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I. INTRODUCTION

Polymer brushes are enormously versatile systems, whose properties can be tailored chemically and physically[1–3]. They were studied with great interest due to their technological significance, biophysical importance and theoretical subtleties and complexity. They find important applications in colloid stabilization, lubrication, and developments where friction[4], adhesion and wetting properties are important[1]. They are also used as coatings in “smart surfaces” to fine tune reversibly some property of an interface upon changes in external stimuli, for example, pH[5], temperature, solvent quality[6, 7]. Brushes are at the forefront of recent developments, ranging from responsive bio-interphases, controlled drug-delivery and release systems, thin films and particles, which act as sensors of minute amounts of analytes[3, 6]. The geometry and curvature of the grafted surface contribute also to distinctive properties of brushes, which can be grown in planar or cylindrical geometries, in the surface of nanoparticles or into the backbone of macromolecules (bottle-brushes)[1, 8]. An important biophysical system of end-grafted semiflexible polymers is the endothelial glyocalyx layer, present in the inner side of vasculature in animals and plants[9]. It has a complex structure formed by a matrix of macromolecular carbohydrates, consisting of proteoglycans and glycoproteins, that coat the surface of endothelial cells[9, 10].

We performed molecular dynamics simulations to study equilibrium and flow properties of a liquid in a nano-channel with confining surfaces coated with a layer of grafted semiflexible polymers. The coverage spans a wide range of grafting densities from essentially isolated chains to dense brushes. The end-grafted polymers were described by a bead spring model with an harmonic potential to include the bond stiffness of the chains. We varied the rigidity of the chains, from fully flexible polymers to rigid rods, in which the configurational entropy of the chains is negligible. The brush-liquid interaction was tuned to obtain a super-hydrophobic channel, in which the liquid did not penetrate the polymer brush, giving rise to a Cassie-Baxter state. Equilibrium properties such as brush height and bending energy were measured, varying the grafting density and the stiffness of the polymers. We studied also the characteristics of the brush-liquid interface and the morphology of the polymers chains supporting the liquid for different bending rigidities. Non-equilibrium simulations were performed, moving the walls of the channel in opposite directions at constant speed, obtaining a Couette velocity profile in the bulk liquid. The molecular degrees of freedom of the polymers were studied as a function of the Weissenberg number. Also, the violation of the no-slip boundary condition and the slip properties were analyzed as a function of the shear rate, grafting density and bending stiffness. At high grafting densities, a finite slip length independent of the shear rate or bending constant was found, while at low grafting densities a very interesting non-monotonic behaviour on the bending constant is observed.

It is known to have important functions as modulator of permeability in the exchange of water, primary molecular sieve of plasma proteins, mechanotransducer of fluid shear stress and as regulator of red and white blood cells[9]. From a mesoscopic point of view, this branched comb-like polymer brush, can be modelled as composed of semiflexible homopolymer chains whose beads have the size of the side-chains[10]. In this model, the glyocalyx is closely related to the geometry we study in this work.

Polymer brushes were studied intensely in equilibrium and non-equilibrium conditions[8, 11, 12] by computer simulations and theory[13]. However, the vast majority of work was devoted to fully flexible polymer chains. Computer simulation of coarse-grained systems studied friction of bearing brushes[14–17] and the flow of simple liquids or polymer melts confined in brush-coated polymer channels as function of grafting density, interface properties and flow intensity[18–22]. Milchev and Binder studied the structure of brushes formed by semiflexible chains by Monte Carlo and molecular dynamics simulations[23]. They found an interesting phase transition of the system under compression, which presents buckling at moderate pressures and bending at higher pressures[24]. The first bonds of the grafted polymers, with this reference, are directed perpendicularly to the growing plane, as it is the case in our work. Using soft potentials within the Dissipative Particle Dynamics (DPD) simulation scheme various groups studied different aspects of a polymer brush exposed to flow[10, 21, 25]. Deng et al.[10] included bending rigidity in their model to account for semiflexible polymers in the context of glyocalyx and studied brush height and slip length for

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Couette (shear-driven) and Poiseuille (pressure-driven) flows. Benetatos et al. explored the morphology and in-plane collapse of a grafted layer of attractive semiflexible chains, analysing the bundling of neighboring chains[26]. Kim et al studied the height behavior of the brush layer as function of chain stiffness, shear rate and grafting density by Brownian hydrodynamics, lattice Boltzmann simulations and mean field theories[27]. More recently, Römer and Fedosov[28] developed a theoretical model and compared with DPD simulations, a brush of stiff polymers under flow. They found a good agreement of model and simulations for brush height, velocity profile and apparent flow viscosity, for a wide range of shear rates and grafting densities. Active semiflexible polymers (filaments) have been studied to analyse efficiency of pumping of liquid[29] and methachronal waves in bidimensional arrays of grafted chains[30]. In these works the focus is in the hydrodynamic coupling of the semiflexible chains and its role in transport efficiency.

Relevant to the present work are also the superhydrophobic surfaces. In these, surface texture or structure at the micro and/or nanoscopic level, is used to enhance the intrinsic hydrophobic chemistry of the surface, to produce highly non-wetting surfaces. Superhydrophobic surfaces typically exhibit very high water-repellency, to produce highly non-wetting surfaces. Superhydrofobic surfaces. In these, surface texture or structure at the micro and/or nanoscopic level, is used to enhance the inherent hydrophobicity. Micrometer-sized surface textures have been shown to increase contact angles by over 30°[20] and to enable the system to maintain water-repellency for a wide range of flow conditions[21, 22]. In these works the focus is in the hydrodynamic coupling of the semiflexible chains and its role in transport efficiency.

In this work, we study in equilibrium and under flow, the interface of a liquid and a layer of end-grafted semiflexible polymers for different bending rigidities and as a function of grafting density. The interaction between grafted chains and liquid is chosen such that the interface is hydrophobic and compatible with super-hydrophobic behavior, for a wide range of grafting densities. A typical snapshot of the system in the superhydrophobic Cassie-Baxter state, for the stiffest polymers considered in this work, can be observed in Fig. 1. We studied a small number of bending rigidities, ranging from fully flexible to highly stiff polymers, such that the polymer resembles a pillar, typical of superhydrophobic surfaces, and its configurational entropy is minimal. This implies also statistical Kuhn segments much smaller, of the order, much larger than the contour length of the chains. In section II we present the simulation technique, details of the model, and the way in which we imposed flow in the system. Simulation results are presented in section III, devoting subsections III A and III B to equilibrium properties and behavior under flow, respectively. We provide a final discussion and conclusions in Section IV.

II. MODEL AND SIMULATION TECHNIQUES

The polymer chains are simulated by the widely utilized coarse-grained beadspring Kremer-Grest model[44, 45]. A finite extensible nonlinear elastic (FENE) potential models the bond interactions between neighboring beads of the same chain.

\[
U_{\text{FENE}} = -\frac{1}{2}kR_0^2\ln \left( 1 - \left( \frac{r}{R_0} \right)^2 \right), \quad r < R_0,
\]

\[
U_{\text{FENE}} = \infty, \quad r \geq R_0
\]

where \(R_0 = 1.5\sigma\), the spring constant \(k = 30\epsilon/\sigma^2\), and \(r = |r_i - r_j|\) is the distance between neighboring monomers. This potential simulates the conformational dynamics of polymers for a variety of thermodynamic conditions[44, 45]. All monomers of the \(N = 10\) polymer chains we simulated, interact with each other through truncated and shifted Lennard-Jones potential

\[
U_{\text{LJ}}(r) = \begin{cases} U_{\text{LJ}}(r) & r < r_C, \\
0 & r \geq r_C \end{cases}
\]

with \(U_{\text{LJ}}\) is the standard Lennard-Jones potential:

\[
U_{\text{LJ}} = 4\epsilon_{\text{ab}} \left( \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right),
\]

where \(\epsilon_{\text{ab}} = 1\) defines the length unit. \(r_C\) is the interaction cut-off radius: if the distance between particles exceed \(r_C\), this interaction is turned off. For the interaction between particles belonging to the polymer, the cut-off radius was set to the minimum of the Lennard-Jones potential \(r_C = \sqrt{2}\sigma \approx 1.12\sigma\), giving a purely repulsive force which is interpreted as a polymer in good solvent conditions. For the liquid-liquid and liquid-polymer interactions, the cut-off radius was set to twice the minimum of the potential \(r_C = 2\sqrt{2}\sigma \approx 2.24\sigma\), to include...
attractive interactions. This allows the formation of liquid and gas phases for temperatures below the evaporation point. The interaction strength $\varepsilon_{ab}$ was set to unity for the polymer-polymer and liquid-liquid interactions $\varepsilon_{pp} = \varepsilon_{ll} = \varepsilon = 1$, and defines the energy unit. The temperature unit is therefore $\varepsilon/k_B$. For polymer-liquid interactions the parameter was set to $\varepsilon_{pl} = 1/3\varepsilon$, representing the chemical incompatibility between species.

Rough surfaces with similar interaction potential and $\varepsilon_{pl}$ parameter have been studied previously[46–49], resulting in highly hydrophobic substrates and yielding to contact angles $\theta \sim 120^\circ$ for droplets deposited over that surfaces.

The rigidity of the polymers is imposed by a harmonic potential

$$U_b(r_{i+1}, r_{i+2}) = \frac{1}{2}k_b\theta_{i+1,i+2}^2,$$  

(4)

where $\theta_{i+1,i+2}$ is the angle formed by two consecutive bonds, and is measured from the polymer backbone. The bending constant $k_b$ is one of the parameters varied in this work to explore different persistence lengths. It is important to note that the sum of bending forces over all the beads is null, so the forces applied on the free end bead are transmitted, in average, to the grafted bead. To induce a privileged orientation of the polymer chains perpendicular to the substrate, the bending force calculations are performed adding a phantom bead exactly below the fixed end. This additional bond generates a force on the second bead, favouring an orientation angle $\theta_1 = 90^\circ$ with respect to the wall.

The walls at $z = 0$ and $z = D$ were modelled as impenetrable flat surfaces, which interact with monomers via a purely repulsive integrated Lennard-Jones potential of the form

$$U_{wall}(z) = |A| \left( \frac{\sigma}{r} \right)^9 + A \left( \frac{\sigma}{r} \right)^3,$$  

(5)

with $A = 3.2\varepsilon[19, 22]$.

The grafted end-beads of the 10-bead polymers are placed at a distance $z = 1.2\sigma$ from the substrate, and the lateral positions is randomly chosen. The number of grafted sites is determined by the grafting density. The initial positions of the liquid particles are uniformly random distributed, and a force switch-on mechanism is used to relax the system towards thermodynamic equilibrium and eliminate the overlap between beads. Production runs of typically $10^7$ steps are done after this relaxation and $10^6$ time steps of thermalization.

We used the dissipative particle dynamics (DPD) scheme[50, 51] to perform the simulations at constant temperature. An important feature of this thermostat is that it accounts for correct hydrodynamic behaviour[52], which is essential to obtain reliable results out of equilibrium[53–55]. The equation of motion for a particle reads:

$$m\frac{d^2r_i}{dt^2} = F_i + F_i^D + F_i^R,$$  

(6)

where $F_i$ represents the conservative forces, $F_i^D$ and $F_i^R$ are the dissipative and the random forces, respectively. The mass $m$ is set to unity for all particles. Random and dissipative forces are applied in pairs, in such a way that local momentum is conserved. The form for these forces is

$$F_i^D = \sum_{j\neq i} F_{ij}^D, F_{ij}^D = -\gamma\omega^D(r_{ij})(\hat{r}_{ij}\cdot \mathbf{v}_{ij})\hat{r}_{ij},$$  

$$F_i^R = \sum_{j\neq i} F_{ij}^R, F_{ij}^R = \zeta\omega^R(r_{ij})\eta_{ij}\hat{r}_{ij},$$  

(7)

where the notation $\mathbf{a}_{ij} = \mathbf{a}_i - \mathbf{a}_j$ is used for vectors $\mathbf{a}$ is the friction constant, $\zeta$ denotes the noise strength, $\omega^D$ and $\omega^R$ are weight functions. These parameters and functions are not completely free to choose in order to satisfy the fluctuation-dissipation theorem. This constrain imposes the condition $\zeta^2 = 2\epsilon T \gamma$, and $[\omega^R]^2 = \omega^D$. The parameter $\gamma$ was set to unity in all simulations, because this value allows to maintain constant temperature and doesn’t produce a significative over-damping of conservative forces[55]. The random variable $\eta_{ij}$ has zero mean and second moment $\langle \eta_{ij}(t)\eta_{ij}(t') \rangle = \delta_{ij} \delta_{ij}(t - t')$. The usual choice of the weight functions for continuous forces was made[53, 55].

$$[\omega^R]^2 = \omega^D = \begin{cases} 1 - \gamma(r_C)^2, & r < r_C \\ 0, & r \geq r_C \end{cases}$$  

with the cut-off radius is chosen equal to that that Lennard-Jones interaction.$r_C = 2\sqrt{2}\sigma \approx 2.24\sigma$.

We define the time units in terms of the Lennard-Jones parameters and the mass $\tau = \sqrt{m/\varepsilon}$. The Verlet[56] integration scheme was used to integrate the equations of motion 6, with time steps in the range $2 \times 10^{-3}\tau \leq dt \leq 5 \times 10^{-4}\tau$. Under these conditions, it is possible to perform equilibrium and non-equilibrium simulations, maintaining a constant temperature over the whole system. We tested this with temperature profiles across the channel, obtained from the mean square velocity of the particles.

The temperature value $T = 0.8\varepsilon/k_B$ was chosen to set the liquid far enough from the condensation point and to have a relatively low vapor density[57]. In thermal equilibrium, the liquid and gas number densities are $n_l = 0.69\sigma^{-3}$ and $n_v = 0.03\sigma^{-3}$, respectively. Several systems were studied, varying the number of grafted chains per surface area and bending constant. In order to compare quantitatively the behaviour of the different systems, it is necessary to maintain some magnification constant for all cases. We chose to maintain the initial density in the bulk equal in all studied systems to provide a consistent thermodynamic condition. Therefore, the number of liquid particles was changed so that in every equilibrated system, the bulk liquid density satisfied...
The relaxation times of the polymers for different grafting densities were also measured in simulations without liquid and at low grafting density ($\rho_g = 0.05 \sigma^{-3}$). Studying the autocorrelation of the components of the end-to-end vector of the polymers it is possible to extract a characteristic correlation time of the chains:

$$C_i(t) = \frac{\langle (R_i(0) - \langle R_i \rangle) (R_i(t) - \langle R_i \rangle) \rangle}{\langle (R_i - \langle R_i \rangle)^2 \rangle},$$

where $R_i(t)$ is the $i$-component of the end-to-end vector. Four values of the bending constant $k_b$ were studied over a wide range of $k_BT$ values ($k_b/k_BT = 0; 1; 10$ and $100$), giving the possibility to explore different rigidity regimes, from fully flexible chains ($k_b = 0 \varepsilon$), to very rigid rods that resemble pillars ($k_b = 80 \varepsilon$). For $k_b = 0 \varepsilon$, $0.8 \varepsilon$, $8 \varepsilon$ the correlation follows an exponential decay and the relaxation time can be obtained by inferring the curve $C(t) = \exp(-t/\tau_R)$. The relaxation of the polymers with $k_b = 80 \varepsilon$ is more complicated and does not follow a simple exponential decay. For short times the behaviour is exponential, but it presents in addition a longtime correlation tail. This suggests complicated dynamics of the internal degrees of freedom of the molecules excited by thermal fluctuations. In order to estimate a relaxation time in this case we fitted an exponential law up to a correlation value of $C(t) = 0.05$. This gives a relaxation time which is roughly the double of the longtime correlation decay ($\tau_R = 125 \tau$).

Relaxation times of the grafted chains were also measured in simulations including the liquid. The results were fitted in less than $30\%$ from those measured without liquid, except for the case of $k_b = 80 \varepsilon$. For the stiffer chains ($k_b = 80 \varepsilon$), we found that the collisions between liquid particles and the free end of the polymers decrease dramatically the relaxation time, turning out in a lower value than that of fully flexible chains. Taking this into account, it is evident that the relaxation times can be highly dependent on the particular system where the polymers are embedded. In order to obtain a characteristic polymer property, we decided to use the relaxation times calculated in absence of liquid, which are shown in Table I. We do not intend to study the complex dynamics of these highly rigid semiflexible polymers, but to estimate an order of magnitude of their relaxation times. A thorough study should be conducted to understand the relaxation kinetics of these semiflexible chains.

![Figure 1: Left panel: Snapshot of the simulated system for a grafting density of $\rho_g = 0.05 \sigma^{-3}$ and bending constant $k_b = 80 \varepsilon$ in equilibrium. Particles of the liquid are shown in light gray, the polymer chains in the brush are presented in red, while their end beads (steadily fixed to the wall) are shown in light-blue. Right panel: Configuration of the polymer brushes for the different bending rigidities studied (the liquid is not shown).](image-url)
We begin our discussion by presenting the monomer density profiles of the system in equilibrium for all the studied bending constants $k_b$ (see Figure 2). Upon increasing the rigidity of the bonds, the polymers stretch in the direction perpendicular to the grafted surface. This builds up the structure of the brush, noticeable in the decay of the bond correlation with the bond number, $\cos(\theta_s) = \exp(-s l_p / \langle |a| \rangle)$, it is possible to estimate the persistence length $l_p$. For the stiffest cases $k_b = 8\sigma$ and $k_b = 80\sigma$, simulations with larger polymers ($N = 100$) were also carried out, to estimate more precisely $l_p$, yielding similar results as those obtained for $N = 10$.

### III. RESULTS

#### A. Static properties

We begin our discussion by presenting the monomer density profiles of the system in equilibrium for all the studied bending constants $k_b$ (see Figure 2). Upon increasing the rigidity of the bonds, the polymers stretch in the direction perpendicular to the grafted surface. This builds up the structure of the brush, noticeable in the sharp peaks of the polymer density profiles. The chemical incompatibility between liquid and polymer particles, expressed by the L-J parameters $\epsilon_{pl} = 1/3 \epsilon_{ll}$, prevents the fluid from penetrating the brush in its liquid state. The liquid density decreases rapidly in the vicinity of the polymers, and only isolated particles enter in the brush, forming a gas in coexistence with the liquid. This gives rise to a well-defined, narrow liquid-brush interface, even for the lowest grafting density studied $\rho_g = 0.02\sigma^{-2}$. We note also that the fact the the liquid does not penetrate in the the brush agrees with the definition of the Cassie-Baxter state, present in nano-structured superhydrophobic surfaces[35, 58, 59]. This is illustrated for a relatively low grafting density $\rho_g = 0.05\sigma^{-2}$ in Fig. 1.

The height of the polymer brush varies significantly upon increase of the bending constant, due to the chains stretching. In Figure 3 the vertical component of the end-to-end vector ($|R_{end}|$) is shown as a function of the grafting density for the different bending constants $k_b$. $R_{end}$ is defined as the vector joining both end-beads of the same polymer. As it was observed for the density profiles (Figure 2), increasing the bending rigidity results in a thicker brush. The behaviour is similar for all bending rigidities: the chains elongate in the $z$ direction upon increasing $\rho_g$, due to excluded volume interactions, as reported in Refs. [18, 19] for brushes composed of fully flexible chains. This effect is less pronounced for the more rigid chains ($k_b = 8\sigma$) and barely noticeable for the stiffest case ($k_b = 80\sigma$), due to the finite extension

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<table>
<thead>
<tr>
<th>$k_b$ [ε]</th>
<th>$l_p/\ell_c$</th>
<th>$l_p/\langle a \rangle$</th>
<th>$k_b/k_BT$</th>
<th>$\tau_\varepsilon$ [τ]</th>
</tr>
</thead>
<tbody>
<tr>
<td>0</td>
<td>0.16 ± 0.01</td>
<td>1.4 ± 0.1</td>
<td>0</td>
<td>17</td>
</tr>
<tr>
<td>0.8</td>
<td>0.22 ± 0.03</td>
<td>2.0 ± 0.3</td>
<td>1</td>
<td>18</td>
</tr>
<tr>
<td>8</td>
<td>1.12 ± 0.06</td>
<td>10.1 ± 0.5</td>
<td>10</td>
<td>50</td>
</tr>
<tr>
<td>80</td>
<td>10.1 ± 1.7</td>
<td>91 ± 15</td>
<td>100</td>
<td>260</td>
</tr>
</tbody>
</table>

Table I: Persistence length ($l_p$) over contour length ($\ell_c$) for every studied bending constant ($k_b$) studied. In the third column is presented the persistence length over mean bond length ($\langle a \rangle$). The persistence length covers a wide range of polymer rigidities, from fully flexible chains, to highly rigid rods. The values were obtained from independent simulations of isolated chains. The last column shows an estimate of the relaxation time of the grafted polymers, obtained from the correlation time of the end-to-end vector autocorrelation function.
The energy (the chains end to end distance. The energy also increases with the proximity of the bond to the free end bead. For $k_b = 8\varepsilon$ the bending energy per bond remains roughly the same for all grafting densities.

The inset of Figure 3 shows the standard deviation of $R_z$, $SDR_z$. This magnitude is associated with the typical chain height variation, i.e. the mobility of the end bead in $z$ direction. Interestingly, there is a non-monotonic behaviour for the most flexible cases ($k_b = 0\varepsilon$ and $k_b = 0.8\varepsilon$). At low $\rho_g$, $R_z$ increases with grafting density due to a larger free space available for the chains. Naturally, this causes an increase in $SDR_z$. At high $\rho_g$, the presence of neighbour polymers hinders the mobility of the end bead, decreasing $SDR_z$. The chains with $k_b = 8\varepsilon$ have the greatest mobility at low $\rho_g$, but it decreases rapidly upon increasing $\rho_g$, because the bending potential allows global buckling of the chains, but no local bending. This means that the mobility of the end bead is highly dependent on the mobility of the entire polymer. $SDR_z$ is fairly constant as function of $\rho_g$ for the stiffest studied polymers ($k_b = 80\varepsilon$).

To investigate the behavior of the internal degrees of freedom of the polymers we measured the bond angles of each joint. Figure 4 shows the mean bending energy per bond $U_b = 1/2k_b\langle \theta_i^2 \rangle$ as a function of $\rho_g$, where $\theta_i$ is the angle formed by two consecutive bonds in the $i$-th joint. Naturally, $\theta$ decreases upon increasing the bending constant $k_b$, but this does not translate immediately in increment of $U_b$. At low $\rho_g$, increasing the bending constant from $k_b = 0.8\varepsilon$ to $k_b = 8\varepsilon$ roughly doubles the bending energy stored in each bond, but increasing the bending constant from $k_b = 8\varepsilon$ to $k_b = 80\varepsilon$ varies $U_b$ in less than 5%. For the lowest grafting density, the increase in the end-to-end distance of the polymers with persistence length $l_p$ is in good agreement with the Worm-like chain or Kratky-Porod model[60].

Upon increasing the grafting density, the excluded volume interactions induce polymer stretching (see Figure 3), which translates in a decrease in the bond angles. It can also be observed that bond angles near the free end bead tend to be larger than those near the grafted end. Particularly, the last bead is only subject to the bending force of the last bond and is constantly colliding with liquid particles, which leads to a greater variance.

The first "orientation" bond (not shown in Figure 4) behaves differently than the other bond angles of the polymers, due to the constraints to which it is subjected to. For low bending constants ($k_b = 0\varepsilon$ and $k_b = 0.8\varepsilon$), the orientation angle is always lower than the chain bond angles, because the wall limits the values this angle can take. This reasoning is not true for the most rigid polymers ($k_b = 8\varepsilon$ and $k_b = 80\varepsilon$), because the strong bending forces bound the internal bond angles as well as the orientation bond, and the behaviour for both, internal and orientation angles, is similar at low $\rho_g$. At $\rho_g \geq 0.2$, the presence of overlapping grafting points force the polymers to take high orientation angles as observed in the inset of Figure 4. This leads to a pronounced increase in the orientation bond energy, which comes only for the way in which the sample is generated. The grafting points are taken from a uniformly random distribution in the $x$-$y$ plane and, consequently, the probability of producing a sample with overlapping grafting beads increases with grafting density.
B. Flow properties

The system was taken out of equilibrium by moving the walls at constant velocity in opposite directions[18, 22]. This induces a flow with a linear velocity profile in the liquid phase indicated in Figure 1. In Figure 5 the symmetrized velocity profile of the liquid is shown for various shear rates. In the bulk of the liquid, a linear dependence can be clearly observed, as expected for a simple liquid under these boundary conditions. In the vicinity of the interface (vertical dotted lines), the behaviour changes and near the wall the free gas particles match the wall velocity. It is interesting to note that the liquid layer in contact with the polymer brush does not match the wall velocity, i.e. the no-slip condition is violated and there is a partial slip of the liquid on the brush. The shear rate \( \dot{\gamma} \) was obtained by fitting the linear velocity profile in the bulk of the liquid. \( \dot{\gamma} \) increases with the wall velocity, but also depends on the grafting density \( \rho_g \) and the bending constant \( k_b \). We define a Weissenberg number \( We \) to quantify the shear rate with the natural relaxation of the semiflexible polymers. \( We \) is defined as the shear rate \( \dot{\gamma} \) times the relaxation time \( \tau_R \) of an isolated chain at the same temperature as the simulated sample: \( We = \dot{\gamma} \tau_R \). As shown in table I, the relaxation time is highly dependent on the bending constant.

At high grafting densities, the density profiles are almost independent of the shear velocity, while at low \( \rho_g \), the brush height diminishes upon increasing shear rate. For lower grafting densities the polymers lean more, because there are less excluded volume interactions with other neighbours. The mean value of the normal component decreases. For high grafting densities, the mobility is less hindered by the bending rigidity. This effect that causes the stretching of fully flexible chains in equilibrium[60]. As expected, increasing the shear rate \( \dot{\gamma} \) induces a larger end-to-end angle, which means that the chains tilt in the direction of the flow. For \( \rho_g \leq 0.05 \),...
the inclination angle is always lower than 20°, and the increment of \( \Theta_{end} \) is linear with the shear rate. For lower grafting densities and high \( We \), there is a non-linear behaviour, due to saturation effects. The dependence of \( \Theta_{end} \) on the rigidity of the polymer is more complex. There are two different regimes according to the relation between persistence and contour lengths: a) \( l_p < l_c \) and b) \( l_p > l_c \). If the persistence length is lower than the contour length, increasing \( b \) (rigidity parameter) produces a higher end-to-end angle of the chains under shear. The consecutive bonds tend to stretch coherently in the shear direction, due to the correlation imposed by the bending potential. The pressure exerted by the simple liquid on the polymers induce chain buckling and they take a banana-like shape, maximizing \( \Theta_{end} \) (see inset in Figure 7). Increasing \( b \) further (regime \( l_p > l_c \)) prevents the chain to achieve high local bending angles, and the polymer stretches in the vertical direction, disfavouring the buckling phenomenon. As seen in Figure 3, the stiffest chains \( b = 80 \varepsilon \) in equilibrium are totally stretched vertically for all \( \rho_g \). It is important to recall that in this model there is a privileged orientation angle \( \Theta_{end} = 0° \), because the first bond \( \theta_1 \) is subjected to the same bending potential as the rest of the polymer with an orientation perpendicular to the wall.

Another interesting quantity, is the standard deviation of the \( R_{end} \) component parallel to the shear direction. This magnitude gives information about the typical mobility of the polymer’s free end-bead, around the mean value, in the shear direction. In Figure 8 is presented the dependence of \( SDR_x \) on \( \rho_g \), for various shear rates and bending rigidities. It can be observed that for \( l_p < l_c \) (\( b = 8 \varepsilon \)), increasing bending rigidity induces a higher displacement of the free end. The pressure exerted by the liquid causes the more rigid polymers (\( b = 8 \varepsilon \)) to bend in a defined direction, due to correlation between bonds, and thus enhances their displacement. Increasing \( b \) above \( l_p \approx l_c \) hinders the mobility of the chain, because of the strong bending potential. This behaviour was previously observed by Milchev and Binder[61] for brushes in equilibrium without explicit solvent. An interesting feature of \( SDR_x \) is that for high \( \gamma \), the dependence on the grafting density is non-monotonic. At low \( \rho_g \) the chains lean in the shear direction and don’t relax to the equilibrium configuration, thus having a small mobility. Increasing the number of neighbouring chains produces a screening effect, allowing the free end of the polymer to gain mobility. \( SDR_x \) reaches a maximum at intermediate \( \rho_g \), because at high grafting densities the mobility of the free end decreases, due to excluded volume interactions.

The typical displacement \( SDR_z \) is compared to the mean distance between grafting points (dashed line in Figure 8). As expected for high \( \rho_g \), the chain movement is limited by the presence of neighbour polymers, and \( SDR_z \) match the mean distance between neighbours. At low grafting densities, the standard deviation of \( R_z \) is limited by the finite length of the polymer, and by anisotropic effects. At intermediate \( \rho_g \), \( SDR_z \) is higher than the mean distance between grafting points for semiflexible chains (\( b = 8 \varepsilon \) and \( b = 8 \varepsilon \)). It may be also observed that for rigid rods (\( b = 80 \varepsilon \)) \( SDR_z \) is always less than the mean distance between grafting points. This indicates that even for high \( \rho_g \), the dynamics of these polymers are greatly affected by the bending force.

We also checked the existence of cyclic motion dynamics in semiflexible chains under shear, found in previous works for fully flexible isolated grafted chains or polymer brushes[20, 22]. We analysed the effect of stiffness in the mechanism of cyclic motion, whose origin is attributed to spontaneous fluctuations of the polymer chains towards regions of higher velocity in the liquid. According to this description, a stiffer chain should have a reduction in the cyclic motion, because the spontaneous fluctuations are reduced. This is in fact what we observe. We computed the mean brush momentum velocity profile \( p(z) = \langle m_p(z) v(z) \rangle \), exactly in the same way, than previous references[22] (not shown). The presence of cyclic motion is evidenced when a positive \( p(z) \) is observed in the region in which the chains are directly exposed and dragged by the fluid and, at the same time, a negative \( p(z) \) is observed in an inner region, inside the brush. We found cyclic motion dynamics for bending constants...
k_b = 0.8\varepsilon \text{ and } k_b = 8\varepsilon. This was the case for low grafting density ($\rho_g = 0.02\sigma^{-2}$ and $0.05\sigma^{-2}$), in which the dynamics of chains can be regarded as independent, but also for higher grafting ($\rho_g = 0.3\sigma^{-2}$) already in the brush regime. The cyclic dynamics is indeed reduced upon increase of the chain stiffness and vanishes completely for the stiffest chains ($k_b = 80\varepsilon$). A more comprehensive characterization of the effect and its consequences for flow inversion[20, 22] could be an interesting future study.

In Figure 9 the probability density function (PDF) of the end to end inclination angle $\Theta_{end}$ is presented for various Weissenberg numbers and bending rigidities. The equilibrium distribution has mean value $\langle \Theta_{end} \rangle = 0$ for all cases, and high variance $\text{var}(\Theta_{end}) = \langle \Theta_{end}^2 \rangle - \langle \Theta_{end} \rangle^2$. Upon increasing the shear rate, the chains lean in the flow direction and decrease their mobility, thus raising the mean value and reducing the variance. The shape of PDF ($\Theta_{end}$) is highly non-symmetrical with respect to the maximum for the samples under shear. This behaviour is observable for all bending rigidities, except for $k_b = 80\varepsilon$ that is less noticeable. It is interesting to note that all the PDF deviate from a normal distribution, and particularly the tails follow a non-gaussian decay. This non-gaussian behaviour was already found in theory and experiments of single chains under shear[62].

The buckling process for the case $k_b = 8\varepsilon$ can also be observed from the equilibrium PDF (black circles, in lower left panel in Fig. 9). There are two maxima, which correspond to $\Theta_{end} = \pm 40^\circ$, and there is a local minimum at $\Theta_{end} = 0$, which corresponds to the free-end of the polymer, vertically aligned with the grafted monomer. The pressure exerted by the liquid induces a buckling in the semiflexible chains, adopting the shape shown in the inset of Fig. 9. This effect is most noticeable for chains whose persistence length is similar to the contour length, due to the global order these polymers achieve. For lower bending, a flat plateau at the maximum, rather than a bimodal distribution is observed. While for high bending rigidity ($k_b = 80\varepsilon$), the equilibrium distribution corresponds clearly to a normal curve.

It’s interesting to analyze how the internal degrees of freedom of the polymers are affected by the shear stress. In Figure 10 is presented the bending energy $U_b$ for each bond as a function of the Weissenberg number $We$, for the lowest grafting density ($\rho_g = 0.02\sigma^{-2}$). Two different regimes can be distinguished: for $l_p \ll l_c$ (panels a and b) the orientation angle (Bond 1) is considerably smaller than the internal angles, and the angles decrease upon increasing shear rate, thus increasing the end to end distance. In this case, the total bending energy is reduced with increasing shear rate. For $l_p \gg l_c$ (panel d), the orientation angle is greater than the internal angles, and the bending energy of all bonds increase upon increasing $We$. It is also observed a decrease in the distance between grafted and end beads. In this regime, the chains tend to take a banana-like shape, leaning with the flowing liquid, and decreasing the end to end distance.

Polymers of $k_b = 8\varepsilon$ are a limit case (panel c), because the persistence length coincides with the contour length. The average Bending energy (solid black dots in Fig. 10) and the end-to-end distance (not shown) changes less than 2%. In this limiting case, there are no important changes in the internal structure of the polymers, which implies that the external perturbation only alters the inclination angle. In addition, it can be observed that the variation of bending energy between internal angles of each...
same chain do not vary more than 30% for $We \leq 4$. This means that the contribution of the internal angles to the total inclination of the chain is roughly equidistributed along the polymer, as already observed by Milchev and Binder[61].

As mentioned above, the no-slip condition was violated in every system studied. The friction reduction in systems with polymer brushes has been a well known interesting feature for over a decade now[4]. We quantify here the effects of stiffness in the semiflexible polymers. To quantify this phenomenon, the slip length ($\delta$) was measured and plotted versus the grafting density $\rho_g$ in Figure 11. $\delta$ is defined as the distance at which the extrapolation of the far field velocity profile of the liquid matches the substrate velocity (see Ref. [59]). The liquid-polymer interface position was defined as the $z$ coordinate at which the function $f(z) = p_l(z) - p_p(z)$ presents a maximum, where $p_l$ is the liquid density, $p_p$ is the polymer density and $z$ is the coordinate perpendicular to the walls. At low $\rho_g$, the slip of liquid over the polymer brush is highly dependent on the properties of the substrate. Varying the number of grafted polymers per surface area and their rigidity yield different results. In general, $\delta$ decreases rapidly upon increasing grafting density, but the dependence with the stiffness ($k_b$) is non-monotonic. At intermediate to high grafting densities ($\rho_g > 0.1\sigma^{-2}$), $\delta$ seems to saturate and takes a constant value, regardless of the flexibility parameter $k_b$, the shear rate ($\dot{\gamma}$), or the grafting density. Only for rigid rods ($k_b = 80\epsilon$), there is a decrease in the slip at high grafting densities, and $\delta$ is a decrease in the slip at high grafting densities, and seems to saturate and takes a constant value, regardless of the stiffness in the semiflexible polymers. To quantify these cases present non-zero slip length, indicating that the no-slip condition between the polymer brush and the free melt does not hold. The rod-like polymers ($k_b = 80\epsilon$) are the ones that offer the most resistance to the liquid flow for every grafting density and shear rate.

Figure 11: Slip length ($\delta$) as a function of the grafting density ($\rho_g$) for the studied shear rates. The slip length decreases rapidly upon increasing $\rho_g$. For intermediate and high grafting densities ($\rho_g \geq 0.1\sigma^{-2}$), $\delta$ takes roughly the same value $\delta \approx 11\sigma$ for every wall velocity and bending rigidity. All the studied cases present non-zero slip length, indicating that the no-slip condition between the polymer brush and the free melt does not hold. The rod-like polymers ($k_b = 80\epsilon$) are the ones that offer the most resistance to the liquid flow for every grafting density and shear rate.

We $\leq 1$. We think that above $\rho_g = 0.1\sigma^{-2}$, the slip length is determined strongly by the brush-liquid compatibility which was fixed in this study with the value $\varepsilon_{lp} = 1/3\varepsilon_{ll} = 1/3\varepsilon_{pp}$.

In Figure 12 is presented the slip length $\delta$ as a function of the Weissenberg number $We$, for a low grafting density $\rho_g = 0.05\sigma^{-2}$. For every bending constant $k_b$, $\delta$ increases upon increasing shear rate, at low grafting densities, as it was already reported for fully flexible chains[18]. A very interesting feature is that semiflexible chains with $k_b = 0.8\epsilon$ ($l_p/c_l \simeq 0.2$) induce higher slip lengths than fully flexible chains, for all the studied shear rates. We think that a moderate stiffness in the polymer could reduce the contacts of the chains with the liquid as compared to the case of fully flexible chains in mushroom regime. The dependence of $\delta$ on $k_b$ is non-monotonic. Increasing further the stiffness ($k_b = 8\epsilon$ to $k_b = 80\epsilon$) decreases the slip length. Even though the slip for $k_b = 8\epsilon$ is smaller than the slip of more flexible chains, the rate of change of $\delta$ with $We$ is roughly the same for $We \geq 0.5$. The rigid-rods brush ($k_b = 80\epsilon$) is harder, and it shows a greater capacity to drag liquid along, than the more flexible brush which has lower slip lengths in the whole range of $We$ studied, and the growing slope is also lower. It is known that the slip length is highly dependent on the normal pressure: the higher the pressure, the smaller the slip[59]. In this case, the normal pressure of the different systems varied less than 15%, and decreased monotonically with the bending constant. This implies that the non-monotonic dependence of the slip length on the bending constant can not be attributed only to variations in the normal pressure.
density $\rho$, shear rate, and the slope is highly dependent on the grafting density $\rho$. It can be observed that $\mu$ is proportional to the shear rate, and the slope is highly dependent on the grafting density $\rho$. For systems with the same $\rho$, the friction coefficient is higher, the higher the bending constant $b$. Inset: double-logarithmic scaling plot for the slope of $\mu$ vs $\dot{\gamma}$, as a function of the grafting density. The dependence of the slope on $\rho$ seems to follow a power law.

The force applied to the system in the direction of the wall displacement as a function of the shear rate was also analyzed. The data from the eight different grafting densities, four bending constants, and five shear rate values collapse to a single straight line (not shown). This means that the relation between the total force exerted on the walls and the shear rate is independent of the polymer brush, and only depends on the liquid. It was possible to extract the viscosity of the simple liquid from the relation: $F_x = \eta \dot{A} \dot{\gamma}$, obtaining a value of $\eta = 1.28 \pm \sigma^{-2}$ [18, 21]. This essentially shows that the liquid behaves as a simple (newtonian) liquid with a well defined viscosity for the whole range of shear rates, and that the different morphologies of brush-liquid interface or gas content inside the brush layer does not have a significant role in the diffusion of momenta in the system.

Another useful quantity to compare different polymer brushes is the kinetic friction coefficient ($\mu$), defined as the quotient between shear and normal forces: $\mu = F_x / F_z$. Figure 13 presents $\mu$ versus $\dot{\gamma}$ as a function of the bending constant $b$, for different grafting densities $\rho$. It may be observed that the friction coefficient does not depend strongly on the stiffness of polymer chains. The variation of $\mu$ between systems with the same grafting density and under the same shear rate is in all cases less than 20%. Upon increasing the grafting density, the friction coefficient decreases. As encountered in other studies [14, 21], $\mu$ increases with shear rate. In this case it is observed that the relation between $\mu$ and $\dot{\gamma}$ is linear for high $\rho$, because the normal force variation is very small with shear rate for intermediate to high grafting densities.

Every curve was fitted, performing a linear regression of the form $\mu = m \dot{\gamma}$, with only one parameter ($m$). The slope ($m$) of each curve is plotted against $\rho$ in the inset of Figure 13, in a log-log scale. This implies that $m$ follows a power law of the form $m = B(k_b) \rho^{-1}$, where $B(k_b)$ is a function that depends on $k_b$ only. The exponent of $\rho$ turns out to be $-1$ for all the bending rigidities. Combining the last two equations is possible to obtain a scaling law for the behaviour of the friction coefficient $\mu$ as a function of the shear rate $\dot{\gamma}$ and the grafting density $\rho$:

$$\mu \sim B(k_b) \dot{\gamma} \rho^{-1}$$

It is interesting to note that the friction coefficient decreases with increasing number of grafted chains per surface area. We think that increasing $\rho$ hinders the liquid from penetrating the polymer brush. Adding polymer chains in this superhydrophobic system could result in a smoother liquid-brush interface, thus decreasing $\mu$. The components of the total force ($F_x$ and $F_z$) acting on the grafted end of the polymers have a different dependence on the grafting density (not shown). While $\mu$ increases linearly with $\rho$, $F_z$ increases rapidly at low $\rho$, and at high grafting densities ($\rho > 0.3 \sigma^{-2}$) seems to reach a saturation value. We note however, that friction phenomena are frequently scale dependent and friction coefficients give different results in experiments measurements at nanoscopic and macroscopic levels [22]. They have usually contributions from dissipative processes at various length scales. Also, macroscopic concepts are not readily extrapolated down to small scale topographies [59].

IV. CONCLUSIONS

In this work, we studied comprehensively the properties in equilibrium and under flow, of an interface of grafted semiflexible chains and a simple liquid. We chose the interaction parameters between liquid and polymer chains such that the liquid is in super-hydrophobic regime on the top of the brush in the well-known Cassie-Baxter state. In such, the liquid does not wet the interior of the grafted layer of polymers and only molecules belonging to the coexisting gas phase enter in the grafted polymer layer. The grafting density was varied in a wide range, from mushroom regime of mostly isolated chains to a dense brush with significant excluded volume interactions among chains. Also the stiffness of the chains were studied in an ample physical range, from fully-flexible chains to highly rigid chains, in which the polymers resemble pillars, which were used as structures to produce super-hydrophobic surfaces. We characterized the properties of the chains by independent simulations in which we calculate the persistence length, and the relaxation dynamics, by means of the time correlation function of the end-to-end vector.
We analysed the equilibrium structure of the grafted polymer layers with density profiles, the behavior of the end-to-end vector components, mean angles with the grafting plane, and bending energy of the chains. The mean bending energy per bond behaves in a non-trivial way, while for low bending constants ($k_B \sim k_BT$) the energy increases with stiffness, for high $k_B$ it reaches a saturation value. We studied also the probability density function of these angles in equilibrium and under flow, obtaining non-gaussian distributions. We found a rich and distinctive behavior of the chains for different bending rigidities. This is also reflected in the typical shape of the end-to-end vector in the direction parallel to the substrate. In the rigid rod limit, the polymers adopt typically a straight vertical configuration, and buckling phenomenon is observed, due to the normal force exerted by the liquid. This leads to a higher displacement that the chains adopt while supporting the liquid layer on the substrate. In the rigid rod limit, the polymers adopt typically a straight vertical configuration, and buckling phenomenon is observed, due to the normal force exerted by the liquid. This leads to a higher displacement that the chains adopt while supporting the liquid layer on the substrate. We measured the slip length of the liquid and found a saturation flow as function of grafting density and shear rates. We measured also the friction coefficient of the interface by means of the mean normal and shear forces, needed to impose a given shear rate. For a given shear rate, we found that friction is inversely proportional to grafting density. It could be interesting to expand these studies analysing the differences in the reported behavior for liquids of lower chemical incompatibility with the brush chains and the dynamics of droplets on top of these polymer-coated surfaces of semiflexible chains. This should be also of relevance for the fields of super-hydrophobic responsive surfaces and microfluidics.

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