liverpool, Crown Street, Liverpool, L69 7ZD, UK. E-mail: frederic.blanc@liverpool.ac.uk

^b Stephenson Institute for Renewable Energy, University of Liverpool, Crown Street, Liverpool, L69 7ZD, UK



Ashlea Hughes

Ashlea Hughes obtained her MChem in Chemistry (with a year in industry) from the University of York (2014). Following a few years in industry, she returned to her studies in 2016 at the University of Liverpool to start a PhD in Frédéric Blanc's group focusing on the structure and dynamics of advanced materials by solid state NMR.



Frédéric Blanc

A portrait of Frédéric Blanc, a man with short dark hair and glasses, wearing a black and white checkered sweater over a dark collared shirt. He is smiling slightly.

Frédéric Blanc

Frédéric Blanc is currently Reader in the Department of Chemistry and Stephenson Institute for Renewable Energy at the University of Liverpool. He received a M.Sc, a MChem (2004) in Physical Chemistry and Catalysis from the Lyon School of Chemistry and a PhD (2007) from the University of Lyon. Following postdoctoral work and a number of Fellowships at the State University of New York at Stony Brook, NY and at the University of Cambridge, he joined the faculty at the University of Liverpool (2012). His current research focuses on the interplay of (MAS DNP) NMR with materials chemistry, in particular energy materials and supramolecular assemblies.

A range of experimental methods has been used to investigate host-guest systems. Fluorescence spectroscopy is widely exploited within this area, due to its ability to yield key information regarding the occurrence of complexation in the form of a spectral shift¹³ whilst contributing rather limited information about the strength or location of the host-guest interaction and is therefore often combined with other methods to achieve this. Calorimetry is able to provide the strength of the host-guest interaction,¹⁴

The two most common methods to explore host-guest interactions by solid state NMR methods exploit chemical shieldings (directly relating to the observed chemical shifts) and dipolar interactions, both of which are highly dependent on local structure and packing arrangements. The chemical shielding interaction can be decomposed into diamagnetic and paramagnetic shielding contributions. The diamagnetic

CrystEngComm, 2021, 23, 2491–2503 | 2493

Deuterium ^2H ($I = 1$) is a quadrupolar nucleus which lineshape is dominated by broadening arising from the quadrupolar interaction. ^2H NMR is a very common tool exploited for the study of molecular dynamics as its lineshape is highly sensitive to a range of different motions and their rates on the kHz–MHz timescale.⁴⁰ The low natural abundance of this isotope (0.0115%) is both a drawback with the need for isotopic enrichment requiring additional

^2H NMR has been used to explore the motion in molecular rotors in a range of systems^{45–47} and it has recently been identified that tubular covalent cages (TCC)⁴⁸ are the fastest exclusively organic molecular rotors found to date. The two TCCs that possess a central phenylene ring between either the trisubstituted aromatic (TCC2) or acetylene moieties (TCC3, Fig. 3(a))⁴⁹ have the potential to rotate in the solid state. Variable temperature ^2H static NMR lineshape analysis of TCCs (that have been deuterated on the phenylene ring) show a typical Pake pattern in the slow motion regime at 105 K and a temperature dependent lineshape that is typical of a fast rotational 180° flip upon heating. Through numerical simulations, the ^2H NMR lineshapes of TCC3-*R* were found to consistently have faster reorientation rates than its smaller counterpart TCC2-*R* and this is ascribed to facile rotation around the acetylene bonds due to

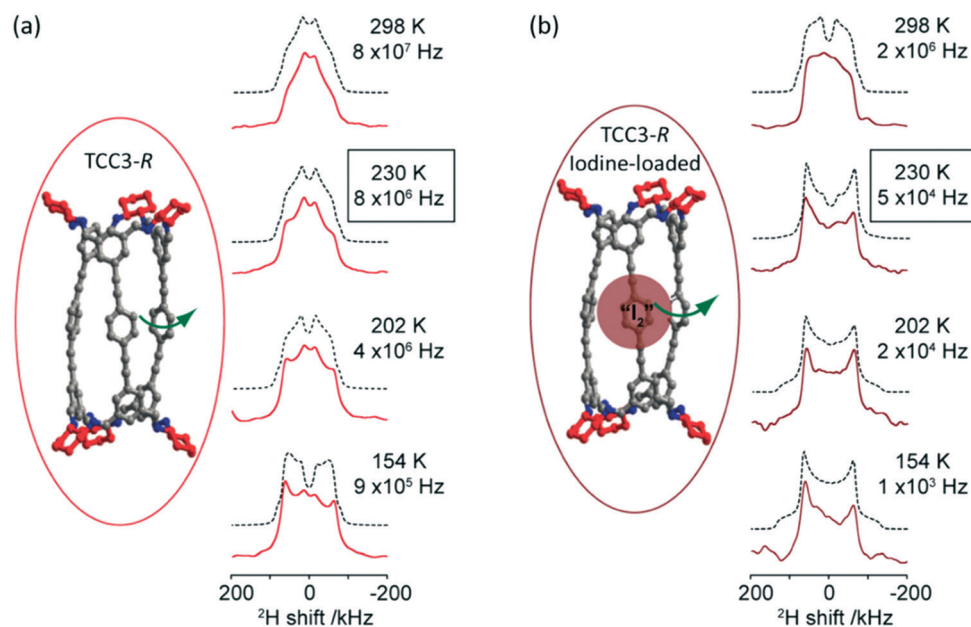
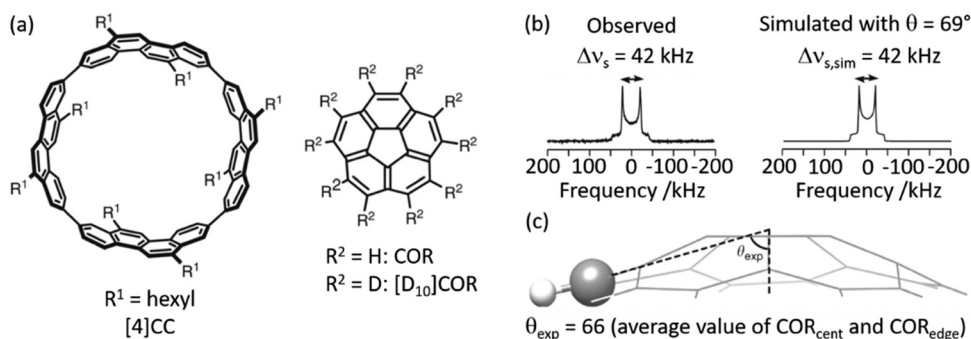


Fig. 3 (a) TCC3-*R* structure (red) and variable temperature ^2H static solid echo NMR spectra of $[\text{D}_{12}]\text{TCC3-}R$. (b) Iodine-loaded TCC3-*R* structure (burgundy) and variable temperature ^2H static solid echo NMR spectra of iodine loaded $[\text{D}_{12}]\text{TCC3-}R$. Corresponding simulated spectra (black dashed lines) and rotational rates obtained from numerical simulations of the NMR lineshape obtained at various temperatures are also given. In the crystal structures, the cyclohexane groups are represented in red; other C, grey; N, blue; H/D omitted for clarity. Ref. 48 is an open access article distributed under the terms of the Creative Commons CC BY license.

Guest capture of CO₂ is of particular interest in supramolecular chemistry^{45,51,52} and has recently been extensively investigated from the analysis of the ¹³C chemical shift anisotropy (CSA) lineshape.^{52,53} ¹³C labelled CO₂ is often used to focus on the loaded CO₂ whilst signals arising from the MOF framework linkers at natural abundance are highly decreased in intensity. One example compares two different metal centered (Al and Ga) MIL-53 metal organic frameworks

Static ^{13}C NMR spectra of $^{13}\text{CO}_2$ loaded $\text{Zn}_2(\text{dobdc})$ recorded as a function of increasing gradient field strengths



CrystEngComm, 2021, 23, 2491-2503 | 2495

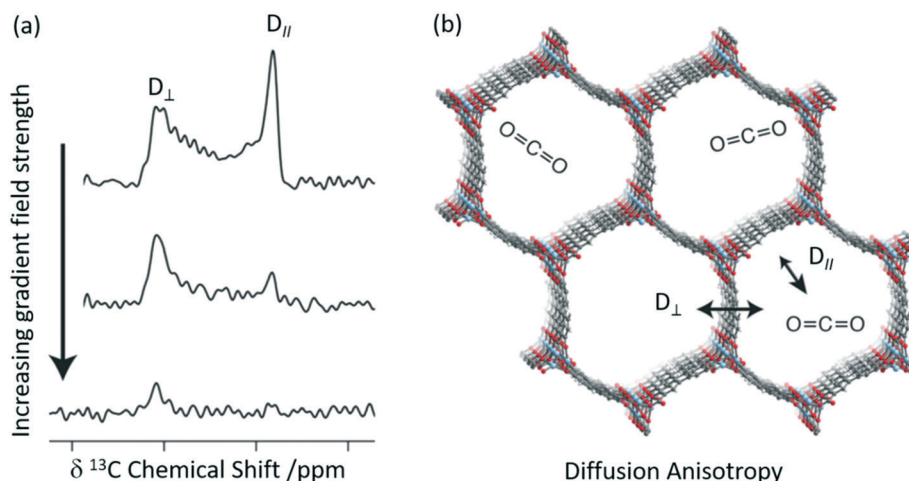


Fig. 5 (a) ^{13}C PFG static NMR spectra for $\text{Zn}_2(\text{dobpdc})$ crystals at a pressure of 625 mbar $^{13}\text{CO}_2$ with different applied gradient strengths and at 298 K. $D_{||}$ is the signal arising from crystals parallel to the external magnetic field at a lower chemical shift, whilst those at higher chemical shift are aligned perpendicular to the field $D_{\perp}$. (b) Cross-section of the crystal structure of the $\text{Zn}_2(\text{dobpdc})$ framework at 298 K showing the two diffusion pathways. Light blue, red, grey, and white spheres represent Zn, O, C, and H atoms, respectively. Reprinted (adapted) with permission from ref. 52. Copyright 2018 American Chemical Society.

(Fig. 5(a)) have revealed faster decay of the deshielded signal than the shielded signal and enabled CO_2 self-diffusion coefficients in both the parallel ($D_{||}$) and perpendicular ($D_{\perp}$) directions to the hexagonal channels to be obtained. CO_2 was observed to diffuse through the channels parallel to the crystallographic c axis ($D_{||}$) as expected but also and, more surprisingly, between the hexagonal channels in the crystallographic ab plane ($D_{\perp}$, Fig. 5(b)) and has been attributed to defects in the MOF structure. It is worth pointing out that neither *in situ* XRD or MD revealed this anisotropic diffusion thereby highlighting the complementarity of these methods.

3.3. NMR relaxometry

NMR relaxometry encompasses a wide range of experiments that access the rates of relaxation (either longitudinal relaxation times in the laboratory frame T_1 or rotating frame $T_{1\rho}$, or transverse relaxation time T_2) which measure polarisation returning to their equilibrium values in a magnetic field. Temperature dependent relaxometry measurements in particular are insightful for the study of molecular dynamics, allowing access to the correlation times of the motion and their activation energies, providing some important information about host-guest interactions.^{22,36}

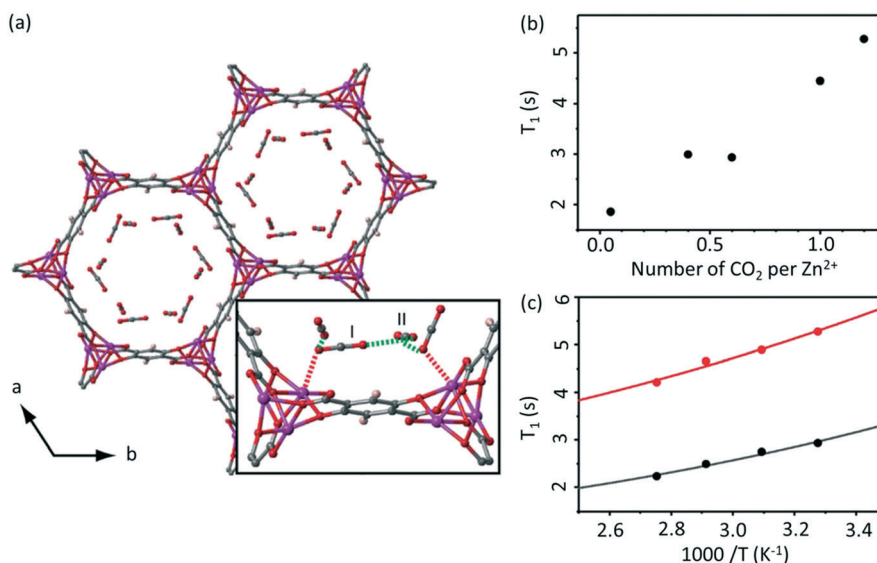


Fig. 6 (a) MOF-74 chemical structure with adsorption sites shown. (b) Pressure and (c) temperature dependence of T_1 values for the adsorbed CO_2 in $\text{Zn}_2(\text{dotp})$ MOF-74 under MAS at $\nu_r = 6$ kHz, 305 K and at 100 mbar (black) and 1000 mbar (red), respectively. Experimental and calculated T_1 values (based on a ^{13}C CSA relaxation mechanism) are shown as data points and solid lines, respectively. Reproduced from ref. 55 with permission from the PCCP Owner Societies.



The adsorption of ^{13}C labelled CO_2 in MOF-74 $\text{Zn}_2(\text{dotp})$ (where $\text{dotp}^{4-} = 2,5\text{-dioxidoterephthate}$, Fig. 6(a)) has recently been studied using a combination of solid state NMR techniques, including ^{13}C CSA lineshape analysis, 2D exchange spectroscopy, ^{13}C MAS and NMR relaxometry.⁵⁵ The T_1 times of the CO_2 guest were found to increase with increasing numbers of CO_2 molecules (Fig. 6(b)), indicating that at higher pressures the CO_2 mobility is hindered due to the $\text{CO}_2\text{-CO}_2$ interactions. The correlation times obtained from T_1 data over the 305–363 K temperature range and assuming a CSA relaxation mechanism were calculated and yielded activation energies of 4.4 and 3.5 kJ mol^{-1} for two samples with different CO_2 pressure of 100 and 1000 mbar, respectively (Fig. 6(c)). This slight difference supports the hypothesis that the $\text{CO}_2\text{-CO}_2$ and $\text{CO}_2\text{-pores}$ surface interactions increase with pressure. The correlation time found was also significantly longer than that of gaseous CO_2 , also providing the evidence for the host-guest interactions occurring at the pore surface.

Analysis of the T_1 data in $\text{Zn}_2(\text{dotp})$ MOF-74 also revealed the occurrence of unusually slow localized wobbling and hopping motions (at a rate of 10^{-8} s) that is much smaller than for the Mg-MOF-74 (10^{-10} s) analogue and further studies are highly anticipated to provide more information on the mobility, adsorption and separation properties of these MOFs.

3.4. Spin diffusion

Spin diffusion NMR experiments observe the transfer of nuclear magnetization *via* through space dipole-dipole interactions and has been used to probe host-guest interactions among a range of systems.³⁹ As spin diffusion coefficients, D , are proportional to the cube root of the nuclei concentration and the square of the nuclei gyromagnetic ratio, nuclei with sufficiently high values (*e.g.* ^1H , ^7Li , *etc.*) are often used in these experiments and provide qualitative and quantitative information on spatial proximities.

Along with T_1 relaxometry data on $\text{Zn}_2(\text{dotp})$ MOF-74, ^{13}C two-dimensional (2D) exchange MAS NMR data was also used to investigate the diffusion of CO_2 enriched in ^{13}C thereby increasing ^{13}C spins concentration.⁵⁵ With the possibility of multiple adsorption sites, the determination of the type of sites is crucial for the understanding of the mechanism of adsorption. Surprisingly and in contrast to previous work assessing the ^{13}C CSA patterns of Mg-MOF-74,⁵⁶ two CO_2 resonances were observed in Zn-MOF-74 (ref. 55) and initially assigned to a primary adsorption site and a secondary adsorption site only accessible under greater pressure. However, the corresponding 2D exchange ^{13}C MAS NMR spectra revealed that these two signals do not exchange (over 10–1000 ms mixing times) and it was therefore concluded that the second resonance is due to mobile CO_2 located in the dead space of the NMR tube rather than a secondary adsorption site.

^1H ^1H spin diffusion experiments were performed on $\text{Zr}_6\text{-O}_4(\text{OH})_4(\text{bpdc})_6$ MOF UiO-67 (where $\text{bpdc}^{2-} = \text{biphenyl-4,4'}$ -

dicarboxylate) loaded with a range of light alkanes (methane, ethane and propane) to explore their interactions into the MOF pores (Fig. 7(a)).⁵⁷ First principle calculations and neutron diffraction studies had previously been used to probe host-guest interaction in methane-loaded $\text{M}_2(\text{dhtp})$ MOFs ($\text{M} = \text{Mg}, \text{Mn}, \text{Co}, \text{Ni}, \text{Zn}$; $\text{dhtp}^{2-} = 2,5\text{-dihydroxyterephthalate}$),⁵⁸ however, due to the motional dynamics and disorder in these systems, significant direct experimental evidence for the interaction at room temperature has been difficult to achieve. The ^1H ^1H spin diffusion spectra of the methane loaded UiO-67 are shown in Fig. 7(b) at various mixing times and revealed cross peaks correlating methane (-0.4 ppm) with aromatics proton belonging to the UiO-67 bpdc organic linker, supporting van der Waals interactions. Experiments conducted with heavier alkanes (ethane and propane) showed similar results, albeit requiring longer spin diffusion times for cross peaks to appear and suggests there is a slight preference for the smaller alkanes, and overall showing evidence that UiO-67 could be utilised for the storage of these light alkanes.

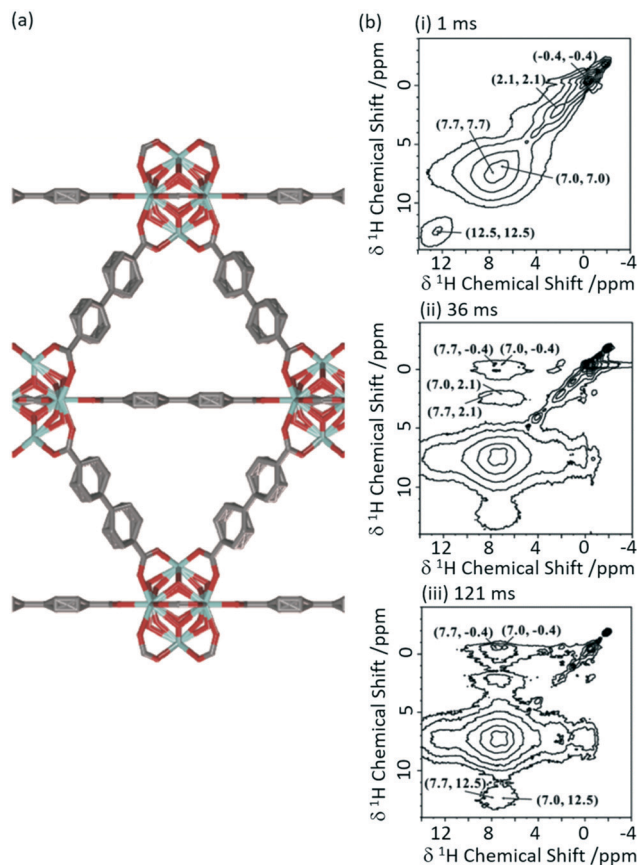


Fig. 7 (a) Schematic structure of UiO-67. (b) 2D ^1H - ^1H spin-diffusion spectra of methane loaded UiO-67 at mixing times of (i) 1 ms, (ii) 36 ms and (iii) 121 ms obtained at a MAS rate $\nu_r = 10$ kHz. Reprinted (adapted) with permission from ref. 57. Copyright 2017 American Chemical Society.



3.5. Xenon NMR

Xenon ^{129}Xe ($I = 1/2$) enables detailed information about the size and shape of cavity in porous solids⁵⁹ and is often used to probe supramolecular assemblies. This is largely due to the inert nature and large electron density of Xe which makes this nucleus particularly sensitive to the atomic scale.^{59,60} Although the high gyromagnetic ratio ($-7.441 \times 10^7 \text{ rad T}^{-1} \text{ s}^{-1}$) and natural abundance (26.4%) seem favourable for NMR detection, the reduced xenon-xenon interactions needed to promote observation of adsorption sites requires low concentration challenging detection. Hyperpolarised ^{129}Xe NMR is therefore used to perform experiments at lower xenon concentration with significant sensitivity. Importantly, changes in the chemical shifts from free gaseous xenon (at 0 ppm) indicate host-guest interactions and result from a combination of the xenon-surface interaction and the pore volume.⁵⁹

^{129}Xe has been extensively used to investigate the pore structure in solids which is largely key to the physical properties. An example of such material includes zeolites, which structural flexibility enabled by structure-directing agent leads to new catalytic applications.^{61,62} Recently, interlayer expansion of zeolites has been proposed to expand the current range of available zeolite topological structures and ^{129}Xe has been shown to determine the order of the interlayer structure in a new ten-membered ring (MR) functionalised skeleton zeolite (named COE-4) prepared by expansion of the 8-MR layered silicate RUB-36 zeolite precursor.⁶³ The ^{129}Xe spectrum of the calcined RUB-36 (RUB-37) phase shows a weak signal at 89 ppm at 213 K (Fig. 8(a)) and has been ascribed to accumulated Xe in the mesopores and not in the 8-MR channels (of size of $3.1 \times 4.7 \text{ \AA}$ and $2.5 \times 4.2 \text{ \AA}$ along the [010] and [001] directions, respectively) which are too small for Xe (dynamic diameter of $4.4 \text{ \AA}$). In COE-4, the ^{129}Xe signal appears at a larger shift of 111 ppm (Fig. 8(b)) indicating the xenon is more “trapped” and these are classed as micropores. Additionally, the ^{129}Xe lineshape of the single resonance observed in COE-4 (Fig. 8(b)) indicates high symmetry and a homogenous distribution of the interlayer expansion. It should also be noted that the ^{129}Xe

NMR spectra of both COE-4 and ZSM-5 zeolites are very similar, indicating that the pore sizes are comparable.

Porous organic cages (POCs) are a relatively new family of supramolecular assemblies, with the above mentioned TCC3 being one of their newest members. CC3 is another POC consisting of tetrahedral cages formed by imine bonds connecting rigid aromatic rings to the more flexible cyclohexane linkers and arranges into an interconnected 3D pore structure when packed together (Fig. 9(a)). Since their discovery in 2009, POCs have been shown to possess wide applications and unprecedented performance in rare gas separation.⁶⁰ The cage structure of CC3 was initially studied by ^{129}Xe NMR and it was observed that the xenon can diffuse in three dimension throughout the crystalline pore structure, moving between the window cavities *via* the cage cavities and indicating interconnectivity of the pore structure.⁶⁴ A follow up study⁶⁰ delved into the CC3 cage and its dynamics revealing only one resonance for both the cage and window sites due to fast thermal motion causing the exchange of nuclei between these sites to occur faster than the NMR timescale (Fig. 9(b)). Upon Xe loading, it was observed that the ^{129}Xe chemical shift increases to 210 ppm as the more shielded window sites of the cage (as opposed to the Xe nuclei located in the less shielded cage cavities which appear at 20 ppm) become more occupied. At a lower temperature of 260 K, the chemical shift is slightly reduced due to the increase in relative occupancy of the window cavities. ^{129}Xe spin lattice relaxation rates were also measured due to their dependence on dynamics. Slower motional correlation times were obtained with increasing Xe loading, suggesting restricted xenon diffusion at high loading due to the cage and window cavities being occupied.

3.6. Cross polarisation

CP is one of the most commonly used experiments in solid state NMR to transfer polarisation from abundant nuclei (typically ^1H) to sparse nuclei (*e.g.* ^{13}C , ^{15}N) in order to gain sensitivity. The polarisation transfer is largely dependent on

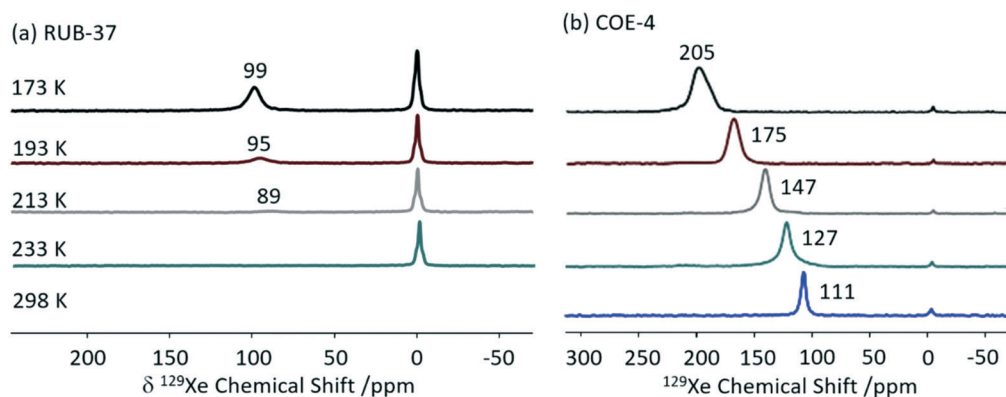


Fig. 8 Variable temperature hyperpolarised static ^{129}Xe NMR spectra of (a) RUB-37 and (b) COE-4. Reprinted from ref. 63: structural investigation of interlayer-expanded zeolite by hyperpolarized ^{129}Xe and ^1H NMR spectroscopy, 109555, Copyright 2019, with permission from Elsevier.



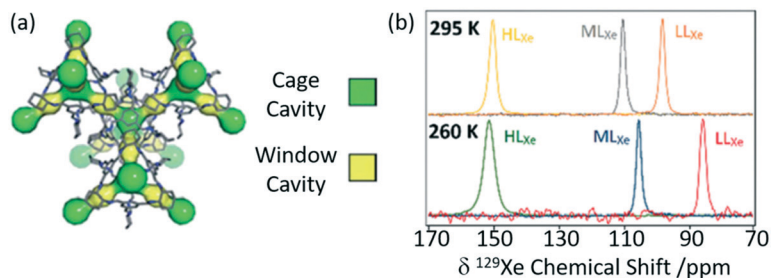


Fig. 9 (a) Crystal structure of CC3 showing both cage cavities (yellow) and window cavities (green). (b) ^{129}Xe NMR spectra of CC3 with varying xenon loadings (HL_{Xe} – high loading, ML_{Xe} – middle loading and LL_{Xe} – low loadings) at 295 and 260 K. Reproduced from ref. 60 with permission from the PCCP Owner Societies.

the through space heteronuclear dipolar interactions, and hence can also be exploited to probe dynamics and measure distances including the observation of internuclear interactions and informing host–guest binding. MIL-53 has also been studied with CP and, recently, ^1H ^{129}Xe CP has been utilised for the first time in this class of supramolecular assemblies to probe the mobility of xenon (Fig. 10).⁶⁵

No ^{129}Xe signal is present in large-pore MIL-53 due to the high mobility completely averaging the ^1H ^{129}Xe heteronuclear dipolar interactions, and hence significantly reducing CP efficiency. The signal intensity obtained for the ^{129}Xe located in the narrow-pore MIL-53 was found to increase with CP contact time until a maximum at 10 ms which is significantly longer than in the absence of motion (2 ms), hence concluding that the xenon atoms still possess significant motions in these pores.

Zeolite H-ZSM-5 catalysed methanol-to-hydrocarbon (MTH) process remains a highly studied industrial reaction due to its mechanistic aspects and complexity.⁶⁶ Solid state NMR has already identified a range of catalytic intermediates aiding the

mechanistic understandings.⁶⁷ Recently, it has also been reported that Lewis acidity is not a spectator in this MTH reaction and that the incorporation of alkaline-earth metals into the zeolite results in a generation of Lewis acid site changing the reactivity of the zeolite catalyst throughout the MTH process.⁶⁸ Using 1D ^1H ^{13}C CP MAS experiments performed at increasing contact times, the post reacted H-ZSM-5 and Ca-ZSM-5 zeolites have been analysed to study diffusion characteristics (Fig. 11).⁶⁹ The peak at 100 ppm is assigned as an acetal (OCH_2O) and its intensity increases with increasing contact time, indicating a more efficient CP transfer at longer contact times, suggesting that acetals are preferentially located within the zeolite framework. Direct spectral comparison of both zeolites shows large signal intensity in the 120–135 ppm region for H-ZSM-5 indicating a higher concentration of unsaturated HCP species (like aromatics/polyaromatics) on its surface in comparison to Ca-ZSM-5 (Fig. 11(a)). Upon increasing the CP contact times to 3 ms, the aromatic peak around 133 ppm in the H-ZSM-5 spectra (Fig. 11(b)(i)) disappears, implying the presence of polyaromatic products exclusively residing on the surface of the zeolite. However, this weak peak at 133 ppm is absent in the post-reacted Ca-ZSM-5 zeolite (Fig. 11(b)(ii)) at short contact times, which indicates the absence of any aromatic species in the calcium modified zeolite, especially on its surface.

3.7. Rotational echo double resonance

Driven by the need for high resolution, MAS is used extensively to average out heteronuclear dipolar interactions to zero, however, these interactions provide access to spatial proximities and hence probe host–guest interactions. Their reintroduction, whilst under MAS to still achieve the desired spectral resolution, is often achieved using recoupling pulse sequences so that the dipolar coupling is not averaged out to zero under MAS of which the rotational echo double resonance (REDOR)⁷¹ is the archetype. Spin echo experiments with (S_0) and without (S') refocusing pulses are acquired as a function of evolution time and enable access to (recoupled) heteronuclear dipolar couplings and interatomic distances. In particular, REDOR type experiments are prevalent in the study of zeolites due to their catalytic performance often being linked to host–guest interactions.

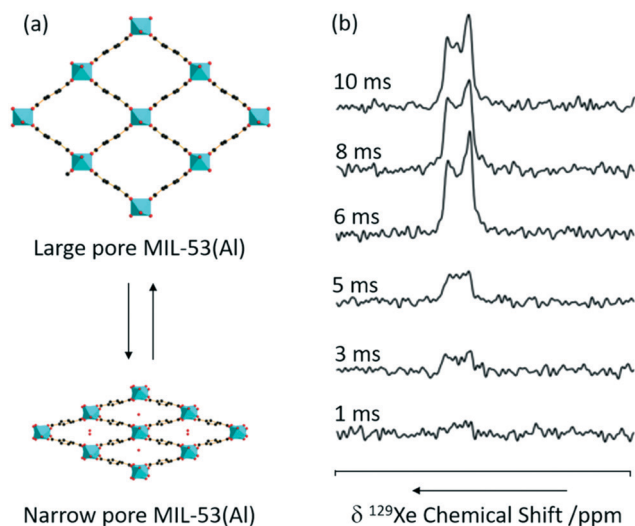


Fig. 10 (a) 3D structure of MIL-53(Al) in exchange between the large and narrow pore form. (b) ^1H ^{129}Xe CP spectra as a function of contact time under a MAS rate $\nu_r = 8$ kHz at a temperature of 153 K and a pressure of 500 mbar. Reprinted (adapted) with permission from ref. 65. Copyright 2016 American Chemical Society.



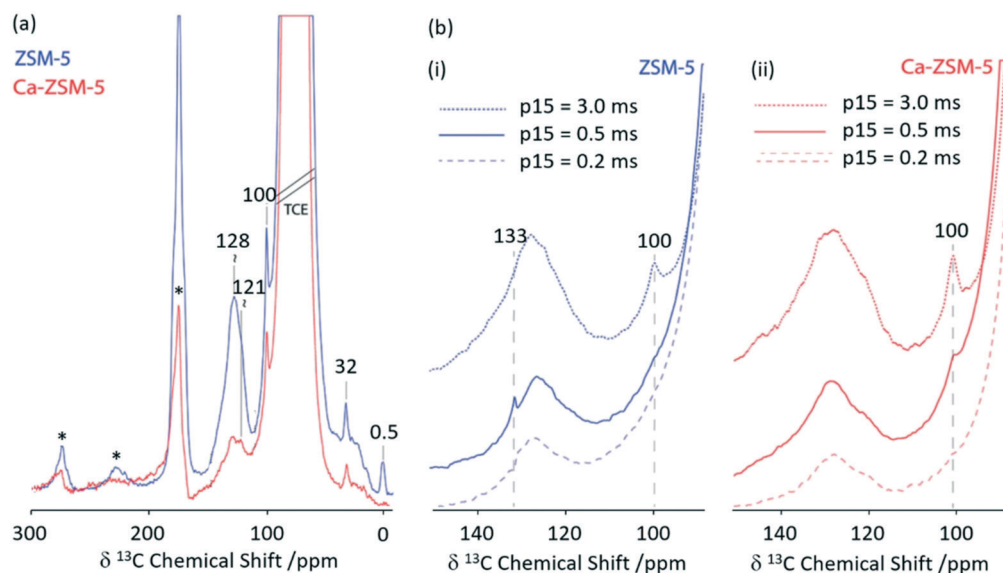


Fig. 11 (a) 1D ^1H ^{13}C CP MAS DNP spectra of post-MTH reacted H-ZSM-5 (blue) and Ca-ZSM-5 (red) at a contact time of 2 ms. (b) 1D ^1H ^{13}C CP MAS DNP spectra at various CP contact times (p15) of (i) H-ZSM-5 (blue) and (ii) Ca-ZSM-5 (red). Samples have been formulated in a 16 mM TEKPol solution⁷⁰ in 1,1,2,2-tetrachloroethane. Data were obtained with a recycle delay of 3 s and a MAS rate $\nu_r = 8$ kHz. Ref. 69 – published by the Royal Society of Chemistry.

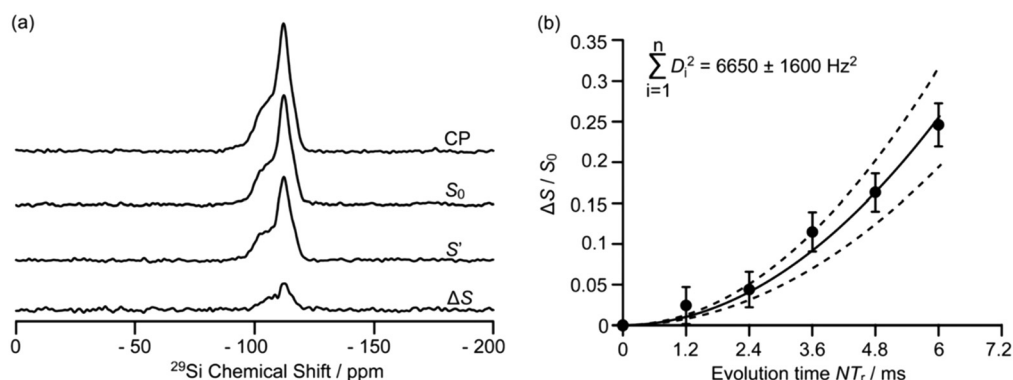


Fig. 12 (a) ^{29}Si CP, ^{29}Si CP spin echo (S_0), and ^{29}Si ^{13}C REDOR signal reintroducing dipolar couplings (S') and the difference ΔS signal ($S_0 - S'$) at a MAS rate of $\nu_r = 5$ kHz for H-ZSM-5 zeolite. (b) Plot of the REDOR fraction $S_0 - S'/S_0$ as a function of evolution time N/ν_r and corresponding best-fit curve (black line) and fit boundaries (dashed lines). Ref. 72 – published by the Royal Society of Chemistry.

Fig. 12 shows such ^{29}Si – ^{13}C REDOR experiments to probe the confined hydrocarbon species in the H-ZSM-5 zeolite framework.⁷² Fitted REDOR fraction (Fig. 12(b)) returns a ^{29}Si – ^{13}C heteronuclear coupling of 80 Hz and an internuclear distance of 4.2 Å which highlights the strong hydrocarbon – zeolite interaction and provide insights into zeolite poisoning by blockage of the channels rather than the acid sites.

An adaptation of the REDOR experiment to measure dipolar interactions between spin 1/2 and quadrupolar nuclei was introduced in 2010 as symmetry-based resonance-echo saturation-pulse double-resonance (S-RESPDOR)⁷³ and recently used to understand spatial proximities between ^{13}C in the MTH hydrocarbon pools and ^{27}Al in zeolites having different topologies (e.g. H-ZSM-5, HSSZ-13 and H-MOR).⁷⁴ Comparison of both S_0 and S' RESPDOR signals (Fig. 13)

indicate the spatial proximities of methylbenzenes (17.0 and 19.3 ppm), cyclopentyl cations (25.2, 45.8 and 48.1 ppm), surface bound methanol (58.7 ppm) and DME (60.0 and 63.4 ppm) with a preference for the latter.

Variable temperature experiments in the 300–400 K range showed minimal change in S_0 and S' spectra for the cyclopentyl carbons suggesting the retention of close spatial proximity even at higher temperatures. This is in sharp contrast to the methylbenzene signals where ΔS becomes 0 at higher temperatures due to a reduction in the spatial interaction of methylbenzene enabling the guest to have more mobility within the framework. Comparison between three topologies indicate no significant difference for the cyclic carbocations, however, methylbenzenes present stronger interactions with H-SSZ-13 and H-MOR implying



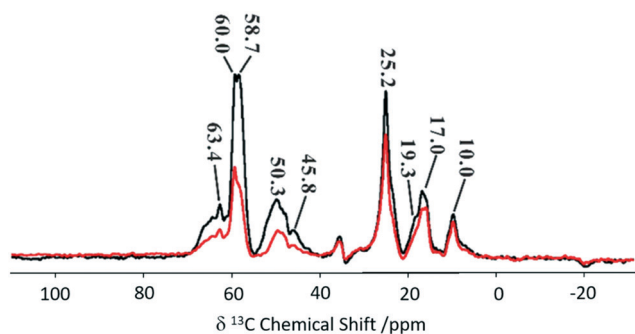


Fig. 13 ^{13}C spectra observed with (S , black) and without (S_0 , red) $\{^{27}\text{Al}\}$ S-RESPDOR for the trapped products retained in H-ZSM-5 after reaction with methanol for 15 minutes at 573 K. NMR spectra obtained at a MAS rate of $\nu_r = 10$ kHz. Reprinted (adapted) with permission from ref. 74. Copyright 2017 American Chemical Society.

that these two zeolites may have a range of reactivity in the methanol to olefin reaction.

3.8. Separated local field NMR

Separated local field (SLF) NMR experiments correlate heteronuclear dipolar couplings with chemical shifts in two dimensional NMR spectra and provide orientation and distance dependent restraints to aid structure determination, accessing dynamics and probing host-guest interactions. These experiments were initially designed for the study of liquid crystals^{75,76} and now find widespread use including in supramolecular assemblies. A particular version of SLF is the

dipolar chemical shift correlation (DipShift) which uses a principle similar to REDOR that recouple heteronuclear dipolar interactions with refocusing π pulses.¹³ ^{13}C - ^1H DipShift NMR was used to obtain the magnitude of the motionally averaged one-bond ^{13}C ^1H dipolar couplings and values of -11.1 and -8.7 kHz were obtained for the aromatic carbons of styrene and ethylbenzene loaded on MIL-53(Al), respectively (Fig. 14).⁷⁷ These values are much slower than the ones expected in the absence of motion (-23 kHz) and determined from CH distances which indicate significant dynamic processes. The lower value obtained for ethylbenzene indicates higher mobility in the MIL-53 pore than for styrene and suggests that this latter aromatic has a stronger host-guest interaction with this MOF. This work has provided direct insights into the structure-selectivity relationship in MIL-53 and has contributed to further understanding of the exceptional property of this MOF to separate both aromatics which is a challenging separation.

4. Conclusion

Solid state NMR approaches that provide detailed insights into host-guest chemistry in supramolecular assemblies have been described. These studies enable understanding of the role of dynamics, binding sites and spatial interactions and often complement other approaches *e.g.* diffraction-based methods and computational investigations, to deliver comprehensive knowledge. The NMR technology plays an important role in the design and development of materials with new and improved adsorption phenomena for a wide range of applications including separation, catalysis and energy conversion and storage.

Conflicts of interest

There are no conflicts of interest to declare.

Acknowledgements

Financial support from the Engineering and Physical Sciences Research Council (EPSRC) for a Doctoral Training Studentship to A. R. H. is acknowledged.

References

- 1 J. W. Steed and J. L. Atwood, *Supramolecular Chemistry*, John Wiley & Sons, Chichester, UK, 2nd edn, 2009.
- 2 C. J. Pedersen, *J. Am. Chem. Soc.*, 1967, **89**, 2495–2496.
- 3 J.-M. Lehn, *Angew. Chem., Int. Ed. Engl.*, 1988, **27**, 89–112.
- 4 C. J. Pedersen, *Angew. Chem., Int. Ed. Engl.*, 1988, **27**, 1021–1027.
- 5 D. J. Cram, *Angew. Chem., Int. Ed. Engl.*, 1988, **27**, 1009–1112.
- 6 J. P. Sauvage, *Angew. Chem., Int. Ed.*, 2017, **56**, 11080–11093.
- 7 J. F. Stoddart, *Angew. Chem., Int. Ed.*, 2017, **56**, 11094–11125.
- 8 B. L. Feringa, *Angew. Chem., Int. Ed.*, 2017, **56**, 11060–11078.
- 9 Y. Xu, S. A. Southern, P. M. J. Szell and D. L. Bryce, *CrystEngComm*, 2016, **18**, 5236–5252.

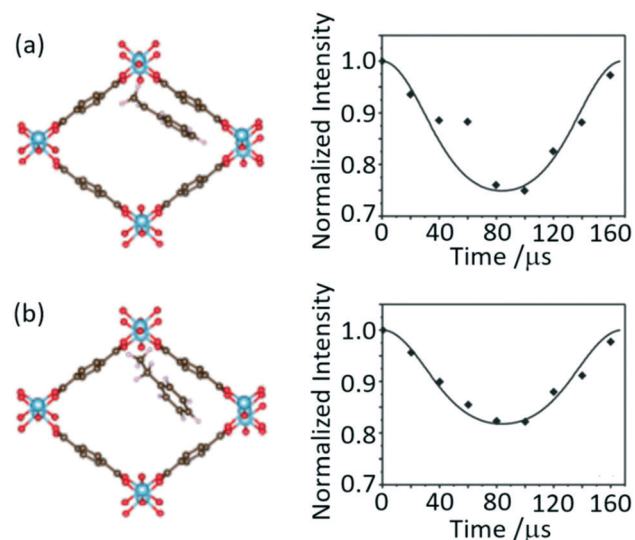


Fig. 14 Schematic of MIL-53(Al) upon adsorption of (a) styrene and (b) ethylbenzene and their corresponding dipolar dephasing curves extracted from ^{13}C - ^1H DipShift MAS NMR dataset obtained at a MAS rate of $\nu_r = 6$ kHz. The CH sites at 128 and 127 ppm were chosen to capture the dynamic in styrene and ethylbenzene, respectively. Reprinted from ref. 77: host-guest interaction of styrene and ethylbenzene in MIL-53 studied by solid-state NMR, Copyright 2018, with permission from Elsevier.



- 10 M. R. Chierotti and R. Gobetto, *Solid-State NMR Studies on Supramolecular Chemistry*, Wiley, Chichester, UK, 2012.
- 11 M. J. Duer, R. K. Harris and R. E. Wasylshen, *NMR Crystallography*, Wiley, Chichester, 2009.
- 12 M. J. Potrzebowski and S. Kazmierski, *Top. Curr. Chem.*, 2004, **246**, 91–140.
- 13 M. Novo, D. Granadero and J. B. W. Al-soufi, *J. Inclusion Phenom. Macrocyclic Chem.*, 2011, **70**, 259–268.
- 14 F. P. Schmidtchen, in *Analytical Methods in Supramolecular Chemistry*, Wiley, Weinheim, 2nd edn, 2012, pp. 67–104.
- 15 K. Jie, M. Liu, Y. Zhou, M. A. Little, A. Pulido, S. Y. Chong, A. Stephenson, A. R. Hughes, F. Sakakibara, T. Ogoshi, F. Blanc, G. M. Day, F. Huang and A. I. Cooper, *J. Am. Chem. Soc.*, 2018, **140**, 6921–6930.
- 16 S. Halder, A. Barma, C. Rizzoli, P. Ghosh and P. Roy, *ACS Appl. Nano Mater.*, 2019, **2**, 5469–5474.
- 17 S. E. Ashbrook and D. M. Dawson, *Acc. Chem. Res.*, 2013, **46**, 1964–1974.
- 18 J. A. Ripmeester and R. E. Wasylshen, *CrystEngComm*, 2013, **15**, 8598.
- 19 B. Elena and L. Emsley, *J. Am. Chem. Soc.*, 2005, **127**, 9140–9146.
- 20 D. C. Apperley, R. K. Harris and P. Hodgkinson, *Solid-State NMR: Basic Principles and Practice*, Momentum Press, New York, 1st edn, 2012.
- 21 R. W. Martin, J. E. Kelly and J. I. Kelz, *J. Struct. Biol.*, 2019, **206**, 73–89.
- 22 P. Hodgkinson and S. Wimperis, *Phys. Chem. Chem. Phys.*, 2009, **11**, 6875.
- 23 A. W. Overhauser, *Phys. Rev.*, 1953, **92**, 411–415.
- 24 T. R. Carver and C. P. Slichter, *Phys. Rev.*, 1956, **102**, 975–981.
- 25 A. S. L. Thankamony, J. J. Wittmann, M. Kaushik and B. Corzilius, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2017, **102–103**, 120–195.
- 26 A. Pines, M. G. Gibby and J. S. Waugh, *J. Chem. Phys.*, 1973, **59**, 569.
- 27 E. R. Andrew, A. Bradbury and R. G. Eades, *Nature*, 1958, **182**, 1659.
- 28 A. Brinkmann and M. H. Levitt, *J. Chem. Phys.*, 2001, **115**, 357–384.
- 29 S. E. Ashbrook, R. G. Griffin and K. E. Johnston, *Annu. Rev. Anal. Chem.*, 2018, **11**, 485–508.
- 30 A. D. Buckingham, *Can. J. Chem.*, 1960, **38**, 300–307.
- 31 M. P. Williamson, *Prog. Nucl. Magn. Reson. Spectrosc.*, 2013, **73**, 1–16.
- 32 D. H. Brouwer, S. Alavi and J. A. Ripmeester, *Phys. Chem. Chem. Phys.*, 2008, **10**, 3857–3860.
- 33 N. Asakawa, S. Kuroki, H. Kurosu, I. Ando, A. Shoji and T. Ozaki, *J. Am. Chem. Soc.*, 1992, **114**, 3261–3265.
- 34 A. Newman, *Pharmaceutical Amorphous Solid Dispersions*, John Wiley & Sons, Hoboken, New Jersey, 1st edn, 2015.
- 35 M. Teodorescu and M. Bercea, *Polym.-Plast. Technol. Eng.*, 2015, **54**, 923–943.
- 36 K. Ueda, K. Higashi and K. Moribe, *Mol. Pharmaceutics*, 2016, **13**, 852–862.
- 37 D. C. Apperley, A. H. Forster, R. Fournier, R. K. Harris, P. Hodgkinson, R. W. Lancaster and T. Rades, *Magn. Reson. Chem.*, 2005, **43**, 881–892.
- 38 X. Yuan, D. Sperger and E. J. Munson, *Mol. Pharmaceutics*, 2014, **11**, 329–337.
- 39 S. P. Brown, *Solid State Nucl. Magn. Reson.*, 2012, **41**, 1–27.
- 40 L. A. O'Dell and C. I. Ratcliffe, in *NMR of Quadrupolar Nuclei in Solid Materials*, John Wiley & Sons, Chichester, 1st edn, 2012, pp. 213–232.
- 41 L. S. Batchelder, in *Encyclopedia of Nuclear Magnetic Resonance*, ed. R. K. Harris and R. E. Wasylshen, John Wiley & Sons, Chichester, 1st edn, 2007.
- 42 R. L. Vold and G. L. Hoatson, *J. Magn. Reson.*, 2009, **198**, 57–72.
- 43 R. R. Vold and R. L. Vold, *Advances in Magnetic and Optical Resonance*, Academic Press, San Diego, 1991.
- 44 V. Macho, L. Brombacher and H. W. Spiess, *Appl. Magn. Reson.*, 2001, **20**, 405–432.
- 45 S. Bracco, F. Castiglioni, A. Comotti, S. Galli, M. Negroni, A. Maspero and P. Sozzani, *Chem. – Eur. J.*, 2017, **23**, 11210–11215.
- 46 S. L. Gould, D. Tranchemontagne, O. M. Yaghi and M. A. Garcia-Garibay, *J. Am. Chem. Soc.*, 2008, **130**, 3246–3247.
- 47 S. Bracco, A. Comotti, P. Valsesia, B. F. Chmelka and P. Sozzani, *Chem. Commun.*, 2008, 4798–4800.
- 48 A. R. Hughes, N. J. Brownbill, R. C. Lalek, M. E. Briggs, A. G. Slater, A. I. Cooper and F. Blanc, *Chem. – Eur. J.*, 2017, **23**, 17217–17221.
- 49 A. G. Slater, M. A. Little, A. Pulido, S. Y. Chong, D. Holden, L. Chen, C. Morgan, X. Wu, G. Cheng, R. Clowes, M. E. Briggs, T. Hasell, K. E. Jelfs, G. M. Day and A. I. Cooper, *Nat. Chem.*, 2017, **9**, 17–25.
- 50 T. Matsuno, M. Fujita, K. Fukunaga, S. Sato and H. Isobe, *Nat. Commun.*, 2018, **9**, 3779.
- 51 V. J. Witherspoon, J. Xu and J. A. Reimer, *Chem. Rev.*, 2018, **118**, 10033–10048.
- 52 A. C. Forse, M. I. Gonzalez, R. L. Siegelman, V. J. Witherspoon, S. Jawahery, R. Mercado, P. J. Milner, J. D. Martell, B. Smit, B. Blümich, J. R. Long and J. A. Reimer, *J. Am. Chem. Soc.*, 2018, **140**, 1663–1673.
- 53 R. M. Marti, J. D. Howe, C. R. Morelock, M. S. Conradi, K. S. Walton, D. S. Sholl and S. E. Hayes, *J. Phys. Chem. C*, 2017, **121**, 25778–25787.
- 54 Y. Zhang, B. E. G. Lucier and Y. Huang, *Phys. Chem. Chem. Phys.*, 2016, **18**, 8327–8341.
- 55 M. Inukai, T. Kurihara, Y. Noda, W. Jiang, K. Takegoshi, N. Ogiwara, H. Kitagawa and K. Nakamura, *Phys. Chem. Chem. Phys.*, 2020, **22**, 14465–14470.
- 56 L. C. Lin, J. Kim, X. Kong, E. Scott, T. M. McDonald, J. R. Long, J. A. Reimer and B. Smit, *Angew. Chem., Int. Ed.*, 2013, **52**, 4410–4413.
- 57 J. Li, S. Li, A. Zheng, X. Liu, N. Yu and F. Deng, *J. Phys. Chem. C*, 2017, **121**, 14261–14268.
- 58 H. Wu, W. Zhou and T. Yildirim, *J. Am. Chem. Soc.*, 2009, **131**, 4995–5000.
- 59 E. Weiland, M. A. Springuel-Huet, A. Nossor and A. Gédéon, *Microporous Mesoporous Mater.*, 2016, **225**, 41–65.



- 60 P. Håkansson, M. A. Javed, S. Komulainen, L. Chen, D. Holden, T. Hasell, A. Cooper, P. Lantto and V. V. Telkki, *Phys. Chem. Chem. Phys.*, 2019, **21**, 24373–24382.
- 61 J. Li, A. Corma and J. Yu, *Chem. Soc. Rev.*, 2015, **44**, 7112–7127.
- 62 J. Xu, Q. Wang, S. Li and F. Deng, *Solid-State NMR in Zeolite Catalysis*, Springer Nature, Singapore, 1st edn, 2019.
- 63 Z. Zhao, X. Li, S. Li, S. Xu, X. Bao, Y. Bilge, P. Andrei-Nicolae, M. Ulrich and W. Zhang, *Microporous Mesoporous Mater.*, 2019, **288**, 109555.
- 64 S. Komulainen, J. Roukala, V. V. Zhivonitko, M. A. Javed, L. Chen, D. Holden, T. Hasell, A. Cooper, P. Lantto and V. V. Telkki, *Chem. Sci.*, 2017, **8**, 5721–5727.
- 65 R. Giovine, C. Volkringer, M. A. Springuel-Huet, A. Nossov, F. Blanc, J. Trébosc, T. Loiseau, J. P. Amoureux, O. Lafon and F. Pourpoint, *J. Phys. Chem. C*, 2017, **121**, 19262–19268.
- 66 P. Tian, Y. Wei, M. Ye and Z. Liu, *ACS Catal.*, 2015, **5**, 1922–1938.
- 67 I. Yarulina, A. D. Chowdhury, F. Meirer, B. M. Weckhuysen and J. Gascon, *Nat. Catal.*, 2018, **1**, 398–411.
- 68 I. Yarulina, S. Bailleul, A. Pustovarenko, J. R. Martinez, K. De Wispelaere, J. Hajek, B. M. Weckhuysen, K. Houben, M. Baldus, V. Van Speybroeck, F. Kapteijn and J. Gascon, *ChemCatChem*, 2016, **8**, 3057–3063.
- 69 A. D. Chowdhury, I. Yarulina, E. Abou-hamad, A. Gurinov and J. Gascon, *Chem. Sci.*, 2019, **10**, 8946–8954.
- 70 A. Zagdoun, G. Casano, O. Ouari, M. Schwarzwälder, A. J. Rossini, F. Aussenac, M. Yulikov, G. Jeschke, C. Copéret, A. Lesage, P. Tordo and L. Emsley, *J. Am. Chem. Soc.*, 2013, **135**, 12790–12797.
- 71 T. Gullion and J. Schaefer, *J. Magn. Reson.*, 1989, **81**, 196–200.
- 72 D. Xiao, S. Xu, X. Han, X. Bao, Z. Liu and F. Blanc, *Chem. Sci.*, 2017, **8**, 8309–8314.
- 73 L. Chen, Q. Wang, B. Hu, O. Lafon, J. Trébosc, F. Deng and J. P. Amoureux, *Phys. Chem. Chem. Phys.*, 2010, **12**, 9395–9405.
- 74 C. Wang, J. Xu, Q. Wang, X. Zhou, G. Qi, N. Feng, X. Liu, X. Meng, F. Xiao and F. Deng, *ACS Catal.*, 2017, **7**, 6094–6103.
- 75 V. Domenici, *Liq. Cryst. Today*, 2017, **26**, 2–10.
- 76 S. V. Dvinskikh, in *Thermotropic Liquid Crystals*, ed. A. Ramamoorthy, Springer, Netherlands, Dordrecht, 2007, pp. 117–140.
- 77 S. Li, J. Li, J. Tang and F. Deng, *Solid State Nucl. Magn. Reson.*, 2018, **90**, 1–6.

