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Effect of a benzothiadiazole spacer on transport properties and N-doping of naphthalene-diimide-based copolymers†

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We report the synthesis of two n-type semiconducting polymers containing a naphthalene diimide (NDI) monomer and thiophene spacers. The two polymers differ by the introduction of an additional 2,1,3-benzothiadiazole (BTD) moiety in the polymer backbone. Although there are examples of polymers combining these two moieties, the effects of introducing the BTD as a third co-monomer on the transport properties and, in particular, the doping mechanisms of these materials have not been elucidated. We describe the optoelectronic, structural organisation and transport properties of these two polymers and investigate the effect of BTD introduction on molecular doping with N-DMBI as a dopant using UV-Vis spectroscopy, DFT modelling, cyclic voltametry, and X-ray diffraction. We show that the incorporation of a BTD unit in the polymer backbone not only improves the charge transport in organic field effect transistors but also the doping efficiency with N-DMBI as well as the conductivity and stability in the conducting state. Our results highlight the need for post-deposition annealing to optimise the N-DMBI doping of the BTD-based polymer and ultimately its conductivity and thermoelectric properties.

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Introduction

Organic semiconducting polymers have been extensively studied for more than three decades and are now emerging as the materials of choice for the development of the next generation of optoelectronic devices.[1–3] Their unique ability to combine the mechanical properties of polymers with greater processability, chemically tuneable optoelectronic properties, potentially higher sustainability[4] and lower production costs compared to inorganic materials has led to their adoption in many application areas.[5,6] Like their inorganic counterparts,

polymer-organic semiconductors can be divided into two categories: p-type materials, which primarily transport holes, and n-type materials, which primarily transport electrons.[7]

Historically, the development of n-type polymers has lagged far behind their p-type counterparts in many organic electronic devices such as organic photovoltaics (OPVs)[8] and organic field-effect transistors (OFETs)[9] and more recently in organic electrochemical transistors (OECTs)[10] and organic thermoelectric generators (OTEGs).[11] This is mainly due to the dominant use of n-type molecular semiconductors such as fullerenes and its derivates and to inherent instability of doped n-type materials under ambient conditions.[12] The synthesis of n-type polymers has gained momentum with the advent of the copolymer approach and the development of novel electron-rich (donor, D) and electron-deficient (acceptor, A) blocks, which allow the modulation of optoelectronic properties and energy levels to obtain more stable n-type materials.[13]

Among these classes of electron deficient moieties, naphthalene diimide (commonly abbreviated to NDI) is probably the most studied.[14,15] Aromatic diimide species are known to undergo two reversible one-electron reductions.[16] The resulting anions and dianions are indeed stabilised by a mesomeric effect. The electron density is mostly distributed over the oxygens of the two carbonyl groups.[17] In addition to its remarkable electrochemical stability, this moiety has an advantageous

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packing ability due to its planar and rigid structure. The solubility and therefore the processability of the NDI derivatives can be tuned by alkylation of the nitrogen atoms. Most of the NDI-based polymers reported so far have β-branched octyldodecyl chains, allowing them to be processed with environmentally friendly solvents.[18] Finally, NDI-based copolymers typically exhibit LUMO energy levels between −3.8 eV and −4.1 eV and some of them have shown remarkable air stability.[19,20] For all these reasons, NDI copolymers have emerged as promising materials for future industrial developments in energy conversion devices.[21,22] One example is their use in organic photovoltaic devices with power conversion efficiencies above 11%.[18] In the field of thermoelectricity, NDI-based polymers have shown good conductivity of 1.8 ± 0.1 S cm^{-1} leading to a power factor of 4.6 ± 0.2 μW m^{-1} K^{-2}.[23] They also show the lowest thermal conductivity values reported for n-type doped conjugated polymers.[24] There is a plethora of n-type materials based on NDI or BTD in the literature, but the combination of these two electron-deficient units within a single structure has rarely been explored.[15,21,22] Although there are examples of polymers combining these two units, the effects of introducing the BTD unit as a third co-monomer on the transport properties and, in particular, the doping mechanisms of these materials have not been elucidated. In this study, we report on the synthesis of two NDI-based polymers, one of which contains a BTD unit, and we investigate in detail their optoelectronic and transport properties. We investigate the effect of BTD on molecular doping using 4-(2,3-dihydro-1,3-dimethyl-1H-benzimidazol-2-yl)-N,N-dimethylbenzenamine (also known as N-DMBI) as a dopant. We show that the introduction of a BTD unit does not significantly affect the optoelectronic properties, but helps to improve the charge transport and doping efficiency with N-DMBI and, hence, the conductivity and thermoelectric properties.

Results and discussion

Polymers synthesis

Fig. 1 shows the synthetic routes to the synthons and the two NDI-based copolymers. The preparation of the copolymers requires the synthesis of a di-brominated NDI precursor **6** starting from naphthalene di-anhydride (Fig. S1, ESI†). This step is limiting since the electrophilic aromatic substitution is hindered by the electron deficiency of the substrate. In addition,
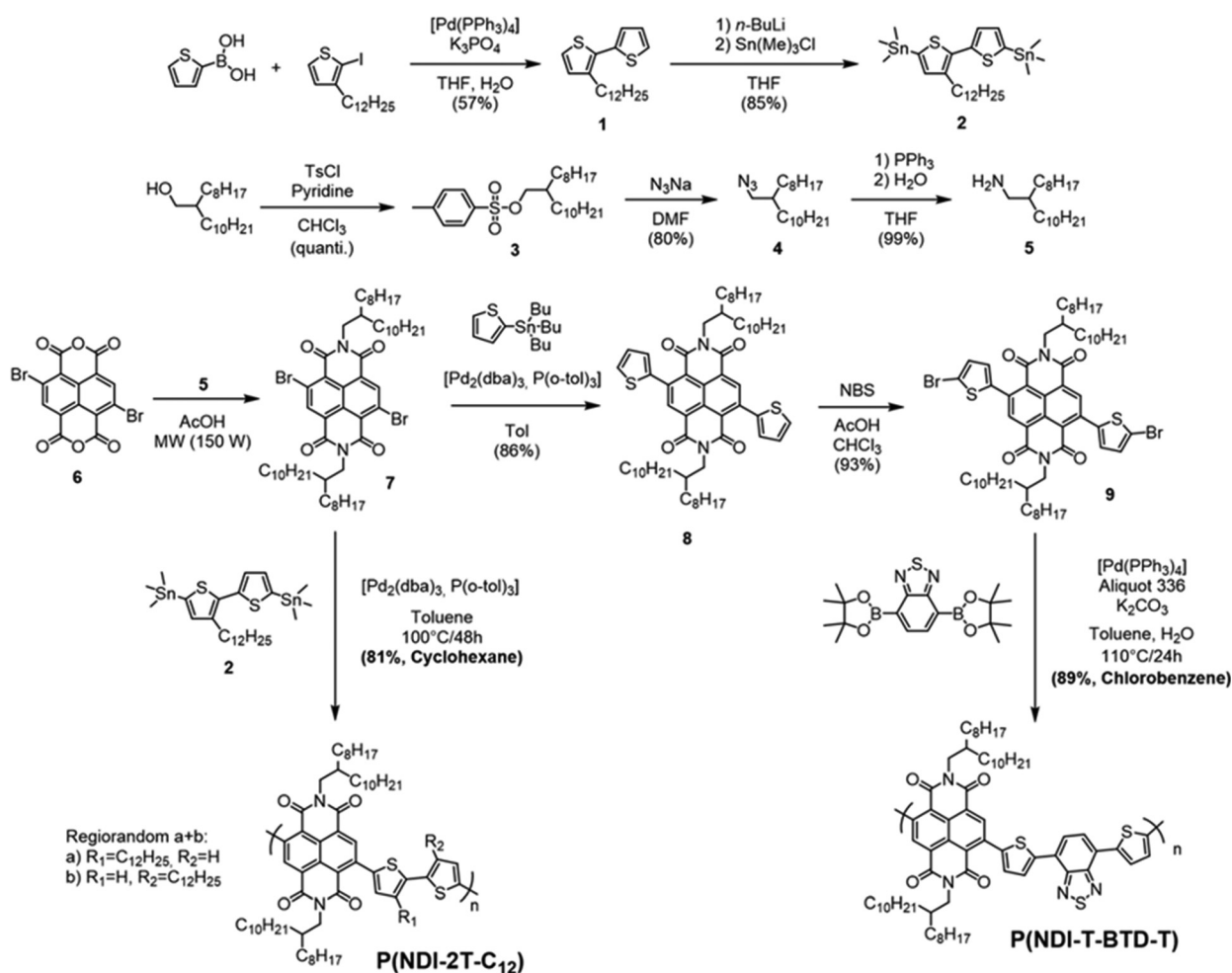

Fig. 1 Synthetic route towards the precursors and the NDI-based copolymers.

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the bromination product is insoluble in common solvents, making it difficult to purify and analyse. Despite the widespread use of NDI-based compounds, the synthesis of 2,6-dibromo-naphthalene dianhydride **6** is rarely reported. Our review of the literature shows a wide variation in reaction yields ranging from 12 to 42% (two steps, after *N*-alkylation) with the best yields described as poorly reproducible.[25] The standard bromination methods are: (i) oleum (fuming sulphuric acid) with Br_2, or (ii) concentrated sulphuric acid (H_2SO_4) with dibromo-isocyanuric acid as the brominating agent. Compared to Br_2, dibromo isocyanuric acid is expensive. As for oleum, it is very toxic by inhalation and reacts *via* a dangerous exothermic reaction (explosively with water). In an attempt to reduce the use of oleum and to find a low cost route to produce gram scale dibrominated naphthalene dianhydride, we undertook a systematic study inspired by the work of Reczek *et al.*[26]

Firstly, we set up the conditions for preparing compound **7** by synthesising an analogue bearing an *n*-octyl chain. The synthesis routes and total yields obtained for the preparation of the compound **7a** are summarised in Fig. S1 (ESI†). In a first step, three brominating agents were tested. The resulting crude products were then *N*-alkylated with commercial *n*-octylamine by microwave irradiation. The product **7a** can be conveniently purified by filtration to quickly obtain a soluble product for routine NMR analysis. Note that a highly protic solvent (glacial AcOH) is required to avoid nucleophilic aromatic substitution of bromine by alkylamine.[27] For reference, we obtained a total yield of 11% **7a** using dibromo-isocyanuric acid. We first tried using NaBr as a brominating agent,[28] which did not give any reaction in H_2SO_4, but gave 14% dibrominated NDI **7a** in oleum (150 °C/16 h) with significant amounts of non-functionalised and mono-brominated compound. At a lower temperature (110 °C), a lower yield was obtained even after 69 hours of reaction. To reduce the hazard, we focused on optimising the conditions at 100 °C with Br_2. Finally, we identified suitable Br_2/oleum conditions (0.15 mol L^{-1}, 100 °C/ 40 h) which gave pure **7a** with a good yield of 37%. These conditions were then used to synthesise the desired monomer **7** with octyl-dodecyl side chains using amine **4**. Prior to optimisation, residual brominated NDI was difficult to remove from the final product (Fig. S2, blue ^{1}H NMR trace, 8.2–8.9 ppm, ESI†). Optimisation of the bromination step enhances the formation of the di-brominated product and thus facilitates the purification of the monomer. With the optimised route, we were able to synthesise compound **7** on a gram scale with a satisfactory yield (33%) and excellent purity (Fig. S2, red trace, ESI†).

Starting from compound **7**, we synthesized the P(NDI-2T-C_{12}) with a regiorandom dodecyl side chain on only one thiophene, being an analogue of the archetypal P(NDI-2T) without side chains on the thiophene, also known commercially as N2200. Interestingly, P(NDI-2T-C_{12}) is highly soluble in common organic solvents and most of the formed copolymer was extracted from the cyclohexane fraction after Soxhlet purification. This fraction shows a moderate number average molar mass (Mn) of 8.6 kg mol^{-1} and a dispersity (*Đ*) of 1.9.

Starting from the compound **7**, we can easily access the compound **8** where two thiophenes are introduced *via* Stille

coupling.[29] This compound can then be brominated and used as a precursor for the preparation of P(NDI-T-BTD-T). In this copolymer, the thiophenes (abbreviated T) act as spacers allowing higher number average molar mass to be achieved. Indeed, low Mn of only 5.2 kg mol^{-1} was reached for P(NDI-BTD)[30] – without thiophene – while Fu and coll. managed to reach Mn up to 31.0 kg mol^{-1} by introducing a T-spacer to form P(NDI-T-BTD-T)[31] (both synthesised by Suzuki coupling). We tentatively attribute the higher degree of polymerisation to a higher reactivity of the donor thiophene unit. We further optimised the polymerisation by reproducing the conditions used by Fu and coll.[31] but with a shorter reaction time of 24 hours instead of 48 hours. We hence obtained P(NDI-T-BTD-T) with an excellent polymerisation yield of 89% (*vs.* 47%,[31] chloroben-zene fraction) and a higher Mn estimated to be 64.1 kg mol^{-1} (*vs.* 31.0 kg mol^{-1}). The dispersity (*Đ*) is 4.6 (see Fig. S11, ESI†). A higher Mn along with a higher polymerization yield, suggest that reducing the reaction time allows better control of the polymer chain length and prevents insoluble polymer aggregates. These differences may not only be due to the shorter reaction time, but also to the purity of the materials and the accuracy of the weighting according to Carothers's equation.[32]

Optical and electronic properties

The UV-Vis absorption spectra of the two polymers were recorded in solution and in thin films (Fig. 2). Spectral data are summarised in Table 1. In solution, P(NDI-2T-C_{12}) and P(NDI-T-BTD-T) exhibit a defined absorption band at high energy (~350 nm), which is attributed to the S_0–S_1 transition. For both D–A copolymers, a broad absorption band at lower energy (~600–785 nm) is observed and attributed to the internal charge transfer (ICT) absorption band between the donor

Fig. 2 UV-Vis absorption spectra of both polymers (a) in chlorobenzene solution at 298 K normalized to the maximum absorption peak, and (b) as as-cast thin films in ambient conditions.

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Table 1 Summary of the optoelectronic and transport properties

Copolymers	$\lambda_{\text{max\_sol}}$ [a] (nm)	$\lambda_{\text{max\_film}}$ [b] (nm)	λ_{onset} [c] (nm)	E_g^{opt} [d] (eV)	E_{LUMO} [e] (eV)	E_{HOMO} [f] (eV)	E_g^{CV} [g] (eV)	E_b [h] (eV)	$\mu_{\text{OFET}}^{\text{max}}$ [i] (cm^2 V^{-1} s^{-1})	σ [j] (S cm^{-1})
P(NDI-2T-C$_{12}$)	365, 605	390 (8.3×10^4) 725 (5.9×10^4)	890	1.4	−3.8 (−3.62)	−5.8 (−5.36)	2.0 (1.74)	0.6	$(2.2 \pm 0.4) \times 10^{-3}$	$(1.3 \pm 0.5) \times 10^{-4}$ [k]
P(NDI-T-BTD-T)	350, 490, 690	370 (9.0×10^4) 700 (1.1×10^5)	810	1.5	−3.8 (−3.63)	−5.7 (−5.39)	1.9 (1.76)	0.4	$(5.3 \pm 0.3) \times 10^{-2}$	$(1.2 \pm 0.4) \times 10^{-3}$

Values extracted from UV-Vis-NIR absorbance spectra (Fig. 2).[abcd]

[a] Wavelength of peak maxima in chlorobenzene solution at 298 K. [b] Wavelength of peak maxima in as-cast thin films. Brackets: absorption coefficients calculated by $\alpha = (2.303 \text{ abs})/\text{thickness}$ in cm^{-1},[39] [c] Wavelength of absorbance onset in films. [d] Corresponding optical band gap calculated by $hc/\lambda_{\text{onset}}$, [e] Values extracted from CV on films (Fig. 3a). Calculated by E_{LUMO} (eV) $\approx -[4.8 + E^{\text{red}}_{\text{onset}}$ (V vs. Fc$^+$/Fc)] [f] Values extracted from CV on films (Fig. 3a). Calculated by E_{LUMO} (eV) $\approx -[4.8 + E^{\text{red}}_{\text{onset}}$ (V vs. Fc$^+$/Fc)]. [g] Values extracted from CV on films (Fig. 3a). Corresponding electrochemical band gap calculated by E_g^{CV} (eV) $= E_{\text{LUMO}} - E_{\text{HOMO}}$. Brackets: DFT results at the B3LYP/TZ2P level. [h] Exciton binding energy calculated by $E_b = E_g^{\text{CV}} - E_g^{\text{opt}}$.[36] [i] Electron field-effect mobilities in saturation regime optimized by post-deposition thermal annealing in inert atmosphere. [j] Optimized electrical conductivity with N-DMBI doping. [k] Without thermal activation of the dopant.

and acceptor units.[33] When deposited as thin films, no obvious change in the S_0–S_1 transition bands is observed. As previously reported,[34] no change in the ICT band is observed for P(NDI-T-BTD-T). However, significant bathochromic shifts of the ICT bands are observed for P(NDI-2T-C$_{12}$) from solution to film. Such a red shift indicates a reduction in the energy gap associated with charge transfer. In solution, the ICT is associated to an intra-chain HOMO–LUMO transition, while in solid state it is attributed to an inter-chain charge transfer from the HOMO of one chain to the LUMO of another π-stacked chain, in accordance with the π-stacking observed in GIWAXS.[35] This results in a strong red-shift of about 150 nm and 40 nm of the absorption onsets for P(NDI-2T-C$_{12}$) and P(NDI-T-BTD-T), respectively.

The difference with the transport bandgap of 2.0 eV (E_g measured by CV) indicates a rather high exciton binding energy (E_b) of about 0.6 eV.[36]

In thin films, P(NDI-2T-C$_{12}$) exhibits two peaks in terms of absorption coefficients at 8.3×10^4 cm^{-1} at 390 nm and 5.9×10^4 cm^{-1} at 725 nm. In comparison, P(NDI-T-BTD-T) shows higher absorption coefficients for almost the entire visible range (450 to 750 nm) with two maxima at 9.0×10^4 cm^{-1} at 370 nm and 1.1×10^5 cm^{-1} at 700 nm (Table 1).

This improvement is tentatively attributed to enhanced polaronic delocalisation due to the extended π conjugation afforded by the incorporation of the BTD spacer.

Cyclic voltammetry (CV) experiments were carried out to elucidate the optoelectronic properties of the two polymers (Fig. 3a and b). The voltammograms were recorded at a scan rate of 50 mV s^{-1} without stirring to limit analyte convection.[37] The polymers were drop-cast films on the working electrode from chloroform solution for P(NDI-2T-C$_{12}$) and chlorobenzene for P(NDI-T-BTD-T). The LUMO and HOMO energy levels were estimated from the reduction onset potential E_{red} and the oxidation onset potential E_{ox}, respectively. For both polymers the oxidation waves show a lack of reversibility, which is not surprising for n-type materials. In agreement with the DFT-calculated values, the estimated HOMO energy levels are in the same range for both polymers, with P(NDI-2T-C$_{12}$) at −5.8 eV and P(NDI-T-BTD-T) at −5.7 eV. When analysing the reduction process, we found that both polymers P(NDI-2T-C$_{12}$) and P(NDI-T-BTD-T) exhibit stepwise one-electron transfers upon reduction.[17]

These remarkably stable and reversible reduction waves are attributed to the two successive reductions of the NDI moiety (Fig. S12, ESI†). Such stability of the anion and di-anion species
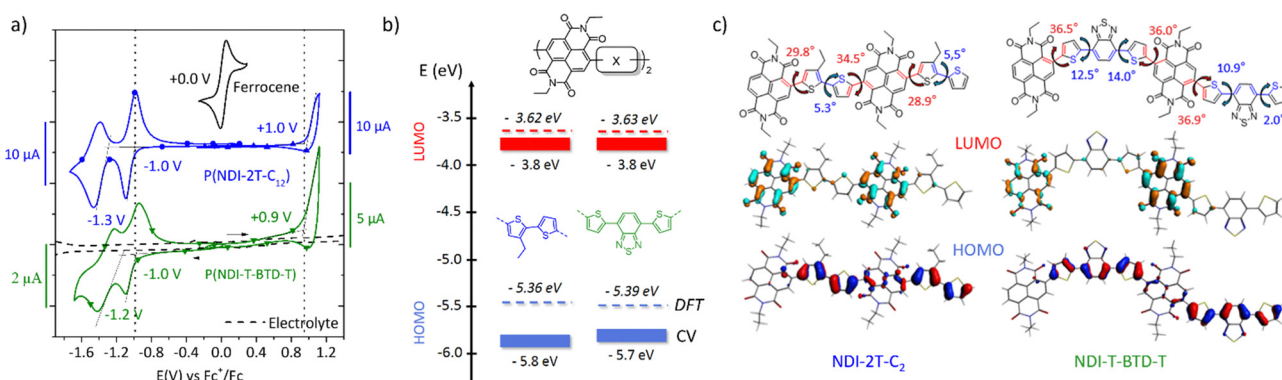

Fig. 3 (a) Cyclic voltammograms of both NDI-based copolymers. Experimental conditions: acetonitrile (MeCN), scan rate = 50 mV s^{-1}, 0.1 M [NBu$_4$][PF$_6$], Pt electrode, 0.01 M Ag/AgNO$_3$ calibrated Fc$^+$/Fc. (b) Comparison of CV-measured (thick lines ±0.1 eV) and DFT-calculated (dashed lines) HOMO and LUMO energy levels. (c) Ground-state DFT calculations: (top) geometry-optimized structures and dihedral angles (VWN + PBE/TZ2P/Grimme 3/MeCN), (middle and bottom) LUMO and HOMO electronic density distributions, respectively (B3LYP/TZ2P/MeCN). The copolymers are approximated to dimers and the alkyl chains are replaced by ethyl chains.

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is valuable for the long-term operation in thermoelectric devices. It should also be noted that for P(NDI-T-BTD-T), the introduction of a BTD unit has a minor influence on the band gap when compared to P(NDI-2T-C$_{12}$). It should be noted that for P(NDI-T-BTD-T), the values reported by Michinobu and coll. are in accordance with our measurements ($E_{LUMO} = -3.8$ *vs.* -3.52 eV,[34] $E_{HOMO} = -5.7$ *vs.* -5.68 eV[34]).

These experimental observations are supported by modelling. density functional theory (DFT) calculations allowed us to estimate the overall molecular geometry, the electronic density distribution along the conjugated backbone and the energy levels of the frontier orbitals (Fig. 3b and c). Since the length of the alkyl side chains has a negligible effect on the position of the energy levels, phenyl-alkyl and alkyl chains have been substituted by phenyl-methyl and ethyl or propyl chains respectively.[38] From the DFT calculations, it can be clearly seen that the HOMO energy level is spatially extended over the two thiophenes in P(NDI-2T-C$_{12}$) and additionally delocalised on the benzene ring of the BTD spacer in P(NDI-T-BTD-T). The large torsional angles between the thiophenes and the NDI unit in the two polymers (comprised between 28° and 38°) prevent the delocalisation of the HOMO over the NDI units. The HOMO energy levels are found at -5.36 eV in P(NDI-2T-C$_{12}$) and -5.39 eV in P(NDI-T-BTD-T). Regarding to the LUMO energy levels, they are strongly localised on the NDI units in both systems, and the energy level are found at around -3.6 eV.

Transport properties in pristine and doped states

Organic field-effect transistor (OFET). With the goal to characterize the intrinsic ability of the copolymers to transport charge carriers, the mobility of electrons was investigated in organic field-effect transistors (OFETs) with neat undoped polymers as channel. The electron-only OFETs made of P(NDI-2T-C$_{12}$) and P(NDI-T-BTD-T) were optimized by step-by-step annealing process ranging from 80 °C to 200 °C. Each step involved a 40 °C increase and lasted for 10 minutes, all conducted within a N$_2$-filled glovebox. In both cases, the transistors exhibited high series resistances independently of the thermal treatment applied. This is illustrated by the low source-to-drain current (I_{DS}) measured at V_{GS} below 60 V in output characteristics (Fig. S15c and d, ESI†). Therefore, the electron mobility values in linear regime are likely underestimated. In addition, no obvious saturation regime is obtained for V_{DS} up to 100 V. As mobility rises exponentially with respect to the root square of the electric field magnitude, we limited V_{DS} to 100 V.[40,41] The electrical properties extracted in linear regime ($V_{DS} = +10$ V) and in saturation regime ($V_{DS} = +100$ V) are summarized in Table S3 (ESI†). First, we found that both materials exhibit relatively good electron mobility in the 10^{-3} cm^2 V^{-1} s^{-1} range for P(NDI-2T-C$_{12}$) and 10^{-2} cm^2 V^{-1} s^{-1} range for P(NDI-T-BTD-T). Such mobility can be correlated to the edge-on orientation that favours in-plane charge conduction *via* inter chain hopping as demonstrated by grazing incidence wide angle X-ray scattering (GIWAXS) analysis (Fig. S16, ESI†).[42] The higher performances obtained for P(NDI-T-BTD-T) are tentatively attributed to both: the extended planarity of the backbone thanks

to the BTD unit which enhances edge-on π-stacking, and to the higher molar mass obtained which favours long-range transport.[43] The mobility values in the saturation regime are about one order of magnitude higher than in the linear regime, which is not surprising given the higher carrier concentration in the channel in the saturation regime. Low threshold voltages (V_{Th}) ranging from -2 V to 9 V are obtained, with a systematic increase upon thermal annealing.

The evolution of the electron mobility *versus* annealing temperature in the saturation regime is shown in Fig. 4. The corresponding output and transfer characteristics are shown in Fig. S15c–f (ESI†). Interestingly, both materials show enhanced field-effect mobility after thermal annealing. P(NDI-2T-C$_{12}$) shows a maximum mobility of $(2.2 \pm 0.4) \times 10^{-3}$ cm^2 V^{-1} s^{-1} after annealing at 120 °C. Above 120 °C, both the electron mobility and I_{on}/I_{off} ratio decrease. This behaviour was confirmed on several devices. Considering the remarkable thermal stability up to 350 °C of P(NDI-2T-C$_{12}$) measured by thermogravimetric analysis on powder (TGA, Fig. S13, ESI†), the decrease in mobility upon annealing therefore suggests a degradation of the OFET device rather than a degradation of the polymer itself. Potential causes of degradation include unfavourable morphological rearrangement of the polymer channel impeding long-range charge transport[44] (*e.g.*, charge trapping, grain boundaries, and lower molecular order) and degradation of the polymer/contact interface resulting in higher resistances to charge injection. Conversely, P(NDI-T-BTD-T) shows a steady improvement with increasing annealing temperature up to stability limit at 200 °C, reaching $(5.3 \pm 0.3) \times 10^{-2}$ cm^2 V^{-1} s^{-1}, in excellent agreement with previously reported mobility.[31] GIWAXS diffractograms show a decrease of the π-stacking signal after annealing at 200 °C, but a favourable enhancement of the edge-on lamellar orientation, which is likely responsible for the improved mobility (Fig. S15b and d,
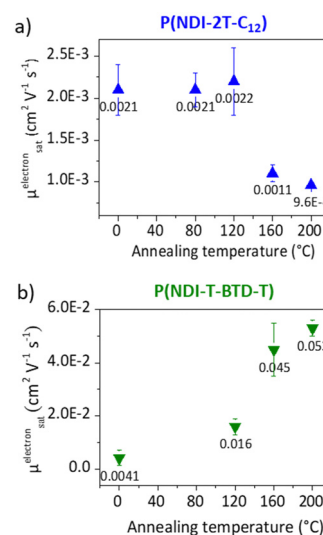

Fig. 4 (a) and (b) Electron mobility in saturation regime ($V_{DS} = +100$ V) as a function of the annealing temperature of P(NDI-2T-C$_{12}$) and P(NDI-T-BTD-T) bottom gate bottom contacts OFETs.

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ESI†). Interestingly, annealing the OFETs also appears to be an effective post-processing method to increase the I_{on}/I_{off} ratio and decrease hysteresis in-line with a charge-carrier trap reduction in the OFET channel after thermal annealing.[45,46]

N-doping of the polymers. With a view to possible application in thermoelectric devices, it is essential to verify that these polymers have good thermal stability. For this reason, a thermo-gravimetric analysis (TGA) was carried out. The materials showed good thermal stability up to 225 °C for P(NDI-T-BTD-T) with less than 5% loss, and good overall stability up to 350 °C for P(NDI-2T-C$_{12}$). (Fig. S13, ESI†).

To reach sufficient electrical conductivity for thermoelectric applications, we have doped the polymers with the widely used n-dopant precursor N-DMBI.[12,47] In our previous work, the HOMO energy level of N-DMBI was measured by CV at −4.4 eV.[48] Using the same setup, the LUMO levels of the polymers were measured at −3.8 eV (Fig. 3a), which should hinder direct electron transfer from N-DMBI HOMO. Doping *via* SOMO-mediated electron and hydride transfers are two possible pathways.[49–51]

Electron paramagnetic resonance (EPR) measurements provide useful insights on the doping events occurring in solution prior to the deposition. EPR in solution allows to easily compare the reactivity of the two polymers in the presence of N-DMBI. The comparison is easier in solution than in thin films as the thickness and morphology of the films do not interfere. P(NDI-2T-C$_{12}$) and P(NDI-T-BTD-T) were stirred with optimum amount of N-DMBI in toluene and chlorobenzene respectively. Both solutions have the same total concentration of polymer (5 mg mL^{-1}). The EPR spectra of the solutions stirred at room temperature (RT) and after heating up at 120 °C/ 15 min are shown in Fig. 5. In both cases, no obvious radical signal is measured for neat undoped polymers. Conversely, doped films exhibit EPR signals with symmetric line shapes (mostly Lorentzian). P(NDI-T-BTD-T):N-DMBI solution stirred at RT (green line) clearly exhibits free radicals (signal at 3440 G, $g \sim 2.0$). This is a clear evidence of its doping at room temperature using N-DMBI.[52] By analogy with PCBM,[47] all spontaneous doping observed at RT in dark conditions can be associated to the spontaneous hydride transfer from N-DMBI-H to P(NDI-T-BTD-T). This leads to the formation of the polymer anion (P–H$^-$), followed by an electron transfer yielding to the neutral radical (P–H$^•$) and the radical anion (P$^{•-}$) stabilized by N-DMBI$^+$. As expected, the quantity of spins is significantly increased (by approximately 650%) after heating the doped P(NDI-T-BTD-T) solution at 120 °C (Fig. 5d green line). Similarly, the thermal activation of the dopant allows efficient doping of P(NDI-2T-C$_{12}$). In such conditions, the SOMO mediated doping pathway is enabled. According to the double integration of the signal, less spins are generated on doped-P(NDI-2T-C$_{12}$), however a rate limited reaction (due to lower concentration of N-DMBI) and the presence of diamagnetic bipolaronic species (inactive in EPR) cannot be excluded. Interestingly, doped NDI-based polymers exhibit both air stability over 30 minutes of exposure (purple lines). The slight signal increase could originate from the kinetic of spin

Fig. 5 EPR spectra of (a) neat P(NDI-2T-C12) and P(NDI-2T-C12):N-DMBI mixed at RT in toluene (inert atmosphere, in dark). (b) Neat P(NDI-T-BTD-T) and P(NDI-T-BTD-T):N-DMBI mixed at RT in chlorobenzene (inert atmosphere, dark). (c) and (d) Same solutions heated up at 120 °C in inert atmosphere and exposed to air. Magnetic field power: 158.63 mW, modulation amplitude: 3.0 G.

generation. These experiments show that: (i) both polymers can be n-doped with N-DMBI in solution, (ii) heating up the solution at 120 °C allows more efficient doping, (iii) the doping is stable in aerated solution over short period of time. The UV-Vis-NIR absorbance spectra of neat and doped films of P(NDI-2T-C12) and P(NDI-T-BTD-T) at various dopant concentrations of N-DMBI are presented in Fig. 6a–d. The new optical transitions can be used as a fingerprint to identify the formation of charge-transfer complexes (CTCs) or ion-pairs by UV-Vis-NIR spectroscopy.[53] Neat films exhibit two neutral peaks with absorbance maxima at 400 nm/725 nm and 370 nm/700 nm for P(NDI-2T-C$_{12}$)/P(NDI-T-BTD-T), assigned to S$_0$–S$_1$/intramolecular charge transfer (ICT) transitions respectively. The S$_0$–S$_1$ transition of the polymers overlap with the absorption feature of neat N-DMBI molecules. For doped P(NDI-T-BTD-T), the increase of dopant concentration induces a bleaching of the neutral absorption band concomitant with a growth of broad sub-gap absorption transitions in the near-infrared domain. Similar behaviour is observed for doped P(NDI-2T-C$_{12}$) films, the bleaching of the neutral peak being accompanied by the emergence of a high-energy signal at *circa* 500 nm and a broadening of the ICT peak in the 900–1100 nm range – in accordance with previous studies on doped N2200 and NDI-based polymers.[23,54] The sub-gap absorbing bands are characteristic of transitions to (bi)polaronic levels, indicating effective doping.[55] The films were then annealed at 120 °C/1 h under inert atmosphere to thermally activate N-DMBI dopant.

Surprisingly, it causes a slight dedoping of P(NDI-2T-C$_{12}$) layers evidenced by a recovery of the neutral absorption band. In opposition, the doping level increases for annealed P(NDI-T-BTD-T) layers at high dopant concentration, as expected. To help analysing the data, we plotted the normalized neutral peak intensity evolution as a function of the dopant concentration
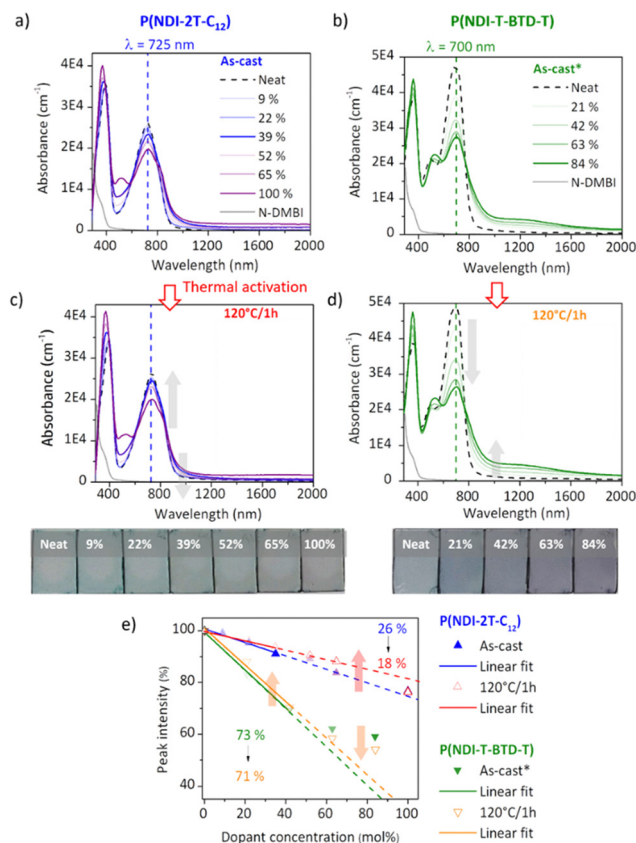
Fig. 6 UV-Vis-NIR absorbance spectra of neutral and doped (+x mol% of N-DMBI) (a) and (c) P(NDI-2T-C$_{12}$) (thickness = 30–40 nm) and (b) and (d) P(NDI-T-BTD-T) thin films (thickness = 40–45 nm). (a) and (b) As-cast (*from 110 °C solution). (c) and (d) After annealing at 120 °C/1h in glovebox. The arrows symbolize the evolution trend after annealing. Inset: Pictures of all the corresponding samples. (e) Absorption peak intensity at the neutral peak (725 nm for P(NDI-2T-C$_{12}$) and 700 nm for P(NDI-T-BTD-T)) normalized to that of the neat film. Inset: Doping efficiency derived from the slope of the linear correlation between the peak intensity and the dopant concentration. The arrows highlight the evolution trends upon thermal activation.

for as-cast and annealed films. It provides useful insights on the doping efficiency.

At low dopant concentration, we observe a linear dependence of the peak intensity with the dopant concentration. It suggests a first order reduction reaction of neutral chromophore to polaronic species, as supported by the isosbestic point at ca. 800 nm.[56] The doping efficiency is approximated to the slope of the linear fit.[23] The sub-linear deviation at high dopant concentration (>50 mol%) observed for P(NDI-T-BTD-T) can be attributed to either the formation of higher reduction states that absorb at the sampled wavelengths (e.g. bipolaronic species),[57] or the aggregation of N-DMBI within the polymer medium[58] which would decrease the ionization efficiency.[59] However, despite this, the doping efficiency of P(NDI-T-BTD-T) is significantly higher compared to P(NDI-2T-C$_{12}$), with a doping yield of around 70% vs. 20–25% (see slope of linear fit of Fig. 6e). One should keep in mind that P(NDI-T-BTD-T) films are casted from hot solutions, hence improving the doping. Yet,

the higher doping efficiency observed by absorbance spectroscopy is in good accordance with the higher spin density found in EPR. For P(NDI-2T-C$_{12}$) the doping level remains low, indicating that most of the added dopants do not contribute to (bi)polaron formation.

Upon annealing, we could expect an increase of the doping efficiency similarly to what was observed in solution in EPR, but the doping yield is decreased from 27% to 18%. It hence seems that, in the solid state, N-DMBI activation is counterbalanced by another effect. We clarified this phenomenon by probing the topography of doped films (22 mol%) before and after annealing using atomic force microscopy (AFM) (Fig. S17, ESI†) and found aggregates on the top surface of the films regardless of the deposition solvent used. It appears that neutral N-DMBI-H and/or the cation N-DMBI$^+$ are partly immiscible with P(NDI-2T-C$_{12}$), similarly to N2200.[60] Consequently, the low doping level evidenced by UV-Vis-NIR can be rationalized by the shielding of dopants, which remain electrically inactive inside the aggregates. Moreover, upon annealing the aggregates tend to cluster together. As a result, the dopant/polymer interface area decreases, leading to a reduction in the effective number of neat dopants and cations available to generate and stabilize doped states, respectively. This explains the unexpected loss of doping spectral signature despite the activation of N-DMBI upon annealing in the solid state. For P(NDI-T-BTD-T), a similar impact of thermal annealing can be predicted for low dopant concentration considering the slight decrease of efficiency from 73% to 71%. In opposition, for dopant concentration above 64 mol%, annealing increases the doping efficiency.

This result implies that the aggregates in as-cast layers – presumably identified by the sub-linear deviation – contain neutral N-DMBI-H molecules that are enabled to participate in doping by morphological rearrangement upon annealing.

These observations shed light on the many intertwined parameters that determine the effectiveness of doping with N-DMBI for this class of polymers. The processing conditions, including stirring time and temperature, should be explicitly reported and should not be overlooked. The conductivity associated with the processing conditions discussed above is presented in Fig. 7.

The mobility of undoped polymers is determined using OFET with gold contacts. For fair comparison, we used gold contacts for conductivity measurements to ensure similar contact resistances. Note that high injection barriers are expected considering the large energy offset between the LUMO levels of the polymers (−3.8 eV) and the work function of gold (−5.1 eV).[61] The measured electrical properties may thus be underestimated.[62] We tried to estimate the conductivity of the doped layers by four-point probe in ambient conditions.

Unfortunately, no reliable values were obtained despite brief air exposure less than 30 seconds and apparent air stability observed in solution using EPR. Consequently, instability caused by diffusion of oxidative ambient species within the solid polymer matrix can be speculated.[63] Conductivity measurements in inert atmosphere are hence performed on both undoped and N-DMBI-doped films using the Transfer Length

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Fig. 7 (a) and (b) Evolution of the conductivity of doped (a) P(NDI-2T-C$_{12}$) and (b) P(NDI-T-BTD-T) films when exposed to ambient and inert atmospheres. Red lines correspond to the samples at optimal dopant concentration after annealing (thickness: 40 nm both). Black lines correspond to samples only aged under N$_2$ atmosphere for reference (thickness: 45 nm and 40 nm respectively). (c) and (d) Evolution of the conductivity of doped (c) P(NDI-2T-C$_{12}$) and (d) P(NDI-T-BTD-T) films. Red lines correspond to the samples at optimal dopant concentration after annealing (thickness: 40 nm both). Black lines correspond to samples only aged under N$_2$ atmosphere for reference (thickness: 45 nm and 40 nm respectively).

Method (TLM, Fig. S18, ESI†). All the neat undoped films exhibit intrinsic resistivity high enough compared to the input impedance of the setup to reach the measurement limit. This is consistent with the low carrier density expected for non-resonant structures. Considering the lowest conductivity measured on this setup, the conductivity of the neat polymers should be inferior to 10^{-8} S cm^{-1}. The electrical conductivities of as-cast and annealed doped films as a function of the dopant concentration are shown in Fig. 7a and b. The samples discussed here were deposited from the same solutions than the UV-Vis-NIR samples on the same day, in the same conditions for direct comparison.

Considering as-cast doped films, we clearly see that in both cases (blue and green lines) the conductivity rises with respect to the dopant concentration before decreasing. This phenomenon is commonly observed for molecular-doped organic layers.[23,60,64,65] The decrease of activation energy with the addition of dopant is usually counterbalanced by a synergetic effect of dopant segregation, in accordance to the AFM and UV-Vis-NIR experiments, and a loss of morphological order due to lattice disruption upon dopant uptake. Upon annealing at 120 °C/1 h (red and orange lines), two distinct behaviours are observed (highlighted by arrows). At low dopant concentration, the conductivity decreases by more than one order of magnitude, whereas at high concentration the conductivity increases. The evolution of the conductivity agrees with the formation of (bi)polaronic absorbance bands previously discussed. Further annealing the films up to 16 h exacerbates these behaviours. Hence, the annealing time should also be controlled with care.

Maximum conductivities of $(1.3 \pm 0.5) \times 10^{-4}$ S cm^{-1} for as-cast P(NDI-2T-C$_{12}$):N-DMBI (39 mol%) film and of $(1.2 \pm 0.4) \times 10^{-3}$ S cm^{-1} for annealed (120 °C/16 h) P(NDI-T-BTD-T):N-DMBI (84 mol%) film are achieved. Note that the optimal dopant concentrations (39 and 84 mol%) are high compared to the high-performing NDI-based polymers reported,[24,60,66] especially for P(NDI-T-BTD-T). Considering the high doping efficiency determined for P(NDI-T-BTD-T), this observation suggests that dopant/polymer miscibility is not the limiting factor but rather that the carriers generated at low doping level are poorly mobile at the macro-scale (µm-mm) probed by TLM.

The origins can be diverse, for instance: low dissociation of charge transfer states leading to bound charges, high density of electron traps, grain boundaries hindering long-range hopping. The conductivities reached for P(NDI-T-BTD-T) are about one order of magnitude higher than the conductivities achieved with P(NDI-2T-C$_{12}$). It corroborates with the higher mobility observed for undoped films in OFET. While the achieved conductivity values may not reach the level of the highest performing thermoelectric materials,[11,15] we have successfully demonstrated an enhancement in conductivity using N-DMBI without the requirement of post-deposition thermal annealing. Avoiding post-deposition annealing is particularly interesting for polymer films deposited on flexible substrates being more sensitive to high temperature.

The values reported for N-DMBI-doped N2200 are in the range of 10^{-4} to 10^{-2} S cm^{-1} after thermal annealing.[67,68] With P(NDI-2T-C$_{12}$) we achieved a maximum conductivity that was only slightly inferior. One avenue for improvement is to optimize the synthesis route to produce higher molar masses. These are expected to improve the long-range interconnectivity of the polymer chains and provide higher charge transport properties. It should be noted that we deposited doped-P(NDI-2T-C$_{12}$) films by a simple spin-coating method, from a non-halogenated solvent (toluene) and without heating at any step.

Stability of the doped films. In views of applicative use, the stability of doped films over time and upon exposure to ambient is investigated. The details of the ageing tests and the resulting conductivity variations are presented in Fig. 7c and d. To facilitate the comparison of the conductivity stability between P(NDI-2T-C$_{12}$) and P(NDI-T-BTD-T), the scales of both graphs have been set to be identical.

For as-cast films, both materials exhibit conductivities of about 10^{-4} S cm^{-1}. After annealing, the conductivity of P(NDI-2T-C$_{12}$) slightly decreases while the conductivity of P(NDI-T-BTD-T) is enhanced of one order of magnitude as explained. The samples were then exposed to air one hour in ambient dark conditions, demonstrating remarkable air-stability. We conclude that the characterization made in air are relevant for the doped state of these polymers with LUMO level at -3.8 eV. After 15 hours of exposure, a five-fold conductivity drop is

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observed for P(NDI-2T-C$_{12}$) *vs.* a two-fold drop for P(NDI-T-BTD-T) under the same conditions. The electron affinity and layer thicknesses being similar (*ca.* 40 nm), the slower degradation kinetic is tentatively assigned to a higher degree of crystallinity of doped-P(NDI-T-BTD-T) films – as observed for neat polymer (see GIWAXS in Fig. S16, ESI†). However, the partial recovery of the conductivity when returned to inert atmosphere (9 days in N$_2$) suggests a reversible process. Moisture and oxygen are known to act as electron traps, thereby reducing the overall conductivity.[69] Consequently, the conductivity loss is mainly assigned to the reversible adsorption and diffusion of moisture and molecular oxygen inside the polymer matrix. To highlight this phenomenon, we exposed the layers to air more than two days (49 h). Most of the mobile charges are trapped in P(NDI-2T-C$_{12}$) film – the conductivity being below the measurement limit – while the conductivity of P(NDI-T-BTD-T) dropped by two orders of magnitudes (Fig. 7c and d). Remarkably, the initial values of conductivity can be recovered after annealing at 120 °C in only one hour in the glovebox. In 2018, one article reported on a similar behavior for N-DMBI-doped polymer.[66] They assigned the conductance recovery to the accelerated desorption of the electron-trapping species upon post annealing.

A similar claim is found for another n-type dopant.[70] Yet, considering the low doping efficiency and the dopant propensity to aggregate, the hypothesis of activation of pristine N-DMBI-H molecules trapped within the aggregates due to morphological rearrangement cannot be ruled out. The black lines correspond to samples only aged in glovebox. After 23 days in the glove box, 78% and 56% of the conductivity was retained for doped-P(NDI-2T-C$_{12}$) and doped-P(NDI-T-BTD-T) films respectively. It is noteworthy that the loss of spontaneous conductivity on ageing under N$_2$ seems to be relatively similar for films exposed to air (red) or not (black) (see between 9 days and 23 days). Identical conclusions were drawn by monitoring the (bi)polaronic absorption band bleaching/growth upon air exposure/re-annealing using UV-Vis-NIR spectroscopy. Overall, good conductivities in the range of 10^{-3} to 10^{-4} S cm^{-1} have been achieved with NDI-based polymers. Improving the dopant:host interaction should be considered to increase the doping efficiency and achieve the carrier concentration required for high power factor.[66]

Thermoelectric properties. In the final part of this work, we set out to determine the thermoelectric properties of our materials. To do this, we deposited by drop-casting 2 μm-thick films (with the same optimum dopant ratios as found above) on chips for measurements with the Linseis thin film apparatus. (Linseis Messgeräte GmbH). Thermal activation of N-DMBI was realized *in situ* at 120 °C/16 h before measurements under vacuum (9 × 10^{-6} mbar). All in-plane thermoelectric parameters *i.e.* Seebeck coefficient (*S*), thermal conductivity (κ∥) and electrical conductivity (σ), were simultaneously extracted from a single film. Unfortunately, the thick films prepared from P(NDI-2T-C$_{12}$) proved too resistive for our measurements. As the P(NDI-T-BTD-T) doped samples were more conductive, we were able to measure their thermoelectric parameters. Note that the electrical conductivity found for these μm-thick samples is comparable to

that measured for the nm-thick samples. This study, although preliminary, allowed us to determine a maximum power factor of 9 × 10^{-3} μW m^{-1} K^{-2} at 115 °C with an *S* coefficient of about −94 μV K^{-1}. (See Fig. S19, ESI†). The κ∥ was found to be 0.38 W m^{-1} K^{-1} giving a *ZT* of 9.2 × 10^{-6} (See ESI† for calculation of power factor and *ZT*). Although these thermoelectric properties are being low compared to the state of the art, they are in very good agreement with those reported by Baran and co-workers for NDI-based copolymers with high N-DMBI ratios.[24]

Conclusions

In this work, we have synthesized and characterized the optoelectronic properties of two NDI- and thiophene-based copolymers differing by the presence of a BTD spacer unit with the aim of elucidating its effect on the transport properties and doping mechanism with N-DMBI. We show that the introduction of the BTD unit does not change the LUMO energy level found at −3.8 eV. The electron mobility of the materials is studied in OFETs and reaches 5.3 ± 0.3 × 10^{-2} cm^2 V^{-1} s^{-1} for the polymer embedding a BTD spacer between the thiophenes. Upon doping with N-DMBI, we found a higher doping efficiency for the BTD-containing copolymer, leading to a better conductivity of up to (1.2 ± 0.4) × 10^{-3} S cm^{-1} for annealed films (120 °C/16 h), which is an order of magnitude higher compared to the copolymer without BTD.

Notably, we have observed spontaneous doping of the copolymers without thermal treatment when cast with N-DMBI at room temperature and we show that post-deposition annealing can be detrimental in certain cases, here for low dopant concentrations. These findings are in stark contrast to the typical approach to n-doping with N-DMBI and provide the field with nuances regarding the need for and effect of thermal activation of the dopant. Therefore, we argue that as-cast N-DMBI doped films should be characterised prior to any thermal treatment. Indeed, a systematic post-deposition annealing may not be optimal depending on the N-DMBI:host miscibility and the tendency of the semiconducting host to form C–H bonds.

The stability of the doped films over time and when exposed to ambient conditions was investigated. We show that the introduction of the BTD units provides the polymer with a remarkable stability in the conducting state, even after a short period of exposure to air. We also found that the degradation of conductivity after prolonged exposure to air is reversible and that the initial values can be recovered after annealing at 120 °C in only one hour in a glove box. A further improvement in stability could be achieved with adapted encapsulation layers. Finally, we have measured the thermoelectric parameters of the doped BTD-based films and found a power factor of 9 × 10^{-3} μW m^{-1} K^{-2} at 115 °C.

Author contributions

O. B.: conceptualization, investigation, formal analysis, data curation, methodology, visualization, writing original draft,

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writing review & editing. Y. K.: investigation. A. A. M.: investigation, methodology, visualization. S. P.: methodology. T. N. D.: investigation, methodology. A. C.: methodology, validation. C. A.: investigation, methodology, validation. P. L.: investigation, methodology. A. C conceptualization, project administration. R. D.: conceptualization, formal analysis, project administration, supervision, resources, writing review & editing.

Conflicts of interest

There are no conflicts to declare.

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Notes and references

1 F. Garnier, G. Horowitz and D. Fichou, *Synth. Met.*, 1989, **28**, 705–714.
2 G. Horowitz, *Adv. Mater.*, 1990, **2**, 287–292.
3 Z. Qin, C. Gao, W. W. H. Wong, M. K. Riede, T. Wang, H. Dong, Y. Zhen and W. Hu, *J. Mater. Chem. C*, 2020, **8**, 14996–15008.
4 A. Marrocchi, A. Facchetti, D. Lanari, C. Petrucci and L. Vaccaro, *Energy Environ. Sci.*, 2016, **9**, 763–786.
5 G. Horowitz, *J. Mater. Res.*, 2004, **19**, 1946–1962.
6 O. Ostroverkhova, *Chem. Rev.*, 2016, **116**, 13279–13412.
7 G. Horowitz, D. Fichou and F. Garnier, *Solid State Commun.*, 1989, **70**, 385–388.
8 C. Gu, X. Su, Y. Li, B. Liu, Y. Tian, W. Tan, J. Ma and X. Bao, *Mol. Syst. Des. Eng.*, 2022, **7**, 1364–1384.
9 J. Quinn, J. Zhu, X. Li, J. Wang and Y. Li, *J. Mater. Chem. C*, 2017, **5**, 8654–8681.
10 S. Yu, C. J. Kousseff and C. B. Nielsen, *Synth. Met.*, 2023, **293**, 117295.
11 A. Tripathi, Y. Lee, S. Lee and H. Young Woo, *J. Mater. Chem. C*, 2022, **10**, 6114–6140.
12 D. Yuan, W. Liu and X. Zhu, *Chem. Soc. Rev.*, 2023, **52**, 3842–3872.
13 Y. Zhang, Y. Wang, C. Gao, Z. Ni, X. Zhang, W. Hu and H. Dong, *Chem. Soc. Rev.*, 2023, **52**, 1331–1381.
14 H. Jia and T. Lei, *J. Mater. Chem. C*, 2019, **7**, 12809–12821.
15 M. Li and Y. Shi, *ChemPlusChem*, 2023, **88**, e202300215.
16 Y. Matsunaga, K. Goto, K. Kubono, K. Sako and T. Shinmyozu, *Chem. – Eur. J.*, 2014, **20**, 7309–7316.
17 S. K. Lee, Y. Zu, A. Herrmann, Y. Geerts, K. Müllen and A. J. Bard, *J. Am. Chem. Soc.*, 1999, **121**, 3513–3520.
18 K. Zhang, R. Xia, B. Fan, X. Liu, Z. Wang, S. Dong, H.-L. Yip, L. Ying, F. Huang and Y. Cao, *Adv. Mater.*, 2018, **30**, 1803166.
19 Y. Xu, J. Yuan, S. Zhou, M. Seifrid, L. Ying, B. Li, F. Huang, G. C. Bazan and W. Ma, *Adv. Funct. Mater.*, 2019, **29**, 1806747.
20 A. F. Paterson, A. Savva, S. Wustoni, L. Tsetseris, B. D. Paulsen, H. Faber, A. H. Emwas, X. Chen, G. Nikiforidis, T. C. Hidalgo, M. Moser, I. P. Maria, J. Rivnay, I. McCulloch, T. D. Anthopoulos and S. Inal, *Nat. Commun.*, 2020, **11**, 3004.
21 N. Zhou and A. Facchetti, *Mater. Today*, 2018, **21**, 377–390.
22 Y. Song, J. Ding, X. Dai, C. Li, C. Di and D. Zhang, *ACS Mater. Lett.*, 2022, **4**, 521–527.
23 J. Liu, G. Ye, B. van der Zee, J. Dong, X. Qiu, Y. Liu, G. Portale, R. C. Chiechi and L. J. A. Koster, *Adv. Mater.*, 2018, **30**, 1804290.
24 D. R. Villalva, S. Singh, L. A. Galuska, A. Sharma, J. Han, J. Liu, M. A. Haque, S. Jang, A. H. Emwas, L. J. A. Koster, X. Gu, B. C. Schroeder and D. Baran, *Mater. Horiz.*, 2022, **9**, 500–508.
25 S. Sharma, N. B. Kolhe, V. Gupta, V. Bharti, A. Sharma, R. Datt, S. Chand and S. K. Asha, *Macromolecules*, 2016, **49**, 8113–8125.
26 A. C. Thompson, H. M. Grimm, A. G. Bé, K. J. McKnight and J. J. Reczek, *Synth. Commun.*, 2015, **45**, 1127–1136.
27 C. Thalacker, C. Röger and F. Würthner, *J. Org. Chem.*, 2006, **71**, 8098–8105.
28 T. D. M. Bell, S. Yap, C. H. Jani, S. V. Bhosale, J. Hofkens, F. C. De Schryver, S. J. Langford and K. P. Ghiggino, *Chem. – Asian J.*, 2009, **4**, 1542–1550.
29 J. K. Stille, *Angew. Chem., Int. Ed. Engl.*, 1986, **25**, 508–524.
30 S. Vasimalla, S. P. Senanayak, M. Sharma, K. S. Narayan and P. K. Iyer, *Chem. Mater.*, 2014, **26**, 4030–4037.
31 C. Gu, W. Hu, J. Yao and H. Fu, *Chem. Mater.*, 2013, **25**, 2178–2183.
32 W. H. Carothers, *Trans. Faraday Soc.*, 1936, **32**, 39.
33 T. E. Anderson, E. W. Culver, I. Badía-Domínguez, W. D. Wilcox, C. E. Buysse, M. C. R. Delgado and S. C. Rasmussen, *Phys. Chem. Chem. Phys.*, 2021, **23**, 26534–26546.
34 Y. Wang, T. Hasegawa, H. Matsumoto, T. Mori and T. Michinobu, *Adv. Mater.*, 2018, **30**, 1707164.
35 C. R. Martinez and B. L. Iverson, *Chem. Sci.*, 2012, **3**, 2191.
36 B. P. Rand, D. P. Burk and S. R. Forrest, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2007, **75**, 115327.
37 N. Elgrishi, K. J. Rountree, B. D. McCarthy, E. S. Rountree, T. T. Eisenhart and J. L. Dempsey, *J. Chem. Educ.*, 2018, **95**, 197–206.
38 G. Kim, A.-R. Han, H. Rang Lee, J. Lee, J. Hak Oh and C. Yang, *Chem. Commun.*, 2014, **50**, 2180–2183.
39 V. S. Sangawar and N. A. Moharil, *Chem. Sci. Trans.*, 2012, **1**, 447–455.
40 W. D. Gill, *J. Appl. Phys.*, 1972, **43**, 5033–5040.
41 J. Veres, S. D. Ogier, S. W. Leeming, D. C. Cupertino and S. M. Khaffaf, *Adv. Funct. Mater.*, 2003, **13**, 199–204.
42 H. Sirringhaus, P. J. Brown, R. H. Friend, M. M. Nielsen, K. Bechgaard, B. M. W. Langeveld-Voss, A. J. H. Spiering,

R. A. J. Janssen, E. W. Meijer, P. Herwig and D. M. de Leeuw, *Nature*, 1999, **401**, 685.

43 J.-M. Verilhac, G. LeBlevennec, D. Djurado, F. Rieutord, M. Chouiki, J.-P. Travers and A. Pron, *Synth. Met.*, 2006, **156**, 815–823.

44 A. Zen, J. Pflaum, S. Hirschmann, W. Zhuang, F. Jaiser, U. Asawapirom, J. P. Rabe, U. Scherf and D. Neher, *Adv. Funct. Mater.*, 2004, **14**, 757–764.

45 W. Shi, Y. Zheng, A. D. Taylor, J. Yu and H. E. Katz, *Appl. Phys. Lett.*, 2017, **111**, 043301.

46 Y.-H. Zhao, W. Li, T. Shen, Y. Zhao, Y. Liu and Y. Wang, *Sci. China: Chem.*, 2023, **66**, 548–561.

47 B. D. Naab, S. Guo, S. Olthof, E. G. B. Evans, P. Wei, G. L. Millhauser, A. Kahn, S. Barlow, S. R. Marder and Z. Bao, *J. Am. Chem. Soc.*, 2013, **135**, 15018–15025.

48 O. Bardagot, P. Kubik, T. Marszalek, P. Veyre, A. A. Medjahed, M. Sandroni, B. Grévin, S. Pouget, T. Nunes Domschke, A. Carella, S. Gambarelli, W. Pisula and R. Demadrille, *Adv. Funct. Mater.*, 2020, **30**, 2000449.

49 O. Bardagot, C. Aumaître, A. Monmagnon, J. Pécaut, P.-A. Bayle and R. Demadrille, *Appl. Phys. Lett.*, 2021, **118**, 203904.

50 S. Jhulki, H.-I. Un, Y.-F. Ding, C. Risko, S. K. Mohapatra, J. Pei, S. Barlow and S. R. Marder, *Chem*, 2021, **7**, 1050–1065.

51 F. Pallini, S. Mattiello, N. Manfredi, S. Mecca, A. Fedorov, M. Sassi, K. Al Kurdi, Y.-F. Ding, C.-K. Pan, J. Pei, S. Barlow, S. R. Marder, T.-Q. Nguyen and L. Beverina, *J. Mater. Chem. A*, 2023, **11**, 8192–8201.

52 J. Han, C. Ganley, Q. Hu, X. Zhao, P. Clancy, T. P. Russell and H. E. Katz, *Adv. Funct. Mater.*, 2021, **31**, 2010567.

53 I. Zozoulenko, A. Singh, S. K. Singh, V. Gueskine, X. Crispin and M. Berggren, *ACS Appl. Poly. Mater.*, 2019, **1**, 83–94.

54 Y. Wang, M. Nakano, T. Michinobu, Y. Kiyota, T. Mori and K. Takimiya, *Macromolecules*, 2017, **50**, 857–864.

55 D. Huang, H. Yao, Y. Cui, Y. Zou, F. Zhang, C. Wang, H. Shen, W. Jin, J. Zhu, Y. Diao, W. Xu, C. Di and D. Zhu, *J. Am. Chem. Soc.*, 2017, **139**, 13013–13023.

56 P. Cavassin, I. Holzer, D. Tsokkou, O. Bardagot, J. Réhault and N. Banerji, *Adv. Mater.*, 2023, 2300308.

57 G. Rebetez, O. Bardagot, J. Affolter, J. Réhault and N. Banerji, *Adv. Funct. Mater.*, 2022, **32**, 2105821.

58 J. Liu, L. Qiu, R. Alessandri, X. Qiu, G. Portale, J. Dong, W. Talsma, G. Ye, A. A. Sengrian, P. C. T. Souza, M. A. Loi, R. C. Chiechi, S. J. Marrink, J. C. Hummelen and L. J. A. Koster, *Adv. Mater.*, 2018, **30**, 1704630.

59 J.-M. Kim, S.-J. Yoo, C.-K. Moon, B. Sim, J.-H. Lee, H. Lim, J. W. Kim and J.-J. Kim, *J. Phys. Chem. C*, 2016, **120**, 9475–9481.

60 R. A. Schlitz, F. G. Brunetti, A. M. Glaudell, P. L. Miller, M. A. Brady, C. J. Takacs, C. J. Hawker and M. L. Chabinyc, *Adv. Mater.*, 2014, **26**, 2825–2830.

61 L.-L. Chua, J. Zaumseil, J.-F. Chang, E. C.-W. Ou, P. K.-H. Ho, H. Sirringhaus and R. H. Friend, *Nature*, 2005, **434**, 194–199.

62 Y. Zhao, Y. Guo and Y. Liu, *Adv. Mater.*, 2013, **25**, 5372–5391.

63 T. Nunes Domschke, O. Bardagot, A. Benayad, R. Demadrille, A. Carella, R. Clerc and A. Pereira, *Synth. Met.*, 2020, **260**, 116251.

64 J. Liu, L. Qiu, G. Portale, S. Torabi, M. C. A. Stuart, X. Qiu, M. Koopmans, R. C. Chiechi, J. C. Hummelen and L. J. Anton Koster, *Nano Energy*, 2018, **52**, 183–191.

65 K. Shi, F. Zhang, C.-A. Di, T-W. Yan, Y. Zou, X. Zhou, D. Zhu, J.-Y. Wang and J. Pei, *J. Am. Chem. Soc.*, 2015, **137**, 6979–6982.

66 D. Kiefer, A. Giovannitti, H. Sun, T. Biskup, A. Hofmann, M. Koopmans, C. Cendra, S. Weber, L. J. Anton Koster, E. Olsson, J. Rivnay, S. Fabiano, I. McCulloch and C. Müller, *ACS Energy Lett.*, 2018, **3**, 278–285.

67 Y. M. Gross, D. Trefz, C. Dingler, D. Bauer, V. Vijayakumar, V. Untilova, L. Biniek, M. Brinkmann and S. Ludwigs, *Chem. Mater.*, 2019, **31**, 3542–3555.

68 S. Wang, D. Fazzi, Y. Puttisong, M. J. Jafari, Z. Chen, T. Ederth, J. W. Andreasen, W. M. Chen, A. Facchetti and S. Fabiano, *Chem. Mater.*, 2019, **31**, 3395–3406.

69 R. T. Weitz, K. Amsharov, U. Zschieschang, E. B. Villas, D. K. Goswami, M. Burghard, H. Dosch, M. Jansen, K. Kern and H. Klauk, *J. Am. Chem. Soc.*, 2008, **130**, 4637–4645.

70 M. L. Tietze, B. D. Rose, M. Schwarze, A. Fischer, S. Runge, J. Blochwitz-Nimoth, B. Lüssem, K. Leo and J.-L. Brédas, *Adv. Funct. Mater.*, 2016, **26**, 3730–3737.

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